

Gravitaional Waves

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1 Tidal Field

In a gravitational field tidal tensor \mathcal{E}_{ij} is given by:

$$\mathcal{E}_{ij} = \frac{\partial^2 \phi}{\partial x^i \partial x^j} \quad (1)$$

For a body falling towards the Earth :

$$\mathcal{E}_{ij} = -\frac{GM_E}{r^5} [3x^i x^j - \delta_{ij} r^2] \quad (2)$$

Power radiated in gravitational waves is given by:

$$\frac{dW}{dt} = -\frac{1}{2} \mathcal{E}_{ij} \frac{dI^{ij}}{dt}, \quad (3)$$

where I^{ij} is called quadrupole tensor:

$$I^{ij} = \int x^i x^j \rho(x) d^3x. \quad (4)$$

Note that the moment of inertia tensor is given by:

$$\mathcal{I}_{ij} = \int (r^2 \delta_{ij} - x_i x_j) \rho(x) d^3x. \quad (5)$$

2 Binary System

Suppose that the two masses are m_1 and m_2 , and they are separated by a distance r . The power given off (radiated) by this system is:

$$\begin{aligned} P = \frac{dE}{dt} &= -\frac{32}{5} \frac{G^4}{c^5} \frac{(m_1 m_2)^2 (m_1 + m_2)}{r^5} \\ &= -\frac{32}{5} \left(\frac{G m_1}{c^2 r} \right)^2 \left(\frac{G m_2}{c^2 r} \right)^2 (m_1 + m_2) c^2 \left(\frac{c}{r} \right) \text{Watt}. \end{aligned} \quad (6)$$

For the Sun-Earth system with the mass of the Suna and earth $m_1 = 2 \times 10^{30}$, $m_2 = 6 \times 10^{24}$ Kg and distance $r = 1.5 \times 10^{11}$ m we get power radiated around 200 W (compare that with the total power radiated by Sun 3.86×10^{26} watts. This emission of gravitational radiation will decay the orbit of earth by 10^{-15} mt per day.

The rate of decrease of distance for a binary system is given by:

$$\begin{aligned} \frac{dr}{dt} &= -\frac{64}{5} \frac{G^3}{c^5} \frac{m_1 m_2 (m_1 + m_2)}{r^3} \\ &= -\frac{64}{5} \left(\frac{G m_1}{c^2 r} \right) \left(\frac{G m_2}{c^2 r} \right) \left(\frac{G (m_1 + m_2)}{c^2 r} \right) c. \end{aligned} \quad (7)$$

The lifetime of an orbit is given by:

$$\begin{aligned} t &= \frac{5}{256} \frac{c^5}{G^3} \frac{r^4}{m_1 m_2 (m_1 + m_2)} \\ &= \frac{5}{256} \left(\frac{G m_1}{c^2 r} \right)^{-1} \left(\frac{G m_2}{c^2 r} \right)^{-1} \left(\frac{G (m_1 + m_2)}{c^2 r} \right)^{-1} \frac{r}{c}. \end{aligned} \quad (8)$$

Gravitational wave strain amplitude for the the plus and cross polarization for a binary system is given by:

$$\begin{aligned} h_+ &= -2 \left(\frac{G m_1}{c^2 R} \right) \left(\frac{G m_2}{c^2 r} \right) (1 + \cos^2 \theta) \cos[2\omega(t - R)] \\ h_X &= -4 \left(\frac{G m_1}{c^2 R} \right) \left(\frac{G m_2}{c^2 r} \right) (\cos \theta) \sin[2\omega(t - R)]. \end{aligned} \quad (9)$$

where the angular velocity is given by:

$$\omega = \sqrt{\frac{G(m_1 + m_2)}{r^3}}. \quad (10)$$

For the Earth-Sun system we get:

$$h_+ = -\frac{1}{R} 1.7 \times 10^{-10} m t. \quad (11)$$

For $R = 1$ lyr we get $h_+ \approx 10^{-26}$.

3 Gravitational waves in Tensor notation

Let us consider a perturbed Minkowski metric:

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu \approx (\eta_{\mu,\nu} + h_{\mu,nu}) dx^\mu dx^\nu. \quad (12)$$

Here $\eta_{\mu,\nu}$ is the usual Minkowski metric and $h_{\mu,nu}$ represents the linearized gravitational field. Upon linalization, the coordinate invariance of full general relativity is replaced by global Lorentz invariance and local gauge invariance under infinitesimal coordinate transformaton $x^\mu \rightarrow x^\mu + \xi^\mu$

$$h_{\mu\nu} \rightarrow h_{\mu\nu} - \partial_\mu \xi_\nu - \partial_\nu \xi_\mu. \quad (13)$$

In Lorentz Gauge Einstein Field equations reduce to:

$$\left(\frac{\partial^2}{\partial t^2} - \nabla^2 \right) \bar{h}_{\mu\nu} = 16 G_N T_{\mu\nu}. \quad (14)$$

with

$$\bar{h}_{\mu\nu} = h_{\mu\nu} - \frac{1}{2} h \eta_{\mu\nu} \quad (15)$$

In transverse-traceless gauge gravitational wave propagating along z-direction can be written as:

$$h_{\mu\nu} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & h_+ & h_\times & 0 \\ 0 & h_\times & -h_+ & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}. \quad (16)$$

The lowest order contribution to gravitational waves comes from the quadruple distribution of matter:

$$h_{ij}(t, \vec{x}) = \frac{2G}{c^4} \frac{\ddot{Q}_{ij}}{r} \quad (17)$$

with

$$Q_{ij} = \int d^3x \left(x_i x_j - \frac{2}{3} x^2 \delta_{ij} \right) \rho(t, \vec{x}), \quad (18)$$

or

$$Q_{ij} = M \mathcal{R}_{ij} \quad (19)$$

or

$$h_{ij}(t, \vec{x}) = \left(\frac{2GM}{c^2 r} \right) \frac{\ddot{R}_{ij}}{c^2}. \quad (20)$$

Note that in above equations we should actually use retarded time $t - r/c$ in place of t .
Newtonian potential is given by:

$$\phi(x, y, z) = \frac{GM}{(x^2 + y^2 + z^2)^{1/2}}, \quad (21)$$

and from which we can get tidal tensor:

$$\mathcal{E}_{ij} = -\frac{GM}{r^5} (3x_i x_j - \delta_{ij} r^2). \quad (22)$$

4 Quadrouple formula

We can get wave equation with:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} \quad \text{with} \quad |h_{\mu\nu}| \ll 1. \quad (23)$$

and so we get :

$$\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \right) \bar{h}_{\mu\nu} = -\frac{8\pi G}{c^2} T_{\mu\nu}. \quad (24)$$