

# A new framework to testing modified gravity using galaxy surveys

Joaquin Armijo-Torres. IPMU postdoc colloquium

# Hello!

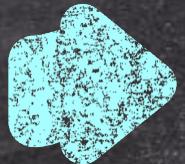
New postdoc started in  
October!

I'm at A049, please pass by.  
Interested in cosmology and  
the large-scale structure of  
the Universe.

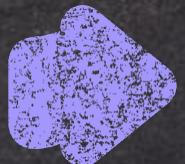


# Outline of my talk.

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Cold dark matter universe with and without cosmological constant:  $\Lambda$ CDM vs Modified gravity.



GR and MG simulations and mock galaxy catalogues.

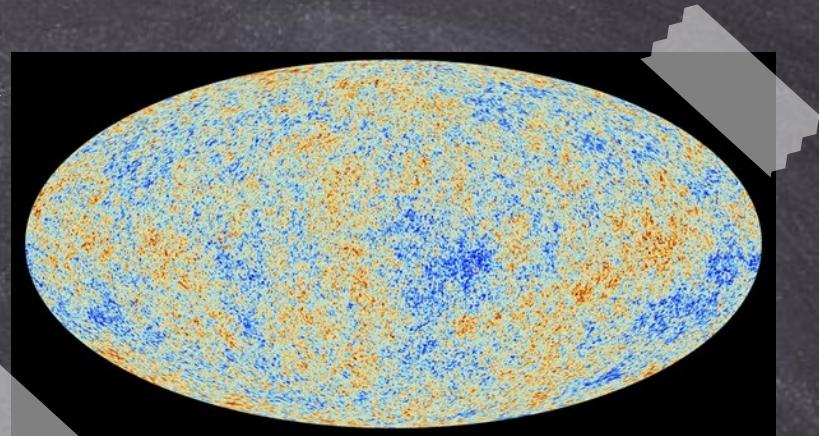
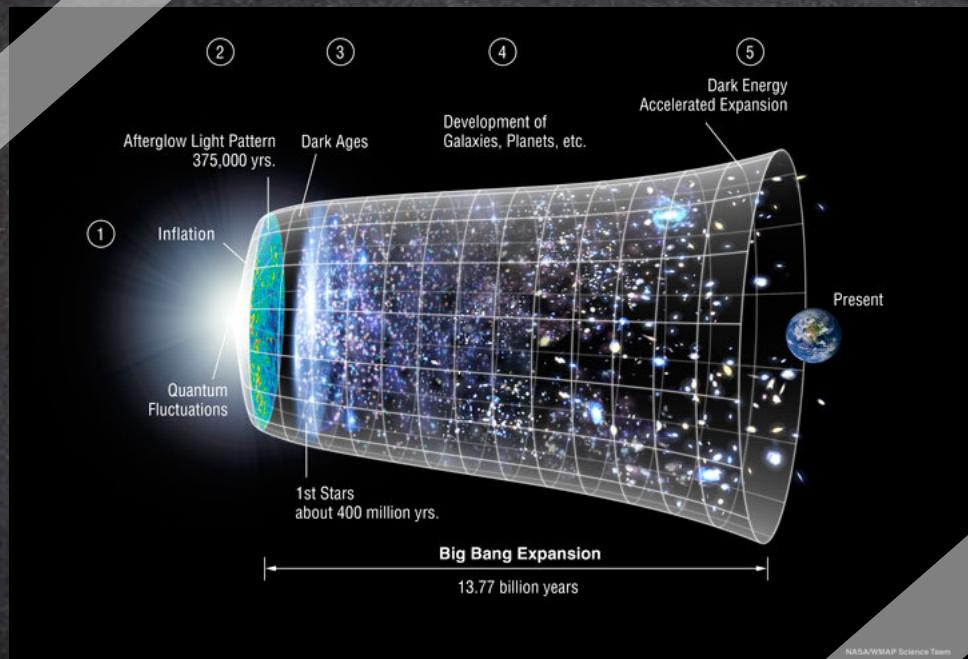


Probes for modified gravity: what does the LSS tell us beyond 2-point statistics? A new approach.

01

# $\Lambda$ CDM vs. Modified gravity

# The $\Lambda$ CDM Universe



$\Lambda$  drives dark energy which accounts for ~70% energy density budget. The rest comes from matter (25% dark matter, 5% baryons).

What about modified gravity?

# Modified gravity

Ezquiaga & Zumalacarregui (2017)

$$c_g = c$$

$$c_g \neq c$$

Horndeski

General Relativity

quintessence/k-essence [46]

Brans-Dicke/ $f(R)$  [47, 48]

Kinetic Gravity Braiding [50]

beyond H.

Derivative Conformal (19) [17]

Disformal Tuning (21)

quadratic DHOST with  $A_1 = 0$

Viable after GW170817

kinetic braiding model



quartic/quintic Galileons [13, 14]

Fab Four [15]

de Sitter Horndeski [49]

$G_{\mu\nu}\phi^\mu\phi^\nu$  [51],  $f(\phi)\cdot$ Gauss-Bonnet [52]

quartic/quintic GLPV [18]

quadratic DHOST [20] with  $A_1 \neq 0$

cubic DHOST [23]

Non-viable after GW170817

symmetron models  
dilaton models

Courtesy Baojiu Li

# Modified gravity

- Modified gravity can explain the cosmic acceleration without a cosmological constant.
- Scalar field coupled to matter or extra term in the Einstein-Hilbert action triggers extra fifth force that enhances gravity.
- Screening mechanism (Chameleon in  $f(R)$ ) hides the fifth force in high density regions. This is needed to make observationally viable theory (solar system scale).
- The fifth force is screened in the early Universe (CMB is unchanged).
- Gravity needs to be probed at cosmic scale using the Large-scale structure.

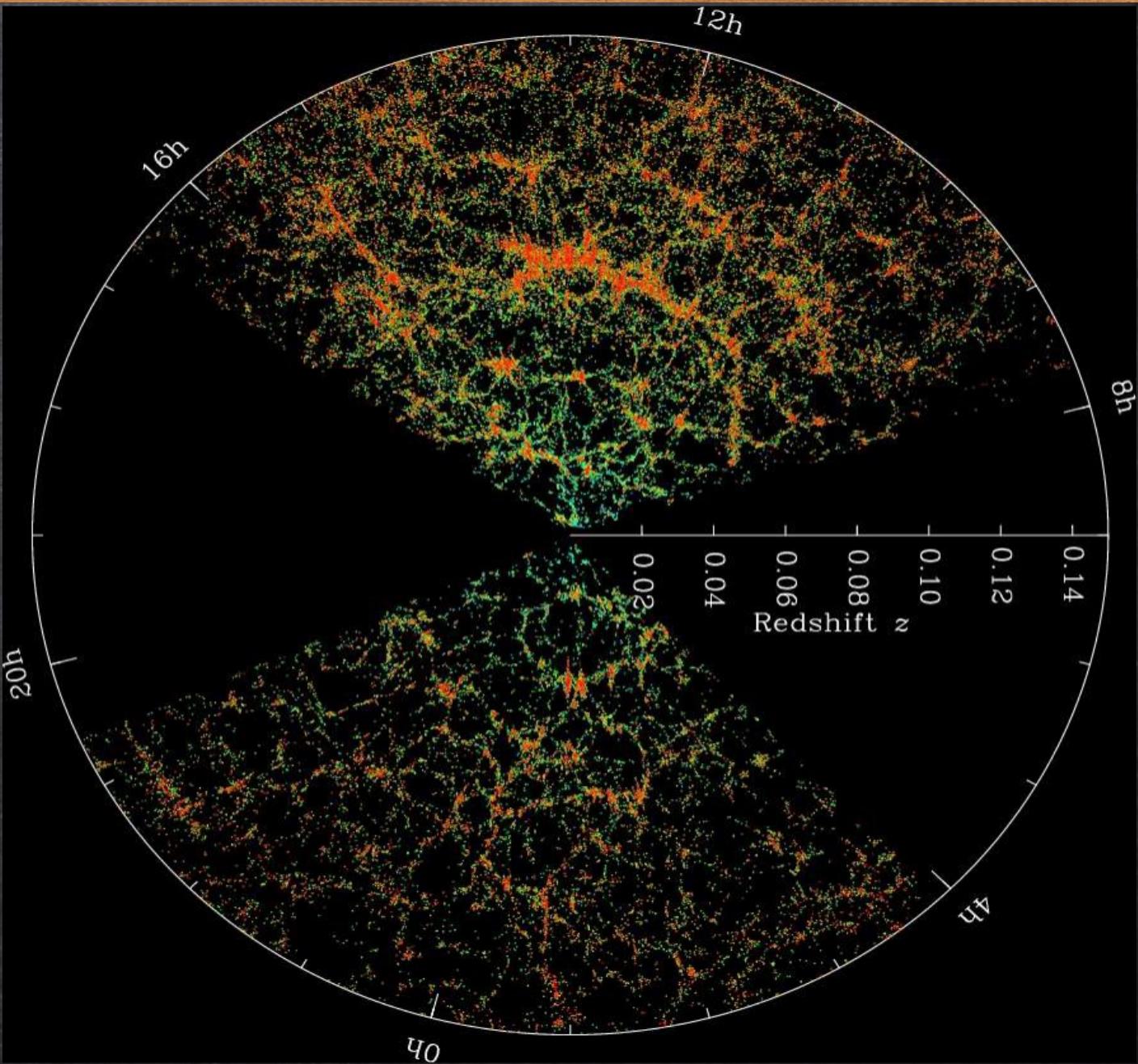


# Modified gravity

Matter density at  $z \sim 0$  is highly non-linear.

It shows galaxies but not mass.

We understand roughly how galaxies populate the dark matter but with some few parameters.



# $f(R)$ gravity

Replace the cosmological constant by  $f(R)$  in the action:

$$S = \int d^4x \sqrt{-g} \left[ \frac{1}{2\kappa^2} (R + f(R)) + \mathcal{L}_m \right]$$

where a modified Poisson equation can be obtained (varying the Action):

$$\vec{\nabla}^2 \Phi = 4\pi G a^2 \delta \rho_m - \frac{1}{2} \vec{\nabla}^2 f_R$$

The new term  $f_R \equiv \frac{df}{dR}$  (the scalaron) is understand a new potential which mediates a new effective "fifth force". In the limit

$$\lim_{R \rightarrow \infty} f(R) = \text{const.}$$

When  $f_R \rightarrow 0$ ,  $f(R) = 2\Lambda$ , which is the first condition for viable  $f(R)$ . The second limit is for  $f_R \rightarrow 0$  when  $\rho_m \rightarrow \infty$  (Chameleon mechanism).

The Hu & Sawicki model satisfy these conditions with  $f(R)$  constant in the background cosmology throughout cosmic history.

# The Hu & Sawicki model

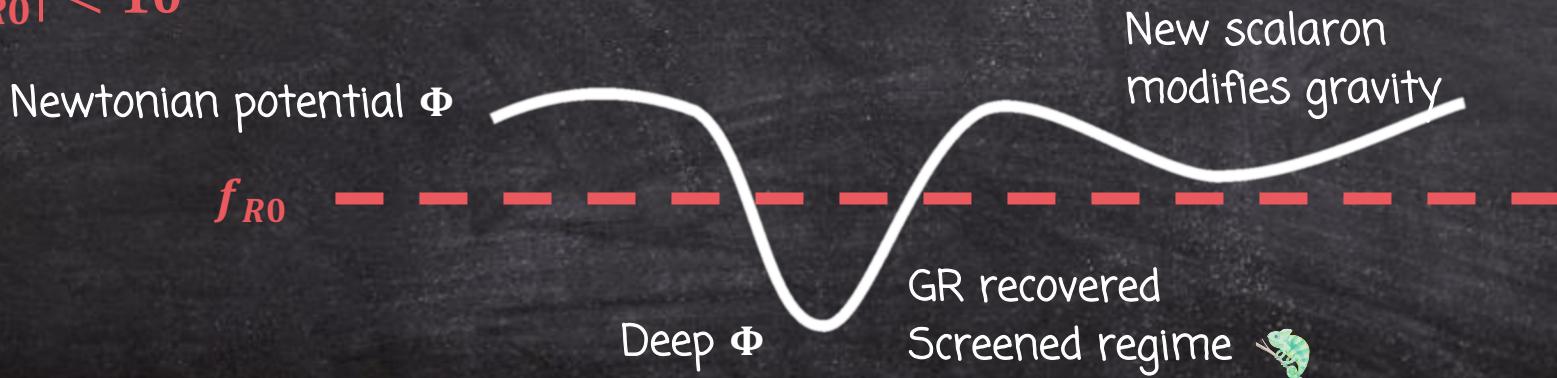
The function  $f(R)$  takes the form:

$$f(R) = -m^2 \frac{c_1 \left(\frac{R}{m^2}\right)^n}{c_2 \left(\frac{R}{m^2}\right)^n + 1}, \text{ (Hu and Sawicki 2007)}$$

Satisfying the previous limit for high curvature we can expand and solve for the background cosmology:

$$f(R) \approx \frac{c_1}{c_2} m^2 + \frac{c_1}{c_2^2} m^2 \left(\frac{m^2}{R}\right)^n, \text{ with } \frac{c_1}{c_2} = \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \text{ and } \frac{c_1}{c_2^2} = -\frac{1}{n} \left[ 3 \left( 1 + 4 \frac{\Omega_{\Lambda,0}}{\Omega_{m,0}} \right) \right]^{n+1} f_{R0}.$$

We can constraint the model based in 2 free parameters. For  $n=1$  we obtain  $|f_{R0}| < 10^{-4}$  (Schmidt et al. 2009). Current constraints using abundance of clusters and weak lensing give  $|f_{R0}| < 10^{-5}$



02

GR and MG simulations

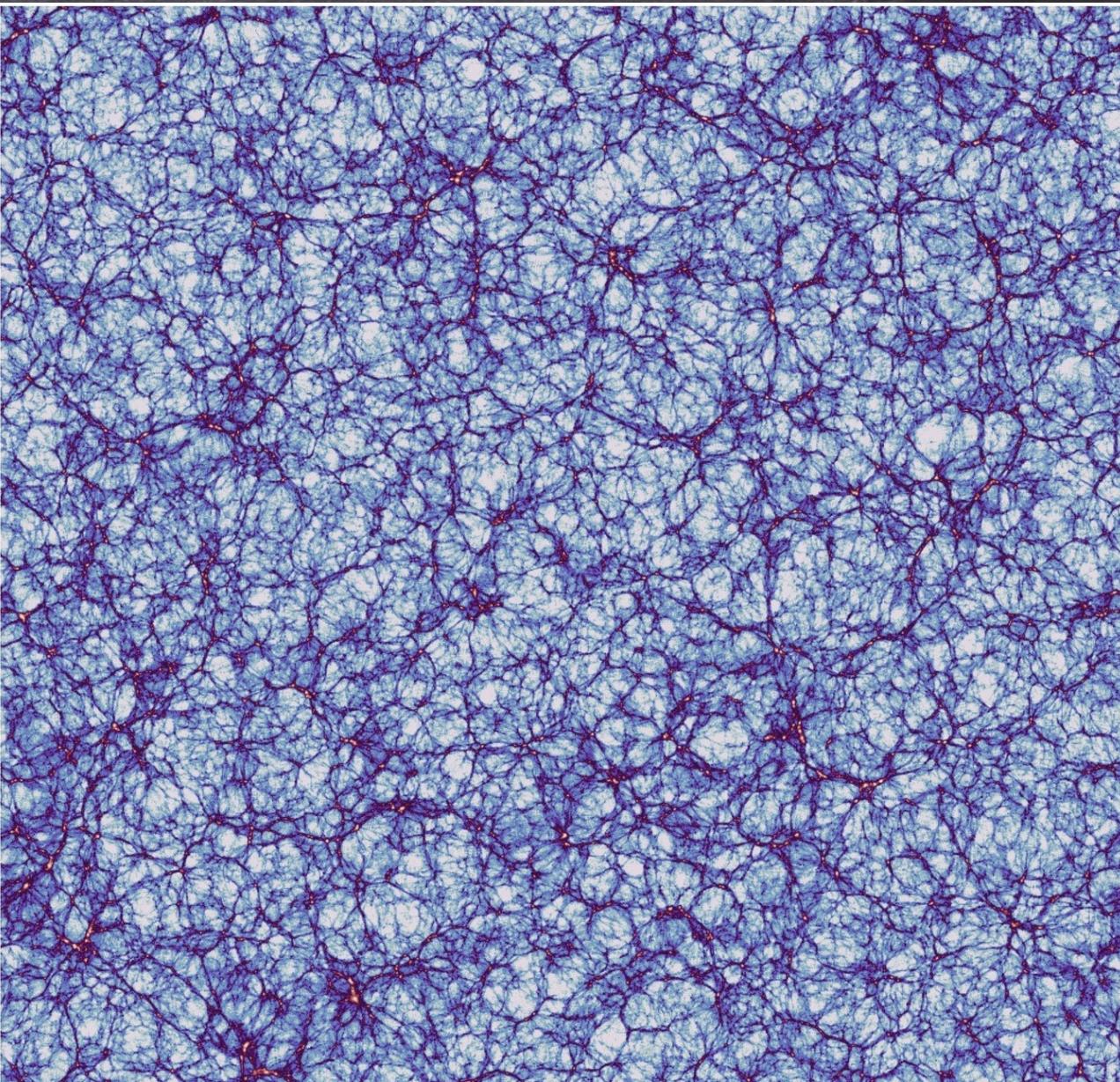
# Modified gravity simulations

Fully non-linear equations of motion for gravity  $\Phi$  and the scalaron  $f_R$  (Poisson equation).

Solutions to these equations using the ECOSMOG code (Li et al. 2012) and MG-GADGET (Arnold et al. 2018).

Simulations for GR and  $f(R)$  models with  $|f_{R0}| < \mathbf{10^{-5}}, \mathbf{10^{-6}}$  (F5, F6).

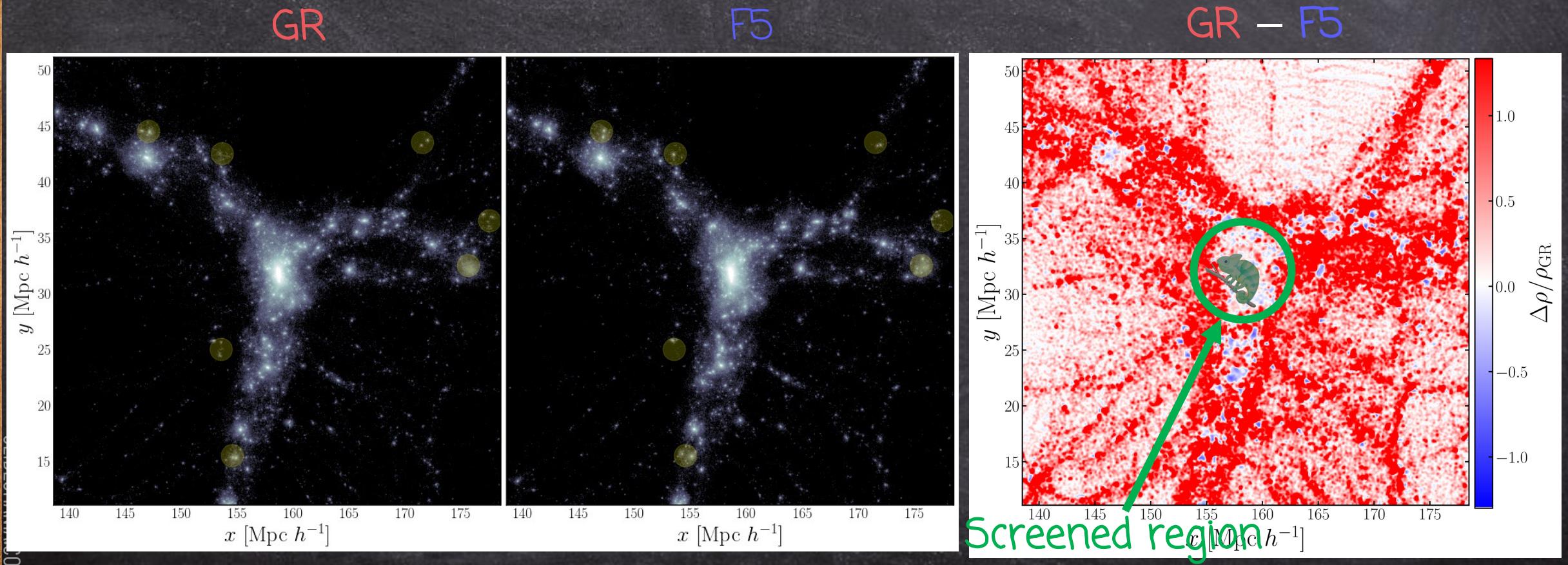
Simulations start from the same initial conditions.



The MG lightcone project (Arnold et al. 2018)  $z = 0, L_{\text{box}} = 768 \text{ Mpc } h^{-1}$

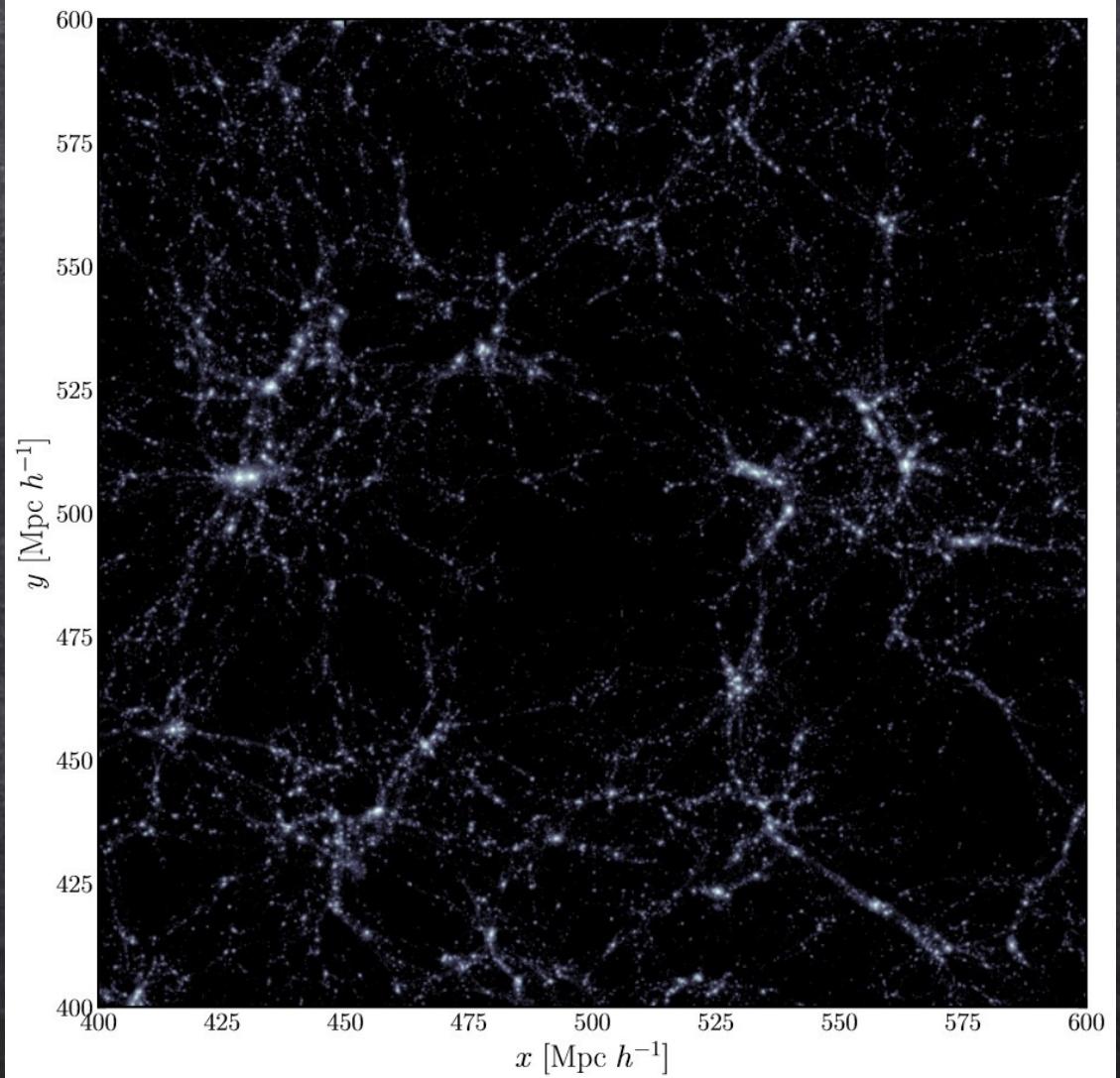
# Modified gravity simulations

Chameleon mechanism: Gravity is modified, however screened in high-density regions (e.g. inside a large mass halo)



# Testing gravity using the cosmic web

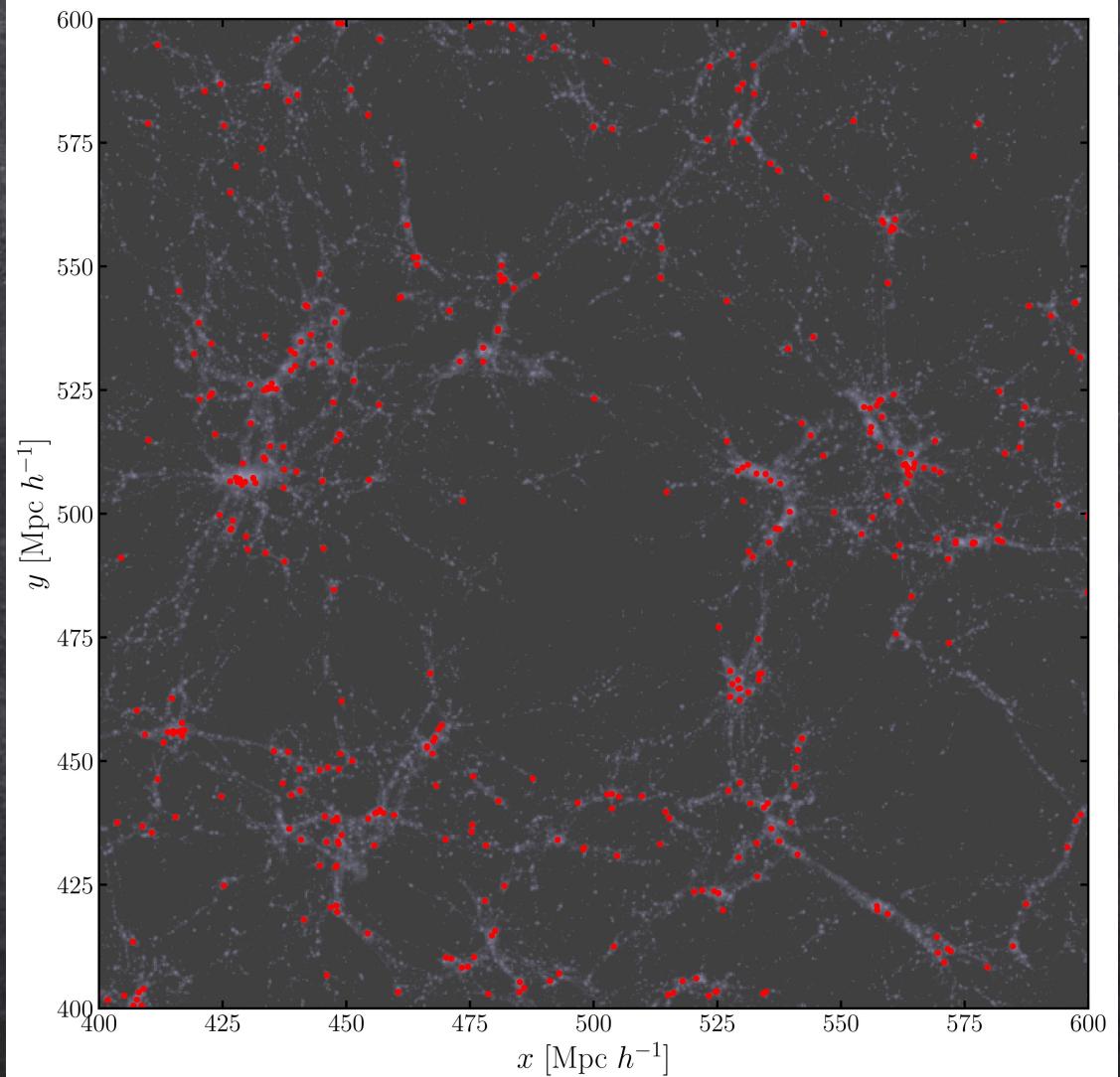
- The distribution of matter at late times forms large-scale structures.
- These structures are shaped mainly by gravity and the late cosmic acceleration, creating different environments where galaxies live.
- Nodes, filaments, walls and voids can be found in the cosmic web.



# Testing gravity using the cosmic web

We observe galaxies:

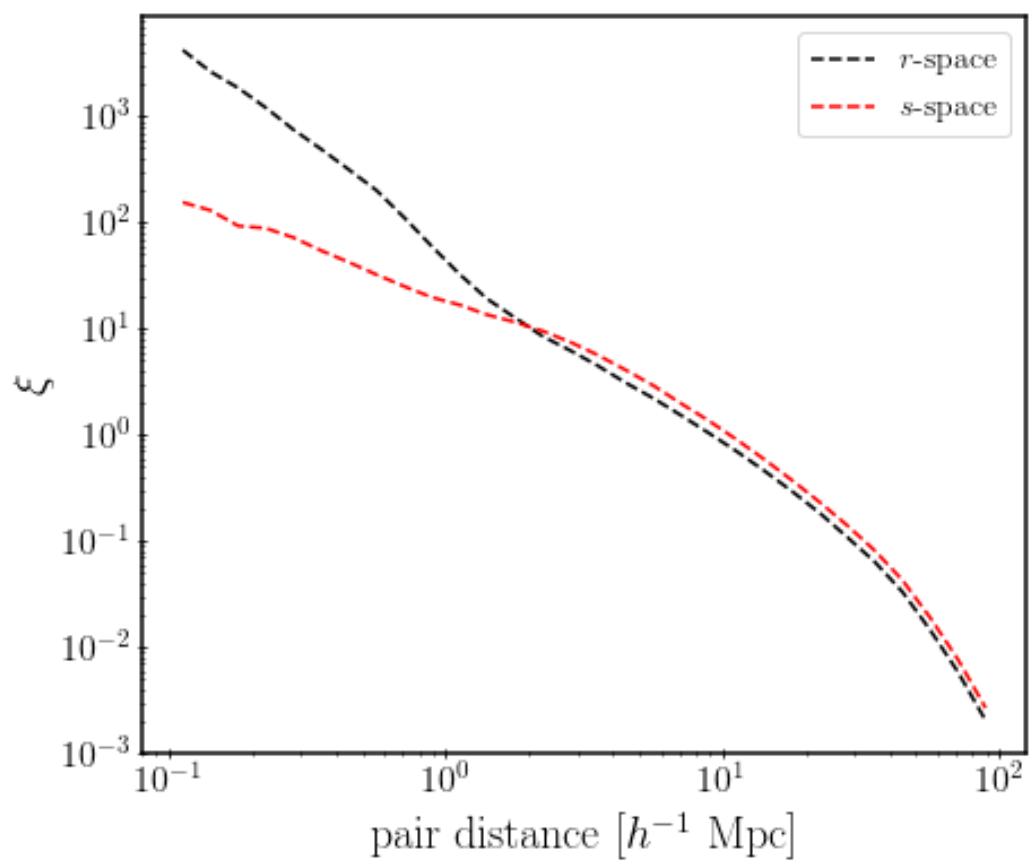
- Galaxies trace the underlying matter distribution assuming a non-linear and stochastic bias relation  $\delta_g = b\delta_m$ .
- They follow the formation of the same structures as the cosmic web formed by matter.
- Studying the connection between dark matter (haloes) and galaxies help us to understand the cosmological model.



# Galaxy clustering: The 2-point correlation function

$$1 + \xi = \frac{DD}{RR}$$

Estimator: relates the number of galaxy pairs DD in comparison to a random distribution of pairs RR

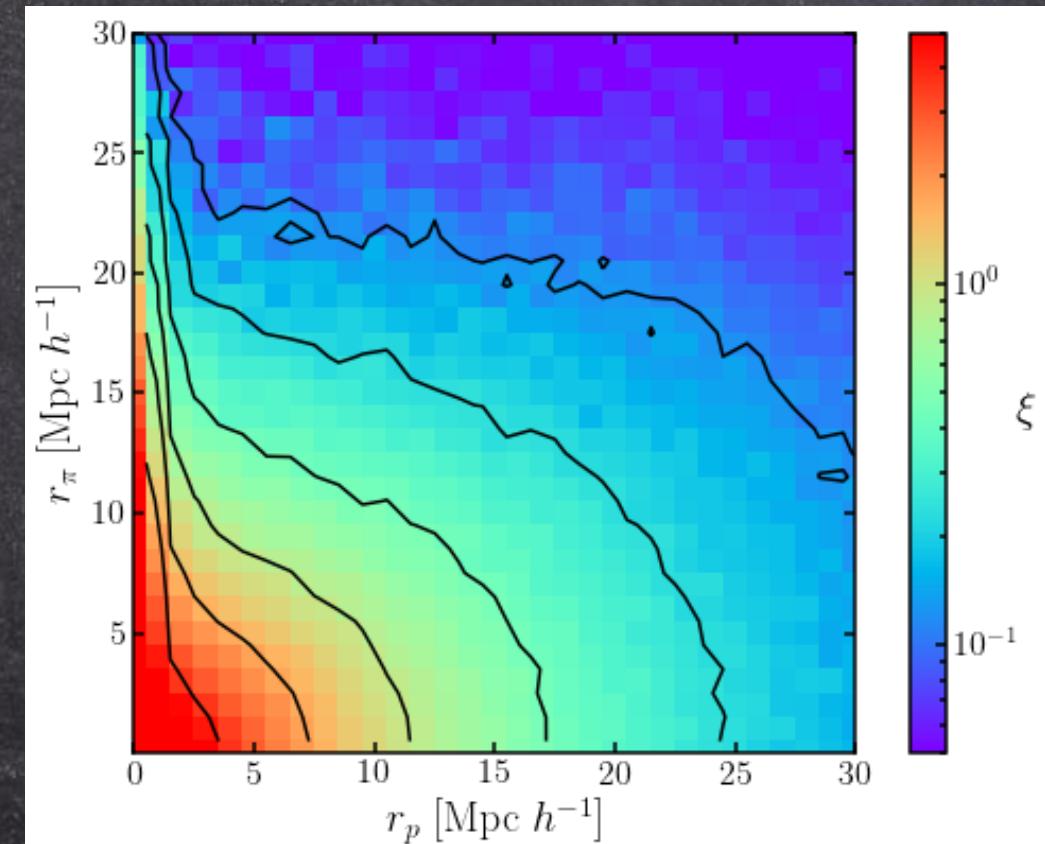


$$s = r + \frac{(1+z)v_{||}}{H(z)} \hat{e}_{||}$$

3D pair distance including pairwise  
'peculiar' velocities

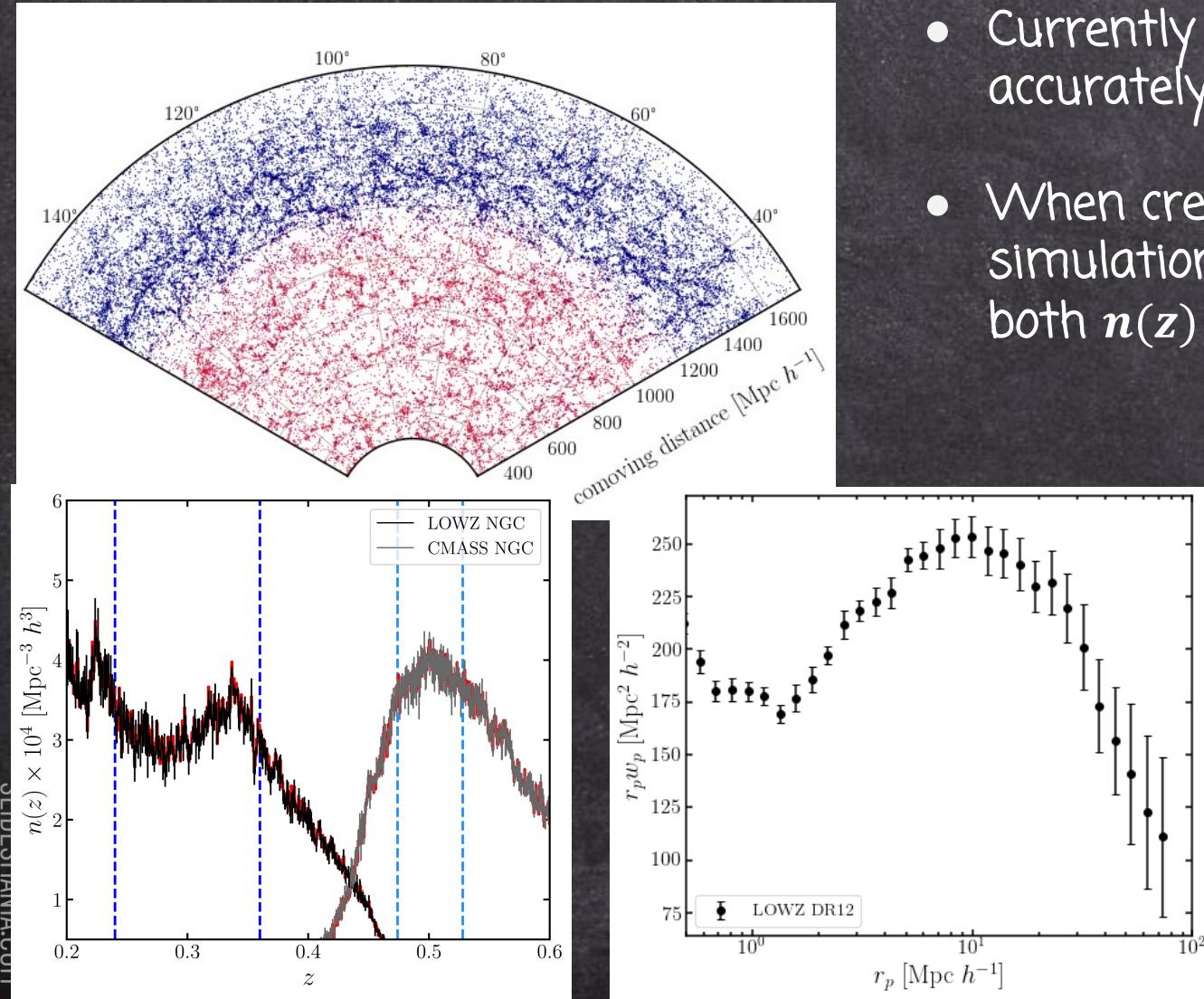
$$\frac{w_p(r_p)}{r_p} = \frac{1}{r_p} \int_{-\infty}^{\infty} dr_{\pi} \xi(r_p, r_{\pi})$$

The projected-correlation function



# Large-scale galaxy surveys

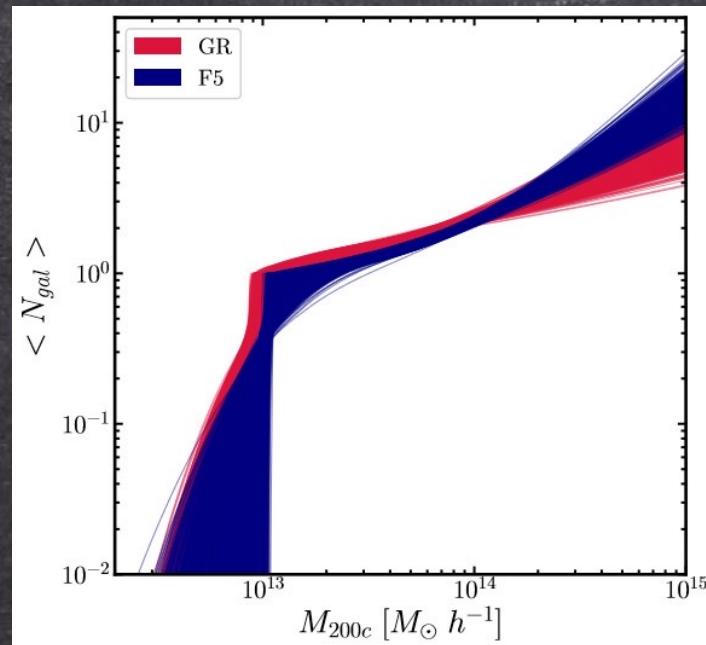
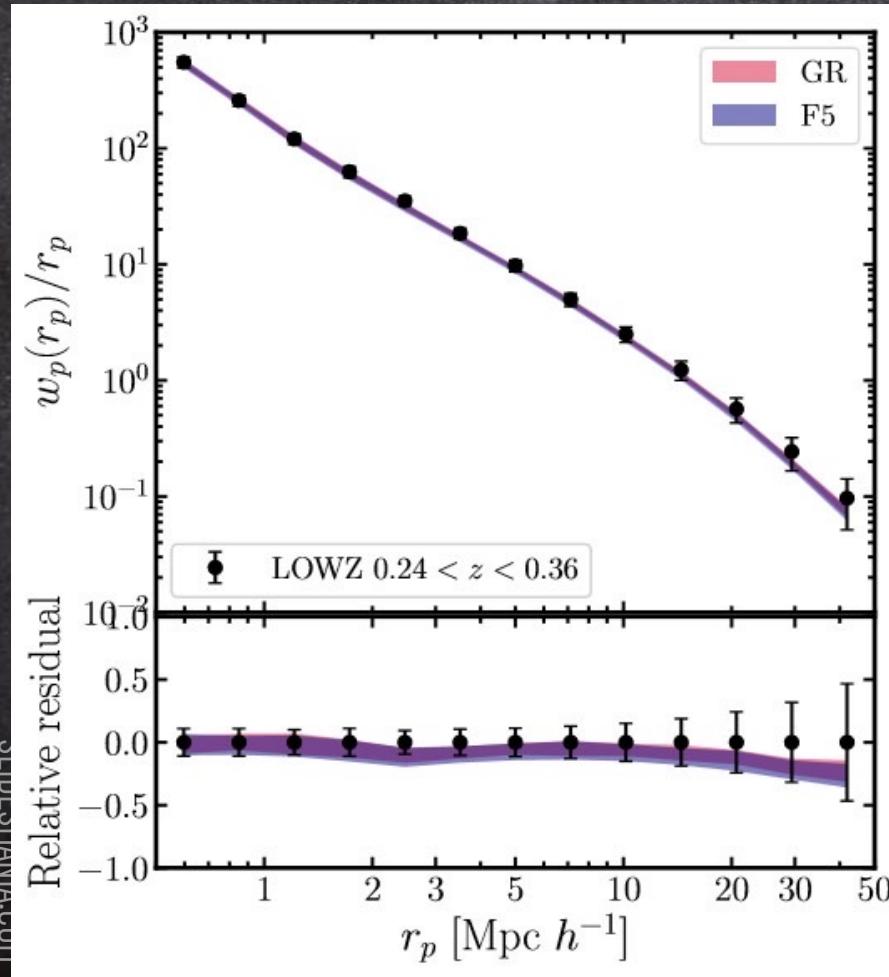
SDSS-III BOSS



- Currently 1- and 2-point functions are measured accurately in a large range of scales.
- When creating mock galaxy catalogues from simulations (GR and MG) we are forced to reproduce both  $n(z)$  and the projected 2-point clustering.
- We can rely in the fine tuning of parameters of the halo-galaxy models to generate such observational constraints. (Cautun et al. 2017, Armijo et al. 2018, Paillas et al. 2019)

# Can we replicate these results directly in the simulations? That's the neat part! Absolutely yes!

We need to create galaxy mock catalogues



The Halo occupation distribution (HOD) model:

Average number of galaxies inside a dark matter halo as function of its mass:

$$\langle N_c \rangle = \frac{1}{2} \left[ 1 + \text{erf} \left( \frac{\log M_{min} - M}{\sigma_{\log M_{min}}} \right) \right]$$

$$\langle N_s \rangle = \langle N_c \rangle \left( \frac{M - M_0}{M_1} \right)^\alpha$$

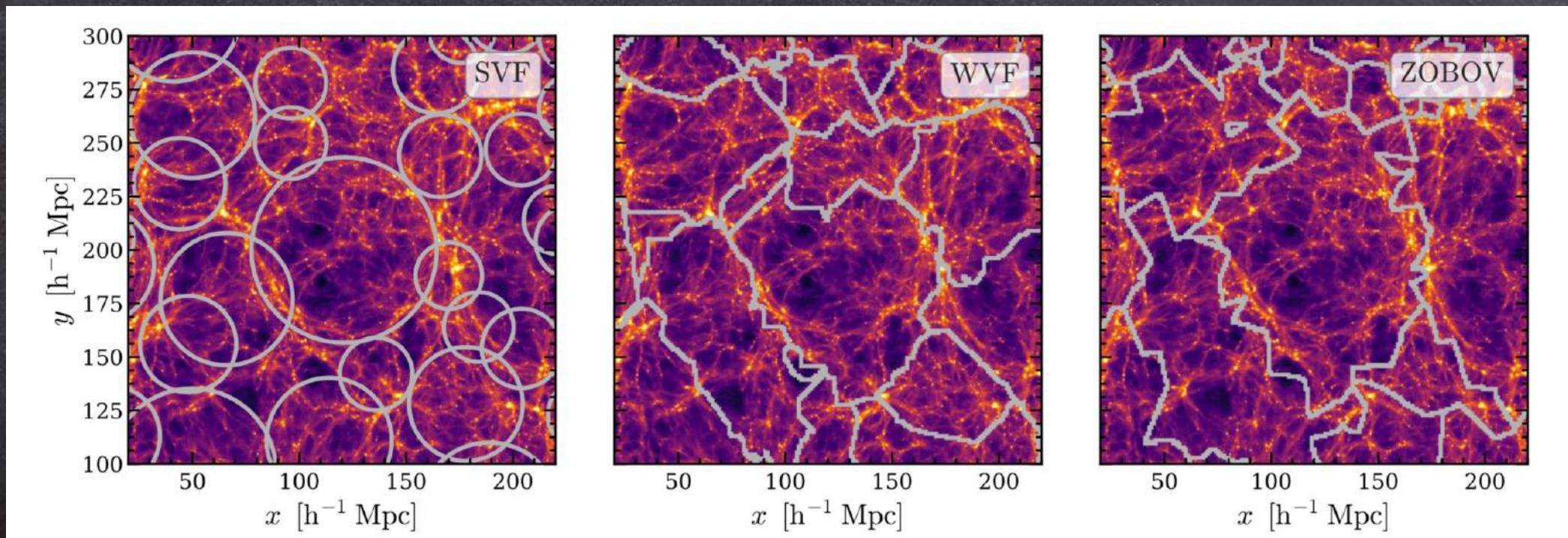
(caveat) the HOD function (5 parameters) is degenerated with the number density of galaxies  $n_{gal}$  and the clustering  $\xi$ .

03

Probes for modified gravity

# Weak lensing in cosmic voids

- Finding void regions in the cosmic web and comparing void finders: 7 void finders (4 3D-VF, 3 2D-VF).
- Using galaxies in voids to measure weak lensing statistics and compare it for different gravity models.

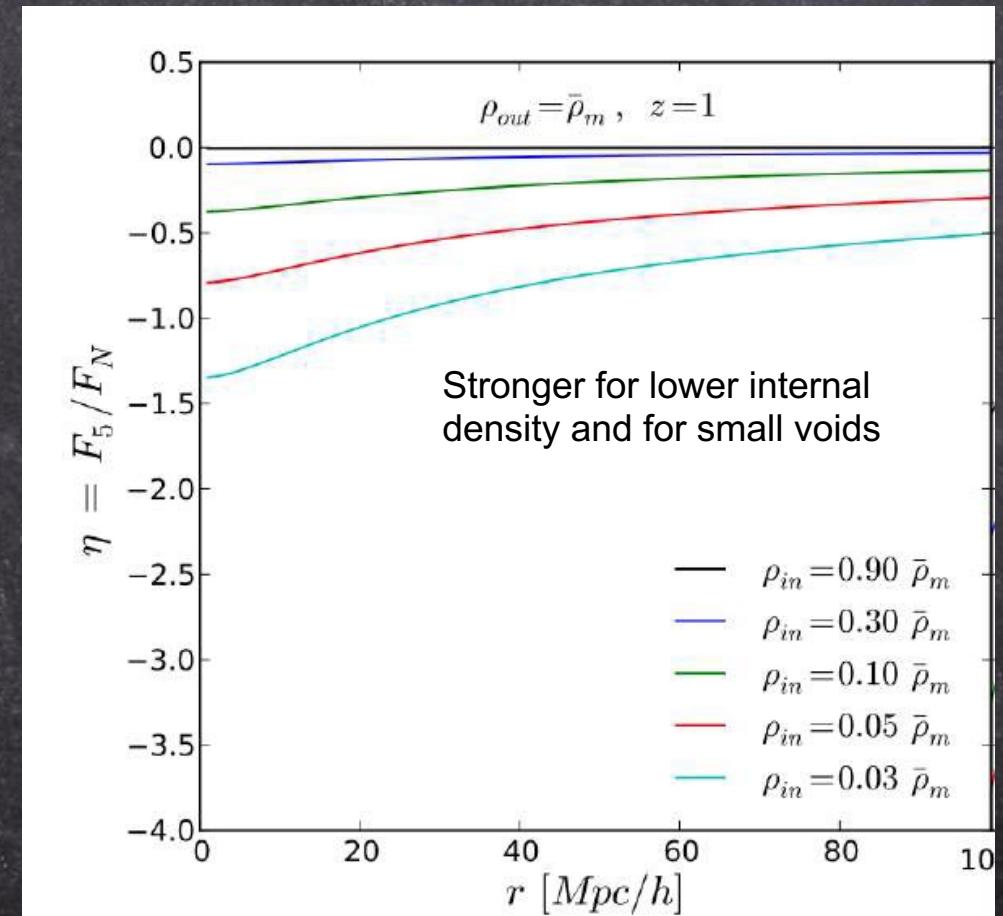
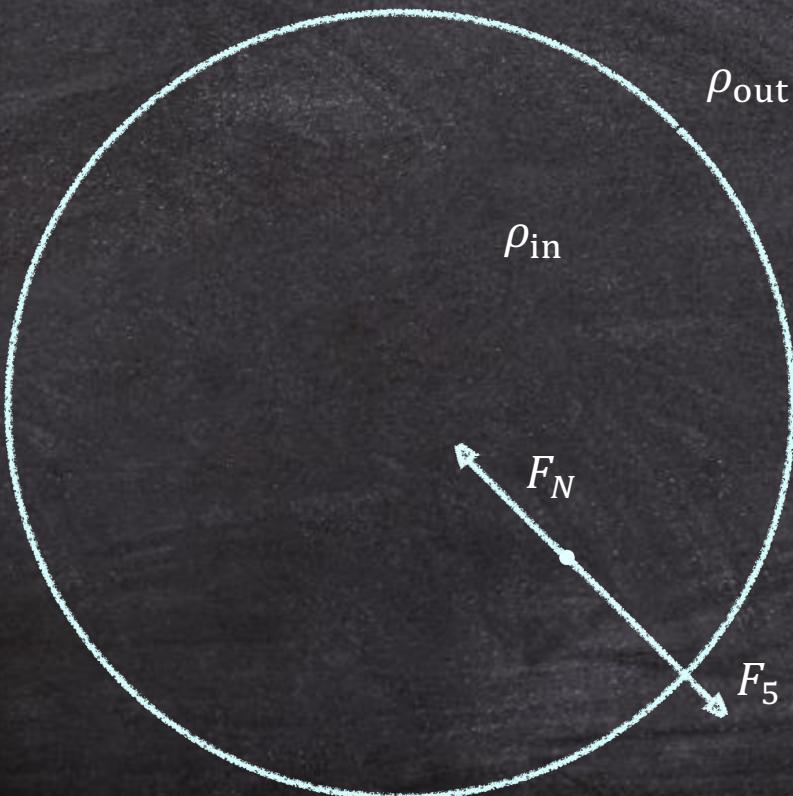


The SHED collaboration: Cautun et al. 1710.01730; Paillas et al. 1810.02864

# $f(R)$ voids: prediction using the spherical top-hat model:

Clampitt et al. 2013 calculate the Newtonian and fifth force for a top-hat empty region

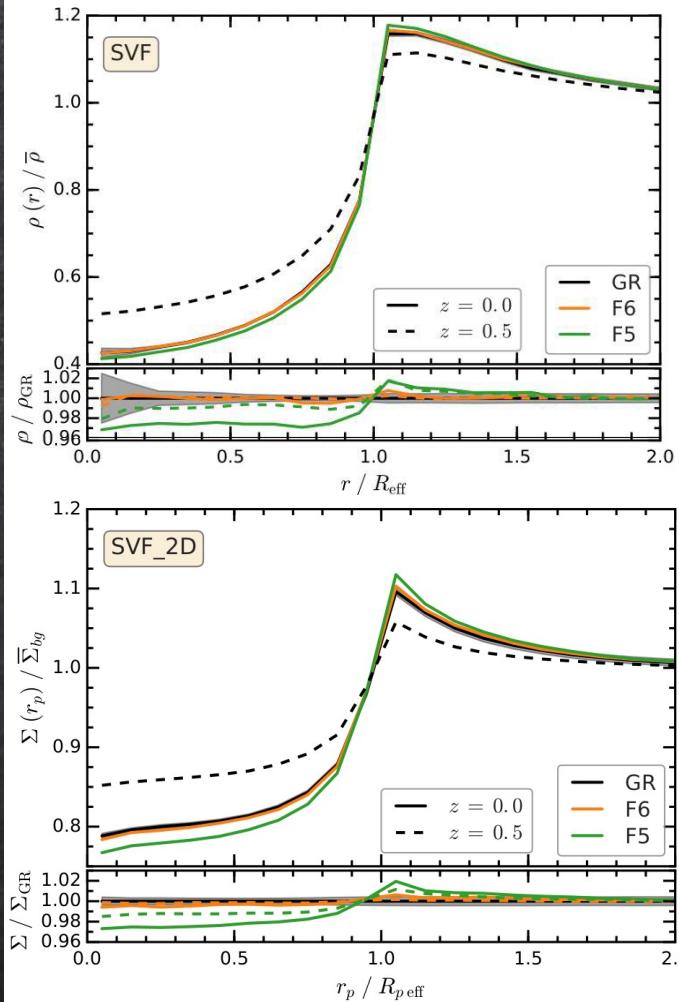
Negative fifth force inside voids acting in opposite  
Direction to Newtonian



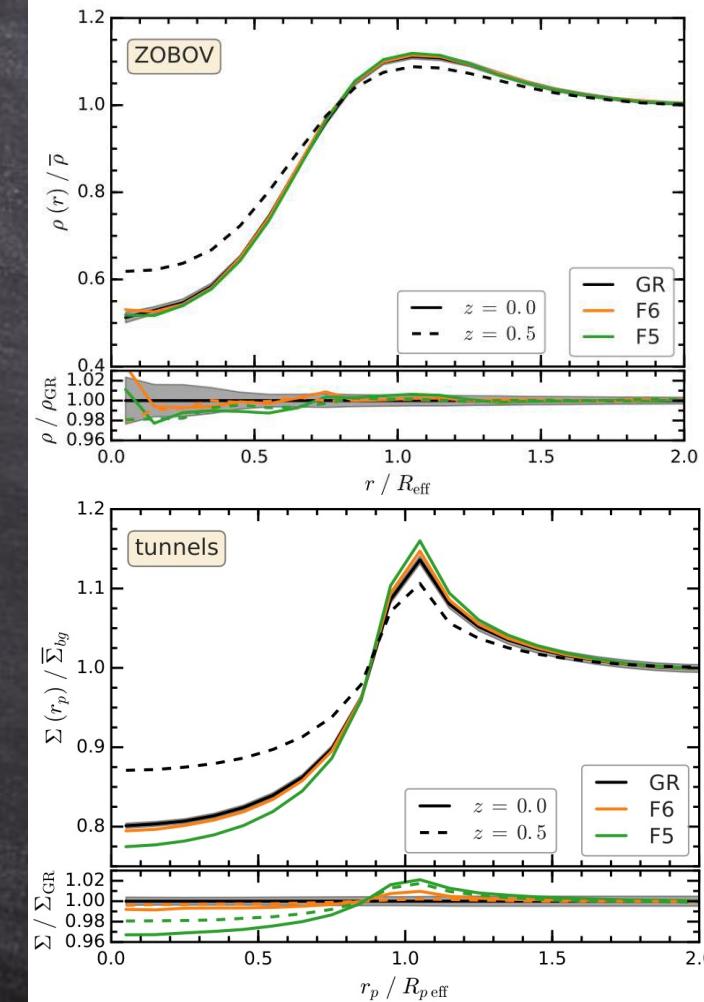
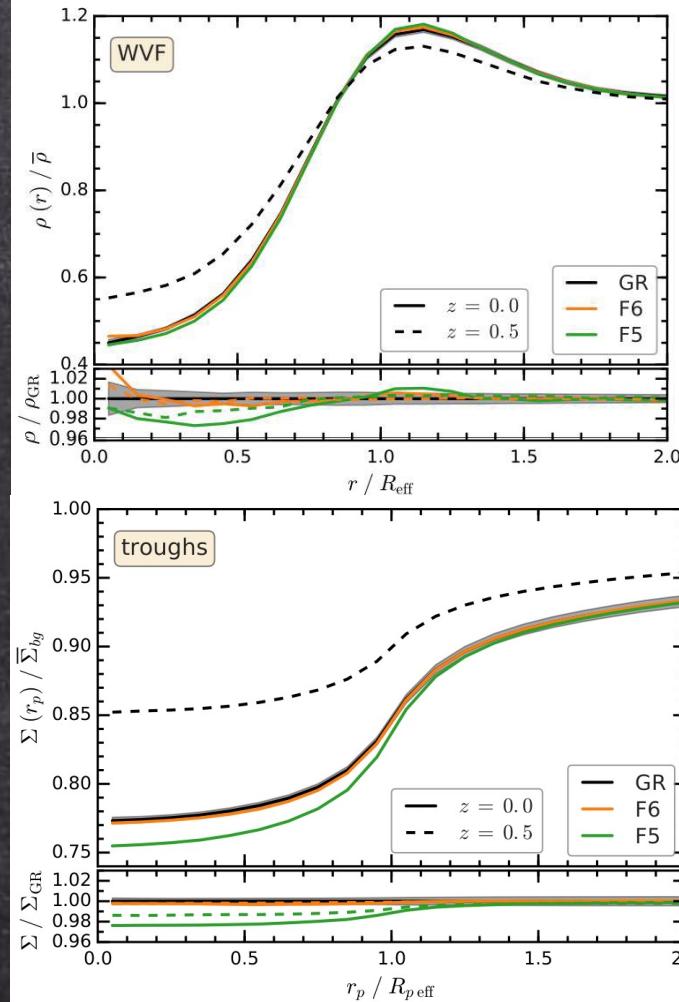
# Void stacked mass profiles

$$\begin{aligned}\rho(r) &= \bar{\rho}_{bg}(1 + \xi_{vm}; 3D) \\ \Sigma(r) &= \bar{\Sigma}_{bg}(1 + \xi_{vm}; 2D),\end{aligned}$$

3D voids

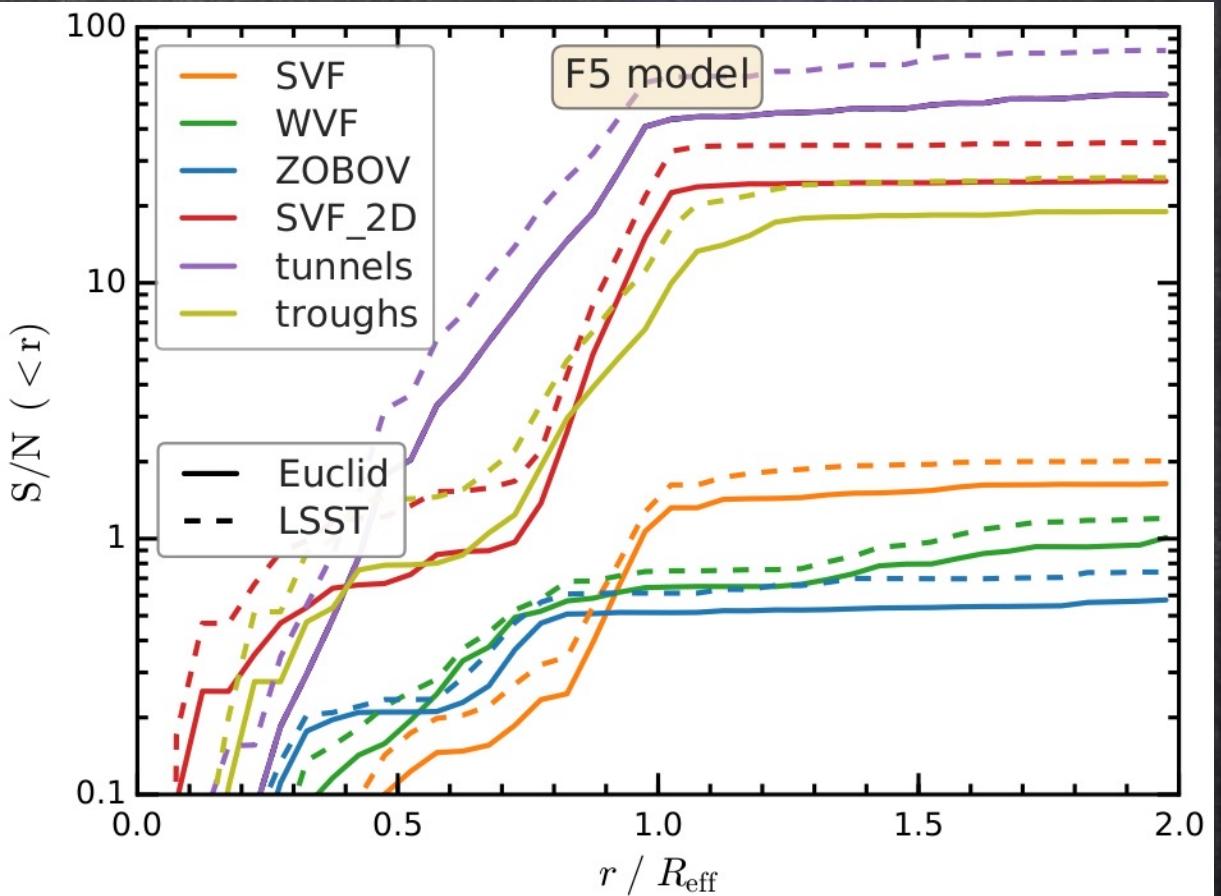


2D voids



# Testing MG with voids in the LSST era

- Lensing can be used to detect  $f(R)$  with LSST especially with 2D VF.
- DM profiles confirm emptier voids in MG models.
- These results are found for models even with the same 2-point clustering.

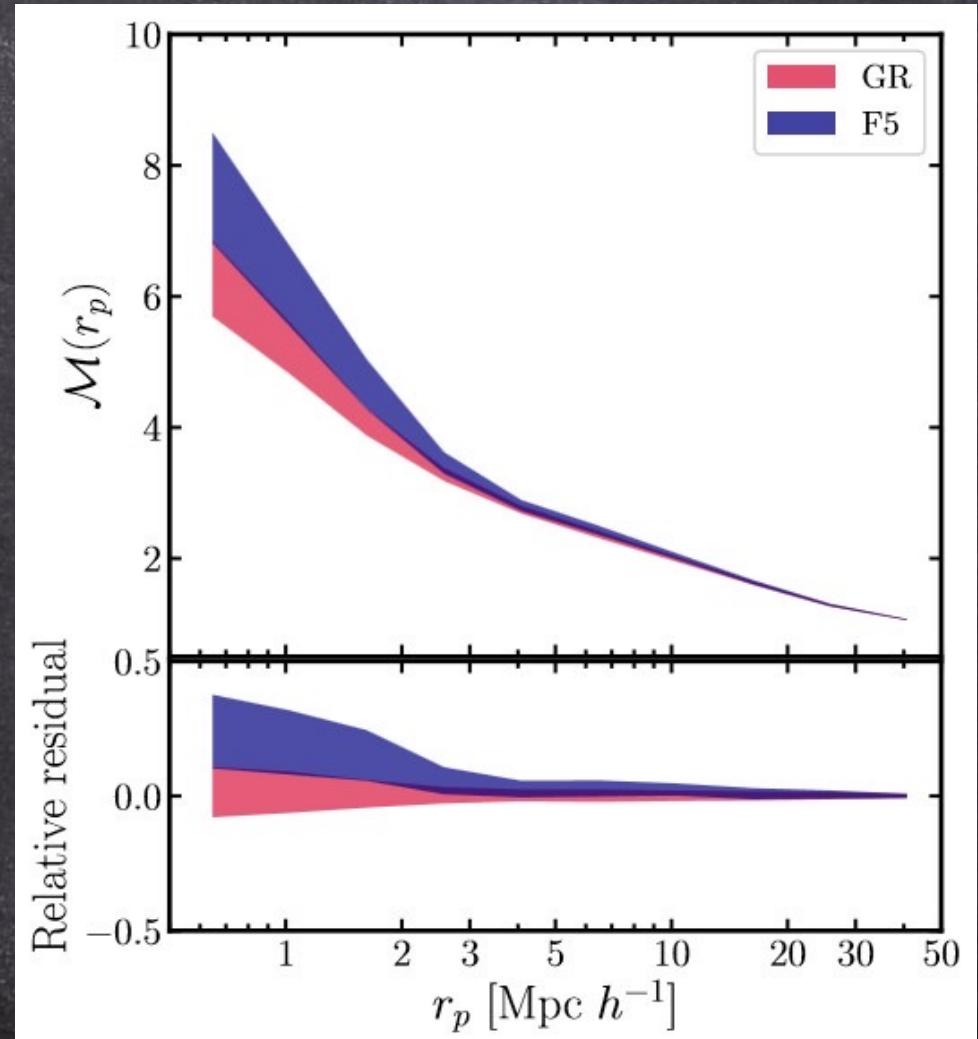


# Marked statistics

$$\mathcal{M}(r) \equiv \frac{1}{n(r)\bar{m}^2} \sum_{ij} m_i m_j = \frac{1+W}{1+\xi}$$

$$m = \left(\frac{\rho}{\bar{\rho}}\right)^p$$

- "Marks" can be used in two-point statistics to upweight models given a density environmental property.
- The marked correlation function can be used to break degeneracies in HOD modeling (White et al. 2008).
- The marked correlation function has been used to distinguish between standard GR and MG (Armijo et al. 2018).

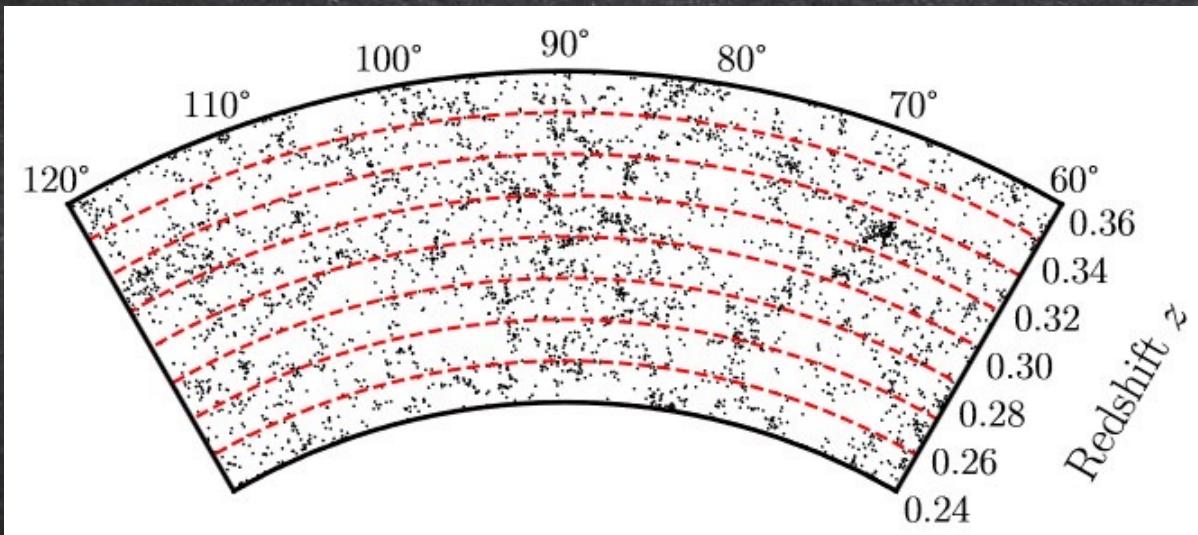


# Testing modified gravity using the marked correlation function

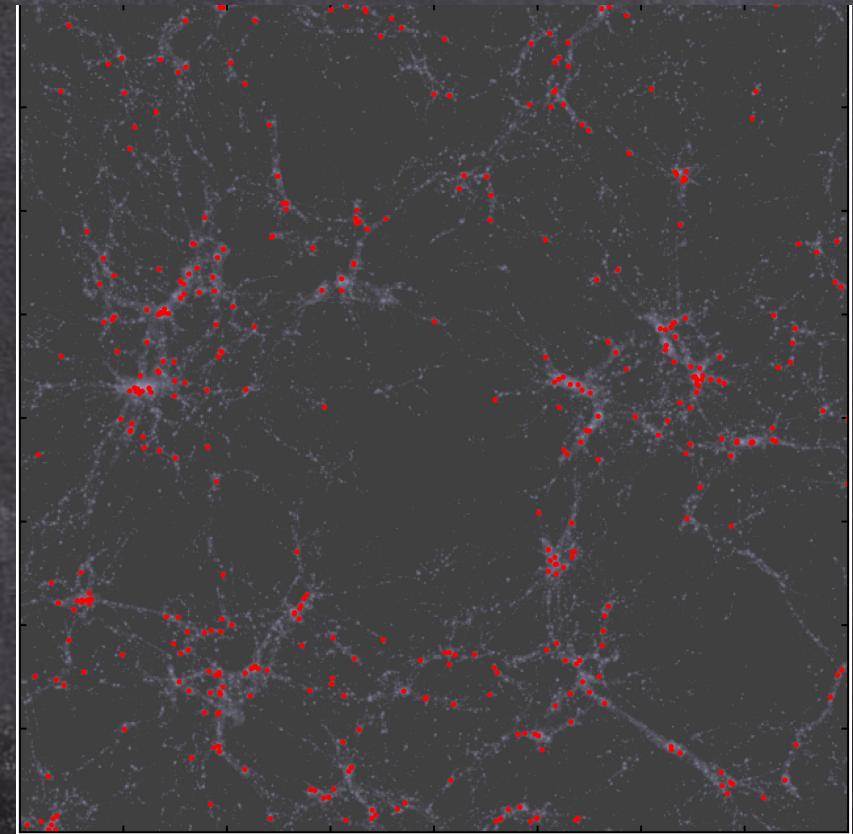
Marks:

based in the local density marks defined in White et al. (2016). We estimate the local density using the Voronoi tessellation method

Survey data: (RA, Dec, z)



Project in 2D-plane

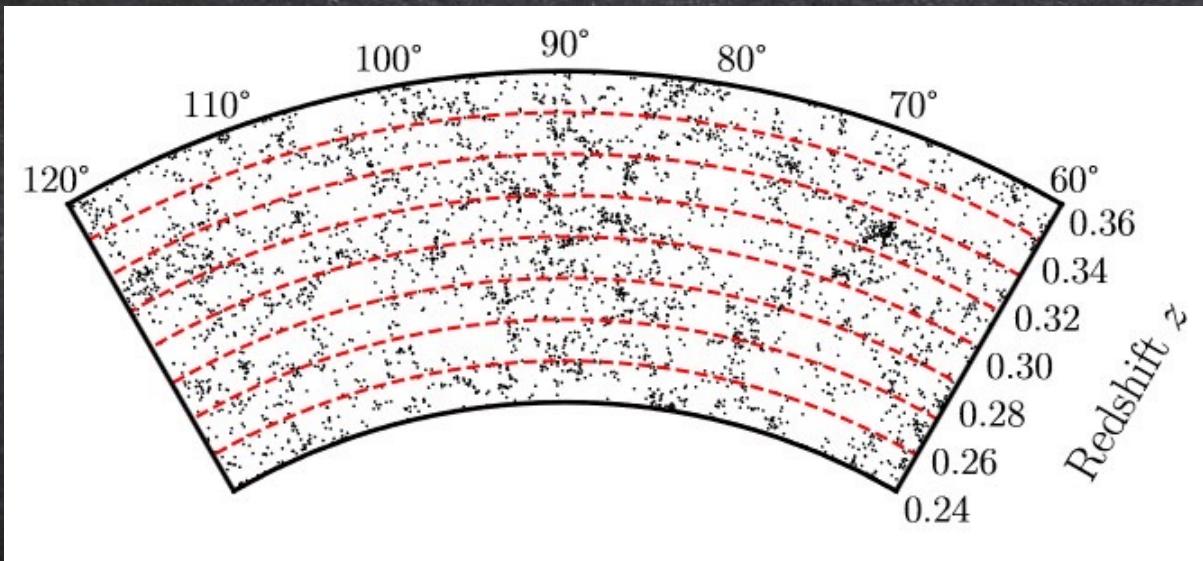


# Testing modified gravity using the marked correlation function

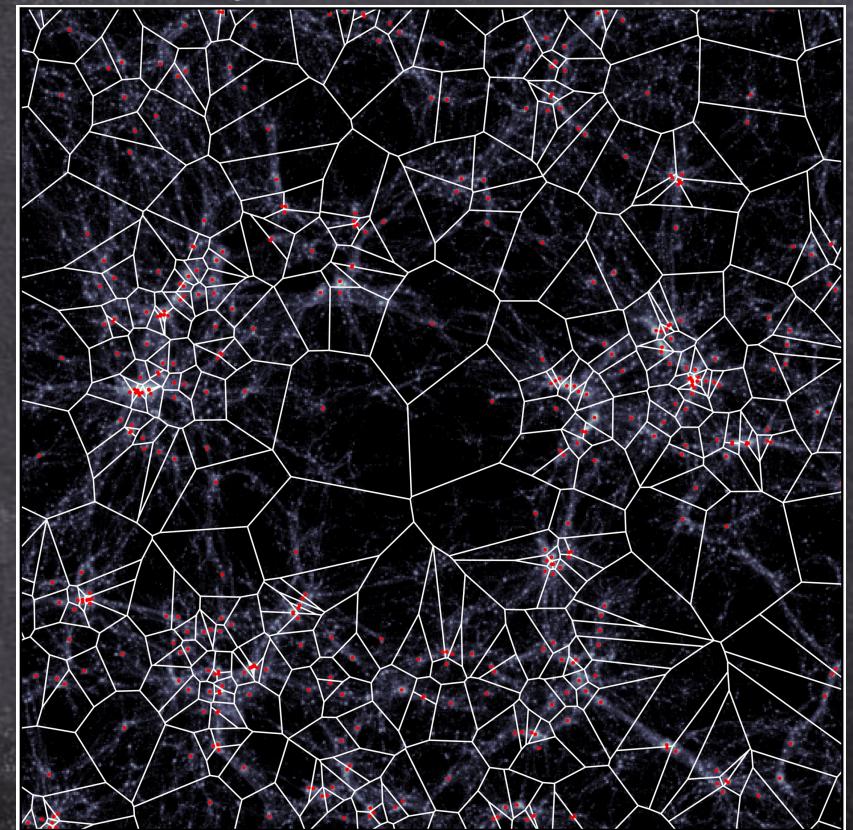
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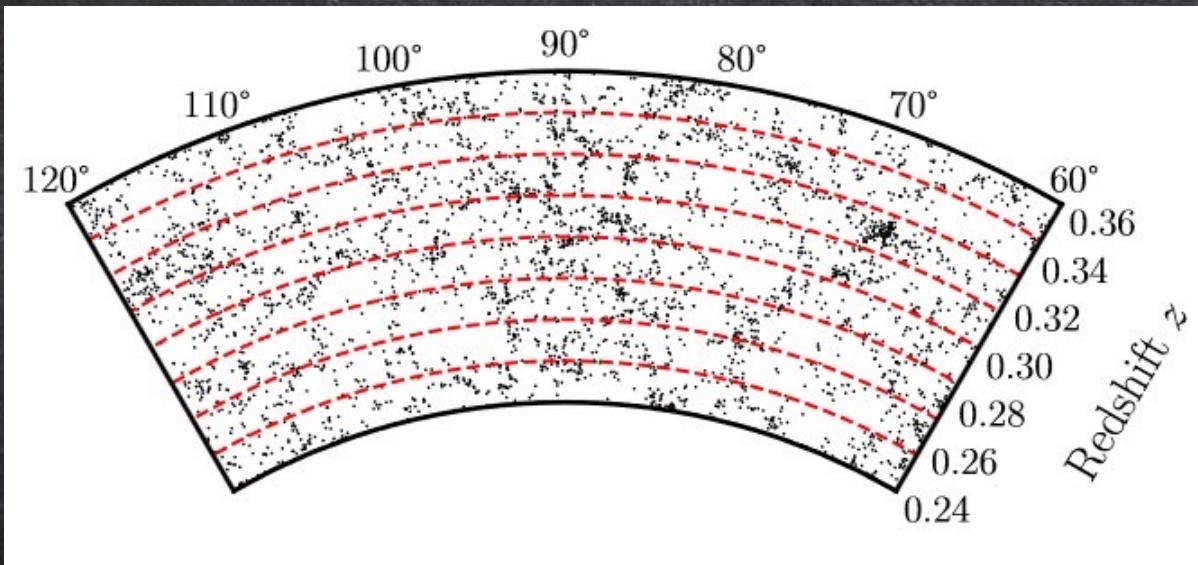


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Marks:

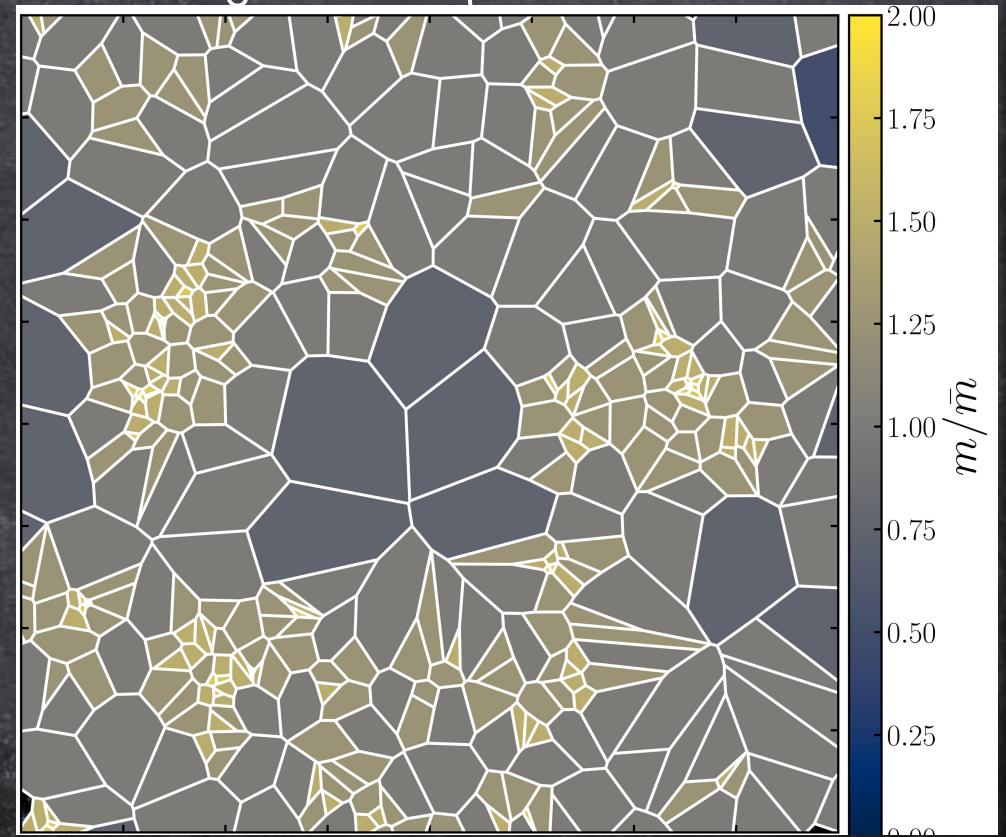
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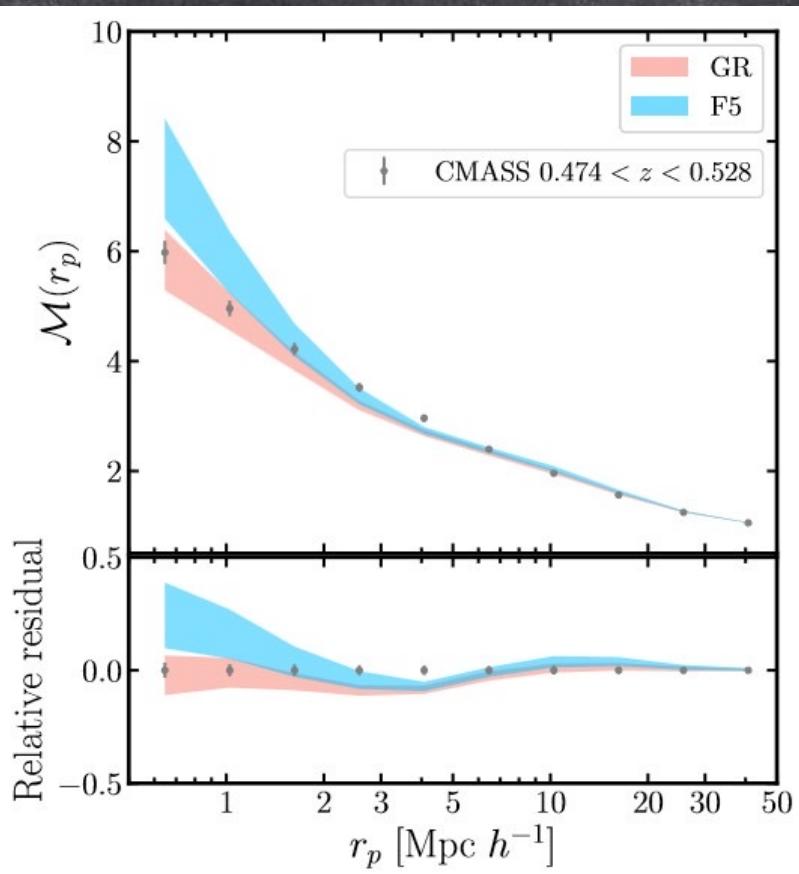
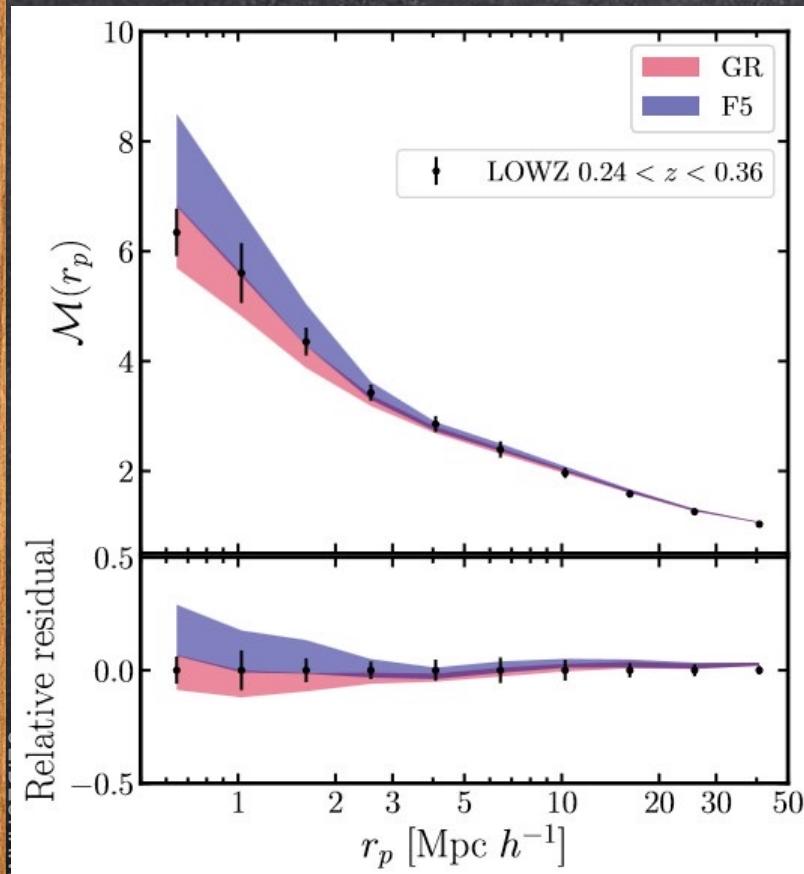


$$m = \left(\frac{\rho}{\bar{\rho}}\right)^p$$

Project in 2D-plane



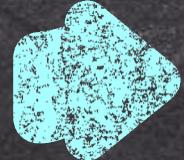
# Testing modified gravity using the marked correlation function



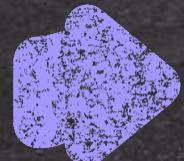
$$\mathcal{M}(r) = \frac{1 + W}{1 + \xi}$$

- We measure  $\mathcal{M}(r_p)$  for LOWZ and CMASS. We can constrain the HOD model with the data but not rule out modified gravity.
- Although CMASS agrees better with GR, some disagreement is found at  $3 < r_p/\text{Mpc } h^{-1} < 5$ .
- Future LRG sample (DESI survey) could improve this measurement.

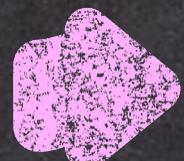
# Summary and Conclusions



In the search for the elusive camouflaged fifth force in modified gravity. Looking into the predicted **modified environments** where  $f(R)$  gravity acts.



Into the void: **emptier voids in MG models** might be a key to constraint the fifth force.



Marked statistic to test gravity on large-scales and understand the halo model. **It helps to break the degeneracy between MG and the HOD modelling.** Future data such DESI-LRG will be relevant for this test.



# Thank you!





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