

# Problem Sets 1 and 2

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## Chapter 2

2.40) Given the traveling wave  $\Psi(x, t) = 5.0\exp(-ax^2 - bt^2 - 2\sqrt{ab}xt)$ , determine its direction of propagation. Calculate a few values of  $\Psi$  and make a sketch of the wave at  $t = 0$ , taking  $a = 25m^{-2}$  and  $b = 9.0s^{-2}$ . What is the speed of the wave?

After factoring notice that  $\Psi = 5.0\exp(-(x + \sqrt{\frac{b}{a}}t)^2)$ . Thus, the direction of propagation is in the **negative x direction**.

The following is  $\Psi$  as  $x$  varies.

x	$\Psi$
0.6	0.0006
0.4	0.09
0.2	1.8
0.0	5.0
-0.2	1.8
-0.4	0.09
-0.6	0.0006

Let  $t = 0$ ,  $a = 25m^{-2}$ , and  $b = 9s^{-2}$ . Then  $\Psi = -5.0\exp 25x^2$ .

$\Psi$  is a function of  $(x + vt)$ , where  $x + vt = (x + \sqrt{\frac{a}{b}}t)$ . Thus, the speed of the wave is  $v = \sqrt{\frac{a}{b}} = \sqrt{\frac{9m^2}{25s^2}} = \frac{3m}{5s} = 0.6ms^{-1}$ .

2.44) Working with exponentials directly, show that the magnitude of  $\Psi = Ae^{i\omega t}$  is  $A$ . Then, rederive the same result using Euler's formula. Prove that  $e^{i\alpha}e^{i\beta} = e^{i(\alpha+\beta)}$

Let  $\tilde{z} = Ae^{i\omega t}$  and its conjugate  $\tilde{z}^* = Ae^{-i\omega t}$ . Let the modulus, or magnitude, of  $\tilde{z}$  be  $\|\tilde{z}\| = (\tilde{z}\tilde{z}^*)^{\frac{1}{2}}$ . Thus, the magnitude of  $\Psi$  is  $\|\tilde{z}\| = (\tilde{z}\tilde{z}^*)^{\frac{1}{2}} = [(Ae^{i\omega t})(Ae^{-i\omega t})]^{\frac{1}{2}} = [A^2e^{i\omega t - i\omega t}]^{\frac{1}{2}} = \sqrt{A^2} = A$ .

Euler's formula states that  $e^{i\omega t} = \cos \omega t + i \sin \omega t$ . Notice that  $\omega = 2\pi\nu = \frac{2\pi}{t}$ . Plugging this in yields  $Ae^{i\frac{2\pi}{t}t} = A[\cos \frac{2\pi}{t}t + i \sin \frac{2\pi}{t}t]$ .

If  $t = \tau$ , then  $Ae^{i\frac{2\pi}{t}t} = A[\cos \frac{2\pi}{t}t + i \sin \frac{2\pi}{t}t] = Ae^{2\pi} = A[\cos 2\pi + i \sin 2\pi] = A[1 + 0] = A$ .

By law of exponents,  $e^x + e^y = e^{x+y}$ .

Thus,  $e^{i\alpha}e^{i\beta} = e^{i\alpha+i\beta} = e^{i(\alpha+\beta)}$ .

2.49) Show that  $\Psi(x, y, z, t) = f(\alpha x + \beta y + \gamma z - vt)$  and  $\Psi(x, y, z, t) = g(\alpha x + \beta y + \gamma z - vt)$  which are plane waves of arbitrary form satisfy the three-dimensional wave equation.

Since  $f$  and  $g$  are functions which are twice differentiable, compute the partial derivatives of  $\Psi$

$$\frac{\partial^2 \Psi}{\partial x^2} = -\alpha^2 k^2 \Psi \quad \frac{\partial^2 \Psi}{\partial y^2} = -\beta^2 k^2 \Psi \quad \frac{\partial^2 \Psi}{\partial z^2} = -\gamma^2 k^2 \Psi \quad \frac{\partial^2 \Psi}{\partial t^2} = -k^2 v^2 \Psi$$

Add the three spatial derivatives  $\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} + \frac{\partial^2 \Psi}{\partial z^2} = -k^2(\alpha^2 + \beta^2 + \gamma^2)\Psi$  and recall that  $\alpha^2 + \beta^2 + \gamma^2 = 1$ . Thus  $-k^2(\alpha^2 + \beta^2 + \gamma^2)\Psi = -k^2\Psi$ .

Plugging the second partial of  $\Psi$  with respect to  $t$  into the 3-d wave function yields  $\frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} = -k^2\Psi$ .

2.46) Take the complex quantities  $\tilde{z}_1 = (x_1 + iy_1)$   $\tilde{z}_2 = (x_2 + iy_2)$  and show that

$$\mathbf{Re}(\tilde{z}_1 + \tilde{z}_2) = \mathbf{Re}(\tilde{z}_1) + \mathbf{Re}(\tilde{z}_2)$$

If  $\mathbf{Re}$  means to extract the real components of a complex number, then since the argument  $\tilde{z}_1 + \tilde{z}_2$  resolves to  $x_1 + iy_1 + x_2 + iy_2$ , another complex number, simply extract the real component. Thus, the real component from  $x_1 + x_2 + (y_1 + y_2)i$  is  $x_1 + x_2$ . This result is the same as  $\mathbf{Re}(\tilde{z}_1) + \mathbf{Re}(\tilde{z}_2) = \mathbf{Re}(x_1 + iy_1) + \mathbf{Re}(x_2 + iy_2)$ .

2.47) Take the complex quantities  $\tilde{z}_1 = (x_1 + iy_1)$  and  $\tilde{z}_2 = (x_2 + iy_2)$  and show that  $\mathbf{Re}(\tilde{z}_1) \times \mathbf{Re}(\tilde{z}_2) \neq \mathbf{Re}(\tilde{z}_1 \times \tilde{z}_2)$ .

Start with the right hand side and notice that

$$\begin{aligned} & \mathbf{Re}(\tilde{z}_1) \times \mathbf{Re}(\tilde{z}_2) \\ &= \mathbf{Re}(\mathbf{x}_1 + i\mathbf{y}_1) \times \mathbf{Re}(\mathbf{x}_2 + i\mathbf{y}_2) \\ &= x_1 \times x_2 = x_1 x_2. \end{aligned}$$

Then the lefthand side resolves to

$$\begin{aligned} & \mathbf{Re}(\tilde{z}_1 \times \tilde{z}_2) \\ &= \mathbf{Re}((x_1 + iy_1) \times (x_2 + iy_2)) \\ &= \mathbf{Re}(x_1 x_2 + x_1 i y_2 + x_2 i y_1 + i^2 y_1 y_2) \\ &= \mathbf{Re}(x_1 x_2 - y_1 y_2 + i(x_1 y_2 + x_2 y_1)) \\ &= x_1 x_2 - y_1 y_2 \neq x_1 x_2. \end{aligned}$$

2.56) Show explicitly that the function  $\Psi(\vec{r}, t) = A \exp[i(\vec{k} \cdot \vec{r} + \omega t + \epsilon)]$  describes a wave provided that  $v = \frac{\omega}{k}$ .

In order for  $\Psi$  to describe a wave, it needs to hold true that  $v = \frac{\omega}{k}$ . Let's assume that  $\Psi$  is a wave function. Utilize the 3-D wave equation.

$$\nabla^2 \Psi - \frac{1}{v^2} \frac{\partial^2 \Psi}{\partial t^2} = 0$$

$$\nabla^2[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]] - \frac{1}{v^2} \frac{\partial^2}{\partial t^2}[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]] = 0$$

$$\frac{\partial}{\partial r}[\frac{\partial}{\partial r}[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]]] - \frac{1}{v^2}[\frac{\partial}{\partial t}[\frac{\partial}{\partial t}[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]]]] = 0$$

$$\frac{\partial}{\partial r}[ik[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]]] - \frac{1}{v^2}[\frac{\partial}{\partial t}[i\omega[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]]]] = 0$$

$$(ik)^2[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]] - \frac{1}{v^2}[(i\omega)^2[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]]] = 0$$

$$-k^2[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]] + \frac{1}{v^2}[\omega^2[A\exp[i[\vec{k} \cdot \vec{r} + \omega t + \epsilon]]]] = 0$$

$$-k^2 = -\frac{1}{v^2}(\omega^2)$$

$$k^2 = \frac{\omega^2}{v^2}$$

$$v = \frac{\omega}{k}$$

### Chapter 3

3.19) A 1.0 mW laser produces a nearly parallel beam  $1.0\text{cm}^2$  in cross-sectional area at a wavelength of 650nm. Determine the amplitude of the electric field in the beam, assuming the wavefronts are homogeneous and the light travels in a vacuum.

$$P = 1.0 \times 10^{-3}\text{W}$$

$$A = 1.0 \times 10^{-4}\text{m}^2$$

$$\lambda = 650 \times 10^{-9}\text{m}$$

$$P = \frac{E_{max}^2}{2\mu_0 c} \rightarrow E = \sqrt{\frac{2\mu_0 c P}{A}} = 2.30 \times 10^{-1}\text{Vm}^{-1}$$

3.21) The following is the expression for the  $\vec{E}$ -field of an electromagnetic wave traveling in a homogeneous dielectric:  $\vec{E} = (-100\text{Vm}^{-1})\hat{i}e^{i(kz-\omega t)}$ . Here  $\omega = 1.80 \times 10^{15}\frac{\text{rad}}{\text{s}}$  and  $k = 1.20 \times 10^7\frac{\text{rad}}{\text{m}}$ .

(a) Determine the associated  $\vec{B}$ -field.

$$B = \frac{Em}{c} = \frac{-100\text{Vm}^{-1}}{3 \times 10^8\text{ms}^{-1}} = -3.33 \times 10^{-7}\text{T}$$

(b) Find the index of refraction

$$v = f\lambda = 2\pi \frac{f\lambda}{2\pi} = \frac{\omega}{k} = \frac{1.8 \times 10^{15}\text{s}^{-1}\text{rad}}{1.2 \times 10^7\text{m}^{-1}} = 1.5 \times 10^8\text{ms}^{-1}$$

$$n = \frac{c}{v} = \frac{3 \times 10^8\text{ms}^{-1}}{1.5 \times 10^8\text{ms}^{-1}} = 2$$

(c) Compute the permittivity

$$\mu = \sqrt{\left(\frac{\mu_0 \epsilon_0}{\mu_0}\right)^{-1}} \rightarrow \epsilon^{\frac{1}{2}} = 2\sqrt{8.85 \times 10^{-12}} \rightarrow k = 3.5 \times 4 \times 10^{-12}$$

(d) Find the irradiance

$$I = \frac{C\epsilon_0 E_0^2}{2} = \frac{3 \times 10^8 \times 8.85 \times 10^{-12} \times (100)^2}{2} = 13.275\text{Wm}^{-2}$$

3.23) Consider a linearly polarized plane electromagnetic wave traveling in the positive x direction in free

space having as its plane of vibration the xy-plane. Given that its frequency is 10MHz and its amplitude is  $E_0 = 0.08Vm^{-1}$ ,

- (a) find the period of the wave
- (b) write an expression for  $E(t)$  and  $B(t)$
- (c) find the flux density,  $\langle S \rangle$ , of the wave.

(a) The period of an electromagnetic wave can be calculated as  $\tau = \frac{1}{\nu} = \frac{1}{10^7s^{-1}} = 10^{-7}s$ .

$$(b)\lambda = \frac{v}{\nu} = \frac{c}{\nu} = \frac{3 \times 10^8 ms^{-1}}{10^{-7}s^{-1}} = 30m.$$

$$\text{Since } E(x, t) = \hat{y}E_0 \cos(k_x x - \omega t) = \hat{y}E_0 \cos 2\pi(\frac{x}{\lambda} - \nu t) = \hat{y}(0.08Vm^{-1}) \cos(2\pi(\frac{x}{30} - 10^7t)).$$

Recall that  $\vec{B} = \frac{1}{c}\hat{k} \times \vec{E} = \frac{1}{c}\hat{x} \times \hat{y}\|\vec{E}\|$ . Since the wave is propagating in the x direction,  $\hat{k} = \hat{x}$ .

$$\begin{aligned} \text{Thus } \vec{B} &= \hat{z} \frac{0.08Vm^{-1}}{2.998 \times 10^8 ms^{-1}} \cos[2\pi(\frac{x}{30} - 10^7t)] \\ &= \hat{z}(2.67 \times 10^{-10} Vsm^{-2}) \cos[2\pi(\frac{x}{30} - 10^7t)] \\ &= \hat{z}(2.67 \times 10^{-10} T) \cos[2\pi(\frac{x}{30} - 10^7t)]. \end{aligned}$$

(c) The flux density is given by  $u = \epsilon_0 E_0^2 = (8.825 \times 10^{-12} \frac{C^2}{Nm^2})(0.08Vm^{-1})^2 = 5.648 \times 10^{-14} Jm^{-3}$ .

The average magnitude of the Poynting vector is

$$\langle S \rangle = \frac{c\epsilon_0}{2} E_0^2 = \frac{(3 \times 10^8 ms^{-1})(5.648 \times 10^{-14} Jm^{-3})}{2} = 8.472 \times 10^{-6} Wm^{-2}.$$