Modelling and Numerical Methods

Lecture 3
Kinematics of Continua
& Conservation Equations

Outline Lecture 3

Part 1: Kinematics

- Material vs. spatial descriptions
- Time derivatives

chapter3.ipynb

- Displacement
- Infinitesimal Strain Tensor

Part 2: Conservation Equations

- Conservation of Mass
- Conservation of Momentum

chapter4.ipynb

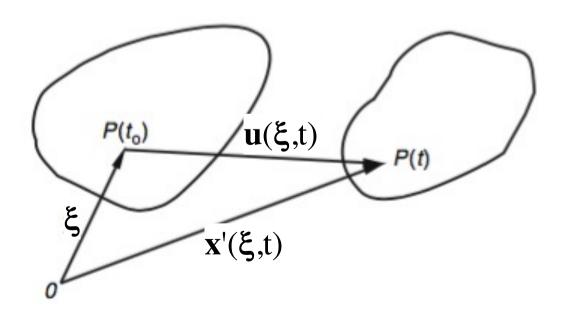
Learning Objectives Kinematics

- Be able to use material and spatial descriptions of variables and their time derivatives.
- Be able to compute infinitesimal strain (strain rate) tensor given a displacement (velocity) field.
- Know meaning of the different components of the infinitesimal strain (rate) tensor
- Be able to find principal strain(rate)s and strain (rate) invariants and know what they represent
- Understand difference between infinitesimal and finite strain

Displacement

Can result in

- (a) Rigid body motion
- (b) Deformation of the body



Rigid body motion

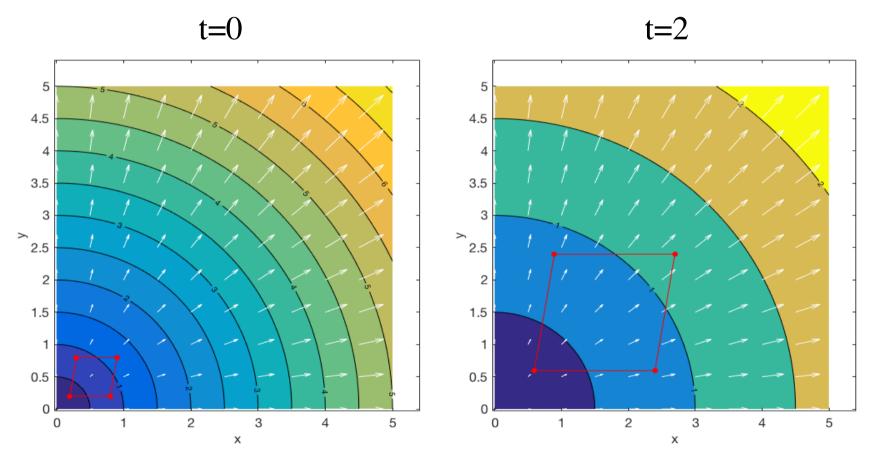
- Translation: $\mathbf{x}' = \boldsymbol{\xi} + \mathbf{c}(t)$, with $\mathbf{c}(0) = \mathbf{0}$ $\Rightarrow \mathbf{u} = \mathbf{x}' - \boldsymbol{\xi}$, each point same $\mathbf{u}(t) = \mathbf{c}(t)$
- Rotation: \mathbf{x}' - \mathbf{b} = $\mathbf{R}(t)(\boldsymbol{\xi}$ - $\mathbf{b})$, where $\mathbf{R}(t)$ is rotation tensor, with $\mathbf{R}(0)$ = \mathbf{I} , \mathbf{b} is the point of rotation. $\mathbf{R}(t)$ is an orthogonal transformation (preserves lengths and angles, $\mathbf{R}^T\mathbf{R}$ = \mathbf{I} , $\det(\mathbf{R})$ = $\mathbf{1}$)

If \mathbf{u} depends on \mathbf{x} and t, then internal deformation

Example:

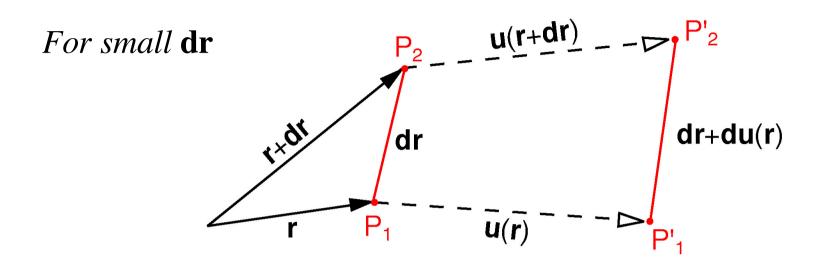
$$v_i = \frac{kx_i}{1 + kt}$$

Displacement



translation & deformation

Deformation tensor



$$P_1$$
 at $r \rightarrow P'_1$ at $r+u(r)$, P_2 at $r+dr \rightarrow P'_2$ at $r+dr+u(r+dr)$.

$$dr' = P'_2 - P'_1 = dr + [u(r+dr) - u(r)] = dr + du(r) = dr + dr \cdot \nabla u(r)$$

deformation of P_2 - P_1 described by: $du_i = dx_j \frac{\partial u_i}{\partial x_j}$

$$\mathbf{d}\mathbf{u} = \mathbf{d}\mathbf{r} \cdot \nabla \mathbf{u} = \nabla^{\mathrm{T}}\mathbf{u} \cdot \mathbf{d}\mathbf{r}$$

$$\mathrm{d} u_i = \mathrm{d} x_j \ \frac{\partial u_i}{\partial x_j} \ : \qquad \begin{pmatrix} \mathrm{d} u_1 \\ \mathrm{d} u_2 \\ \mathrm{d} u_3 \end{pmatrix} = \left(\mathrm{d} x_1 \ \mathrm{d} x_2 \ \mathrm{d} x_3\right) \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_2}{\partial x_1} & \frac{\partial u_3}{\partial x_1} \\ \frac{\partial u_1}{\partial x_2} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_3}{\partial x_2} \\ \frac{\partial u_1}{\partial x_3} & \frac{\partial u_2}{\partial x_3} & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

$$\frac{\partial u_{i}}{\partial x_{j}} = \frac{1}{2} \left(\frac{\partial u_{i}}{\partial x_{j}} + \frac{\partial u_{j}}{\partial x_{i}} \right) + \frac{1}{2} \left(\frac{\partial u_{i}}{\partial x_{j}} - \frac{\partial u_{j}}{\partial x_{i}} \right)$$

$$\nabla \mathbf{u}^{T} = \frac{1}{2} (\nabla \mathbf{u} + \nabla \mathbf{u}^{T}) + \frac{1}{2} (\nabla \mathbf{u} - \nabla \mathbf{u}^{T})$$

Total deformation is:

- rigid body translation $\mathbf{u}(\mathbf{r})$
- rigid body rotation $\omega \cdot dr$
- internal deformation, strain **\varepsilon dr** result of stresses

Infinitesimal strain and rotation tensors

$$\boldsymbol{\varepsilon} = \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_1} + \frac{\partial u_1}{\partial x_2} \right) & \frac{1}{2} \left(\frac{\partial u_3}{\partial x_1} + \frac{\partial u_1}{\partial x_3} \right) \\ \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{\partial u_2}{\partial x_2} & \frac{1}{2} \left(\frac{\partial u_3}{\partial x_2} + \frac{\partial u_2}{\partial x_3} \right) \\ \frac{1}{2} \left(\frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

$$\omega = \begin{bmatrix} 0 & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} \right) & \frac{1}{2} \left(\frac{\partial u_3}{\partial x_1} - \frac{\partial u_1}{\partial x_3} \right) \\ \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} - \frac{\partial u_2}{\partial x_1} \right) & 0 & \frac{1}{2} \left(\frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3} \right) \\ \frac{1}{2} \left(\frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1} \right) & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_3} - \frac{\partial u_3}{\partial x_2} \right) & 0 \end{bmatrix}$$

diagonal infinitesimal strain tensor elements

For a line segment $\mathbf{dr} = (dx_1,0,0)$ deforming in displacement field $\mathbf{u} = (u_1,0,0)$:

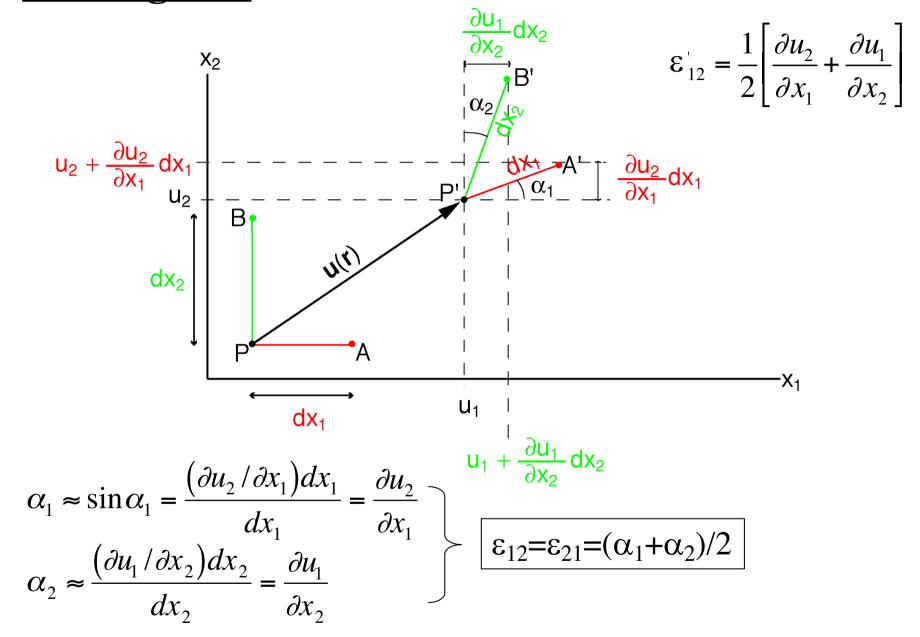
$$x_1 + u_1(x_1)$$
 $x_1 + dx_1 + u_1(x_1) + \frac{\partial u_1}{\partial x_1} dx_1$
 x_1 $x_1 + dx_1$

the new length $dx'_1 \approx dx_1 + (\partial u_1/\partial x_1)dx_1 = dx_1 + \varepsilon_{11}dx_1$

 $\Rightarrow \varepsilon_{11} = [dx'_1 - dx_1]/dx_1 =$ the relative change in length of a line element, originally in x_1 direction.

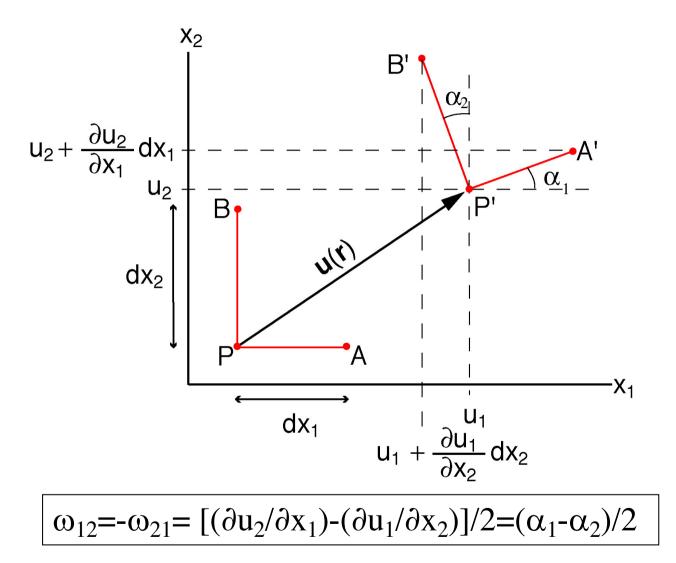
The relative change in volume (V'-V)/V of a cube $V=dx_1dx_2dx_3 \approx \varepsilon_{11}+\varepsilon_{22}+\varepsilon_{33}=\varepsilon_{ii}=tr(\varepsilon)=\nabla\cdot\mathbf{u}$.

off-diagonal infinitesimal strain tensor elements



 $2\epsilon_{12}$ is the change in angle of an originally 90° angle between dx_1 and dx_2

infinitesimal rotation tensor elements



 ω_{12} is common rigid rotation angle of vectors in the dx₁ - dx₂ plane (around x₃)

Extra: Rotation tensor and rotation vector

For any antisymmetric tensor **W**, a corresponding *dual* or *axial vector* **w** can be found so that

$$\mathbf{W} \cdot \mathbf{a} = \mathbf{w} \times \mathbf{a}$$

Vector w relates to the components of W as:

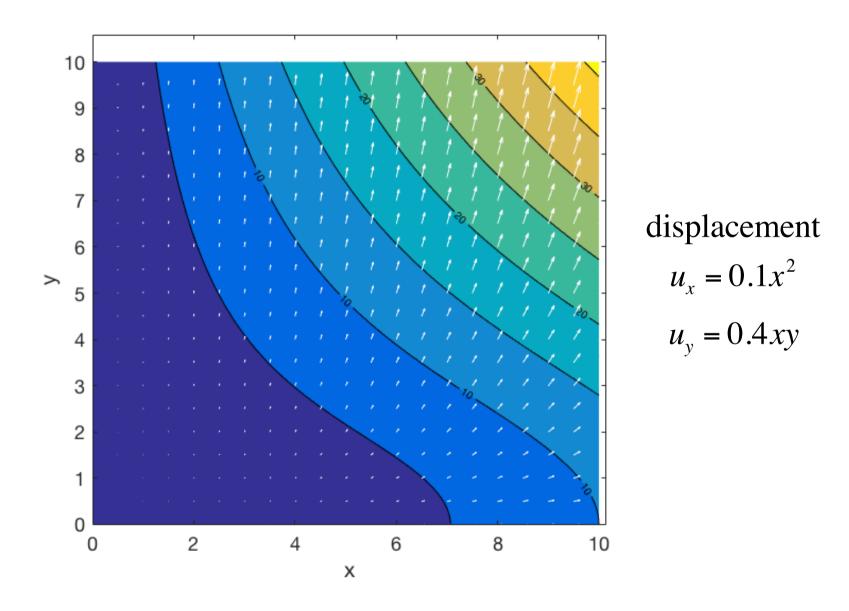
$$\mathbf{w} = -\mathbf{W}_{23}\hat{\mathbf{e}}_1 + \mathbf{W}_{13}\hat{\mathbf{e}}_2 - \mathbf{W}_{12}\hat{\mathbf{e}}_3$$

For the rotation tensor, an equivalent rotation vector exists:

$$\boldsymbol{\omega} \cdot \mathbf{dx} = \mathbf{r}_{\omega} \times \mathbf{dx}$$
 where: $\mathbf{r}_{\omega} = \frac{1}{2} \nabla \times \mathbf{u}$

Note that ω only describes the overall rigid body rotation, not the total rotation of each individual segment dx, which is also influenced by ε

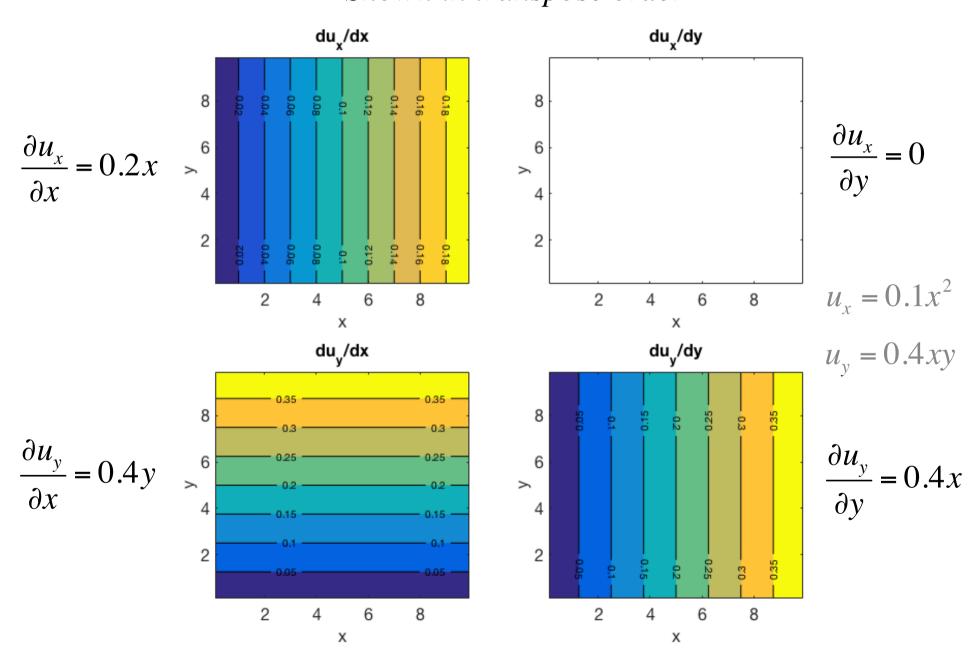
Example displacement – infinitesimal strain



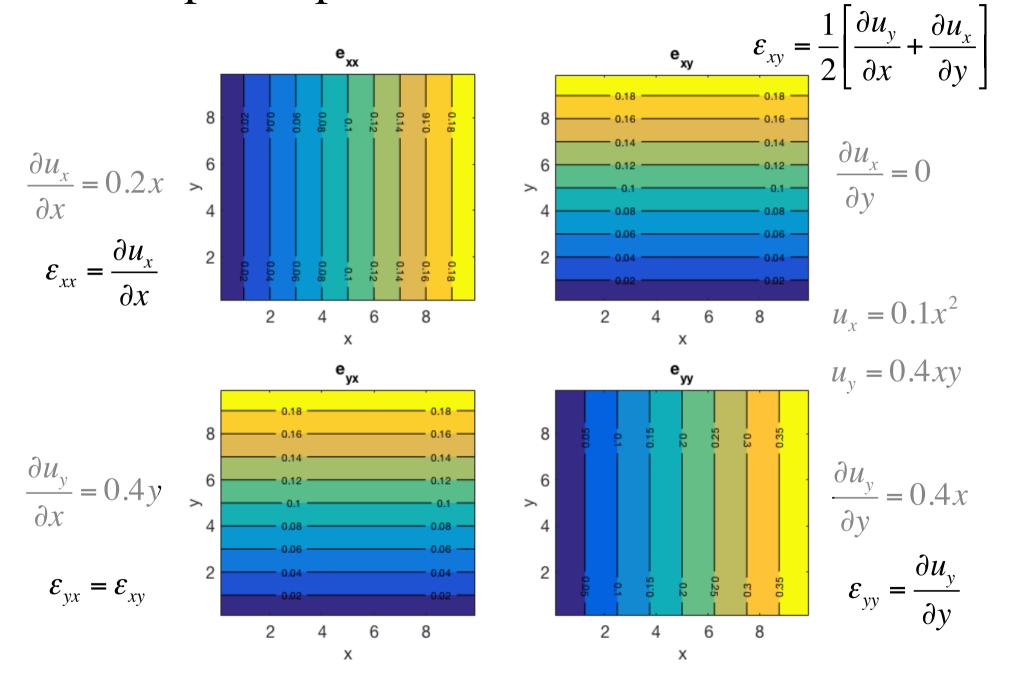
Please take a break

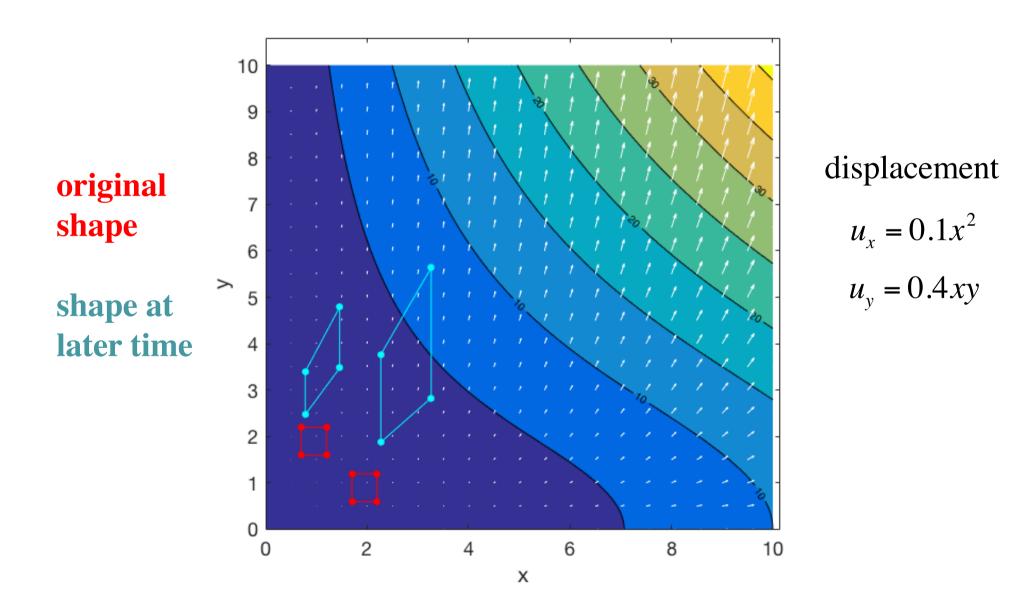
And then look at the "Class Exercise" on Displacement and Strain fields and Exercise 8

Example displacement – infinitesimal strain Shown in transpose order

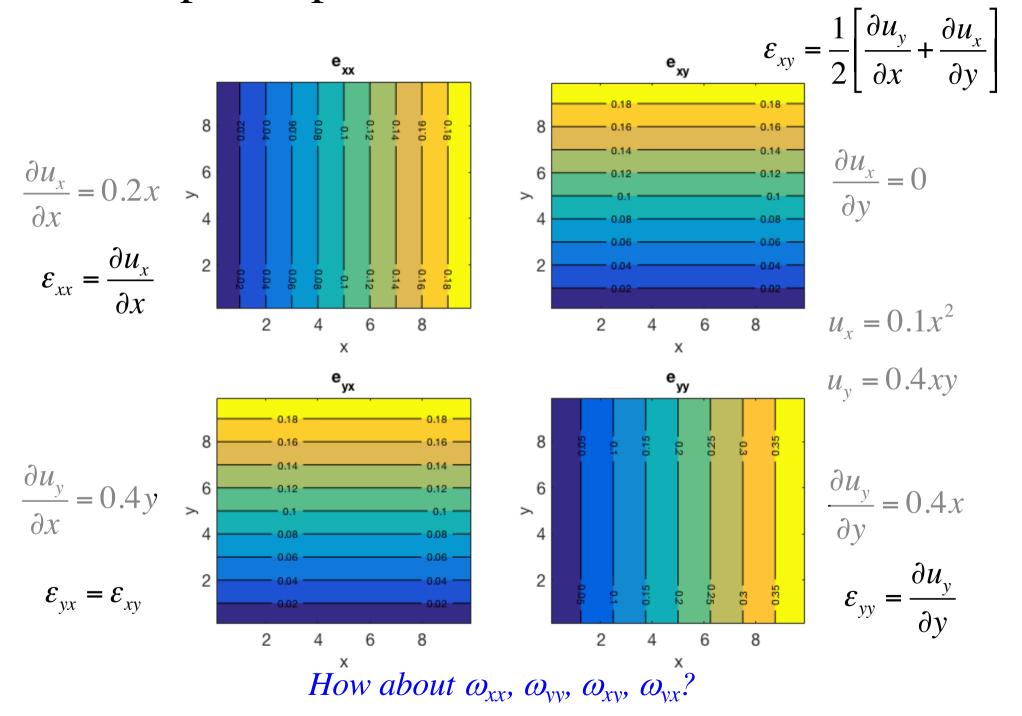


Example displacement – infinitesimal strain

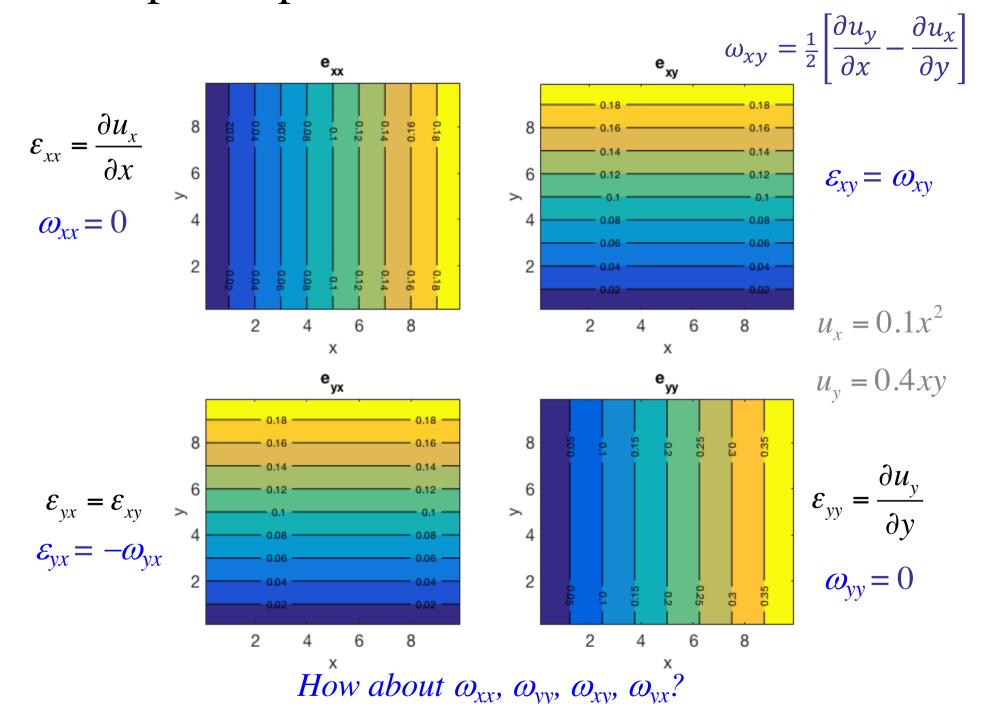




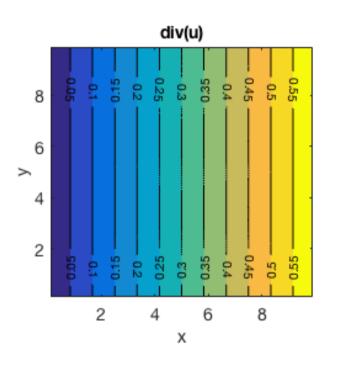
Example displacement – infinitesimal strain

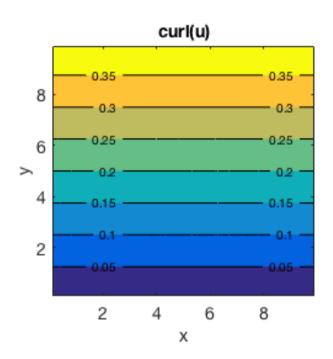


Example displacement – infinitesimal rotation



Example displacement – infinitesimal strain





$$\nabla \cdot \mathbf{u} = \frac{\partial u_x}{\partial x} + \frac{\partial u_y}{\partial y}$$

$$\nabla \times \mathbf{u} = \left(\frac{\partial u_y}{\partial x} - \frac{\partial u_x}{\partial y}\right) \hat{\mathbf{e}}_z$$

infinitesimal strain tensor properties

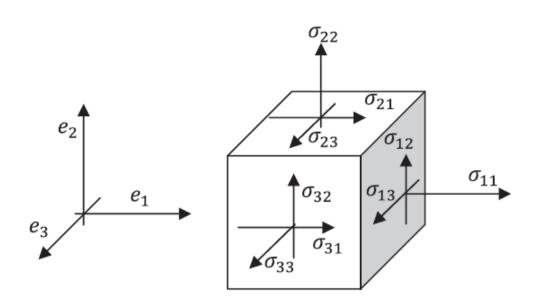
transform to arbitrary coordinate frame:

$$\begin{array}{ll} \epsilon_{nn} & = \epsilon_{11} cos^2 \phi + \epsilon_{21} sin \phi cos \phi + \epsilon_{12} sin \phi cos \phi + \epsilon_{22} sin^2 \phi \\ \epsilon_{ns} & = \epsilon_{11} sin \phi cos \phi + \epsilon_{21} sin^2 \phi - \epsilon_{12} cos^2 \phi - \epsilon_{22} sin \phi cos \phi \end{array}$$

 $\epsilon_1, \epsilon_2, \epsilon_3$ - principal strains : minimum, maximum and intermediate fractional length changes

isotropic, deviatoric strain: $\varepsilon_{ij} = (\theta/3)\delta_{ij} + \varepsilon'_{ij}$

- $tr(\varepsilon) = \theta = sum of normal strains = volume change$
- ϵ'_{ij} is deviatoric strain, change in shape; no change in volume
- $tr(\varepsilon') = 0$, does not imply $\varepsilon'_{ij} = 0$ for i = j
- $\varepsilon_{ij} = 0$ for $i \neq j$ does not ensure no shape change



Stress components

Reminder

traction on a plane

$$\mathbf{t} = \begin{bmatrix} \sigma_{11} & \sigma_{21} & \sigma_{31} \\ \sigma_{12} & \sigma_{22} & \sigma_{32} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{bmatrix} \cdot \hat{\mathbf{n}}$$

what is $\hat{\mathbf{e}}_1 \cdot \mathbf{t} = \hat{\mathbf{e}}_1 \cdot \boldsymbol{\sigma}^T \cdot \hat{\mathbf{n}}$? t_1 on plane with normal $\hat{\bf n}$

what is $\hat{\mathbf{e}}_1 \cdot \boldsymbol{\sigma}^T \cdot \hat{\mathbf{e}}_1$? σ_{11} what is $\hat{\mathbf{e}}_1 \cdot \boldsymbol{\sigma}^T \cdot \hat{\mathbf{e}}_2$?

 σ_{21}

Strain components

$$\mathbf{\varepsilon} = \begin{bmatrix} \varepsilon_{11} & \varepsilon_{12} & \varepsilon_{13} \\ \varepsilon_{21} & \varepsilon_{22} & \varepsilon_{23} \\ \varepsilon_{31} & \varepsilon_{32} & \varepsilon_{33} \end{bmatrix} \qquad \hat{\mathbf{e}}_{1} \cdot \mathbf{\varepsilon} \cdot \hat{\mathbf{e}}_{1} = \varepsilon_{11} \\ \hat{\mathbf{e}}_{1} \cdot \mathbf{\varepsilon} \cdot \hat{\mathbf{e}}_{2} = \varepsilon_{12}$$

$$\hat{\mathbf{e}}_1 \cdot \mathbf{\epsilon} \cdot \hat{\mathbf{e}}_1 = \varepsilon_{11}$$

$$\hat{\mathbf{e}}_1 \cdot \mathbf{\epsilon} \cdot \hat{\mathbf{e}}_2 = \varepsilon_{12}$$

 $\mathbf{\epsilon} \cdot \hat{\mathbf{p}} = \mathbf{dp}'$ the change in unit vector $\hat{\mathbf{p}}$ after deformation by $\mathbf{\epsilon}$

 $\hat{\mathbf{p}} \cdot \mathbf{\epsilon} \cdot \hat{\mathbf{p}} = \text{elongation by } \mathbf{\epsilon} \text{ of unit vector } \hat{\mathbf{p}} \text{ in direction } \hat{\mathbf{p}}$

$$= \widehat{\mathbf{p}} \cdot \mathbf{dp}' = |\mathbf{dp}'| \cos \alpha$$

Where α is angle between dp' and \hat{p}

Strain Rate Tensor

In similar way as strain tensor, a tensor that describes the rate of change of deformation can be defined from **velocity gradient**:

$$\frac{\mathbf{D}}{\mathbf{D}t}\mathbf{dr} = \nabla \mathbf{v}$$

$$\nabla \mathbf{v} = \begin{bmatrix} \frac{\partial v_1}{\partial x_1} & \frac{\partial v_2}{\partial x_1} & \frac{\partial v_3}{\partial x_1} \\ \frac{\partial v_1}{\partial x_2} & \frac{\partial v_2}{\partial x_2} & \frac{\partial v_3}{\partial x_2} \\ \frac{\partial v_1}{\partial x_3} & \frac{\partial v_2}{\partial x_3} & \frac{\partial v_3}{\partial x_3} \end{bmatrix}$$

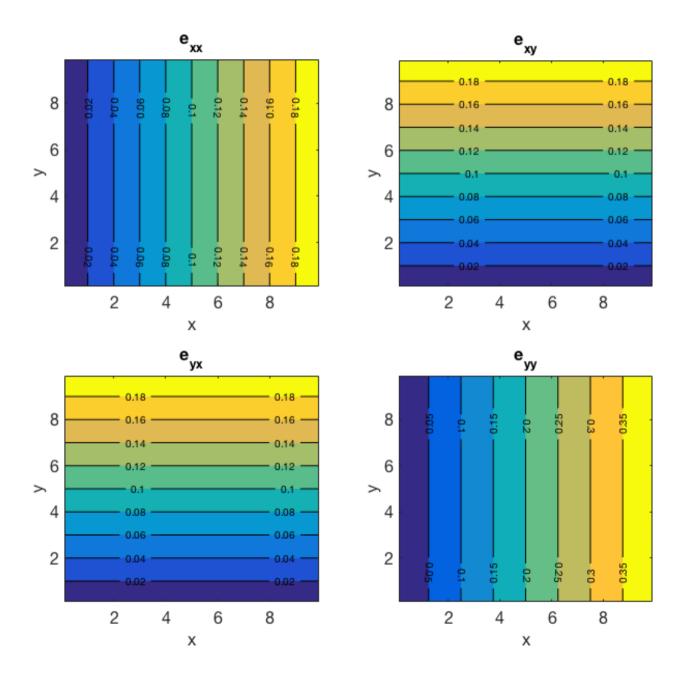
$$\nabla \mathbf{v} = \frac{1}{2} (\nabla \mathbf{v} + \nabla \mathbf{v}^T) + \frac{1}{2} (\nabla \mathbf{v} - \nabla \mathbf{v}^T)$$

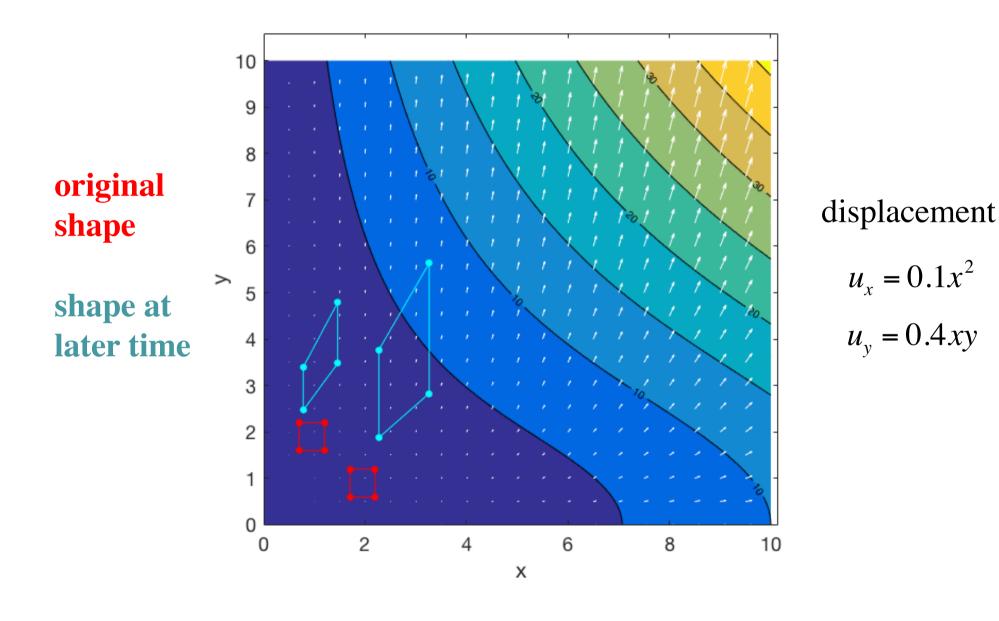
$$\nabla \mathbf{v} = \mathbf{D} + \mathbf{W}$$

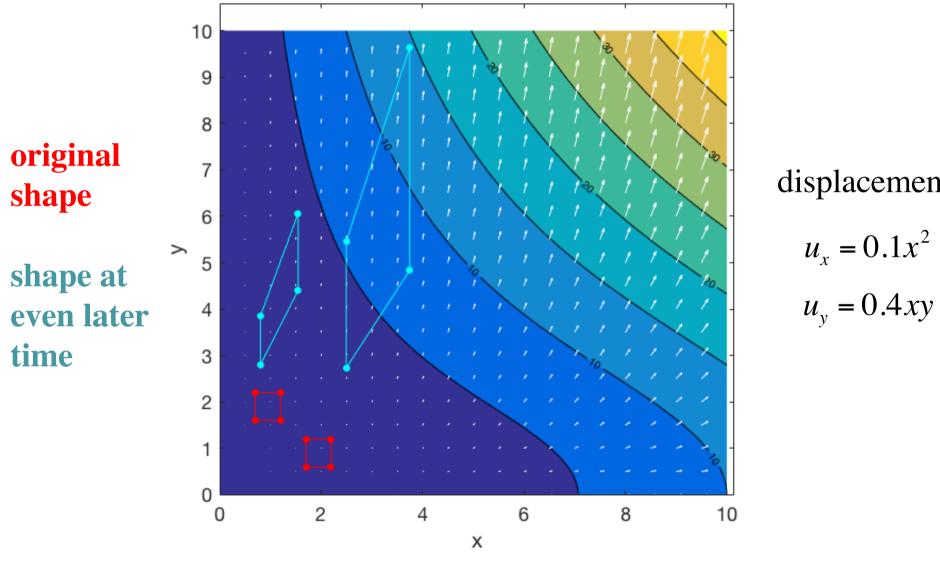
Velocity gradient tensor is the sum of **strain rate** and **vorticity** tensors

Infinitesimal strain

Over small increments, can assume constant displacement gradient encountered



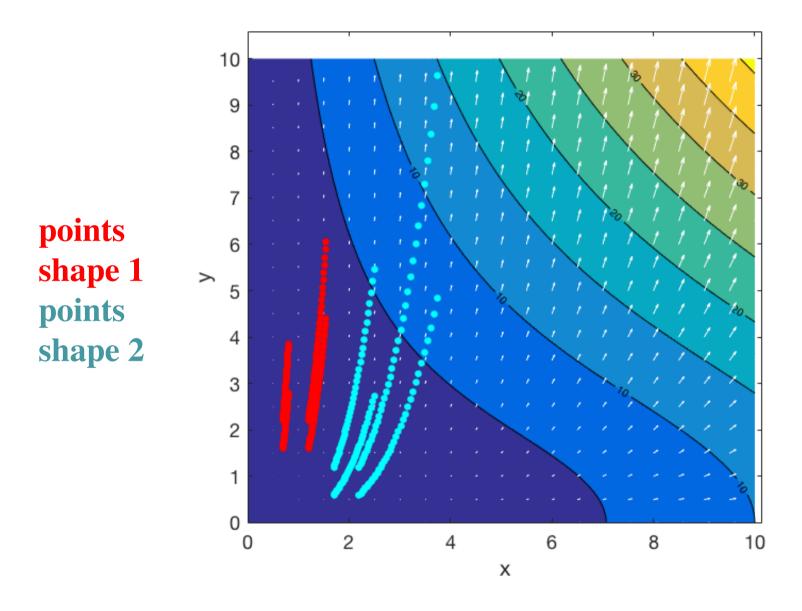




displacement

$$u_{x} = 0.1x^{2}$$

$$u_{y} = 0.4xy$$



displacement

$$u_{r} = 0.1x^{2}$$

$$u_x = 0.1x^2$$
$$u_y = 0.4xy$$

Learning Objectives

- Be able to use material and spatial descriptions of variables and their time derivatives.
- Be able to compute infinitesimal strain (strain rate) tensor given a displacement (velocity) field.
- Know meaning of the different components of the infinitesimal strain (rate) tensor
- Be able to find principal strain(rate)s and strain (rate) invariants and know what they represent
- Understand difference between infinitesimal and finite strain

Outline Lecture 3

- Material vs. spatial descriptions
- Time derivatives
- Displacement
- Infinitesimal Deformation

Further reading on the topics in the lecture can be done in for example: Lai, Rubin, Kremple (2010): Ch. 3-1 through 3-15 and we covered some of the basics discussed in 3-20 to 3-26

Try yourself

Complete any of the remaining exercises in *chapter3.ipynb*:

Exercise 1, 2, 5, 7, 9a

- Additional practise: Exercise 3, 6, 8
- Advanced practise: Exercise 4, 9b, 10