

## Mathematics 12: Quantum Spins

Quantum mechanics for a particle in space is controlled by the commutation relations:

$$[\hat{x}_i, \hat{p}_j] \equiv \hat{x}_i \hat{p}_j - \hat{p}_j \hat{x}_i = i\hbar \delta_{ij}$$

for the different spatial directions labelled by  $i$ . If we consider *orbital* motion using the orbital angular momentum:

$$\begin{aligned}\hat{\mathbf{L}} &= \hat{\mathbf{x}} \times \hat{\mathbf{p}} \\ \hat{L}_i &= \epsilon_{ijk} \hat{x}_j \hat{p}_k\end{aligned}$$

then the consequential commutation relations are:

$$[\hat{L}_i, \hat{L}_j] \equiv \hat{L}_i \hat{L}_j - \hat{L}_j \hat{L}_i = i\hbar \epsilon_{ijk} \hat{L}_k$$

Although orbital motion involves integer values for angular momentum, the electronic *spin* involves half-integral values. We will deal with spin-half here, the smallest permissible value.

The natural representation is  $2 \times 2$ , with:

$$\hat{S}^x = \frac{\hbar}{2} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad \hat{S}^y = \frac{\hbar}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \quad \hat{S}^z = \frac{\hbar}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

acting on a basis,  $\{|\uparrow\rangle, |\downarrow\rangle\}$ . You can easily verify the previous commutation relations, and further:

$$\hat{S}^i \hat{S}^j = \frac{\hbar^2}{4} \delta_{ij} \mathbf{1} + \frac{i\hbar}{2} \epsilon_{ijk} \hat{S}_k$$

The ‘best’ representation is in terms of:

$$\hat{S}^z = \frac{\hbar}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad \hat{S}^+ = \hat{S}^x + i\hat{S}^y = \hbar \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \quad \hat{S}^- = \hat{S}^x - i\hat{S}^y = \hbar \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}$$

the so-called *raising* and *lowering* operators. In terms of our basis,  $\hat{S}^+$  ‘raises’ the spin (ie turns a down into an up) and  $\hat{S}^-$  ‘lowers’ the spin (ie turns an up into a down).

We are interested in interactions between spins:

$$\hat{H}_1 = \sum_{ij} J_{ij} \hat{S}_j^z \hat{S}_i^z = \sum_{ij} J_{ij} \hat{\mathbf{S}}_i^{\parallel} \cdot \hat{\mathbf{S}}_j^{\parallel}$$

the **Ising model**,

$$\hat{H}_2 = \sum_{ij} J_{ij} \{\hat{S}_i^x \hat{S}_j^x + \hat{S}_i^y \hat{S}_j^y\} = \frac{1}{2} \sum_{ij} J_{ij} \{\hat{S}_i^+ \hat{S}_j^- + \hat{S}_i^- \hat{S}_j^+\} = \sum_{ij} J_{ij} \hat{\mathbf{S}}_i^{\perp} \cdot \hat{\mathbf{S}}_j^{\perp}$$

the **x-y model**, and

$$\hat{H}_3 = \sum_{ij} J_{ij} \{\hat{S}_i^x \hat{S}_j^x + \hat{S}_i^y \hat{S}_j^y + \hat{S}_i^z \hat{S}_j^z\} = \frac{1}{2} \sum_{ij} J_{ij} \{\hat{S}_i^+ \hat{S}_j^- + \hat{S}_i^- \hat{S}_j^+\} + \sum_{ij} J_{ij} \hat{S}_i^z \hat{S}_j^z = \sum_{ij} J_{ij} \hat{\mathbf{S}}_i \cdot \hat{\mathbf{S}}_j$$

the **Heisenberg model**, where  $J_{ij}$  are the bond strengths between the  $i$ 'th and  $j$ 'th spin. If  $J_{ij} < 0$  the spins want to be parallel at low temperature, and if  $J_{ij} > 0$  the spins want to be anti-parallel at low temperature.

To understand the action of these interactions we need to consider the basis for a pair of spins:

$$\{|\uparrow\uparrow\rangle, |\uparrow\downarrow\rangle, |\downarrow\uparrow\rangle, |\downarrow\downarrow\rangle\}$$

in terms of which:

$$\hat{S}^z \hat{S}'^z |\uparrow\uparrow\rangle = \frac{1}{4} |\uparrow\uparrow\rangle \quad \hat{S}^z \hat{S}'^z |\uparrow\downarrow\rangle = -\frac{1}{4} |\uparrow\downarrow\rangle$$

$$\hat{S}^z \hat{S}'^z |\downarrow\uparrow\rangle = -\frac{1}{4} |\downarrow\uparrow\rangle \quad \hat{S}^z \hat{S}'^z |\downarrow\downarrow\rangle = \frac{1}{4} |\downarrow\downarrow\rangle$$

for the  $z$ -components and we have chosen to scale  $\hbar \mapsto 1$ . Hence  $\hat{H}_1$  is *diagonal* in our chosen basis, for which all spins are either up or down. Also:

$$\hat{\mathbf{S}}^\perp \hat{\mathbf{S}}'^\perp |\uparrow\uparrow\rangle = 0 \quad \hat{\mathbf{S}}^\perp \hat{\mathbf{S}}'^\perp |\uparrow\downarrow\rangle = \frac{1}{2} |\downarrow\uparrow\rangle$$

$$\hat{\mathbf{S}}^\perp \hat{\mathbf{S}}'^\perp |\downarrow\uparrow\rangle = \frac{1}{2} |\uparrow\downarrow\rangle \quad \hat{\mathbf{S}}^\perp \hat{\mathbf{S}}'^\perp |\downarrow\downarrow\rangle = 0$$

and so  $\hat{H}_2$  and  $\hat{H}_3$  makes the spins 'fluctuate'.

In project 2 you can analyse the role of *thermal* fluctuations on the Ising model,  $\hat{H}_1$ . In project 3 you can analyse the role of *quantum* fluctuations at zero temperature in the  $x$ - $y$  model,  $\hat{H}_2$ , or the Heisenberg model,  $\hat{H}_3$ .