

Habitable Evaporated Cores

Converting Mini-Neptunes into Super-Earths in the Habitable Zone of M Dwarfs

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Parameter	Range	Default
M_* (M_\odot)	0.08 – 0.4	-
M_p (M_\oplus)	1 – 10	-
R_{XUV} (R_p)	1.0 – 1.2	1.2
a	IHZ - OHZ	-
e	0.0 - 0.9	-
$P_{0,*}$ (days)	1.0 – 100	30.0
f_H	10^{-6} – 0.5	-
ϵ_{XUV}	0.1 – 0.4	0.3
ξ_{min}	$1 + 10^{-5}$ – 3	3
Atmos. esc.	R/R-Lim / E-Lim	-
Tidal model	CPL/CTL	CTL
Q_*	10^5 – 10^6	10^5
Q_p	10^1 – 10^5	10^4
τ_* (s)	10^{-2} – 10^{-1}	10^{-1}
τ_p (s)	10^{-3} – 10^3	10^{-1}
β	0.7 – 1.23	1.23
t_{sat} (Gyr)	0.1 – 1.0	0.1
t_0 (Myr)	10.0 – 100.0	10.0
t_{stop} (Gyr)	0.01 – 5.0	5.0

Table 1: A list of free parameters and their ranges, as well as default values used in the plots.

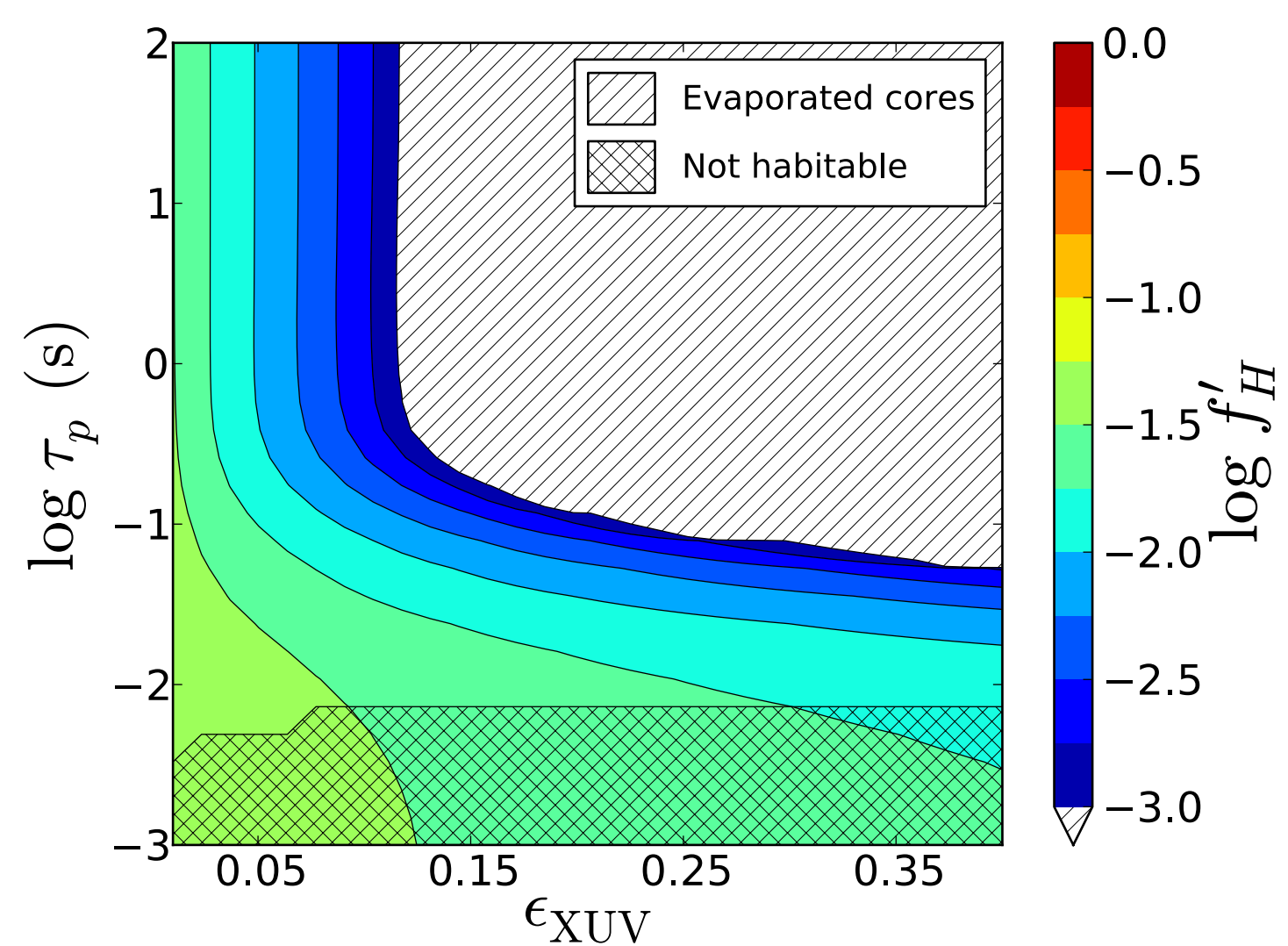


Figure 5: Final H fraction as a function of XUV absorption efficiency and tidal time lag for a sample run of our code. Mass loss and tidal evolution are strongly coupled in these systems.

References

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The low masses and luminosities of M dwarfs make them ideal targets for the detection of terrestrial planets in the habitable zone (HZ¹). However, studies suggest that planets formed only from material in the HZs of these stars are likely to be small and dry^{2,3}. Moreover, the long contraction phase of M dwarfs results in these stars being superluminous for as long as 1 Gyr (Figure 3), exposing planets that will eventually be in the HZ to extreme levels of radiation. Earth-like planets in this regime will enter a runaway greenhouse and may be completely desiccated by H loss to space, particularly around late M dwarfs (Figure 4).

As a result, habitable planets around M dwarfs may preferentially form from bodies that migrate into the HZ from beyond the snow line. These planets are enhanced in volatiles and may have accreted large H/He envelopes, which could help shield the surface from the extreme radiation environment early on. In this study, we show that photoevaporation and Roche lobe overflow of migrating “mini-Neptunes” can lead to the complete loss of their gaseous envelopes, transforming them into potentially habitable worlds, which we call **habitable evaporated cores (HECs)**. We couple planet structure evolution models⁴ with a simple Roche lobe overflow scheme and an extreme ultraviolet (XUV)-induced mass loss model^{5,6}. We also couple the orbital effects of anisotropic mass loss with tidal evolution⁷ and show that this coupling can lead to orbital changes that significantly enhance the mass loss rate.

HECs are most likely to form from small mini-Neptunes ($\lesssim 4 M_\oplus$) with large ($\lesssim 50\%$) initial hydrogen fractions orbiting M4 stars and later. Given the steep decrease in stellar XUV flux with time⁸, mass loss is negligible after ~ 1 Gyr, at which point a habitable evaporated core is capable of degassing and maintaining a secondary atmosphere. This process may be one of the dominant formation mechanism for habitable planets around M dwarfs, and may be discovered by missions such as TESS and PLATO.

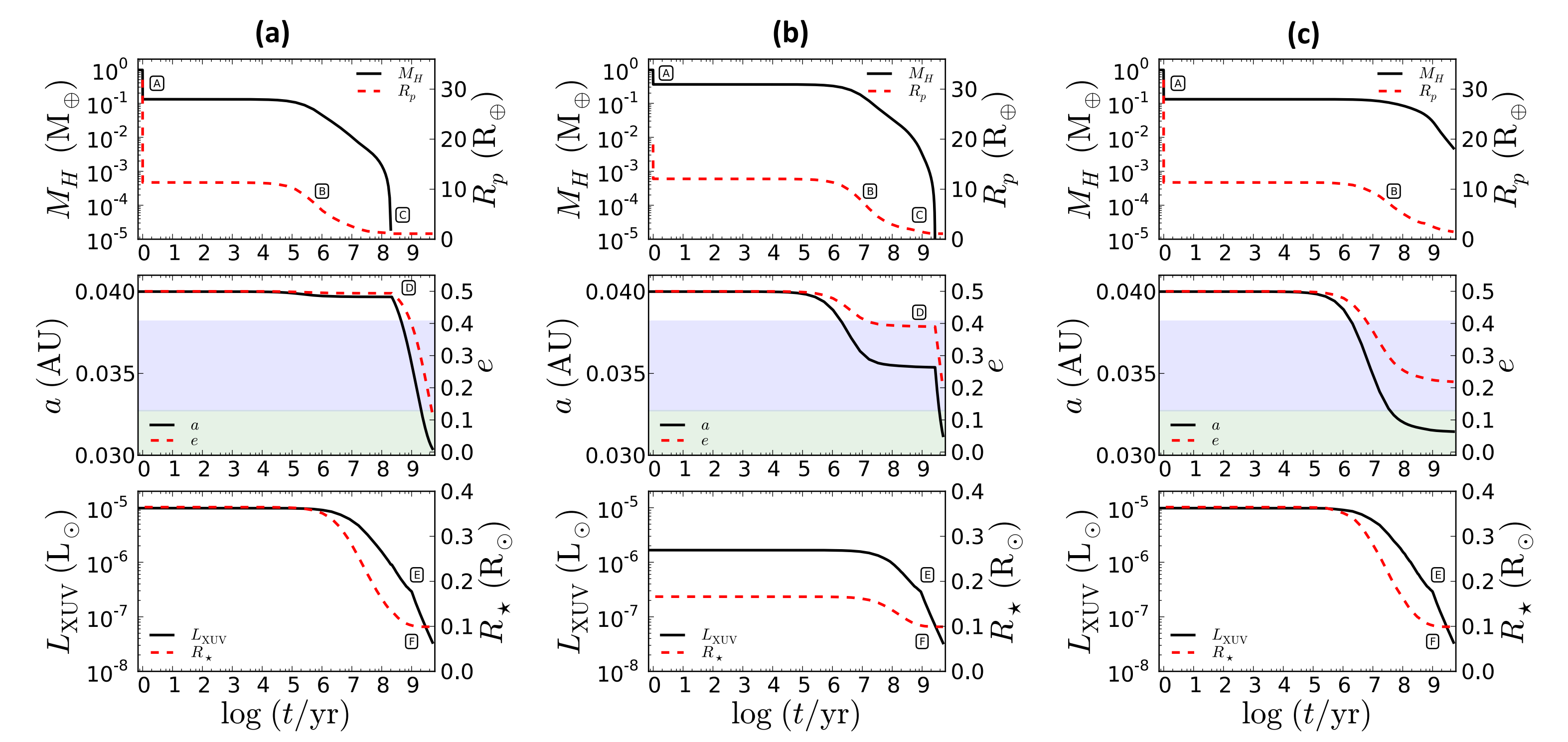
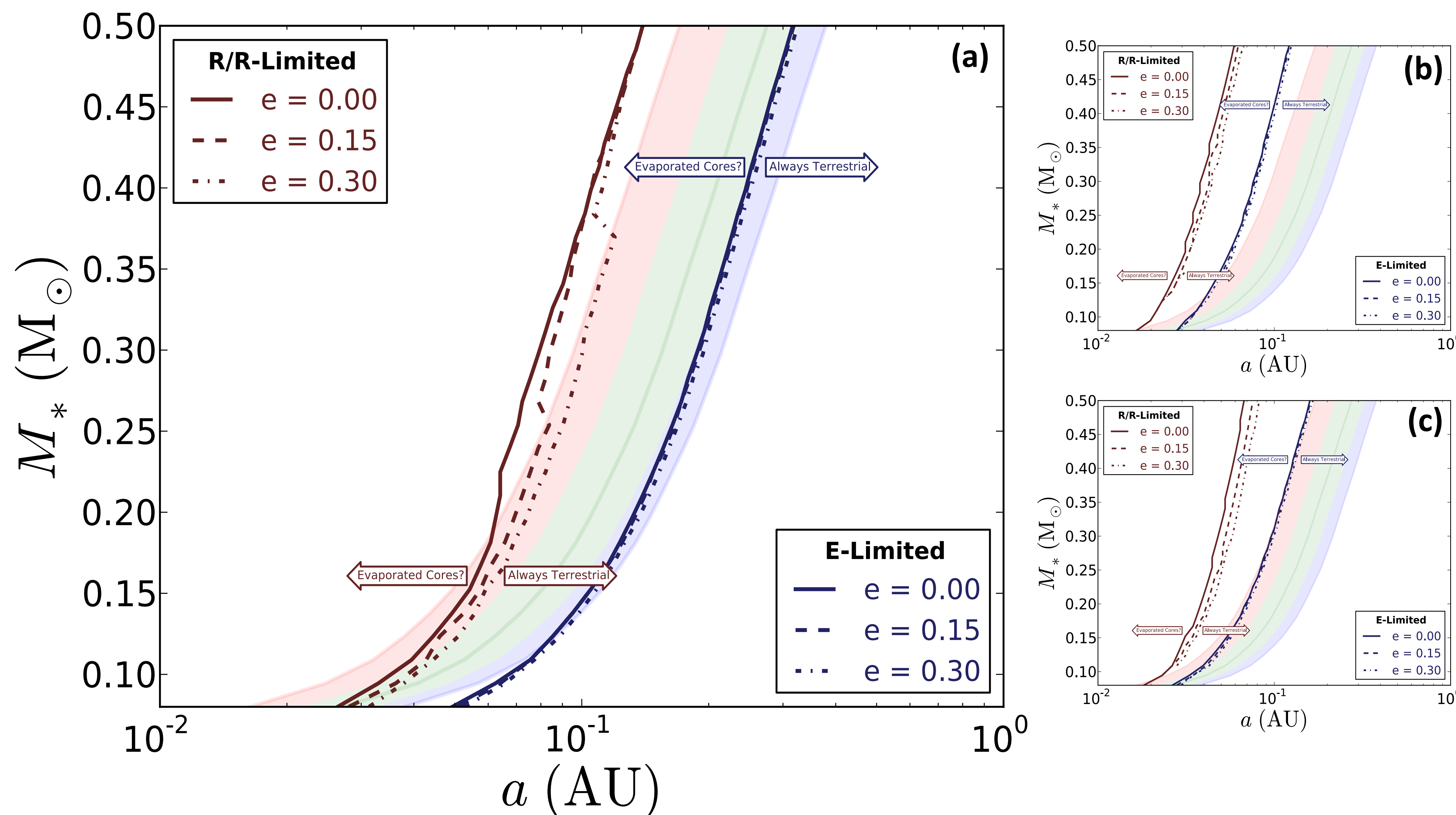


Figure 2: Three sample integrations of our code, showing envelope mass (M_H), planet radius (R_p), semi-major axis (a), eccentricity (e), stellar XUV luminosity (L_{XUV}) and radius (R_*) as a function of the time t since formation. The planet is initially a $1 M_\oplus$ core with a $1 M_\oplus$ envelope orbiting a $0.08 M_\odot$ M-dwarf with $e = 0.5$ at a semi-major axis of 0.04 AU, just outside the OHZ (outer HZ; light blue shading). As it loses mass and tidally evolves, it migrates into the CHZ (center HZ; light green shading). **(a): Energy-limited escape.** The planet loses its entire envelope at $t \approx 100$ Myr. **(b): Late migration**, with $t_0 = 100$ Myr. While the envelope still completely evaporates, this occurs at a much later time, $t \approx 2$ Gyr. **(c): Radiation/recombination-limited escape.** Here, the envelope does not fully evaporate and tidal migration is noticeably weaker.

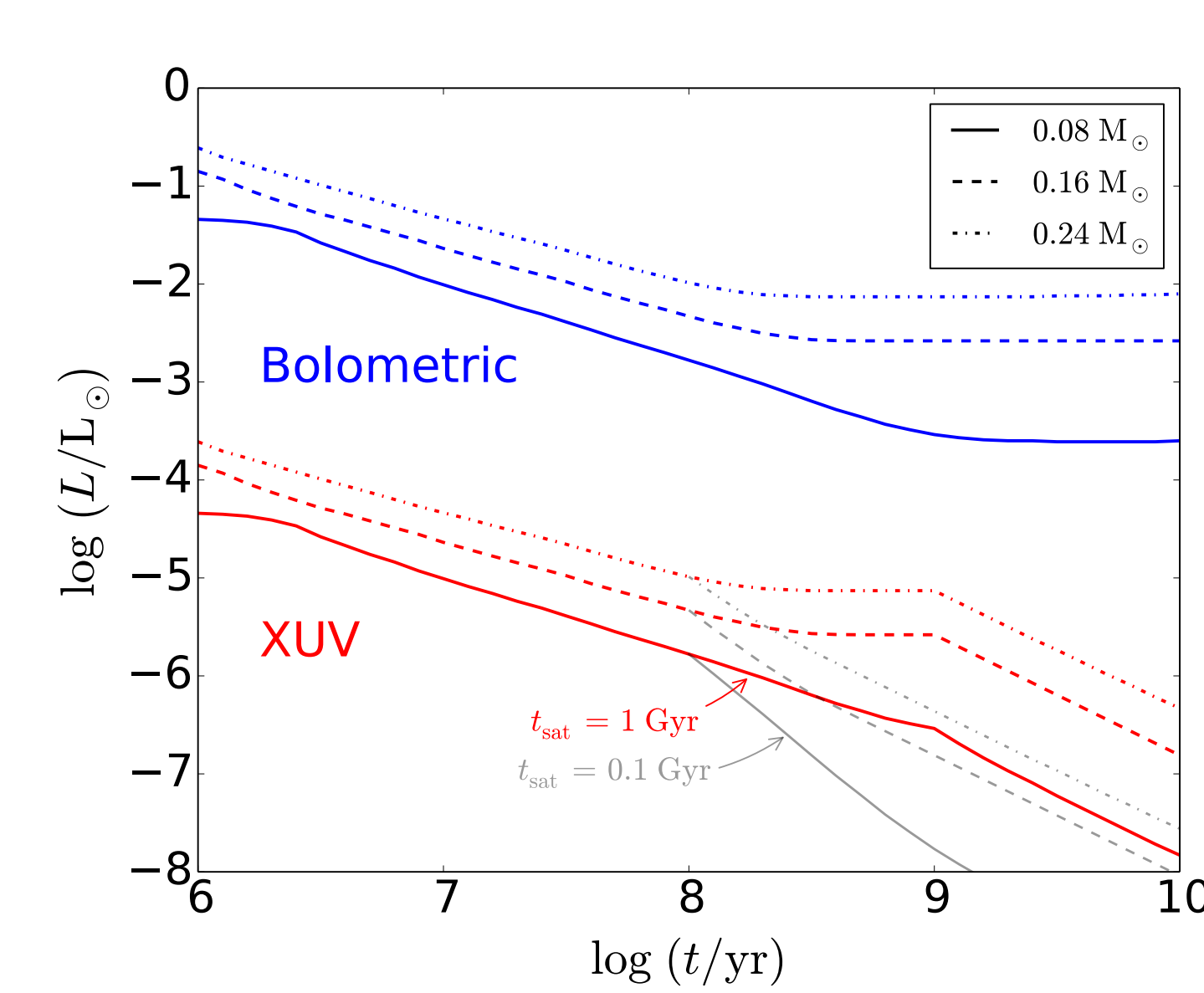


Figure 3: Evolution of the bolometric (blue) and XUV (red) luminosity for three different stellar masses (solid: $0.08 M_\odot$; dashed: $0.16 M_\odot$; dash-dotted: $0.24 M_\odot$) as a function of time. Two XUV saturation times are shown: 1 Gyr (red) and 0.1 Gyr (grey). Note the several order-of-magnitude drop in the first ~ 1 Gyr for the latest M dwarfs.

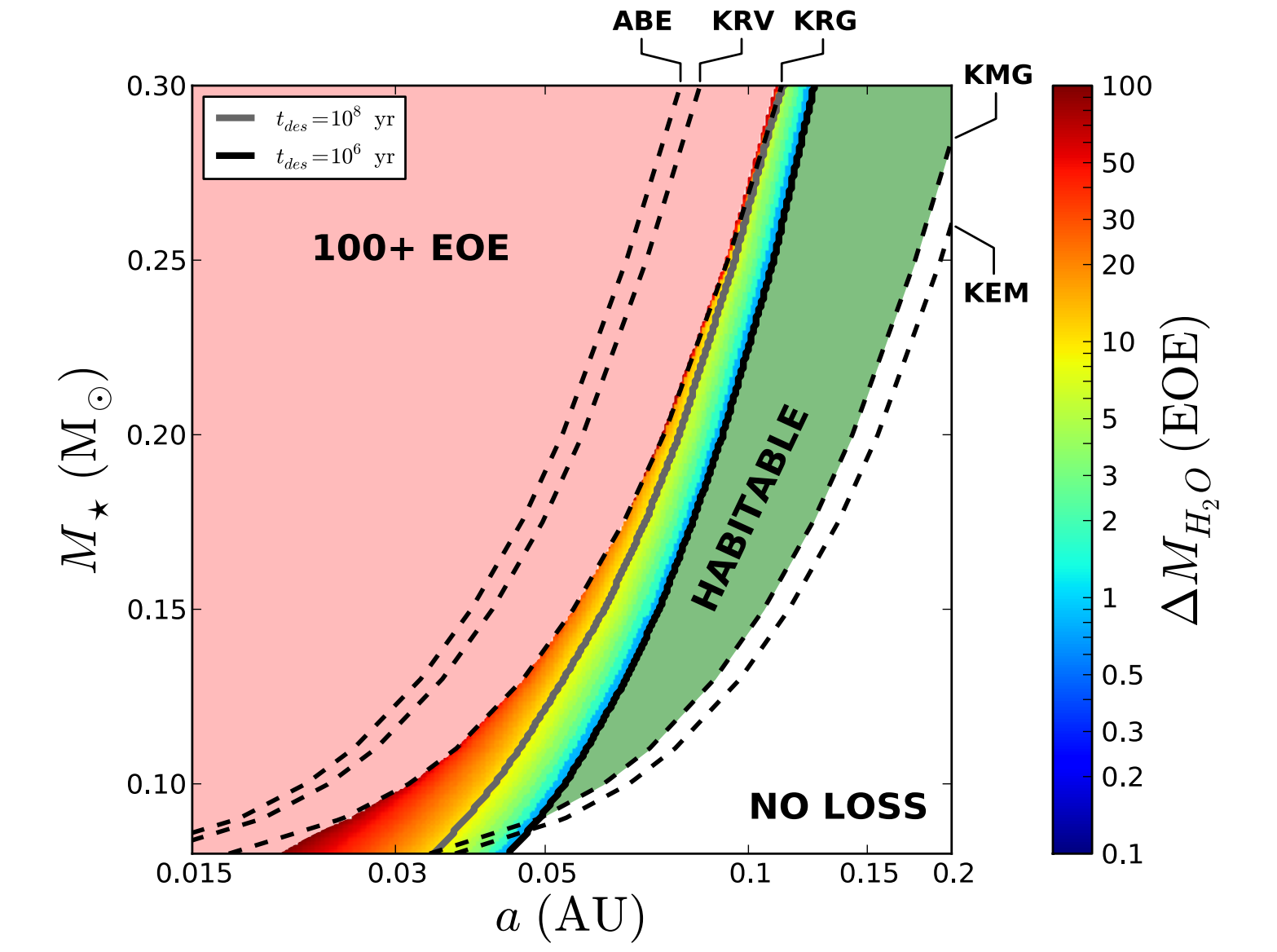


Figure 4: Mirage Earths? Because of the long contraction phase of late M dwarfs, Earths formed in situ in the HZ may lose several Earth Ocean Equivalents (EOE) of water in the first ~ 100 Myr, rendering them dry and uninhabitable. The latest M dwarfs may not have a HZ for Earths formed in situ.

Figure 1: Regions of parameter space that may be populated by HECs. Terrestrial planets detected today occupying the space to the *left* of each contour line could be the evaporated cores of gaseous planets. Planets detected to the *right* of the contour lines have always been terrestrial/gaseous. Dark red lines correspond to mass loss dominated by radiation/recombination-limited escape, while dark blue lines correspond to the more vigorous energy-limited mechanism. Which mechanism operates will depend on the ratio of stellar EUV to X-ray luminosity⁹. Different line styles correspond to different eccentricities today. Terrestrial planets detected at higher eccentricity (dashed and dash-dotted lines) could be evaporated cores at slightly larger orbital separations than planets detected on circular orbits (solid lines). In the energy-limited regime, all $1 M_\oplus$ terrestrial planets in the HZ of low-mass M dwarfs could be habitable evaporated cores. **(a): The default run**, for $M_{core} = 1 M_\oplus$ and $M_H \leq 1 M_\oplus$. **(b): Weak XUV run**, for $t_{sat} = 0.1$ Gyr and $\epsilon_{XUV} = 0.15$. **(c): Super-Earth**, with $M_{core} = 2 M_\oplus$ and $M_H \leq 2 M_\oplus$.