

Notes on Physics from Symmetry

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This document contains my personal notes on Jakob Schwichtenberg's Physics from Symmetry ([Schwichtenberg, 2015](#)).

1. Special relativity

1.1. Definitions and postulates

In special relativity, **inertial frames of reference** are coordinate systems moving with constant velocity relative to each other. Special relativity has two basic postulates:

1. **The principal of relativity:** The laws of physics are the same in all inertial frames of reference.
2. **The invariance of the speed of light:** The velocity of light has the same value c in all inertial frames of reference.

Theorem 1.1 (Invariant of special relativity). *Consider two events A and B in an inertial observer O 's frame of reference. Let the time interval measured by O between the two events be (Δt) , and the three spatial intervals be (Δx) , (Δy) , (Δz) . Then, the quantity*

$$(\Delta s)^2 := (\Delta ct)^2 - (\Delta x)^2 - (\Delta y)^2 - (\Delta z)^2 \quad (1.1)$$

is invariant between all frames of reference. I.e.

$$(\Delta s')^2 = (\Delta s)^2 \quad (1.2)$$

for any inertial frame of reference O' .

Theorem 1.1 follows directly from the invariance of the speed of light (consider a pair of mirrors, for two observers with relative velocity).

Definition 1.1 (Proper time). *Proper time, τ , is the time measured by an observer in the special frame of reference where the object in question is at rest. In this frame of reference,*

$$(\Delta s)^2 = (c\Delta\tau)^2. \quad (1.3)$$

In the infinitesimal limit

$$(ds)^2 = (cd\tau)^2. \quad (1.4)$$

Physically, Defn. 1.1 means that all observers agree on the time interval between events for an observer who travels with the object in question. However, different observers **do not** in general agree on the time interval between events generally: $(\Delta t) \neq (\Delta t')$ – this is called **time dilation**.

1.2. c is an upper speed limit

All observers agree on the value of $(ds)^2 = (cd\tau)^2$. Furthermore, we commonly assume that there exists a minimal proper time of $\tau = 0$ for two events if $\Delta s^2 = 0$. We can therefore write that when $\tau = 0$

$$c^2 = \frac{(dx)^2 + (dy)^2 + (dz)^2}{(dt)^2} \quad (1.5)$$

between two events with an infinitesimal distance. We can equate the right-hand side with a squared velocity, and hence

$$\tau = 0 \implies c^2 = v^2 \quad (1.6)$$

so

$$(ds)^2 \geq 0 \implies c^2 \geq v^2 \quad (1.7)$$

for **any** pair of events (which are causally connected, although how this follows is not immediately clear to me right now).

1.3. Minkowski Notation

Definition 1.2 (Four-vector (contravariant)). A position four-vector is defined as

$$x^\mu = \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix} \equiv \begin{pmatrix} x^0 \\ x^1 \\ x^2 \\ x^3 \end{pmatrix}. \quad (1.8)$$

Definition 1.3 (Minkowski metric). The Minkowski metric is defined as

$$\eta_{\mu\nu} = \text{diag}(1, -1, -1, -1). \quad (1.9)$$

We define $\eta^{\mu\nu}$ through the relation

$$\eta^{\mu\nu} \eta_{\nu\sigma} = \delta^\mu_\sigma \quad (1.10)$$

where we have applied the **Einstein summation convention**, where a repeated Greek index implies a summation from 0 to 3 (where the zeroth index is time), and a repeated Roman index is summed from 1 to 3. Hence, for a matrix multiplication between two 3×3 matrices A and B , $(AB)_{ij} = A_{ik}B_{kj}$, and $(A^T)_{ij} = A_{ji}$.

Definition 1.4 (One-form (covariant vector)). We define a one-form as

$$x_\mu = \eta_{\mu\nu} x^\nu. \quad (1.11)$$

Thus,

$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu. \quad (1.12)$$

Definition 1.5 (Scalar product). A scalar product between four-vectors x and y is defined as

$$x \cdot y := x^\mu y^\nu \eta_{\mu\nu} = x_\mu y_\nu \eta^{\mu\nu} = x^\mu y_\mu = x_\nu y^\nu \quad (1.13)$$

due to the symmetry of the metric: $\eta_{\mu\nu} = \eta_{\nu\mu}$.

In general, a (q, r) tensor field can have q upper indices and r lower indices, defined through the following transformation law:

Definition 1.6 ((q, r) tensor). A (q, r) tensor is defined through the transformation law

$$T^{i'_1 i'_2 \dots i'_q}_{j'_1 j'_2 \dots j'_r}(x') = T^{i_1 i_2 \dots i_q}_{j_1 j_2 \dots j_r}(x) (M^{i'_1}_{i_1} \dots M^{i'_q}_{i_q}) (M^{j_1}_{j'_1} \dots M^{j_r}_{j'_r}). \quad (1.14)$$

1.4. Lorentz transformations

From the invariant of SR (Theorem 1.1), we have

$$ds'^2 = dx'_\mu dx'_\nu \eta^{\mu\nu} = dx_\mu dx_\nu \eta^{\mu\nu} \quad (1.15)$$

for all reference frames. We denote Λ as a transformation to move from one reference frame to another

$$dx'^\mu = \Lambda^\mu_\nu dx^\nu \quad (1.16)$$

which leaves the ds^2 invariant, i.e. $ds'^2 = ds^2$. It follows that

$$\begin{aligned} \eta_{\mu\nu} &= \Lambda^\sigma_\mu \Lambda^\delta_\nu \eta_{\sigma\delta} \\ \eta &= \Lambda^T \eta \Lambda \end{aligned} \quad (1.17)$$

References

Schwichtenberg, J., 2015 *Physics from symmetry*. Springer.