

UNIVERSITY OF CALIFORNIA
Los Angeles

**Correlations with Non-Photonic Electrons in
 $\sqrt{s_{NN}} = 200$ GeV AuAu Collisions in STAR**

A dissertation submitted in partial satisfaction
of the requirements for the degree
Doctor of Philosophy in Physics

by

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The dissertation of Lloyd Edward Dunkelberger Jr. is approved.

University of California, Los Angeles

2015

TABLE OF CONTENTS

1	Identification of Non-photonic Electrons	1
1.1	Outline of the NPE Identification	1
1.2	Dataset and Event Selection	1
1.3	Track Reconstruction and TPC Cuts	1
1.4	BEMC Points and Matching	1
1.5	Electron Purity	1
1.6	Photonic Electron Identification	1
1.7	Photonic Electron Reconstruction Efficiency	1
2	Azimuthal Correlations of Non-Photonic Electrons to Hadrons	2
2.1	Overview of Constructing the NPE-hadron Correlation	2
2.2	Acceptance Corrections	3
2.2.1	Single Particle ϕ -weighting	3
2.2.2	Mixed Event Background	6
2.3	Background from Flow	6
2.3.1	Measurements of Flow	6
2.3.2	Background Normalization	7
2.4	Correlations in Au+Au	10
2.4.1	Associated Hadrons	10
2.4.2	Constructing the NPE-hadron correlation	10
2.4.3	Raw Correlations	13
2.4.4	Subtracted Distributions and Yields	13
2.5	Correlations in p+p	20

2.5.1	Data and Correlations	21
2.5.2	Charm to Bottom Ratios	21
2.6	Comparisons of Yields	25
2.6.1	Away Side Shape	25
2.6.2	I_{AA}	25
2.7	Event-Plane Dependent Correlations	25
2.7.1	Event Plane Reconstruction	28
2.7.2	Correlations	31

LIST OF FIGURES

2.1	Phi distribution for all tracks in TPC.	4
2.2	p_T dependence of ϕ acceptance	5
2.3	STAR measured hadron v_2	8
2.4	STAR NPE v_2	9
2.5	Associated hadron efficiency	11
2.6	Raw Correlations 40-60% Centrality	14
2.7	Raw Correlations 20-40% Centrality	15
2.8	Raw Correlations 0-10% Centrality	16
2.9	Subtracted Correlations 40-60% Centrality	17
2.10	Subtracted Correlations 20-40% Centrality	18
2.11	Subtracted Correlations 0-10% Centrality	19
2.12	NPE-hadron correlations in p+p	22
2.13	Raw and Weighted Pythia Generated Hard Processes	24
2.14	D/B Ratio Fit Example	26
2.15	B/D Ratio	27
2.16	v_2 and Event Plane Diagram	29
2.17	Corrected Event Plane Distribution	30
2.18	Event Plane Resolutions	32

LIST OF TABLES

2.1	Associated hadron cuts	10
2.2	Yields and Errors in p+p Correlations, Low Trigger	23
2.3	Yields and Errors in p+p Correlations, High Trigger	23

CHAPTER 1

Identification of Non-photonic Electrons

We discuss the procedure for identifying electrons in events at STAR and how we remove photonic background. We show the event and track selection criteria and then lastly we will analyze the efficiency for identifying background photonic electrons. The identification of non-photonic electrons (NPE) and efficiency thereof will be critical factors when we construct the NPE-hadron correlations in later chapters.

1.1 Outline of the NPE Identification

1.2 Dataset and Event Selection

1.3 Track Reconstruction and TPC Cuts

1.4 BEMC Points and Matching

1.5 Electron Purity

1.6 Photonic Electron Identification

1.7 Photonic Electron Reconstruction Efficiency

CHAPTER 2

Azimuthal Correlations of Non-Photonic Electrons to Hadrons

We will now investigate the correlations of triggered non-photonic electrons to hadrons in $Au + Au$ and $p + p$ collisions at 200 GeV. Hard processes in these collisions will produce back to back jets in the azimuthal angle ϕ . We search for potential modification of the jet in $Au + Au$ collisions compared to $p + p$.

2.1 Overview of Constructing the NPE-hadron Correlation

Several steps are needed to produce the NPE-h correlation. The trigger particle electrons are identified by the procedure described in the previous chapter. The nonuniform acceptance of detector results in false correlations which are not a result of the underlying physics. This is corrected in two ways, the ϕ distribution of all particles is flattened and then the correlations from mixed events are calculated and then a weighting is determined so as to flatten these as well.

In correlations from $Au + Au$ collisions there is an underlying background correlation from the flow of the both the trigger electron and the associated hadron. In this analysis we only consider the second order harmonic of flow, v_2 . For hadrons, v_2 is very accurately measured across a wide range centralities and p_T . For non-photonic electrons, the measurements of v_2 are not so precise, thus we can only estimate its contribution to the background. This uncertainty will be reflected in the analysis of systematic error.

We will also look at the dependence of the correlation on the angle between the trig-

gered electron and the event plane. A dependence on this angle could point to path length dependence on the jet suppression in QGP.

2.2 Acceptance Corrections

The STAR detector give full 2π azimuthal coverage, however there are still regions of the detector which have noticeably poorer efficiency. This causes an uneven azimuthal acceptance which in turn lead to spurious correlations between detected particles. To correct this we need to apply a weighting to each track depending on where it is in the detector.

2.2.1 Single Particle ϕ -weighting

We begin by correcting for the single particle acceptance in ϕ . The boundaries between sectors produce regions of lower efficiency, also in Run11 one sector of the TPC had noticeably lower efficiency than the rest. These effects are both shown in Figure 2.1.

The dependence of the acceptance on ϕ however is not the same for all tracks. Whether a track crosses a sector boundary or passes through the dead sector will depend on that particular track's geometry. Tracks at low p_T curve more in the magnetic field and thus the effects of these lower efficiency areas apply to wider regions in track ϕ . The dependence of acceptance on p_T is shown in Figure 2.2. At low p_T the dependence is especially strong thus for $p_T \leq 1$ GeV/c we divide tracks into p_T bins of .1 GeV/c, which is near the limit of the momentum resolution of the TPC. Above 1 GeV/c the tracks are roughly straight so the effects on acceptance from the sector boundaries and dead sector are consistent bin-to-bin up to arbitrarily large p_T .

While the dependence of acceptance on track p_T is by far the largest effect, we still further subdivide the tracks to make acceptance corrections. It is possible for the acceptance to depend on η , and we are especially concerned with edge effects when $|\eta| \sim 1$, thus we divide into 4 even bins in pseudorapidity ranging from -1 to 1.

Likewise we account for dependence on the event vertex (in both $p + p$ and $Au + Au$)

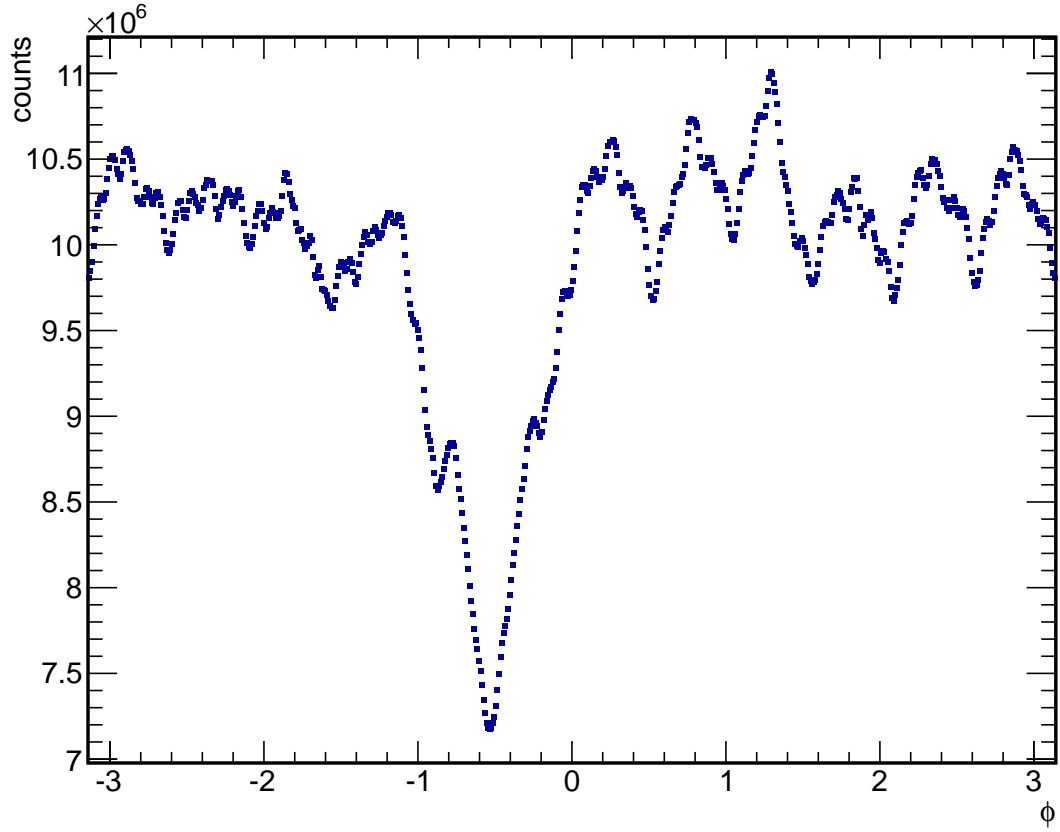


Figure 2.1: The azimuthal angular distribution of all tracks in Run11 $Au + Au$ collisions at 200 GeV. Periodic bumps can be seen from the sector boundaries, as well as a dip in the poorly performing sector.

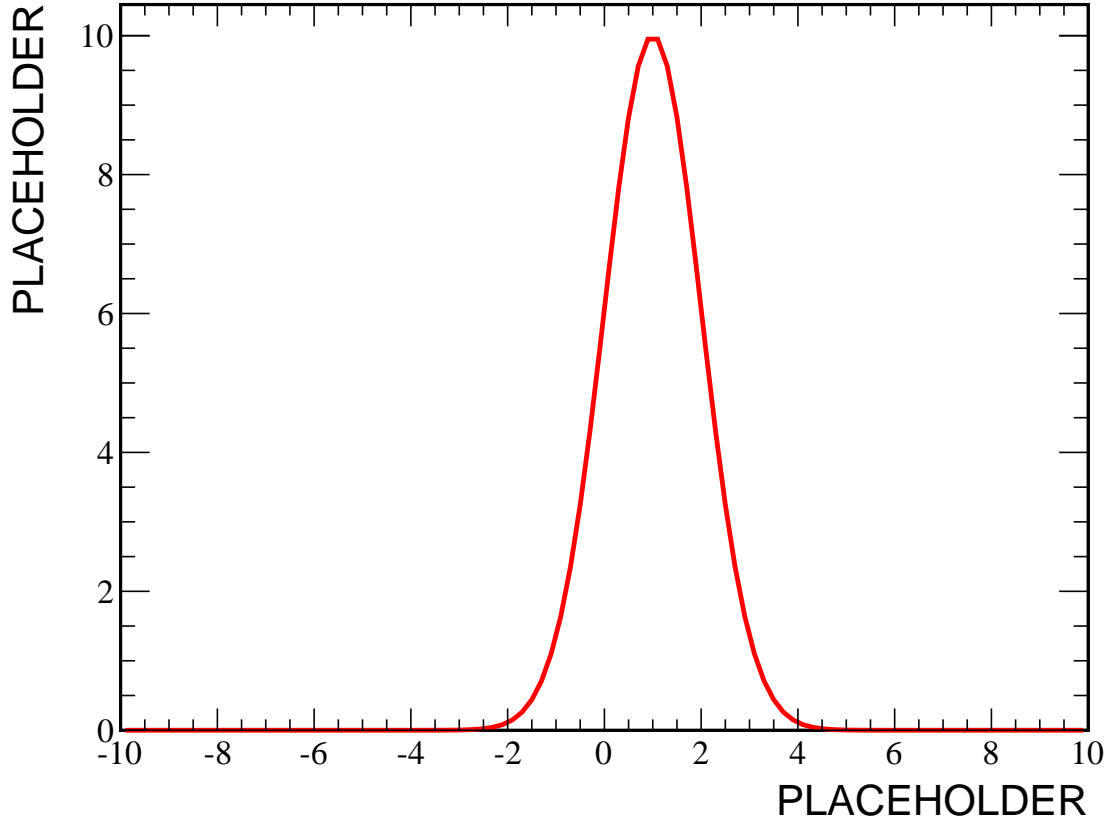


Figure 2.2: ϕ distributions for single particles in different p_T bins. Strong p_T dependence is seen especially below 1 GeV/c due to the different track geometries.

and multiplicity (only for $Au + Au$) by dividing into bins of vertex-z and centrality. For the centrality bin divisions, all centrality bins from 30% – 80% are taken together since in the peripheral bins the statistics are too low to get a reliable acceptance correction.

Finally, since the tracks in the TPC are curved, there will be a dependence on which direction the track curves. For example, two particles may start on opposite sides of a sector boundary separated by some distance in ϕ but both may cross the boundary if they curve in opposite directions. So we need to take separate weightings based on the product of the magnetic field and the particle’s charge, $B \cdot q$.

After calculating the single ϕ correction we apply it to each track in the analysis whenever we calculate event planes or 2-particle correlations. Since some areas of the detector have very low efficiencies they can introduce huge weights for a small number of particles. This can destabilize results, so we cap the weight an individual particle can get at 5.0.

2.2.2 Mixed Event Background

To further correct for nonuniformities in detector acceptance we use a mixed event weighting. In an ideal detector the correlations of trigger particles to associated hadrons from a different event should be flat, however acceptance effects will result in nonphysical correlations which need to be removed.

Similar to the single particle corrections we divide the mixed event corrections into bins to account for various systematic differences. In mixed event we bin according to associated particle p_T , triggered particle p_T , centrality, vertex z position, and η . As in single particle corrections, the most extreme bin to bin variations occur between the low associated p_T bins.

2.3 Background from Flow

2.3.1 Measurements of Flow

The motivations behind two-particle correlation studies are typically the investigation of jet modification in QGP and the response of the medium to jets. But even in the absence of jets

we still expect to see some correlation within events from flow. The azimuthal anisotropy resulting from the second order flow harmonic, v_2 , of both the trigger and associated particles produces a background shape with the form:

$$B[1 + v_2^{trig}v_2^{asso} \cos(2\Delta\phi)] \quad (2.1)$$

where B is an overall constant factor. Higher order harmonics v_3 , v_4 , etc. can also contribute to the background. Large v_3 in particular is a potential explanation for some of the results in dihadron correlations, but these effects are not considered for this analysis.

Hadron v_2 has been measured to high precision in a wide range of p_T bins at STAR. Figure 2.3 shows the results of STAR v_2 measurements using an event plane method and illustrates the general dependence on p_T and centrality. To calculate the hadron v_2 we extrapolate the v_2 measurement to the center of the associated hadron p_T bin. Then when looking at correlations across multiple hadron p_T bins we use the weighted average of v_2 based on the number of hadrons in each p_T bin.

Measurements of electron v_2 at STAR have shown that non-photonic electrons also have large elliptic flow. Because of limited statistics electron v_2 is measured in much larger p_T and centrality bins. Various measurements of NPE v_2 are seen in Figure 2.4, showing that they tend to fall in a range between .05 and .15 depending on the measurement procedure. For this analysis we assume that NPE v_2 is .1 in all bins, we then vary the NPE v_2 between .05 and .15 and take the difference in final correlations as a systematic error.

2.3.2 Background Normalization

Knowing the values of v_2 for hadrons and non-photonic electrons, we then need to determine the overall normalization constant B as in Equation 2.1. There are two simple ways of estimating this, both relying on the assumption that the jet like contributions to the azimuthal correlation are concentrated in peaks around 0 and π , and that any remaining correlations there are the result of the underlying v_2 background.

In one case we can simply pick a point between the near and away sides and then set B so

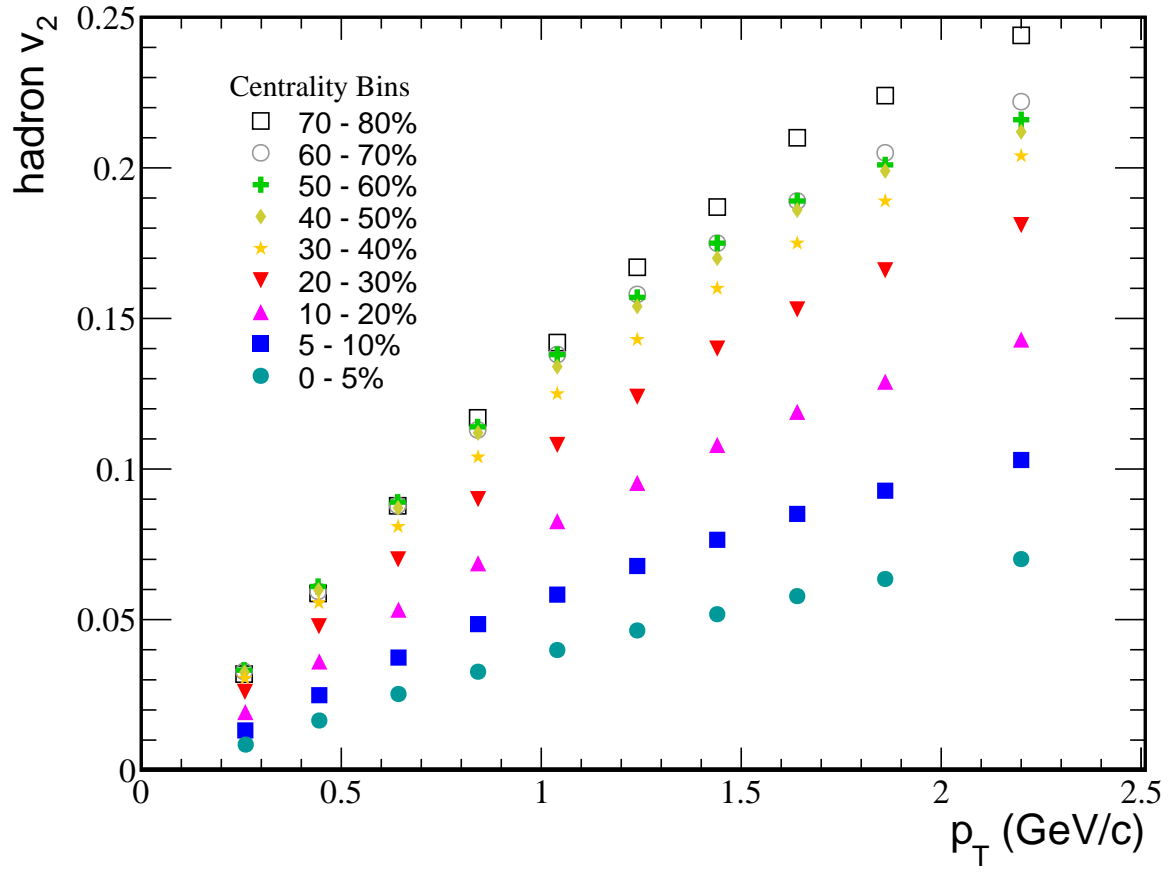


Figure 2.3: Measured v_2 values for hadrons across a range of p_T and centralities.

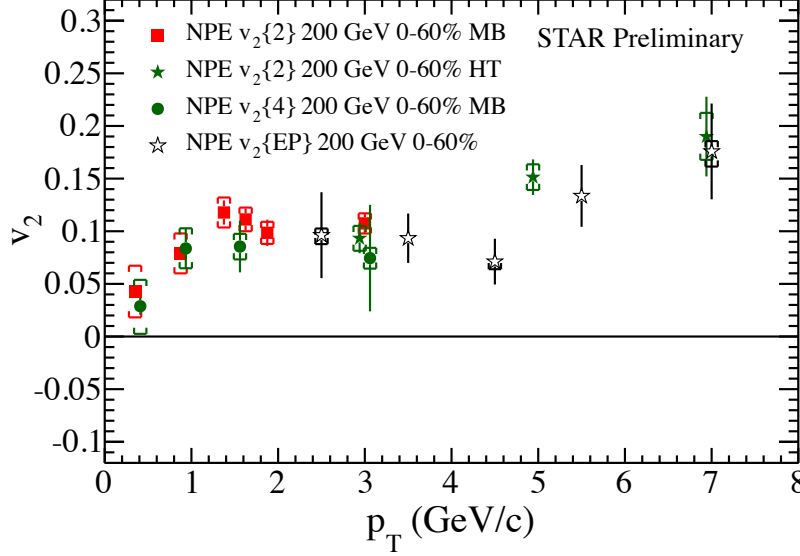


Figure 2.4: Various measurements of NPE v_2 in STAR. Going forward we assume .1 to be the value for NPE v_2 in all bins.

that the overall yield of particles above background at that point is 0. This point is typically taken to be around 1 radian and thus this method is called the zero yield at 1 (ZYA1) normalization. Although when we implment ZYA1 normalization we take the lowest absolute yield of the 3 points closest to 1 radian. Alternatively we can instead pick the point in the raw correlation with lowest value and normalize so that that point produces zero yield. This is the zero yield at minimum (ZYAM) method. These methods tend to coincide in practice and unless otherwise noted we use ZYAM normalization. There is another technique called absolute background subtraction used by PHENIX in their NPE-hadron correlation measurement but we do not use this method.

When using ZYAM or ZYA1 normalization our background subtracted yield can be very suceptible to downward fluctuations of points causing an abnormally high yield. To account for this we also look at the effect of normalizing to the next highest point in the correlation. We then compare the values of B that we get and then quote the difference as the systematic error of background normalization.

Variable	Cut
Track Type	$< .5$ (Primary)
Global DCA	$< 2.0\text{cm}$
η	$\in (-1.0, 1.0)$
p_T	$\geq .2 \text{ GeV}$

Table 2.1: Cuts for associated hadrons used in e-h correlations

2.4 Correlations in Au+Au

We will now look at putting together the results of the previous sections and creating the NPE-h correlation in Au+Au collisions. We will then discuss the results in Au+Au before moving on to p+p and event plane dependent correlations.

2.4.1 Associated Hadrons

The basic quantity we will measure is the yield $\frac{dN}{d\Delta\phi}$ of associated hadrons at various relative to some triggered electron. For the associated hadrons the cuts we use are summarized in Table 2.1.

The correlations are further broken up into bins in event centrality and associated hadron p_T . This is the point at which we apply the acceptance corrections from the single particle ϕ weighting as well as the mixed event weighting. Additionally we also correct the yield for the efficiency of the associated hadron yield. The TPC efficiency is lower for the high occupancy events in central collisions, and efficiency is also significantly worse for very low p_T hadrons. The efficiency is calculated from embedding and the results are summarized in Figure 2.5.

2.4.2 Constructing the NPE-hadron correlation

Now with azimuthal electron-hadron correlation functions we look at how we create the NPE-h correlation. The definition of the NPE-h correlation is:

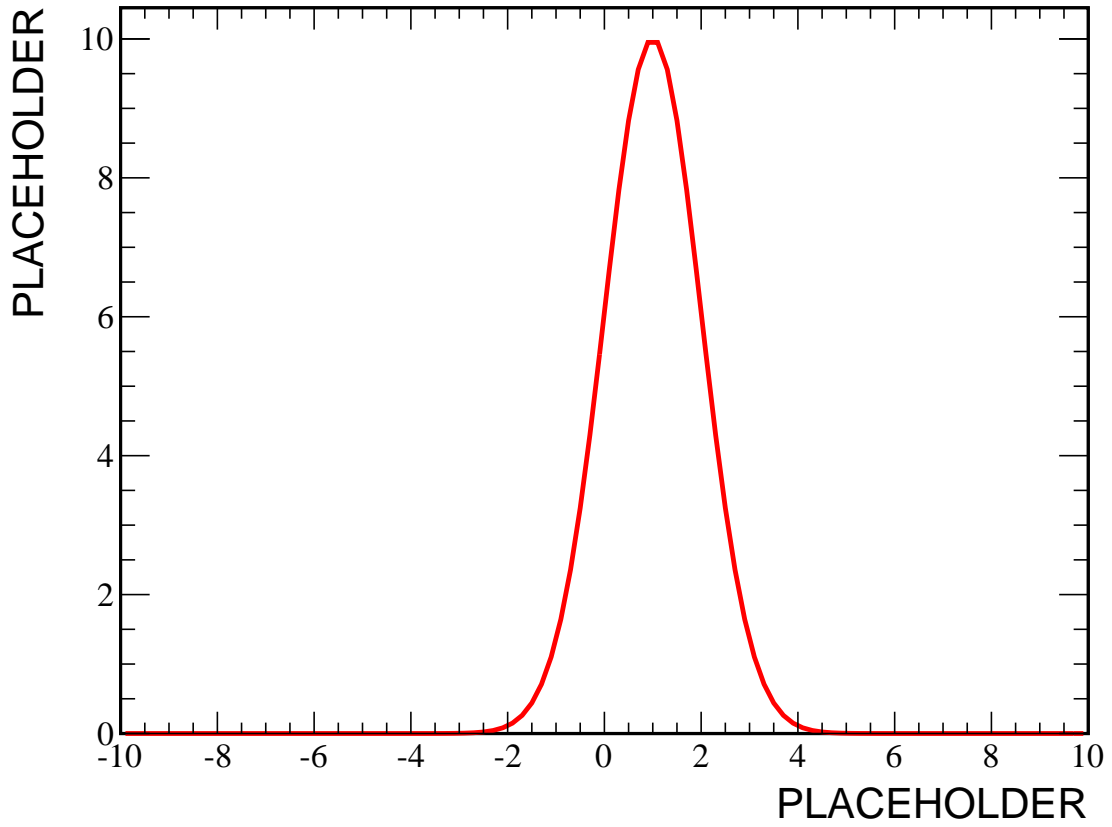


Figure 2.5: The TPC efficiency for hadrons as a function of hadron p_T . The different plots are for different centralities which correspond to 0-5%, 5-10%, 10-20%, 20-30%, 30-40%, 40-50%, and 50-60%.

$$\frac{dN_{NPE-h}}{d\Delta\phi} = \frac{dN_{semi-h}}{d\Delta\phi} - \left(\frac{1}{\epsilon_\gamma} - 1\right) \frac{dN_{photonic-h}}{d\Delta\phi} + \frac{dN_{same-h}}{d\Delta\phi} \quad (2.2)$$

An explanation of these terms:

- **Semi-inclusive electrons:** This is the correlation of inclusive electrons for which no photonic partner track could be found. This sample will include many non-photonic electrons as well as some photonic background for which we could not find a partner track.
- **Unidentified photonic electrons:** the term:

$$\left(\frac{1}{\epsilon_\gamma} - 1\right) \frac{dN_{photonic-h}}{d\Delta\phi}$$

is intended to remove the remaining photonic background triggers from the semi-inclusive sample. To do this we take the correlation for identified photonic electrons to hadrons and scale it up according to the estimated photonic electrons reconstruction efficiency, ϵ_γ . The reconstruction efficiency is determined by embedding simulations.

- **Same-sign electrons:** The method for identifying photonic electrons, pairing all tracks and calculating DCAs and invariant masses, will result in some oversubtraction of NPE signal. We account for the combinatorially removed points by looking at the results of same sign pairing tracks, i.e. the tracks which pass all of the photonic partner cuts except that they have the same sign. We add this term back make up for the NPE signal which was removed by the previous two terms.

There is also the potential for contamination of the triggered electrons with hadrons. This would require the subtraction of a dihadron correlation term:

$$\frac{dN_{h-h}}{d\Delta\phi}$$

We expect the purity of our triggered electrons to be high in the relevant p_T ranges so for this analysis we will not include it.

2.4.3 Raw Correlations

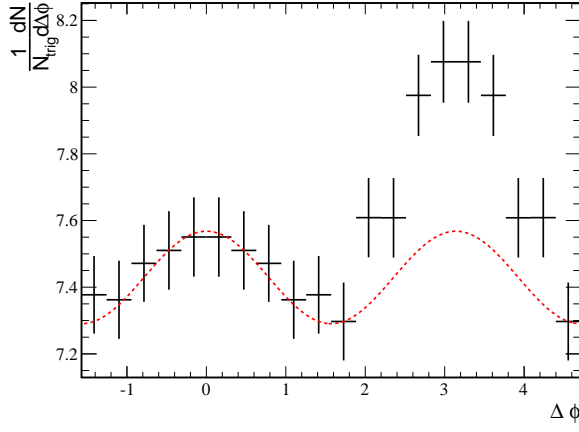
The raw correlation is the distribution $\frac{dN_{NPE-h}}{d\Delta\phi}$ before we subtract the background from v_2 . The subtraction and correction spelled out in Equation 2.2 has already been performed and what is shown in the following figures is the NPE-h correlation with no background subtraction. The raw correlations serve as an initial check of the correlation method to spot any problems with our procedure.

Figures 2.6, 2.7 and 2.8 show the raw correlations in 200 GeV AuAu collisions and that they conform to our rough expectations. Overall particle yields are also higher at lower p_T and are much higher in central events where multiplicity is higher. The general trend is for particle yields to be higher around 0 angle relative to the triggered NPE and at π , this is normal dijet distribution which is seen in hard processes. We also see that these dijets sit on top of a modulated background from v_2 . We can see that the calculated backgrounds are reasonable and we also get a sense of the performance and limitations of the ZYAM method. For example in Figure 2.8a we see that a low fluctuation in one bin may have pulled down the normalization causing the near side peak to sit farther above the background. We will account for these types when we estimate the systematic uncertainties.

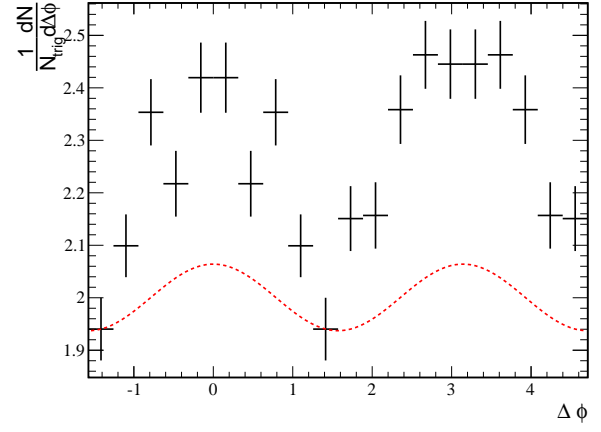
2.4.4 Subtracted Distributions and Yields

We now want to look at how the jet-like distributions of particles changes as a function of collision centrality and trigger particle p_T . We subtract off the background from the underlying event and v_2 to examine effects of the heavy quark fragmentation and propagation through the medium. The subtracted plots are summarized in Figures 2.9, 2.10 and 2.11.

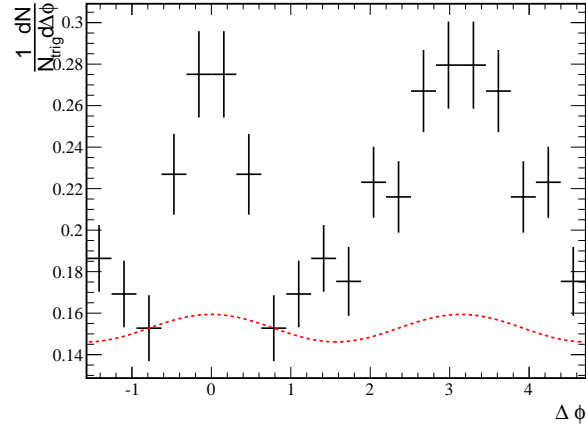
For these NPE-h correlations we also consider three sources of systematic error: Uncertainty from NPE v_2 , uncertainty in photonic electron reconstruction efficiency, and background normalization. Results NPE v_2 are over wide ranges in p_T and centrality and are roughly around .1 we take that value when calculating the background but we also calculate backgrounds with v_2 of .05 and .15. We then take the difference between these extremes as the uncertainty.



(a) $0.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$

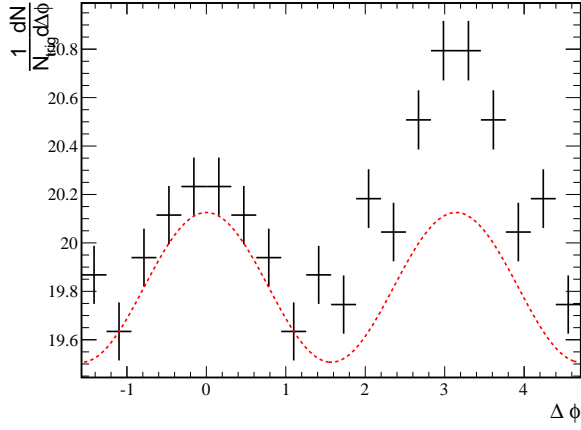


(b) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$

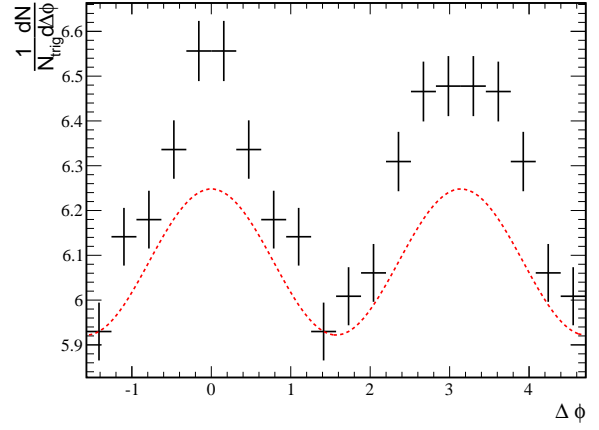


(c) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

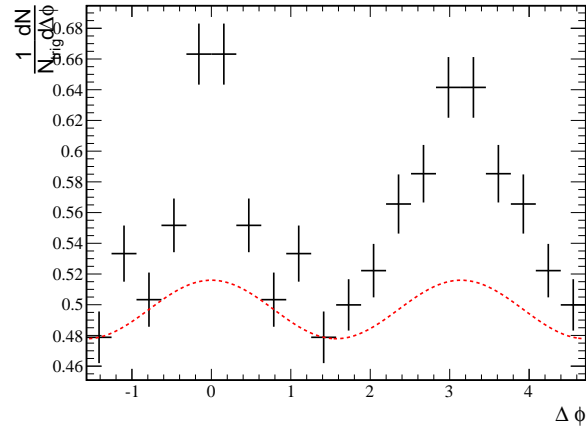
Figure 2.6: Raw NPE-h Correlations for 40-60% centrality events. Trigger p_T is $4.0 \text{ GeV}/c \leq p_{T,trig} \leq 6.0 \text{ GeV}/c$



(a) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$

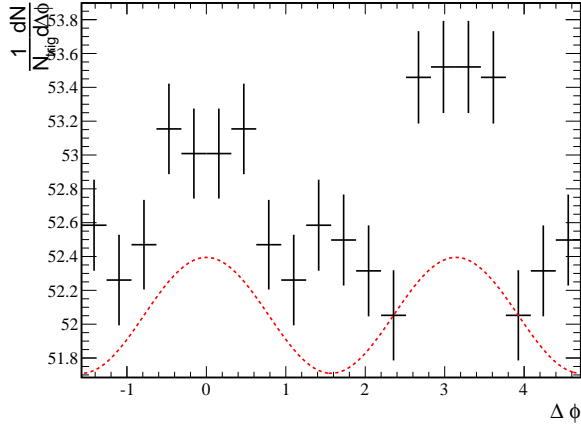


(b) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$

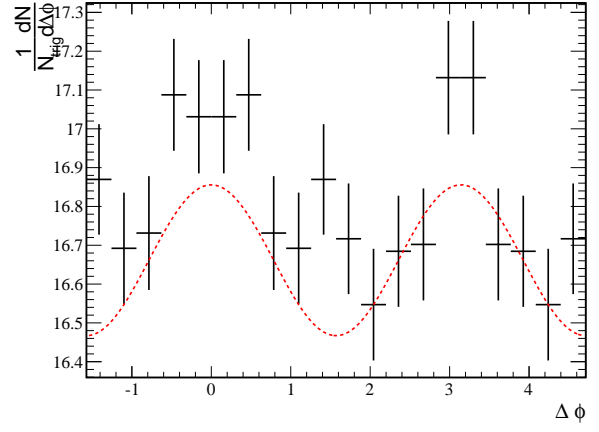


(c) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

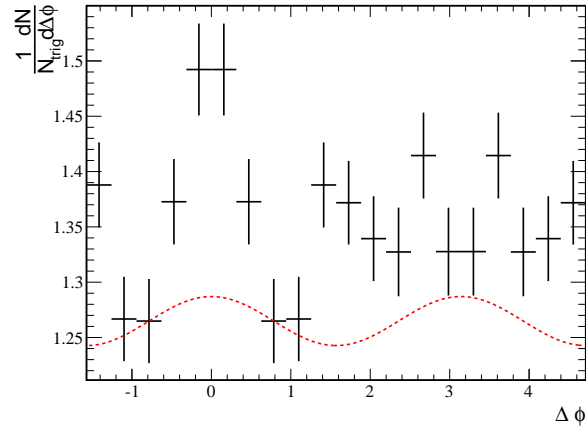
Figure 2.7: Raw NPE-h Correlations for 20-40% centrality events. Trigger p_T is $4.0 \text{ GeV}/c \leq p_{T,trig} \leq 6.0 \text{ GeV}/c$



(a) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$

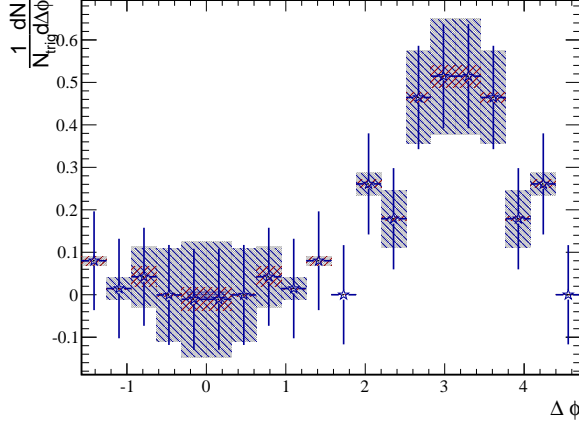


(b) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$

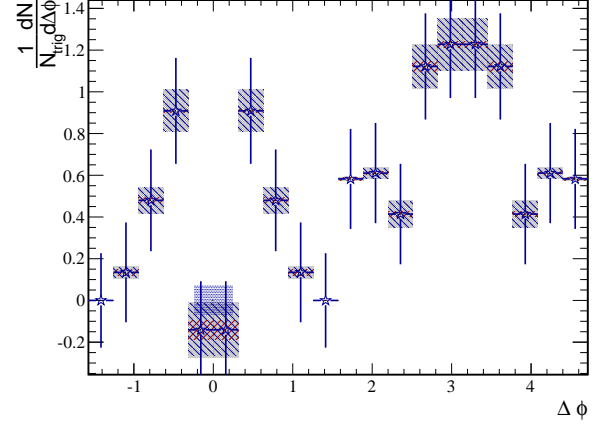


(c) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

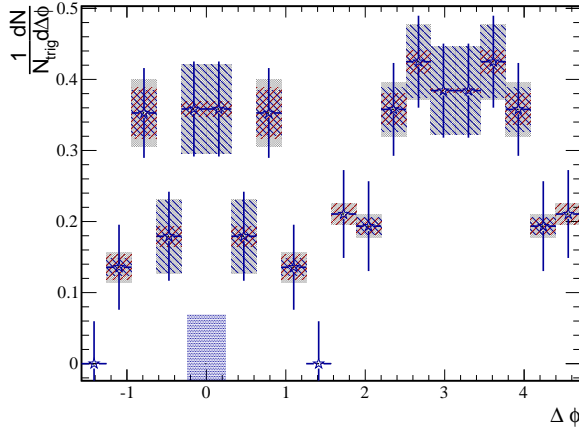
Figure 2.8: Raw NPE-h Correlations for 0-10% centrality events. Trigger p_T is $4.0 \text{ GeV}/c \leq p_{T,trig} \leq 6.0 \text{ GeV}/c$



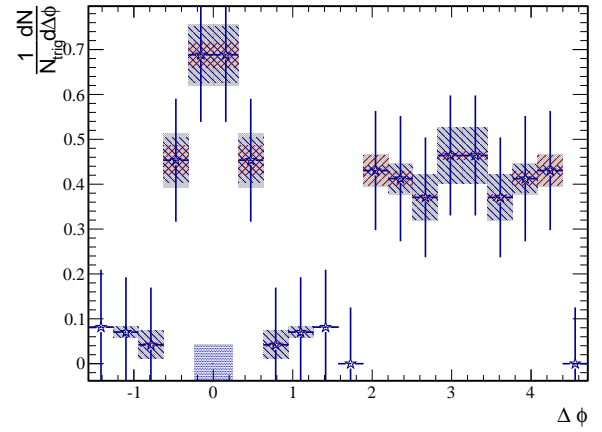
(a) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



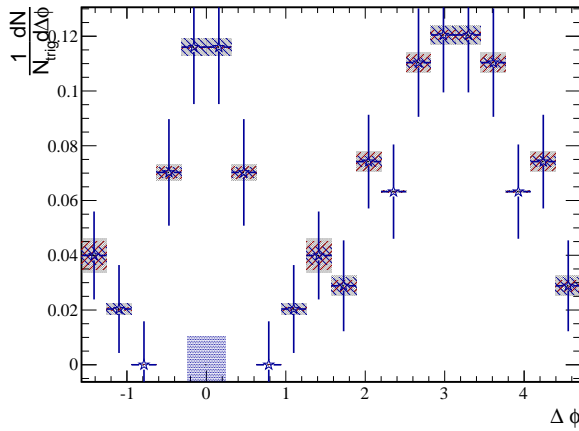
(b) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



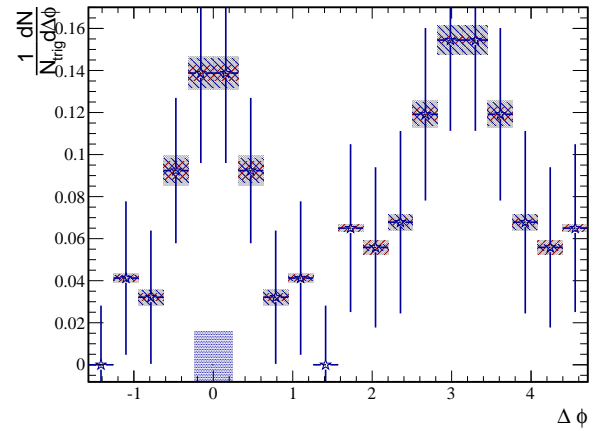
(c) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$



(d) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$

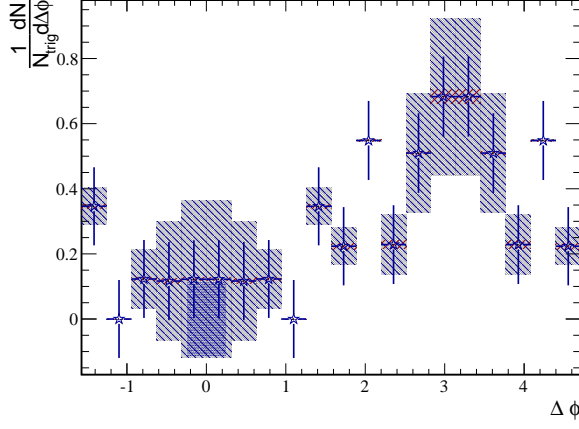


(e) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

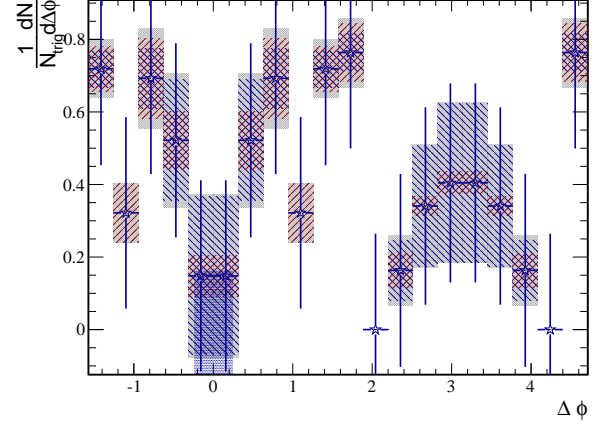


(f) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

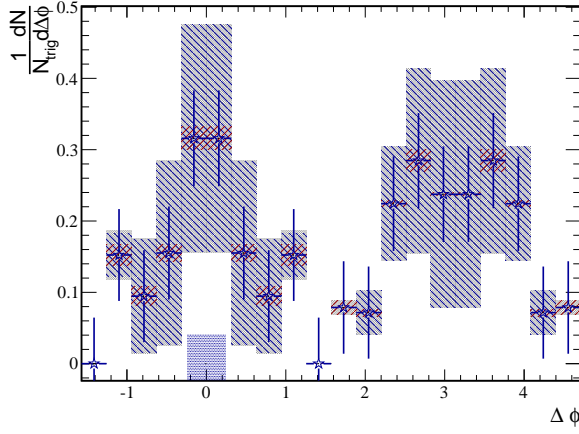
Figure 2.9: Background subtracted NPE-h correlations for 40-60% centrality events. Trigger p_T is $4.0 \text{ GeV}/c \leq p_{T, \text{trig}} \leq 6.0 \text{ GeV}/c$



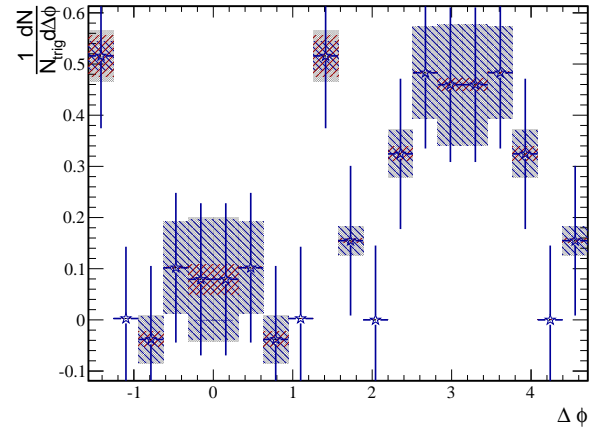
(a) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



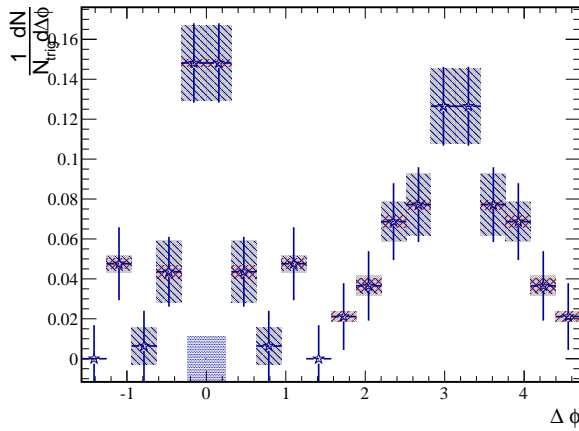
(b) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



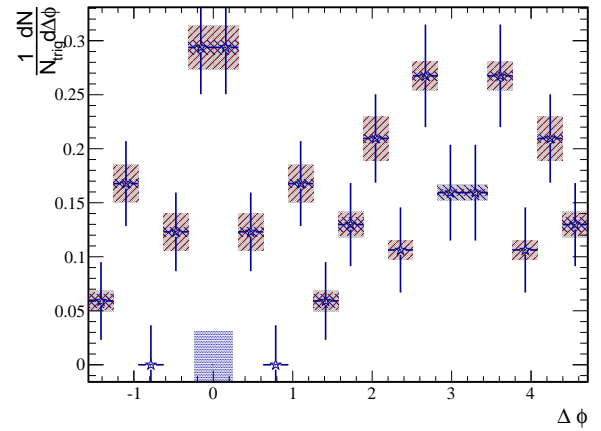
(c) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$



(d) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$

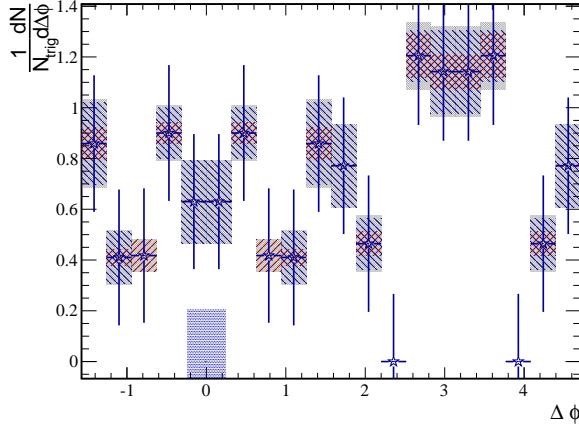


(e) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

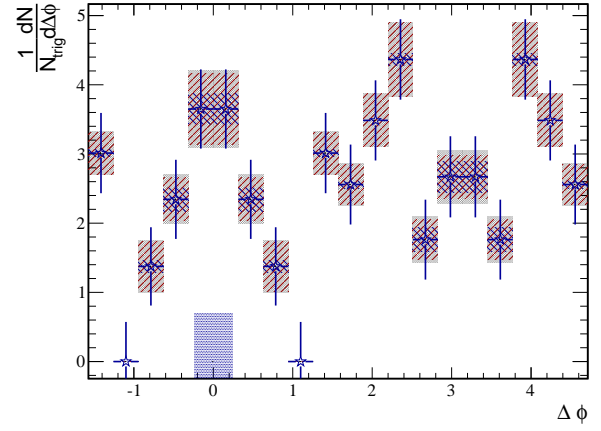


(f) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

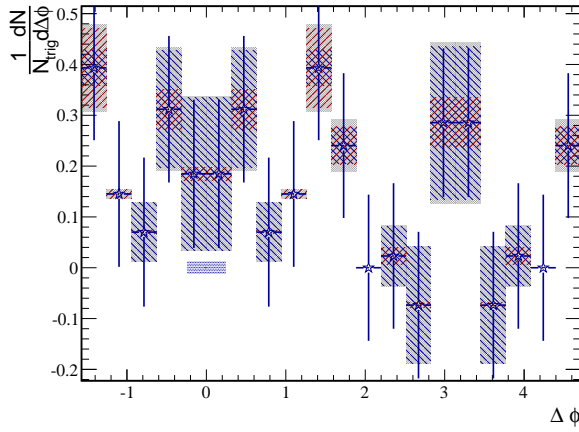
Figure 2.10: Background subtracted NPE-h correlations for 20-40% centrality events. Trigger p_T is $4.0 \text{ GeV}/c \leq p_{T, \text{trig}} \leq 6.0 \text{ GeV}/c$



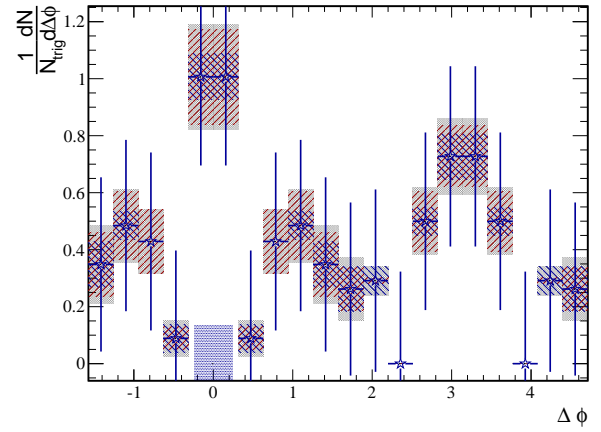
(a) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



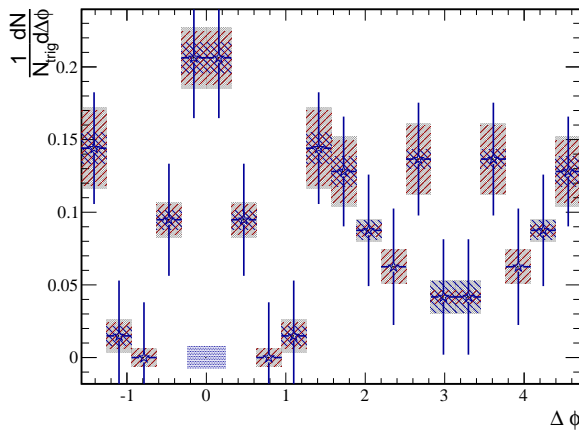
(b) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



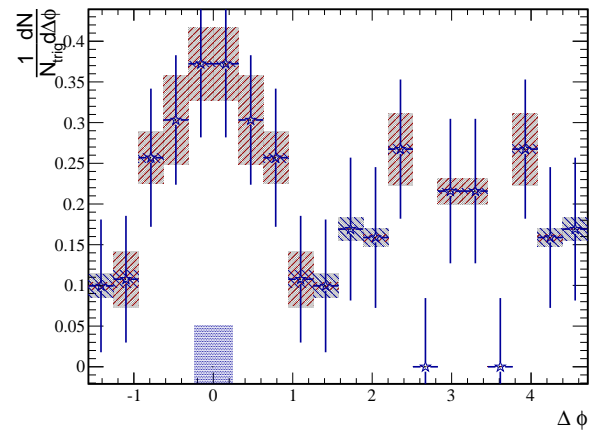
(c) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$



(d) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$



(e) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$



(f) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

Figure 2.11: Background subtracted NPE-h correlations for 0-10% centrality events. Trigger p_T is $4.0 \text{ GeV}/c \leq p_{T,trig} \leq 6.0 \text{ GeV}/c$

The photonic electron reconstruction efficiency (ϵ_γ) is determined from embedding simulations but the extracted values tend to vary from analysis to analysis. To be safe we allow the efficiency to vary by 10% and then take the difference in distributions as the error. This is done point by point. The combined NPE v_2 and ϵ_γ systematics are represented on the plots by the shaded region around the points. The NPE v_2 error tends to be the dominant source of uncertainty and the systematics are much larger for lower associated hadron p_T .

The systematic uncertainty from background normalization is calculated by performing the ZYAM procedure on the two lowest points in the correlation. The difference in normalization factors is taken as the uncertainty and we display this as a shaded bar at 0 yield and 0 angle. This uncertainty would move all points together in a uniform manner.

The subtracted distributions give some insight into the interactions of heavy quarks with the QGP medium. For all trigger p_T shown the direction of trigger electron is well correlated to the direction of the parent B or D meson. Thus we look to the near and away side yields for clues to the nature of the initially created heavy quarks interactions. We calculate the background subtracted yield for the near side region $\Delta\phi \leq .942$, in the away side “head” $\Delta\phi \geq 2.2$, and in the away side shoulder $1.25 \leq \Delta\phi \leq 2.2$. For the separated shoulder and head regions we are looking for signs of away side broadening in the correlation. We might expect to find that the ratio of the shoulder to head yields is larger in more central collisions as a result of jets being diverted or smeared out as a result of interactions with the QGP. We can also look for evidence of medium responses to a heavy quark traversing it. The yields from these plots are listed in Table AAA. We will summarize these results once we have the correlations from p+p as well.

2.5 Correlations in p+p

With correlations from Au+Au collisions we can study the effects on observed particles resulting from heavy quark interactions with the medium. By looking across centralities we can select different fireball sizes and durations to see how the presence of QGP affects the formation of dijets. Now we can also look at p+p collisions also at $\sqrt{s_{NN}} = 200$ GeV to see

the correlation without any QGP and use this as a baseline for comparison with our Au+Au results.

NPE-h correlations have also been used to study the charm to bottom produced in these collisions. This is done by fitting the observed correlations to Pythia simulations of NPE-h correlations from charm and bottom decays. Those resulting from bottom will have a broader distribution because of the higher mass of the B mesons compared to D . We will show a calculation of this as a consistency check with previous NPE-h analyses.

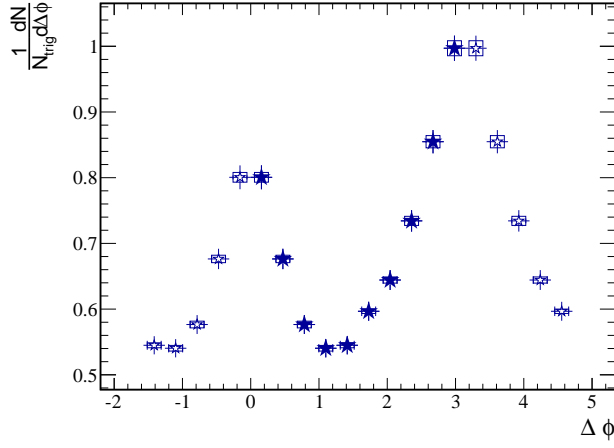
2.5.1 Data and Correlations

The dataset for the p+p correlations is the BHT triggered events in STAR run 12 p+p 200 GeV. The procedure for identifying non-photonic electrons and constructing the NPE-h correlation is nearly identical to Au+Au. We still need to perform the acceptance corrections as in Au+Au, however because of the lower multiplicities in p+p collisions it is difficult to get enough statistics for mixed event correlations so we will rely only on the single particle ϕ weighting. In run 12 the TPC performed much better and has a more uniform acceptance than in run 11 so practically these correction are far less important.

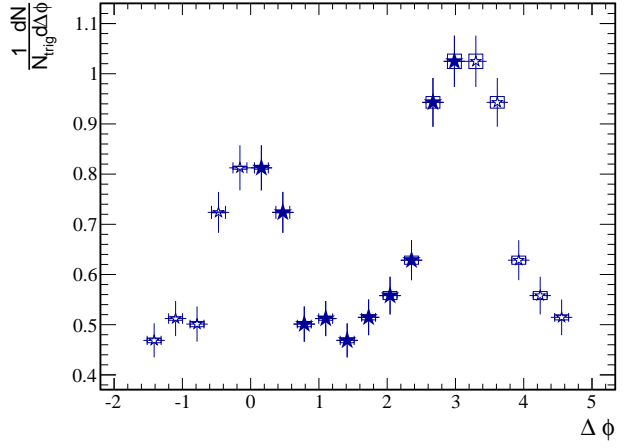
In p+p correlations there is no need to perform the background subtraction as in Equation 2.1 since there is no elliptic flow in p+p collisions. So we no longer need to consider raw correlations, we can just take the results from Equation 2.2 and use those as our correlations. Since there is no NPE-h v_2 and no need to normalize to some background distribution we no longer have to consider 2 of the 3 sources of systematic uncertainty present in Au+Au. We only need to account for uncertainty in ϵ_γ , the photonic electron reconstruction efficiency. We do this as in Au+Au collisions, allowing the efficiency to vary by 10%. Tables 2.2–2.3 summarize the yields with errors obtained from p+p collisions.

2.5.2 Charm to Bottom Ratios

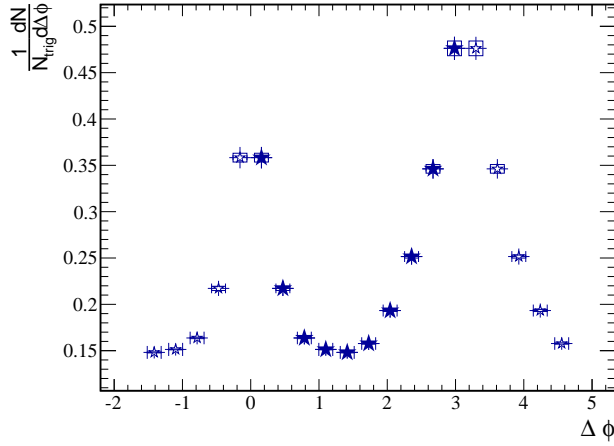
We can use the p+p NPE-h correlations to investigate the relative contributions of charm and bottom to non-photonic electrons by fitting the observed correlations with those from



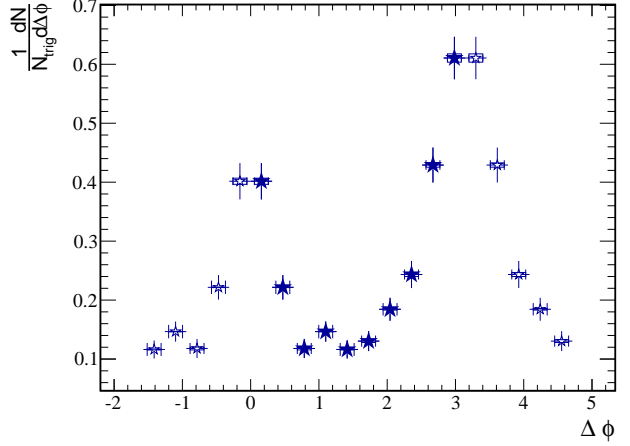
(a) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



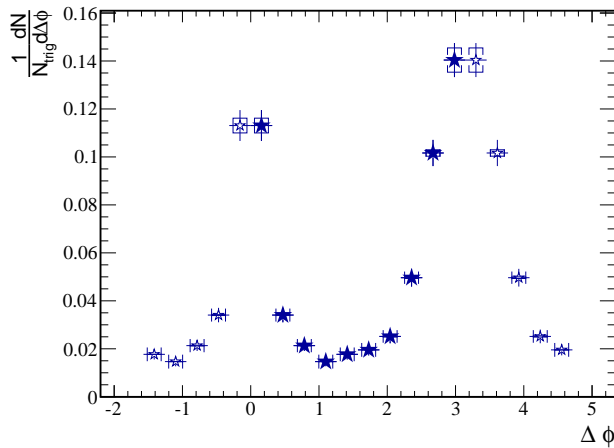
(b) $.5 \text{ GeV}/c \leq p_{T,h} \leq 1.0 \text{ GeV}/c$



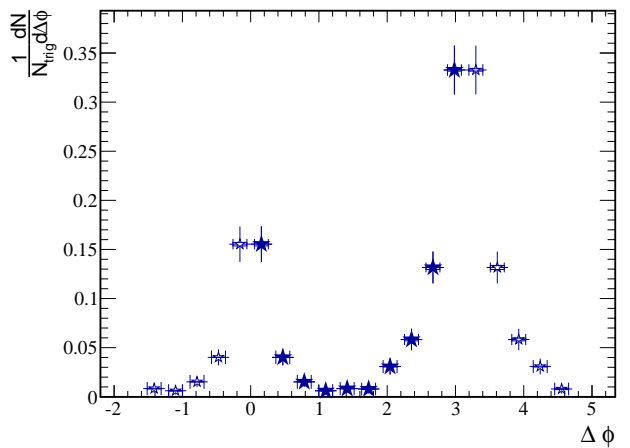
(c) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$



(d) $1.0 \text{ GeV}/c \leq p_{T,h} \leq 2.0 \text{ GeV}/c$



(e) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$



(f) $2.0 \text{ GeV}/c \leq p_{T,h} \leq 4.0 \text{ GeV}/c$

Figure 2.12: NPE-h correlations for p+p collisions at 200 GeV left column shows triggers with $4.0 \text{ GeV}/c \leq p_T \leq 6.0 \text{ GeV}/c$ and right column is $6.0 \text{ GeV}/c \leq p_T \leq 9.0 \text{ GeV}/c$.

Associated p_T	$\Delta\phi$ region	Yield $\frac{1}{N_{trigger}\Delta\phi}$	Stat. Error	Sys. Error
$p_{T,asso} \in (.5, 1.0)$ GeV/c	Near-side	0.814888	0.00977926	0.00360017
$p_{T,asso} \in (.5, 1.0)$ GeV/c	Head	0.812439	0.00965688	0.00511522
$p_{T,asso} \in (.5, 1.0)$ GeV/c	Shoulder	0.561056	0.00786714	0.00216749
$p_{T,asso} \in (1.0, 2.0)$ GeV/c	Near-side	0.279777	0.00536135	0.00167489
$p_{T,asso} \in (1.0, 2.0)$ GeV/c	Head	0.337404	0.00580794	0.00293765
$p_{T,asso} \in (1.0, 2.0)$ GeV/c	Shoulder	0.156817	0.00379766	0.000439399
$p_{T,asso} \in (2.0, 4.0)$ GeV/c	Near-side	0.057533	0.00242196	0.00096921
$p_{T,asso} \in (2.0, 4.0)$ GeV/c	Head	0.0916335	0.0030426	0.00164648
$p_{T,asso} \in (2.0, 4.0)$ GeV/c	Shoulder	0.019603	0.00132994	8.75715e-05

Table 2.2: Yields and Errors from NPE-h correlations in p+p collisions with trigger 4.0GeV/c $\leq p_t \leq 6.0$ GeV/c.

Associated p_T	$\Delta\phi$ region	Yield $\frac{1}{N_{trigger}\Delta\phi}$	Stat. Error	Sys. Error
$p_{T,asso} \in (.5, 1.0)$ GeV/c	Near-side	0.801004	0.0245826	0.00217934
$p_{T,asso} \in (.5, 1.0)$ GeV/c	Head	0.815639	0.025429	0.00632914
$p_{T,asso} \in (.5, 1.0)$ GeV/c	Shoulder	0.48424	0.0194036	0.00332109
$p_{T,asso} \in (1.0, 2.0)$ GeV/c	Near-side	0.278887	0.0138082	0.00181987
$p_{T,asso} \in (1.0, 2.0)$ GeV/c	Head	0.403106	0.0163408	0.00270691
$p_{T,asso} \in (1.0, 2.0)$ GeV/c	Shoulder	0.135421	0.00940179	0.000890293
$p_{T,asso} \in (2.0, 4.0)$ GeV/c	Near-side	0.0681146	0.00652803	0.000581799
$p_{T,asso} \in (2.0, 4.0)$ GeV/c	Head	0.164206	0.00992462	0.000868219
$p_{T,asso} \in (2.0, 4.0)$ GeV/c	Shoulder	0.0147791	0.00321696	0.000280081

Table 2.3: Yields and Errors from NPE-h correlations in p+p collisions with trigger 6.0GeV/c $\leq p_t \leq 9.0$ GeV/c.

B and D to find the ratio as a function of electron p_T . Several experiments have performed this measurement, making this a reasonable check that our method for constructing NPE-h correlations is working as intended.

The simulations are done with Pythia 8.2 with the standard STAR heavy flavor tune. To get reasonable samples across all electron p_T and to avoid the low p_T divergence in heavy flavor processes in Pythia we generate the correlations in several bins in `ptHat` and then piece them together according to a weighting, this closely follows the Pythia example `main08.cc`. For $\text{ptHat} \leq 3.0$ GeV/c we use the “minbias” process `SoftQCD:nonDiffractive = on` for the bins above this we use the hard QCD processes `HardQCD:all = on`. Then to patch the different bins together we scale each bin by the generated cross section, the raw and scaled `ptHat` spectra can be seen in Figure 2.13.

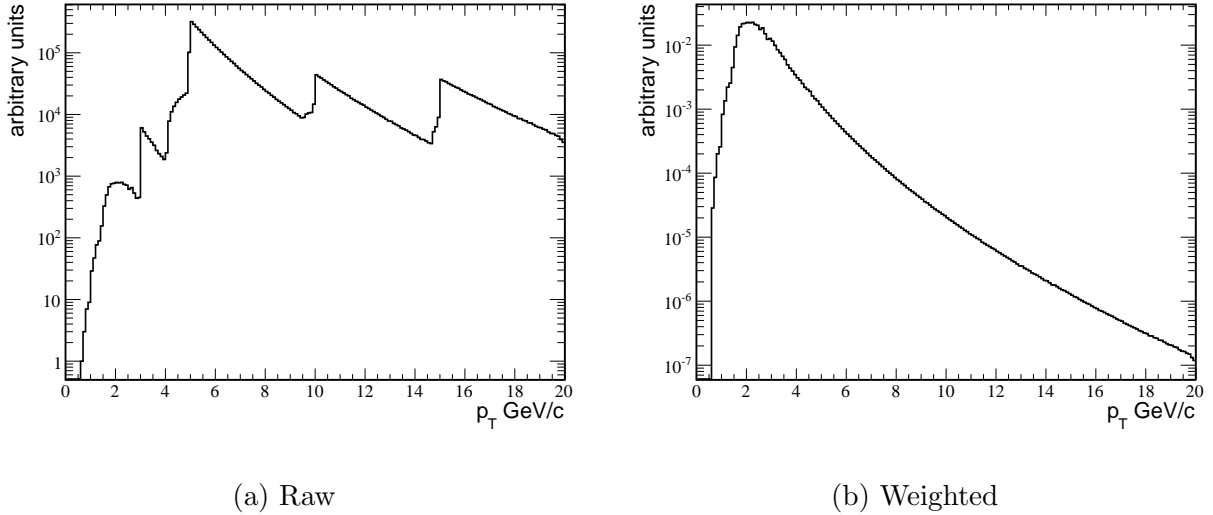


Figure 2.13: The raw and weighted `ptHat` spectra from Pythia.

Then we construct the electron-hadron correlations from the Pythia simulations in two separate cases one for initial b quarks and one for initial c quarks. We check that the final state contains an electron and that the electron has a parent B or D meson. To increase statistics, if there is no final state electron we pick on of the heavy mesons, undo and then redo the decay until we get an electron effectively setting the branching ratio for semileptonic decays to 100%. Then we apply acceptance cuts that closely match the STAR acceptance:

$\eta_e \in (-.7, .7)$, $p_{T,e} > 2.0$ GeV/c, $\eta_h \in (1.0, 1.0)$, and $p_{T,h} > 0.2$ GeV/c. This gives us the electron-hadron correlation that we use in the fit.

For the fit we compare the correlations for $p_{T,h} > .2$ GeV/c and vary the p_T of the trigger particle. We fit the correlation with the function:

$$\frac{dN_{NPE-h}}{d\Delta\phi} = r_B f_B(\Delta\phi) + (1 - r_B) f_D(\Delta\phi) \quad (2.3)$$

Where f_B and f_D are the correlations for bottom and charm electron-hadron correlations from Pythia, and $r_B = \frac{e_B}{e_B + e_D}$. The fit is done in a range around the near side peak $\Delta\phi < 1.5$ due to difficulties in Pythia accurately recreating away side behavior. Figure 2.14 shows an example fit, plus the correlations from Pythia.

Figure 2.15 shows the results for obtaining r_B from fits of the p+p correlation plus a comparison to previous published results. Errors are large at high trigger p_T due to the distributions from B and D becoming similar thus reducing the ability to distinguish between the contributions to the overall p+p correlation shape. Discrepancies in the mid p_T range may be due to difficulties in patching together the `ptHat` bins in the Pythia simulations and this could be fixed by only using `SoftQCD:nonDiffractive` to generate the simulated correlations, but this would be far more time consuming.

2.6 Comparisons of Yields

2.6.1 Away Side Shape

2.6.2 I_{AA}

2.7 Event-Plane Dependent Correlations

The motivation for measurements of two particle correlations in Au+Au comes from the fact that partons are expected to strongly interact with the color charges present in QGP and that the jet like correlations will depend on the path length traversed by the parton in QGP.

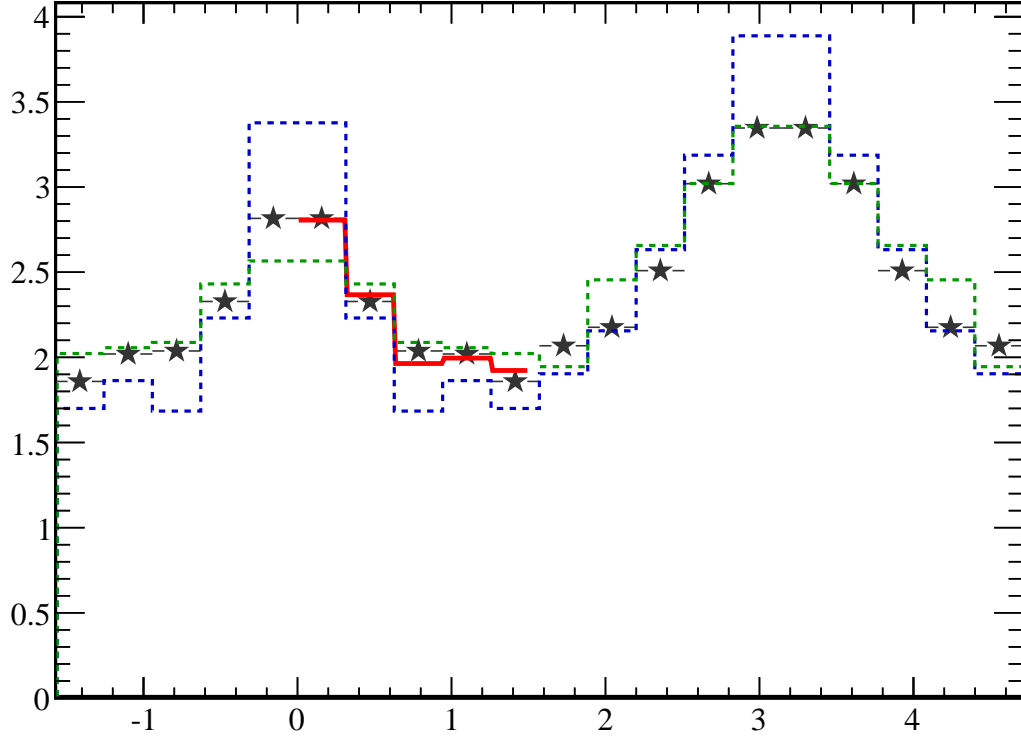


Figure 2.14: NPE-h correlation with $5.0 \text{ GeV}/c < p_{T,e} < 6.0 \text{ GeV}/c$ and $p_{T,h} > 0.2 \text{ GeV}/c$. Pythia distributions are $f_D(\Delta\phi)$ (blue), $f_B(\Delta\phi)$ (green), and the distribution with best fit value of r_B (red).

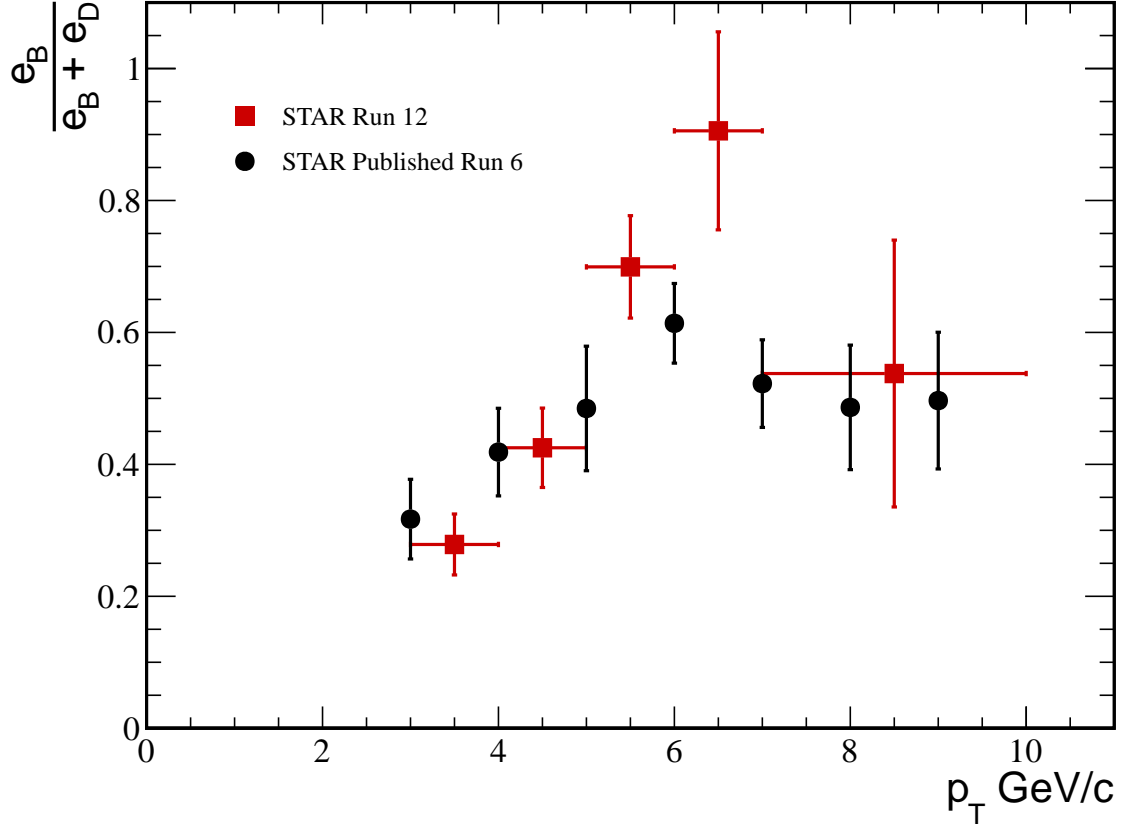


Figure 2.15: Measured value of r_B as a function of trigger p_T . Red points are for the current analysis and the black are from STAR published in Ref. Only statistical errors are shown.

Previously we have looked at comparisons between central and peripheral events as well as p+p collisions as a way of investigating the dependence on the presence of QGP. Now we will try to look at a more direct dependence on path length by measuring the correlations relative to emission in or out of the event plane.

In non-central Au+Au collisions the collision region formed from the overlap of the two incident nuclei is ellipsoidal. This initial anisotropy in the medium results in elliptic flow, v_2 , and in the final state we observe anisotropic emissions relative to the reaction plane, defined by the momenta of the initial beam particles (see Figure 2.16).

Elliptic flow is defined relative to the reaction plane, however there is no way to experimentally determine the reaction plane in a given event. Instead we use the azimuthal distribution of observed particles to estimate the reaction plane. We get an angle for the estimated reaction plane which we call the event plane. Due to the fact that we calculate the event plane from a finite number of particles we also have some resolution of our event plane angle.

The event plane angle gives us an estimate for the reaction plane angle, we will then look at our trigger particle's angle relative to the event plane. As seen in Figure 2.16 particles traversing the medium perpendicular to the event plane should move through a larger part of the medium. We might guess that out-of-plane correlations show larger away side suppression compared to in-plane, which has already been observed in dihadron correlations. A large portion of our high tower trigger data comes from mid-centralities (30-60% central) which is useful for this analysis since that is where we have the best event plane resolution, but due to higher hadron v_2 this region also has larger systematic uncertainties which may make drawing conclusions difficult.

2.7.1 Event Plane Reconstruction

We use the azimuthal distribution of the particles in an event to calculate the event plane, Ψ_{EP} . Event planes can be calculated for any order of harmonic in the decomposition of the particle distribution, but since we are interested in v_2 and the second order event plane

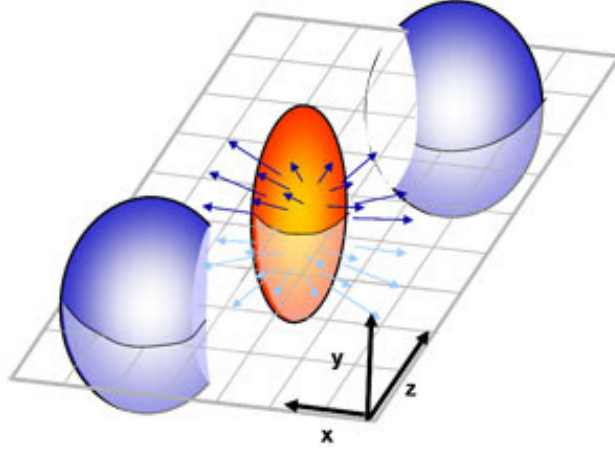


Figure 2.16: Illustration of non-central heavy ion collision and the resulting v_2 . The reaction plane corresponds to the xz -plane in the diagram.

everything that follows is specifically the $n = 2$ case. For a more general treatment see ref. First we apply single particle ϕ -weighting to correct for the acceptance of the detector. This is the same procedure as used for the two-particle correlation.

We use the hadrons in the event with $0.2 \text{ GeV}/c \leq p_T \leq 2.0 \text{ GeV}/c$ in the event plane calculation. Since these events are high tower triggered events there is also a high p_T leading particle in each event. Since the presence of a jet may bias the calculation of the event plane, we should try to remove this contribution. To do this we exclude the particles in the event with $|\Delta\eta| < .5$, where $\Delta\eta$ is the pseudorapidity difference between the hadron and the leading particle in the event (in events used in NPE analysis this track is usually the electron). This procedure is called the Modified Reaction Plane (MRP) method.

With our sample of hadrons we then construct a flow vector \mathbf{Q} for the event. The components of this vector are:

$$X = \sum_i w_i \cos(2\phi_i) \quad (2.4)$$

$$Y = \sum_i w_i \sin(2\phi_i) \quad (2.5)$$

where w_i is some weight given to the particle, in this analysis we use the particle's p_T as the weight. The second order event plane angle is then given by:

$$\Psi_{EP} = \frac{1}{2} \tan^{-1}\left(\frac{Y}{X}\right) \quad (2.6)$$

The true distribution of the reaction planes should be uniform across all events, however there will still be some distortion in the event plane distribution which we need to correct. To do this we use a shifting method where the harmonics of the uncorrected distribution are used to flatten it. Figure ?? shows the distribution of the event plane angle after applying the shifting correction. We see that the resulting distribution is flat to within 1%.

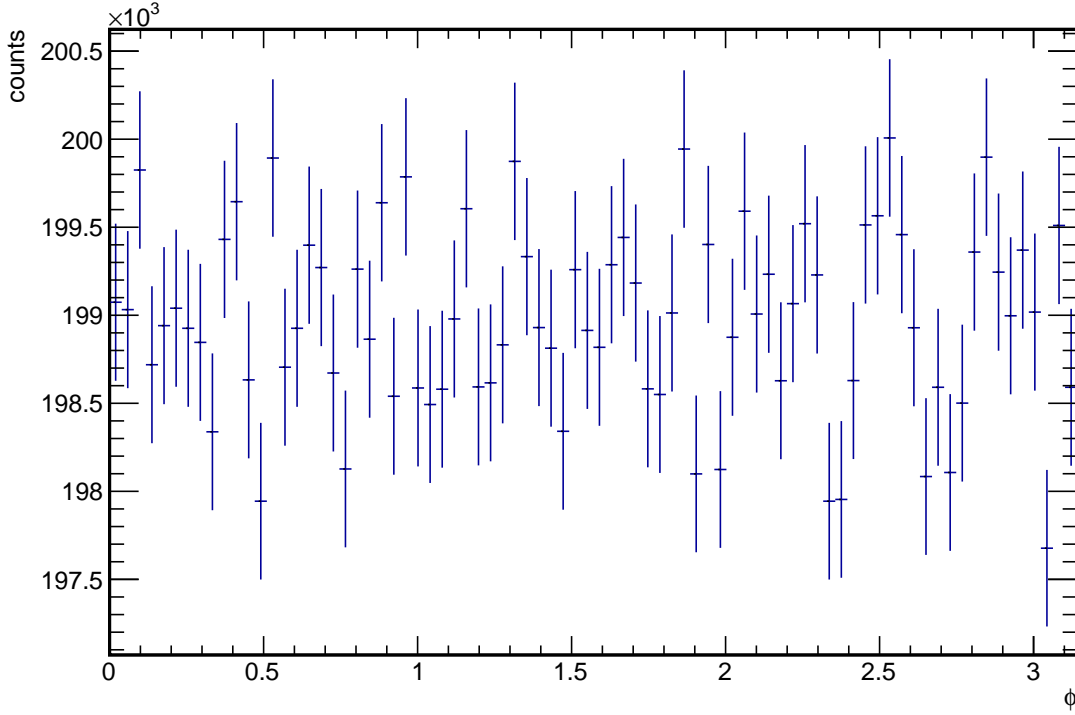


Figure 2.17: The ϕ distribution of the event plane angle Ψ_{EP} after applying the shifting correction.

When we calculate the event plane dependent NPE-h correlations we will potentially be correlating the trigger electron with some of the particles used in the event plane calculation.

We would like the calculated event plane to be independent of the particles in the correlation so we will actually calculate 5 separate event planes per event. We calculate one for all hadrons in the range $.2 \text{ GeV}/c \leq p_T \leq 2.0 \text{ GeV}/c$ then we calculate one for each case where the ranges $.2 \text{ GeV}/c \leq p_T \leq .5 \text{ GeV}/c$, $.5 \text{ GeV}/c \leq p_T \leq 1.0 \text{ GeV}/c$, $1.0 \text{ GeV}/c \leq p_T \leq 1.5 \text{ GeV}/c$, and $1.5 \text{ GeV}/c \leq p_T \leq 2.0 \text{ GeV}/c$ are excluded. Each bin with an excluded p_T range will have a lower resolution due to fewer particles used in the event plane calculation.

Using a finite number of particles to calculate the event plane leaves us with an event plane resolution by which we must scale up our azimuthal anisotropy measurements to get the anisotropy relative to the true reaction plane. This is done to get the correct value of v_2 in flow analyses but we will also need the event plane resolution for calculating the event plane dependent v_2 background for electron-hadron correlations. We calculate the resolution by subevent planes. We divide each event randomly into two equally sized sub events then independently calculate the event plane for each. The differences between the subevent planes can be used to measure the event plane resolution. Figure 2.18 summarizes the resolution as a function of event centrality as well for different hadron p_T slices. Resolution is best around 30% centrality and falls off in central events due to lower azimuthal anisotropy and in peripheral events because of lower multiplicity.

2.7.2 Correlations

Now we would like to calculate the NPE-h correlation and look at the dependence based on how the trigger particle is oriented relative to the event plane. Since for trigger particles which are out of the event plane the heavy quark traversed a longer path in medium, we might wonder if this could effect the correlation.

The correlation is constructed as previously described for Au+Au collisions with one important difference. Before the background from flow assumed no dependence of the trigger particle on the reaction plane. The form of this background resulted from the dependence on the reaction plane cancelling leaving only the difference $\Delta\phi$ between the correlated particles. Now we are choosing specific orientations of the trigger particle relative to the event plane

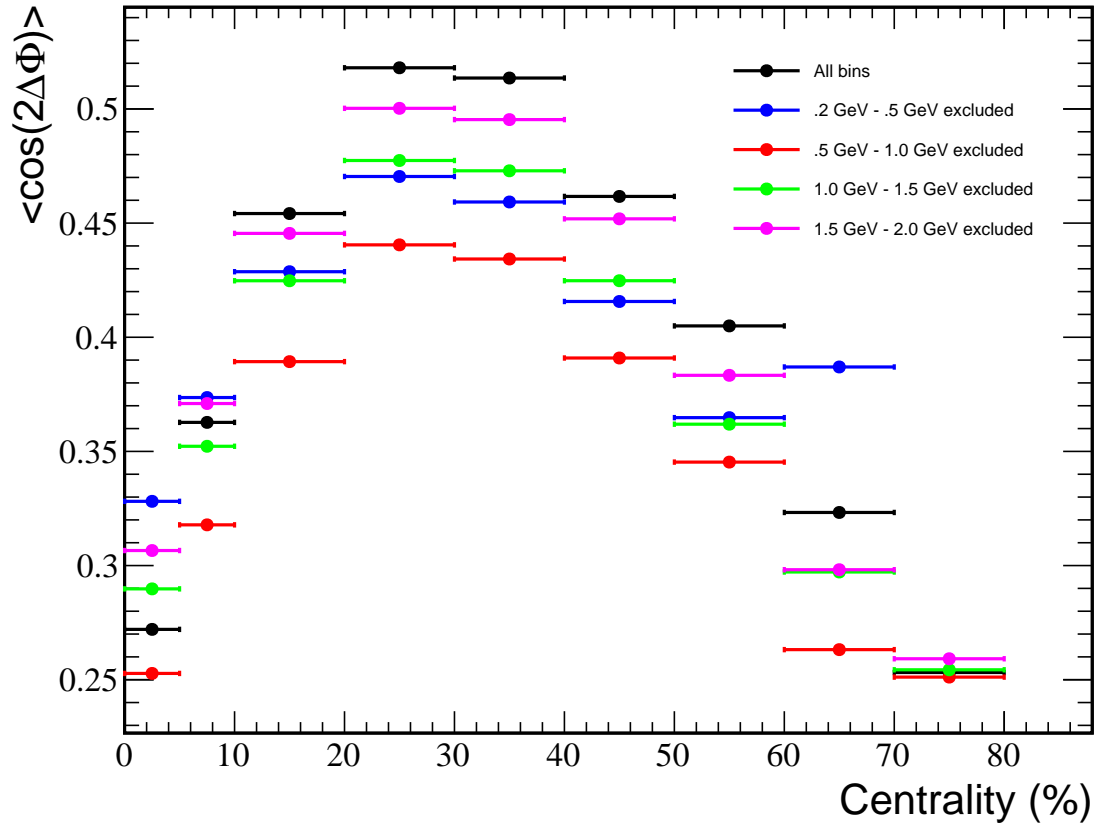


Figure 2.18: Event plane resolution as a function of centrality.

and thus the background will be different and depend on that orientation.

We use the same functional form for the background, Equation 2.1, as when we originally calculated the NPE-h correlations in Au+Au, but now we replace the v_2 of the electron with an ‘effective’ v_2 , denoted \tilde{v}_2 , which will depend on the electrons angle relative to the event plane as well as the event plane resolution. The effective v_2 can be calculated for any arbitrary slice relative to the event plane, but we will only consider the case of two equal slices: in-plane and out-of-plane. For this case \tilde{v}_2 is given by:

$$0 \leq |\phi_{trig} - \Psi_{EP}| \leq \frac{\pi}{4} : \tilde{v}_2 = \frac{\pi v_2 + 2\langle \cos(2\Delta\Psi) \rangle}{\pi + 4v_2\langle \cos(2\Delta\Psi) \rangle} \quad (2.7)$$

$$\frac{\pi}{4} \leq |\phi_{trig} - \Psi_{EP}| \leq \frac{\pi}{2} : \tilde{v}_2 = \frac{\pi v_2 - 2\langle \cos(2\Delta\Psi) \rangle}{\pi - 4v_2\langle \cos(2\Delta\Psi) \rangle} \quad (2.8)$$

where Equation 2.7 is for in-plane triggers and Equation 2.8 is for out-of-plane triggers.