

Worksheet 4.

Mathematical Foundations of Quantum Mechanics

José Antonio QUIÑONERO GRIS

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Statement: Calculate the Clebsh-Gordan coefficients corresponding to the composition of the angular momenta $l_1 = 1$ and $l_2 = 2$ by using the step ladder angular momentum operators.

Solution: Given the step ladder angular momentum operators

$$\hat{L}_+ = \hat{L}_x + i\hat{L}_y, \quad \hat{L}_- = \hat{L}_x - i\hat{L}_y, \quad (1)$$

that act on the eigenfunctions of \hat{L}^2 , the spherical harmonics $|Y_{lm}\rangle \equiv |l, m\rangle$, as

$$\hat{L}_+ |l, m\rangle = \sqrt{l(l+1) - m(m+1)}\hbar |l, m+1\rangle, \quad (2)$$

$$\hat{L}_- |l, m\rangle = \sqrt{l(l+1) - m(m-1)}\hbar |l, m-1\rangle, \quad (3)$$

and can be expressed, for each electron, as

$$\hat{L}_+ = \hat{L}_{1+} + \hat{L}_{2+}, \quad \hat{L}_- = \hat{L}_{1-} + \hat{L}_{2-}. \quad (4)$$

The possible values of the composition of the angular momenta l and m_l for the first electron, $|l, m\rangle_1$, and for the second electron, $|l, m\rangle_2$, are

$$l_1 = 2 \begin{cases} m_{l_1} = +2 & |2, 2\rangle_1 \\ m_{l_1} = +1 & |2, 1\rangle_1 \\ m_{l_1} = 0 & |2, 0\rangle_1 \\ m_{l_1} = -1 & |2, -1\rangle_1 \\ m_{l_1} = -2 & |2, -2\rangle_1 \end{cases} ; \quad l_2 = 1 \begin{cases} m_{l_2} = +1 & |1, 1\rangle_2 \\ m_{l_2} = 0 & |1, 0\rangle_2 \\ m_{l_2} = -1 & |1, -1\rangle_2 \end{cases} . \quad (5)$$

Then, the total angular momentum $L = l_1 + l_2, \dots, |l_1 - l_2|$, is

$$L = 3, 2, 1. \quad (6)$$

Now, for each component of the total angular momentum, the possible combinations of $|L, M\rangle$ are

$$L = 3 \begin{cases} M = +3 & |3, 3\rangle \\ M = +2 & |3, 2\rangle \\ M = +1 & |3, 1\rangle \\ M = 0 & |3, 0\rangle \\ M = -1 & |3, -1\rangle \\ M = -2 & |3, -2\rangle \\ M = -3 & |3, -3\rangle \end{cases} ; \quad L = 2 \begin{cases} M = +2 & |2, 2\rangle \\ M = +1 & |2, 1\rangle \\ M = 0 & |2, 0\rangle \\ M = -1 & |2, -1\rangle \\ M = -2 & |2, -2\rangle \end{cases} ; \quad L = 1 \begin{cases} M = +1 & |1, 1\rangle \\ M = 0 & |1, 0\rangle \\ M = -1 & |1, -1\rangle \end{cases} . \quad (7)$$

The Clebsh-Gordan coefficients must be calculated for each of the combinations above.

1 $L = 3$

Starting with the $L = 3, M = 3$ state, there is only one possible combination of states for electron 1 and electron 2, and that is

$$\boxed{|3, 3\rangle = |2, 2\rangle_1 |1, 1\rangle_2} . \quad (8)$$

(Note: the results are summarized at the end of the document)

In order to construct the rest of combinations for $L = 3$, we can apply the lower operator \hat{L}_- [eq. (3)] to the state $|3, 3\rangle$ to get $|3, 2\rangle$

$$\hat{L}_- |3, 3\rangle = \sqrt{3(3+1) - 3(3-1)}\hbar |3, 2\rangle = \sqrt{6}\hbar |3, 2\rangle . \quad (= \text{LHS}) \quad (9)$$

The ladder operator must be also applied to the right hand side (RHS) of the eq. (8). Considering the separation of the ladder operator as [eq. (4)]:

$$\hat{L}_- = \hat{L}_{1-} + \hat{L}_{2-}, \quad (10)$$

then

$$\hat{L}_- |2, 2\rangle_1 |1, 1\rangle_2 = (\hat{L}_{1-} + \hat{L}_{2-}) |2, 2\rangle_1 |1, 1\rangle_2 = \hat{L}_{1-} |2, 2\rangle_1 |1, 1\rangle_2 + \hat{L}_{2-} |2, 2\rangle_1 |1, 1\rangle_2, \quad (11)$$

and the fact that each component of the ladder operator acts only on the corresponding electron

$$(\hat{L}_{1-} |2, 2\rangle_1) |1, 1\rangle_2 = (\sqrt{2(2+1) - 2(2-1)}\hbar |2, 1\rangle_1) |1, 1\rangle_2 = 2\hbar |2, 1\rangle_1 |1, 1\rangle_2, \quad (12)$$

$$(\hat{L}_{2-} |1, 1\rangle_2) |2, 2\rangle_1 = (\sqrt{1(1+1) - 1(1-1)}\hbar |1, 0\rangle_2) |2, 2\rangle_1 = \sqrt{2}\hbar |2, 2\rangle_1 |1, 0\rangle_2, \quad (13)$$

then

$$\hat{L}_- |2, 2\rangle_1 |1, 1\rangle_2 = 2\hbar |2, 1\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |2, 2\rangle_1 |1, 0\rangle_2 . \quad (= \text{RHS}) \quad (14)$$

Therefore, equating LHS = RHS [equations eq. (9) and eq. (14)]

$$\sqrt{6}\hbar |3, 2\rangle = 2\hbar |2, 1\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |2, 2\rangle_1 |1, 0\rangle_2, \quad (15)$$

and factorizing $\sqrt{2}\hbar$ on the RHS

$$\sqrt{6}\hbar |3, 2\rangle = \sqrt{2}\hbar \left(\sqrt{2} |2, 1\rangle_1 |1, 1\rangle_2 + |2, 2\rangle_1 |1, 0\rangle_2 \right), \quad (16)$$

then

$$\boxed{|3, 2\rangle = \frac{1}{\sqrt{3}} \left(|2, 2\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 1\rangle_1 |1, 1\rangle_2 \right)}. \quad (17)$$

Just as before, the down operator \hat{L}_- can be applied again over $|3, 2\rangle$ to get $|3, 1\rangle$

$$\hat{L}_- |3, 2\rangle = \sqrt{3(3+1) - 2(2-1)}\hbar |3, 1\rangle = \sqrt{10}\hbar |3, 1\rangle. \quad (= \text{LHS}) \quad (18)$$

and to the RHS

$$\hat{L}_- \underbrace{\left\{ \frac{1}{\sqrt{3}} \left(\sqrt{2} |2, 1\rangle_1 |1, 1\rangle_2 + |2, 2\rangle_1 |1, 0\rangle_2 \right) \right\}}_{\{\text{RHS}\}} = \sqrt{\frac{1}{3}} \left(\sqrt{2} \underbrace{\hat{L}_- |2, 1\rangle_1 |1, 1\rangle_2}_{(*)} + \underbrace{\hat{L}_- |2, 2\rangle_1 |1, 0\rangle_2}_{(**)} \right). \quad (19)$$

The result of the operations are calculated separately

$$\begin{aligned} (*) \hat{L}_- |2, 1\rangle_1 |1, 1\rangle_2 &= \left(\hat{L}_{1-} |2, 1\rangle_1 \right) |1, 1\rangle_2 + \left(\hat{L}_{2-} |1, 1\rangle_2 \right) |2, 1\rangle_1 = \\ &= \sqrt{2(2+1) - 1(1-1)}\hbar |2, 0\rangle_1 |1, 1\rangle_2 + \sqrt{1(1+1) - 1(1-1)}\hbar |1, 0\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{6}\hbar |2, 0\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |1, 0\rangle_2 |2, 1\rangle_1, \end{aligned} \quad (20)$$

and

$$\begin{aligned} (**) \hat{L}_- |2, 2\rangle_1 |1, 0\rangle_2 &= \left(\hat{L}_{1-} |2, 2\rangle_1 \right) |1, 0\rangle_2 + \left(\hat{L}_{2-} |1, 0\rangle_2 \right) |2, 2\rangle_1 = \\ &= \sqrt{2(2+1) - 2(2-1)}\hbar |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{1(1+1) - 0(0-1)}\hbar |1, -1\rangle_2 |2, 2\rangle_1 = \\ &= 2\hbar |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, -1\rangle_2 |2, 2\rangle_1. \end{aligned} \quad (21)$$

In eq. (19)

$$\begin{aligned} \hat{L}_- \{\text{RHS}\} &= \frac{\hbar}{\sqrt{3}} \left\{ \sqrt{2} \left(\sqrt{6} |2, 0\rangle_1 |1, 1\rangle_2 + \sqrt{2} |1, 0\rangle_2 |2, 1\rangle_1 \right) + \right. \\ &\quad \left. + 2 |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, -1\rangle_2 |2, 2\rangle_1 \right\}. \end{aligned} \quad (22)$$

It can be rewritten as

$$\begin{aligned} \hat{L}_- \{\text{RHS}\} &= \frac{\hbar}{\sqrt{3}} \left(2\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 + 2 |2, 1\rangle_1 |1, 0\rangle_2 + 2 |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 \right) = \\ &= \frac{\hbar}{\sqrt{3}} \left(2\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 + 4 |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 \right). \end{aligned} \quad (23)$$

Then, with eq. (18)

$$\sqrt{10}\hbar |3, 1\rangle = \frac{\hbar}{\sqrt{3}} \left(2\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 + 4 |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 \right), \quad (24)$$

therefore

$$\boxed{|3, 1\rangle = \sqrt{\frac{1}{30}} \left(\sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 + 4 |2, 1\rangle_1 |1, 0\rangle_2 + 2\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right).} \quad (25)$$

Finally, applying again \hat{L}_- to $|3, 1\rangle$ to obtain $|3, 0\rangle$

$$\hat{L}_- |3, 1\rangle = \sqrt{3(3+1) - 1(1-1)}\hbar |3, 0\rangle = \sqrt{12}\hbar |3, 0\rangle = 2\sqrt{3}\hbar |3, 0\rangle, \quad (= \text{LHS}) \quad (26)$$

and to the RHS of the eq. (25)

$$\hat{L}_- \{\text{RHS}\} = \sqrt{\frac{1}{30}} \left(2\sqrt{3} \underbrace{\hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2}_{(*)} + 4 \underbrace{\hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2}_{(**)} + \sqrt{2} \underbrace{\hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2}_{(***)} \right) \quad (27)$$

Calculating separately $(*)$

$$\begin{aligned} (*) \quad \hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2 &= \hat{L}_{1-} |2, 0\rangle_1 |1, 1\rangle_2 + \hat{L}_{2-} |1, 1\rangle_2 |2, 0\rangle_1 = \\ &= \sqrt{6}\hbar |2, -1\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |2, 0\rangle_1 |1, 0\rangle_2, \end{aligned} \quad (28)$$

and $(**)$

$$\begin{aligned} (**) \quad \hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2 &= \hat{L}_{1-} |2, 1\rangle_1 |1, 0\rangle_2 + \hat{L}_{2-} |1, 0\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{6}\hbar |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |2, 1\rangle_1 |1, -1\rangle_2, \end{aligned} \quad (29)$$

and, lastly, $(***)$

$$\begin{aligned} (***) \quad \hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2 &= \hat{L}_{1-} |2, 2\rangle_1 |1, -1\rangle_2 + \hat{L}_{2-} |1, -1\rangle_2 |2, 2\rangle_1 = \\ &= 2\hbar |2, 1\rangle_1 |1, -1\rangle_2. \end{aligned} \quad (30)$$

Then, in eq. (27)

$$\begin{aligned} \hat{L}_- \{\text{RHS}\} &= \sqrt{\frac{1}{30}} \left\{ 2\sqrt{3} \left(\sqrt{6}\hbar |2, -1\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |2, 0\rangle_1 |1, 0\rangle_2 \right) + \right. \\ &\quad + 4 \left(\sqrt{6}\hbar |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |2, 1\rangle_1 |1, -1\rangle_2 \right) + \\ &\quad \left. + \sqrt{2} (2\hbar |2, 1\rangle_1 |1, -1\rangle_2) \right\}, \end{aligned} \quad (31)$$

that, factorizing, can be rewritten as

$$\hat{L}_- \{\text{RHS}\} = 6\hbar \sqrt{\frac{1}{15}} \left(|2, -1\rangle_1 |1, 1\rangle_2 + \sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, 1\rangle_1 |1, -1\rangle_2 \right). \quad (32)$$

Doing LHS = RHS with eq. (26)

$$2\sqrt{3}\hbar |3, 0\rangle = 6\hbar \sqrt{\frac{1}{15}} \left(|2, -1\rangle_1 |1, 1\rangle_2 + \sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, 1\rangle_1 |1, -1\rangle_2 \right), \quad (33)$$

reads

$$\boxed{|3, 0\rangle = \sqrt{\frac{1}{5}} \left(|2, 1\rangle_1 |1, -1\rangle_2 + \sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, -1\rangle_1 |1, 1\rangle_2 \right).} \quad (34)$$

Now, the rest of combinations for $L = 3$ can be calculated analogously starting with $|3, -3\rangle$ and applying the rising ladder operator \hat{L}_+ [eq. (2)]. The only combination for $|3, -3\rangle$ is

$$\boxed{|3, -3\rangle = |2, -2\rangle_1 |1, -1\rangle_2.} \quad (35)$$

Applying \hat{L}_+ to the LHS

$$\hat{L}_+ |3, -3\rangle = \sqrt{3(3+1) - (-3)((-3)+1)}\hbar |3, -2\rangle = \sqrt{6}\hbar |3, -2\rangle, \quad (= \text{LHS}) \quad (36)$$

and to the RHS

$$\hat{L}_+ |2, -2\rangle_1 |1, -1\rangle_2 = \hat{L}_{1-} |2, -2\rangle_1 |1, -1\rangle_2 + \hat{L}_{2-} |1, -1\rangle_2 |2, -2\rangle_1. \quad (37)$$

As

$$\hat{L}_{1-} |2, -2\rangle_1 = \sqrt{2(2+1) - (-2)((-2)+1)}\hbar |2, -1\rangle_1 = 2\hbar |2, -1\rangle_1, \quad (38)$$

$$\hat{L}_{2-} |1, -1\rangle_2 = \sqrt{1(1+1) - (-1)((-1)+1)}\hbar |1, 0\rangle_2 = \sqrt{2}\hbar |1, 0\rangle_2, \quad (39)$$

then

$$\hat{L}_+ |2, -2\rangle_1 |1, -1\rangle_2 = 2\hbar |2, -1\rangle_1 |1, -1\rangle_2 + \sqrt{2}\hbar |2, -2\rangle_1 |1, 0\rangle_2. \quad (= \text{RHS}) \quad (40)$$

Equating LHS = RHS

$$\begin{aligned} \sqrt{6}\hbar |3, -2\rangle &= 2\hbar |2, -1\rangle_1 |1, -1\rangle_2 + \sqrt{2}\hbar |2, -2\rangle_1 |1, 0\rangle_2 = \\ &= \sqrt{2}\hbar \left(\sqrt{2} |2, -1\rangle_1 |1, -1\rangle_2 + |2, -2\rangle_1 |1, 0\rangle_2 \right), \end{aligned} \quad (41)$$

$|3, -2\rangle$ is found and reads

$$\boxed{|3, -2\rangle = \sqrt{\frac{1}{3}} \left(\sqrt{2} |2, -1\rangle_1 |1, -1\rangle_2 + |2, -2\rangle_1 |1, 0\rangle_2 \right).} \quad (42)$$

Now, \hat{L}_+ can be applied again over $|3, -2\rangle$

$$\hat{L}_+ |3, -2\rangle = \sqrt{3(3+1) - (-2)((-2)+1)}\hbar |3, -1\rangle = \sqrt{10}\hbar |3, -1\rangle, \quad (= \text{LHS}) \quad (43)$$

and for the RHS

$$\hat{L}_+ \{\text{RHS}\} = \sqrt{\frac{1}{3}} \left(\sqrt{2}\hat{L}_+ |2, -1\rangle_1 |1, -1\rangle_2 + \hat{L}_+ |2, -2\rangle_1 |1, 0\rangle_2 \right), \quad (44)$$

as

$$\begin{aligned}
\hat{L}_+ |2, -1\rangle_1 |1, -1\rangle_2 &= \hat{L}_{1-} |2, -1\rangle_1 |1, -1\rangle_2 + \hat{L}_{2-} |1, -1\rangle_2 |2, -1\rangle_1 = \\
&= \sqrt{6}\hbar |2, 0\rangle_1 |1, -1\rangle_2 + \sqrt{2}\hbar |1, 0\rangle_2 |2, -1\rangle_1 \\
&= \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 \right), \tag{45}
\end{aligned}$$

and

$$\begin{aligned}
\hat{L}_+ |2, -2\rangle_1 |1, 0\rangle_2 &= \hat{L}_{1-} |2, -2\rangle_1 |1, 0\rangle_2 + \hat{L}_{2-} |1, 0\rangle_2 |2, -2\rangle_1 = \\
&= 2\hbar |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, 1\rangle_2 |2, -2\rangle_1 \\
&= \sqrt{2}\hbar \left(\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2 \right), \tag{46}
\end{aligned}$$

then

$$\begin{aligned}
\hat{L}_+ \{\text{RHS}\} &= \sqrt{\frac{1}{3}} \left\{ \sqrt{2}\sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 \right) + \right. \\
&\quad \left. + \sqrt{2}\hbar \left(\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2 \right) \right\} \tag{47}
\end{aligned}$$

$$= \sqrt{\frac{1}{3}}\hbar \left(2\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + 4 |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right). \quad (= \text{RHS}) \tag{48}$$

Doing LHS = RHS

$$\sqrt{10}\hbar |3, -1\rangle = \sqrt{\frac{1}{3}}\hbar \left(2\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + 4 |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right), \tag{49}$$

$|3, -1\rangle$ reads

$$|3, -1\rangle = \sqrt{\frac{1}{30}} \left(2\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + 4 |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right), \tag{50}$$

or

$$\boxed{|3, -1\rangle = \sqrt{\frac{1}{15}} \left(\sqrt{6} |2, 0\rangle_1 |1, -1\rangle_2 + 2\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2 \right).} \tag{51}$$

2 $L = 2$

Now, the ladder operators can't be applied ver the already found combinations, since it is not the same value of L . Therefore, the new states for $L = 2$ are found imposing orthogonality conditions

$$\langle L, M_L | L', M_{L'} \rangle = 0 \quad L' \neq L. \tag{52}$$

Then, $|2, 2\rangle$ can be found imposing

$$\langle 3, 2 | 2, 2 \rangle = 0. \tag{53}$$

Rewriting $|3, 2\rangle$ from eq. (17) as

$$|3, 2\rangle = \sqrt{\frac{1}{3}} \left(\sqrt{2} |2, 1\rangle_1 |1, 1\rangle_2 + |2, 2\rangle_1 |1, 0\rangle_2 \right) = \sqrt{\frac{1}{3}} \left(\sqrt{2} |\psi_1\rangle + |\psi_2\rangle \right), \quad (54)$$

$|2, 2\rangle$ can be written as

$$|2, 2\rangle = c_1 |\psi_1\rangle + c_2 |\psi_2\rangle. \quad (55)$$

Imposing now the orthogonality condition from eq. (53)

$$\langle 3, 2 | 2, 2 \rangle = \sqrt{\frac{1}{3}} \left(\sqrt{2} \langle \psi_1 | + \langle \psi_2 | \right) (c_1 |\psi_1\rangle + c_2 |\psi_2\rangle) = 0, \quad (56)$$

and operating considering that $\langle \psi_i | \psi_j \rangle = \delta_{ij}$

$$\begin{aligned} \langle 3, 2 | 2, 2 \rangle &= \sqrt{\frac{1}{3}} \left(\sqrt{2} c_1 \langle \psi_1 | \psi_1 \rangle + c_1 \langle \psi_2 | \psi_1 \rangle + \sqrt{2} c_2 \langle \psi_1 | \psi_2 \rangle + c_2 \langle \psi_2 | \psi_2 \rangle \right) = \\ &= \sqrt{\frac{1}{3}} \left(\sqrt{2} c_1 + c_2 \right), \end{aligned} \quad (57)$$

then

$$\sqrt{\frac{1}{3}} \left(\sqrt{2} c_1 + c_2 \right) = 0 \implies c_1 = -\frac{1}{\sqrt{2}} c_2. \quad (58)$$

Imposing now the normalisation condition

$$\sum_{i=1} |c_i|^2 = 1. \quad (59)$$

Then

$$|c_1|^2 + |c_2|^2 = 1, \quad (60)$$

substituting c_1 from eq. (58)

$$\left| -\frac{1}{\sqrt{2}} c_2 \right|^2 + |c_2|^2 = \frac{3}{2} |c_2|^2 = 1 \implies c_2 = \sqrt{\frac{2}{3}}, \quad (61)$$

and

$$c_1 = -\frac{1}{\sqrt{2}} c_2 = -\frac{1}{\sqrt{2}} \sqrt{\frac{2}{3}} = -\frac{1}{\sqrt{3}} \quad (62)$$

Then, eq. (55) can be written as

$$|2, 2\rangle = \frac{1}{\sqrt{3}} \left(-|\psi_1\rangle + \sqrt{2} |\psi_2\rangle \right) = \frac{1}{\sqrt{3}} \left(-|2, 1\rangle_1 |1, 1\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, 0\rangle_2 \right). \quad (63)$$

Therefore, $|2, 2\rangle$ reads

$$\boxed{|2, 2\rangle = \sqrt{\frac{1}{3}} \left(\sqrt{2} |2, 2\rangle_1 |1, 0\rangle_2 - |2, 1\rangle_1 |1, 1\rangle_2 \right).} \quad (64)$$

Now, the ladder operator can be applied to get the rest of combinations for $L = 2$. Applying \hat{L}_-

over $|2, 2\rangle$ [eq. (64)]

$$\hat{L}_- |2, 2\rangle = \sqrt{2(2+1) - 2(2-1)}\hbar |2, 1\rangle = 2\hbar |2, 1\rangle, \quad (= \text{LHS}) \quad (65)$$

and to the RHS

$$\hat{L}_- \{\text{RHS}\} = \sqrt{\frac{1}{3}} \left(\sqrt{2}\hat{L}_- |2, 2\rangle_1 |1, 0\rangle_2 - \hat{L}_- |2, 1\rangle_1 |1, 1\rangle_2 \right). \quad (66)$$

Calculating separately $\hat{L}_- |2, 2\rangle_1 |1, 0\rangle_2$

$$\begin{aligned} \hat{L}_- |2, 2\rangle_1 |1, 0\rangle_2 &= \hat{L}_{1-} |2, 2\rangle_1 |1, 0\rangle_2 + \hat{L}_{2-} |1, 0\rangle_2 |2, 2\rangle_1 = \\ &= 2\hbar |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, -1\rangle_2 |2, 2\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{2} |2, 1\rangle_1 |1, 0\rangle_2 + |2, 2\rangle_1 |1, -1\rangle_2 \right), \end{aligned} \quad (67)$$

and $\hat{L}_- |2, 1\rangle_1 |1, 1\rangle_2$

$$\begin{aligned} \hat{L}_- |2, 1\rangle_1 |1, 1\rangle_2 &= \hat{L}_{1-} |2, 1\rangle_1 |1, 1\rangle_2 + \hat{L}_{2-} |1, 1\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{6}\hbar |2, 0\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |1, 0\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 + |2, 1\rangle_1 |1, 0\rangle_2 \right), \end{aligned} \quad (68)$$

then, on eq. (66)

$$\begin{aligned} \hat{L}_- \{\text{RHS}\} &= \sqrt{\frac{1}{3}} \left\{ \sqrt{2} \left[\sqrt{2}\hbar \left(\sqrt{2} |2, 1\rangle_1 |1, 0\rangle_2 + |2, 2\rangle_1 |1, -1\rangle_2 \right) \right] \right. \\ &\quad \left. - \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 + |2, 1\rangle_1 |1, 0\rangle_2 \right) \right\} = \\ &= \sqrt{\frac{2}{3}}\hbar \left\{ |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right\}. \quad (= \text{RHS}) \end{aligned} \quad (69)$$

Equating LHS = RHS with eq. (65) and eq. (69)

$$2\hbar |2, 1\rangle = \sqrt{\frac{2}{3}}\hbar \left(|2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right), \quad (70)$$

then

$$|2, 1\rangle = \sqrt{\frac{1}{6}} \left(|2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right). \quad (71)$$

Therefore, $|2, 1\rangle$ reads

$$\boxed{|2, 1\rangle = \sqrt{\frac{1}{6}} \left(\sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 + |2, 1\rangle_1 |1, 0\rangle_2 - \sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right).} \quad (72)$$

As before, applying \hat{L}_- over $|2, 1\rangle$

$$\hat{L}_- |2, 1\rangle = \sqrt{2(2+1) - 1(1-1)}\hbar |2, 0\rangle = \sqrt{6}\hbar |2, 0\rangle, \quad (= \text{LHS}) \quad (73)$$

and to the RHS

$$\hat{L}_- \{\text{RHS}\} = \sqrt{\frac{1}{6}} \left(\sqrt{2} \hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2 + \hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2 - \sqrt{3} \hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2 \right). \quad (74)$$

Calculating separately $\hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2$

$$\begin{aligned} \hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2 &= \hat{L}_{1-} |2, 2\rangle_1 |1, -1\rangle_2 + \hat{L}_{2-} |1, -1\rangle_2 |2, 2\rangle_1 = \\ &= 2\hbar |2, 1\rangle_1 |1, -1\rangle_2, \end{aligned} \quad (75)$$

and $\hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2$

$$\begin{aligned} \hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2 &= \hat{L}_{1-} |2, 1\rangle_1 |1, 0\rangle_2 + \hat{L}_{2-} |1, 0\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{6}\hbar |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, -1\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, 1\rangle_1 |1, -1\rangle_2 \right), \end{aligned} \quad (76)$$

and, lastly, $\hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2$

$$\begin{aligned} \hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2 &= \hat{L}_{1-} |2, 0\rangle_1 |1, 1\rangle_2 + \hat{L}_{2-} |1, 1\rangle_2 |2, 0\rangle_1 = \\ &= \sqrt{6}\hbar |2, -1\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |1, 0\rangle_2 |2, 0\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 + |2, 0\rangle_1 |1, 0\rangle_2 \right). \end{aligned} \quad (77)$$

Then, in eq. (74)

$$\begin{aligned} \hat{L}_- \{\text{RHS}\} &= \sqrt{\frac{1}{6}} \left\{ \sqrt{2} (2\hbar |2, 1\rangle_1 |1, -1\rangle_2) \right. \\ &\quad + \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, 1\rangle_1 |1, -1\rangle_2 \right) \\ &\quad \left. - \sqrt{3} \left[\sqrt{2}\hbar \left(\sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 + |2, 0\rangle_1 |1, 0\rangle_2 \right) \right] \right\} = \\ &= \sqrt{3}\hbar (|2, 1\rangle_1 |1, -1\rangle_2 - |2, -1\rangle_1 |1, 1\rangle_2). \quad (= \text{RHS}) \end{aligned} \quad (78)$$

Equating LHS = RHS with eq. (73) and eq. (78)

$$\sqrt{6}\hbar |2, 0\rangle = \sqrt{3}\hbar (|2, 1\rangle_1 |1, -1\rangle_2 - |2, -1\rangle_1 |1, 1\rangle_2). \quad (79)$$

Then, $|2, 0\rangle$ reads

$$\boxed{|2, 0\rangle = \sqrt{\frac{1}{2}} (|2, 1\rangle_1 |1, -1\rangle_2 - |2, -1\rangle_1 |1, 1\rangle_2).} \quad (80)$$

Now, instead applying again \hat{L}_- to $|2, 0\rangle$, it is simpler to apply \hat{L}_+ to $|2, -2\rangle$. Imposing the orthogonality condition [eq. (52)] between $|3, -2\rangle$ and $|2, -2\rangle$

$$\langle 3, -2 | 2, -2 \rangle = 0, \quad (81)$$

rewriting $|3, -2\rangle$ as

$$|3, -2\rangle = \sqrt{\frac{1}{3}} \left(\sqrt{2} |2, -1\rangle_1 |1, -1\rangle_2 + |2, -2\rangle_1 |1, 0\rangle_2 \right) = \sqrt{\frac{1}{3}} \left(\sqrt{2} |\psi_1\rangle + |\psi_2\rangle \right), \quad (82)$$

and writing $|2, -2\rangle$ as

$$|2, -2\rangle = c_1 |\psi_1\rangle + c_2 |\psi_2\rangle. \quad (83)$$

In eq. (81)

$$\begin{aligned} \langle 3, -2 | 2, -2 \rangle &= \sqrt{\frac{1}{3}} \left[\left(\sqrt{2} \langle \psi_1 | + \langle \psi_2 | \right) (c_1 |\psi_1\rangle + c_2 |\psi_2\rangle) \right] = \\ &= \sqrt{\frac{1}{3}} \left[\sqrt{2} c_1 \langle \psi_1 | \psi_1 \rangle + \sqrt{2} c_2 \langle \psi_1 | \psi_2 \rangle + c_1 \langle \psi_2 | \psi_1 \rangle + c_2 \langle \psi_2 | \psi_2 \rangle \right] = \\ &= \sqrt{\frac{1}{3}} \left(\sqrt{2} c_1 + c_2 \right), \end{aligned} \quad (84)$$

then

$$\sqrt{\frac{1}{3}} \left(\sqrt{2} c_1 + c_2 \right) = 0 \implies \sqrt{2} c_1 + c_2 = 0 \implies c_1 = -\frac{1}{\sqrt{2}} c_2. \quad (85)$$

Imposing also the normalisation condition [eq. (59)]

$$|c_1|^2 + |c_2|^2 = 1, \quad (86)$$

and substituting c_1 from eq. (85)

$$\left| -\frac{1}{\sqrt{2}} c_2 \right|^2 + |c_2|^2 = 1 \implies \left(\frac{1}{2} + 1 \right) |c_2|^2 = 1 \implies c_2 = \sqrt{\frac{2}{3}}. \quad (87)$$

Now, on eq. (85)

$$c_1 = -\frac{1}{\sqrt{2}} c_2 = -\frac{1}{\sqrt{2}} \sqrt{\frac{2}{3}} = -\sqrt{\frac{1}{3}}. \quad (88)$$

Therefore, in eq. (83)

$$|2, -2\rangle = -\sqrt{\frac{1}{3}} |\psi_1\rangle + \sqrt{\frac{2}{3}} |\psi_2\rangle, \quad (89)$$

$|2, -2\rangle$ reads

$$\boxed{|2, -2\rangle = \sqrt{\frac{1}{3}} \left(-|2, -1\rangle_1 |1, -1\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 0\rangle_2 \right)}. \quad (90)$$

Applying now \hat{L}_+ to $|2, -2\rangle$

$$\hat{L}_+ |2, -2\rangle = \sqrt{2(2+1) - (-2)((-2)+1)} \hbar |2, -1\rangle = 2\hbar |2, -1\rangle, \quad (= \text{LHS}) \quad (91)$$

and to the RHS

$$\hat{L}_+ \{\text{RHS}\} = \sqrt{\frac{1}{3}} \left(-\hat{L}_+ |2, -1\rangle_1 |1, -1\rangle_2 + \sqrt{2} \hat{L}_+ |2, -2\rangle_1 |1, 0\rangle_2 \right). \quad (92)$$

Calculating separately $\hat{L}_+ |2, -1\rangle_1 |1, -1\rangle_2$

$$\begin{aligned}\hat{L}_+ |2, -1\rangle_1 |1, -1\rangle_2 &= \hat{L}_{1-} |2, -1\rangle_1 |1, -1\rangle_2 + \hat{L}_{2-} |1, -1\rangle_2 |2, -1\rangle_1 = \\ &= \sqrt{6}\hbar |2, 0\rangle_1 |1, -1\rangle_2 + \sqrt{2}\hbar |1, 0\rangle_2 |2, -1\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 \right),\end{aligned}\quad (93)$$

and $\hat{L}_+ |2, -2\rangle_1 |1, 0\rangle_2$

$$\begin{aligned}\hat{L}_+ |2, -2\rangle_1 |1, 0\rangle_2 &= \hat{L}_{1-} |2, -2\rangle_1 |1, 0\rangle_2 + \hat{L}_{2-} |1, 0\rangle_2 |2, -2\rangle_1 = \\ &= 2\hbar |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, 1\rangle_2 |2, -2\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2 \right).\end{aligned}\quad (94)$$

Then, in eq. (92)

$$\begin{aligned}\hat{L}_+ \{\text{RHS}\} &= \sqrt{\frac{1}{3}} \left\{ -\sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 \right) \right. \\ &\quad \left. + \sqrt{2}\sqrt{2}\hbar \left(\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2 \right) \right\} \\ &= \sqrt{\frac{2}{3}}\hbar \left(-\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right). \quad (= \text{RHS})\end{aligned}\quad (95)$$

Doing LHS = RHS

$$2\hbar |2, -1\rangle = \sqrt{\frac{2}{3}}\hbar \left(-\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right), \quad (96)$$

$|2, -1\rangle$ reads

$$\boxed{|2, -1\rangle = \sqrt{\frac{1}{6}} \left(-\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right).} \quad (97)$$

3 $L = 1$

Now, for $L = 1$, $|1, 1\rangle$ can be constructed imposing orthogonality conditions as before, but now two orthogonality conditions are needed

$$\langle 3, 1 | 1, 1 \rangle = 0, \quad \langle 2, 1 | 1, 1 \rangle = 0. \quad (98)$$

Rewriting $|3, 1\rangle$ as

$$\begin{aligned}|3, 1\rangle &= \sqrt{\frac{1}{30}} \left(\sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 + 4 |2, 1\rangle_1 |1, 0\rangle_2 + 2\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right) = \\ &= \sqrt{\frac{1}{30}} \left(\sqrt{2} |\psi_1\rangle + 4 |\psi_2\rangle + 2\sqrt{3} |\psi_3\rangle \right),\end{aligned}\quad (99)$$

and $|2, 1\rangle$ as

$$\begin{aligned} |2, 1\rangle &= \sqrt{\frac{1}{6}} \left(\sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 + |2, 1\rangle_1 |1, 0\rangle_2 - \sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2 \right) = \\ &= \sqrt{\frac{1}{6}} \left(\sqrt{2} |\psi_1\rangle + |\psi_2\rangle - \sqrt{3} |\psi_3\rangle \right). \end{aligned} \quad (100)$$

$|1, 1\rangle$ can be written as

$$|1, 1\rangle = c_1 |\psi_1\rangle + c_2 |\psi_2\rangle + c_3 |\psi_3\rangle. \quad (101)$$

Imposing the first condition from eq. (98) the first equation (I) from the system of equations is found

$$\begin{aligned} \langle 3, 1 | 1, 1\rangle &= \sqrt{\frac{1}{30}} \left[\left(\sqrt{2} \langle \psi_1 | + 4 \langle \psi_2 | + 2\sqrt{3} \langle \psi_3 | \right) (c_1 |\psi_1\rangle + c_2 |\psi_2\rangle + c_3 |\psi_3\rangle) \right] = \\ &= \sqrt{\frac{1}{30}} \left(\sqrt{2} c_1 + 4c_2 + 2\sqrt{3} c_3 \right) = 0, \end{aligned} \quad (102)$$

then

$$\sqrt{2} c_1 + 4c_2 + 2\sqrt{3} c_3 = 0. \quad (\text{I}) \quad (103)$$

Repeating for the second condition from eq. (98)

$$\begin{aligned} \langle 2, 1 | 1, 1\rangle &= \sqrt{\frac{1}{6}} \left[\left(\sqrt{2} \langle \psi_1 | + \langle \psi_2 | - \sqrt{3} \langle \psi_3 | \right) (c_1 |\psi_1\rangle + c_2 |\psi_2\rangle + c_3 |\psi_3\rangle) \right] = \\ &= \sqrt{\frac{1}{6}} \left(\sqrt{2} c_1 + c_2 - \sqrt{3} c_3 \right) = 0, \end{aligned} \quad (104)$$

then

$$\sqrt{2} c_1 + c_2 - \sqrt{3} c_3 = 0. \quad (\text{II}) \quad (105)$$

The third equation is found from the normalisation condition eq. (59)

$$|c_1|^2 + |c_2|^2 + |c_3|^2 = 1. \quad (\text{III}) \quad (106)$$

Then, the system of equations to solve reads

$$\begin{cases} \sqrt{2} c_1 + 4c_2 + 2\sqrt{3} c_3 = 0 & (\text{I}) \\ \sqrt{2} c_1 + c_2 - \sqrt{3} c_3 = 0 & (\text{II}) \\ |c_1|^2 + |c_2|^2 + |c_3|^2 = 1 & (\text{III}) \end{cases} \quad (107)$$

Isolating c_1 from the first equation

$$c_1 = -\frac{1}{\sqrt{2}} \left(4c_2 + 2\sqrt{3} c_3 \right), \quad (108)$$

and c_2 from the second equation

$$c_2 = \sqrt{3} c_3 - \sqrt{2} c_1, \quad (109)$$

substituting the expression for c_1 given in eq. (108)

$$c_2 = \sqrt{3}c_3 - \sqrt{2} \left[-\frac{1}{\sqrt{2}} (4c_2 + 2\sqrt{3}c_3) \right] = 4c_2 + 3\sqrt{3}c_3 \implies c_2 = -\sqrt{3}c_3, \quad (110)$$

and substituting in eq. (108)

$$c_1 = -\frac{1}{\sqrt{2}} \left[4(-\sqrt{3}c_3) + 2\sqrt{3}c_3 \right] = \sqrt{6}c_3. \quad (111)$$

Isolating now c_3 from the third equation

$$\left| \sqrt{6}c_3 \right|^2 + \left| -\sqrt{3}c_3 \right|^2 + |c_3|^2 = 1 \implies 10|c_3|^2 = 1 \implies c_3 = \sqrt{\frac{1}{10}}. \quad (112)$$

Then, this expression for c_3 can be substituted in eqs. (110) and (111), as

$$c_1 = \sqrt{6}\sqrt{\frac{1}{10}} = \sqrt{\frac{6}{10}} = \sqrt{\frac{3}{5}}, \quad (113)$$

$$c_2 = -\sqrt{3}\sqrt{\frac{1}{10}} = -\sqrt{\frac{3}{10}}. \quad (114)$$

Substituting in eq. (101) for $|1, 1\rangle$

$$\begin{aligned} |1, 1\rangle &= \sqrt{\frac{3}{5}} |\psi_1\rangle - \sqrt{\frac{3}{10}} |\psi_2\rangle + \sqrt{\frac{1}{10}} |\psi_3\rangle \\ &= \sqrt{\frac{3}{5}} |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{\frac{3}{10}} |2, 1\rangle_1 |1, 0\rangle_2 + \sqrt{\frac{1}{10}} |2, 0\rangle_1 |1, 1\rangle_2. \end{aligned} \quad (115)$$

Therefore, $|1, 1\rangle$ reads

$$\boxed{|1, 1\rangle = \sqrt{\frac{1}{10}} \left(\sqrt{6} |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, 1\rangle_1 |1, 0\rangle_2 + |2, 0\rangle_1 |1, 1\rangle_2 \right)}. \quad (116)$$

Applying \hat{L}_- to $|1, 1\rangle$

$$\hat{L}_- |1, 1\rangle = \sqrt{1(1+1) - 1(1-1)} \hbar |1, 0\rangle = \sqrt{2} \hbar |1, 0\rangle, \quad (= \text{LHS}) \quad (117)$$

and to the RHS

$$\hat{L}_- \{\text{RHS}\} = \sqrt{\frac{1}{10}} \left(\sqrt{6} \hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{3} \hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2 + \hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2 \right). \quad (118)$$

Calculating separately $\hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2$

$$\begin{aligned} \hat{L}_- |2, 2\rangle_1 |1, -1\rangle_2 &= \hat{L}_{1-} |2, 2\rangle_1 |1, -1\rangle_2 + \hat{L}_{2-} |1, -1\rangle_2 |2, 2\rangle_1 = \\ &= 2\hbar |2, 1\rangle_1 |1, -1\rangle_2, \end{aligned} \quad (119)$$

and $\hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2$

$$\begin{aligned}\hat{L}_- |2, 1\rangle_1 |1, 0\rangle_2 &= \hat{L}_{1-} |2, 1\rangle_1 |1, 0\rangle_2 + \hat{L}_{2-} |1, 0\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{6}\hbar |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{2}\hbar |1, -1\rangle_2 |2, 1\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, 1\rangle_1 |1, -1\rangle_2 \right),\end{aligned}\quad (120)$$

and, lastly, $\hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2$

$$\begin{aligned}\hat{L}_- |2, 0\rangle_1 |1, 1\rangle_2 &= \hat{L}_{1-} |2, 0\rangle_1 |1, 1\rangle_2 + \hat{L}_{2-} |1, 1\rangle_2 |2, 0\rangle_1 = \\ &= \sqrt{6}\hbar |2, -1\rangle_1 |1, 1\rangle_2 + \sqrt{2}\hbar |1, 0\rangle_2 |2, 0\rangle_1 = \\ &= \sqrt{2}\hbar \left(\sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 + |2, 0\rangle_1 |1, 0\rangle_2 \right).\end{aligned}\quad (121)$$

Substituting in eq. (118)

$$\begin{aligned}\hat{L}_- \{\text{RHS}\} &= \sqrt{\frac{1}{10}} \left\{ \sqrt{6}2\hbar |2, 1\rangle_1 |1, -1\rangle_2 \right. \\ &\quad - \sqrt{3}\sqrt{2}\hbar \left(\sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, 1\rangle_1 |1, -1\rangle_2 \right) \\ &\quad \left. + \sqrt{2}\hbar \left(\sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 + |2, 0\rangle_1 |1, 0\rangle_2 \right) \right\} = \\ &= \sqrt{\frac{1}{5}}\hbar \left(\sqrt{3} |2, 1\rangle_1 |1, -1\rangle_2 - 2 |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 \right). \quad (= \text{RHS})\end{aligned}\quad (122)$$

Equating LHS = RHS with eqs. (117) and (122)

$$\sqrt{2}\hbar |1, 0\rangle = \sqrt{\frac{1}{5}}\hbar \left(\sqrt{3} |2, 1\rangle_1 |1, -1\rangle_2 - 2 |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 \right), \quad (123)$$

$|1, 0\rangle$ reads

$$\boxed{|1, 0\rangle = \sqrt{\frac{1}{10}} \left(\sqrt{3} |2, 1\rangle_1 |1, -1\rangle_2 - 2 |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2 \right).} \quad (124)$$

As for $|1, 1\rangle$, the combination $|1, -1\rangle$ can be obtained imposing orthogonality conditions between

$$\langle 3, -1 | 1, -1 \rangle = 0, \quad \langle 2, -1 | 1, -1 \rangle = 0. \quad (125)$$

Rewriting $|3, -1\rangle$ as

$$\begin{aligned}|3, -1\rangle &= \sqrt{\frac{1}{15}} \left(\sqrt{6} |2, 0\rangle_1 |1, -1\rangle_2 + 2\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2 \right) = \\ &= \sqrt{\frac{1}{15}} \left(\sqrt{6} |\psi_1\rangle + 2\sqrt{2} |\psi_2\rangle + |\psi_3\rangle \right),\end{aligned}\quad (126)$$

and $|2, -1\rangle$ as

$$\begin{aligned} |2, -1\rangle &= \sqrt{\frac{1}{6}} \left(-\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2 \right) = \\ &= \sqrt{\frac{1}{6}} \left(-\sqrt{3} |\psi_1\rangle + |\psi_2\rangle + \sqrt{2} |\psi_3\rangle \right). \end{aligned} \quad (127)$$

$|1, -1\rangle$ can be written as

$$|1, -1\rangle = c_1 |\psi_1\rangle + c_2 |\psi_2\rangle + c_3 |\psi_3\rangle. \quad (128)$$

Imposing the orthogonality conditions from eq. (125)

$$\begin{aligned} \langle 3, -1 | 1, -1 \rangle &= \sqrt{\frac{1}{15}} \left\{ \left(\sqrt{6} \langle \psi_1 | + 2\sqrt{2} \langle \psi_2 | + \langle \psi_3 | \right) (c_1 |\psi_1\rangle + c_2 |\psi_2\rangle + c_3 |\psi_3\rangle) \right\} = \\ &= \sqrt{\frac{1}{15}} \left(\sqrt{6}c_1 + 2\sqrt{2}c_2 + c_3 \right) = 0 \\ &\implies \sqrt{6}c_1 + 2\sqrt{2}c_2 + c_3 = 0, \quad (\text{I}) \end{aligned} \quad (129)$$

and

$$\begin{aligned} \langle 2, -1 | 1, -1 \rangle &= \sqrt{\frac{1}{6}} \left\{ \left(-\sqrt{3} |\psi_1\rangle + |\psi_2\rangle + \sqrt{2} |\psi_3\rangle \right) (c_1 |\psi_1\rangle + c_2 |\psi_2\rangle + c_3 |\psi_3\rangle) \right\} = \\ &= \sqrt{\frac{1}{6}} \left(-\sqrt{3}c_1 + c_2 + \sqrt{2}c_3 \right) = 0 \\ &\implies -\sqrt{3}c_1 + c_2 + \sqrt{2}c_3 = 0. \quad (\text{II}) \end{aligned} \quad (130)$$

The last equation can be obtained from the normalisation condition eq. (59)

$$|c_1|^2 + |c_2|^2 + |c_3|^2 = 1. \quad (\text{III}) \quad (131)$$

Then, the system of equations to solve reads

$$\begin{cases} \sqrt{6}c_1 + 2\sqrt{2}c_2 + c_3 = 0 & (\text{I}) \\ -\sqrt{3}c_1 + c_2 + \sqrt{2}c_3 = 0 & (\text{II}) \\ |c_1|^2 + |c_2|^2 + |c_3|^2 = 1 & (\text{III}) \end{cases} \quad (132)$$

Isolating c_1 from the first equation

$$c_1 = -\frac{1}{\sqrt{6}} \left(2\sqrt{2}c_2 + c_3 \right), \quad (133)$$

and c_2 from the second equation

$$c_2 = \sqrt{3}c_1 - \sqrt{2}c_3. \quad (134)$$

The expression for c_1 given in eq. (137) can be substituted in eq. (136), as

$$c_2 = \sqrt{3} \left[-\frac{1}{\sqrt{6}} \left(2\sqrt{2}c_2 + c_3 \right) \right] - \sqrt{2}c_3 = -2c_2 - \frac{3}{\sqrt{2}}c_3, \quad (135)$$

so

$$3c_2 = -\frac{3}{\sqrt{2}}c_3 \implies c_2 = -\frac{1}{\sqrt{2}}c_3. \quad (136)$$

Substituting now this expresion for c_2 in eq. (137)

$$c_1 = -\frac{1}{\sqrt{6}} \left[2\sqrt{2} \left(-\frac{1}{\sqrt{2}}c_3 \right) + c_3 \right] = -\frac{1}{\sqrt{6}} (-2c_3 + c_3) \implies c_1 = \frac{1}{\sqrt{6}}c_3. \quad (137)$$

And substituting this expresions for c_2 and c_3 in the third equation of the system

$$\left| \frac{1}{\sqrt{6}}c_3 \right|^2 + \left| -\frac{1}{\sqrt{2}}c_3 \right|^2 + |c_3|^2 = 1 \implies \left(\frac{1}{6} + \frac{1}{2} + 1 \right) |c_3|^2 = 1 \implies c_3 = \sqrt{\frac{3}{5}}. \quad (138)$$

Then, in eqs. (133) and (134), the coefficients read

$$c_1 = \frac{1}{\sqrt{6}}c_3 = \frac{1}{\sqrt{6}}\sqrt{\frac{3}{5}} = \sqrt{\frac{1}{10}}, \quad (139)$$

$$c_2 = -\frac{1}{\sqrt{2}}c_3 = -\frac{1}{\sqrt{2}}\sqrt{\frac{3}{5}} = -\sqrt{\frac{3}{10}}, \quad (140)$$

$$c_3 = \sqrt{\frac{3}{5}} = \sqrt{\frac{6}{10}}. \quad (141)$$

Therefore, $|1, -1\rangle$ from eq. (128) reads

$$|1, -1\rangle = \sqrt{\frac{1}{10}} |\psi_1\rangle - \sqrt{\frac{3}{10}} |\psi_2\rangle + \sqrt{\frac{6}{10}} |\psi_3\rangle = \sqrt{\frac{1}{10}} \left(|\psi_1\rangle - \sqrt{3} |\psi_2\rangle + \sqrt{6} |\psi_3\rangle \right), \quad (142)$$

then

$$\boxed{|1, -1\rangle = \sqrt{\frac{1}{10}} \left(|2, 0\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{6} |2, -2\rangle_1 |1, 1\rangle_2 \right).} \quad (143)$$

4 Summary

1. $L = 3$

- $|3, 3\rangle = |2, 2\rangle_1 |1, 1\rangle_2$
- $|3, 2\rangle = \sqrt{\frac{1}{3}} (|2, 2\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, 1\rangle_1 |1, 1\rangle_2)$
- $|3, 1\rangle = \sqrt{\frac{1}{30}} (\sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 + 4 |2, 1\rangle_1 |1, 0\rangle_2 + 2\sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2)$
- $|3, 0\rangle = \sqrt{\frac{1}{5}} (|2, 1\rangle_1 |1, -1\rangle_2 + \sqrt{3} |2, 0\rangle_1 |1, 0\rangle_2 + |2, -1\rangle_1 |1, 1\rangle_2)$
- $|3, -1\rangle = \sqrt{\frac{1}{15}} (\sqrt{6} |2, 0\rangle_1 |1, -1\rangle_2 + 2\sqrt{2} |2, -1\rangle_1 |1, 0\rangle_2 + |2, -2\rangle_1 |1, 1\rangle_2)$
- $|3, -2\rangle = \sqrt{\frac{1}{3}} (\sqrt{2} |2, -1\rangle_1 |1, -1\rangle_2 + |2, -2\rangle_1 |1, 0\rangle_2)$
- $|3, -3\rangle = |2, -2\rangle_1 |1, -1\rangle_2$

2. $L = 2$

- $|2, 2\rangle = \sqrt{\frac{1}{3}} (\sqrt{2} |2, 2\rangle_1 |1, 0\rangle_2 - |2, 1\rangle_1 |1, 1\rangle_2)$
- $|2, 1\rangle = \sqrt{\frac{1}{6}} (\sqrt{2} |2, 2\rangle_1 |1, -1\rangle_2 + |2, 1\rangle_1 |1, 0\rangle_2 - \sqrt{3} |2, 0\rangle_1 |1, 1\rangle_2)$
- $|2, 0\rangle = \sqrt{\frac{1}{2}} (|2, 1\rangle_1 |1, -1\rangle_2 - |2, -1\rangle_1 |1, 1\rangle_2)$
- $|2, -1\rangle = \sqrt{\frac{1}{6}} (-\sqrt{3} |2, 0\rangle_1 |1, -1\rangle_2 + |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 1\rangle_2)$
- $|2, -2\rangle = \sqrt{\frac{1}{3}} (-|2, -1\rangle_1 |1, -1\rangle_2 + \sqrt{2} |2, -2\rangle_1 |1, 0\rangle_2)$

3. $L = 1$

- $|1, 1\rangle = \sqrt{\frac{1}{10}} (\sqrt{6} |2, 2\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, 1\rangle_1 |1, 0\rangle_2 + |2, 0\rangle_1 |1, 1\rangle_2)$
- $|1, 0\rangle = \sqrt{\frac{1}{10}} (\sqrt{3} |2, 1\rangle_1 |1, -1\rangle_2 - 2 |2, 0\rangle_1 |1, 0\rangle_2 + \sqrt{3} |2, -1\rangle_1 |1, 1\rangle_2)$
- $|1, -1\rangle = \sqrt{\frac{1}{10}} (|2, 0\rangle_1 |1, -1\rangle_2 - \sqrt{3} |2, -1\rangle_1 |1, 0\rangle_2 + \sqrt{6} |2, -2\rangle_1 |1, 1\rangle_2)$

2×1		3		
	+3	3	2	
+2 +1	1	+2	+2	
	+2 0	1/3 2/3	3 2 1	
	+1 +1	2/3 -1/3	+1 +1 +1	
		+2 -1	1/15 1/3 3/5	
		+1 0	8/15 1/6 -3/10	3 2 1
		0 +1	6/15 -1/2 1/10	0 0 0
			+1 -1	1/5 1/2 3/10
			0 0	3/5 0 -2/5
			-1 +1	1/5 -1/2 3/10
				3 2 1
				-1 -1 -1
				0 -1
				6/15 1/2 1/10
				-1 0
				8/15 -1/6 -3/10
				-2 +1
				1/15 -1/3 3/5
				3 2
				-2 -2
				-1 -1
				2/3 1/3
				-2 0
				1/3 -2/3
				3
				-3
				-2 -1
				1