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New parametrization of the Second-order Reduced Density Matrix (2-RDDM)

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1 Introduction

Functionals of one-electron reduced density matrix (1-RDDM), γ , defined for an N -electron wavefunction, Ψ , as

$$\gamma(\mathbf{x}, \mathbf{x}') = N \int \cdots \int \Psi(\mathbf{x}, \mathbf{x}_2, \dots, \mathbf{x}_N) \Psi^*(\mathbf{x}', \mathbf{x}_2, \dots, \mathbf{x}_N) d\mathbf{x}_2 \cdots d\mathbf{x}_N, \quad (1)$$

where $\mathbf{x} = (\mathbf{r}, s)$ is a combined spatial and spin coordinate.

As immediate advantage of using 1-RDDM instead of the electron density, ρ , is that the kinetic energy is an explicit functional of γ but not of ρ . Then, there is no need to introduce a fictitious noninteracting system. Moreover, orbitals present in RDMFT are fractionally occupied, so functionals of γ seem to be better suited to account for static correlation.

Self-adjointness of γ , as defined in eq. (1), allows for its spectral representation [1]

$$\gamma(\mathbf{x}, \mathbf{x}') = \sum_p n_p \varphi_p(\mathbf{x}) \varphi_p^*(\mathbf{x}'). \quad (2)$$

The eigenvalues of 1-RDM are called natural occupation numbers, $\{n_p\}$, and the eigenfunctions are known as natural spinorbitals, $\{\varphi_p\}$. For convention, the indices p, q, r, s are referred to natural spinorbitals and a, b, c, d to arbitrary one-electron functions.

References

1. Löwdin, PO [Quantum Theory of Many-Particle Systems. I. Physical Interpretations by Means of Density Matrices, Natural Spin-Orbitals, and Convergence Problems in the Method of Configurational Interaction](#). *Physical Review* **1955**, *97*, 1474–1489.