Solutions for assignment 2:

1. Beiser; Chap. 5: Prob 9: Show that the expectation values  $\langle xp \rangle$  and  $\langle px \rangle$  are related by

$$\langle px \rangle - \langle xp \rangle = \frac{\hbar}{i}$$

This result is described by saying that p and x do not commute and it is intimately related to the uncertainty principle,

Ans. The quantity shown is the expectation value of a special operator called a commutator of operators  $\hat{p}$  and  $\hat{x}$  and is defined as

$$[\hat{p}, \ \hat{x}] = \hat{p}\hat{x} - \hat{x}\hat{p}$$

Let us calculate the expectation values for a normalised wavefunction  $\psi(x)$ ,

$$\langle px \rangle - \langle xp \rangle = \int \psi(x)^* \hat{p}\hat{x}\psi(x)dx - \int \psi(x)^* \hat{x}\hat{p}\psi(x)dx$$

Using  $\hat{p} = \frac{\hbar}{i} \frac{\partial}{\partial x}$ , we have,

$$\langle px \rangle - \langle xp \rangle = \frac{\hbar}{i} \left[ \int \psi(x)^* \frac{\partial}{\partial x} [x\psi(x)] dx - \int \psi(x)^* x \frac{\partial}{\partial x} \psi(x) dx \right]$$
$$= \frac{\hbar}{i} \left[ \int \psi(x)^* x \frac{\partial}{\partial x} \psi(x) + \int \psi(x)^* \psi(x) dx - \int \psi(x)^* x \frac{\partial}{\partial x} \psi(x) dx \right] = \frac{\hbar}{i}$$

The first and the third term cancel out and the second term is equal to one (since the function is normalised.

2. Atkins: 9.12: Confirm that a function of the form  $e^{-gx^2}$  is a solution of the Schrodinger equation for the ground state of a harmonic oscillator and find an expression for g in terms of the mass and force constant of the oscillator.

Ans. We need to show :  $\hat{H}e^{-gx^2} = (\text{eigen value})e^{-gx^2}$ 

$$\begin{split} \hat{H}e^{-gx^2} &= \left(-\frac{\hbar^2}{2\mu}\frac{d^2}{dx^2} + \frac{1}{2}kx^2\right)e^{-gx^2} = -\frac{\hbar^2}{2\mu}\frac{d^2}{dx^2}\left(e^{-gx^2}\right) + \frac{1}{2}kx^2e^{-gx^2} \\ &= -\frac{\hbar^2}{2\mu}(-2g)\left[\frac{d}{dx}\left(xe^{-gx^2}\right)\right] + \frac{1}{2}kx^2e^{-gx^2} = -\frac{\hbar^2}{2\mu}\left[4g^2x^2\left(e^{-gx^2}\right) - 2g\left(e^{-gx^2}\right)\right] + \frac{1}{2}kx^2e^{-gx^2} \end{split}$$

If we can get rid of the  $x^2 \left(e^{-gx^2}\right)$  term, then we have proved our case. For this

$$\frac{4g^2\hbar^2}{2\mu} = \frac{1}{2}k \implies g = \frac{1}{2}\sqrt{\frac{\mu k}{\hbar^2}}$$

Verify if this is correct by looking back at what we got in the class.

3. Atkins: 9.20: Evaluate the z-component of the angular momentum and the kinetic energy of a particle on a ring that is described by the (unnormalized) wavefunctions (a)  $e^{i\phi}$  (b)  $e^{-2i\phi}$ , (c)  $\cos\phi$ , and (d)  $\cos\chi e^{i\phi} + \sin\chi e^{-i\phi}$ .

Ans. The operator for z-component of the angular momentum,  $\hat{l}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$ 

The values for a system with wavefunction  $\psi$ ,  $\langle \hat{l}_z \rangle = \frac{\int \psi^*(\phi) \hat{l}_z \psi(\phi) d\phi}{\int \psi^*(\phi) \psi(\phi) d\phi}$  are as follows:

(a) 
$$\frac{\int_0^{2\pi} e^{-i\phi} \frac{\hbar}{i} \frac{\partial}{\partial \phi} \left(e^{i\phi}\right) d\phi}{\int_0^{2\pi} e^{-i\phi} e^{i\phi} d\phi} = \hbar; \text{ (b) } \frac{\int_0^{2\pi} e^{2i\phi} \frac{\hbar}{i} \frac{\partial}{\partial \phi} \left(e^{-2i\phi}\right) d\phi}{\int_0^{2\pi} e^{2i\phi} e^{-2i\phi} d\phi} = -2\hbar;$$

(c)  $\cos\phi=\frac{e^{i\phi}+e^{-i\phi}}{2}$ , Thus we can see  $\cos\phi$  as a superposition of two different functions. Even though we have an unnormalized function, we can see easily that the contribution of the two components is equal, since each appears with a coefficient of  $\frac{1}{2}$ . We can also see from part (a) that for the state  $e^{-i\phi}$ , the value of  $\langle \hat{l}_z \rangle = -\hbar$ . Thus, for  $\cos\phi$ , we can write,

$$\langle \hat{l}_z \rangle = \frac{1}{2} \left[ \langle \hat{l}_z \rangle_{e^{i\phi}} + \langle \hat{l}_z \rangle_{e^{-i\phi}} \right] = 0.$$

Part (d) is a generalised situation from part (c), where the the contributions of the two components  $e^{i\phi}$  and  $e^{-i\phi}$  are not equal, but is the ratio of  $\sin^2 \chi$  and  $\cos^2 \chi$  (remember the property depends on the probability density and the density is the square of the function).

$$\therefore \langle \hat{l}_z \rangle = \cos^2 \chi \langle \hat{l}_z \rangle_{e^{i\phi}} + \sin^2 \chi \langle \hat{l}_z \rangle_{e^{-i\phi}} = \hbar (\cos^2 \chi - \sin^2 \chi) = \hbar \cos 2\chi.$$

- 4. Atkins: 9.22: This problem was about writing/using a code. We can discuss this in class.
- 5. Atkins: 9.26: Show that the function  $f = \cos ax \cos by \cos cz$  is an eigenfunction of  $\nabla^2$  and determine its eigenvalue.

Ans. 
$$\nabla^2 f = \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right) \cos ax \cdot \cos by \cdot \cos cz = -(a^2 + b^2 + c^2) \cos ax \cdot \cos by \cdot \cos cz$$
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