

Enhanced Power Factor in Nanocomposite Chalcogenides

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L.M. Woods, N. Crane and G.S. Nolas**

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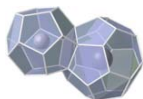
USAMRMC



Marlow Industries



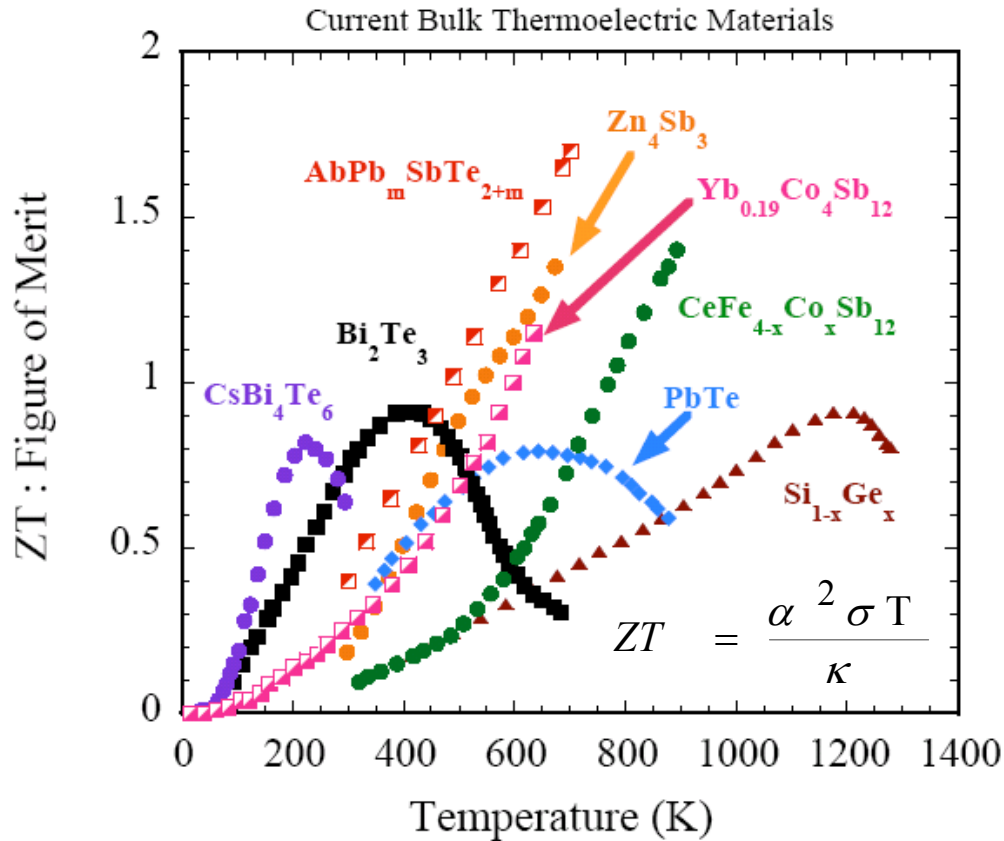
2009 DOE Thermoelectric Applications Workshop



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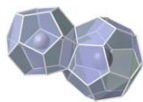
How can we increase ZT?



Thus far two approaches have shown the greatest promise:

1. **Phonon-Glass
Electron-Crystal**
(*G. Slack*)
2. **Low Dimensional
Materials**
(*M. Dresselhaus*)

T.M. Tritt and M.A. Subramanian,
MRS Bulletin **31**, 188 (March 2006)



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TE Enhancement : Nanostructured Materials

- o **Multilayered Structures**
first proposed in 1987^A
- o **One dimensional quantum wires^B**
- o **Bi₂Te₃/Sb₂Te₃ supperlattice: ZT=2.4^C**
- o **Quantum dot supperlattice: ZT =1.6^D**
- o **Metal-based supperlattices^E**
- o **Ag_mPb_mSbTe_{m+2} with nanoscale PbTe²: κ reduction**
- o **Nanostructured SiGe bulk alloys³: κ reduction**
- o **Nanostructured Bi₂Te₃ bulk alloys⁴: κ reduction**
- o **Nanocrystalline CoSb₃⁵: κ reduction**
- o **PbTe with Pb nanoprecipitates⁶ : S increase**
- o **PbTe nanocomposites^{7, 8}: S increase**
- o **PbTe ‘dimensional’ nanocomposites^{9, 10}: S increase**

^A T.E. Whall, *Proc. First Eur. Conf. on Therm.*,
D.M. Rowe (Peter Peregrinus Lt., London, 1987)

^B L.D. Hicks & M.S. Dresselhaus, PRB **47**,
16631 (1993)

^C R. Venkatasubramanian et al, Nature **413**,
597 (2001)

^D O.L. Lazarenkova & A.A. Balandin, PRB **66**,
245319 (2002)

^E D. Vashaee and A. Shakouri, PRL **92**,
106103 (2004)

¹ K.F. Hsu *et al*, Science **303**, 818 (2004)

² G. Joshi *et al*, Nano Lett. **8**, 4670 (2008)

³ M.S. Toprak *et al*, Adv. Func. Mat. **14**, 1189 (2004)

⁴ B. Poudel *et al*, Science **320**, 634 (2008)

⁵ X. Ji, T.M. Tritt *et al*, Phys. Stat. Sol. RRL **1**, 229 (2009)

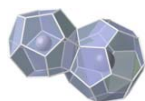
⁶ J.P. Heremans *et al*, JAP **98**, 063703 (2005)

⁷ J.P. Heremans *et al*, Phys. Rev. B **70**, 115334 (2004)

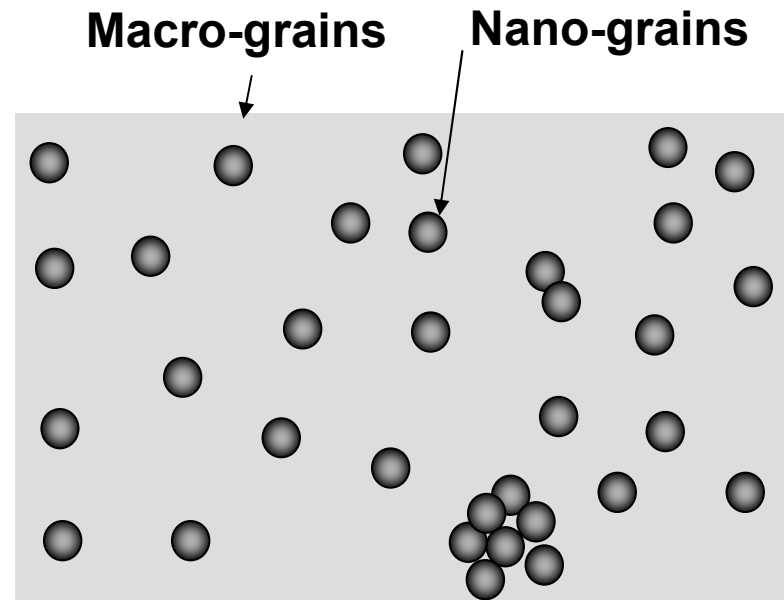
⁸ K. Kishimoto *et al*, JAP **92**, 2544 (2002)

⁹ J. Martin *et al*, Appl. Phys. Lett. **90**, 222112 (2007)

¹⁰ J. Martin *et al*, Phys. Rev. B **79**, 115311 (2009)



‘Dimensional’ Nanocomposites

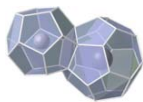


Expectations:

- Reduced thermal conductivity?
- Enhanced thermopower?
- Increased ZT?
- Cheap, self-assembly method

First Steps:

- Synthesis of high yield nanocrystals
- Densify into nanocomposites
- Measure transport properties

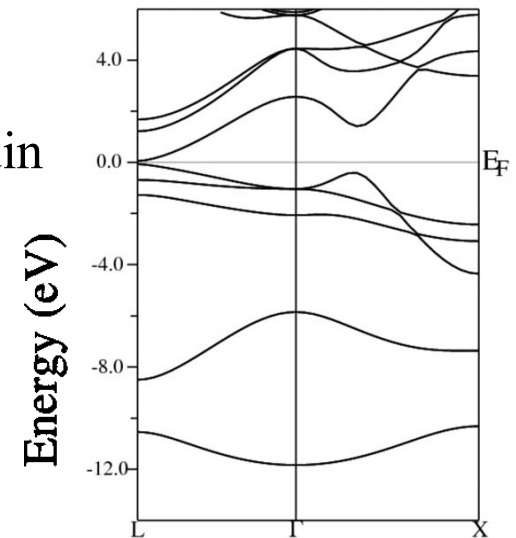
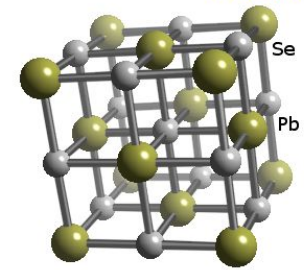
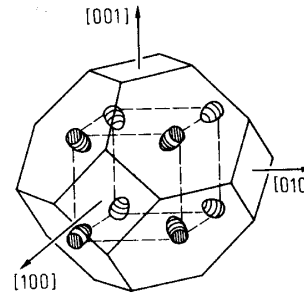


Lead Chalcogenides

PbS, PbSe and PbTe

- NaCl-like FCC lattice structure
- Relatively low thermal conductivity
- Direct-gap semiconductor at the *L*-point of the Brillouin zone
- Bandgap between conduction and valence bands:
PbS > PbTe > PbSe
- Lattice parameter:
PbS < PbSe < PbTe

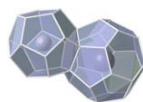
$$\frac{\hbar^2 k_l^2}{2m_l^*} + \frac{\hbar^2 k_t^2}{2m_t^*} = E + E^2 / E_g$$



PbTe energy bandstructure

X. Gao and M.S. Daw,
Phys Rev B 77, 033103 (2008)

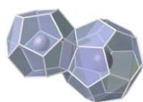
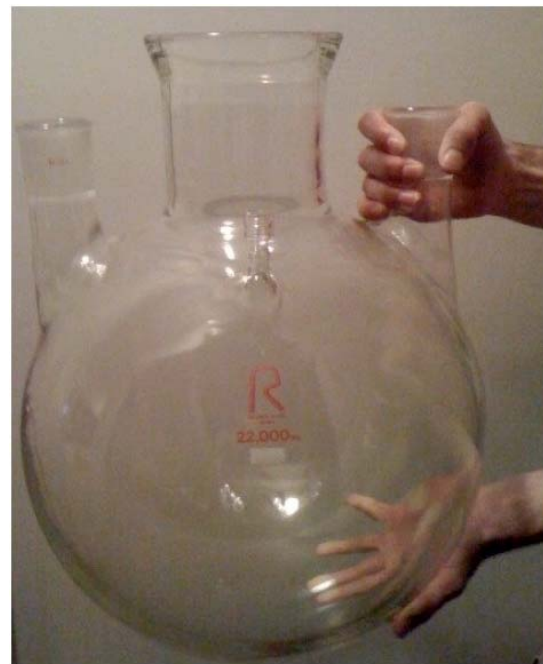
Bulk Property	PbS	PbSe	PbTe
Lattice Parameter (Å)	5.94	6.12	6.46
E_g at 77 K (meV)	307	168	215



Requirement: High Yield Nanocrystal Synthesis



Scalable!

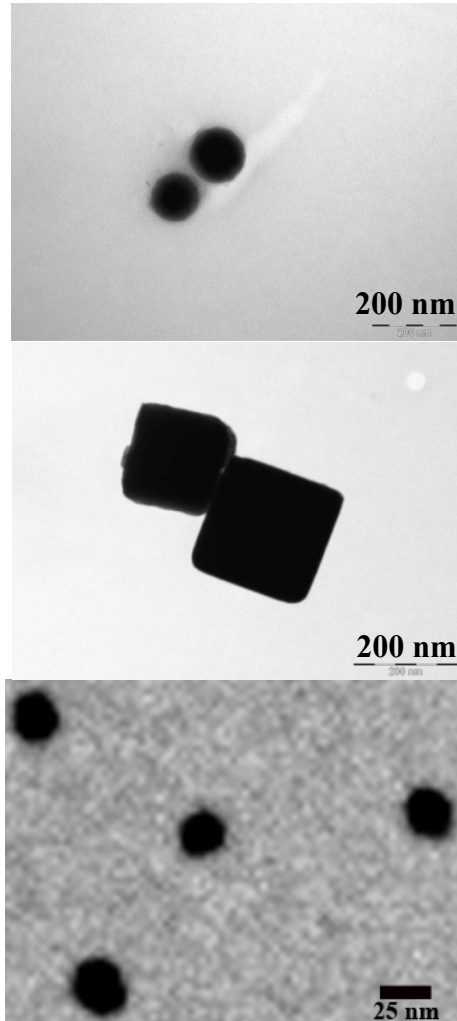


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PbTe Nanocomposites

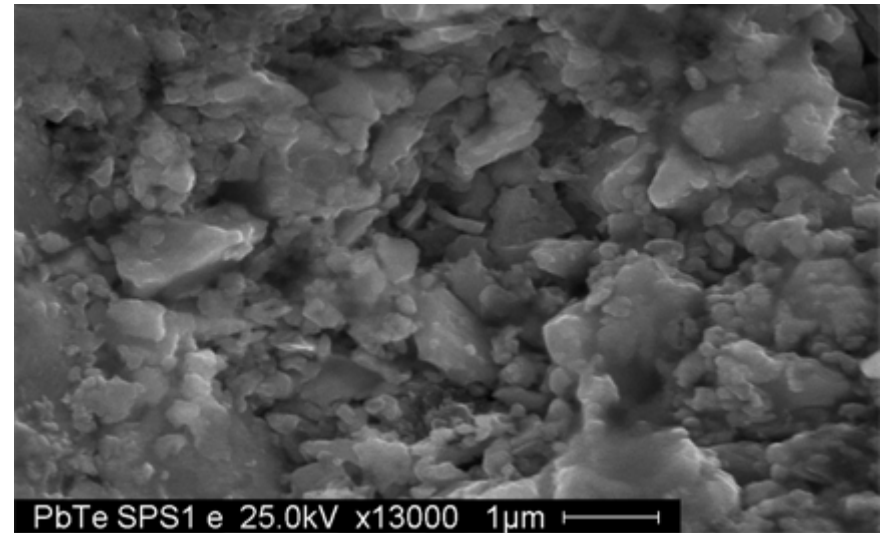
Size & Shape Selectivity



SPS



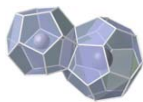
Densification



PbTe nanocrystals within the bulk matrix

J. Martin, G. S. Nolas, W. Zhang, and L. Chen,
Appl. Phys. **90**, 222112 (2007)

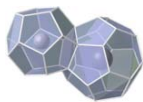
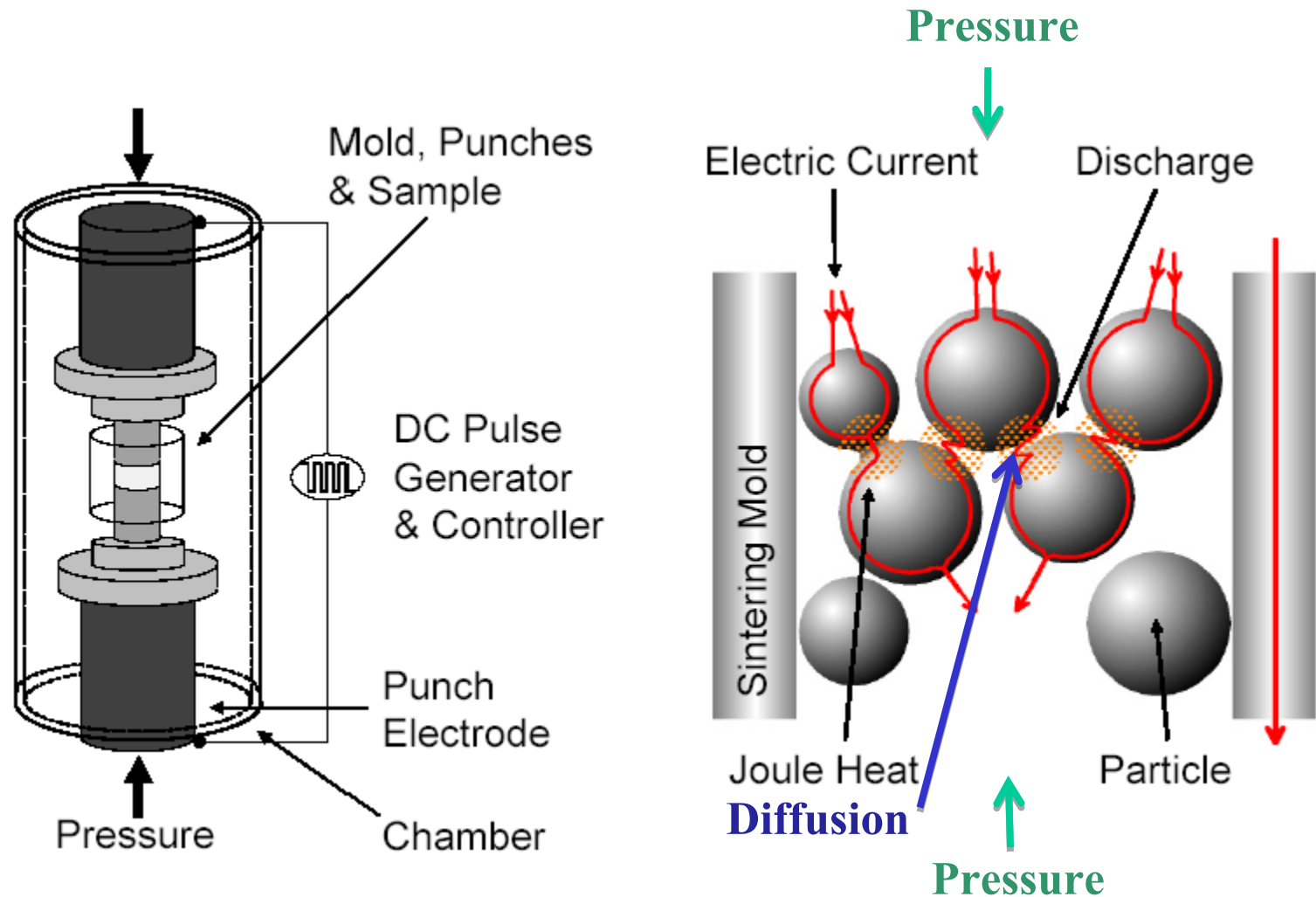
J. Martin, L. Wang, L. Chen, and G.S. Nolas,
Phys. Rev. B **79**, 115311 (2009)



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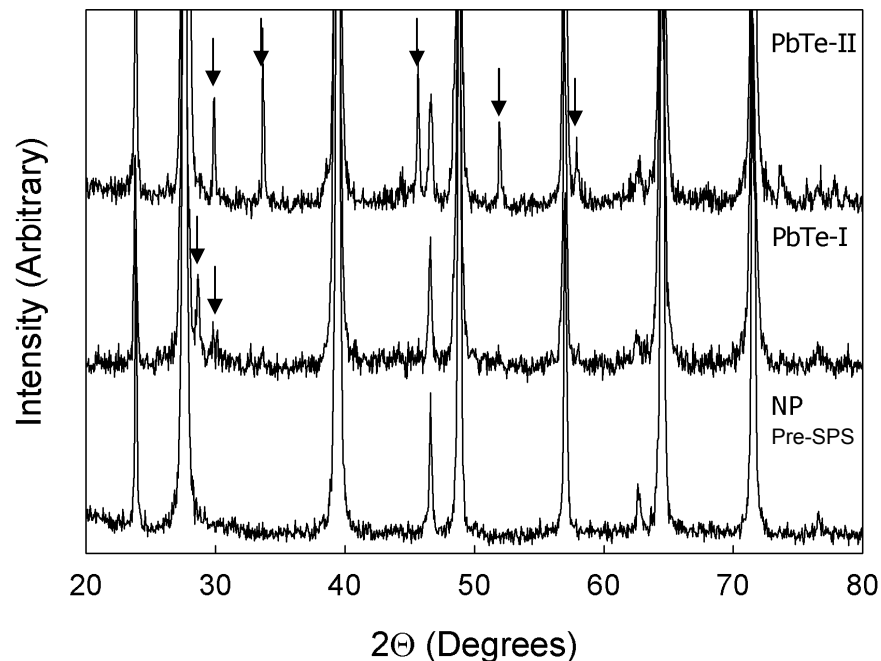


Spark Plasma Sintering

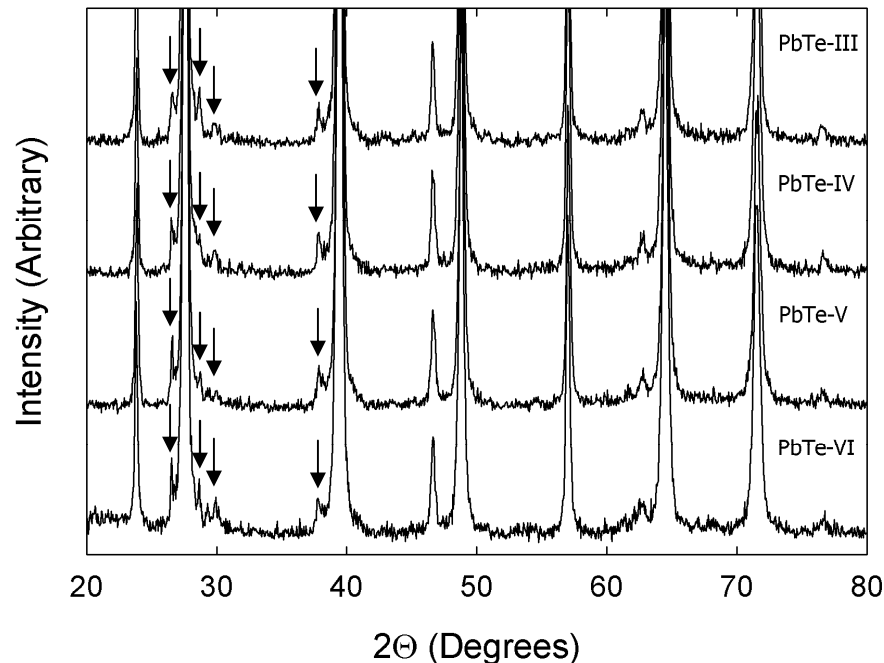


Powder and Polycrystalline XRD

PbTe



Ag-doped PbTe

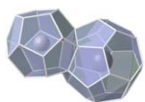


PbTeO₃ impurity

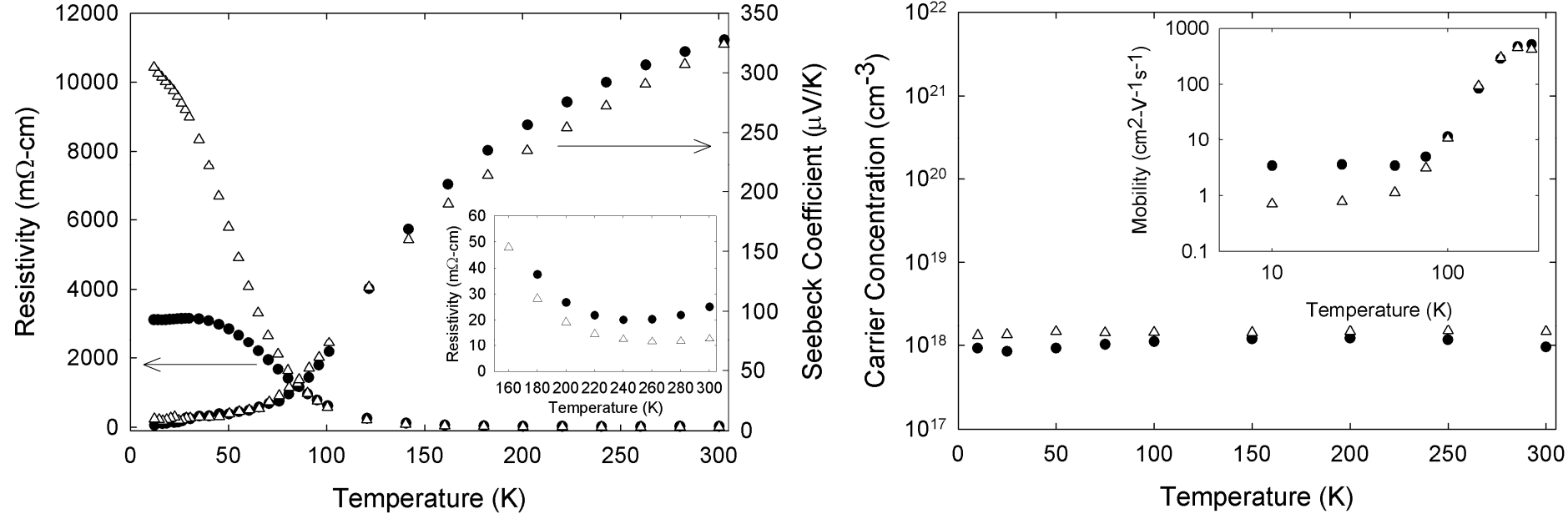
TeO₂ impurity

J. Martin, G. S. Nolas, W. Zhang, and L. Chen, Appl. Phys. **90**, 222112 (2007)

J. Martin, L. Wang, L. Chen, and G.S. Nolas, Phys. Rev. B **79**, 115311 (2009)



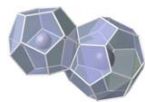
PbTe Nanocomposites



@ 300 K

Specimen	D (%)	ρ (m Ω -cm)	S (μ V/K)	κ (W-m ⁻¹ K ⁻¹)	p (cm ⁻³)	$S^2\sigma$ (μ W/K ² cm)	Pb:Te	
PbTe-I ●	94	24.9	328	2.2	9.5×10^{17}	4.3	49.91 : 50.09	
PbTe-II Δ	95	12.6	324	2.5	1.5×10^{18}	8.3	50.42 : 49.58	
B-I	Bulk	97	37	325	-	8.0×10^{17}	2.9	-
B-II		97	19	250	-	9.5×10^{17}	3.3	-

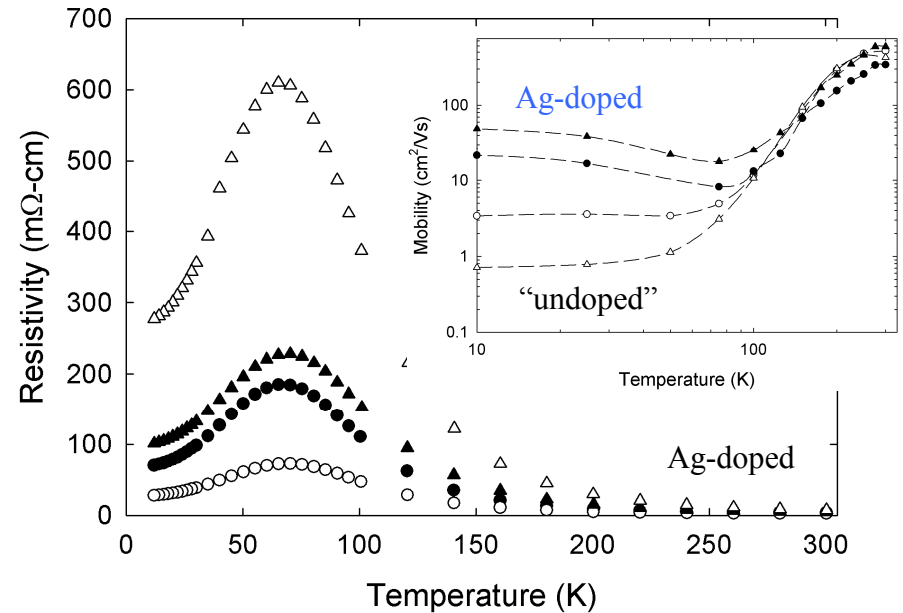
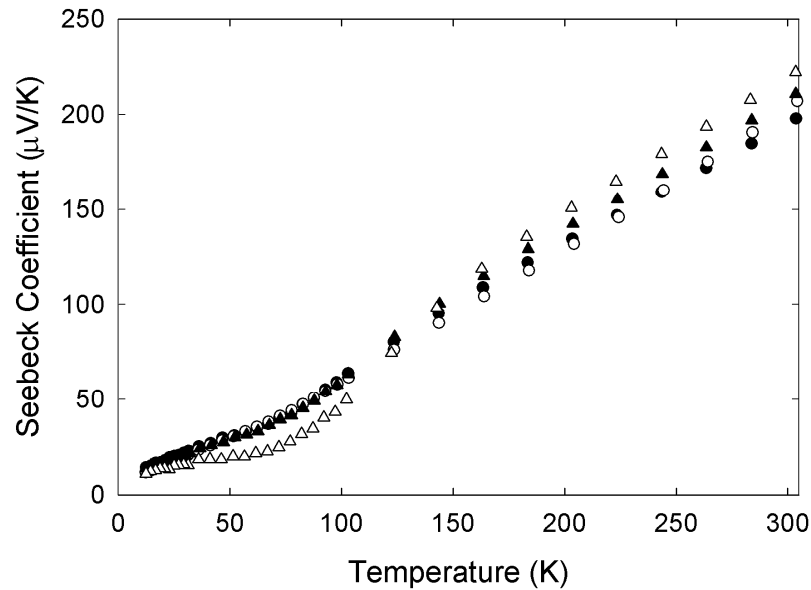
J. Martin, G. S. Nolas, W. Zhang, and L. Chen, Appl. Phys. **90**, 222112 (2007)



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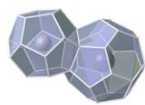
PbTe Nanocomposites



@ 300 K

Specimen	D (%)	ρ ($\text{m}\Omega\text{-cm}$)	S ($\mu\text{V/K}$)	z ($\text{W-m}^{-1}\text{K}^{-1}$)	p (cm^{-3})	$S^2\sigma$ ($\mu\text{W/K}^2\text{cm}$)	Pb:Te
PbTe-I	94	24.9	328	2.2	9.5×10^{17}	4.3	49.91 : 50.09
PbTe-II	95	12.6	324	2.5	1.5×10^{18}	8.3	50.42 : 49.58
PbTe-III	95	3.9	198	2.8	5.1×10^{18}	10.0	-
PbTe-IV	94	2.9	207	2.7	6.2×10^{18}	14.8	-

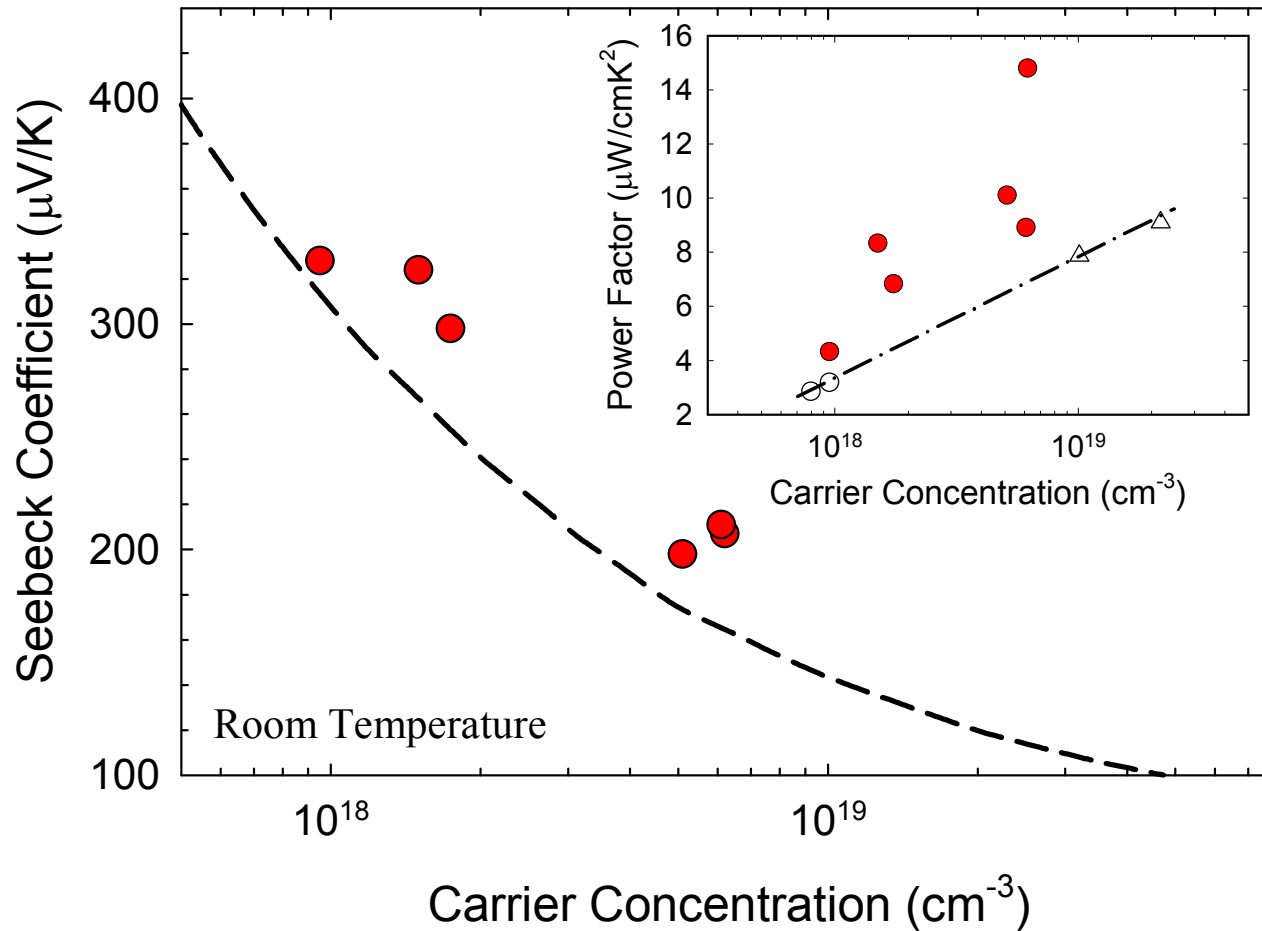
J. Martin, L. Wang, L. Chen, and G.S. Nolas, Phys. Rev. B **79**, 115311 (2009)



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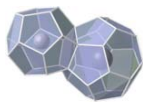
Thermopower Enhancement in Nanocomposites



red PbTe Nanocomposites \triangle , \circ Bulk PbTe

----- A.J. Crocker *et al.*, Brit. J. Appl. Phys. **18**, 563 (1967)

H. Kirby, J. Martin, L. Chen, and G.S. Nolas,
Proc. Mater. Res. Soc. **1166** (2009)



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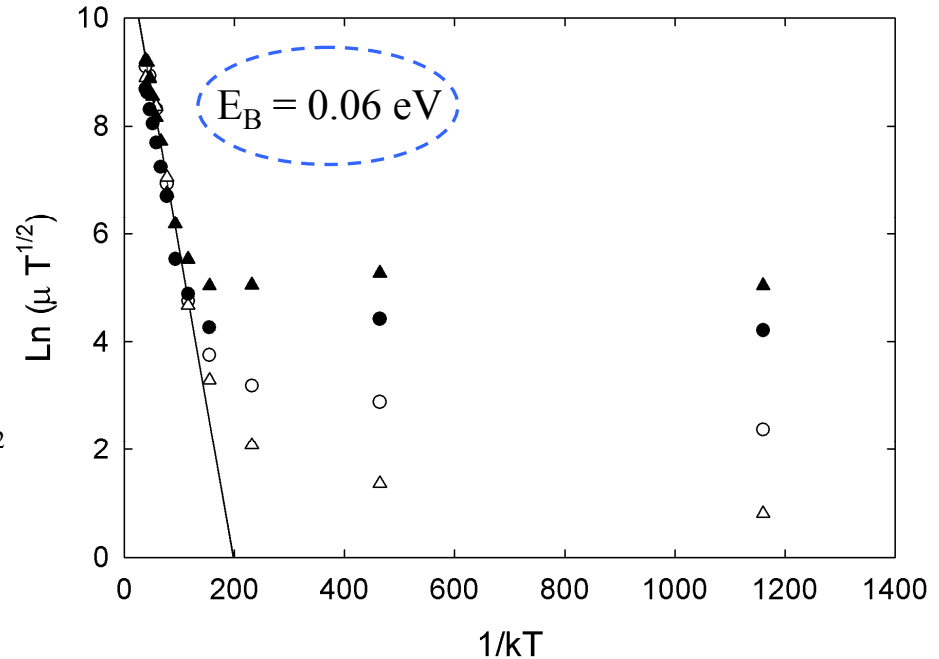


PbTe Nanocomposites

$$\mu_{eff} = Lq \left(\frac{1}{2\pi m^* kT} \right)^{1/2} \exp\left(-\frac{E_B}{kT}\right)$$

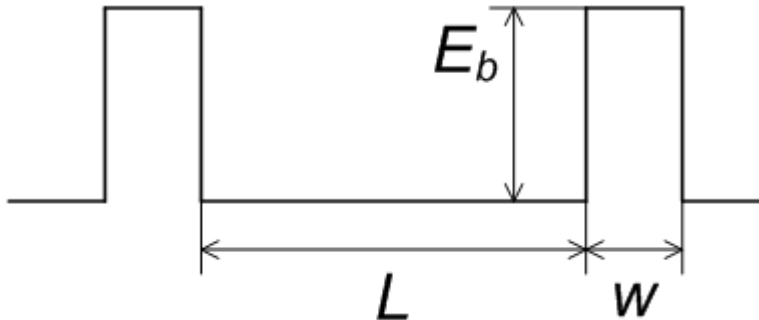
Density of trapping states: $E_B = \frac{q^2 N_t^2}{8\epsilon\epsilon_0 p}$

Width of space-charge region: $W = \left(\frac{2\epsilon\epsilon_0 E_B}{q^2 p} \right)^{1/2}$

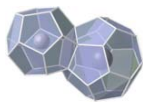


@ 300 K

Specimen		ρ (m Ω -cm)	p (cm $^{-3}$)	E_B (meV)	N_t (cm $^{-2}$)	W (nm)	L (nm)
PbTe I	○	24.9	9.5×10^{17}	60	1.0×10^{13}	54	316
PbTe II	△	12.6	1.5×10^{18}	60	1.3×10^{13}	43	396
PbTe III	●	3.9	5.1×10^{18}	60	2.4×10^{13}	23	376
PbTe IV	▲	2.9	6.2×10^{18}	60	2.6×10^{13}	21	416



C. H. Seager, J. Appl. Phys. **52**, 3960 (1981)
 O. Vigil-Galan, J. Appl. Phys. **90**, 3427 (2001)
 J. Y. W. Seto, J. Appl. Phys. **46**, 5247 (1975)



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Model : Granular Nanocomposites

Goal - describe electrical conductivity and Seebeck coefficient
- understand the role of grain interface scattering

$$\sigma = \frac{2e^2}{3m^*} \int_0^\infty \tau(E) g(E) E \left(-\frac{\partial f(E)}{\partial E} \right) dE$$

$$S = \frac{1}{eT} \left[\frac{\int_0^\infty \tau(E) g(E) E^2 \left(-\frac{\partial f(E)}{\partial E} \right) dE}{\int_0^\infty \tau(E) g(E) E \left(-\frac{\partial f(E)}{\partial E} \right) dE} - \mu \right]$$

e – electron charge

m^* - effective mass

T – temperature

μ – chemical potential

$\tau(E)$ – relaxation rate

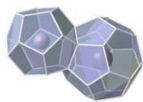
$g(E)$ – total density of states (DOS)

$f(E)$ – energy distribution function

diffusive, quasi-equilibrium transport

σ , S – derived using linear response theory

A. Popescu, L.M. Woods, J. Martin, and G.S. Nolas, Phys. Rev. B **79**, 305302 (2009)



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Scattering Mechanisms

$$\frac{1}{\tau(E)} = \sum_n \frac{1}{\tau_n(E)}$$

Scattering due to acoustic phonons

$$\tau_{a-ph}(E) = \frac{h^4}{8\pi^3} \frac{\rho v_L^2}{k_B T} \frac{1}{(2m^*)^{3/2} D^2} E^{-1/2}$$

Scattering due to optical phonons

$$\tau(E)_{o-ph} = \frac{h^2}{2^{1/2} m^{*1/2} e^2 k_B T (\varepsilon_\infty^{-1} - \varepsilon_0^{-1})} E^{1/2}$$

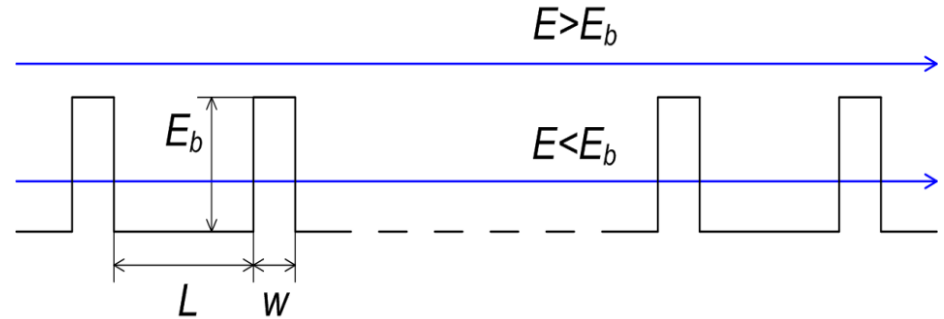
Scattering due to ionized impurities

$$\tau_{imp}(E) = \left[\frac{Z^2 e^4 N_i}{16\pi (2m^*)^{1/2} \varepsilon^2} \ln \left[1 + \left(\frac{2E}{E_m} \right)^2 \right] \right]^{-1} E^{3/2}$$

Interface grain barriers

$$\tau_b(E) = \lambda / v \quad v = \sqrt{2E / m^*}$$

$$\lambda = \sum_{n=1}^{\infty} T^n(E) (1 - T(E)) n L = \frac{T(E) L}{1 - T(E)}$$



Infinite number of barriers:

E_b – barrier height

w – barrier width

L – distance between the barriers

$T(E)$ – quantum mechanical transmission through one barrier

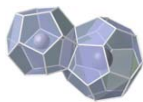
ρ – mass density

v_L – longitudinal speed of sound

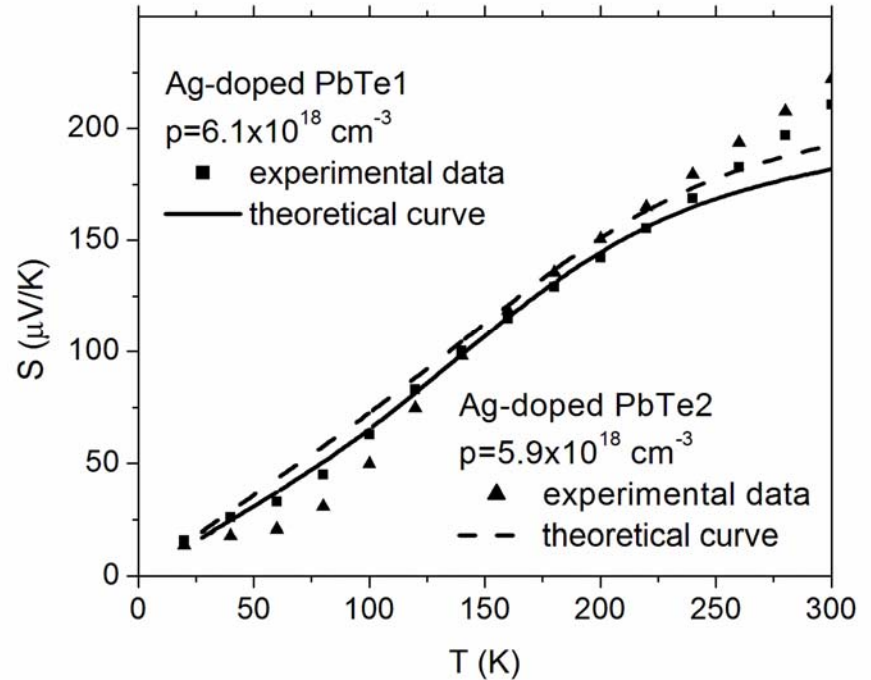
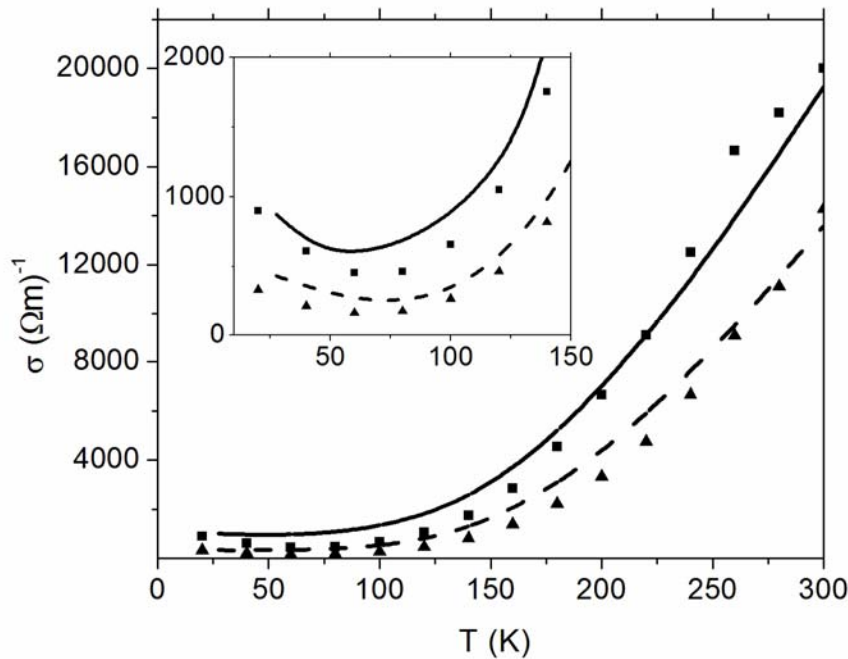
D – deformation potential constant

N_i – concentration of impurities

$\varepsilon_{0,\infty}$ – dielectric constants

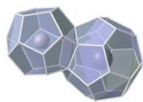


Comparison: Experiment and Theory



comp	ρ (g/cm ³)	v_L (m/s)	D (eV)	m^*/m_0	$\epsilon_\infty^{-1} - \epsilon_0^{-1}$
PbTe	8.16	1730	5 – 15	0.16	0.0022

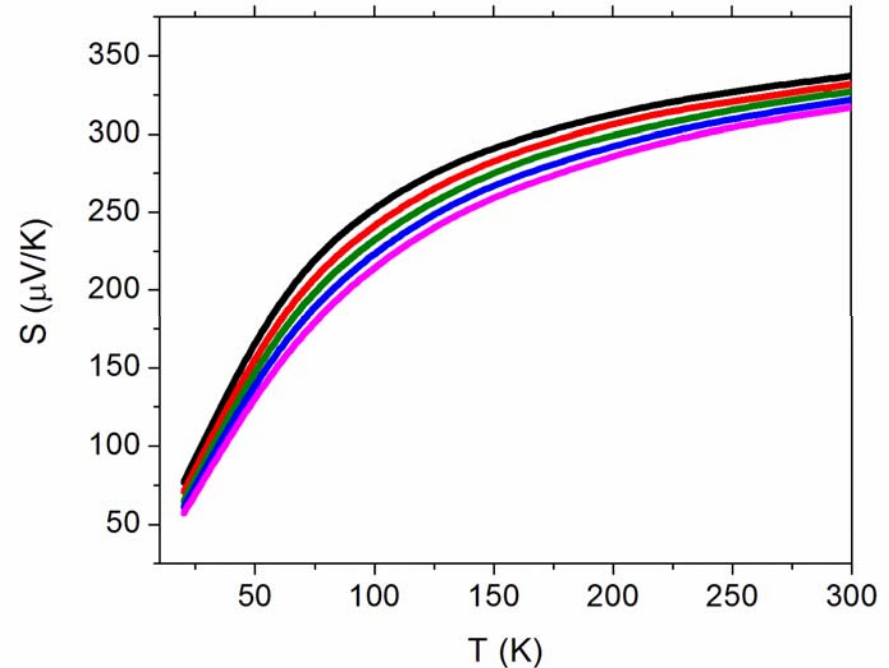
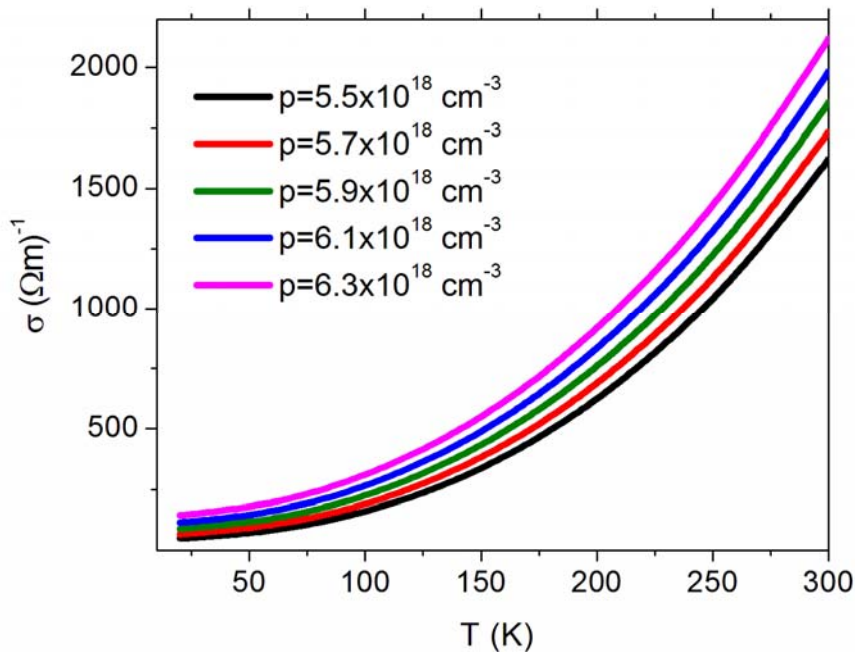
A. Popescu, L.M. Woods, J. Martin, and G.S. Nolas, Phys. Rev. B **79**, 305302 (2009)



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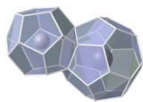
TE Transport Properties: The role of Carrier Concentration



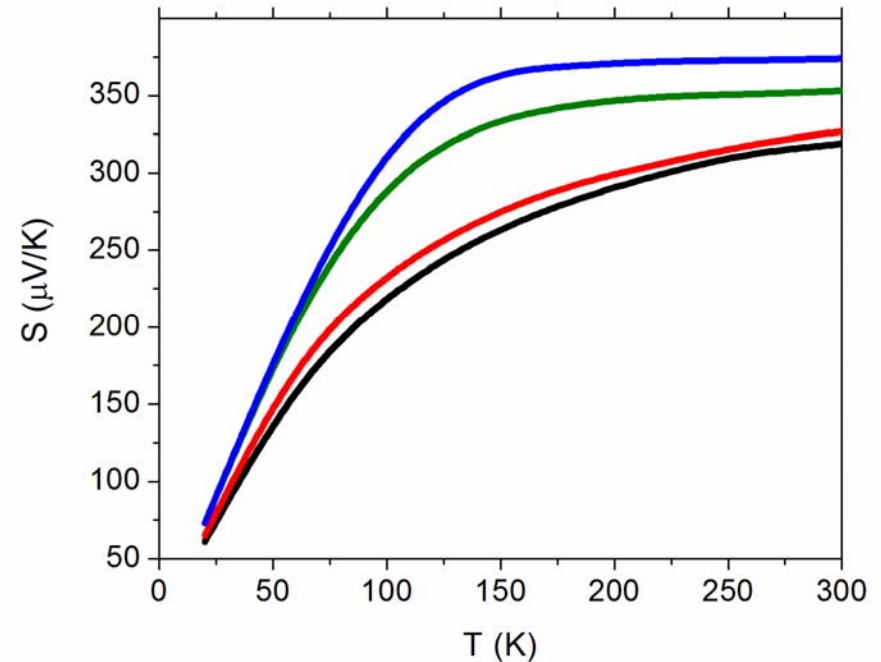
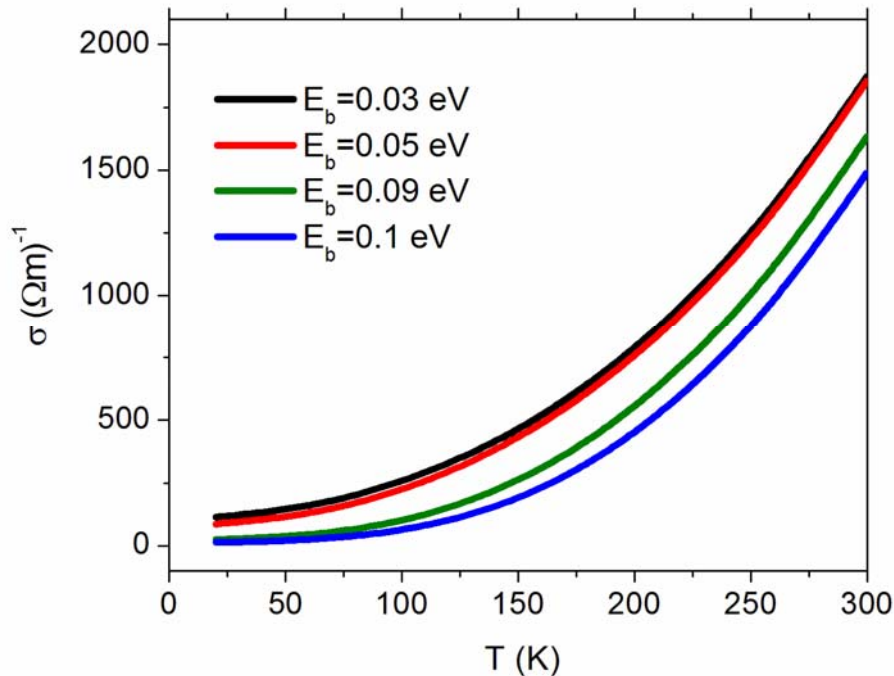
*Higher concentration results
in more transport carriers*



*σ increases &
 S decreases.*

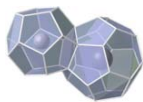


Grain Boundary Height

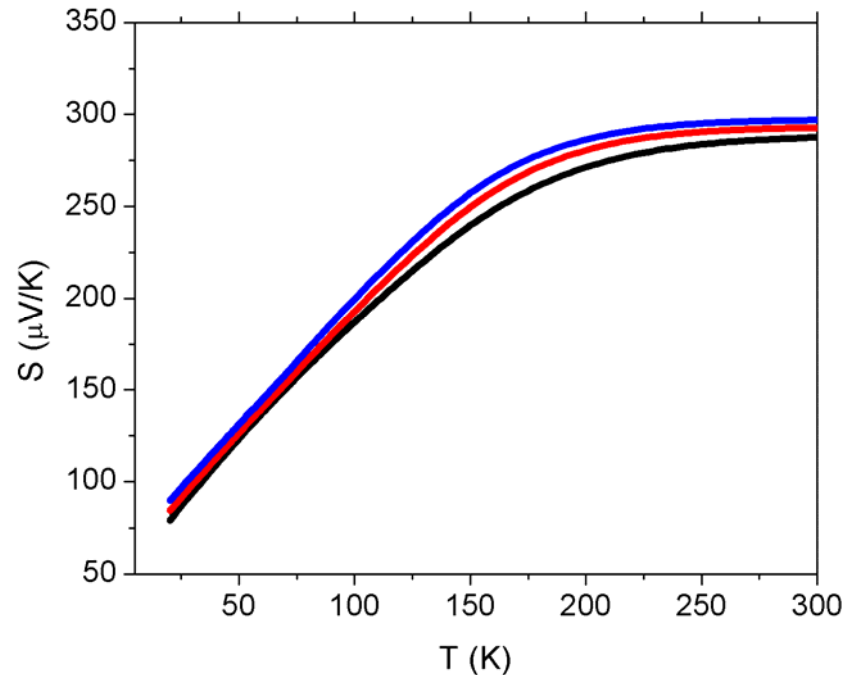
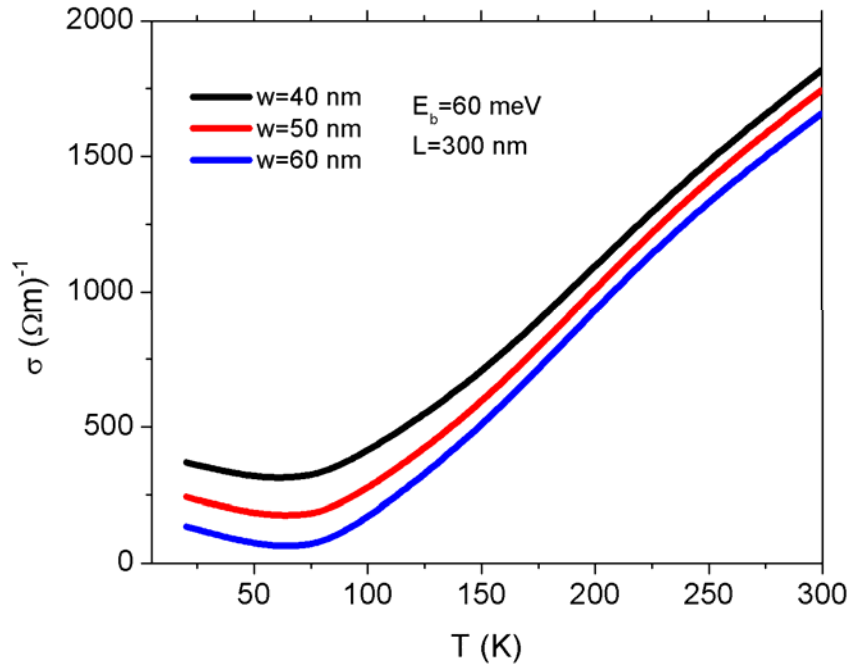


Higher E_b – more carriers are scattered \longrightarrow *σ decreases & S increases*

Lower E_b – carrier energy larger than barrier \longrightarrow *σ increases & S decreases*

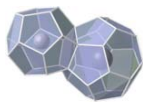


Grain Boundary Width

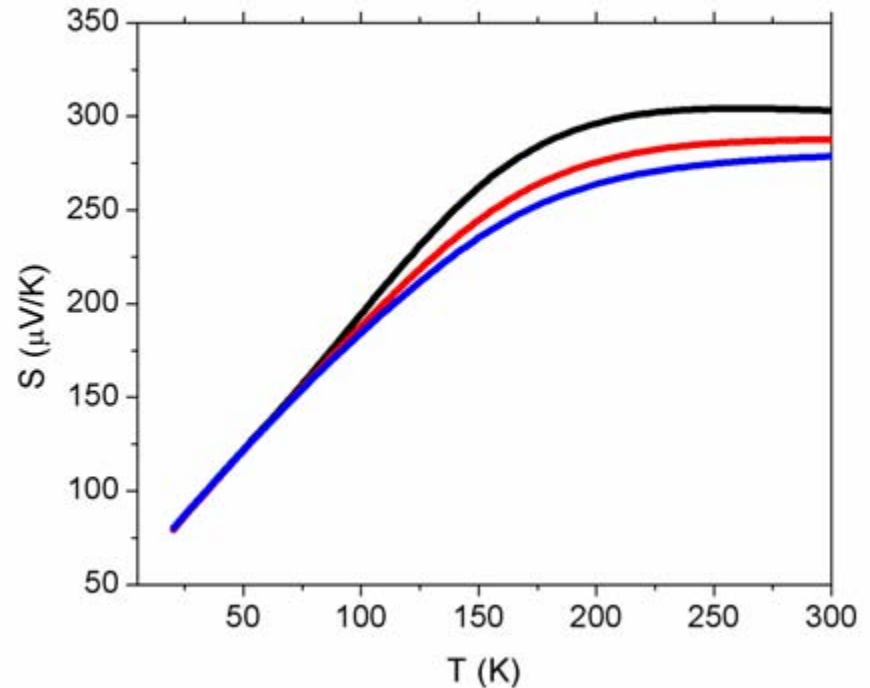
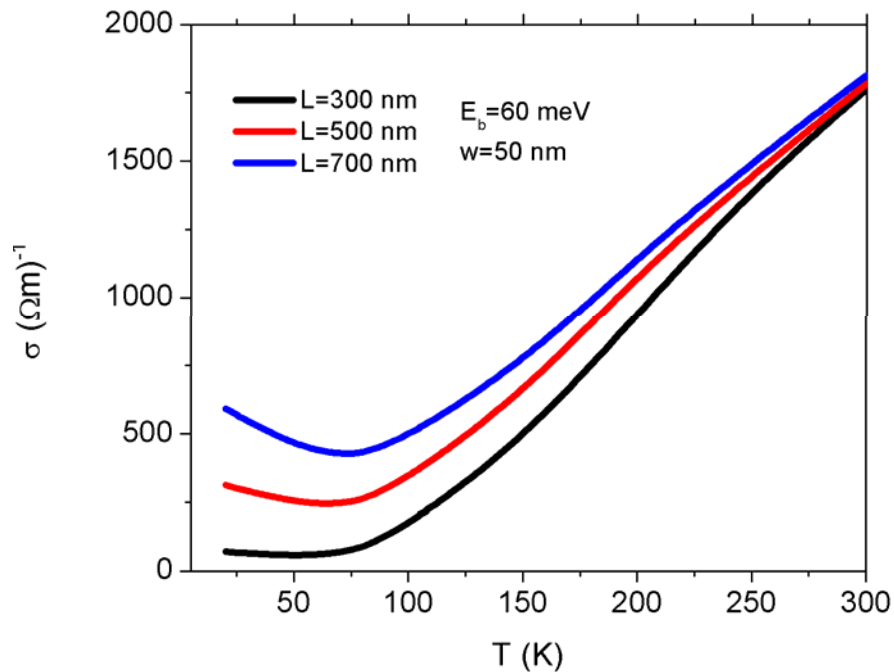


Larger w – smaller transmission probability $T(E)$ \longrightarrow σ decreases & S increases

Smaller w – larger transmission probability $T(E)$ \longrightarrow σ increases & S decreases

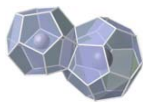


Distance (L) Between Barriers

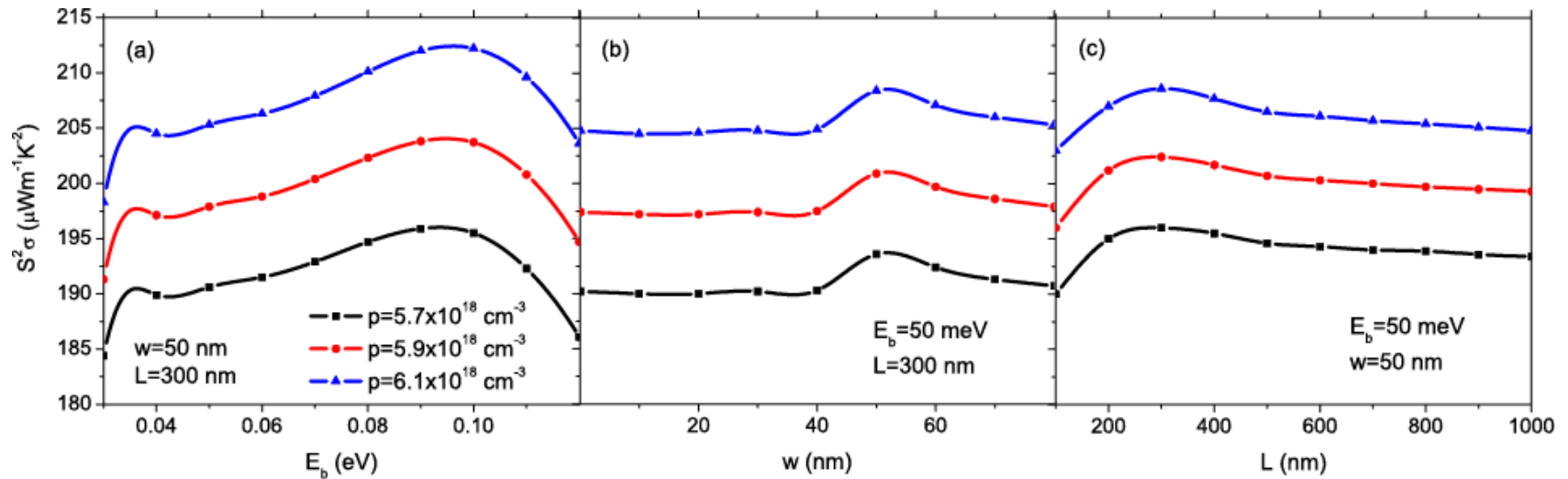


Smaller L – more frequent carrier/interface scattering $\longrightarrow \sigma$ decreases & S increases

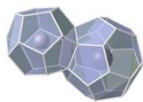
Larger L – less frequent carrier/interface scattering $\longrightarrow \sigma$ increases & S decreases



$S^2\sigma$ for PbTe granular nanocomposites



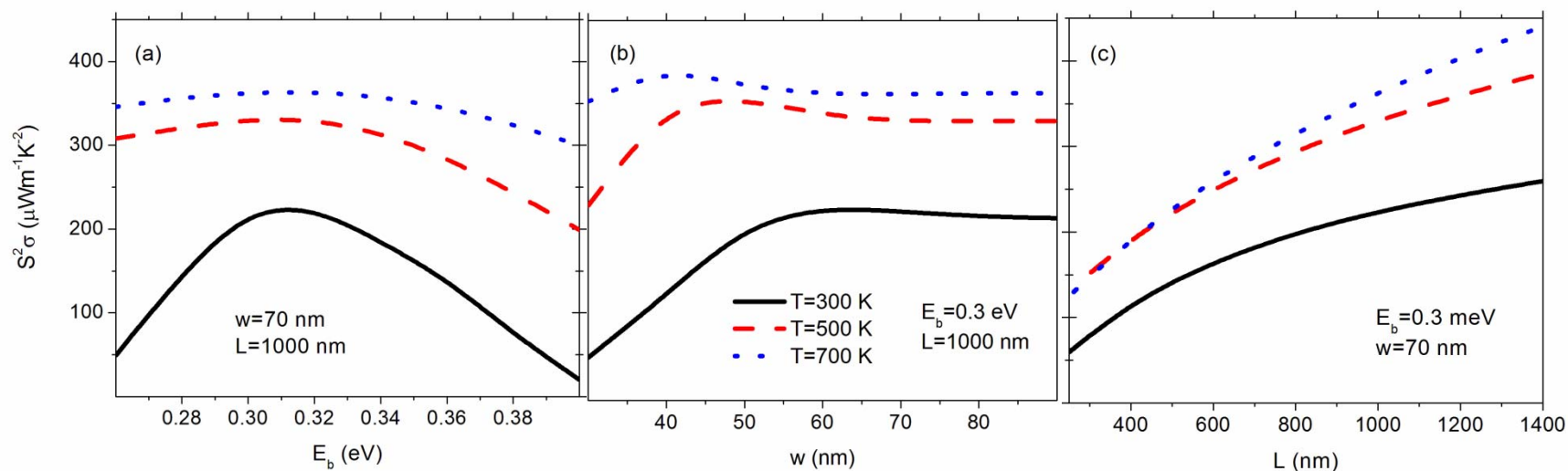
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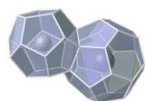
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CoSb₃ for PbTe granular nanocomposites



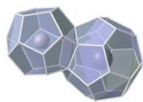
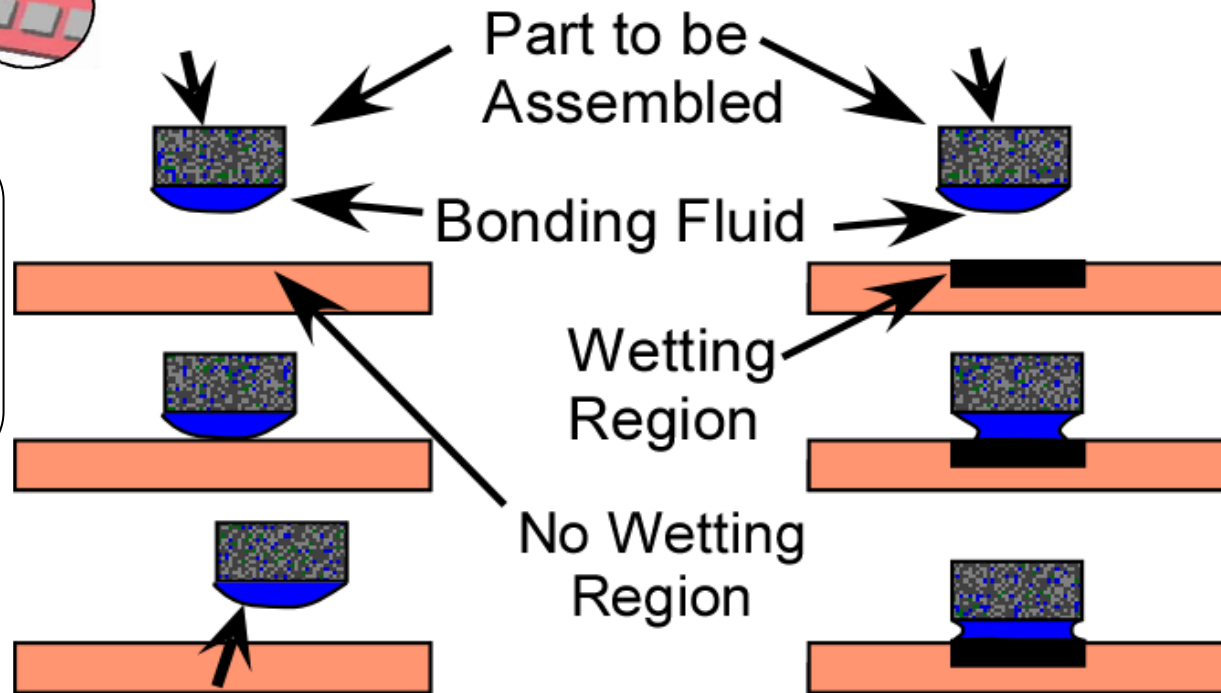
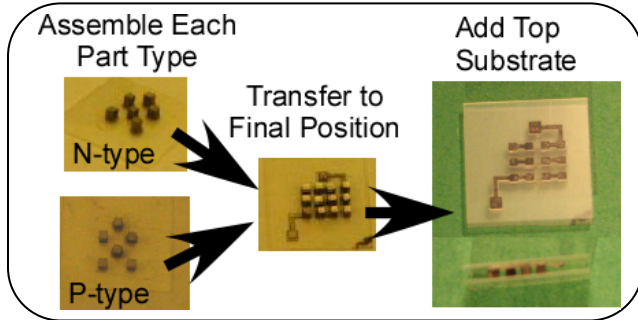
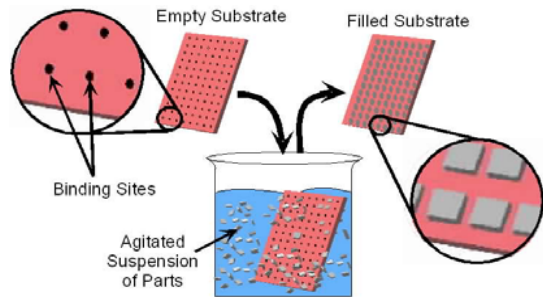
L.M. Woods, A. Popescu, J. Martin, and G.S. Nolas, Proc. Mater. Rec. Soc. **1166** (2009)



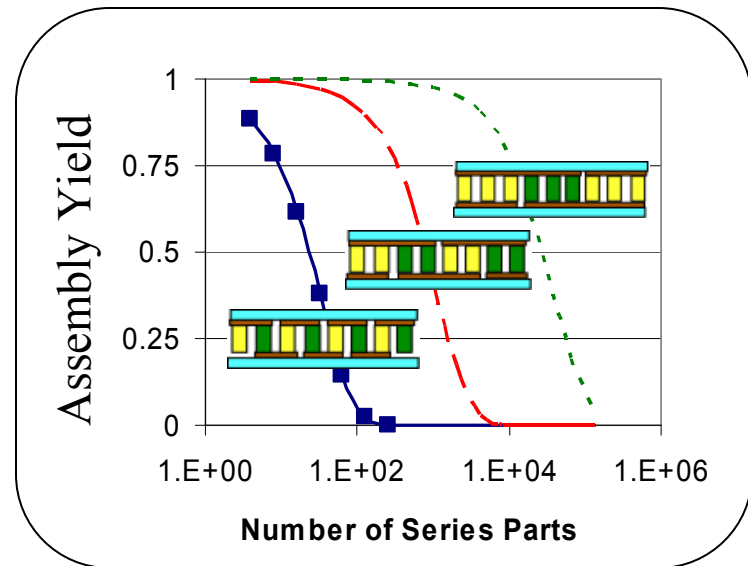
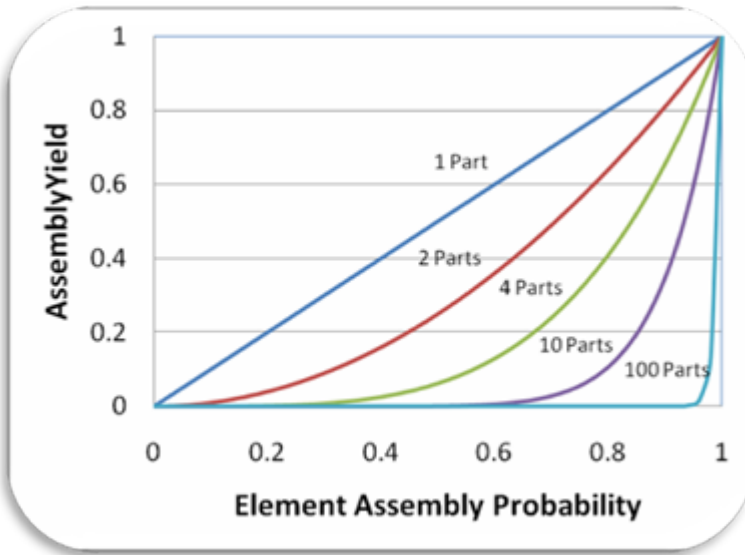
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Device manufacturing: Self Assembly



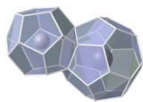
Challenge of Self Assembly: Rate & Yield



Assembly yield depend on process and design parameters.

Goal: Develop predictive models of self assembly process rate and yield to facilitate design of high rate and yield processes at industrial scales.

N. Crane, P. Mishra, J. Murray, G.S. Nolas, J. Electronic Mater **38**, 1252 (2009)



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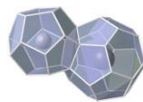
Acknowledgements

Lidong Chen

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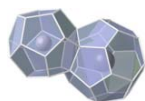
Jim Salvador & Jihui Yang

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