

INF5620 Project

Extended Navier-Stokes solver for Platelet Aggregation

Jarle Sogn
Institutt for matematikk
Universitetet i Oslo
jarlesog@math.no

November 28, 2012

Abstract

The aim of the project is to solve an extended set of Navier-Stokes equations by the finite element method and *fenics*. The model I'm using is a simplification of model found in Aaron L. Fogelson article: **Continuum models of platelet aggregation: Formulation and mechanical properties**. I will focus mainly on equation 1 - 3 in this article. My solver is an extantion of *fenics* Navier-Stoke demo.

The Model

I focus on equation 1 - 3 from the article, witch are:

$$\rho (\mathbf{u}_t + \mathbf{u} \cdot \nabla \mathbf{u}) = -\nabla p + \mu \Delta \mathbf{u} + \mathbf{f}^g \quad (1)$$

$$\nabla \cdot \mathbf{u} = 0 \quad (2)$$

$$\frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla \phi = C \Delta \phi - R(c) \phi \quad (3)$$

The first two equation are Navier-Stokes equations for incompressible fluid. $\mathbf{u}(t, \mathbf{x})$ is the fluid velocity field, $p(t, \mathbf{x})$ is the pressure, ρ is the fluid's mass density(I'll sett this equal to 1, for simplicity.) and μ is the fluids viscosity(assumed to be constant). $\epsilon^{-3}\phi$ is the concentration of non-activated

platelet(ϵ^{-3} is a scaling constant). C is a diffusion constant¹. The term $R(c)\phi$ is the rate in which the non-activated platelets are converted to active platelets. This depends on the ADP concentration c , which vary. Since I'm not using the equation for c , I will have to manufacture a suitable function $c(\mathbf{x}, t)$ ².

Solving Navier-Stokes equation

Now I'm in detail going to talk about one way of solving the Navier-Stokes equations(1 and 2) with *fenics* and the finite element method. I will use Chorin's projection method. The idea is first to compute a tentative velocity(\mathbf{u}^*) by ignoring the pressure in equation 1 and then project the velocity onto the space of divergence free vector fields. Equation 1 then becomes:

$$\frac{\partial}{\partial t} \mathbf{u}^* + \mathbf{u} \cdot \nabla \mathbf{u} - \mu \Delta \mathbf{u}^* = \mathbf{f}^g$$

Let $V = H_0^1(\Omega)$ be a Sobolev space and Ω the domain. I want to write the equation above as a weak formulation. I multiply with the function $\mathbf{v}(\mathbf{x})$ and integrate the space.

$$\int_{\Omega} \frac{\partial \mathbf{u}^*}{\partial t} \cdot \mathbf{v} + (\mathbf{u} \cdot \nabla \mathbf{u}) \cdot \mathbf{v} - \mu \Delta \mathbf{u}^* \cdot \mathbf{v} d\mathbf{x} = \int_{\Omega} \mathbf{f}^g \cdot \mathbf{v} d\mathbf{x} \quad \forall \mathbf{v} \in V$$

Now if I integrate by parts the last term in the left hand side(or use Green's formula), I obtain:

$$\int_{\Omega} \frac{\partial \mathbf{u}^*}{\partial t} \cdot \mathbf{v} + (\mathbf{u} \cdot \nabla \mathbf{u}) \cdot \mathbf{v} + \mu \nabla \mathbf{u}^* \cdot \nabla \mathbf{v} d\mathbf{x} = \int_{\Omega} \mathbf{f}^g \cdot \mathbf{v} d\mathbf{x} \quad \forall \mathbf{v} \in V$$

Finely I discretize the time(n) and space(h), $V_h \subset V$. Then \mathbf{u} becomes \mathbf{u}_h^n . I define $\langle \cdot, \cdot \rangle$ as the $L^2(\Omega)$ norm. The above equation can be written as:

$$\langle D_t^n \mathbf{u}_h^*, \mathbf{v} \rangle + \langle \mathbf{u}_h^{n-1} \cdot \nabla \mathbf{u}_h^{n-1}, \mathbf{v} \rangle + \langle \mu \nabla \mathbf{u}_h^*, \nabla \mathbf{v} \rangle = \langle \mathbf{f}^{g,n}, \mathbf{v} \rangle \quad \forall \mathbf{v} \in V_h \quad (4)$$

The projection give us these two equation:

$$\frac{\mathbf{u}_h^n - \mathbf{u}_h^*}{k_n} + \nabla p_h^n = 0 \quad (5)$$

¹Fogelson article uses D_n and ϕ_n for C and ϕ .

²To understand the equations and symbols better, I advise you to take a look at the first few first pages of Fogelson article

$$\nabla \cdot \mathbf{u}_h^n = 0 \quad (6)$$

where k_h is the size of the local time step. I now multiply 5 with ∇q where $q \in Q_h \subset L^2(\Omega)$ and integrate.

$$\frac{1}{k_n} \int_{\Omega} \mathbf{u}_h^n \cdot \nabla q - \mathbf{u}_h^* \cdot \nabla q + k_n \nabla p_h^n \cdot \nabla q dx = 0 \quad \forall q \in Q_h$$

I integrate by parts the first two terms. The first term becomes zero from equation 6 and $\mathbf{u}|_{\partial\Omega} = 0$.

$$\int_{\Omega} (\nabla \cdot \mathbf{u}_h^*) q / k_n + (\nabla p_h^n \cdot \nabla q) dx = 0 \quad \forall q \in Q_h$$

This can be written as:

$$\langle \nabla p_h^n, \nabla q \rangle = - \langle \nabla \cdot \mathbf{u}_h^*, q \rangle / k_n \quad \forall q \in Q_h \quad (7)$$

Finally I multiply equation 5 with $\mathbf{v} \in V_h$ and integrate. I obtain:

$$\langle \mathbf{u}_h^n, v \rangle = \langle \mathbf{u}_h^*, v \rangle - k_n \langle \nabla p_h^n, v \rangle \quad \forall v \in V_h \quad (8)$$

The three equation 4, 7 and 8 is the Chorin's projection method scheme for solving Navier-Stokes equation. First solve for \mathbf{u}_h^* in equation 4, the solve the pressure p_h^n in equation 7 and finally find the velocity \mathbf{u}_h^n in equation 8. Repeat this for each time-step.

Stability

When running the demo program I note that the speed is less then $1mm/s$. My model is suppose to model blood flow throw arteries. The velocities throw arteries is about $1000mm/s$, hence I need to turn up the pressure to acquire this velocity.

After increasing the velocity in the demo program, it became unstable. The velocity started fluxing and became non-physically high. It turned out to be advection-diffusion that cause the unstable behavior: $\mathbf{u} \cdot \nabla \mathbf{u} \gg \mu \Delta \mathbf{u}$. Similar problem will probably happen for the extension: $\mathbf{u} \cdot \nabla \phi \gg \mu \Delta \phi$.

The extension

As mention earlier I have extended an already existing program. Making it also solve equation 3. In order to do this I have to rewrite 3 to variation form. I multiply the test function $\psi \in \Psi \subset H_0^1(\Omega)$ and integrate.

$$\int_{\Omega} (D_t \phi) \psi + \mathbf{u} \cdot (\nabla \phi) \psi - C (\Delta \phi) \psi dx = - \int_{\Omega} R(\mathbf{x}, t) \phi \psi dx \quad \forall \psi \in \Psi$$

after integrating by part and using backward euler as D_t , I obtain:

$$\int_{\Omega} \left(\frac{\phi^n - \phi^{n-1}}{k_n} \right) \psi + \mathbf{u} \cdot (\nabla \phi^n) \psi + C (\nabla \phi^n \cdot \nabla \psi) + R(\mathbf{x}, t^n) \phi^n \psi dx = 0 \quad \forall \psi \in \Psi$$

I end up with the scheme:

$$\langle \phi^n - \phi^{n-1}, \psi \rangle \frac{1}{k_n} + \langle \psi \nabla \phi^n, \mathbf{u}^n \rangle + C \langle \nabla \phi^n, \nabla \psi \rangle + \langle R \phi^n, \psi \rangle = 0 \quad \forall \psi \in \Psi_h \quad (9)$$