

Mathematics 10: The Brillouin Zone

Any Bravais lattice has a so-called *reciprocal lattice*, defined by the requirement that for all \mathbf{R} in the Bravais lattice then \mathbf{G} is in the reciprocal lattice provided that:

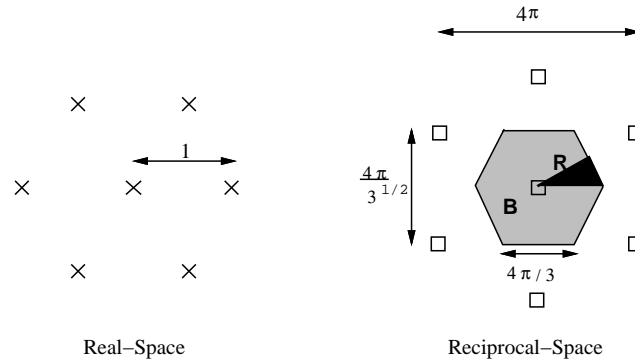
$$e^{i\mathbf{G} \cdot \mathbf{R}} = 1$$

The *Brillouin zone* is then the region in reciprocal space closer to the origin than any other reciprocal lattice vector.

Example: The triangular lattice in two-dimensions.

$$\text{Real-Space: } \mathbf{R} = n_1 \hat{\mathbf{x}}_1 + n_2 \left[\frac{1}{2} \hat{\mathbf{x}}_1 + \frac{\sqrt{3}}{2} \hat{\mathbf{x}}_2 \right]$$

$$\text{Reciprocal-Space: } \mathbf{G} = m_1 \frac{4\pi}{3} \left[\frac{1}{2} \hat{\mathbf{x}}_2 + \frac{\sqrt{3}}{2} \hat{\mathbf{x}}_1 \right] + m_2 \frac{4\pi}{3} \hat{\mathbf{x}}_2$$



The Brillouin zone is marked, **B**. When a system has *point-group symmetry*, one only need calculate in the irreducible zone, which is marked **R** and constitutes a twelfth of the full zone.

Any three-dimensional band-structure calculation involves finding the ‘electron-bands’ in an irreducible Brillouin zone.

The simplest is:

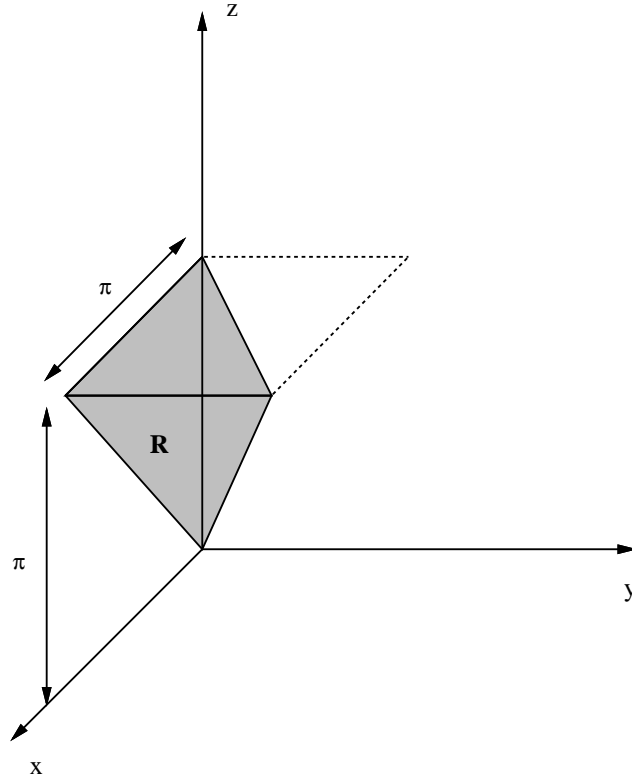
$$\gamma_{\mathbf{k}} = \frac{1}{Z} \sum_{\mathbf{n}.n.} e^{i\mathbf{k} \cdot \mathbf{R}_{\mathbf{n}.n.}}$$

where $\mathbf{R}_{\mathbf{n}.n.}$ are the nearest-neighbours to the origin, and Z is the number of such neighbours.

(1) Simple Cubic:

Real-Space: $\mathbf{R} = n_1 \hat{\mathbf{x}}_1 + n_2 \hat{\mathbf{x}}_2 + n_3 \hat{\mathbf{x}}_3$

Reciprocal-Space: $\mathbf{G} = 2\pi[m_1 \hat{\mathbf{x}}_1 + m_2 \hat{\mathbf{x}}_2 + m_3 \hat{\mathbf{x}}_3]$



$k_2 < k_1$, $k_3 > k_1$, $k_3 > k_2$, $k_3 \in (0, \pi)$ and the simplest dispersion from one degree of freedom per atom and nearest-neighbour hopping yields:

$$\gamma_{\mathbf{k}} = \frac{1}{3} [\cos k_1 + \cos k_2 + \cos k_3]$$

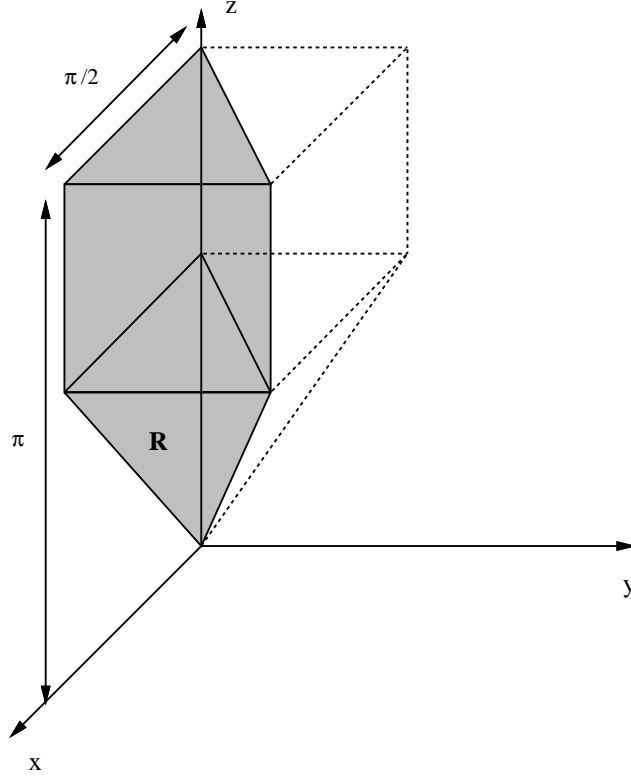
(2) Face-centre cubic:

Real-Space: $\mathbf{R} = n_1 \hat{\mathbf{x}}_1 + n_2 \hat{\mathbf{x}}_2 + n_3 \hat{\mathbf{x}}_3$

with $n_1 + n_2 + n_3 = 2i$ is even

Reciprocal-Space: $\mathbf{G} = \pi[m_1 \hat{\mathbf{x}}_1 + m_2 \hat{\mathbf{x}}_2 + m_3 \hat{\mathbf{x}}_3]$

with $m_2 + m_3 = 2i_1$, $m_3 + m_1 = 2i_2$, $m_1 + m_2 = 2i_3$, are all even



$k_2 < k_1$, $k_1 < \min(k_3, \pi/2)$, $k_2 < \min(k_3, \pi/2)$, $k_3 \in (0, \pi)$ and the simplest dispersion from one degree of freedom per atom and nearest-neighbour hopping yields:

$$\gamma_{\mathbf{k}} = \frac{1}{3} [\cos k_2 \cos k_3 + \cos k_3 \cos k_1 + \cos k_1 \cos k_2]$$

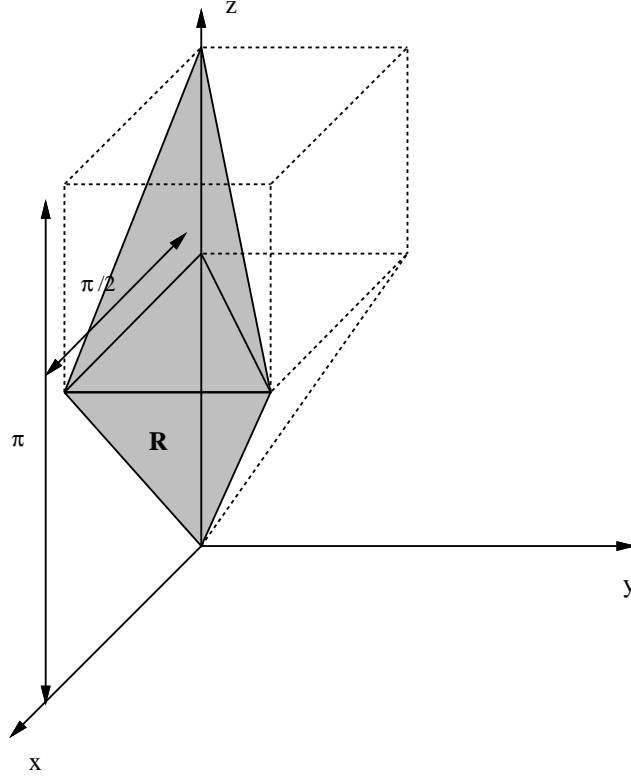
(3) Body-centre cubic:

Real-Space: $\mathbf{R} = n_1 \hat{\mathbf{x}}_1 + n_2 \hat{\mathbf{x}}_2 + n_3 \hat{\mathbf{x}}_3$

with $n_2 + n_3 = 2i_1$, $n_3 + n_1 = 2i_2$, $n_1 + n_2 = 2i_3$, are all even

Reciprocal-Space: $\mathbf{G} = \pi[m_1 \hat{\mathbf{x}}_1 + m_2 \hat{\mathbf{x}}_2 + m_3 \hat{\mathbf{x}}_3]$

with $m_1 + m_2 + m_3 = 2i$ is even



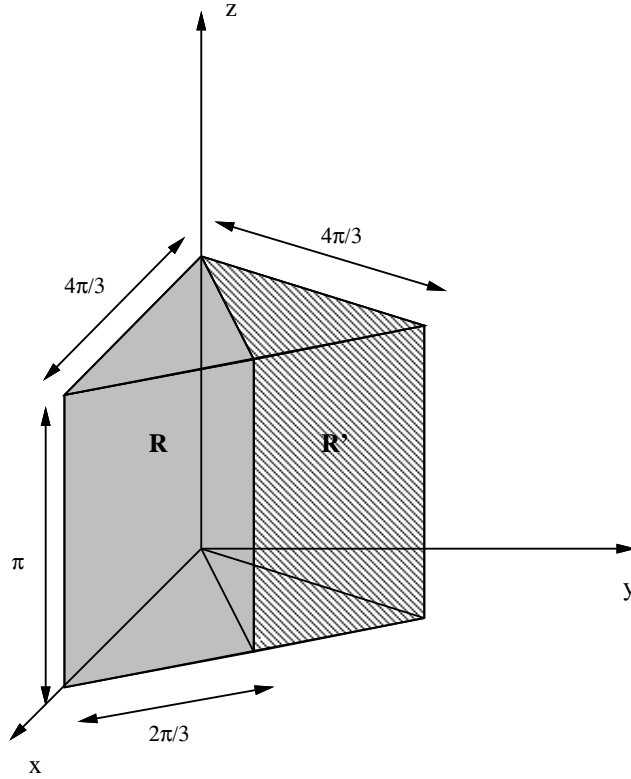
$k_2 < k_1$, $k_1 < \min(k_3, \pi - k_3)$, $k_2 < \min(k_3, \pi - k_3)$, $k_3 \in (0, \pi)$ and the simplest dispersion from one degree of freedom per atom and nearest-neighbour hopping yields:

$$\gamma_{\mathbf{k}} = \cos k_1 \cos k_2 \cos k_3$$

(4) Simple Hexagonal and Hexagonal close-pack:

Real-Space: $\mathbf{R} = n_1 \hat{\mathbf{x}}_1 + n_2 \left[\frac{1}{2} \hat{\mathbf{x}}_1 + \frac{\sqrt{3}}{2} \hat{\mathbf{x}}_2 \right] + n_3 \hat{\mathbf{x}}_3$

Reciprocal-Space: $\mathbf{G} = \frac{4\pi}{\sqrt{3}} m_1 \left[\frac{\sqrt{3}}{2} \hat{\mathbf{x}}_1 + \frac{1}{2} \hat{\mathbf{x}}_2 \right] + \frac{4\pi}{\sqrt{3}} m_2 \hat{\mathbf{x}}_2 + 2\pi m_3 \hat{\mathbf{x}}_3$



It may be best to use $R \cup R'$ and a triangular grid. $k_2 > 0$, $k_2 < \frac{k_1}{\sqrt{3}}$, $k_2 < \frac{4\pi}{\sqrt{3}} - \sqrt{3}k_1$, $k_1 \in (0, 4\pi/3)$ $k_3 \in (0, \pi)$ and the simplest dispersion from one degree of freedom per atom and nearest-neighbour hopping yields:

$$\gamma_{\mathbf{k}\pm} = \frac{\gamma_{\mathbf{k}}^T}{2} \pm \frac{\sqrt{[1 + 2\gamma_{\mathbf{k}}^T]}}{2\sqrt{3}} \cos \frac{k_3}{2}$$

in terms of the triangular lattice structure-factor:

$$\gamma_{\mathbf{k}}^T = \frac{1}{3} \left[2 \cos \frac{k_1}{2} \left[\cos \frac{k_1}{2} + \cos \frac{\sqrt{3}k_2}{2} \right] - 1 \right]$$