

# Notes on Physics from Symmetry

Author: Juvid Aryaman

Last compiled: May 14, 2021

This document contains my personal notes on Jakob Schwichtenberg's Physics from Symmetry ([Schwichtenberg, 2015](#)), with a sprinkling of notes from my undergraduate physics course in quantum field theory (and, to a lesser extent, general relativity).

## 1. Special relativity

### 1.1. Definitions and postulates

In special relativity, **inertial frames of reference** are coordinate systems moving with constant velocity relative to each other. Special relativity has two basic postulates:

1. **The principal of relativity:** The laws of physics are the same in all inertial frames of reference.
2. **The invariance of the speed of light:** The velocity of light has the same value  $c$  in all inertial frames of reference.

**Theorem 1.1 (Invariant of special relativity).** *Consider two events  $A$  and  $B$  in an inertial observer  $O$ 's frame of reference. Let the time interval measured by  $O$  between the two events be  $(\Delta t)$ , and the three spatial intervals be  $(\Delta x)$ ,  $(\Delta y)$ ,  $(\Delta z)$ . Then, the quantity*

$$(\Delta s)^2 := (\Delta ct)^2 - (\Delta x)^2 - (\Delta y)^2 - (\Delta z)^2 \quad (1.1)$$

*is invariant between all frames of reference. I.e.*

$$(\Delta s')^2 = (\Delta s)^2 \quad (1.2)$$

*for any inertial frame of reference  $O'$ .*

Theorem 1.1 follows directly from the invariance of the speed of light (consider a pair of mirrors, for two observers with relative velocity).

**Definition 1.1 (Proper time).** *Proper time,  $\tau$ , is the time measured by an observer in the special frame of reference where the object in question is at rest. In this frame of reference,*

$$(\Delta s)^2 = (c\Delta\tau)^2. \quad (1.3)$$

*In the infinitesimal limit*

$$(ds)^2 = (cd\tau)^2. \quad (1.4)$$

Physically, Defn. 1.1 means that all observers agree on the time interval between events for an observer who travels with the object in question. However, different observers **do not** in general agree on the time interval between events generally:  $(\Delta t) \neq (\Delta t')$  – this is called **time dilation**.

### 1.2. $c$ is an upper speed limit

All observers agree on the value of  $(ds)^2 = (cd\tau)^2$ . Furthermore, we commonly assume that there exists a minimal proper time of  $\tau = 0$  for two events if  $\Delta s^2 = 0$ . We can therefore write that when  $\tau = 0$

$$c^2 = \frac{(dx)^2 + (dy)^2 + (dz)^2}{(dt)^2} \quad (1.5)$$

between two events with an infinitesimal distance. We can equate the right-hand side with a squared velocity, and hence

$$\tau = 0 \implies c^2 = v^2 \quad (1.6)$$

so

$$(ds)^2 \geq 0 \implies c^2 \geq v^2 \quad (1.7)$$

for **any** pair of events (which are causally connected, although how this follows is not immediately clear to me right now).

### 1.3. Minkowski Notation

**Definition 1.2 (Four-vector (contravariant)).** A position four-vector is defined as

$$x^\mu = \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix} \equiv \begin{pmatrix} x^0 \\ x^1 \\ x^2 \\ x^3 \end{pmatrix}. \quad (1.8)$$

**Definition 1.3 (Minkowski metric).** The Minkowski metric is defined as

$$\eta_{\mu\nu} = \text{diag}(1, -1, -1, -1). \quad (1.9)$$

We define  $\eta^{\mu\nu}$  through the relation

$$\eta^{\mu\nu} \eta_{\nu\sigma} = \delta^\mu_\sigma \quad (1.10)$$

where we have applied the **Einstein summation convention**, where a repeated Greek index implies a summation from 0 to 3 (where the zeroth index is time), and a repeated Roman index is summed from 1 to 3. Hence, for a matrix multiplication between two  $3 \times 3$  matrices  $A$  and  $B$ ,  $(AB)_{ij} = A_{ik}B_{kj}$ , and  $(A^T)_{ij} = A_{ji}$ .

**Definition 1.4 (One-form (covariant vector)).** We define a one-form as

$$x_\mu = \eta_{\mu\nu} x^\nu. \quad (1.11)$$

Thus,

$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu. \quad (1.12)$$

**Definition 1.5 (Scalar product).** A scalar product between four-vectors  $x$  and  $y$  is defined as

$$x \cdot y := x^\mu y^\nu \eta_{\mu\nu} = x_\mu y_\nu \eta^{\mu\nu} = x^\mu y_\mu = x_\nu y^\nu \quad (1.13)$$

due to the symmetry of the metric:  $\eta_{\mu\nu} = \eta_{\nu\mu}$ .

In general, a  $(q, r)$  tensor field can have  $q$  upper indices and  $r$  lower indices, defined through the following transformation law:

**Definition 1.6 ( $(q, r)$  tensor).** A  $(q, r)$  tensor is defined through the transformation law

$$T^{i'_1 i'_2 \dots i'_q}_{j'_1 j'_2 \dots j'_r}(x') = T^{i_1 i_2 \dots i_q}_{j_1 j_2 \dots j_r}(x) (M^{i'_1}_{i_1} \dots M^{i'_q}_{i_q}) (M^{j_1}_{j'_1} \dots M^{j_r}_{j'_r}). \quad (1.14)$$

#### 1.4. Lorentz transformations

From the invariant of SR (Theorem 1.1), we have

$$ds'^2 = dx'_\mu dx'_\nu \eta^{\mu\nu} = dx_\mu dx_\nu \eta^{\mu\nu} \quad (1.15)$$

for all reference frames. We denote  $\Lambda$  as a (1,1) tensor field, which transforms a four-vector from one reference frame to another:

$$dx'^\mu = \Lambda^\mu_\nu dx^\nu \quad (1.16)$$

which leaves the  $ds^2$  invariant, i.e.  $ds'^2 = ds^2$ . It follows that

$$\begin{aligned} \eta_{\mu\nu} &= \Lambda^\sigma_\mu \Lambda^\delta_\nu \eta_{\sigma\delta} \\ \eta &= \Lambda^T \eta \Lambda \end{aligned} \quad (1.17)$$

## References

Schwichtenberg, J., 2015 *Physics from symmetry*. Springer.