

*Chandler*

PLANETARY EPHEMERIS PROGRAM

Contract # NOOO14-89-C-2432

Final Report  
For the period 29 September 1989 through 28 October 1990

Principal Investigator

John F. Chandler

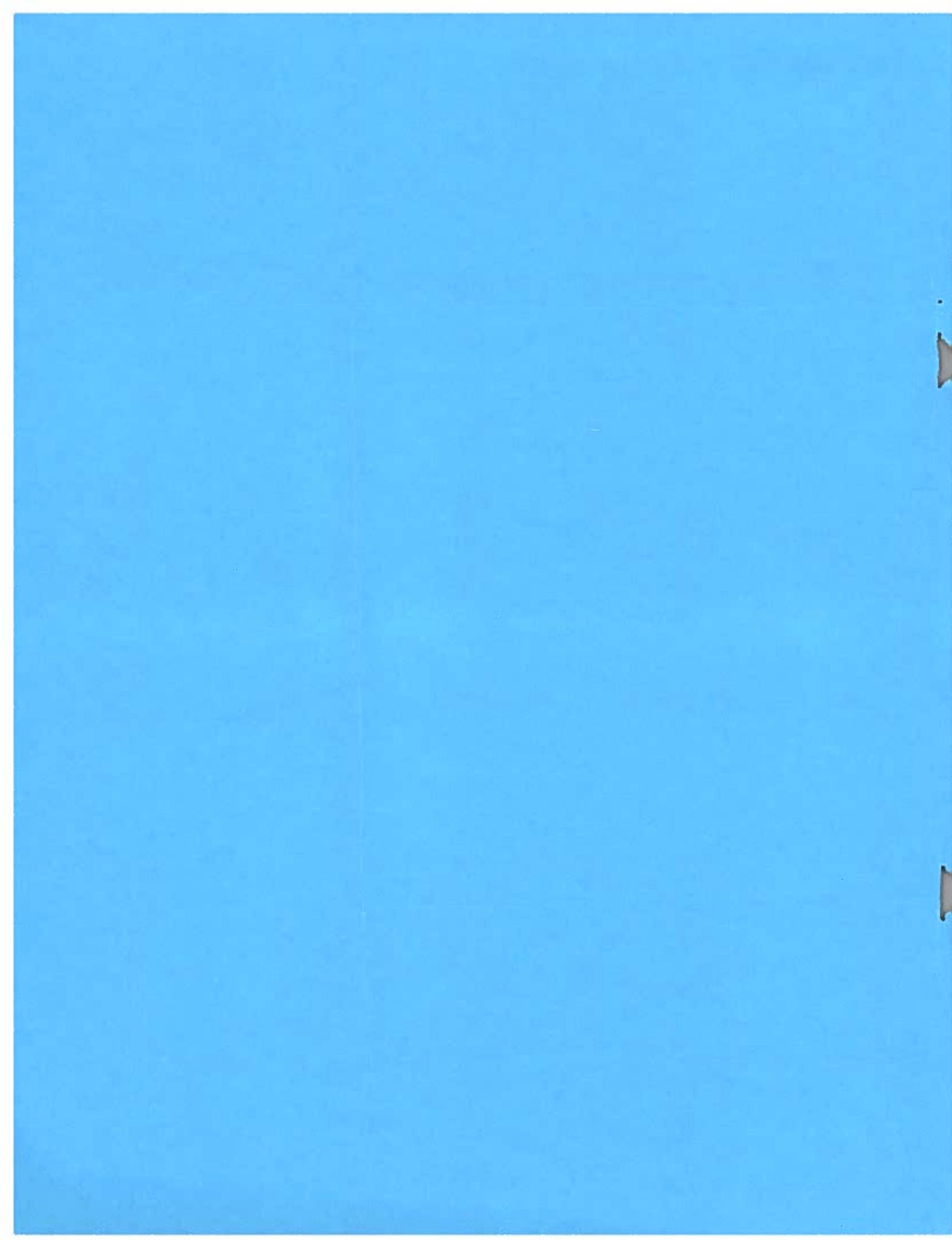
November 1990

Prepared for  
Naval Research Laboratory  
4555 Overlook Avenue, SW  
Washington, DC 20375-5000

Smithsonian Institution  
Astrophysical Observatory  
Cambridge, MA 02138

The Smithsonian Astrophysical Observatory  
is a member of the  
Harvard-Smithsonian Center for Astrophysics

The United States Naval Observatory Technical Officer for this grant is Dr. P. Kenneth Seidelmann, Nautical Almanac Office, 34th Street & Mass. Avenue, NW., Washington, DC 20390.



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Principal Investigator

BODFN

John F. Chandler

SBFN1

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# Harvard-Smithsonian Center for Astrophysics

60 Garden Street, Cambridge, MA 02138

(617) 495-7000  
Telex # 921428



1990 September 26

Dr. P. Kenneth Seidelmann  
United States Naval Observatory  
Nautical Almanac Office  
34th St. & Mass. Ave., NW.  
Washington, DC 20390

Dear Ken,

This letter and its enclosure constitute the report on our activities under contract N00014-89-C-2432 (effective dates 1989 September 29 through 1990 October 28). The focus for this year has been the numerical integration code in PEP, and the enclosure includes documentation of both the logic flow and underlying physical models in the principal subroutines of the n-body and planet integrators. The models as documented are purely descriptive, *i.e.*, they embody what the PEP code actually does. Not surprisingly, we have verified that the documentation included in this report agrees with standard physics. However, a small surprise turned up in the preliminary comparison: a bug in a section of code that had not been used. Partly as an exercise for Mark Murison and partly as the first step in the code upgrade, this bug was fixed.

This report illustrates the basic properties of the eventual "comprehensive" PEP documentation that will result from our project. There are varying degrees of detail in the descriptions, ranging from thorough exposition of the physics to simple display of the subroutine calling structure. These correspond to the relative importance of the subroutines and to the amount of physics they contain.

Not yet available, but expected soon, is a report from Ken Nordtvedt describing his development of a more general framework for tests of general relativity. Bob Reasenberg spoke to Ken at the Fairbank meeting in Rome. Ken has apparently been making good progress, and Bob and I look forward to the completion of his work, which will be the basis for the most significant changes we expect to make.

For future work, we have several more pieces of the integration code to document (for example, the subroutines responsible for the post-Newtonian terms in the equations of motion), and the sections of code dealing with observations. Our goal is to compare the resulting documentation with Ken Nordtvedt's framework at post-Newtonian order and thereby reorganize PEP to include a more general model and, at the same time, to be more maintainable.

Yours truly,

A handwritten signature in cursive script that appears to read 'John'.

John F. Chandler

enc.  
JFC/sas

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## Subroutine Charts

## Preface

## Preface

The following documentation of PEP, the Planetary Ephemeris Program, represents an in-progress snapshot of current work. We have decided that firm understanding of the *physics* embedded in PEP is the highest priority. Later we will concentrate on other goals, including incorporation of a new theoretical parameterized post-Newtonian framework being devised by K.H. Nordtvedt, documentation of the observables sections of PEP, and documentation aimed at the user of PEP.

The first routines subjected to detailed analysis are those involved in the integration of the equations of motion and of the partial derivatives. For each routine, the associated documentation will consist, as needed, of 1) a summary sheet, describing the routine at a glance; 2) a memorandum explaining in detail the physics in the routine; 3) a subroutine call tree; 4) a process structure chart; and 5) a detailed flowchart. An explanation of each of these follows.

The summary sheet for subroutine X contains the call syntax, followed by a brief description of the purpose of the routine. Then each calling argument is explained in enough detail to allow one to use X without having to look at the actual code. Next, under the heading "Calls:", is listed the subprogram dependency; X requires these subprograms in order to run properly. Following this is a list of all the PEP routines which call X. After that is a list of include files required by X. Finally, if X involves any of the global variables in the include files in a fortran EQUIVALENCE, that global variable is listed, as well as the common block it belongs to and the name of the associated include file.

The primary purpose of the memorandum is to document the equations used in the subroutine. This is especially important for the parts of PEP containing physics (as opposed to service routines, numerical analysis routines, *etc.*). For clarity, the order of the presentation does not necessarily follow the sequence of calculation in the subroutine. In general, derivations of the equations are not shown. In addition, not all of the

routines require a memorandum.

As part of the overall PEP documentation, subroutine call tree charts are being created. If a routine happens to be the top node of a particular tree, the chart is included with the other documents describing the routine (this is in addition to the presence of the chart in the call tree documentation). In the call tree charts, each routine is contained in one of three box styles. These correspond to position in the tree of the routine: 1) internal node with further branches, 2) "leaf" with no further subprogram dependencies, and 3) internal node which is a top node and is displayed on another chart. In addition, within a "leaf" box style there is a separate notation for program entry points. These styles are illustrated at the top of the call tree chart for PEPMAIN. All PEP routines will be included somewhere in the charts. The top-level nodes that begin each chart were chosen such that each chart fits onto approximately one page.

Each routine that has a summary sheet will also have a process structure chart. This chart is easily recognized by the shaded box at the upper left corner, containing the name of the routine. The structure chart presents a hierarchical overview of the processes that make up the subroutine.

Following the structure chart is a detailed flowchart. The purpose of the flowchart is to document the logic of the routine, which is often complicated in a program as large as PEP. The flowcharts are hierarchical. The high-level charts contain symbols which indicate a further expansion on another (lower-level) chart. Chart expansion is indicated by a large circle containing the label of the expanded chart in capital letters, and often also including some indication of the purpose of the section of code represented by the expanded chart. For example, the flowchart for subroutine PLANET contains an expansion symbol labeled PLNT1, with the indication "N-body integration". The next flowchart is labeled PLNT1 and contains the N-body integration logical flow. For very complicated routines (SBFN1 is a good example), a chart is also included which details the flowchart hierarchy. This chart of flowcharts is a "directory" of the flowcharts for that routine.

**PEPMAIN**

## **program PEPMAIN**

Description: This is the driver program for PEP.

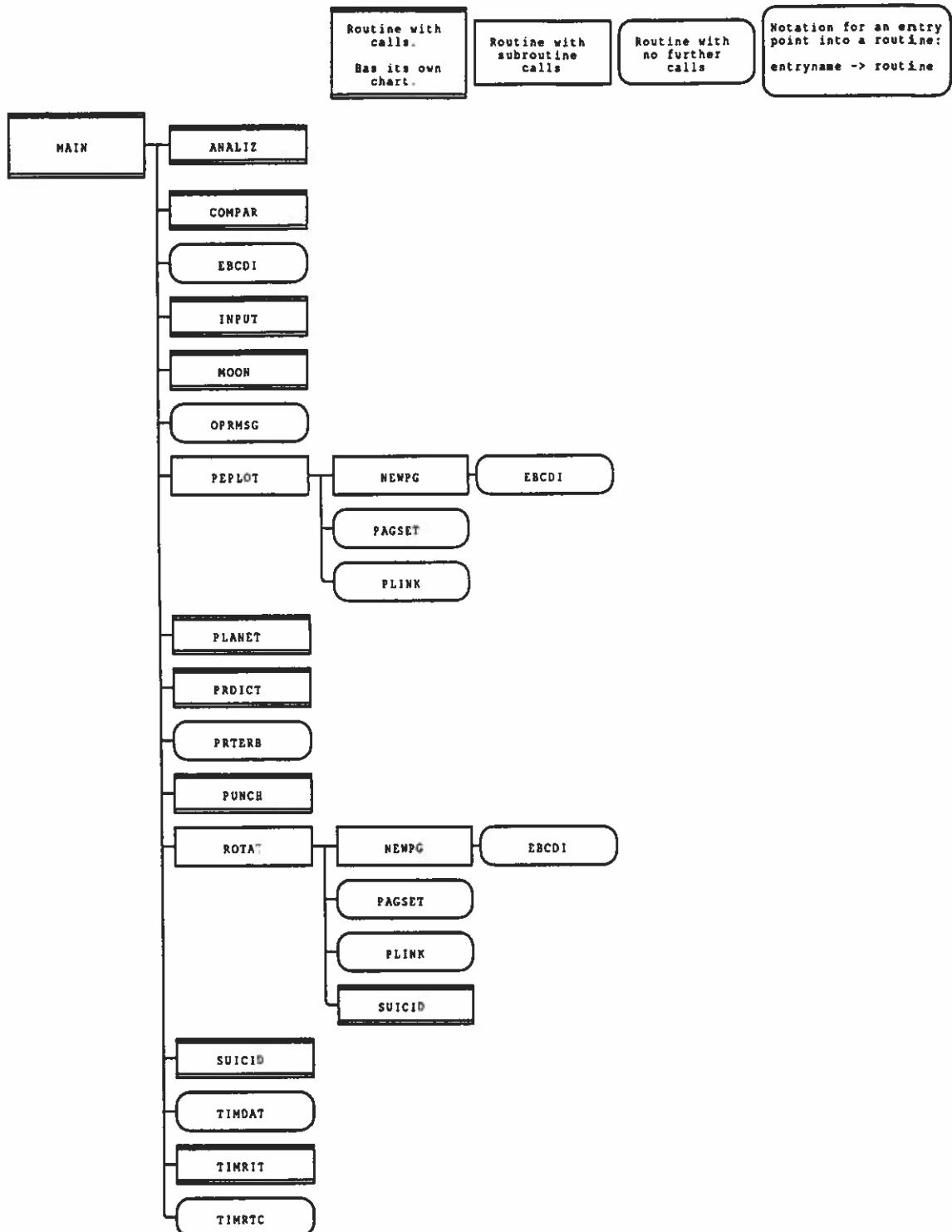
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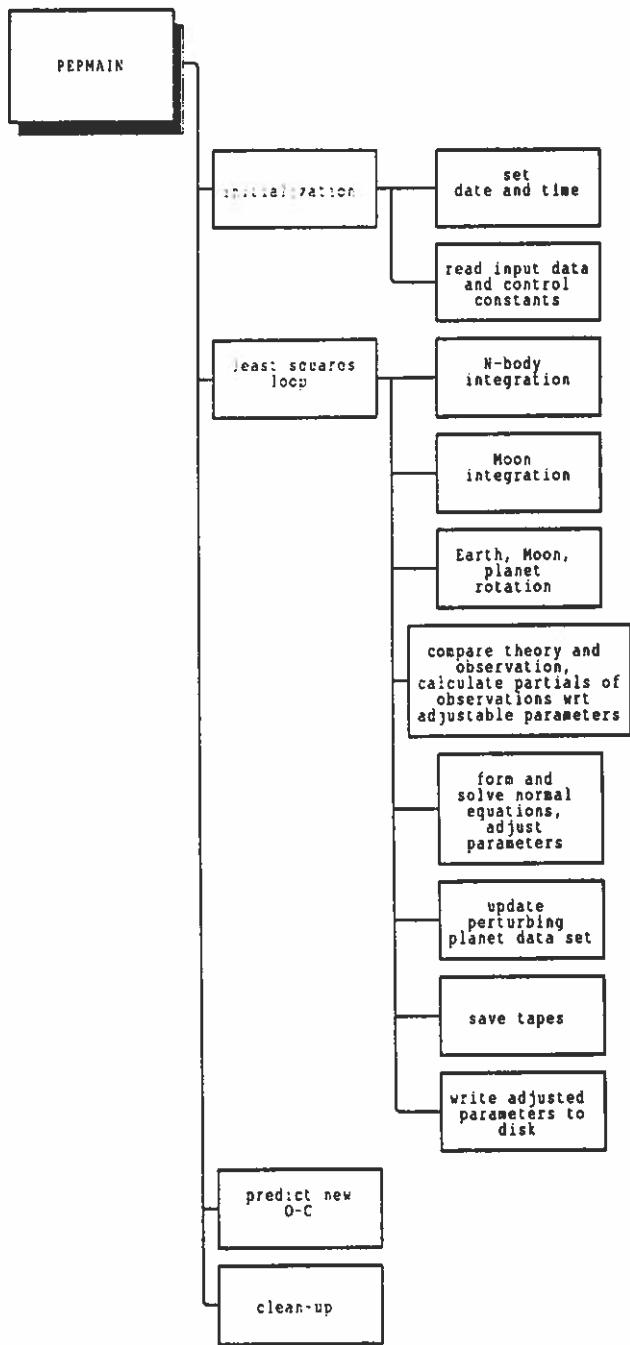
Calls: ANALIZ, COMPAR, EBCDI, INPUT, MOON, OPRMSG, PEPLIT,  
PLANET, PRDICT, PRTERB, PUNCH, ROTAT, SUICID, TIMDAT,  
TIMRIT, TIMRTC

Called by: none

Include Files: BDCTRL, CRDBUF, FCNTRL, INODTA, NAMTIM, TIMSTF

Equivalence: none





**BODFN**

**real\*8 function BODFN( k, j, s )**

Description: BODFN is the function called to evaluate the right-hand sides of the N-body equations of motion. Includes asteroid perturbations, time-variable gravitational constant, 2<sup>nd</sup> harmonic of Sun's gravitational potential, GR motion factors, and ability to treat Earth-Moon system as a single point mass or as two separate bodies for perturbations.

Arguments: k      Equation number. For a particular body i, the state vector is stored in a one-dimensional array y starting at y(m), where m = 6\*(i-1)+1. Thus, for example, k = 15 references the equation for z, and k = 17 references the equation for dy/dt, of body i = 3.

j      Integrator iteration number (generally takes on values from 1 to 3, and is often used as a logic flag).

s      Current time = julian date + 0.5 day.

Calls: BDASTF, BDAST0, PRTCRD

Called by: ADAM, EVAL, NINT, RROAD

Include Files: BDCTRL, BODET, BODSTF, FUNCON, OUTPUT, PARAM, PRTCOD

Equivalence: none

Affected Variables: link      /bodet/      BODET  
                      mp      /bodet/      BODET

Restrictions: none

## BODFN equations:

Upon entry into BODFN, the body currently being integrated will be denoted as body N. Other bodies will be denoted by subscript j, with the exceptions of Earth (subscript E), Moon (subscript M), Earth-Moon barycenter (subscript B), and Sun (subscript S).

Subscripts and their corresponding bodies are

N	Body currently being integrated
S	Sun
E	Earth
M	Moon
B	Earth-moon barycenter
j	bodies other than bodies N, E, M, B

This notation is in keeping with that of the accompanying flowchart.

### 1. Setup

In the initial setup section, first part, compute the velocity of light in astronomical units per day,

$$c_{AU} = \frac{86400}{AU \text{ in light sec}}$$

Define the vector P:

$$\begin{aligned}P_x &= \sin \Omega_0 \sin i_0 \\P_y &= -\cos i_0 \sin \epsilon - \cos \Omega_0 \sin i_0 \cos \epsilon \\P_z &= \cos i_0 \cos \epsilon - \cos \Omega_0 \sin i_0 \sin \epsilon\end{aligned}$$

where

- $\Omega_{\odot}$  Longitude of the ascending node of Sun's equator on the ecliptic, measured from the mean equinox of 1950.0.
- $i_{\odot}$  Inclination of Sun's equator to the ecliptic.
- $\epsilon$  Mean inclination of the ecliptic to the celestial equator of 1950.0 (*i.e.*, the obliquity).

The north pole of the Sun points in the direction of P. The x, y, and z components of P are the cosines of the angles between the Solar north pole and the axes of the usual right-handed orthogonal coordinate system whose X axis is in the direction of the equinox of 1950.0 and whose Z axis is in the direction of the north celestial pole.

## 2. Gravitational Constant

Calculate the gravitational constant according to

$$\gamma(t) = k_G^2 [1 + \alpha_G(t - t_0)]$$

where  $\alpha_G$  is the time variation factor for the gravitational constant, called gmvary in BODFN; t is the current time, called s in BODFN;  $t_0$  is the initial time, called ta in BODFN; and  $k_G$  is the Gaussian gravitational constant, called gauss in BODFN.

In subroutine BODSET (also in SETUP), the Gaussian constant is set to a value of

$$k_G = 0.01720209895 (\text{AU})^{3/2} (\text{ephemeris day})^{-1} (M_{\odot})^{-1/2}$$

## 3. Position Vectors involving Earth and Moon

Position vectors are defined

$$\mathbf{R}_{ik} \equiv \mathbf{R}_k - \mathbf{R}_i$$

The Earth and Moon mass factors are

$$\mu_E \equiv \frac{M_E}{M_E + M_M} \quad \mu_M \equiv \frac{M_M}{M_E + M_M}$$

Loop through all bodies  $j = 1, \dots, \text{nbody}$ . The following relative position vectors may,

depending on the logic flow, be calculated:

$$\begin{aligned} \mathbf{R}_{SE} &= \mathbf{R}_{SB} - M \mathbf{R}_{EM} \\ \mathbf{R}_{SM} &= \mathbf{R}_{SB} + E \mathbf{R}_{EM} \end{aligned}$$

$$\begin{aligned} \mathbf{R}_{Ej} &= \mathbf{R}_j - \mathbf{R}_{SE} \\ \mathbf{R}_{Mj} &= \mathbf{R}_j - \mathbf{R}_{SM} \end{aligned}$$

$$\mathbf{R}_{Bj} = \mathbf{R}_j - \mathbf{R}_{SB}$$

$$\begin{aligned} \mathbf{R}_{jE} &= \mathbf{R}_{SE} - \mathbf{R}_j \\ \mathbf{R}_{jM} &= \mathbf{R}_{SM} - \mathbf{R}_j \end{aligned}$$

$$\begin{aligned} \mathbf{R}_{Ej} &= -\mathbf{R}_{jE} \\ \mathbf{R}_{Mj} &= -\mathbf{R}_{jM} \end{aligned}$$

$$\mathbf{R}_{jB} = \mathbf{R}_{SB} - \mathbf{R}_j$$

$$\mathbf{R}_{Bj} = -\mathbf{R}_{jB}$$

#### 4. Calculate body-body and Sun-body Distances

For each body  $j = 1, \dots, n_{body}$ , excluding the Earth, calculate

Distance from Sun:  $\mathbf{R}_j$  and  $R_j^3$

and, for  $k = 1, \dots, n_{body}$ , again excluding Earth, the position vectors

$$\begin{aligned} \mathbf{R}_{jk} &= \mathbf{R}_k - \mathbf{R}_j \\ \mathbf{R}_{kj} &= -\mathbf{R}_{jk} \end{aligned}$$

#### 5. Accelerations due to Sun and Planets

Acceleration due to Sun:

$$\mathbf{q}_\odot = -\gamma(t) \left( 1 + \frac{M_N}{M_\odot} \right) \frac{\mathbf{R}_N}{R_N^3}$$

Add  $\mathbf{q}_\odot$  to BODFN.

Acceleration with respect to Sun due to planets:

For each body  $j = 1, \dots, \text{nbody}$ , excluding the Earth, sum the vector

$$\mathbf{\Omega}_j = M_j \left( \frac{\mathbf{R}_{Nj}}{R_{Nj}^3} - \frac{\mathbf{R}_j}{R_j^3} \right)$$

Then add  $\gamma(t) \sum_j \mathbf{\Omega}_j$  to BODFN.

Acceleration due to Earth-Moon system:

If Earth and Moon are treated as a single point mass, add

$$\mathbf{q}_{EM} = \gamma(t) M_{EM} \left( \frac{\mathbf{R}_{NB}}{R_{NB}^3} - \frac{\mathbf{R}_B}{R_B^3} \right)$$

to BODFN.

If Earth and Moon are treated as separate masses, add instead the quantity

$$\mathbf{q}_{EM} = \gamma(t) M_{EM} \left[ \mu_E \left( \frac{\mathbf{R}_{NE}}{R_{NE}^3} - \frac{\mathbf{R}_{SE}}{R_{SE}^3} \right) + \mu_M \left( \frac{\mathbf{R}_{NM}}{R_{NM}^3} - \frac{\mathbf{R}_{SM}}{R_{SM}^3} \right) \right]$$

to BODFN.

## 6. Accelerations on Earth and Moon

This section, corresponding to flowchart fragment BODFN4, is executed only if body N

is Earth-Moon ( $npl=3$ ) and Earth and Moon are treated as separate masses ( $kbdy(3)\geq 0$ ).

Acceleration on Earth-Moon due to Sun:

If the gravitational constant does not vary in time, the acceleration vector to be added to BODFN is

$$\mathbf{q} = -\gamma(t) \left( 1 + \frac{M_N}{M_\odot} \right) \left( \mu_E \frac{\mathbf{R}_{SE}}{R_{SE}^3} + \mu_M \frac{\mathbf{R}_{SM}}{R_{SM}^3} \right)$$

where  $M_N = M_E + M_M$ .

Acceleration on Earth-Moon due to planets:

For each body  $j = 1, \dots, nbody$  (excluding  $j=3$ ), add to BODFN the acceleration vector

$$\mathbf{\Omega}_j = -\gamma(t) M_j \left( \mu_E \frac{\mathbf{R}_{Ej}}{R_{Ej}^3} + \mu_M \frac{\mathbf{R}_{Mj}}{R_{Mj}^3} - \frac{\mathbf{R}_j}{R_j^3} \right)$$

## 7. "Extra" Forces

The "extra" forces currently consist of the effects due to general relativity and the Sun's gravitational second harmonic.

If the GR motion factors are to be included in the equations of motion ( $kbdy(21)\geq 0$ ), define the quantities

$$\alpha = \mathbf{V}_N \cdot \mathbf{V}_N \quad \text{and} \quad \beta = \mathbf{R}_N \cdot \mathbf{V}_N$$

Also, let  $\zeta_N$  denote the general relativity motion factor for body N (called `relftb` in BODFN). Add to BODFN the vector

$$\mathbf{q}_{GR} = \zeta_N \frac{\gamma(t)}{c_{AU}^2 R_N^3} \left[ \mathbf{R}_N \left( \frac{4\gamma(t)}{R_N^3} - \alpha \right) + 4V_N \beta \right]$$

Let  $J_2$  be the second harmonic of the gravitational potential of the Sun. If this is to be included in the equations of motion ( $\text{kbdy}(23) \geq 0$ ), then add to **BODFN** the vector

$$\mathbf{h}_{N\odot} = \mathbf{p}_N \left( \frac{\mathbf{R}_N}{R_N} (5g_N^2 - 1) - 2g_N \mathbf{P} \right)$$

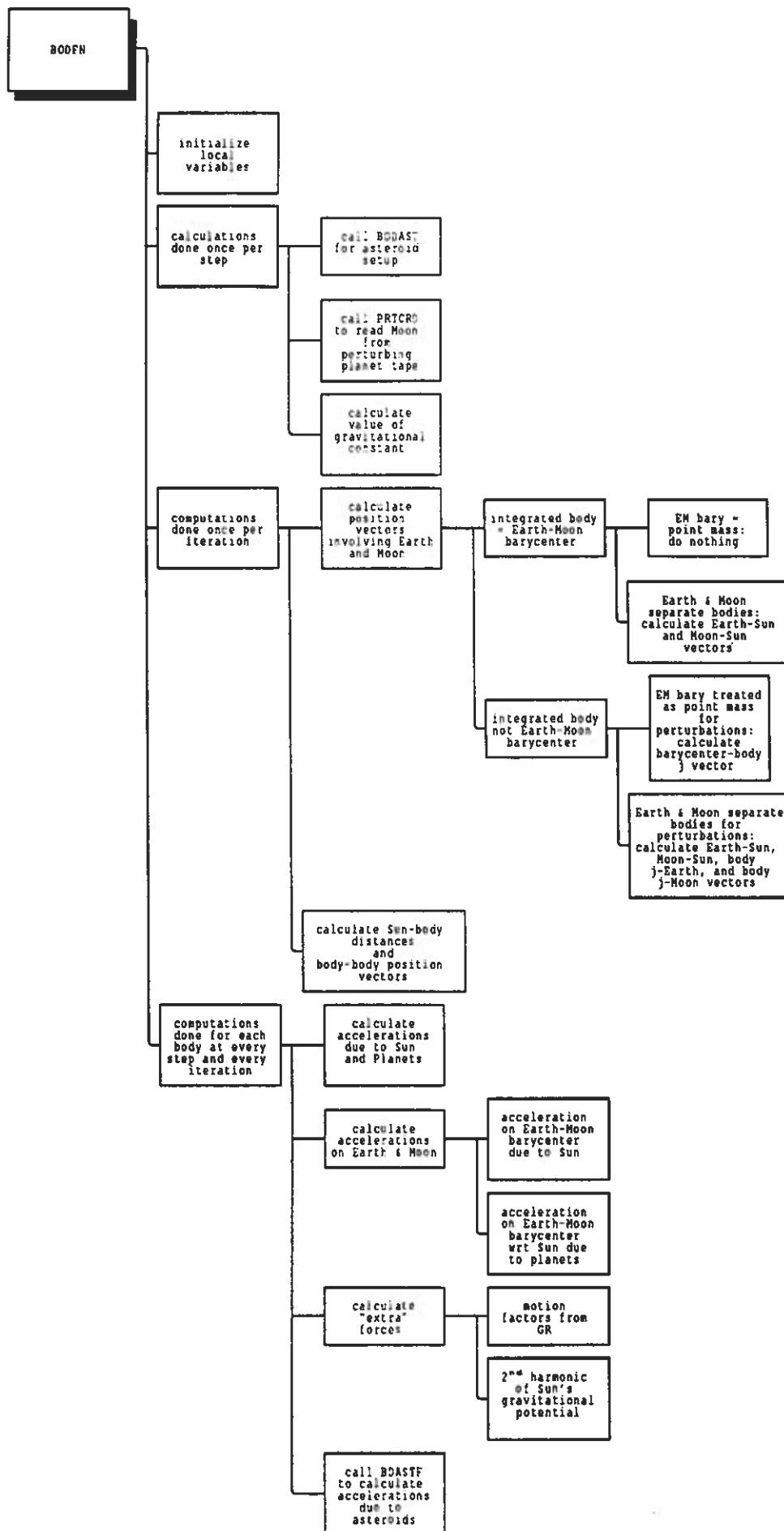
where

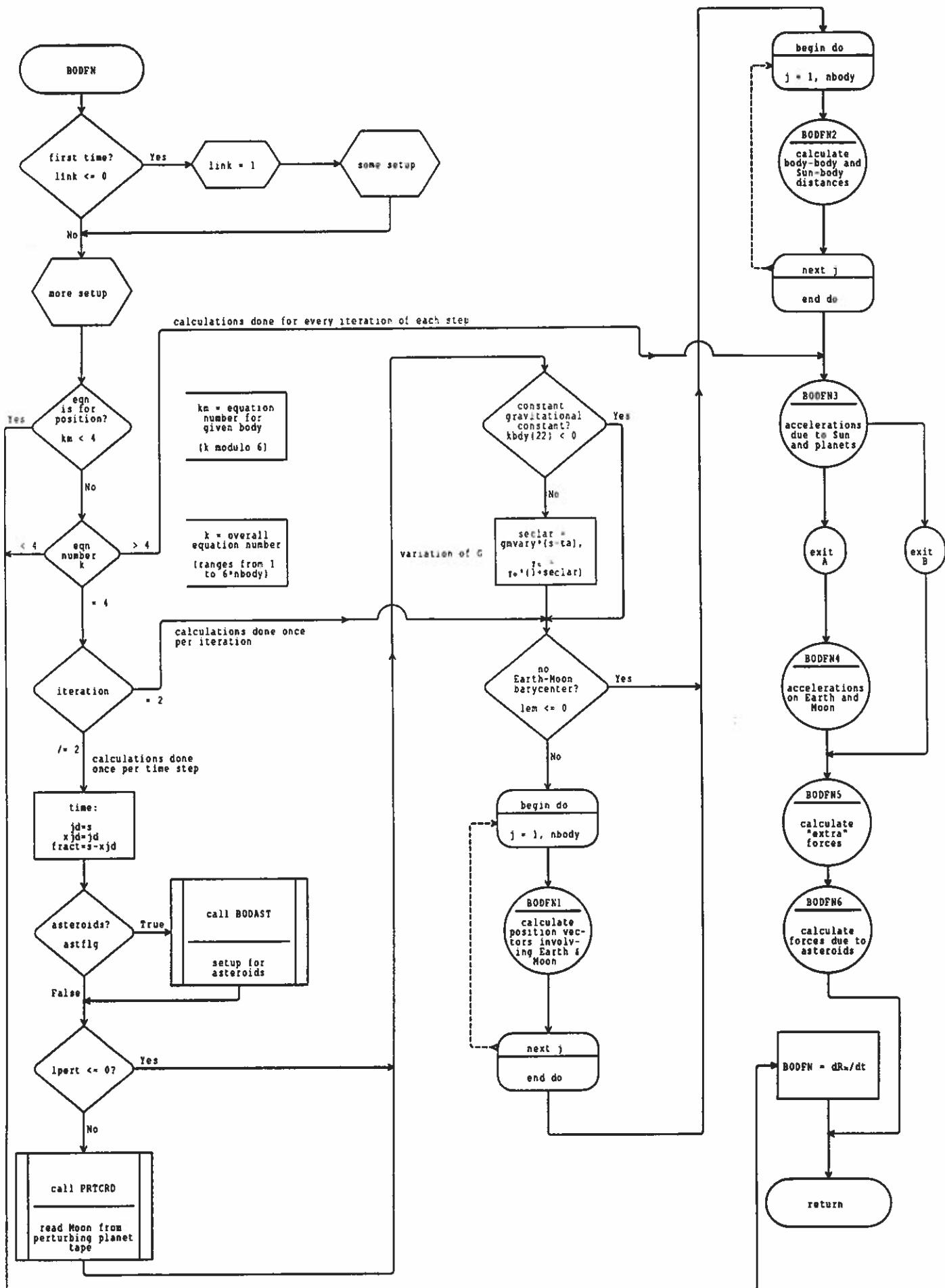
$$g_N \equiv \mathbf{P} \cdot \mathbf{R}_N$$

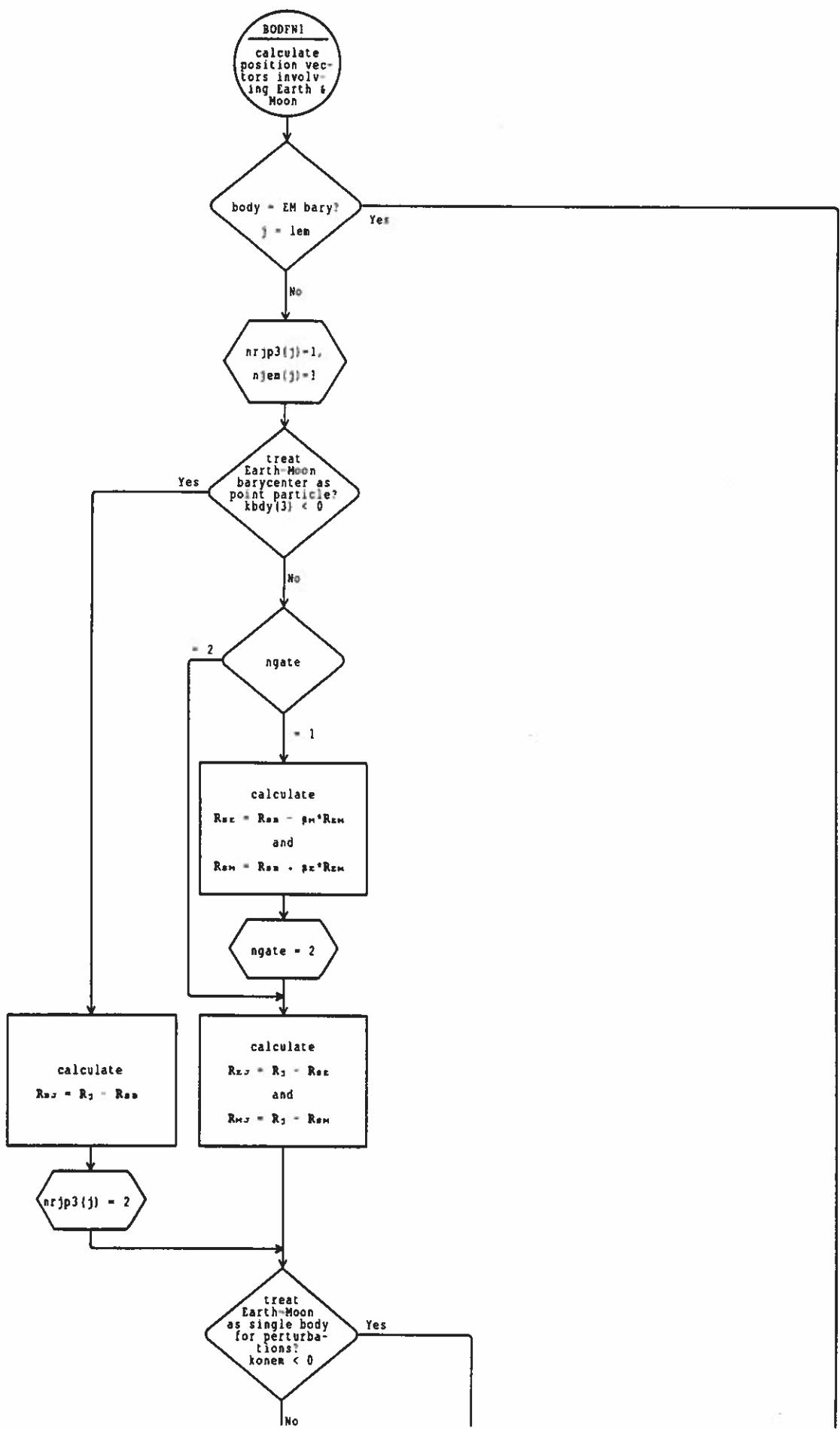
$$p_N \equiv 1.5 \left( 1 + \frac{M_N}{M_\odot} \right) R_\odot^2 J_2$$

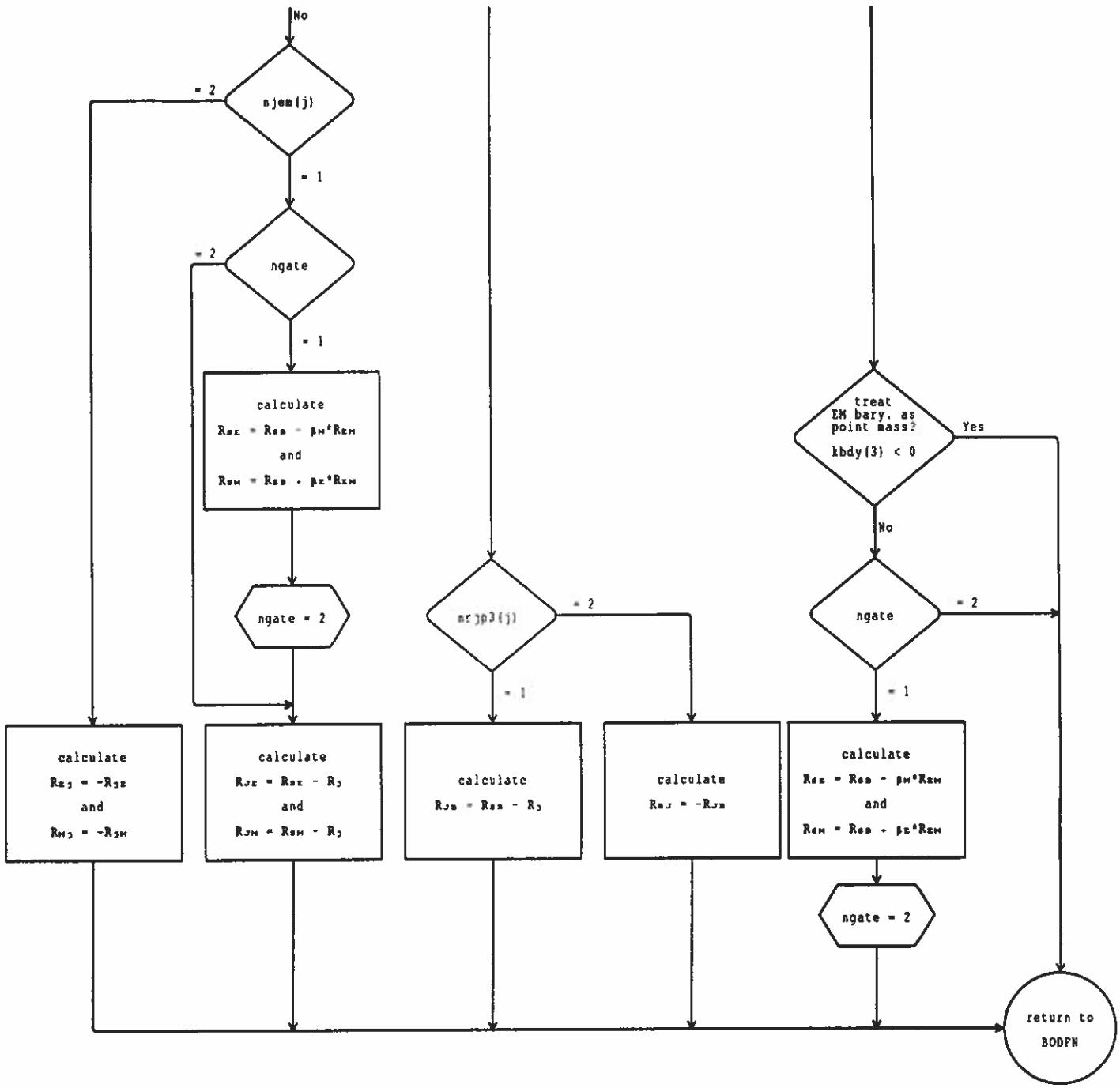
#### 8. Accelerations due to Asteroids

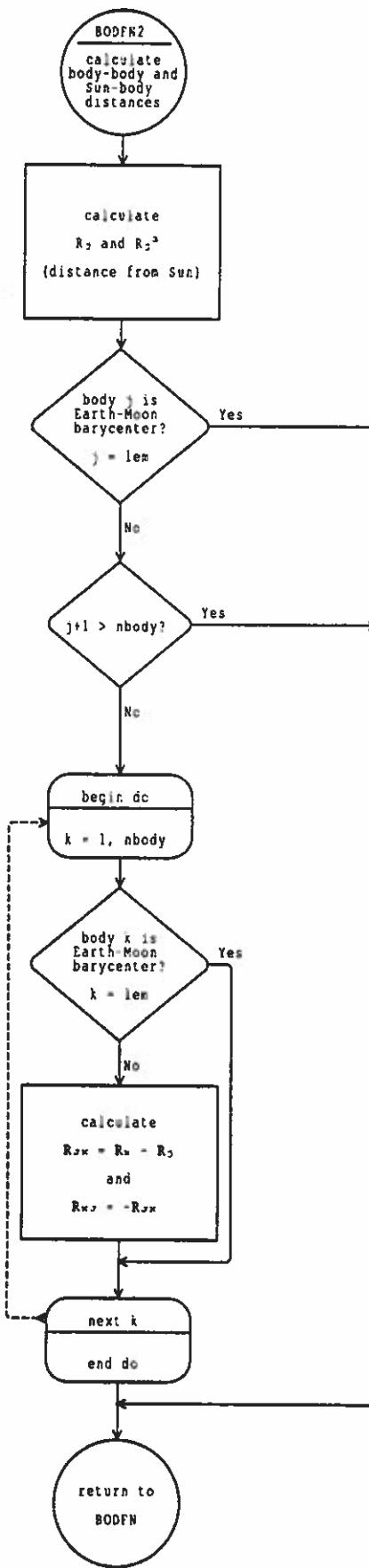
Finally, if asteroids are to be included (`astflg=true`), the **BDASTF** routine is called, returning the components of acceleration due to the asteroids. Call this acceleration  $\mathbf{q}_{\text{ast}}$ . Then  $\gamma(t)\mathbf{q}_{\text{ast}}$  is added to **BODFN**.

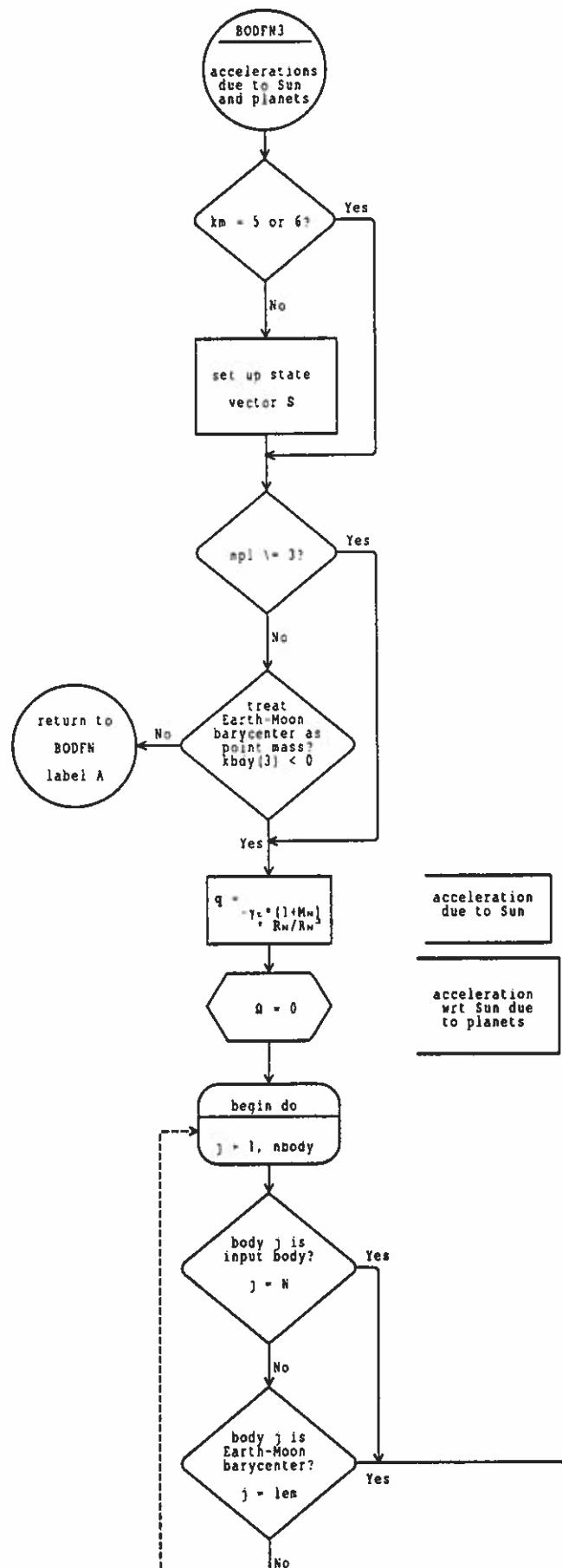


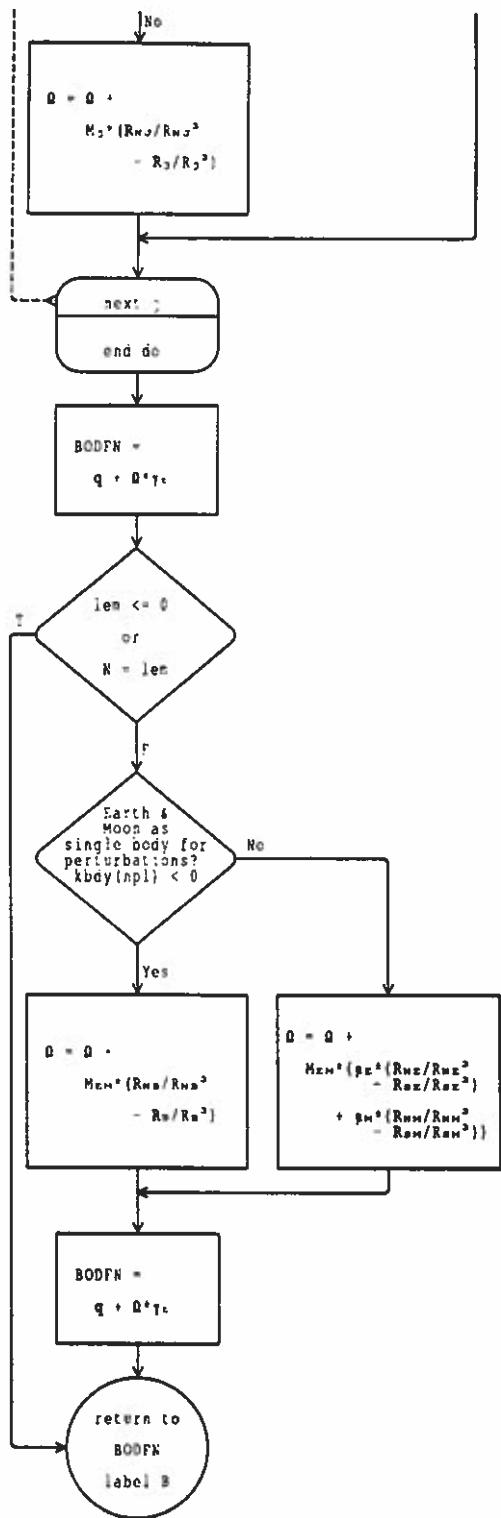


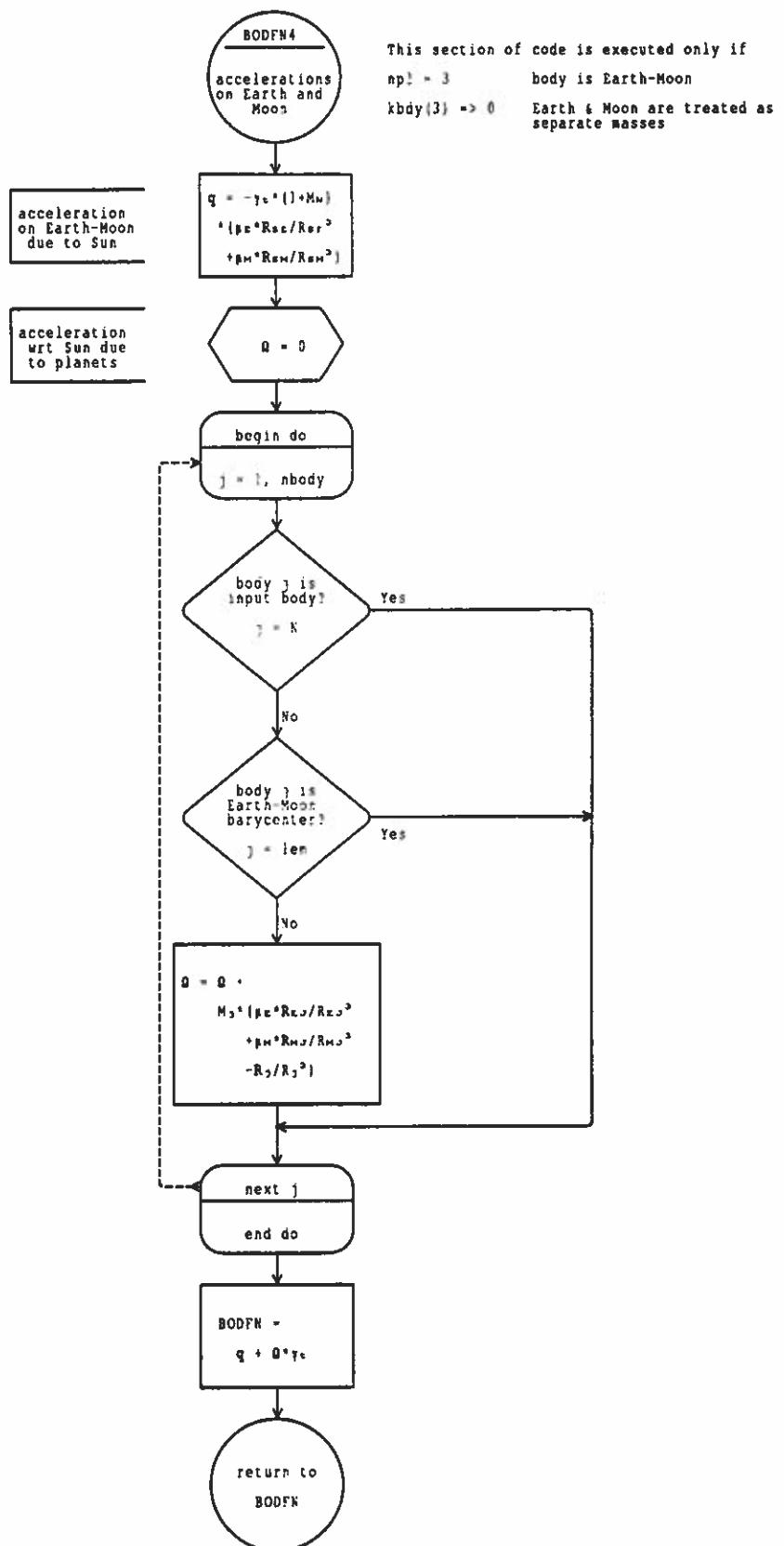


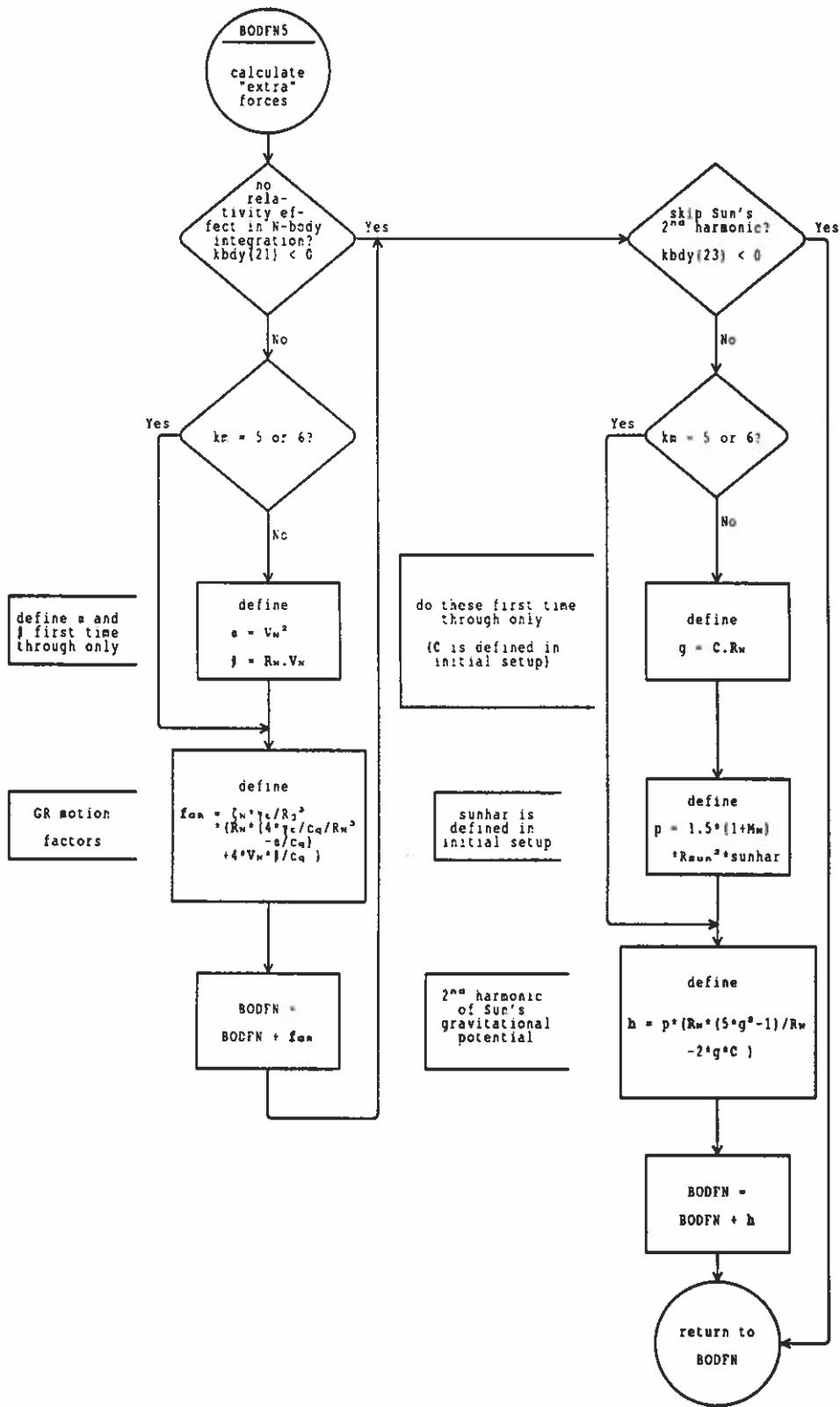


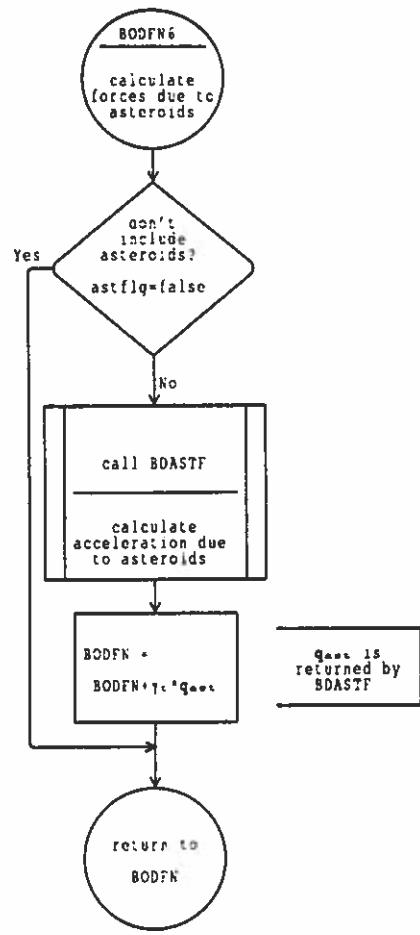












**LUNORB**

**subroutine LUNORB ( jd, fract, nn )**

Description: LUNORB calculates the position, velocity, and acceleration of the mean lunar orbit, using Brown's theory. LUNORB also calculates the partial derivatives of these quantities with respect to the initial mean orbital elements.

Arguments:

jd	Julian Day Number
fract	fraction of day from midnight
nn	control flag: 0 calculate position, velocity, acceleration +1 also calculate derivatives of position, velocity, acceleration -1 also calculate derivatives of position and velocity

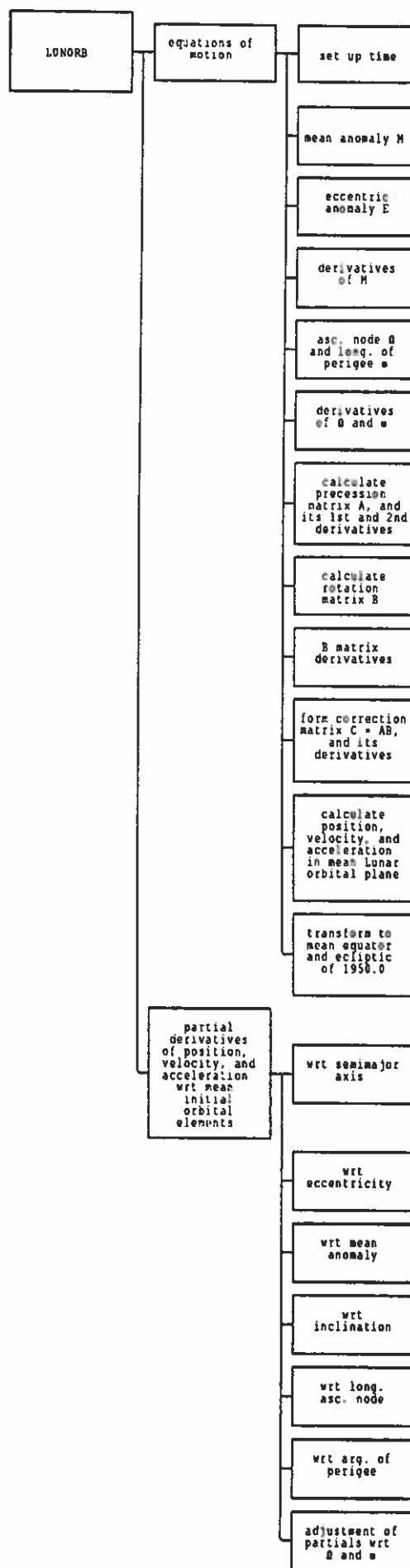
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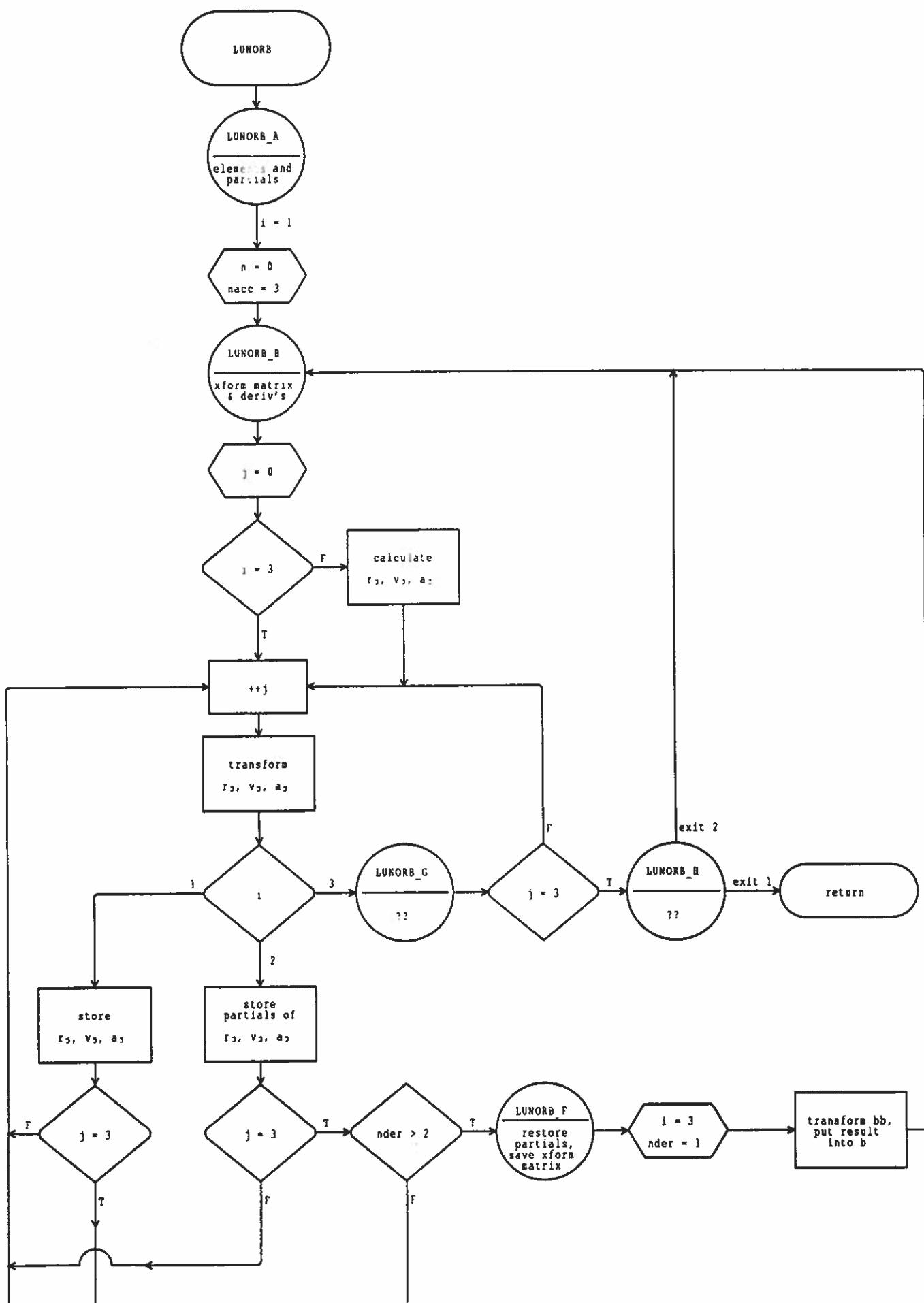
Called by: LUNSET, MORFN, MOROUT, MORSET, PRTC RD, SBFN

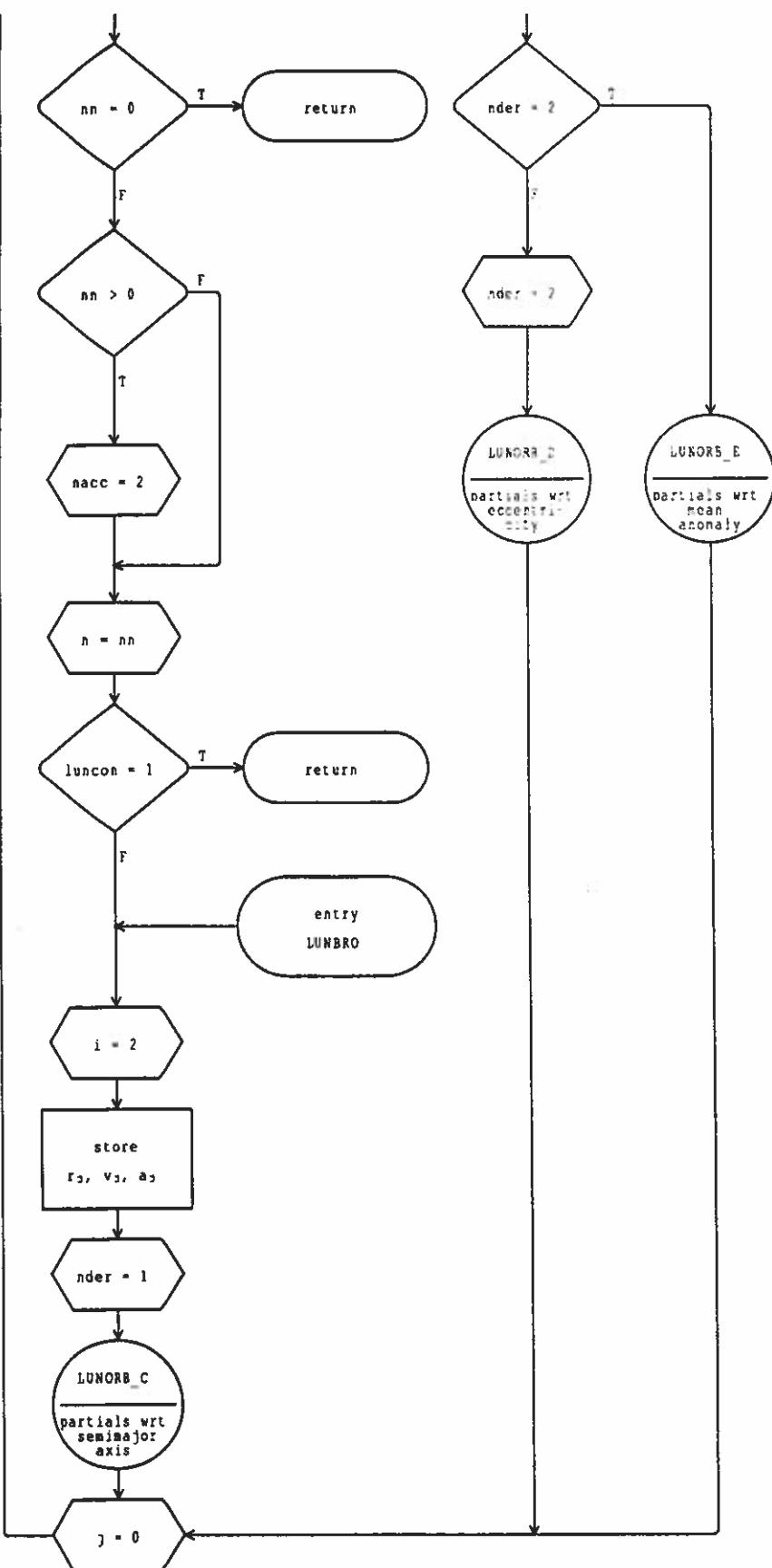
Entry Points: LUNORB, LUNBRO

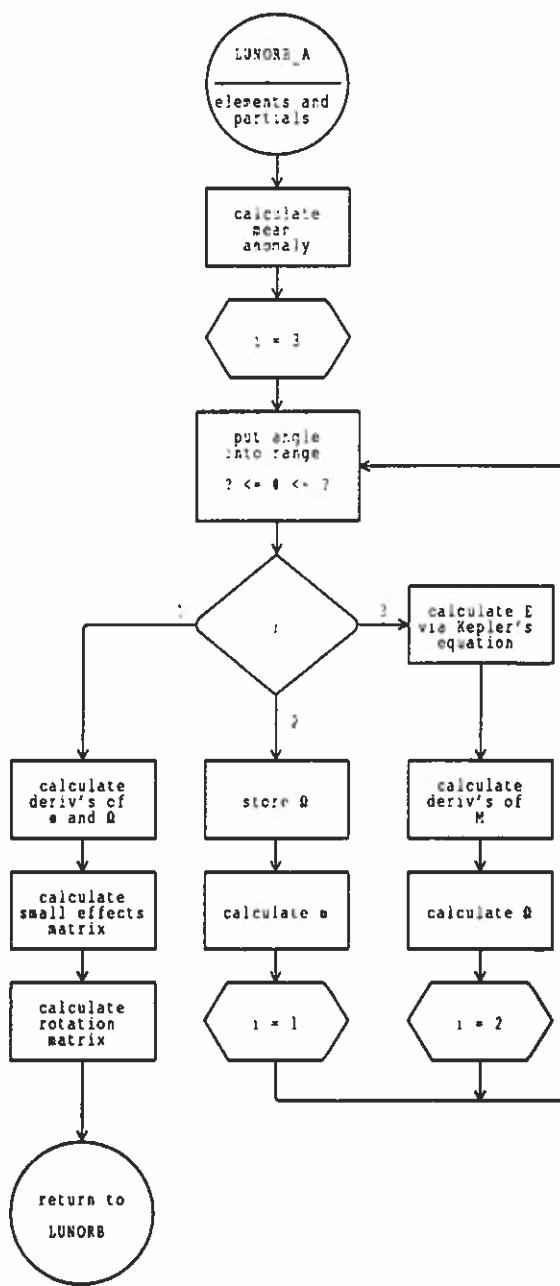
Includes: AACOFF, FUNCON, ORBLUN, ORBSTF, PRECMN

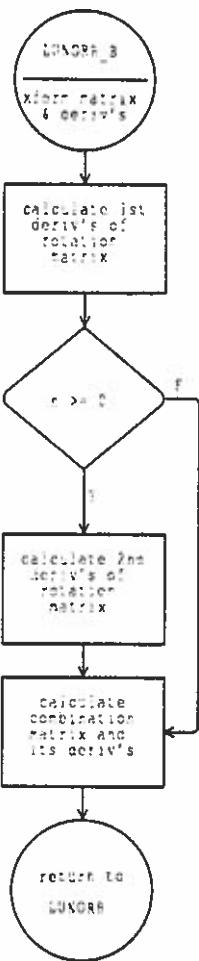
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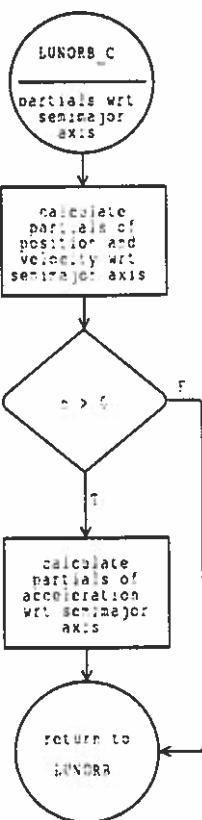


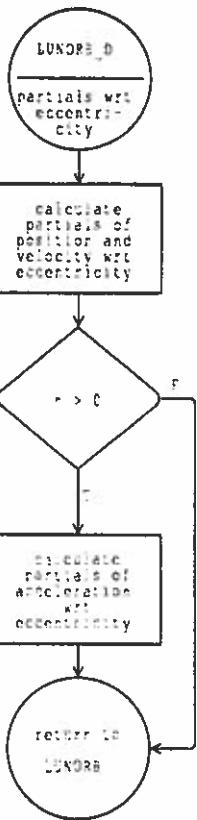


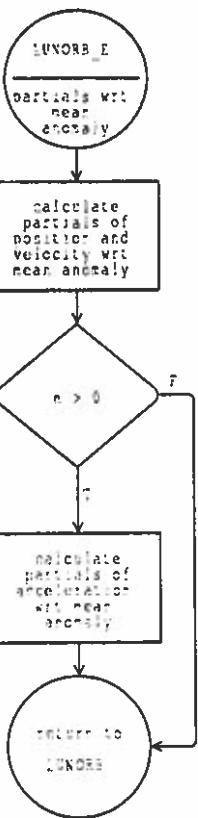


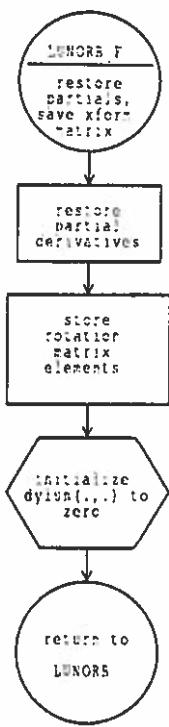


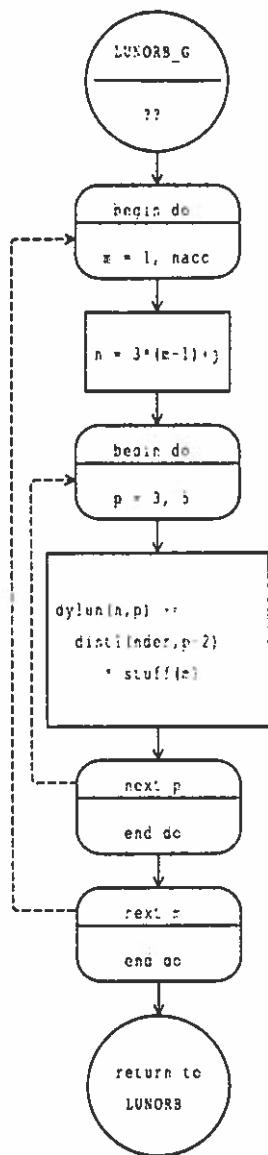


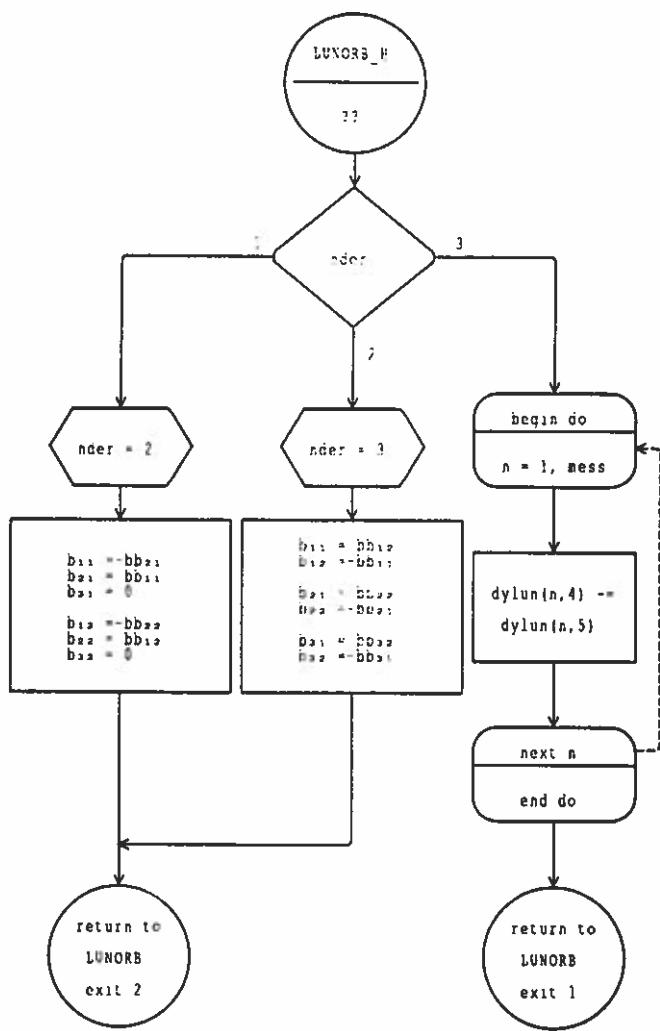












**PLANET**

**subroutine PLANET ( kfit )**

Description: PLANET is the driver routine for numerical integration of the equations of motion.

Arguments: kfit control integer:

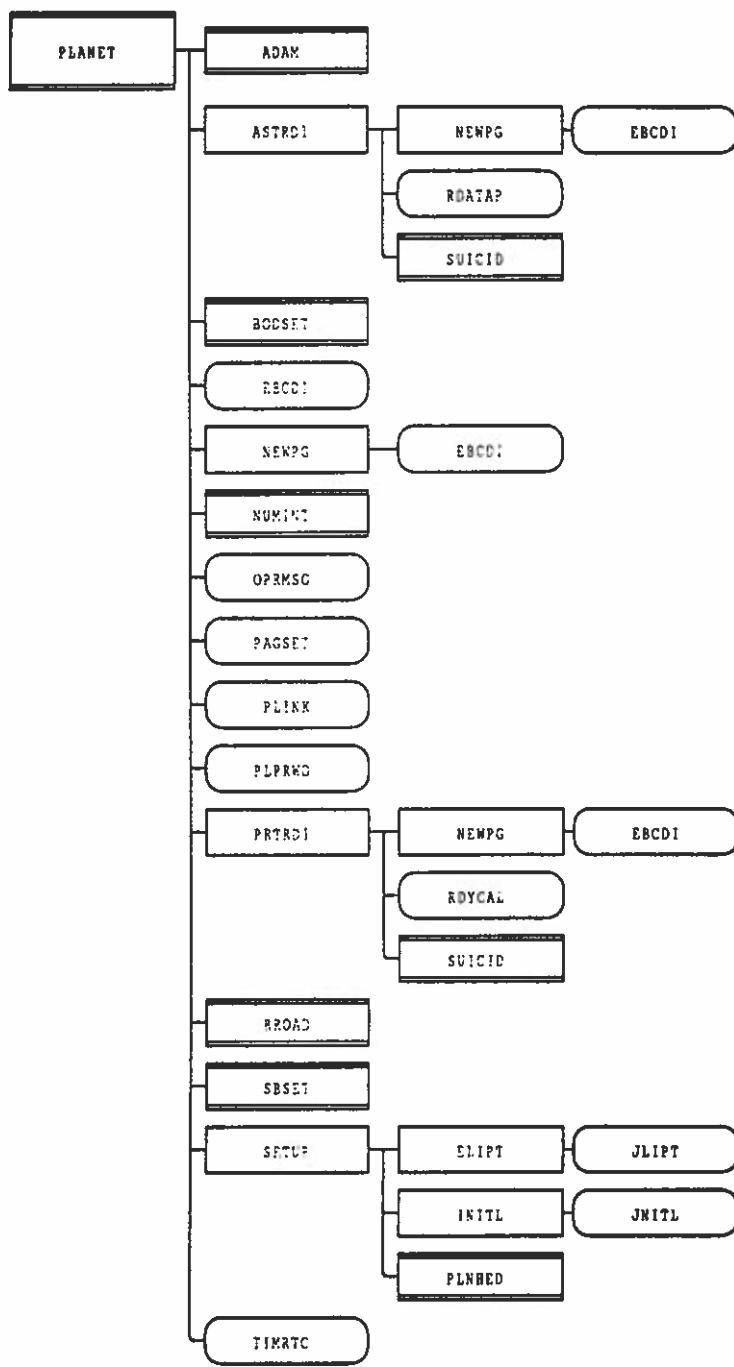
- 0 Integrations are in the midst of orbit fitting iterations.
- 1 Integrations are after orbit fitting convergence. The integration interval will be extended, but without partials calculations.

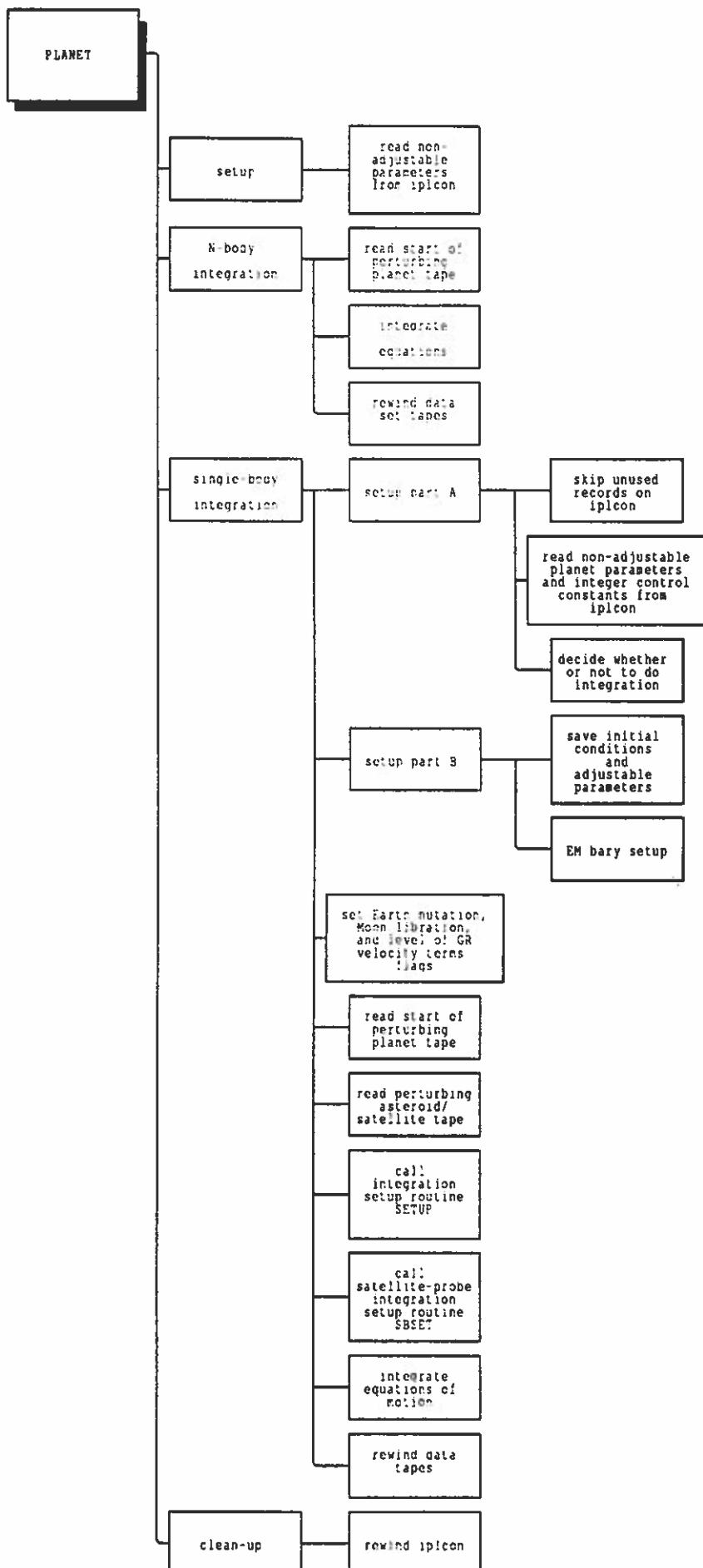
Calls: ADAM, ASTRD1, BODSET, EBCDI, NEWPG, NUMINT, OPRMSG, PAGSET, PLINK, PSRSD, PRTRD1, RROAD, SBSET, SETUP, TIMRTC

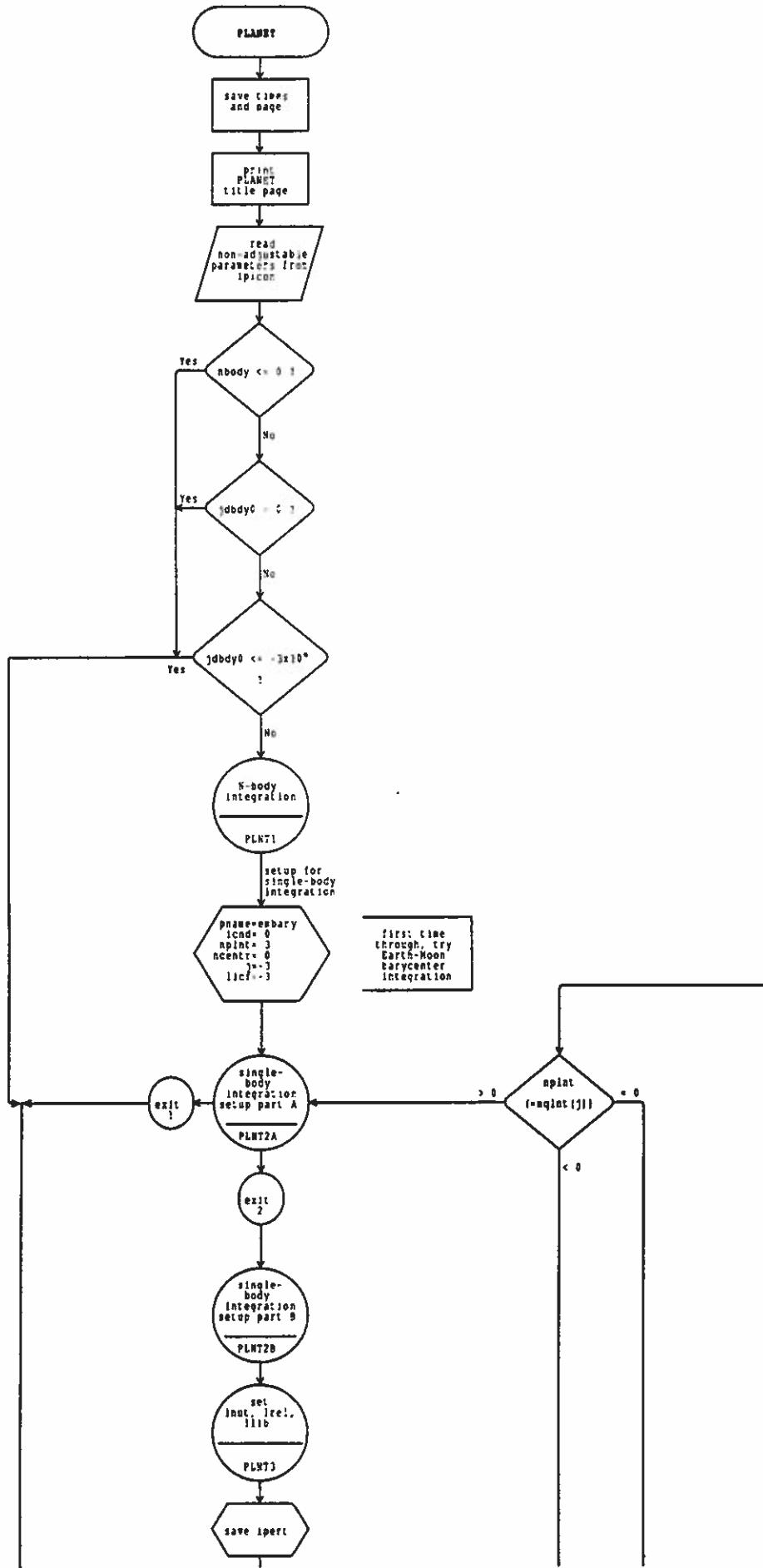
Called by: MAIN

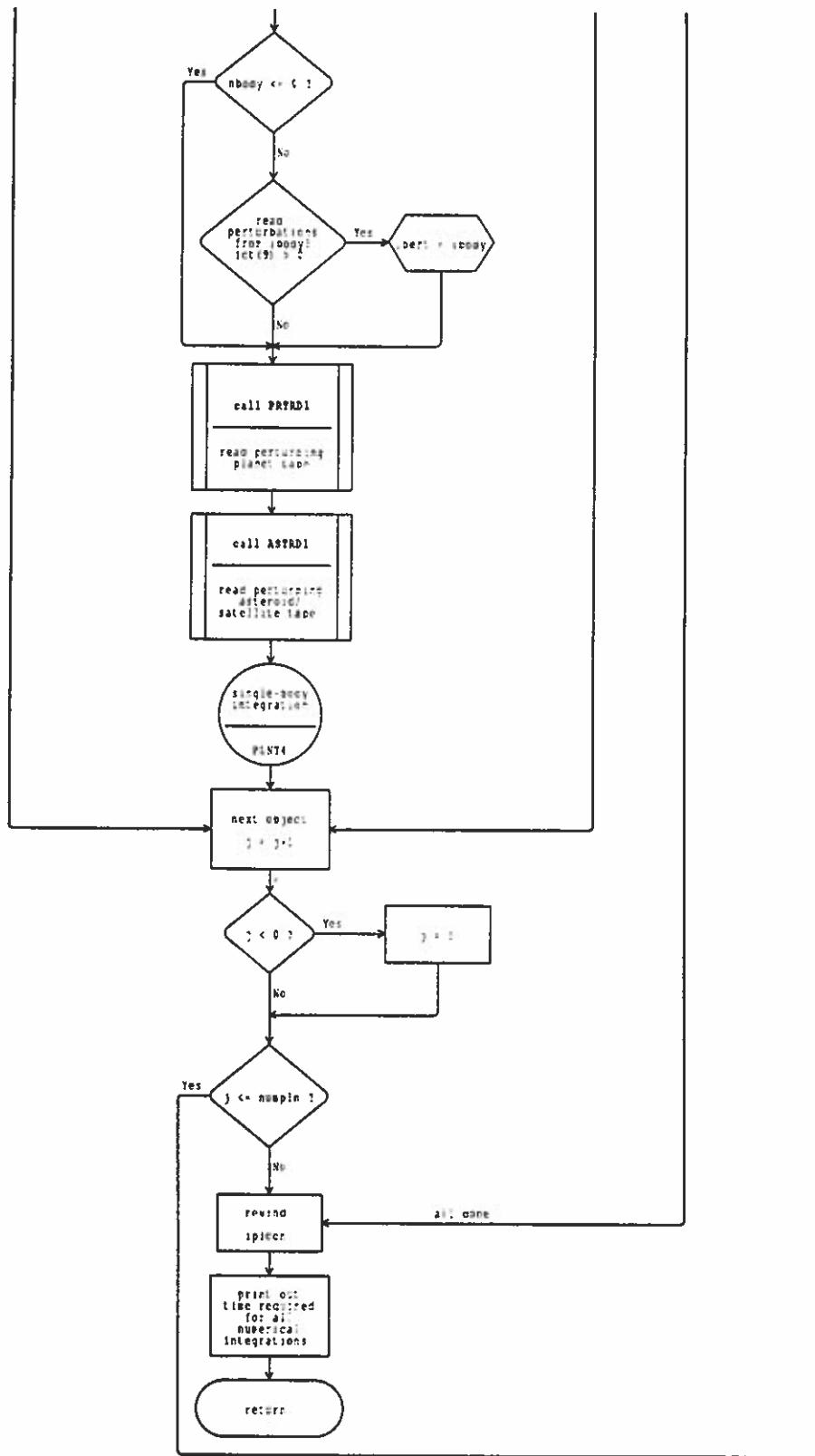
Include Files: BDCTRL, EMPCND, FCNTRL, INODTA, NAMTIMQ, PARAM, PETINA, PETUNA, PLNDDA, STBTST, STINT, TIMSTF

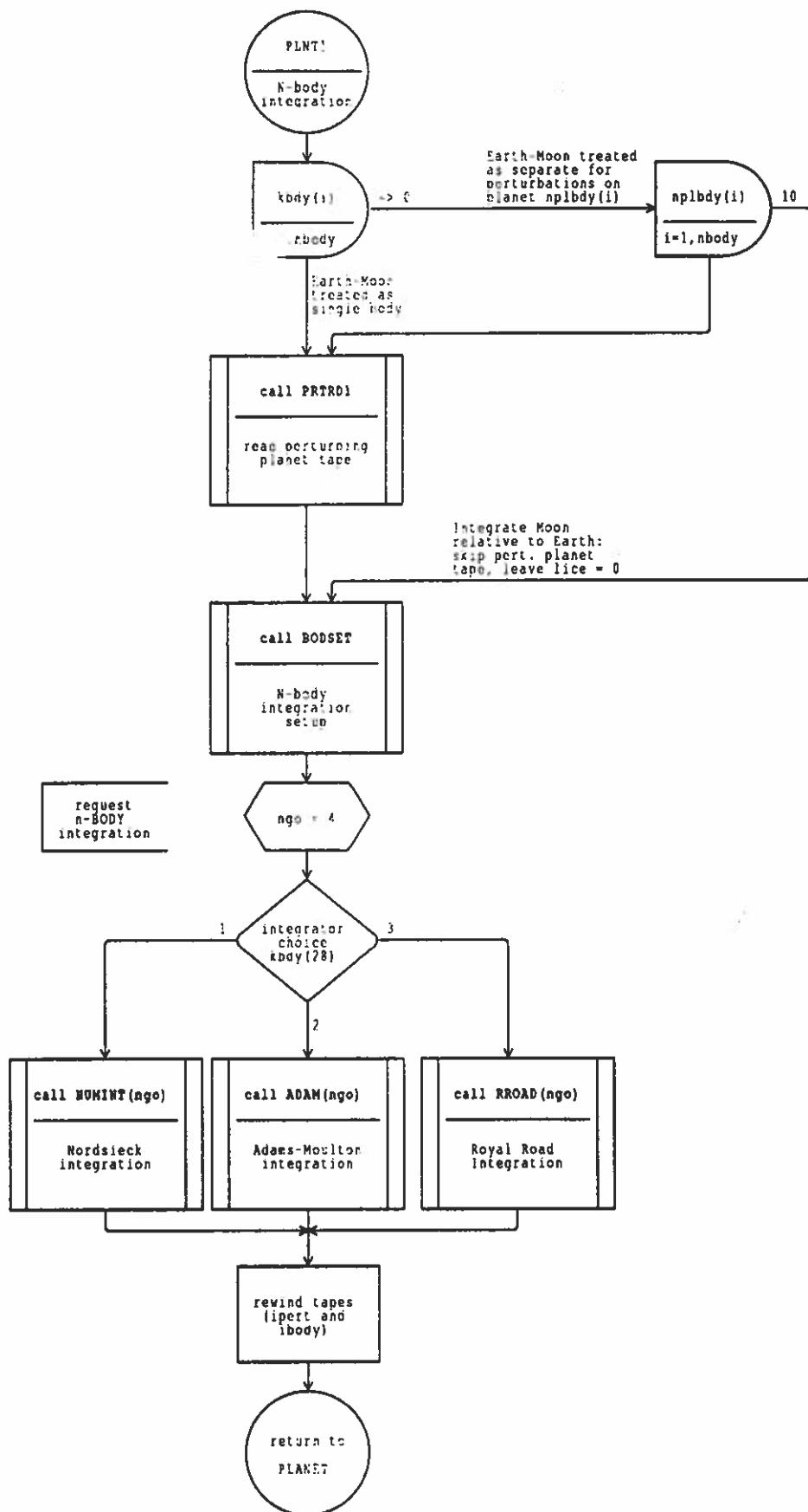
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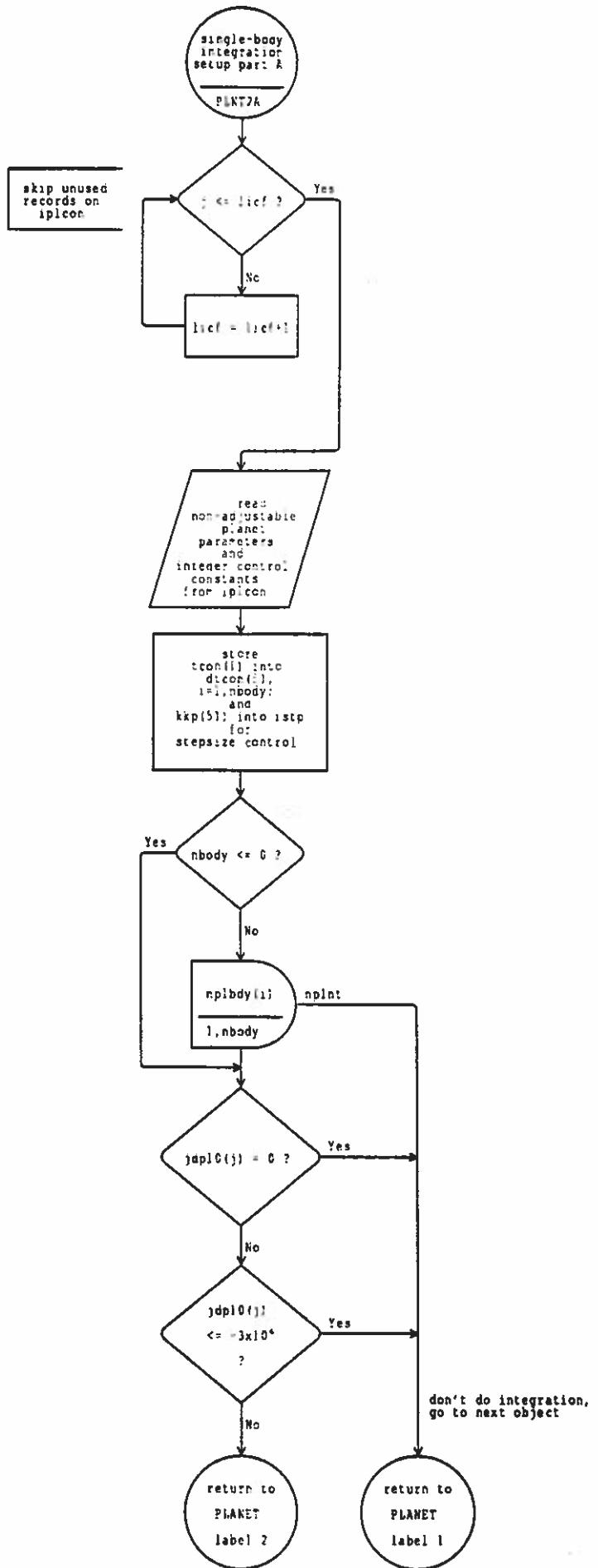












single-body  
integration  
setup part B

PINT75

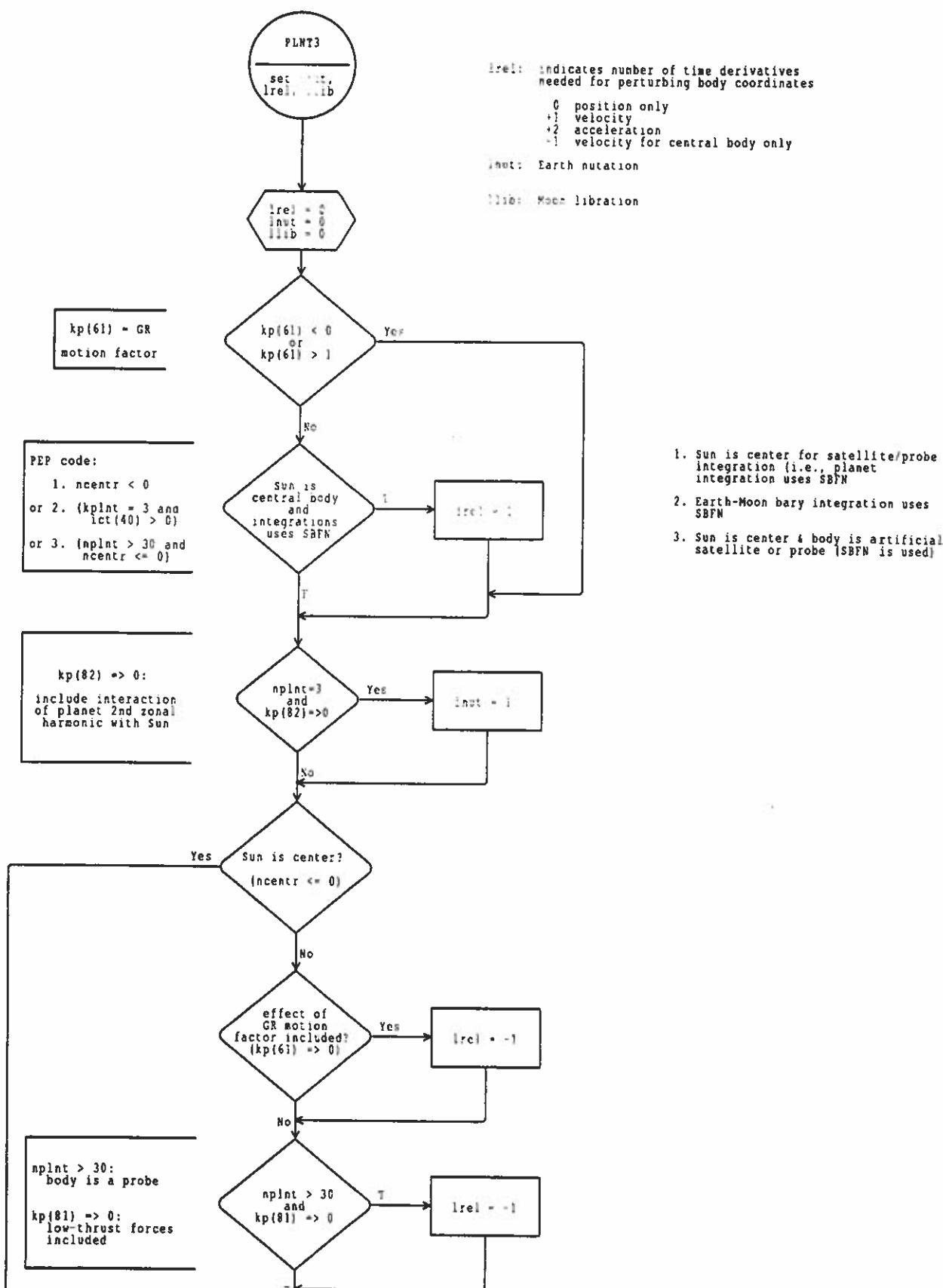
store pcond(i) into  
cond(i) for i = 1,30  
(initial conditions  
and adjustable  
parameters)

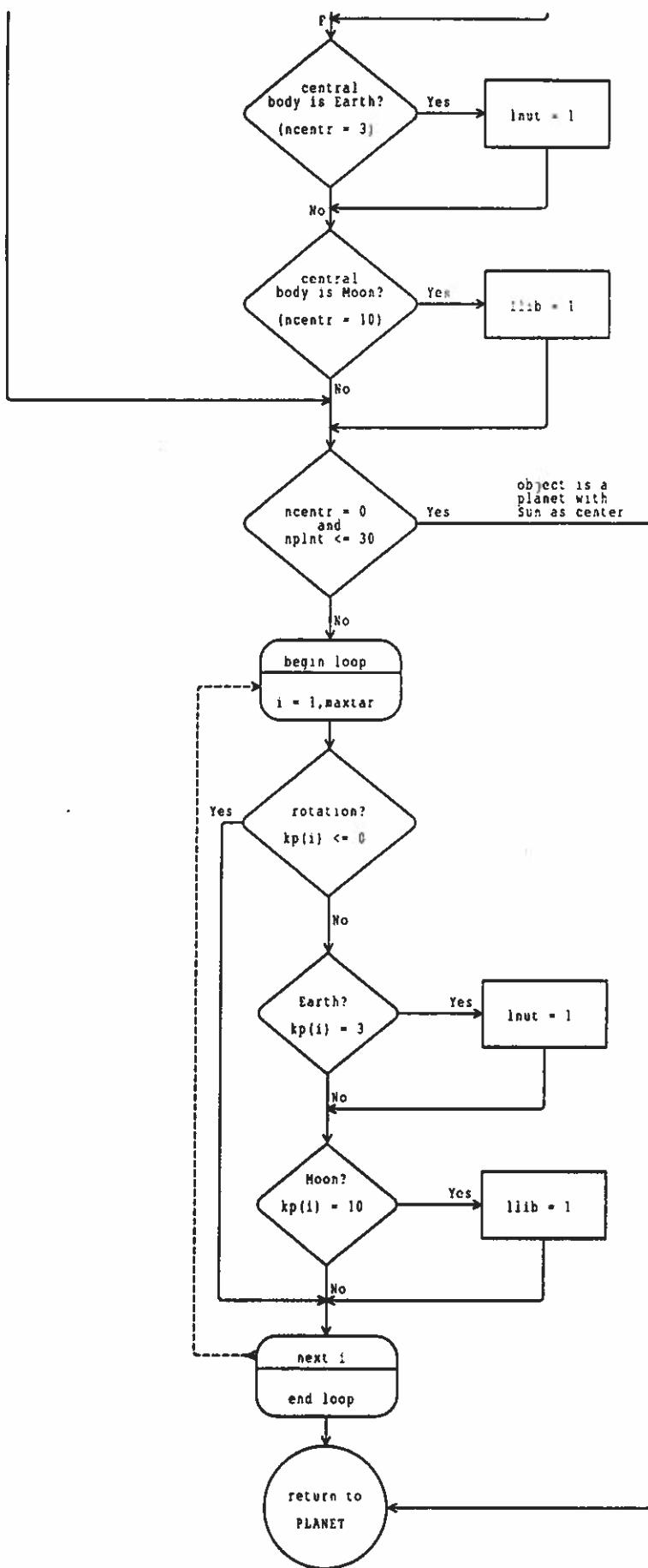
Earth-Moon  
barycenter?  
aplat = 3

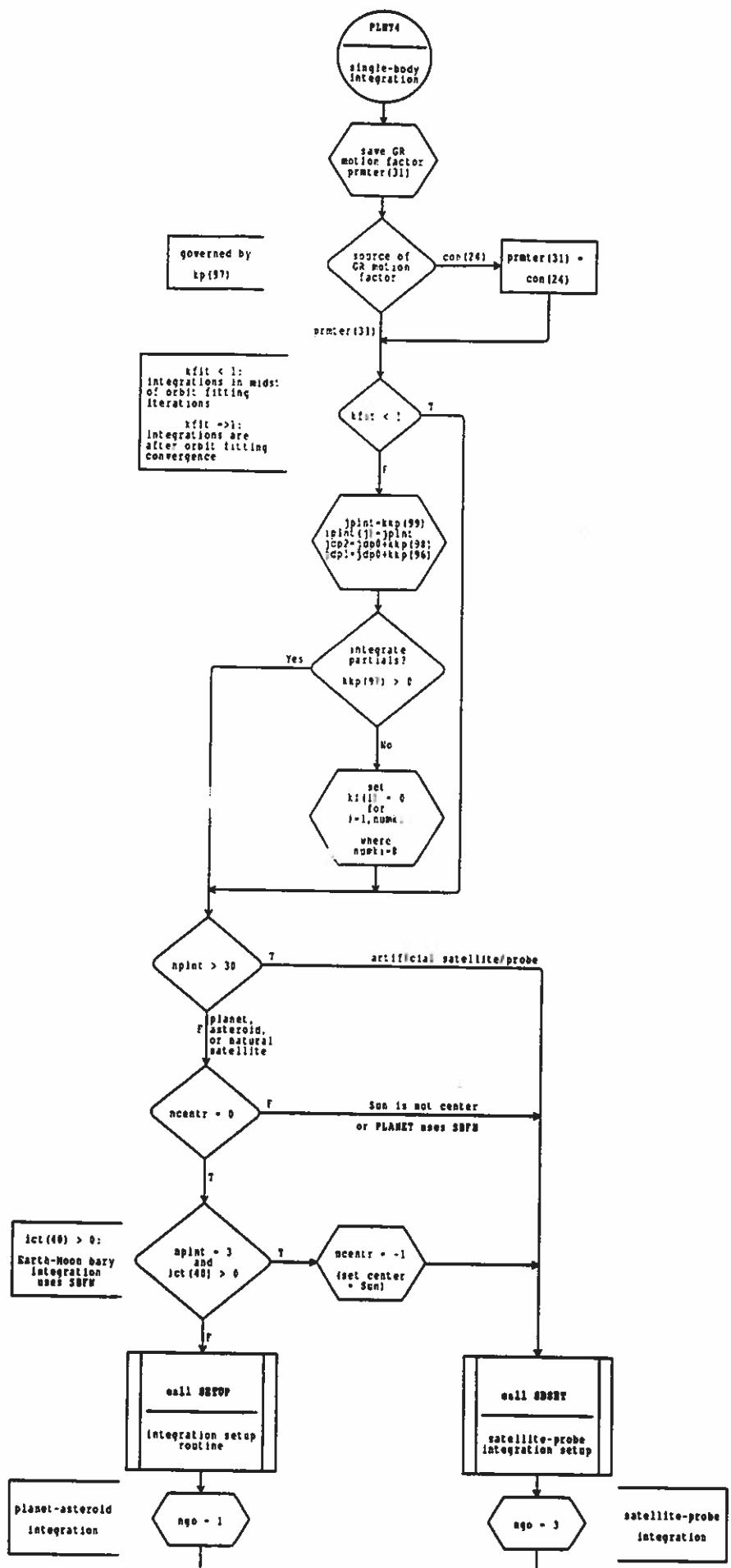
No

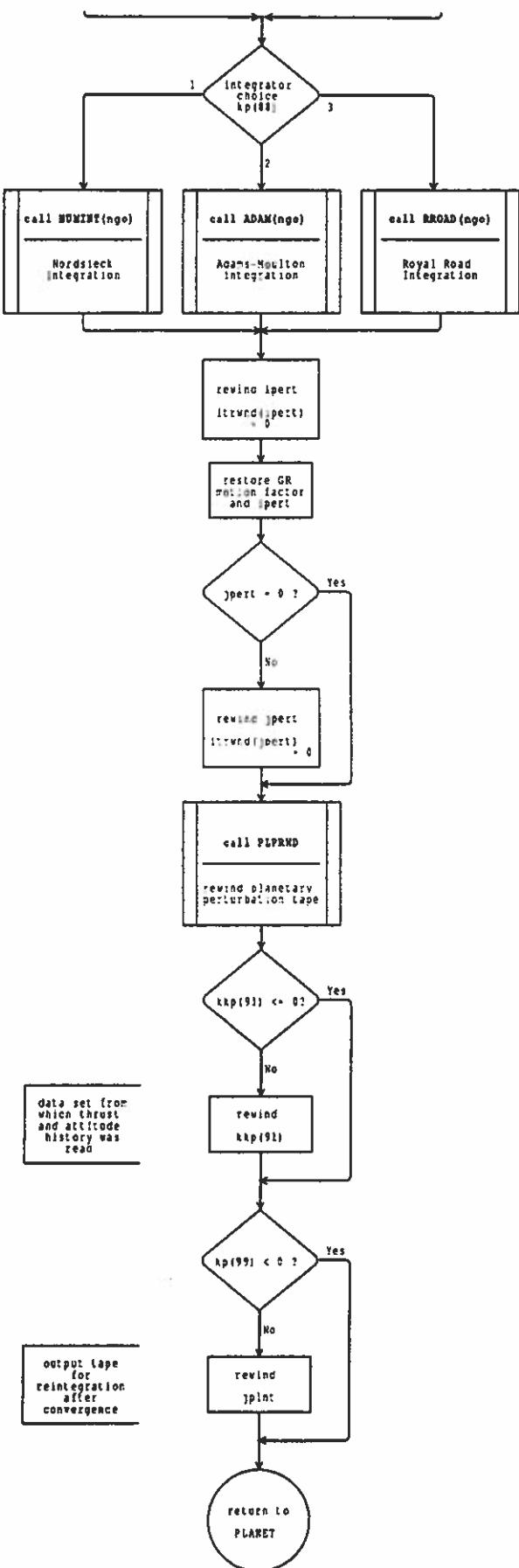
prare = aplat();  
ncenter = npcent();  
icnd = jcond();  
clat = ;

return to  
PLANET









**SBFN**

**real\*8 function SBFN( k, j, s )**

Description: SBFN evaluates (mostly via a call to SBFN1) the equations of motion and partial derivatives for satellites and probes. It does a fair amount of setting up for calculations performed in SBFN1, which contains most of the right hand sides of the differential equations being solved.

Arguments: k      Equation number. There are six equations for the motion and six for each partial derivative. A detailed description of k is provided elsewhere.

j      Integrator iteration number (generally takes on values from 1 to 3, and is often used as a logic flag).

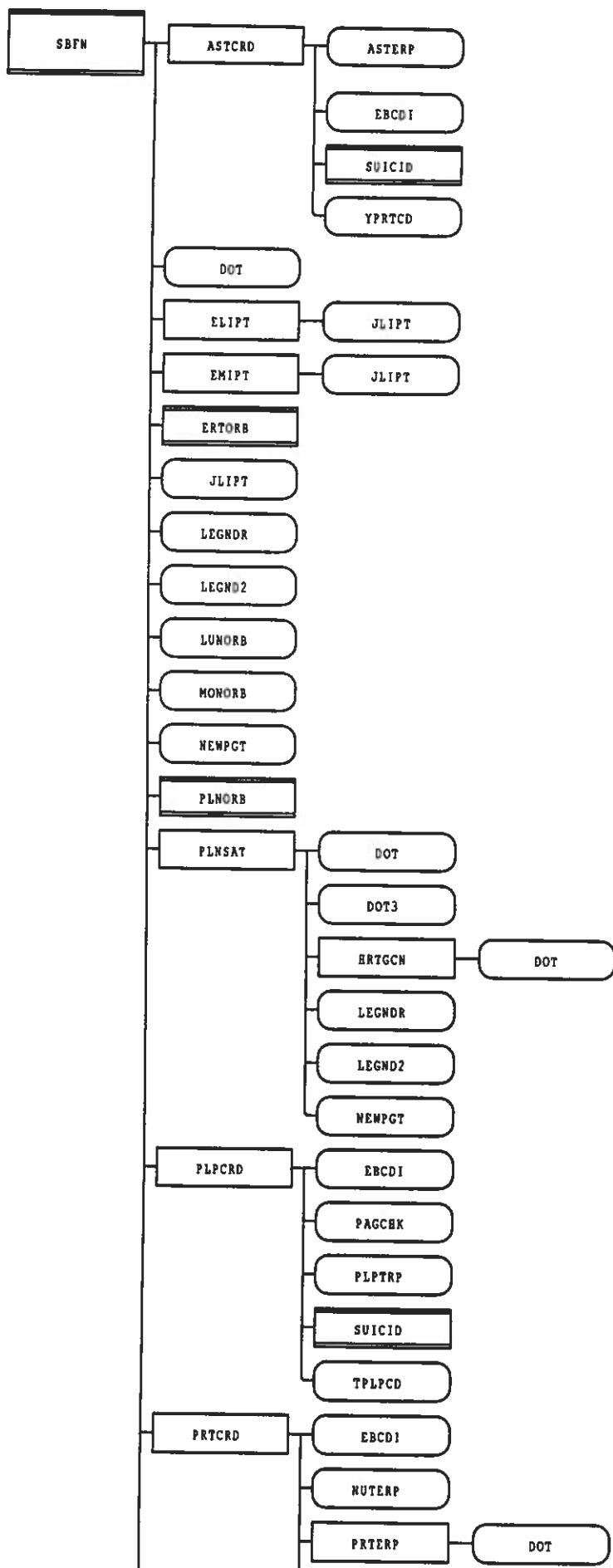
s      Current time = julian date + 0.5 day.

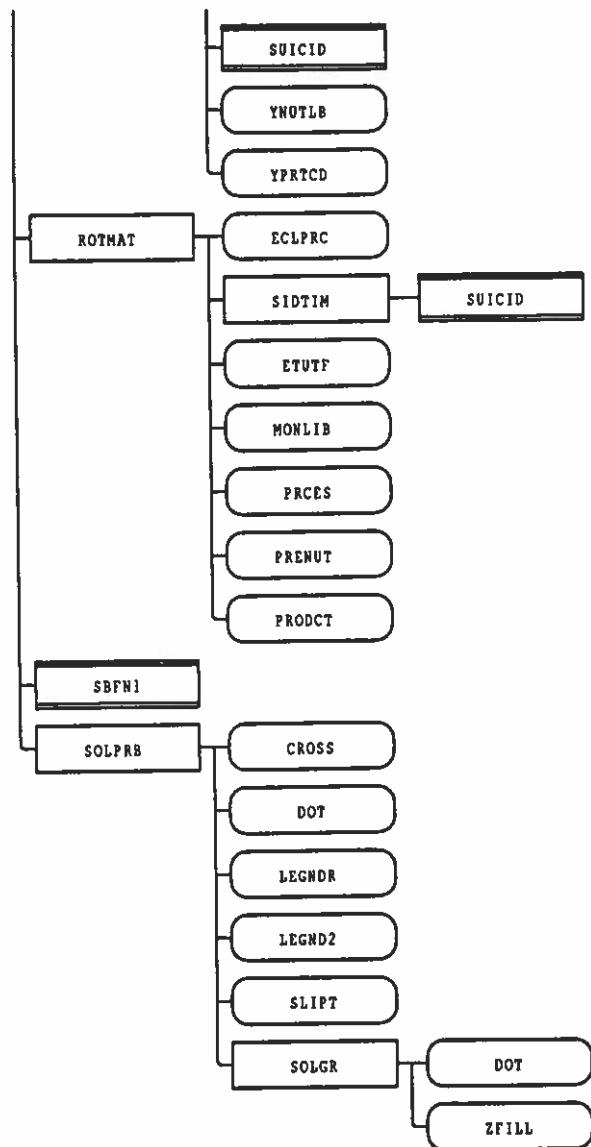
Calls: ASTCRD, DOT, ELLIPT, EMIPT, ERTORB, JLIPT, LEGNDR, LEGND2, LUNORB, MONORB, NEWPGT, PLNORB, PLNSAT, PLPCRD, PRTC RD, ROTMAT, SBFN1, SOLPRB

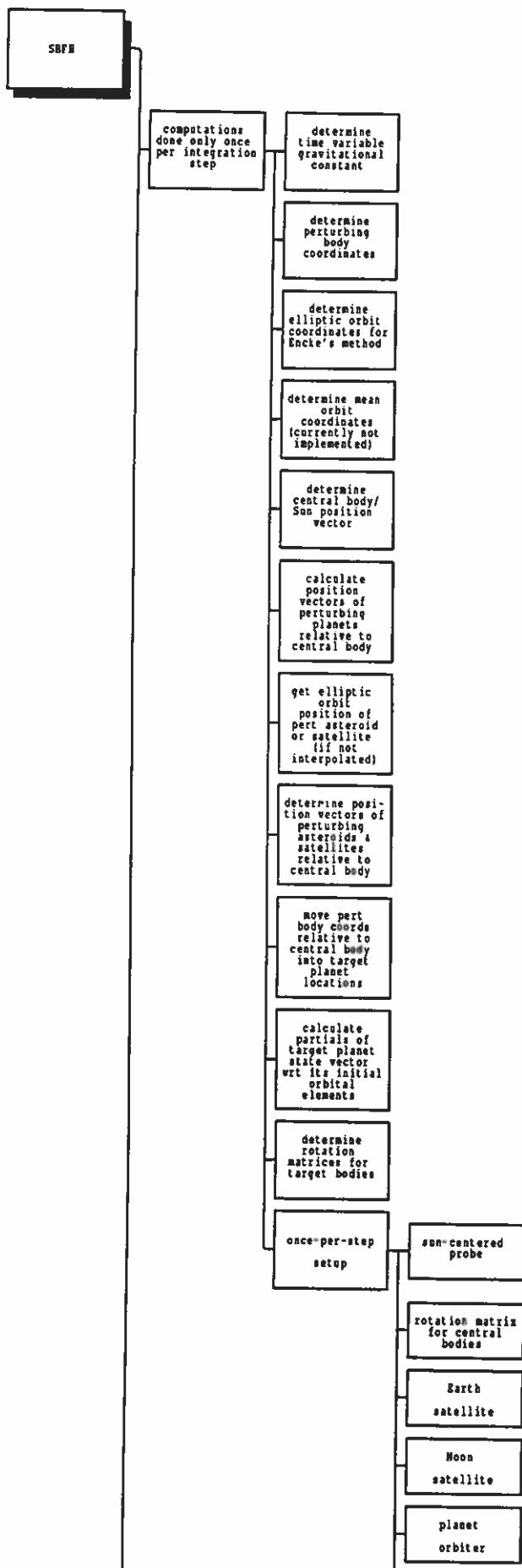
Called by: ADAM, EVAL, NINT, RROAD

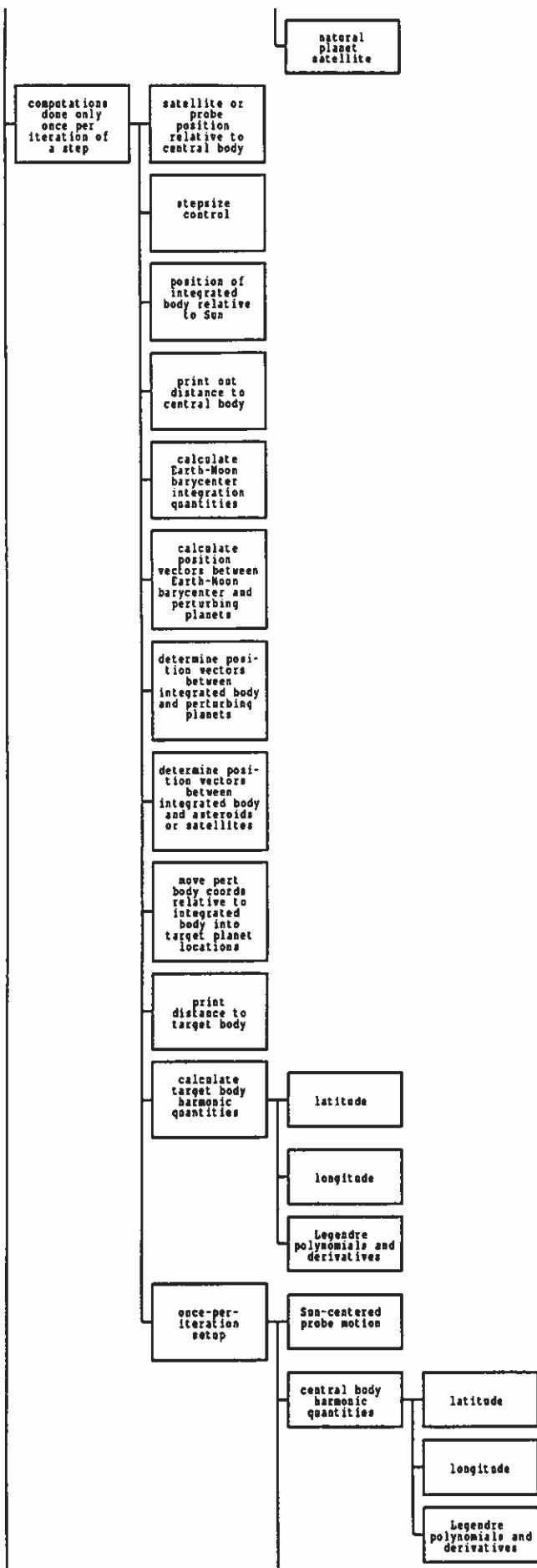
Include Files: CNTCHAR, ELLIPS, EMMIPS, FCNTRL, FMMIPS, INODTA, ORBLUN, OUTPUT, PARAM, PETUNA, PRTC COD, SBEMBR, SBROT, SBSTUF, SBTHNG, STINT, TRGHAR

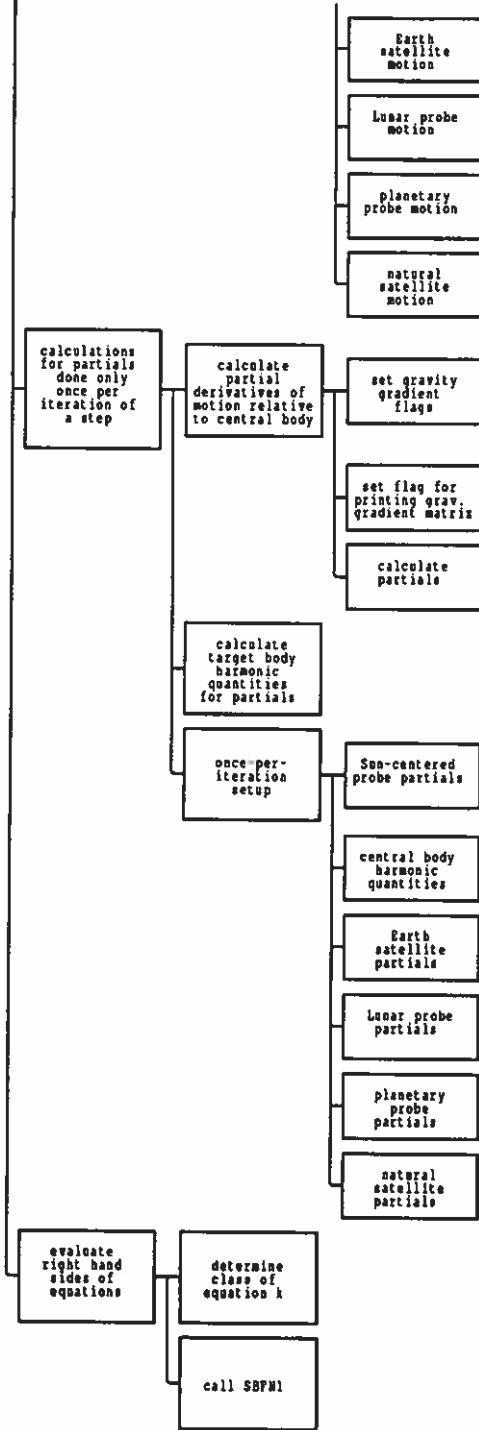
Equivalence: ggfgs(2)      /sbstuf/      SBSTUF

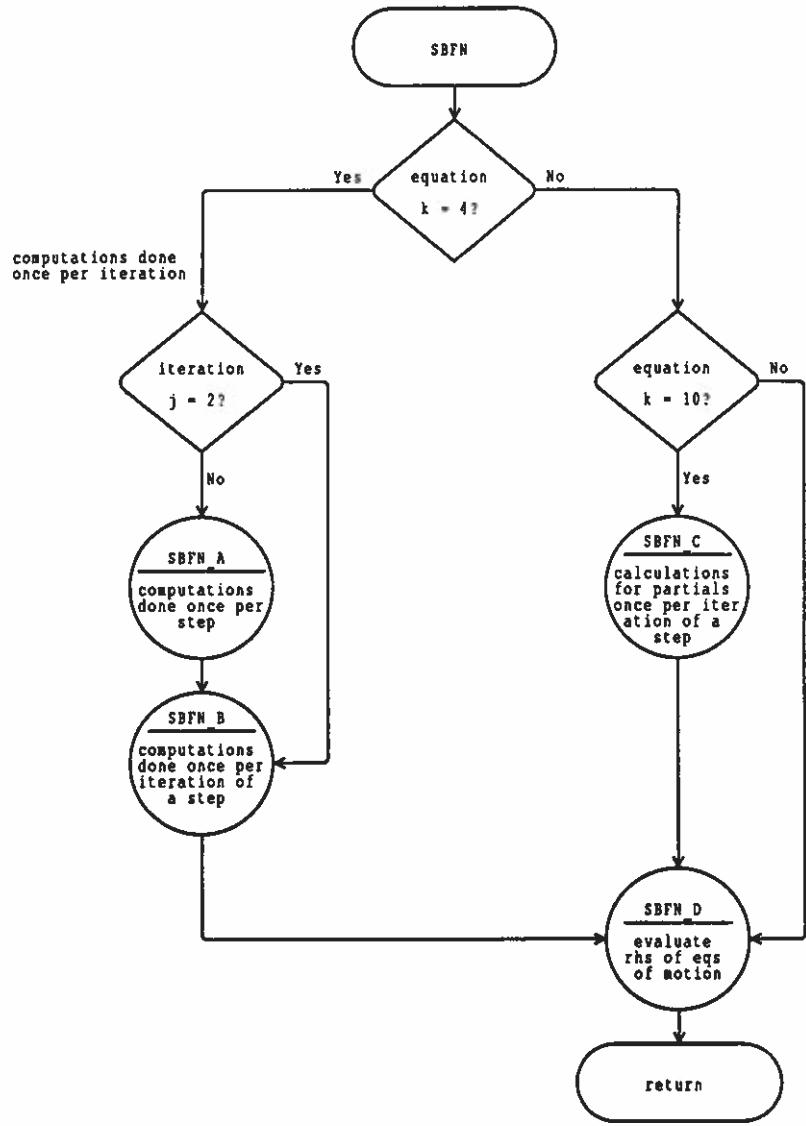


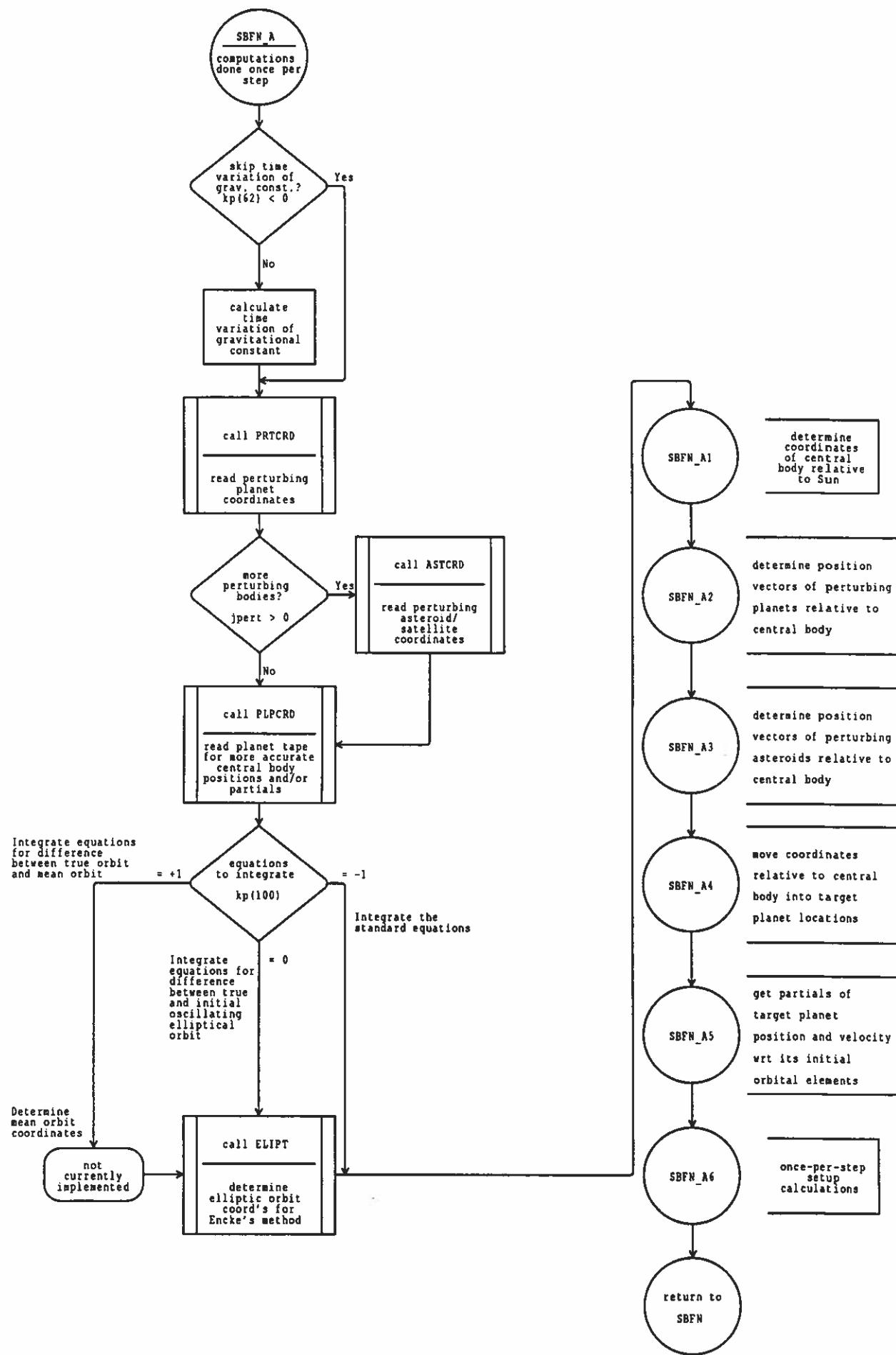


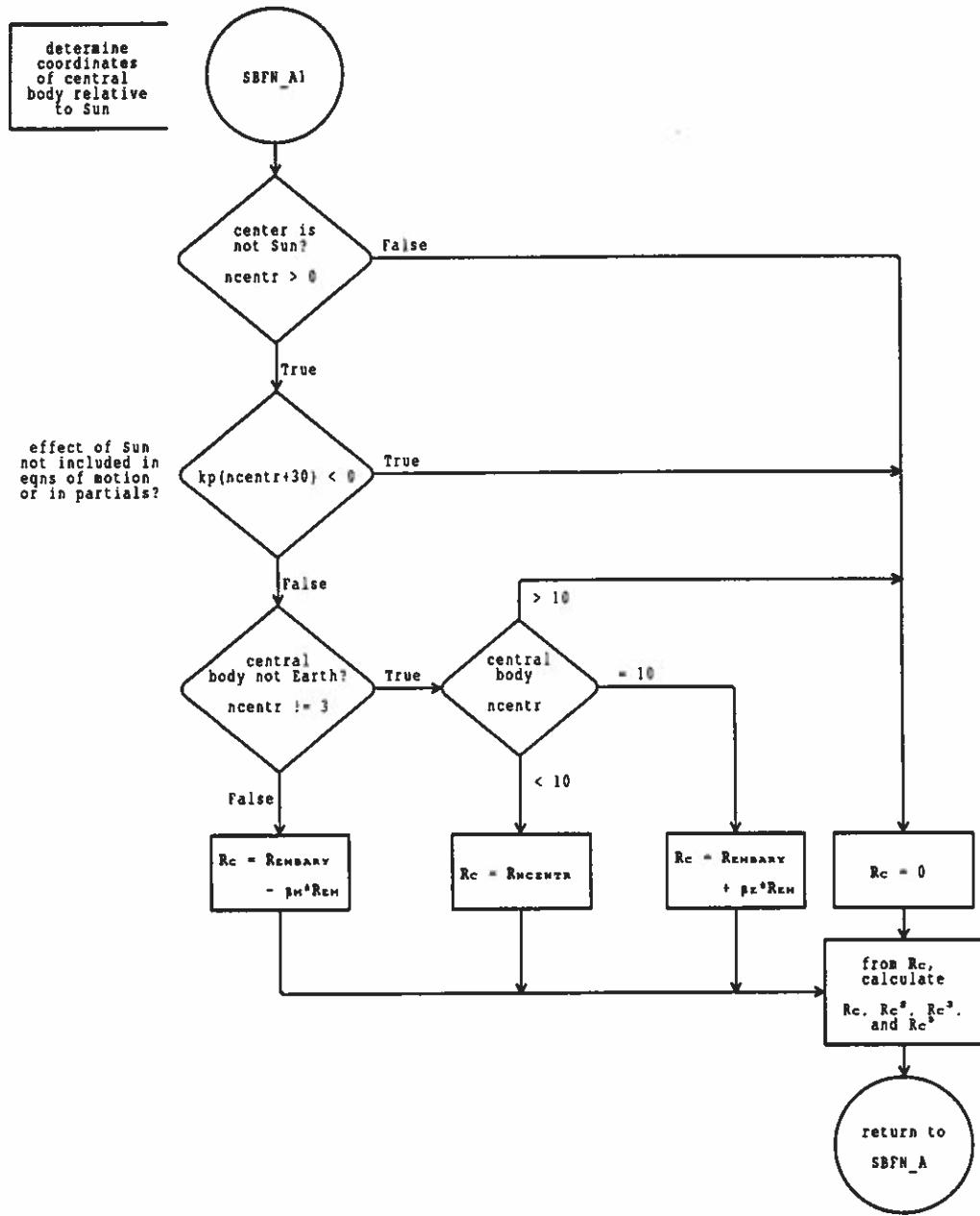


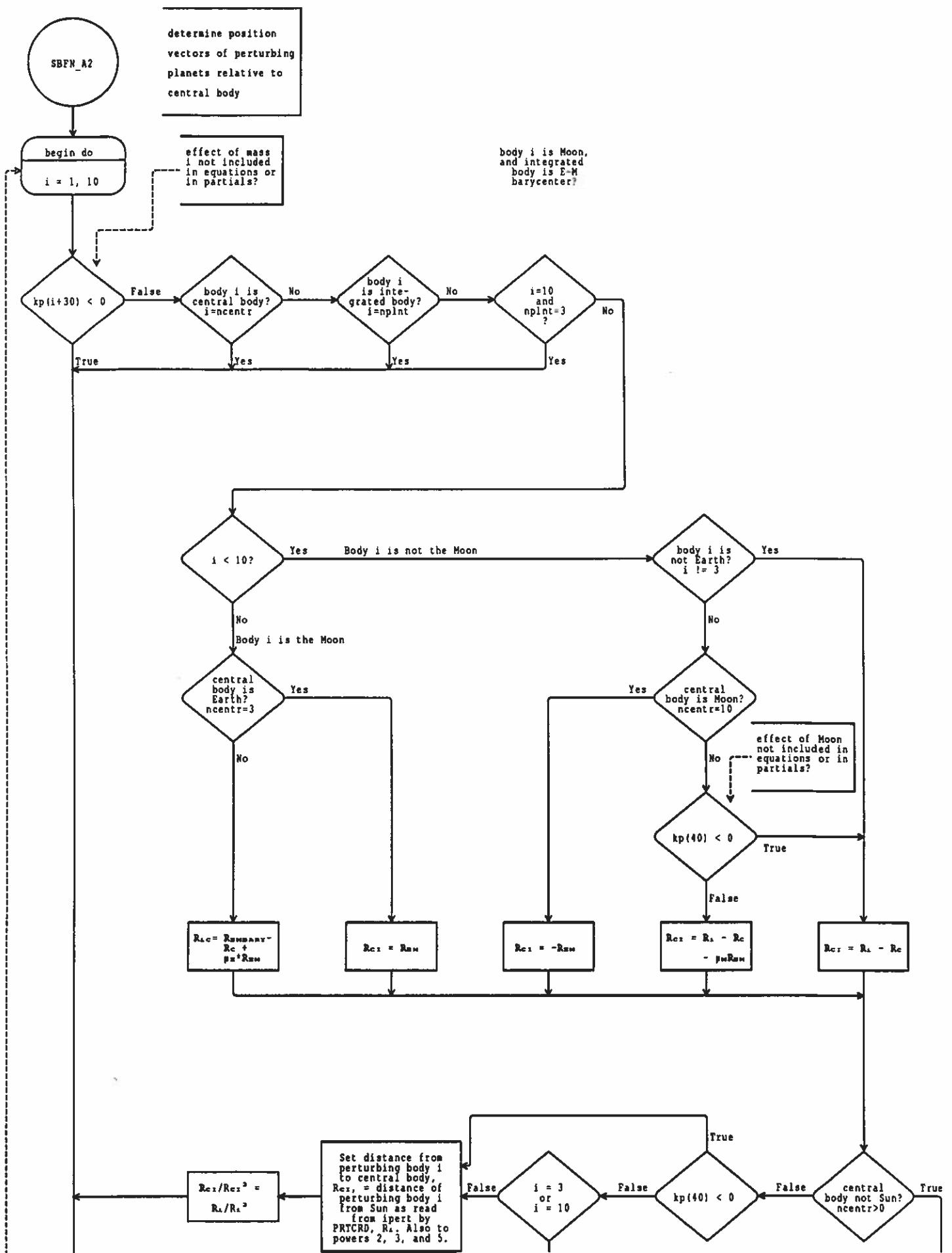


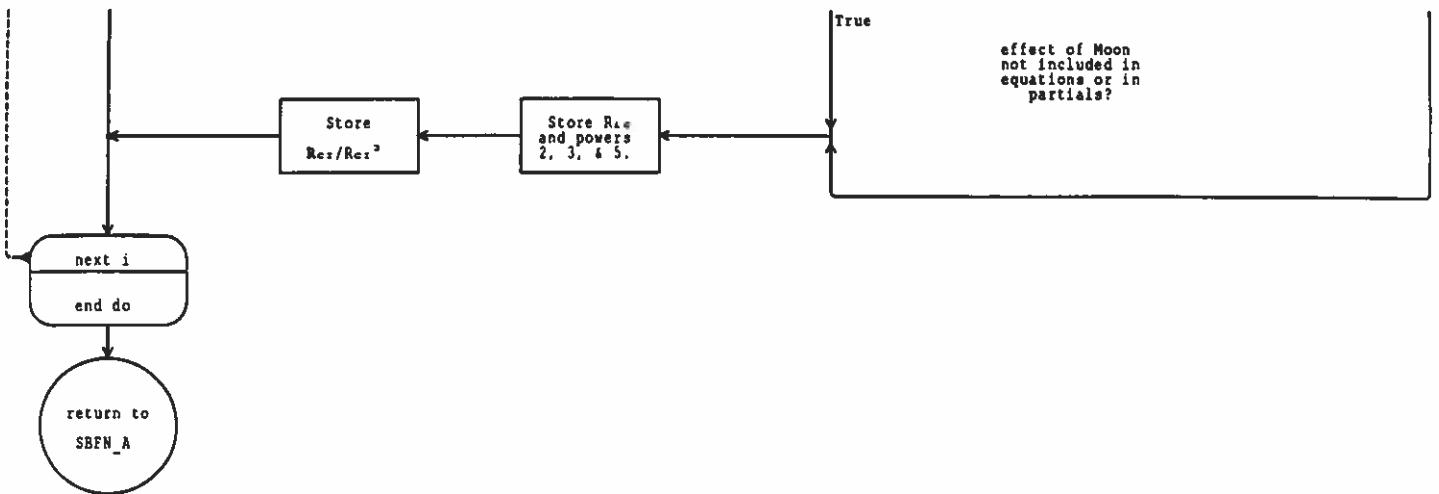


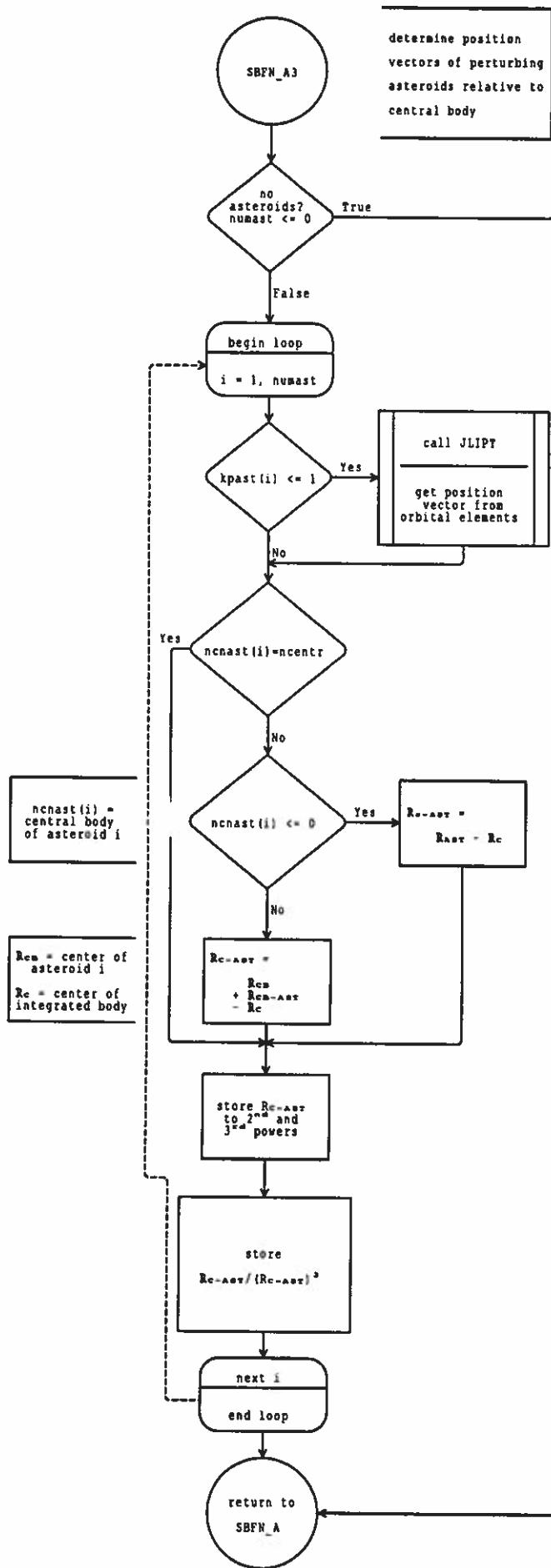


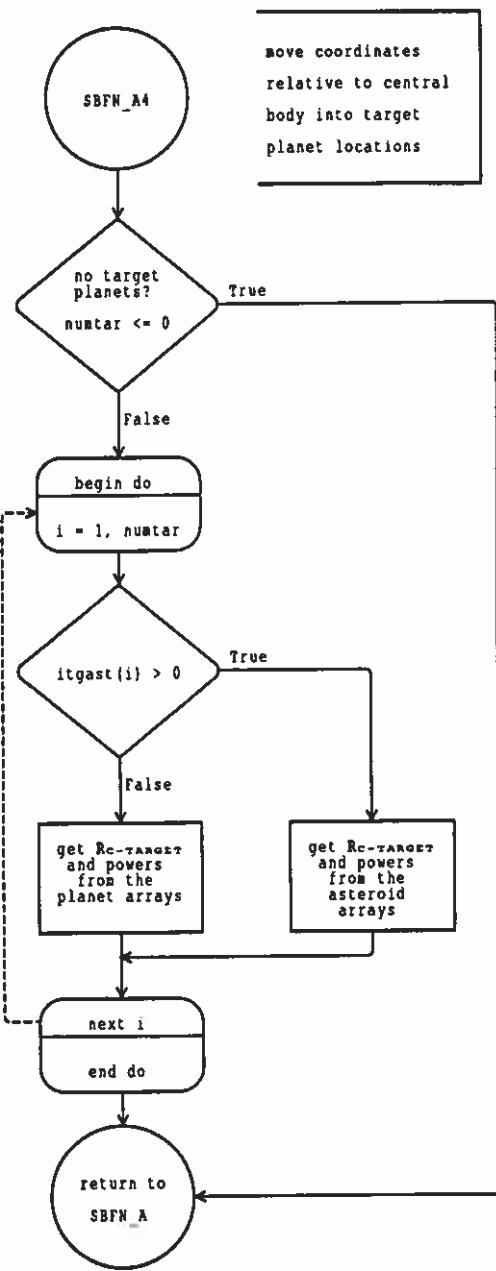


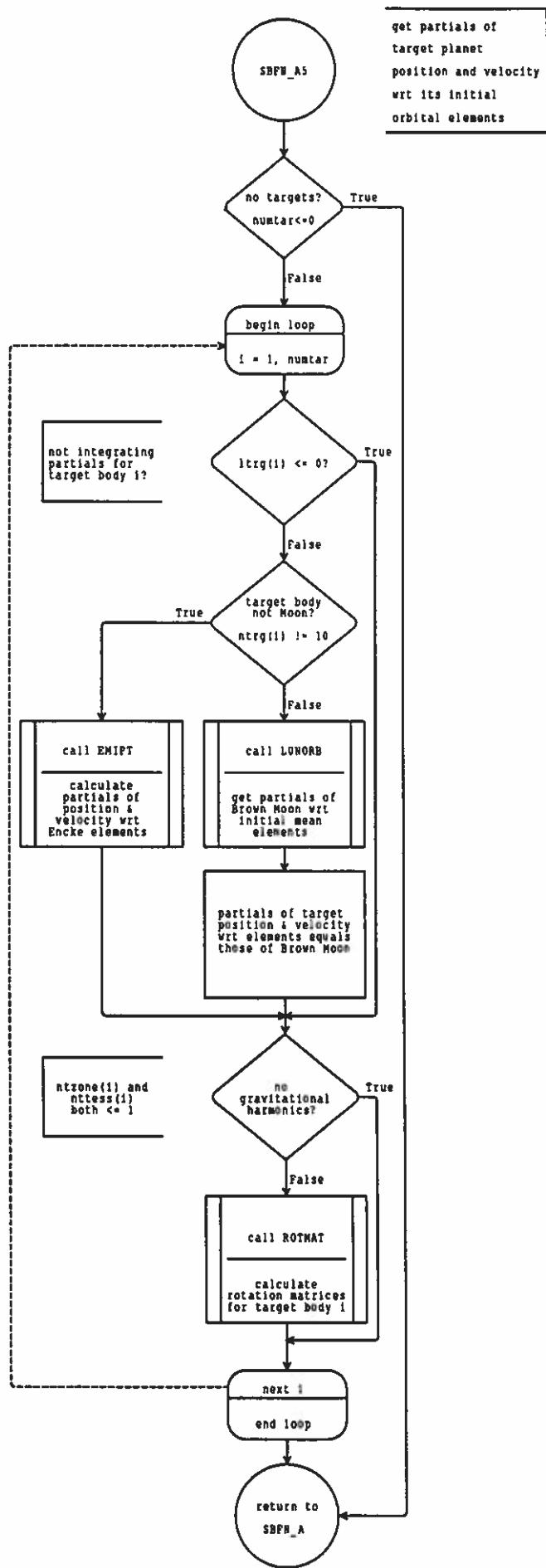


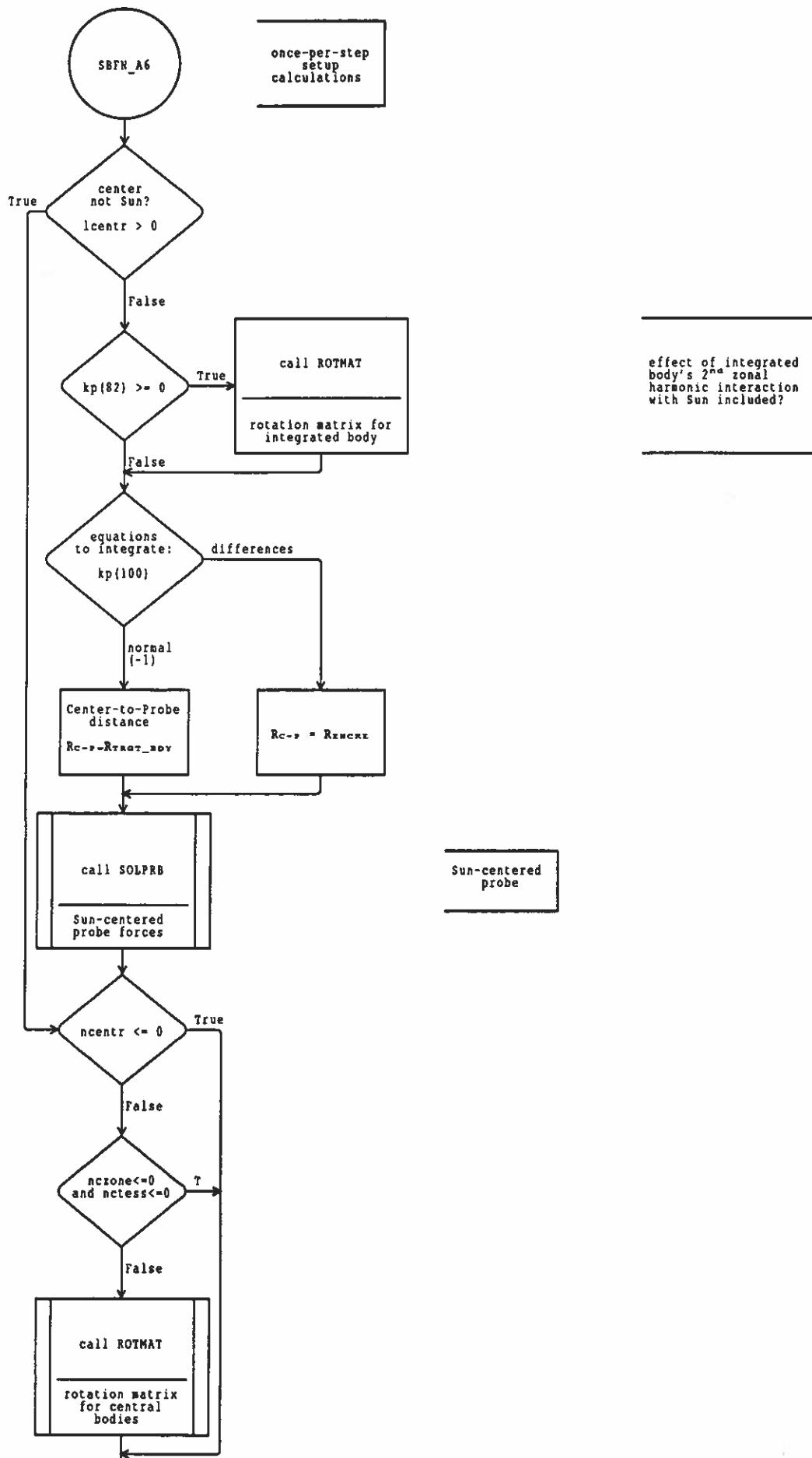


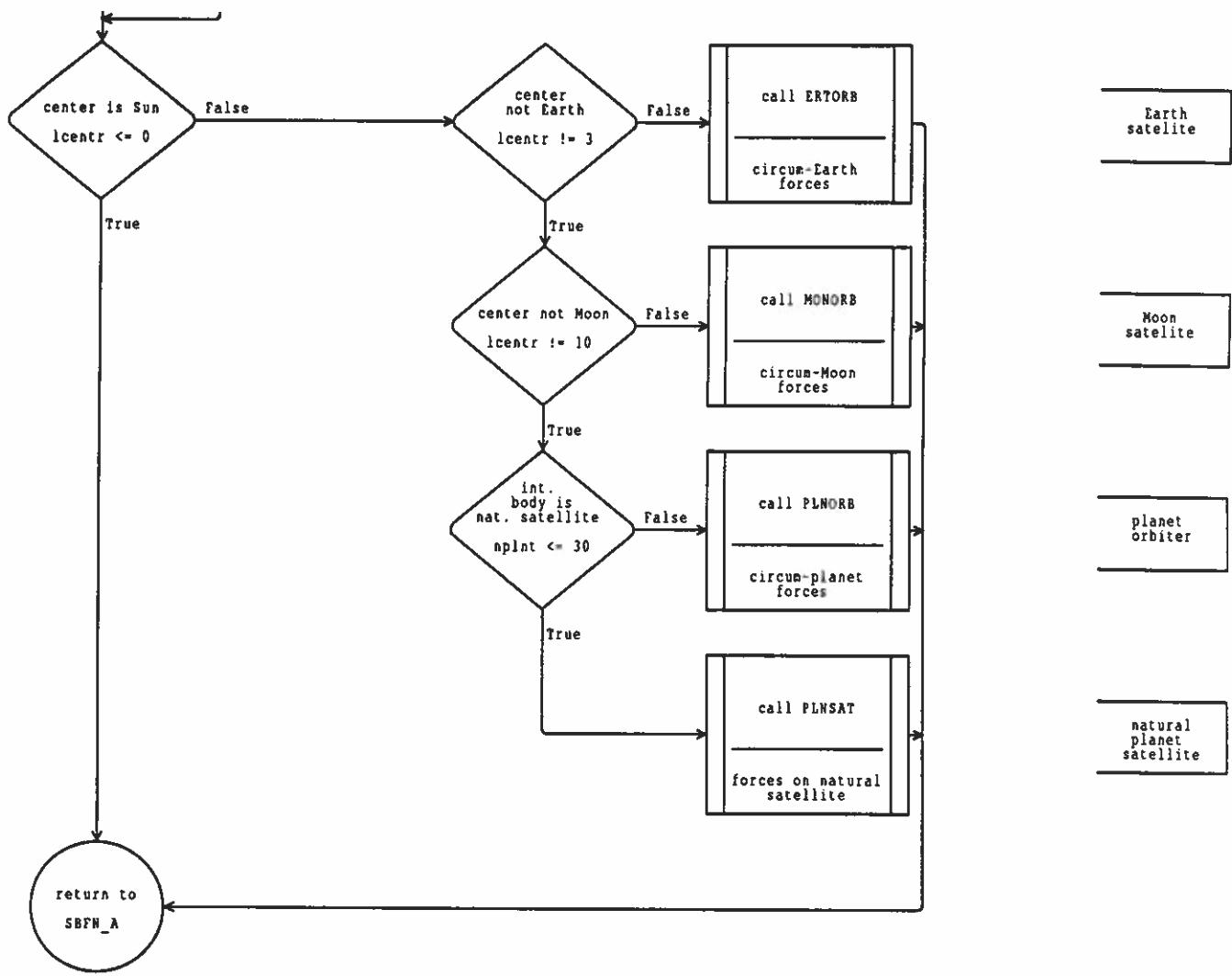


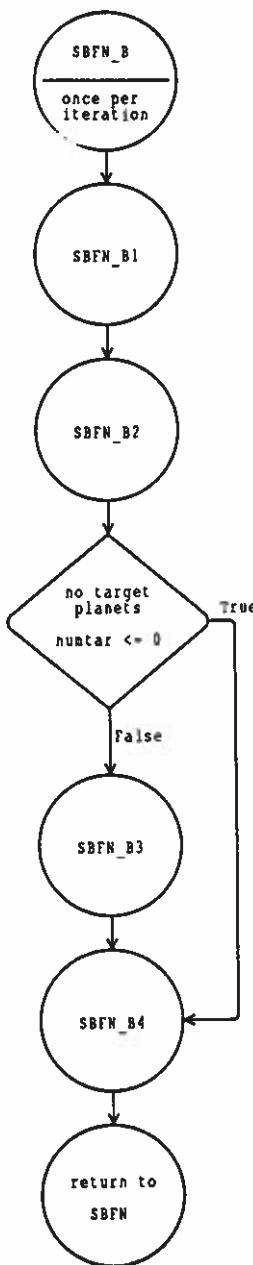


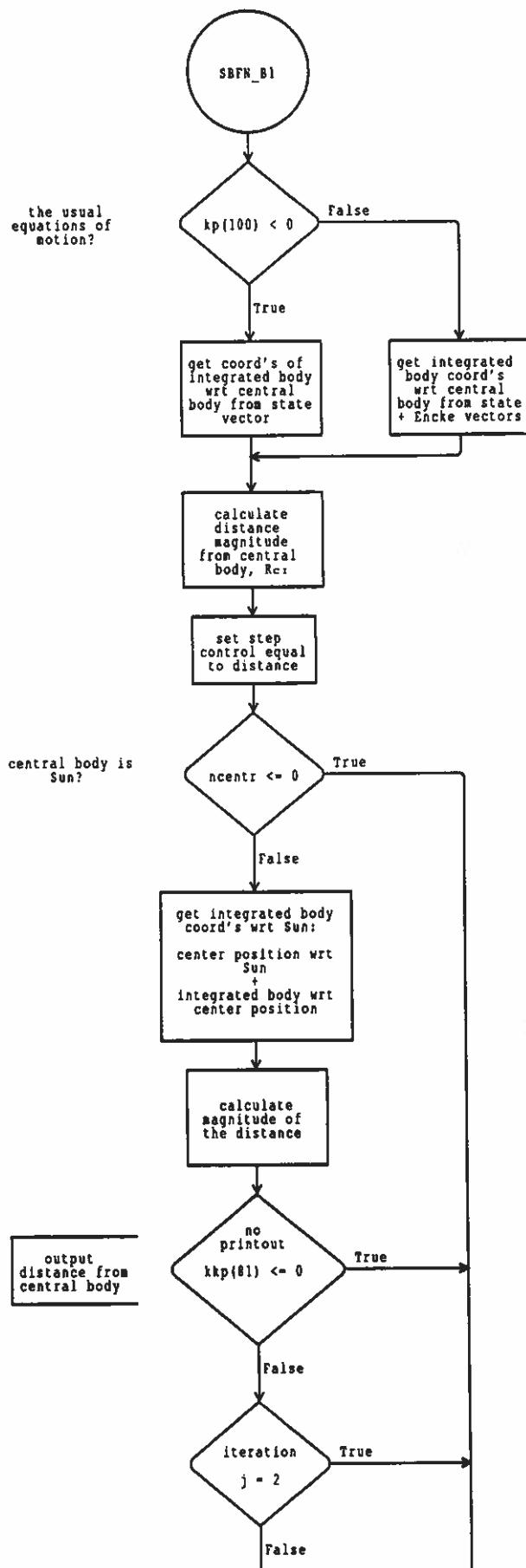


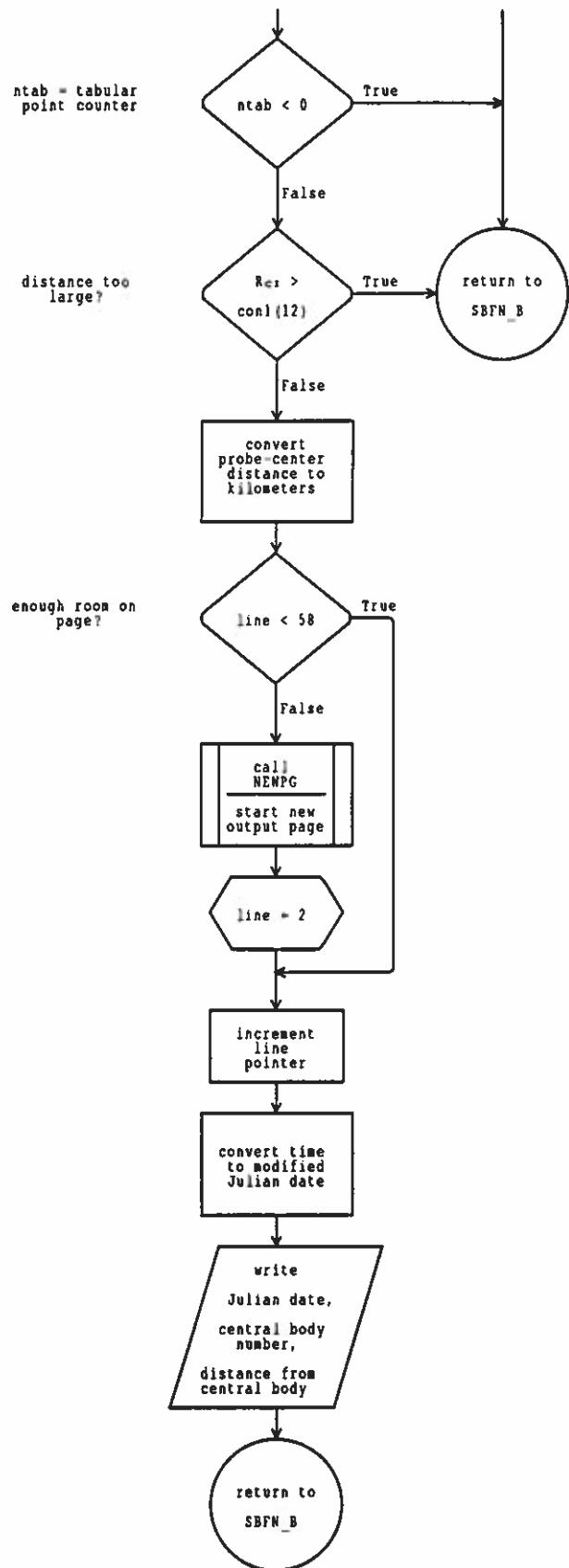


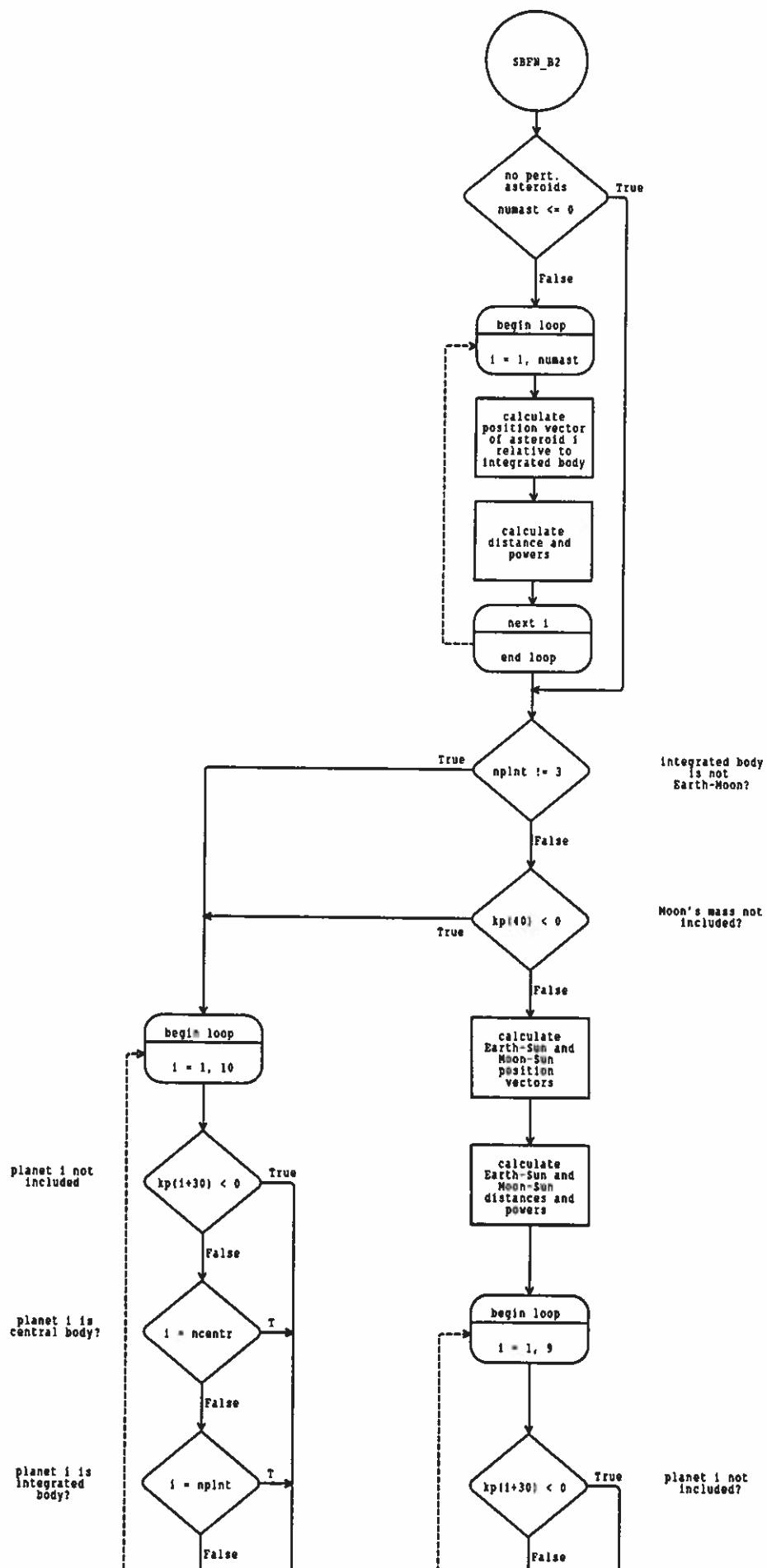


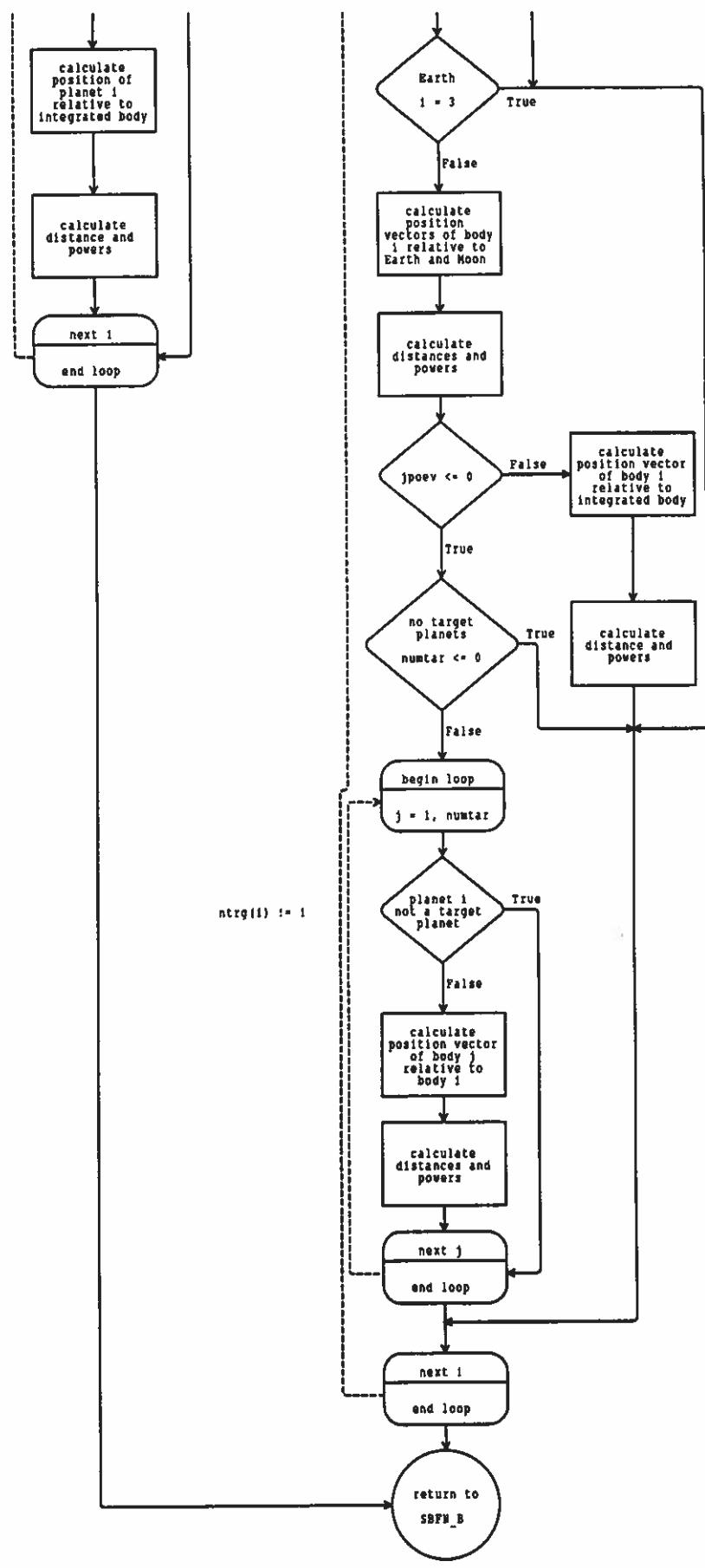


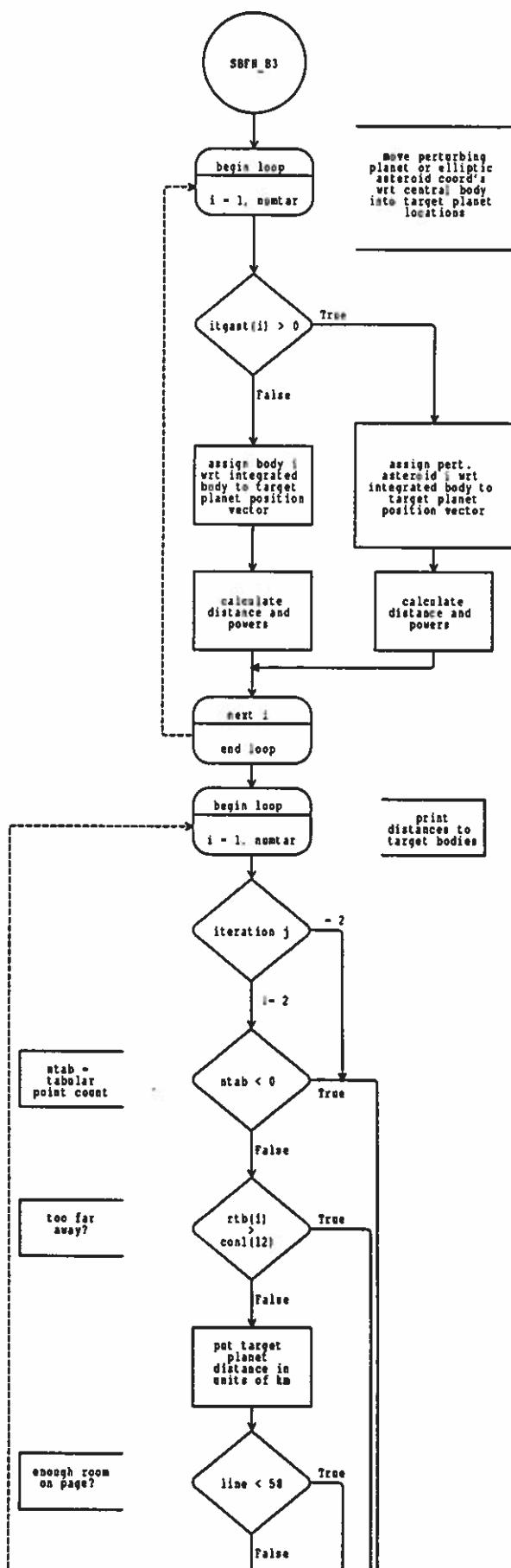


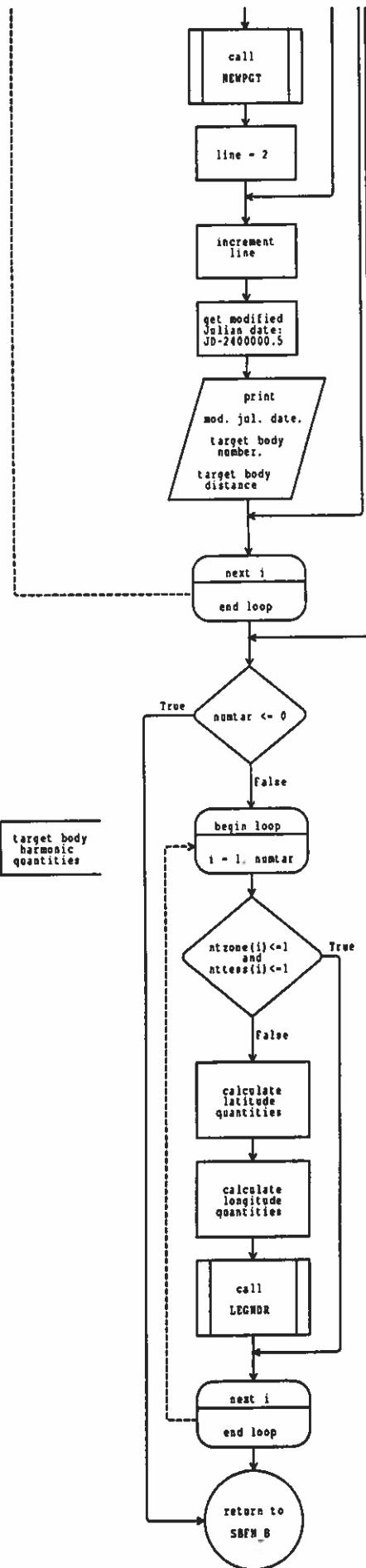


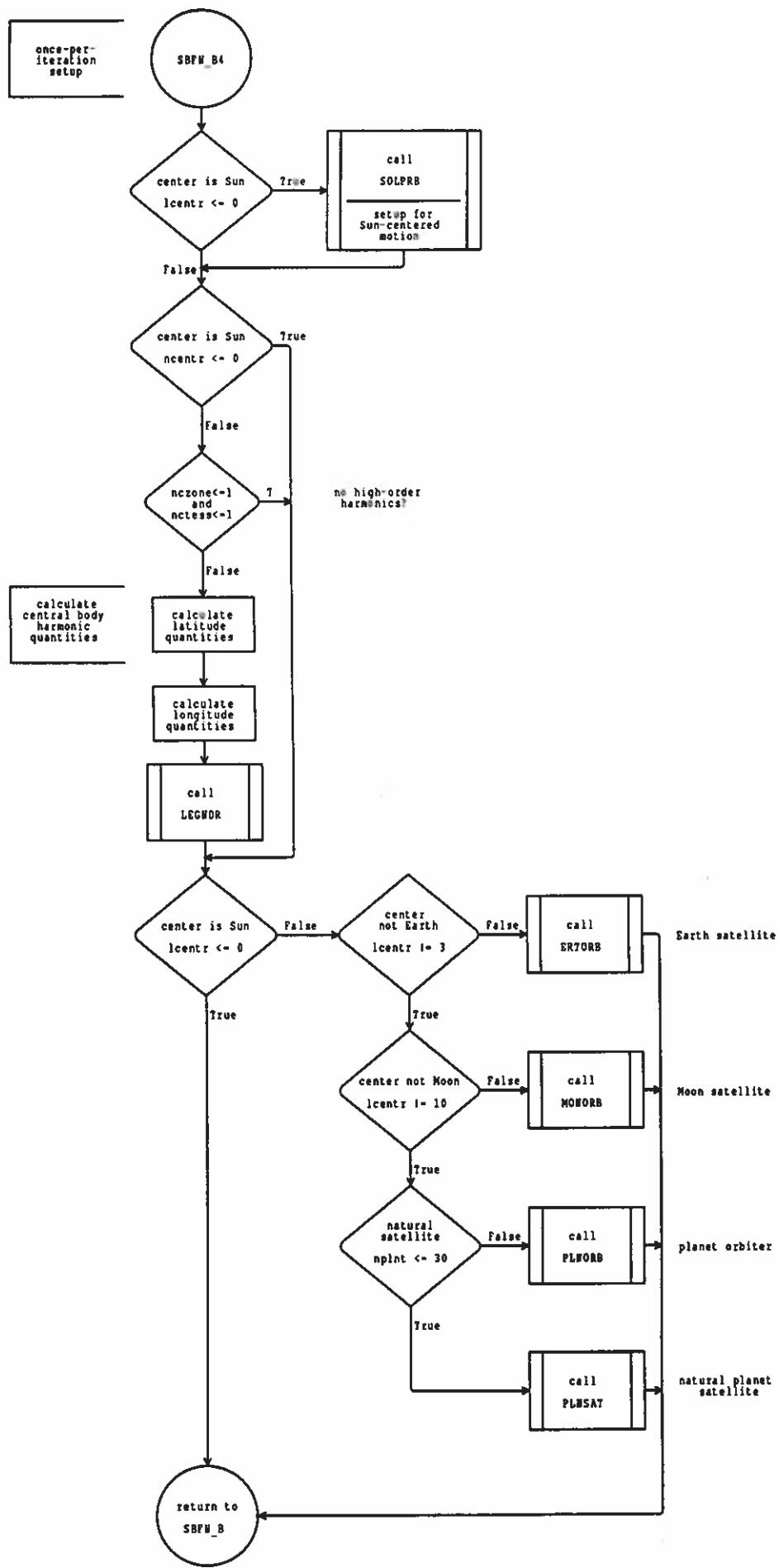


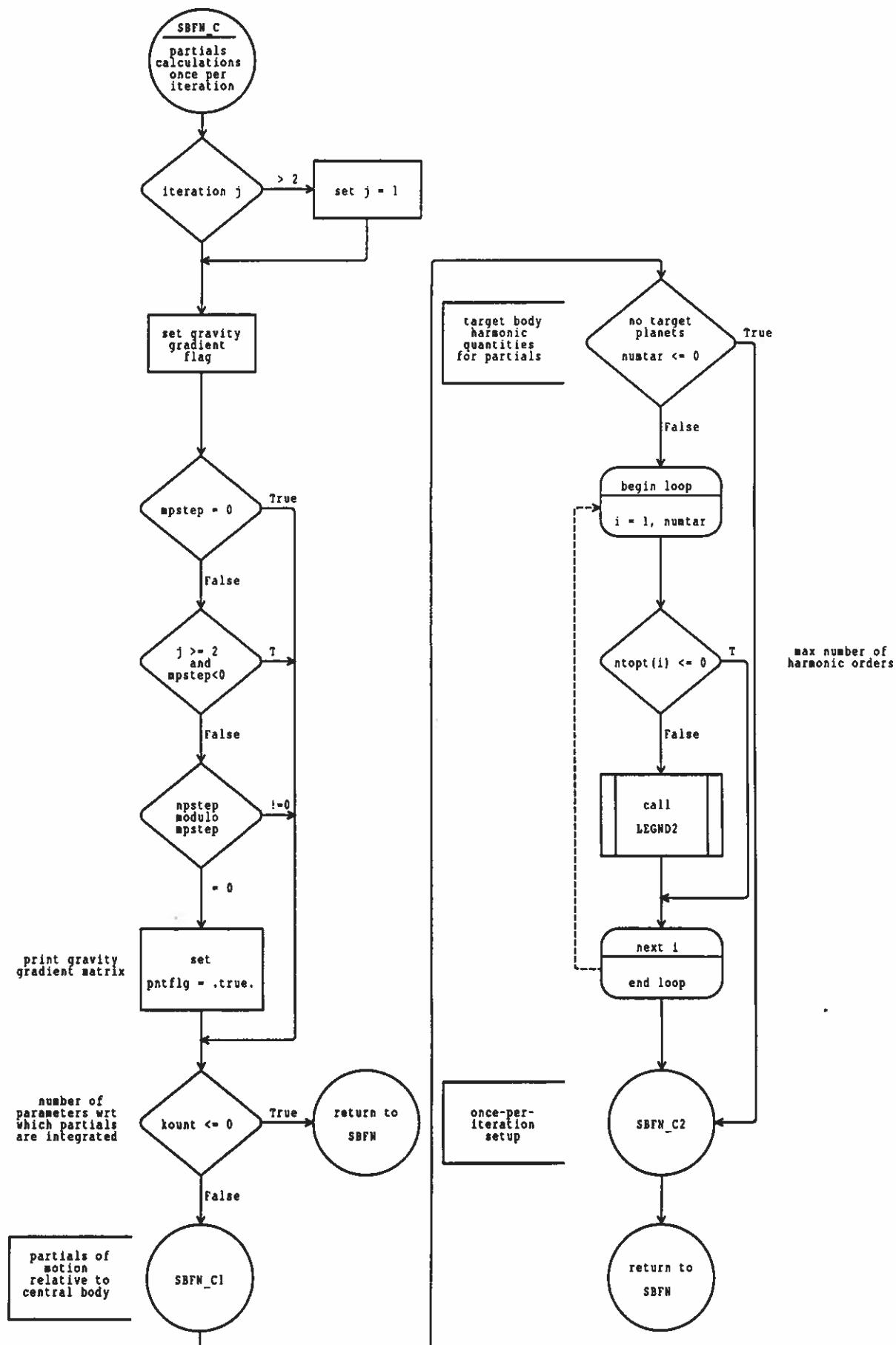


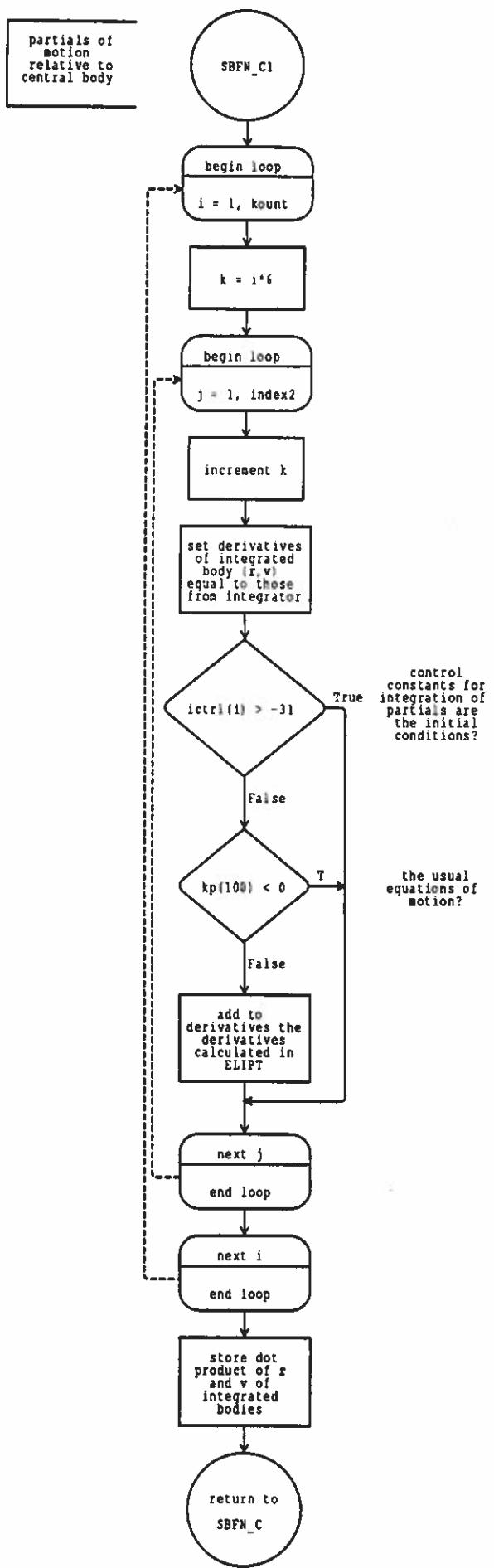


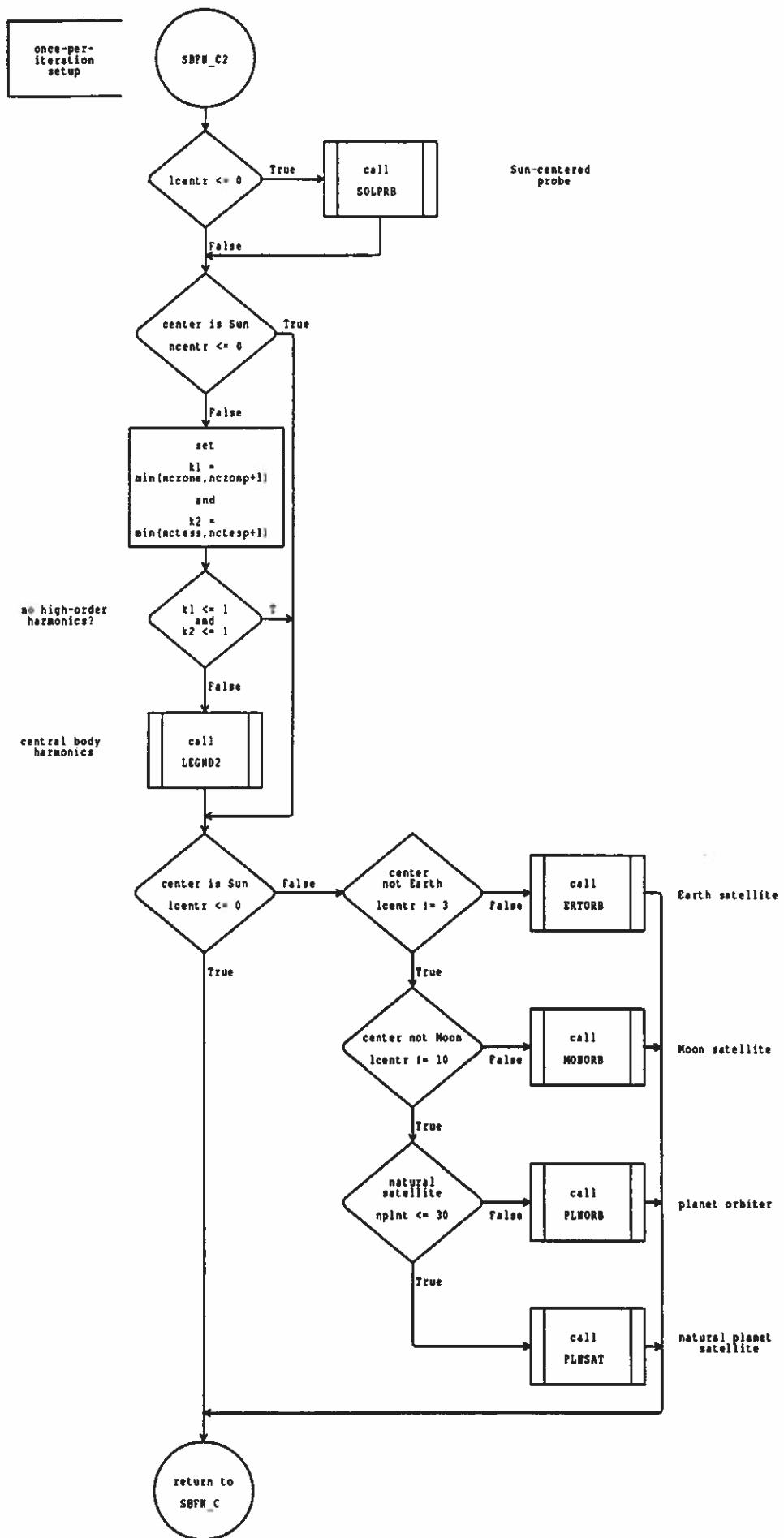


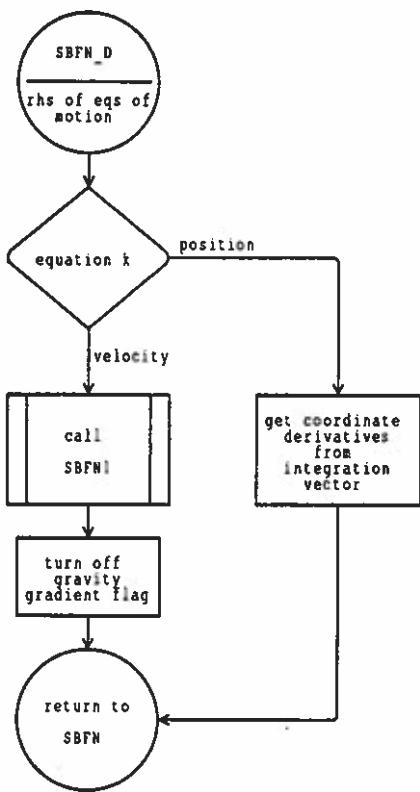












**SBFN1**

**subroutine SBFN1 ( k )**

Description:      SBFN1 evaluates the equations of motion and partial derivatives during numerical integrations. This routine is the natural extension of the setup computations performed in SBFN.

Arguments:      k      Equation number. There are six equations for the motion and six for each partial derivative. A detailed description of k is provided elsewhere.

Calls:            DOT, ERTORB, MONORB, NEWPGT, PLNORB, PLNSAT,  
                  PRODCT, SOLPRB

Called by:        SBFN

Includes:        CNTHAR, DRGPRT, ELLIPS, EMMIPS, FCNTRL, FMMIPS,  
                  INCON, INODTA, NAMTIMQ, PARAM, PETUNA, RTPARS,  
                  SBEMBR, SBROT, SBSTUF, SBTHNG, TRGHAR, YVECTPLP

Equivalence:

dadxc(3,3)	/sbstuf/	SBSTUF
ggfg	/sbstuf/	SBSTUF
ggfgs(2)	/sbstuf/	SBSTUF

## **SBFN1 Equations**

Subroutine SBFN1 does most of the rhs calculations of the equations of motion, both for the position accelerations and for the partials "accelerations". The equations of motion, both in this routine and in PEP in general, are integrated in an inertial frame of reference. For output purposes, though, coordinates are expressed as differences, usually with respect to the position of the central body.

This document is in three parts. The first part describes the equations of motion integrated by PEP for integrations other than the N-body integration performed by BODFN. The effects included in the equations of motion fall into three categories. First, there is the acceleration due to the central body. Second, we have various perturbing accelerations, including perturbing planets, perturbing asteroids, harmonics of the gravitational potentials of the target and central bodies, and acceleration arising from violation of the principle of equivalence. Third, there are various accelerations whose nature depends on the particular environment through which the integrated body finds itself traveling; these are dealt with via calls to appropriate special-case routines. The order in which the following is presented is that found in the source code itself.

The second part of the document deals with the calculation of the "gravity gradient" matrices. The physical effects giving rise to the contributions to these matrices include the perturbing planets, perturbing asteroids, target and central body harmonics, and the Sun and the central body. These are explained in detail in section II.

Finally, section III presents the equations that describe the acceleration of the derivatives of the motion with respect to various parameters. PEP uses these derivatives to calculate the sensitivity of results to changes in the parameters. These parameters include the masses of the perturbing planets, masses of the perturbing asteroids, time variation of the gravitational constant, the principle of equivalence violation parameter, target and central body zonal and tesseral harmonic coefficients, target and central body masses, the masses of the target planet central bodies, the central body rotation parameters, parameters of the equations in the special-cases routines referred

to above, and the initial conditions of the central and target bodies.

In PEP, planetary masses are in units of the mass of the Sun. To emphasize this, the equations described here explicitly show a division by the Sun's mass. These mass ratios are the parameters actually used in PEP, and partial derivatives with respect to planet masses are actually with respect to the mass ratios. Similarly, the masses of satellites (including the Moon) are parameterized as a fraction of the system mass (planet + satellites); this also is explicitly shown in this document.

In several of the equations which follow, the term  $\left(1 + \frac{M_B}{M_C}\right)$  appears.  $M_C$  is the central mass, either the mass of the Sun or the mass of the system containing  $M_C$ ; and  $M_B$  is the mass of the integrated body. This term is actually a shorthand for a more general expression which reduces to this one in the usual case. The exception is the case of an integrated natural satellite, for which the general expression is  $\left(\frac{M_C - \sum M_{sat}}{M_C} + \frac{M_B}{M_C}\right)$

Throughout this document, the term "scaled acceleration" will refer to the actual acceleration divided by the (possibly time-varying) gravitational constant.

## I. Equations of Motion

### A. Perturbing Accelerations

#### *1. Perturbing Planets*

The first acceleration of the integrated body (planet) that is taken into account is that due to the other 8 planets. If the integrated body is the Earth-Moon barycenter, then, its scaled acceleration towards its center (the Sun) due to the other planets is

$$\tilde{a}_{CP} = \sum_{k=1}^9 \frac{M_k}{M_\odot} \left( \mu_E \frac{\vec{r}_{\oplus k}}{r_{\oplus k}^3} + \mu_M \frac{\vec{r}_{Mk}}{r_{Mk}^3} - \frac{\vec{r}_{Ck}}{r_{Ck}^3} \right)$$

where  $\vec{r}_{\alpha\beta} = \vec{r}_\beta - \vec{r}_\alpha$ ,  $\vec{r}_\alpha$  is in general the position vector of body  $\alpha$  with respect to the central body, and

$$\mu_E = \frac{M_\oplus}{M_\oplus + M_M} \quad \mu_M = \frac{M_M}{M_\oplus + M_M}$$

For any other integrated body (planet) B, the scaled acceleration with respect to the central body due to all other planets is

$$\tilde{a}_{CP} = \sum_{\substack{k=1 \\ k \neq B}}^{10} \frac{M_k}{M_C} \left( \frac{\vec{r}_{Bk}}{r_{Bk}^3} - \frac{\vec{r}_{Ck}}{r_{Ck}^3} \right)$$

where body 10 is the Moon.

## 2. Principle of Equivalence Violation

The scaled acceleration of the integrated body B with respect to the central body C due to violation of the principle of equivalence for the planets is

$$\tilde{a}_{C\eta P} = \sum_{\substack{k=1 \\ k \neq B}}^{10} \frac{M_k}{M_\odot} \left[ (\eta \Delta_B + \eta' \Delta'_B) \frac{\vec{r}_{Bk}}{r_{Bk}^3} - (\eta \Delta_\odot + \eta' \Delta'_\odot) \frac{\vec{r}_{Ck}}{r_{Ck}^3} \right]$$

where  $\eta = 4\beta - \gamma - 3$  is the Nordtvedt parameter, and  $\Delta$  is the ratio of internal gravitational energy to total energy. The code is written as above, but it is possible that  $\Delta_\odot$  refers to the central body, which is not necessarily the Sun. The formalism incorporated by PEP includes only the PPN parameters  $\beta$  and  $\gamma$ . The parameter  $\gamma$  is a measure of the curvature induced per unit mass, and  $\beta$  measures the nonlinearity in the super-

position law for gravity. In GR, both  $\beta$  and  $\gamma$  are equal to one. The other parameters in the PPN formalism involve preferred-frame and nonconservation effects, and are zero in conservative theories of gravity, such as GR, with no preferred-frame effects. The primed quantities in the above equation are included in order to allow these effects in PEP; they are to be thought of as an additional contribution to  $\eta$ .

### *3. Perturbing Asteroids*

The scaled acceleration of the integrated body with respect to the central body due to perturbing asteroids is

$$\vec{a}_{C_{ast}} = \sum_{k=1}^{N_{ast}} \frac{M_k}{M_\odot} \left( \frac{\vec{r}_{Bk}}{r_{Bk}^3} - \frac{\vec{r}_{Ck}}{r_{Ck}^3} \right)$$

where  $N_{ast}$  is the number of perturbing asteroids taken into account.

The accelerations corresponding to  $a_{CP}$ ,  $a_{C\eta P}$ , and  $a_{Cast}$  are found by multiplying by

$$\gamma(t) = k_G^2 [1 + \alpha_G(t - t_0)]$$

where  $\alpha_G$  is the time variation factor for  $k_G$ , and  $k_G$  is the Gaussian gravitational constant. Thus,

$$\vec{F}_P + \vec{F}_\eta + \vec{F}_{ast} = \gamma(t) (\vec{a}_{CP} + \vec{a}_{C\eta P} + \vec{a}_{Cast})$$

### *4. Target Body Gravitational Harmonics*

Before illustrating the implementation of gravitational harmonics in PEP, it is worthwhile to develop the necessary formalism. Consider the potential of a spheroidal body of mass  $M$  and equatorial radius  $R_e$ . The point-mass potential plus the expansions in zonal and tesseral harmonics can be written in the standard form

$$\begin{aligned}
 U &= U_0 + U_{\text{zone}} + U_{\text{tess}} \\
 &= -\frac{GM}{r} \left[ 1 - \sum_{n=2}^{\infty} J_n \left( \frac{R_e}{r} \right)^n P_n(\sin\phi) \right. \\
 &\quad \left. + \sum_{n=2}^{\infty} \sum_{m=1}^n (C_{nm} \cos m\lambda + S_{nm} \sin m\lambda) \left( \frac{R_e}{r} \right)^n P_n^m(\sin\phi) \right]
 \end{aligned}$$

axisymmetric  
 non-axisymmetric

where  $\mathbf{r}$  is the position vector from the center of mass of M to the field point,  $J_n$  is the zonal harmonic coefficient,  $P_n(z)$  is the Legendre polynomial of degree  $n$  of  $z$ ,  $\phi$  is the latitude of the field point in the body coordinates,  $\lambda$  is the longitude in these same coordinates,  $P_n^m(z)$  is the associated Legendre function of degree  $n$  and order  $m$  of  $z$ , and G is the gravitational constant. The first sum is the expansion in zonal harmonics, representing an axisymmetric field. The double sum is the expansion in tesseral harmonics, allowing nonaxisymmetric modes. For the tesseral expansion, it is convenient (and standard practice) to split the coefficients into sine and cosine terms, as shown.

The body coordinates are such that the origin is at the center of mass and the z axis is coincident with the symmetry axis of M. Due to the location of the origin at the center of mass, the  $n=1$  term is zero. The acceleration due to this potential field is found by taking the gradient of  $U$ ,

$$\vec{F} = -\vec{\nabla}U = \vec{F}_0 + \vec{F}_{\text{zone}} + \vec{F}_{\text{tess}}$$

where

$$\begin{aligned}
 \vec{F}_0 &= -\frac{GM}{r^2} \frac{\vec{r}}{r} \\
 \vec{F}_{\text{zone}} &= +\frac{GM}{r^2} \sum_{n=2}^{\infty} J_n \left( \frac{R_e}{r} \right)^n \left[ (n+1) P_n(\sin\phi) \frac{\vec{r}}{r} - r \frac{\partial \sin\phi}{\partial \vec{r}} \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} \right] \\
 \vec{F}_{\text{tess}} &= -\frac{GM}{r^2} \sum_{n=2}^{\infty} \sum_{m=1}^n \left( \frac{R_e}{r} \right)^n (T_{nm} \vec{A}_{nm} + T'_{nm} \vec{B}_{nm})
 \end{aligned}$$

and

$$T_{nm} = C_{nm} \cos m\lambda + S_{nm} \sin m\lambda$$

$$T'_{nm} = -C_{nm} \sin m\lambda + S_{nm} \cos m\lambda$$

$$\vec{A}_{nm} = (n+1) P_n^m(\sin\phi) \frac{\vec{r}}{r} - r \frac{\partial \sin\phi}{\partial \vec{r}} \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi}$$

$$\vec{B}_{nm} = -m P_n^m(\sin\phi) r \frac{\partial \lambda}{\partial \vec{r}}$$

Now, if a different coordinate frame is being used than the local body frame just mentioned, an appropriate transformation must be made. Let  $\mathbf{u}$  be the body frame position vector. Then

$$\vec{u} = \mathbf{R} \vec{r}$$

where  $\mathbf{R}$  is the transformation matrix from the external coordinates  $\mathbf{r}$  to the body coordinates (In PEP,  $\mathbf{R}$  includes rotations from the equator and equinox of 1950.0, planet wobble [Earth only], planet nutation [Earth only], and precession). The sine of  $\phi$  is then

$$\sin\phi = \frac{u_3}{r} = \frac{1}{r} \sum_{i=1}^3 \mathbf{R}_{3i} x_i$$

and the components of the gradient required for the force equations above become

$$r \frac{\partial \sin\phi}{\partial x_i} = \mathbf{R}_{3i} - \frac{x_i}{r} \sin\phi$$

The components of the longitude gradient can be found as follows. Using  $\mathbf{R}$  to transform to body coordinates, we have

$$\sin\lambda = \frac{1}{r \cos\phi} \sum_{i=1}^3 R_{2i} x_i$$

$$\cos\lambda = \frac{1}{r \cos\phi} \sum_{i=1}^3 R_{1i} x_i$$

From the sine term, we find

$$r \frac{\partial \lambda}{\partial x_i} = \frac{1}{\cos\phi \cos\lambda} R_{2i} - \frac{x_i}{r} \tan\lambda - \frac{1}{\cos\phi} r \frac{\partial \cos\phi}{\partial x_i} \tan\lambda$$

From the cosine term,

$$r \frac{\partial \lambda}{\partial x_i} = -\frac{1}{\cos\phi \sin\lambda} R_{1i} + \frac{x_i}{r} \cot\lambda + \frac{1}{\cos\phi} r \frac{\partial \cos\phi}{\partial x_i} \cot\lambda$$

Combine these to get

$$r \frac{\partial \lambda}{\partial x_i} = \frac{1}{\cos\phi} (R_{2i} \cos\lambda - R_{1i} \sin\lambda)$$

which is the equation used by PEP.

Now we turn to the implementation of the force terms in PEP. The acceleration of the integrated body with respect to the central body due to gravitational harmonics of the target bodies is

$$\vec{F}_{TH} = \gamma(t) \sum_{k=1}^{N_{\text{Harm}}} \frac{M_k}{M_\odot} (\vec{a}_{CTZ}^k + \vec{a}_{CTR}^k)$$

where the acceleration terms are given below.

a) *Zonal Harmonics*

The acceleration of the integrated body with respect to the central body due to the zonal harmonics of target body k is coded in the form

$$\vec{a}_{CTZ}^k = \frac{1}{|\vec{r}_{Bk}|^2} \sum_{n=1}^{N_k-1} J_{n+1}^k \left( \frac{R_k}{|\vec{r}_{Bk}|} \right)^{n+1} \left[ -(n+2) P_{n+1}(\sin\phi_k) \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|} - |\vec{r}_{Bk}| \frac{\partial \sin\phi_k}{\partial (-\vec{r}_{Bk})} \frac{\partial P_{n+1}(\sin\phi_k)}{\partial z_k} \right]$$

where  $J_n^k$  is the zonal harmonic coefficient for the  $n^{\text{th}}$  zone of target body k, and  $\phi_k$  is the latitude of the integrated body as seen in the body-fixed coordinate frame of body k. Let  $\mathbf{u}$  be the position vector in this frame. The origin is at the center of mass, the  $u_3$  axis is aligned with the rotation axis, and the  $u_1$  axis points toward the intersection of the equator of target body k and the mean ecliptic of 1950.0.

### b) Tesselal Harmonics

As implemented in PEP, the acceleration of the integrated body with respect to the central body due to the target body tesseral harmonics is

$$\begin{aligned} \vec{a}_{CTT}^k = & \frac{1}{|\vec{r}_{Bk}|^2} \sum_{n=1}^{N_k-1} \sum_{m=1}^{n+1} \left[ A_{nm}^k (C_{nm}^k \cos m\lambda_k + S_{nm}^k \sin m\lambda_k) \right. \\ & \left. + B_{nm}^k (-C_{nm}^k \sin m\lambda_k + S_{nm}^k \cos m\lambda_k) \right] \end{aligned}$$

where  $S_{nm}^k$  is the sine coefficient and  $C_{nm}^k$  is the cosine coefficient for the tesseral harmonic of degree n and order m for body k, and where

$$A_{nm}^k = \left( \frac{R_k}{|\vec{r}_{Bk}|} \right)^{n+1} \left[ (n+2) P_{n+1}^m(\sin\phi_k) \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|} + |\vec{r}_{Bk}| \frac{\partial \sin\phi_k}{\partial (-\vec{r}_{Bk})} \frac{\partial P_{n+1}^m(\sin\phi_k)}{\partial \sin\phi_k} \right]$$

$$B_{nm}^k = \left( \frac{R_k}{|\vec{r}_{Bk}|} \right)^{n+1} m P_{n+1}^m(\sin\phi_k) |\vec{r}_{Bk}| \frac{\partial \lambda_k}{\partial (-\vec{r}_{Bk})}$$

which agrees with what we expect from  $F_{\text{tess}}$  above.

### 5. Sun

The acceleration due to the Sun for non-Sun-centered bodies is

$$\vec{F}_\odot = \gamma(t) \vec{a}_{C\odot} = \gamma(t) \left( \frac{\vec{r}_{OC}}{r_{OC}^3} - \frac{\vec{r}_{OB}}{r_{OB}^3} \right)$$

### 6. Central Body Gravitational Harmonics

The acceleration of the integrated body with respect to the central body due to gravitational harmonics of the central body itself is

$$\vec{F}_{CH} = \gamma(t) \frac{M_C}{M_\odot} \left( 1 + \frac{M_B}{M_\odot} \right) (\vec{a}_{CCz} + \vec{a}_{CCR})$$

where the acceleration terms are as follows. The equations have the same derivation as shown above for the target body harmonics.

#### a. Zonal Harmonics

The acceleration term of the integrated body with respect to the central body due to the zonal harmonics of the central body is implemented in PEP as

$$\vec{a}_{CCz} = \frac{1}{|\vec{r}_{CB}|^2} \sum_{n=1}^{N_c-1} J_n^C \left( \frac{R_C}{|\vec{r}_{CB}|} \right)^{n+1} \left[ (n+2) P_{n+1}(\sin\phi_C) \frac{\vec{r}_{CB}}{|\vec{r}_{CB}|} - |\vec{r}_{CB}| \frac{\partial \sin\phi_C}{\partial \vec{r}_{CB}} \frac{\partial P_{n+1}(\sin\phi_C)}{\partial \sin\phi_C} \right]$$

where  $J_n^C$  is the central body zonal harmonic coefficient for the  $n^{\text{th}}$  zone, and  $R_C$  is the equatorial radius of the central body.

### b. Tesserel Harmonics

The acceleration of the integrated body with respect to the central body due to the central body tesseral harmonics is coded as

$$\vec{a}_{CCT} = \frac{1}{|\vec{r}_{CB}|^2} \sum_{n=1}^{N_r-1} \sum_{m=1}^{n+1} [ A_{nm}^C ( C_{nm}^C \cos m\lambda_C + S_{nm}^C \sin m\lambda_C ) + B_{nm}^C ( -C_{nm}^C \sin m\lambda_C + S_{nm}^C \cos m\lambda_C ) ]$$

where  $S_{nm}^C$  is the sine coefficient and  $C_{nm}^C$  is the cosine coefficient for the tesseral harmonic of degree n and order m, and where

$$A_{nm}^C = \left( \frac{R_C}{|\vec{r}_{CB}|} \right)^{n+1} \left[ - (n+2) P_{n+1}^m(\sin\phi_C) \frac{\vec{r}_{CB}}{|\vec{r}_{CB}|} + |\vec{r}_{CB}| \frac{\partial \sin\phi_C}{\partial \vec{r}_{CB}} \frac{\partial P_{n+1}^m(\sin\phi_C)}{\partial \sin\phi_C} \right]$$

$$B_{nm}^C = \left( \frac{R_C}{|\vec{r}_{CB}|} \right)^{n+1} m P_{n+1}^m(\sin\phi_C) |\vec{r}_{CB}| \frac{\partial \lambda_C}{\partial \vec{r}_{CB}}$$

## B. Special-Case Accelerations

Accelerations specific to a

Sun-centered probe

Earth satellite

Lunar orbiter

planetary orbiter

natural planetary satellite

are evaluated via calls to subroutines SOLPRB, ERTORB, MONORB, PLNORB, and PLNSAT.

### C. Central Force

#### 1. Earth-Moon Barycenter

The acceleration, due to the Sun, of the Earth-Moon barycenter with respect to the Sun is

$$\bar{F}_c = \gamma(t) \bar{a}_{cc} = -\gamma(t) \left( 1 + \frac{M_\oplus + M_M}{M_\odot} \right) \left( \mu_E \frac{\vec{r}_{C\oplus}}{r_{C\oplus}^3} + \mu_M \frac{\vec{r}_{CM}}{r_{CM}^3} \right)$$

#### 2. All Other Bodies

For any body but the Earth-Moon barycenter as integrated body, the force with respect to the central body due to the central body is

$$\bar{F}_c = \gamma(t) \frac{M_c}{M_\odot} \left( 1 + \frac{M_B}{M_c} \right) \bar{a}_{cc} = -\gamma(t) \frac{M_c}{M_\odot} \left( 1 + \frac{M_B}{M_c} \right) \frac{\vec{r}_{CB}}{r_{CB}^3}$$

#### 3. Subtraction of Mean Orbit Acceleration

Currently not implemented.

#### 4. Subtraction of Elliptic Orbit Acceleration

The scaled acceleration to be added is that due to the central body of an elliptical orbit,

$$\bar{a}_e = -\frac{\vec{r}_e}{r_e^3}$$

where  $r_e$  is the Encke orbit position with respect to the central body. The acceleration is then

$$\vec{F}_e = \gamma(t) \frac{M_C}{M_\odot} \left( 1 + \frac{M_B}{M_C} \right) \vec{a}_{Ce}$$

The Encke integration was added to PEP in order to provide more accuracy.

### *5. Violation of Principle of Equivalence*

The acceleration with respect to the central body due to violation of the equivalence principle is

$$\vec{F}_{\eta C} = \gamma(t) \vec{a}_{C\eta} = -\gamma(t) \left[ (\eta \Delta_B + \eta' \Delta'_B) + \left( \eta \Delta_\odot + \frac{M_B}{M_\odot} \eta' \Delta'_\odot \right) \right] \cdot \vec{a}_{CC}$$

## II. Partial Derivatives--Gravity Gradient Matrices

The derivatives of the motion (position) with respect to parameters  $\alpha$  can be found by integrating the derivatives of the acceleration along with the equations of motion. Thus, for a position vector  $\mathbf{x}$ ,

$$\frac{d\mathbf{x}}{d\alpha} = \int \int \frac{d\ddot{\mathbf{x}}}{dt} dt dt$$

The integrand is

$$\frac{d\ddot{\mathbf{x}}}{dt} = \frac{\partial \ddot{\mathbf{x}}}{\partial \mathbf{x}} \cdot \frac{\partial \mathbf{x}}{\partial \alpha} + \frac{\partial \ddot{\mathbf{x}}}{\partial \dot{\mathbf{x}}} \cdot \frac{\partial \dot{\mathbf{x}}}{\partial \alpha} + \frac{\partial \ddot{\mathbf{x}}}{\partial \alpha}$$

Or, in component form,

$$\frac{d\ddot{x}_j}{d\alpha} = \frac{\partial \ddot{x}_j}{\partial x_k} \frac{\partial x_k}{\partial \alpha} + \frac{\partial \ddot{x}_j}{\partial \dot{x}_k} \frac{\partial \dot{x}_k}{\partial \alpha} + \frac{\partial \ddot{x}_j}{\partial \alpha}$$

with implied summation over repeated index  $k$ . The third term is zero if  $\alpha$  is an initial condition, since the acceleration functions are independent of the initial conditions. If there is no velocity dependence in the acceleration functions, then the second term is also zero. We are then left with just the first term, the first part of which we will refer to as the gravity gradient matrix:

$$G_{jk} = \frac{\partial \ddot{x}_j}{\partial x_k}$$

Since  $G$  is the second derivative of the gravitational potential, which satisfies Laplace's equation, it is therefore a symmetric matrix.

In this section we list all the gravity gradient matrix contributions calculated by PEP. The gravity gradient matrix is composed of the sum of the individual contributions  $G^\alpha$ :

$$\mathbf{G} = \sum_{\alpha} \mathbf{G}^{\alpha}$$

$$\alpha \in \{P, ast, TZ, TT, \odot, CZ, CT, C\}$$

where

- P = perturbing planets
- ast = perturbing asteroids
- TZ = target body zonal harmonics
- TT = target body tesseral harmonics
- $\odot$  = Sun
- CZ = central body zonal harmonics
- CT = central body tesseral harmonics
- C = central body

### A. Perturbing Bodies

#### *1. Earth-Moon Barycenter*

The contribution to the gravitational gradient matrix for the Earth-Moon barycenter due to the other planets is

$$G_{jk}^P = -\gamma(t) \sum_{\lambda=1}^9 \frac{M_{\lambda}}{M_{\odot}} \left\{ \left[ \mu_E \frac{3x'_{\oplus\lambda} x_{\oplus\lambda}^k}{|\vec{r}_{\oplus\lambda}|^5} + \mu_M \frac{3x'_{M\lambda} x_{M\lambda}^k}{|\vec{r}_{M\lambda}|^5} \right] + \delta_{jk} \left[ \mu_E \frac{1}{|\vec{r}_{\oplus\lambda}|^3} + \mu_M \frac{1}{|\vec{r}_{M\lambda}|^3} \right] \right\}$$

where  $x_{\alpha\beta}^k$  is the  $k^{\text{th}}$  component of the position vector  $\vec{r}_{\alpha\beta}$ , and  $\gamma(t) = k_G^2 [1 + \alpha_G (t - t_0)]$

where  $\alpha_G$  is the time variation factor for  $k_G$ , and  $k_G$  is the Gaussian gravitational constant. In PEP, we have

$$k_G = 0.01720209895 \ (AU)^{\frac{3}{2}} \ (ephemeris\ day)^{-1} \ (M_{\odot})^{-\frac{1}{2}}$$

which is the standard value as specified by the IAU (1976) System of Astronomical Constants.

### 2. All Other Planets

For all planets other than the Earth-Moon system, we have, for integrated planet B,

$$G_{jk}^P = \gamma(t) \sum_{\substack{\lambda=1 \\ \lambda \neq B}}^{10} \frac{M_{\lambda}}{M_{\odot}} \left\{ \frac{3x_{B\lambda}^j x_{B\lambda}^k}{|\vec{r}_{B\lambda}|^5} - \delta_{jk} \frac{1}{|\vec{r}_{B\lambda}|^3} \right\}$$

### 3. Perturbing Asteroids

The contribution due to perturbing asteroids is, for planet B,

$$G_{jk}^{ast} = \gamma(t) \sum_{\lambda=1}^{N_{ast}} \frac{M_{\lambda}}{M_{\odot}} \left\{ \frac{3x_{B\lambda}^j x_{B\lambda}^k}{|\vec{r}_{B\lambda}|^5} - \frac{\delta_{jk}}{|\vec{r}_{B\lambda}|^3} \right\}$$

## B. Target Bodies

Before showing the PEP implementation of the gravity gradient matrices for the target and central body harmonics, we briefly review the formalism. Remember the acceleration vectors due to a spheroid,

$$\begin{aligned}\vec{F}_0 &= -\frac{GM}{r^2} \frac{\vec{r}}{r} \\ \vec{F}_{zone} &= +\frac{GM}{r^2} \sum_{n=2}^{\infty} J_n \left( \frac{R_e}{r} \right)^n \left[ (n+1) P_n(\sin\phi) \frac{\vec{r}}{r} - r \frac{\partial \sin\phi}{\partial \vec{r}} \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} \right] \\ \vec{F}_{tess} &= -\frac{GM}{r^2} \sum_{n=2}^{\infty} \sum_{m=1}^n \left( \frac{R_e}{r} \right)^n \left( T_{nm} \vec{A}_{nm} + T'_{nm} \vec{B}_{nm} \right)\end{aligned}$$

where

$$\begin{aligned}T_{nm} &= C_{nm} \cos m\lambda + S_{nm} \sin m\lambda \\ T'_{nm} &= -C_{nm} \sin m\lambda + S_{nm} \cos m\lambda \\ \vec{A}_{nm} &= (n+1) P_n''(\sin\phi) \frac{\vec{r}}{r} - r \frac{\partial \sin\phi}{\partial \vec{r}} \frac{\partial P_n''(\sin\phi)}{\partial \sin\phi} \\ \vec{B}_{nm} &= -m P_n''(\sin\phi) r \frac{\partial \lambda}{\partial \vec{r}}\end{aligned}$$

and  $\mathbf{r}$  is the position vector referenced from the center of mass of the spheroid.  
The gradient of the spherical contribution produces the matrix

$$G_{lk}^0 = \frac{GM}{r^3} \left( \frac{3x_l x_k}{r^2} - \delta_{lk} \right)$$

The gradient of the zonal harmonic term results in

$$\begin{aligned}
G_{ik}^{zone} = & \frac{GM}{r^3} \sum_{n=2}^{\infty} J_n \left( \frac{R_e}{r} \right)^n \left\{ \left[ \delta_{ik} - (n+3) \frac{x_i x_k}{r^2} \right] (n+1) P_n(\sin\phi) \right. \\
& + \left( \frac{\partial \sin\phi}{\partial x_i} x_k + \frac{\partial \sin\phi}{\partial x_k} x_i \right) (n+1) \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} \\
& \left. - r^2 \left( \frac{\partial \sin\phi}{\partial x_i} \frac{\partial \sin\phi}{\partial x_k} \frac{\partial^2 P_n(\sin\phi)}{(\partial \sin\phi)^2} + \frac{\partial^2 \sin\phi}{\partial x_i \partial x_k} \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} \right) \right\}
\end{aligned}$$

The gradient of the tesseral harmonic term is

$$\begin{aligned}
G_{ik}^{tess} = & - \frac{GM}{r^3} \sum_{n=2}^{\infty} \sum_{m=1}^n \left( \frac{R_e}{r} \right)^n \left\{ T_{nm} \left[ \delta_{ik} - (n+3) \frac{x_i x_k}{r^2} \right] (n+1) P_n^m(\sin\phi) \right. \\
& + T_{nm} \left( \frac{\partial \sin\phi}{\partial x_i} x_k + \frac{\partial \sin\phi}{\partial x_k} x_i \right) (n+1) \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} \\
& + m T'_{nm} \left( \frac{\partial \lambda}{\partial x_i} x_k + \frac{\partial \lambda}{\partial x_k} x_i \right) (n+1) P_n^m(\sin\phi) \\
& - T_{nm} r^2 \left( \frac{\partial \sin\phi}{\partial x_i} \frac{\partial \sin\phi}{\partial x_k} \frac{\partial^2 P_n^m(\sin\phi)}{(\partial \sin\phi)^2} + \frac{\partial^2 \sin\phi}{\partial x_i \partial x_k} \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} \right) - m^2 \frac{\partial \lambda}{\partial x_i} \frac{\partial \lambda}{\partial x_k} P_n^m(\sin\phi) \\
& \left. - m T'_{nm} r^2 \left[ \left( \frac{\partial \sin\phi}{\partial x_k} \frac{\partial \lambda}{\partial x_i} + \frac{\partial \sin\phi}{\partial x_i} \frac{\partial \lambda}{\partial x_k} \right) \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 \lambda}{\partial x_i \partial x_k} P_n^m(\sin\phi) \right] \right\}
\end{aligned}$$

Using the coordinate transformation matrix, we find that the second derivatives of the latitude and longitude can be written

$$\begin{aligned} r^2 \frac{\partial^2 \sin\phi}{\partial x_i \partial x_k} &= \left( \frac{x_i x_k}{r^2} - \delta_{ik} \right) \sin\phi - \left( R_{3k} \frac{x_i}{r} + R_{3i} \frac{x_k}{r} \right) \\ r^2 \frac{\partial^2 \lambda}{\partial x_i \partial x_k} &= -\frac{1}{\cos^2 \phi} [ (R_{2i} \cos\lambda - R_{1i} \sin\lambda)(R_{2k} \sin\lambda + R_{1k} \cos\lambda) \\ &\quad + (R_{2k} \cos\lambda - R_{1k} \sin\lambda)(R_{2i} \sin\lambda + R_{1i} \cos\lambda) ] \end{aligned}$$

Insertion into the expressions for the zonal and tesseral components of the gravity gradient matrices results in the expressions

$$\begin{aligned} G_{ik}^{\text{zone}} &= \frac{GM}{r^3} \sum_{n=2} J_n \left( \frac{R_e}{r} \right)^n \left\{ \delta_{ik} \left[ (n+1) P_n(\sin\phi) + \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} \sin\phi \right] \right. \\ &\quad - \frac{x_i x_k}{r^2} \left[ (n+1)(n+2) P_n(\sin\phi) + (2n+5) \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} \sin\phi + \frac{\partial^2 P_n(\sin\phi)}{(\partial \sin\phi)^2} \sin^2 \phi \right] \\ &\quad + \left( R_{3i} \frac{x_k}{r} + R_{3k} \frac{x_i}{r} \right) \left[ (n+2) \frac{\partial P_n(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 P_n(\sin\phi)}{(\partial \sin\phi)^2} \sin\phi \right] \\ &\quad \left. - R_{3i} R_{3k} \frac{\partial^2 P_n(\sin\phi)}{(\partial \sin\phi)^2} \right\} \end{aligned}$$

and

$$\begin{aligned}
G_{ik}^{tess} = & - \frac{GM}{r^3} \sum_{n=2}^{\infty} \sum_{m=1}^n \left( \frac{R_e}{r} \right)^n \left\{ T_{nm} \delta_{ik} \left[ (n+1) P_n^m(\sin\phi) + \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} \sin\phi \right] \right. \\
& - T_{nm} \frac{x_i x_k}{r^2} \left[ (n+1)(n+3) P_n^m(\sin\phi) + (2n+5) \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} \sin\phi + \frac{\partial^2 P_n^m(\sin\phi)}{(\partial \sin\phi)^2} \sin^2\phi \right] \\
& + T_{nm} \left( \mathbf{R}_{3i} \frac{x_k}{r} + \mathbf{R}_{3k} \frac{x_i}{r} \right) \left[ (n+2) \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 P_n^m(\sin\phi)}{(\partial \sin\phi)^2} \sin\phi \right] \\
& - T_{nm} \mathbf{R}_{3i} \mathbf{R}_{3k} \frac{\partial^2 P_n^m(\sin\phi)}{(\partial \sin\phi)^2} \\
& + T_{nm} \frac{m^2}{\cos^2\phi} (\mathbf{R}_{2i} \cos\lambda - \mathbf{R}_{1i} \sin\lambda) (\mathbf{R}_{2k} \cos\lambda - \mathbf{R}_{1k} \sin\lambda) P_n^m(\sin\phi) \\
& + T'_{nm} \frac{m}{\cos\phi} (\mathbf{R}_{2i} \cos\lambda - \mathbf{R}_{1i} \sin\lambda) \left[ \frac{x_k}{r} (n+1) P_n^m(\sin\phi) - \left( \mathbf{R}_{3k} - \frac{x_k}{r} \sin\phi \right) \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} \right] \\
& + T'_{nm} \frac{m}{\cos\phi} (\mathbf{R}_{2k} \cos\lambda - \mathbf{R}_{1k} \sin\lambda) \left[ \frac{x_i}{r} (n+1) P_n^m(\sin\phi) - \left( \mathbf{R}_{3i} - \frac{x_i}{r} \sin\phi \right) \frac{\partial P_n^m(\sin\phi)}{\partial \sin\phi} \right] \\
& + T'_{nm} \frac{m}{\cos^2\phi} [(\mathbf{R}_{2i} \cos\lambda - \mathbf{R}_{1i} \sin\lambda) (\mathbf{R}_{2k} \sin\lambda + \mathbf{R}_{1k} \cos\lambda) \\
& \quad + (\mathbf{R}_{2k} \cos\lambda - \mathbf{R}_{1k} \sin\lambda) (\mathbf{R}_{2i} \sin\lambda + \mathbf{R}_{1i} \cos\lambda)] P_n^m(\sin\phi) \quad \left. \right\}
\end{aligned}$$

### 1. Zonal Harmonics

In PEP, the matrix components from the zonal gravitational harmonics of target body  $\alpha$  are coded as

$$\begin{aligned}
 G_{ik}^{TZ(\alpha)} = & \gamma(t) \frac{M_\alpha}{M_\odot} \frac{1}{r^3} \sum_{n=1}^{N_c-1} J_n \left( \frac{R_\alpha}{r} \right)^{n+1} \left\{ \delta_{ik} \left[ (n+2) P_{n+1}(\sin\phi) + \frac{\partial P_{n+1}(\sin\phi)}{\partial \sin\phi} \sin\phi \right] \right. \\
 & + \frac{x_i x_k}{r^2} \left[ (n+2)(n+4) P_{n+1}(\sin\phi) + (2n+7) \frac{\partial P_{n+1}(\sin\phi)}{\partial \sin\phi} \sin\phi + \frac{\partial^2 P_{n+1}(\sin\phi)}{(\partial \sin\phi)^2} \sin^2\phi \right] \\
 & + \left( R_{3i} \frac{x_k}{r} + R_{3k} \frac{x_i}{r} \right) \left[ (n+3) \frac{\partial P_{n+1}(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 P_{n+1}(\sin\phi)}{(\partial \sin\phi)^2} \sin\phi \right] \\
 & \left. + R_{3i} R_{3k} \frac{\partial^2 P_{n+1}(\sin\phi)}{(\partial \sin\phi)^2} \right\}
 \end{aligned}$$

where, in terms of previous notation,

$$r \equiv |\vec{r}_{B\alpha}| \quad x_i \equiv x_{B\alpha}^i$$

## 2. Tesserel Harmonics

The tesserel harmonics for target body  $\alpha$  lead to the gravity gradient matrix components

$$\begin{aligned}
 G_{ik}^{tess} = & -\gamma(t) \frac{M_\alpha}{M_\odot} \frac{1}{r^3} \sum_{n=1}^{\infty} \sum_{m=1}^n \left( \frac{R_\alpha}{r} \right)^n \left\{ T_{nm} \delta_{ik} \left[ (n+2) P_{n+1}^m(\sin\phi) + \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \sin\phi \right] \right. \\
 & - T_{nm} \frac{x_i x_k}{r^2} \left[ (n+2)(n+4) P_{n+1}^m(\sin\phi) + (2n+7) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \sin\phi + \frac{\partial^2 P_{n+1}^m(\sin\phi)}{(\partial \sin\phi)^2} \sin^2\phi \right] \\
 & - T_{nm} \left( R_{3i} \frac{x_k}{r} + R_{3k} \frac{x_i}{r} \right) \left[ (n+3) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 P_{n+1}^m(\sin\phi)}{(\partial \sin\phi)^2} \sin\phi \right] \\
 & - T_{nm} R_{3i} R_{3k} \frac{\partial^2 P_{n+1}^m(\sin\phi)}{(\partial \sin\phi)^2} \\
 & + T_{nm} \frac{m^2}{\cos^2\phi} (R_{2i} \cos\lambda - R_{1i} \sin\lambda) (R_{2k} \cos\lambda - R_{1k} \sin\lambda) P_{n+1}^m(\sin\phi) \\
 & - T'_{nm} \frac{m}{\cos\phi} (R_{2i} \cos\lambda - R_{1i} \sin\lambda) \left[ \frac{x_k}{r} (n+2) P_{n+1}^m(\sin\phi) + \left( R_{3k} + \frac{x_k}{r} \sin\phi \right) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \right] \\
 & - T'_{nm} \frac{m}{\cos\phi} (R_{2k} \cos\lambda - R_{1k} \sin\lambda) \left[ \frac{x_i}{r} (n+2) P_{n+1}^m(\sin\phi) + \left( R_{3i} + \frac{x_i}{r} \sin\phi \right) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \right] \\
 & + T'_{nm} \frac{m}{\cos^2\phi} [(R_{2i} \cos\lambda - R_{1i} \sin\lambda) (R_{2k} \sin\lambda + R_{1k} \cos\lambda) \\
 & \quad + (R_{2k} \cos\lambda - R_{1k} \sin\lambda) (R_{2i} \sin\lambda + R_{1i} \cos\lambda)] P_{n+1}^m(\sin\phi) \quad \left. \right\}
 \end{aligned}$$

### C. The Sun

The effect of the Sun on the partial derivatives for non-Sun-centered bodies is

$$G_{jk}^{\odot} = \gamma(t) \left\{ \frac{3x_{\odot B}^j x_{\odot B}^k}{|\vec{r}_{\odot B}|^5} - \delta_{jk} \frac{1}{|\vec{r}_{\odot B}|^3} \right\}$$

### D. Central Body

#### 1. Zonal Harmonics

The zonal harmonics for the central body result in the matrix components

$$\begin{aligned} G_{ik}^{cz} = & -\gamma(t) \frac{M_C}{M_C} \left( 1 + \frac{M_B}{M_\odot} \right) \frac{1}{r^3} \sum_{n=1}^{N_c-1} J_{n+1} \left( \frac{R_C}{r} \right)^{n+1} \left\{ \delta_{ik} \left[ (n+2) P_{n+1}(\sin\phi) + \frac{\partial P_{n+1}(\sin\phi)}{\partial \sin\phi} \sin\phi \right] \right. \\ & - \frac{x_i x_k}{r^2} \left[ (n+2)(n+3) P_{n+1}(\sin\phi) + (2n+7) \frac{\partial P_{n+1}(\sin\phi)}{\partial \sin\phi} \sin\phi + \frac{\partial^2 P_{n+1}(\sin\phi)}{(\partial \sin\phi)^2} \sin^2\phi \right] \\ & + \left( R_{3i} \frac{x_k}{r} + R_{3k} \frac{x_i}{r} \right) \left[ (n+3) \frac{\partial P_{n+1}(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 P_{n+1}(\sin\phi)}{(\partial \sin\phi)^2} \sin\phi \right] \\ & \left. - R_{3i} R_{3k} \frac{\partial^2 P_{n+1}(\sin\phi)}{(\partial \sin\phi)^2} \right\} \end{aligned}$$

where

$$r = |\vec{r}_{CB}| \quad x_i = x_{CB}^i$$

## 2. Tesseral Harmonics

For the central body tesseral harmonics, the gravity gradient matrix components are

$$\begin{aligned}
 G_{ik}^{tess} = & -\gamma(t) \frac{M_C}{M_C} \left(1 + \frac{M_B}{M_\odot}\right) \frac{1}{r^3} \sum_{n=1}^{\infty} \sum_{m=1}^n \left(\frac{R_C}{r}\right)^n \left\{ T_{nm} \delta_{ik} \left[ (n+2) P_{n+1}^m(\sin\phi) + \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \sin\phi \right] \right. \\
 & - T_{nm} \frac{x_i x_k}{r^2} \left[ (n+2)(n+4) P_{n+1}^m(\sin\phi) + (2n+7) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \sin\phi + \frac{\partial^2 P_{n+1}^m(\sin\phi)}{(\partial \sin\phi)^2} \sin^2\phi \right] \\
 & + T_{nm} \left( R_{3i} \frac{x_k}{r} + R_{3k} \frac{x_i}{r} \right) \left[ (n+3) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} + \frac{\partial^2 P_{n+1}^m(\sin\phi)}{(\partial \sin\phi)^2} \sin\phi \right] \\
 & - T_{nm} R_{3i} R_{3k} \frac{\partial^2 P_{n+1}^m(\sin\phi)}{(\partial \sin\phi)^2} \\
 & + T_{nm} \frac{m^2}{\cos^2\phi} (R_{2i} \cos\lambda - R_{1i} \sin\lambda) (R_{2k} \cos\lambda - R_{1k} \sin\lambda) P_{n+1}^m(\sin\phi) \\
 & + T'_{nm} \frac{m}{\cos\phi} (R_{2i} \cos\lambda - R_{1i} \sin\lambda) \left[ \frac{x_k}{r} (n+2) P_{n+1}^m(\sin\phi) - \left( R_{3k} - \frac{x_k}{r} \sin\phi \right) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \right] \\
 & + T'_{nm} \frac{m}{\cos\phi} (R_{2k} \cos\lambda - R_{1k} \sin\lambda) \left[ \frac{x_i}{r} (n+2) P_{n+1}^m(\sin\phi) - \left( R_{3i} - \frac{x_i}{r} \sin\phi \right) \frac{\partial P_{n+1}^m(\sin\phi)}{\partial \sin\phi} \right] \\
 & + T'_{nm} \frac{m}{\cos^2\phi} [(R_{2i} \cos\lambda - R_{1i} \sin\lambda) (R_{2k} \sin\lambda + R_{1k} \cos\lambda) \\
 & \quad + (R_{2k} \cos\lambda - R_{1k} \sin\lambda) (R_{2i} \sin\lambda + R_{1i} \cos\lambda)] P_{n+1}^m(\sin\phi) \quad \left. \right\}
 \end{aligned}$$

## E. Effect of the Central Force

### 1. Earth-Moon System

The effect of the force due to the central body on the partial derivatives for the Earth-Moon barycenter is

$$G_{jk}^c = \gamma(t) \frac{M_c}{M_\odot} \left( 1 + \frac{M_B}{M_c} \right) \left\{ \left[ \mu_E \frac{3x_{C\oplus}^j x_{C\oplus}^k}{|\vec{r}_{C\oplus}|^5} + \mu_M \frac{3x_{CM}^j x_{CM}^k}{|\vec{r}_{CM}|^5} \right] - \delta_{jk} \left[ \mu_E \frac{1}{|\vec{r}_{C\oplus}|^3} + \mu_M \frac{1}{|\vec{r}_{CM}|^3} \right] \right\}$$

### 2. All Other Planets

The effect of the force due to the central body on the partial derivatives for the other bodies is

$$G_{jk}^c = \gamma(t) \frac{M_c}{M_\odot} \left( 1 + \frac{M_B}{M_c} \right) \left\{ \frac{3x_{CB}^j x_{CB}^k}{|\vec{r}_{CB}|^5} - \frac{\delta_{jk}}{|\vec{r}_{CB}|^3} \right\}$$

### III. Partial Derivatives Accelerations

This section documents the accelerations arising from the time derivatives of the derivatives of the motion with respect to the parameters  $\alpha$ ,

$$\vec{f}_\alpha = \frac{d\vec{x}}{d\alpha} = \frac{d^2}{dt^2} \left( \frac{dx}{d\alpha} \right)$$

#### A. Perturbing Bodies

##### 1. $\alpha$ is the Mass of a Perturbing Planet

The acceleration of the derivative with respect to the mass of perturbing body P is as follows. If body P is the Earth and the central body C is not the Moon, the acceleration is

$$\vec{f}_{PP} = \gamma(t) \left[ \mu_E \left( \frac{\vec{r}_{B\oplus}}{|\vec{r}_{B\oplus}|^3} - \frac{\vec{r}_{C\oplus}}{|\vec{r}_{C\oplus}|^3} \right) + \mu_M \left( \frac{\vec{r}_{BM}}{|\vec{r}_{BM}|^3} - \frac{\vec{r}_{CM}}{|\vec{r}_{CM}|^3} \right) \right]$$

while if the central body is the Moon,

$$\vec{f}_{PP} = \gamma(t) \left[ \mu_E \left( \frac{\vec{r}_{B\oplus}}{|\vec{r}_{B\oplus}|^3} - \frac{\vec{r}_{C\oplus}}{|\vec{r}_{C\oplus}|^3} \right) - \mu_M \frac{\vec{r}_{BM}}{|\vec{r}_{BM}|^3} \right]$$

If the perturbing planet is the Moon, then if the central body is the Earth the acceleration is

$$\vec{f}_{PP} = \gamma(t) \frac{M_\oplus + M_M}{M_\odot} \left( \frac{\vec{r}_{BM}}{|\vec{r}_{BM}|^3} - \frac{\vec{r}_{EM}}{|\vec{r}_{EM}|^3} - \frac{\vec{r}_{B\oplus}}{|\vec{r}_{B\oplus}|^3} \right)$$

while if the central body is not the Earth it is

$$\vec{f}_{PP} = \gamma(t) \frac{M_{\oplus} + M_M}{M_{\odot}} \left( \frac{\vec{r}_{BM}}{|\vec{r}_{BM}|^3} - \frac{\vec{r}_{CM}}{|\vec{r}_{CM}|^3} - \frac{\vec{r}_{B\oplus}}{|\vec{r}_{B\oplus}|^3} + \frac{\vec{r}_{C\oplus}}{|\vec{r}_{C\oplus}|^3} \right)$$

For all other perturbing bodies,

$$\vec{f}_{PP} = \gamma(t) \left( \frac{\vec{r}_{BP}}{|\vec{r}_{BP}|^3} - \frac{\vec{r}_{CP}}{|\vec{r}_{CP}|^3} \right)$$

## 2. $\alpha$ is a System Mass

When the derivative is with respect to the mass of a system S which includes perturbing satellites, we have

$$\vec{f}_{PS} = \gamma(t) \sum_{k=1}^{Nast} \frac{M_k}{M_S} \left( \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|^3} - \frac{\vec{r}_{Ck}}{|\vec{r}_{Ck}|^3} \right)$$

## 3. $\alpha$ is an Asteroid Mass

For the derivative with respect to the mass of asteroid k,

$$\vec{f}_{Pa}^k = \gamma(t) \left( \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|^3} - \frac{\vec{r}_{Ck}}{|\vec{r}_{Ck}|^3} \right)$$

## 4. $\alpha$ is the Time Variation of G

When the derivative is with respect to time variation of the gravitational con-

stant, the force is

$$\vec{f}_{PG} = k_G^2 (t - t_0) \cdot (\vec{a}_{CP} + \vec{a}_{Cn\gamma P} + \vec{a}_{Cast})$$

### 5. $\alpha$ is the Principle of Equivalence Violation Parameter

If the derivative is with respect to the Solar ratio of internal gravitational to total energy, the force is

$$\vec{f}_{P_{n\gamma\odot}} = -\gamma(t) \sum_{k=1}^{10} \frac{M_k}{M_\odot} \frac{\vec{r}_{Ck}}{|\vec{r}_{Ck}|^3}$$

If the derivative is with respect to the integrated body ratio of internal gravitational to total energy, the force is

$$\vec{f}_{P_{n\gamma_B}} = +\gamma(t) \sum_{k=1}^{10} \frac{M_k}{M_\odot} \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|^3}$$

## B. Target Bodies

### 1. $\alpha$ is a Zonal Harmonic Coefficient

When the derivative is with respect to the zonal harmonic coefficient of degree n of target body k, we have the corresponding acceleration

$$\begin{aligned} \vec{f}_{TZ}^{k,n} &= \gamma(t) \frac{M_k}{M_\odot} \frac{1}{|\vec{r}_{Bk}|^2} \left( \frac{R_k}{|\vec{r}_{Bk}|} \right)^n \\ &\cdot \left[ -(n+1) P_n(\sin\phi_k) \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|} - |\vec{r}_{Bk}| \frac{\partial \sin\phi_k}{\partial (-\vec{r}_{Bk})} \frac{\partial P_n(\sin\phi_k)}{\partial \sin\phi_k} \right] \end{aligned}$$

where  $\phi_k$  is the latitude of the integrated body with respect to target body k.

### 2. $\alpha$ is a Tesselal Harmonic Cosine Coefficient

When the derivative is with respect to the tesselal harmonic cosine coefficient of degree n and order m of target body k, we have

$$\vec{f}_{TTC}^{k,n,m} = \gamma(t) \frac{M_k}{M_\odot |\vec{r}_{Bk}|^2} (A_{nm}^k \cos m\lambda_k - B_{nm}^k \sin m\lambda_k)$$

where

$$A_{nm}^k = \left( \frac{R_k}{|\vec{r}_{Bk}|} \right)^n \left[ (n+1) P_n^m(\sin\phi_k) \frac{\vec{r}_{Bk}}{|\vec{r}_{Bk}|} + |\vec{r}_{Bk}| \frac{\partial \sin\phi_k}{\partial (-\vec{r}_{Bk})} \frac{\partial P_n^m(\sin\phi_k)}{\partial \sin\phi_k} \right]$$

$$B_{nm}^k = \left( \frac{R_k}{|\vec{r}_{Bk}|} \right)^n m P_n^m(\sin\phi_k) |\vec{r}_{Bk}| \frac{\partial \lambda_k}{\partial (-\vec{r}_{Bk})}$$

### 3. $\alpha$ is a Tesselal Harmonic Sine Coefficient

When the derivative is with respect to the tesselal harmonic sine coefficient of degree n and order m of target body k,

$$\vec{f}_{TTS}^{k,n,m} = \gamma(t) \frac{M_k}{M_\odot |\vec{r}_{Bk}|^2} [A_{nm}^k \sin m\lambda_k + B_{nm}^k \cos m\lambda_k]$$

### 4. $\alpha$ is a Target Planet Mass

When the parameter is the mass of target planet k,

$$\vec{f}_{TM}^k = \gamma(t) \frac{M_k^{TC}}{M_\odot} (\bar{a}_{CTZ}^k + \bar{a}_{CTT}^k)$$

where the acceleration towards the central body due to target body k harmonics has been given previously in section I. Here,  $M_k^{TC}$  is the mass of the central body of target body k.

#### *5. $\alpha$ is the Mass of a Target Planet Central Body*

When the parameter is the mass of the central body of a target body,

$$\vec{f}_{tc_m} = \gamma(t) \sum_{k=1}^{N_{arg}} \frac{M_k^{ast}}{M_\odot} \vec{a}_{cth}^k$$

where  $M_k^{ast}$  is the mass of target satellite k.

#### *6. $\alpha$ is the Time Variation of the Gravitational Constant*

When the parameter is the time variation of G,

$$\vec{f}_{tG} = k_G^2 (t - t_0) \sum_{k=1}^{N_{arg}} \frac{M_k^{arg}}{M_\odot} \vec{a}_{cth}^k$$

### C. The Sun

#### *1. $\alpha$ is the Time Variation of the Gravitational Constant*

When the parameter is the time variation of G,

$$\vec{f}_{OG} = k_G^2 (t - t_0) \vec{a}_{CO}$$

### D. The Central Body

#### *1. $\alpha$ is a Zonal Harmonic Coefficient*

When  $\alpha$  is the zonal harmonic coefficient of degree n for the central body, we have

$$\vec{f}_{CZ} = \gamma(t) \frac{M_C}{M_\odot} \left(1 + \frac{M_B}{M_C}\right) \frac{1}{|\vec{r}_{CB}|^2} \left(\frac{R_C}{|\vec{r}_{CB}|}\right)^n \left[ (n+1) P_n(\sin\phi_C) \frac{\vec{r}_{CB}}{|\vec{r}_{CB}|} - |\vec{r}_{CB}| \frac{\partial \sin\phi_C}{\partial \vec{r}_{CB}} \frac{\partial P_n(\sin\phi_C)}{\partial \sin\phi_C} \right]$$

where  $R_C$  is the equatorial radius of the central body.

### 2. $\alpha$ is a Tesselal Harmonic Cosine Coefficient

When  $\alpha$  is the tesselal harmonic cosine coefficient of degree  $n$  and order  $m$  for the central body, we have

$$\vec{f}_{CTC} = \gamma(t) \frac{M_C}{M_\odot} \left(1 + \frac{M_B}{M_C}\right) \frac{1}{|\vec{r}_{CB}|^2} (A_{nm}^C \cos m\lambda_C - B_{nm}^C \sin m\lambda_C)$$

where

$$A_{nm}^C = \left(\frac{R_C}{|\vec{r}_{CB}|}\right)^n \left[ -(n+1) P_n^m(\sin\phi_C) \frac{\vec{r}_{CB}}{|\vec{r}_{CB}|} + |\vec{r}_{CB}| \frac{\partial \sin\phi_C}{\partial \vec{r}_{CB}} \frac{\partial P_n^m(\sin\phi_C)}{\partial \sin\phi_C} \right]$$

$$B_{nm}^C = \left(\frac{R_C}{|\vec{r}_{CB}|}\right)^n m P_n^m(\sin\phi_C) |\vec{r}_{CB}| \frac{\partial \lambda_C}{\partial \vec{r}_{CB}}$$

### 3. $\alpha$ is a Tesselal Harmonic Sine Coefficient

When  $\alpha$  is the tesselal harmonic sine coefficient of degree  $n$  and order  $m$  for the central body, we have

$$\vec{f}_{CTS} = \gamma(t) \frac{M_C}{M_\odot} \left(1 + \frac{M_B}{M_C}\right) \frac{1}{|\vec{r}_{CB}|^2} (A_{nm}^C \sin m\lambda_C + B_{nm}^C \cos m\lambda_C)$$

*4.  $\alpha$  is one of the Central Body Rotation Parameters*

When  $\alpha$  is one of the central body rotation parameters,

$$\vec{f}_{crot}^{\alpha} = -\gamma(t) \frac{M_C}{M_{\odot}} \left( 1 + \frac{M_B}{M_C} \right) \frac{\partial \mathbf{R}^C}{\partial \alpha} \cdot \vec{a}_{CCH} + (\mathbf{G}^{CH})^T \cdot \frac{\partial \mathbf{R}^C}{\partial \alpha} \cdot \vec{r}_{CB}$$

for rotation parameter  $\alpha$ . Here,  $\mathbf{R}^C$  is the rotation matrix that transforms from central body coordinates to the fixed frame with celestial equator of 1950.0;  $\mathbf{a}_{CCH}$  is the acceleration with respect to the central body due to central body gravitational harmonics; and

$$(\mathbf{G}^{Ch})^T \equiv (\mathbf{G}^{CT} + \mathbf{G}^{CZ})^T$$

where superscript T stands for matrix transpose.

E. Additional Accelerations

The effects on the partial derivatives due to accelerations specific to

Sun-centered probe

Earth satellite

Lunar orbiter

planetary orbiter

natural planetary satellite

are determined via calls to subroutines SOLPRB, ERTORB, MONORB, PLNORB, and PLNSAT, respectively.

F. Initial Conditions

*1.  $\alpha$  is an Initial Condition of a Target Planet*

When the parameter is an initial condition of target planet k, the partials acceleration is

$$\vec{f}_{TI}^k = \gamma(t) \left\{ \frac{M_k}{M_\odot} \left[ \frac{1}{|\vec{r}_{Bk}|^3} \left( \frac{\partial \vec{r}_k}{\partial \alpha_k} - 3\vec{r}_{Bk} \frac{A_{Bk}}{|\vec{r}_{Bk}|^2} \right) \right. \right. \\ \left. \left. - \frac{1}{|\vec{r}_{Ck}|^3} \left( \frac{\partial \vec{r}_k}{\partial \alpha_k} - 3\vec{r}_{Ck} \frac{A_{Ck}}{|\vec{r}_{Ck}|^2} \right) \right] - \frac{\partial \vec{r}_k}{\partial \alpha_k} \cdot (G_k^{TZ} + G_k^{TT}) \right\}$$

where

$$A_{Bk} \equiv \vec{r}_{Bk} \cdot \frac{\partial \vec{r}_k}{\partial \alpha_k} \quad A_{Ck} \equiv \vec{r}_{Ck} \cdot \frac{\partial \vec{r}_k}{\partial \alpha_k}$$

and  $\alpha_k$  is an initial condition (orbital element) of the target planet k.

## 2. *$\alpha$ is an Initial Condition of the Central Body*

When  $\alpha$  is an initial condition of the central body, the following effects contribute to the acceleration of the partial derivative of the central body position with respect to  $\alpha$ .

### a. *The Sun's Effect on Integrated Initial Condition Partial Derivative*

$$\vec{f}_{CI}^\odot = \frac{\partial \vec{r}_C}{\partial \alpha_C} \cdot G^\odot + \gamma(t) \frac{1}{|\vec{r}_C|^3} \left( \frac{\partial \vec{r}_C}{\partial \alpha_C} - 3 \frac{\vec{r}_C}{|\vec{r}_C|^2} \left( \vec{r}_C \cdot \frac{\partial \vec{r}_C}{\partial \alpha_C} \right) \right)$$

### b. *Effects of Perturbing Planets*

$$\vec{f}_{CI}^P = \frac{\partial \vec{r}_C}{\partial \alpha_C} \cdot G^P + \gamma(t) \sum_{k=1}^{10} \frac{M_k}{M_\odot} \frac{1}{|\vec{r}_{Ck}|^3} \left[ \frac{\partial \vec{r}_{Ck}}{\partial \alpha_C} - 3 \frac{\vec{r}_{Ck}}{|\vec{r}_{Ck}|^2} \left( \vec{r}_{Ck} \cdot \frac{\partial \vec{r}_{Ck}}{\partial \alpha_C} \right) \right]$$

*c. Effects of Perturbing Asteroids*

Currently not implemented.

*d. Effects of the Target Body Harmonics*

$$\vec{f}_{CI}^{TH} = \left[ \sum_{k=1}^{N_{arg}} (\mathbf{G}^{TZ} + \mathbf{G}^{TT}) \right] \cdot \frac{\partial \vec{r}_C}{\partial \alpha_C}$$

G. Effect of Central Forces on Partial Derivatives

*1. Earth-Moon Barycenter*

The gravity gradient matrix for the Earth-Moon system orbiting a central body is

$$G_{jk}^C = \gamma(t) \left( 1 + \frac{M_B}{M_\odot} \right) \left\{ \left[ \mu_E \frac{3x_{C\oplus}^j x_{C\oplus}^k}{|\vec{r}_{C\oplus}|^5} + \mu_M \frac{3x_{CM}^j x_{CM}^k}{|\vec{r}_{CM}|^5} \right] - \delta_{jk} \left[ \mu_E \frac{1}{|\vec{r}_{C\oplus}|^3} + \mu_M \frac{1}{|\vec{r}_{CM}|^3} \right] \right\}$$

*2. All Other Bodies*

The gravity gradient matrix for any other body is

$$G_{jk}^C = \gamma(t) \frac{M_C}{M_\odot} \left( 1 + \frac{M_B}{M_C} \right) \left[ \frac{3x_{CB}^j x_{CB}^k}{|\vec{r}_{CB}|^5} - \frac{\delta_{jk}}{|\vec{r}_{CB}|^3} \right]$$

In the special case of the Jovian system, the expression for the mass in the above equation is more complicated.

## H. Gravity Gradient Matrix

The total gravity gradient matrix consists of the sum of all the preceding contributions:

$$\mathbf{G} = \sum \mathbf{G}^\alpha$$

where

$$\alpha \in \{P, ast, TZ, TT, \odot, CZ, CT, C\}$$

as listed at the beginning of section II.

## I. Partial Derivative is with respect to an Initial Condition

### *1. Subtract Mean Orbit Acceleration Partial*

Currently not implemented.

### *2. Subtract Elliptic Orbit Acceleration Partial*

The Encke orbit force to be subtracted:

$$\vec{f}_{IE}^B = -\gamma(t) \frac{M_C}{M_\odot} \left(1 + \frac{M_B}{M_C}\right) \frac{1}{|\vec{r}_E|^3} \left[ 3 \frac{\vec{r}_E}{|\vec{r}_E|} \left( \vec{r}_E \cdot \frac{\partial \vec{r}_E}{\partial \alpha} \right) \right]$$

### *3. $\alpha$ is Time Variation of G*

If the parameter is the variation of the gravitational constant, then

$$\vec{f}_{IG} = -k_G^2 (t-t_0) \frac{M_C}{M_\odot} \left(1 + \frac{M_B}{M_C}\right) \vec{a}_{CC}$$

where  $\vec{a}_{CC}$  is the acceleration due to the central body, given in section I.C.

*4.  $\alpha$  is Mass of Another Satellite*

If the parameter is the mass of another satellite,

$$\vec{f}_{Isat} = \gamma(t) \frac{M_C}{M_\odot} \vec{a}_{cc}$$

*5.  $\alpha$  is Asteroid Mass*

If the parameter is the mass of an asteroid,

$$\vec{f}_{Iast} = -\gamma(t) \frac{M_C}{M_\odot} \vec{a}_{cc}$$

*6.  $\alpha$  is Mass of Central Body*

If the parameter is the mass of the central body, then

$$\vec{f}_{IC} = -\gamma(t) \left( 1 + \frac{M_B}{M_C} \right) \vec{a}_{cc}$$

*7.  $\alpha$  is the Principle of Equivalence Violation Parameter*

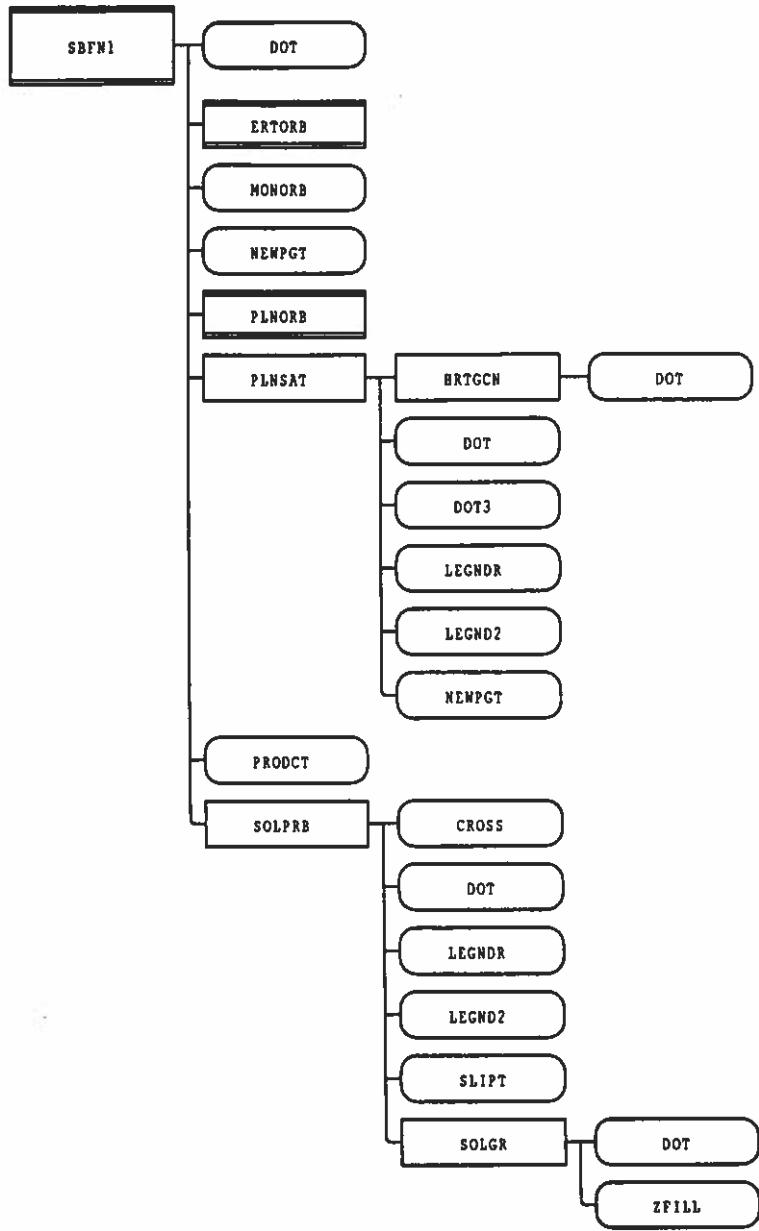
Finally, if  $\alpha$  is the principle of equivalence violation parameter,

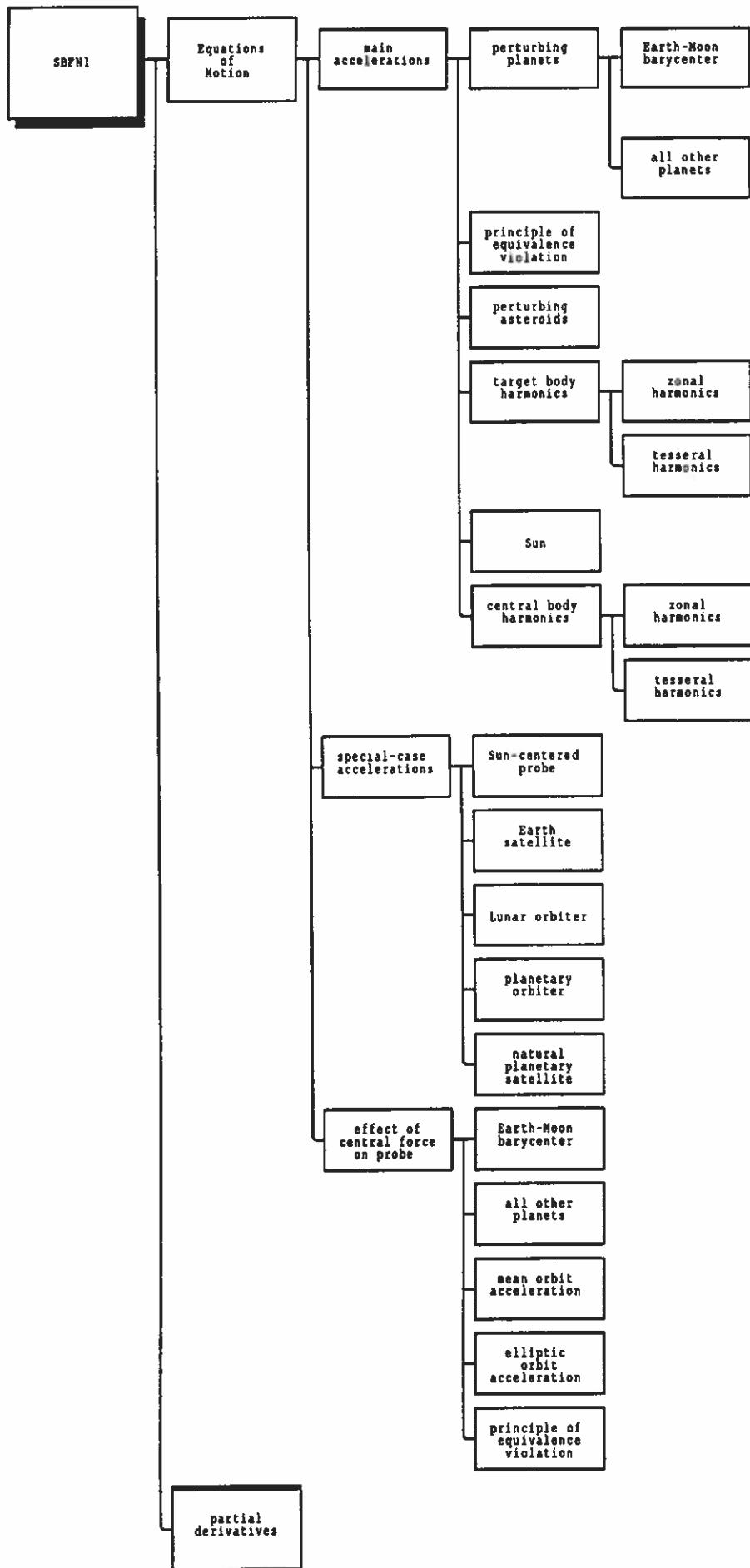
$$\vec{f}_{In} = -\gamma(t) \frac{M_B}{M_C} \vec{a}_{cc}$$

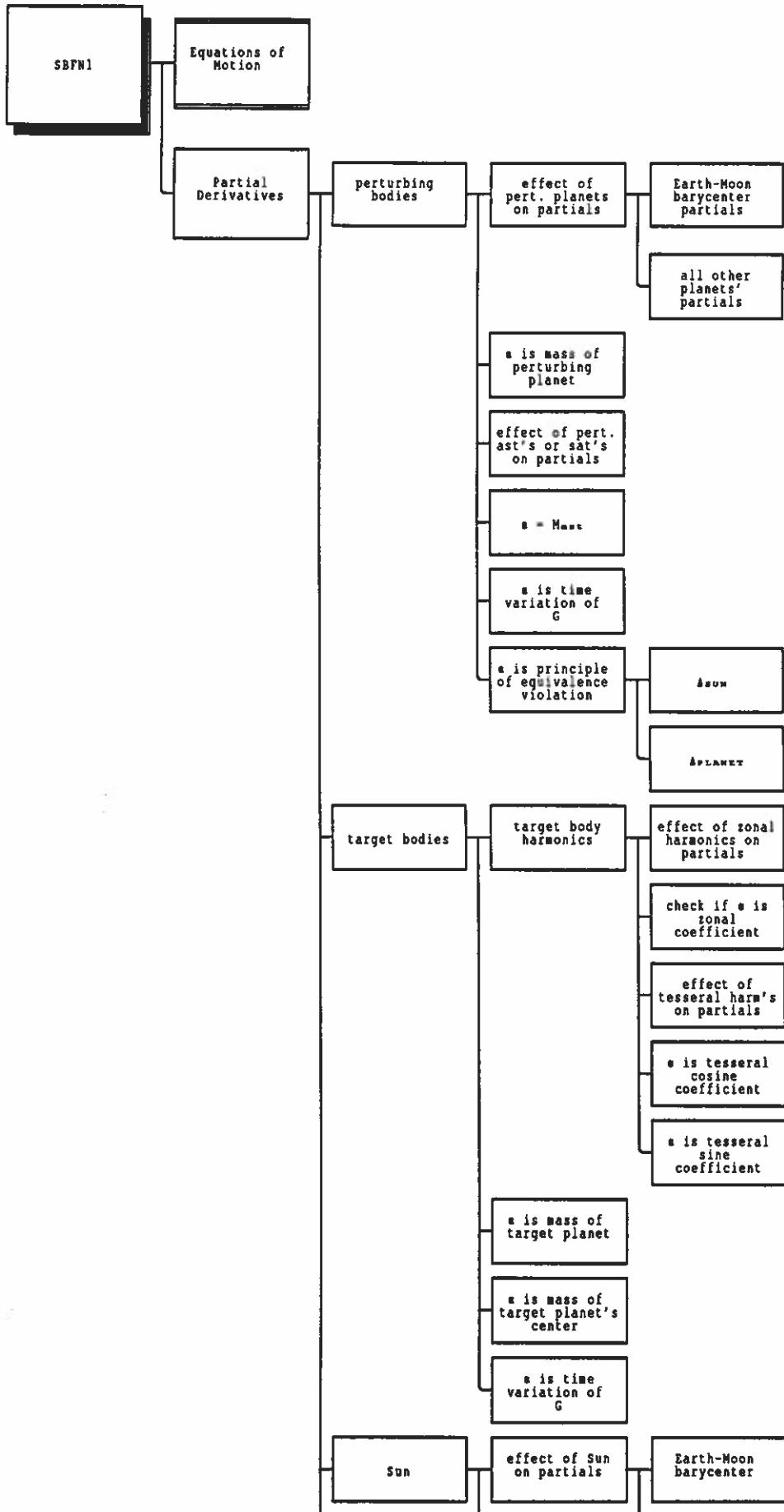
for the Sun, and

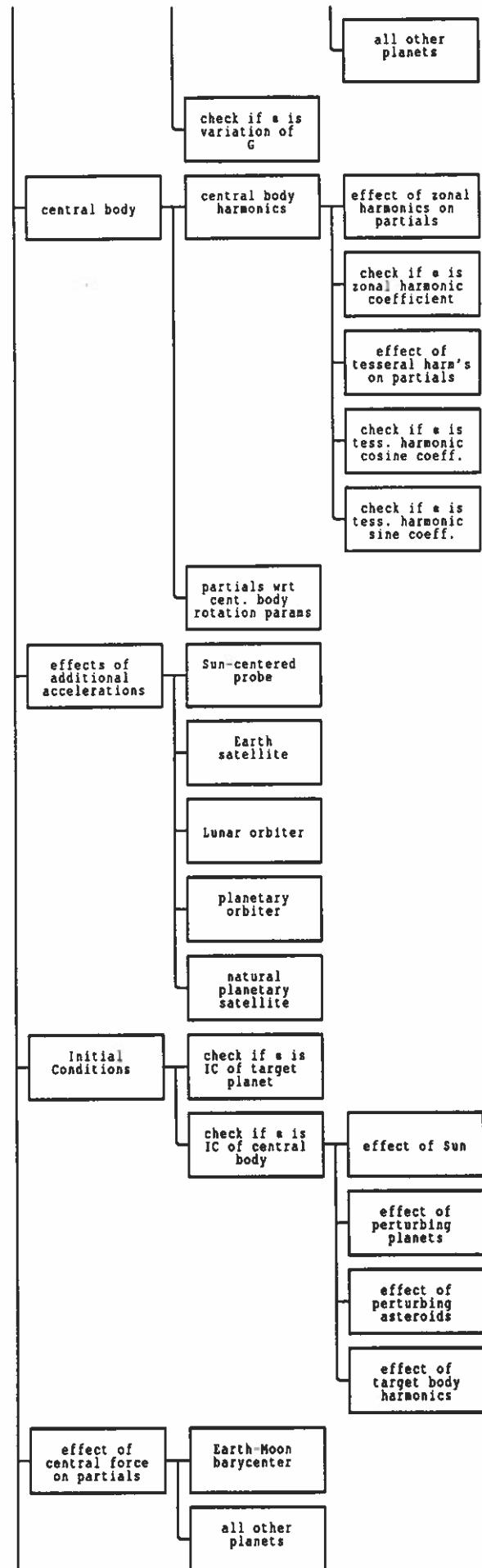
$$\vec{f}_{In} = -\gamma(t) \vec{a}_{cc}$$

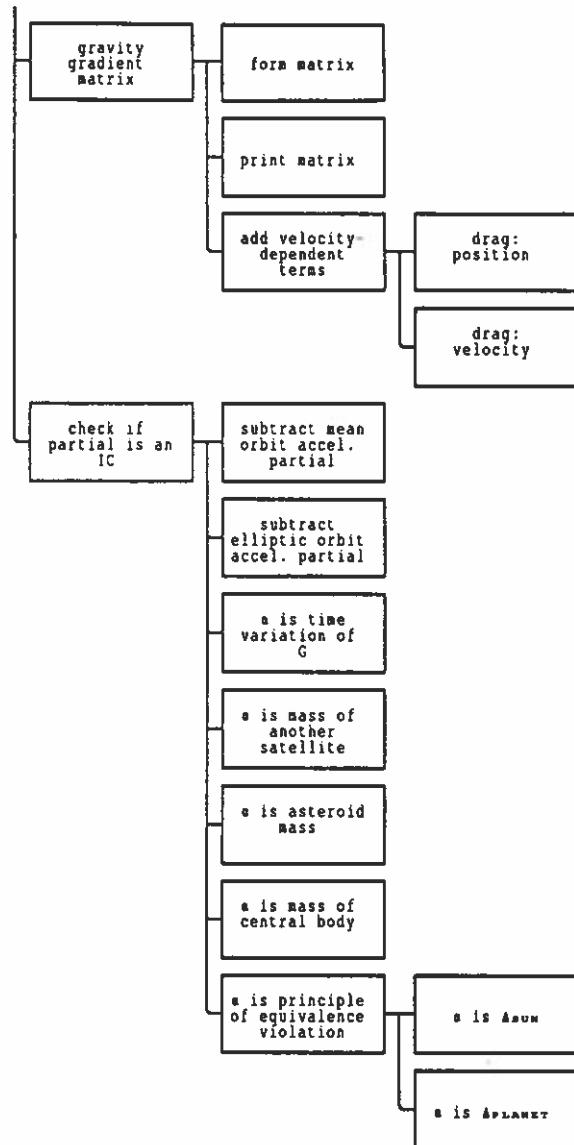
for the integrated body.



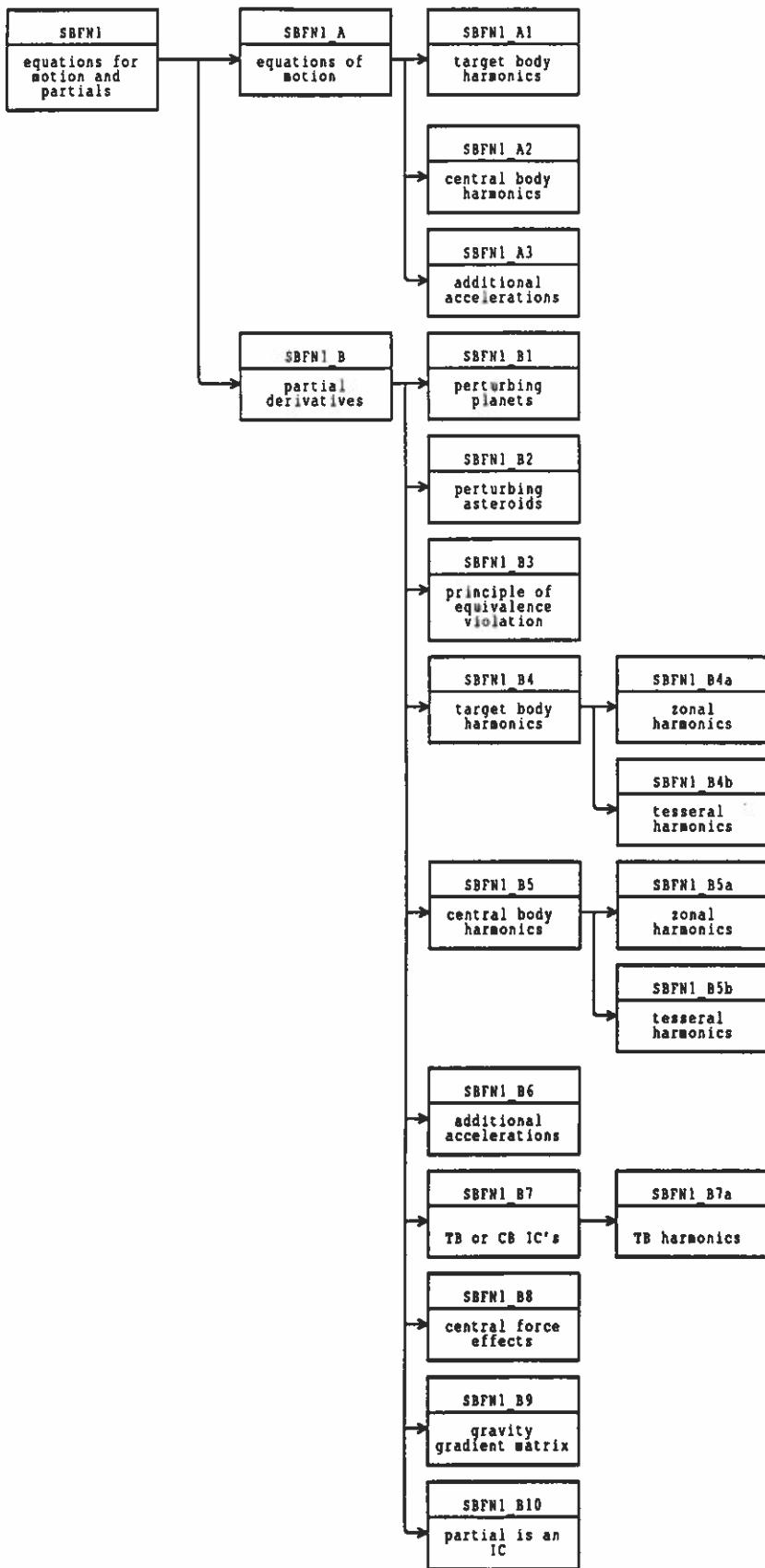








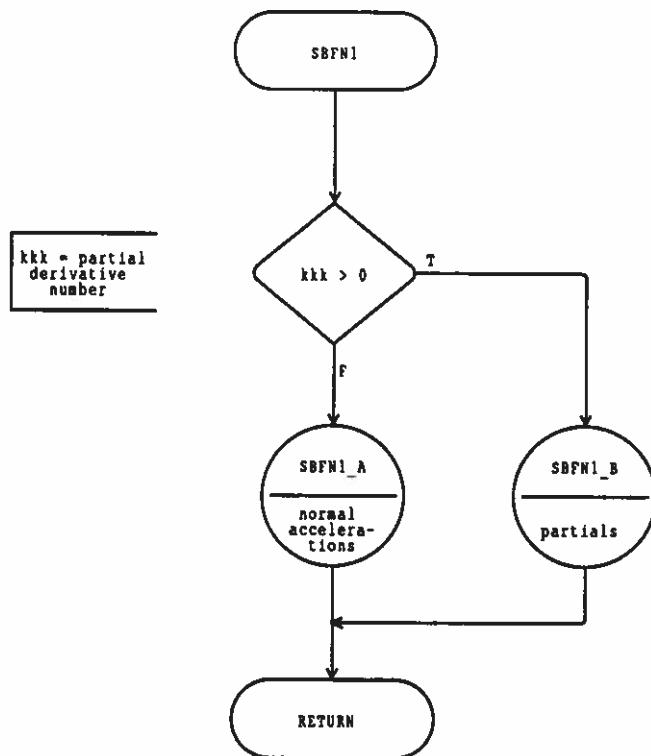
This chart shows the tree structure  
of the SBFN1 flow charts.

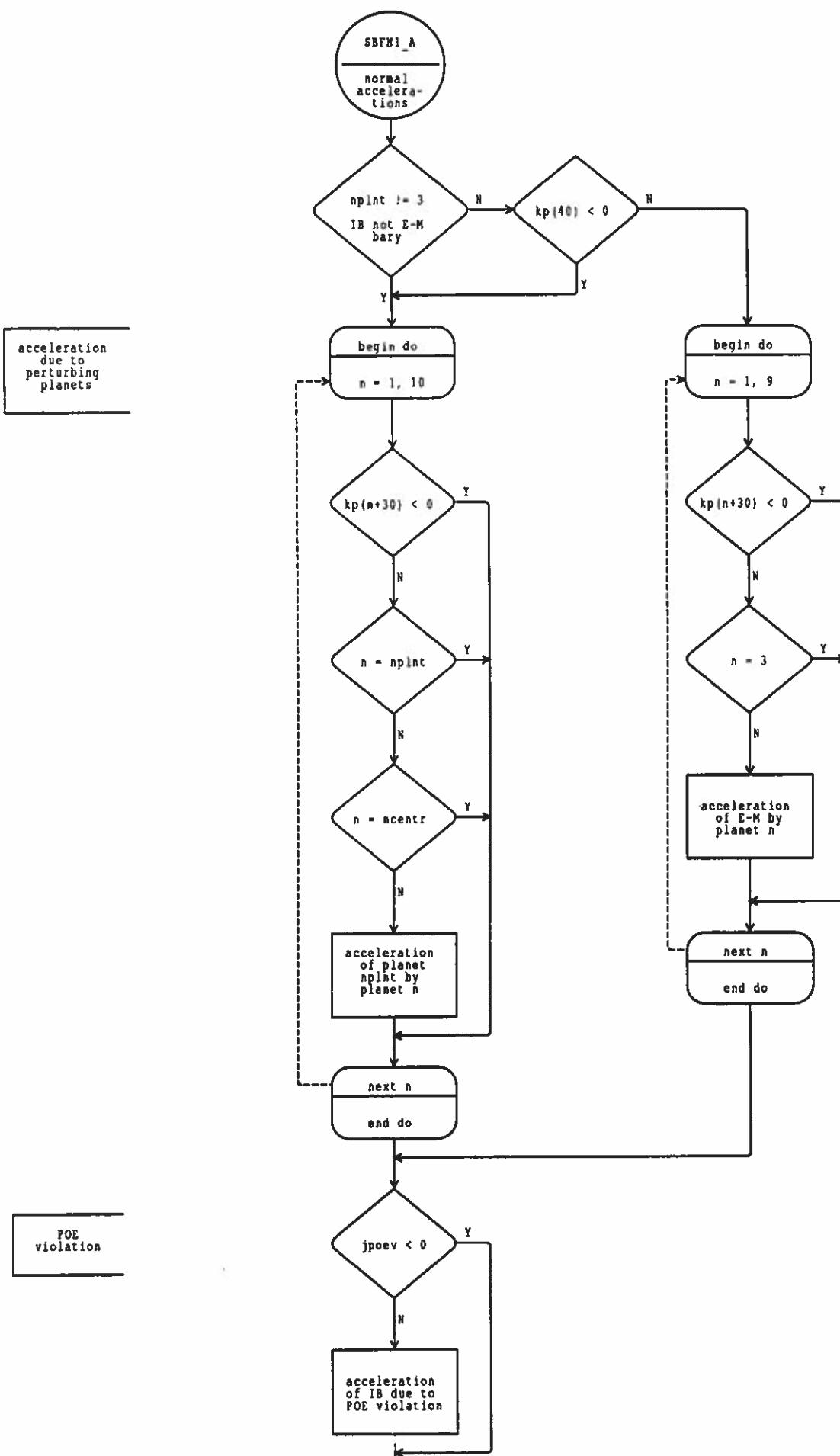


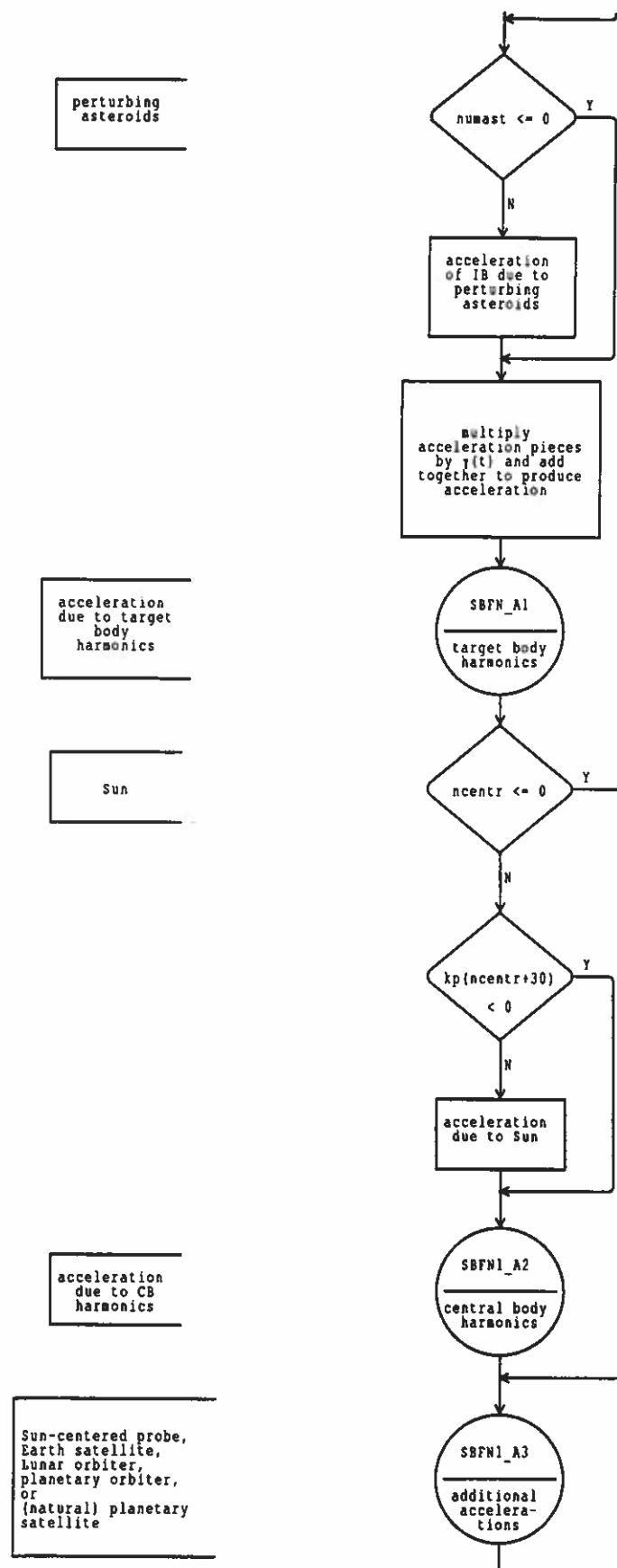
ABBREVIATIONS:

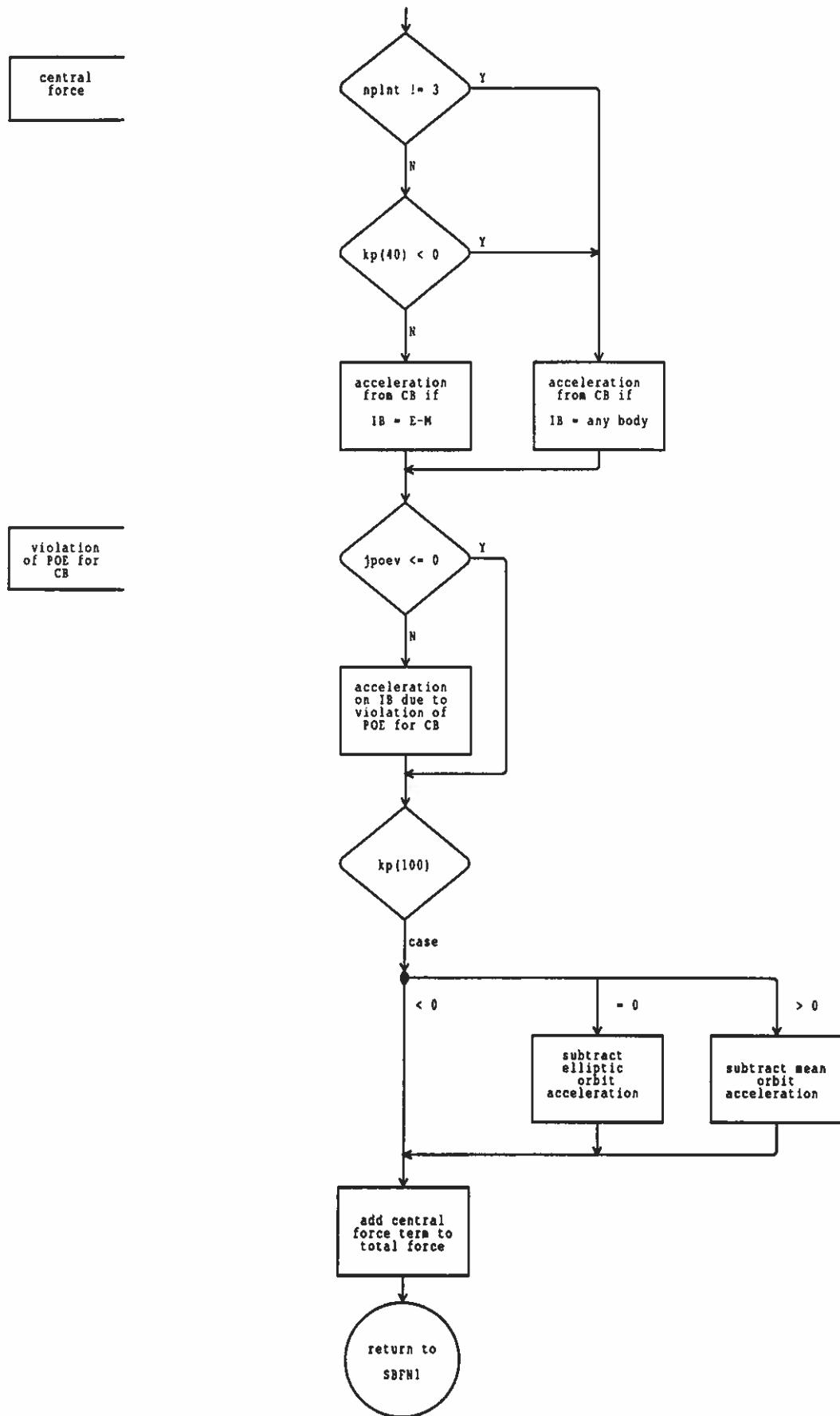
IB Integrated Body  
CB Central Body  
TB Target Body  
E-M Earth-Moon barycenter  
POE Principle Of Equivalence  
 $N_a$  number of zonal harmonics  
 $J_n$  zonal harmonic coefficient of degree  $n$   
 $S_{nm}$  tesseral harmonic sine coefficient of degree  $n$  and order  $m$   
 $C_{nm}$  tesseral harmonic cosine coefficient of degree  $n$  and order  $m$   
 $N_t$  number of tesseral harmonics  
 $N_{tano}$  number of target bodies  
GG Gravity Gradient matrix  
 $k$  partial derivative parameter

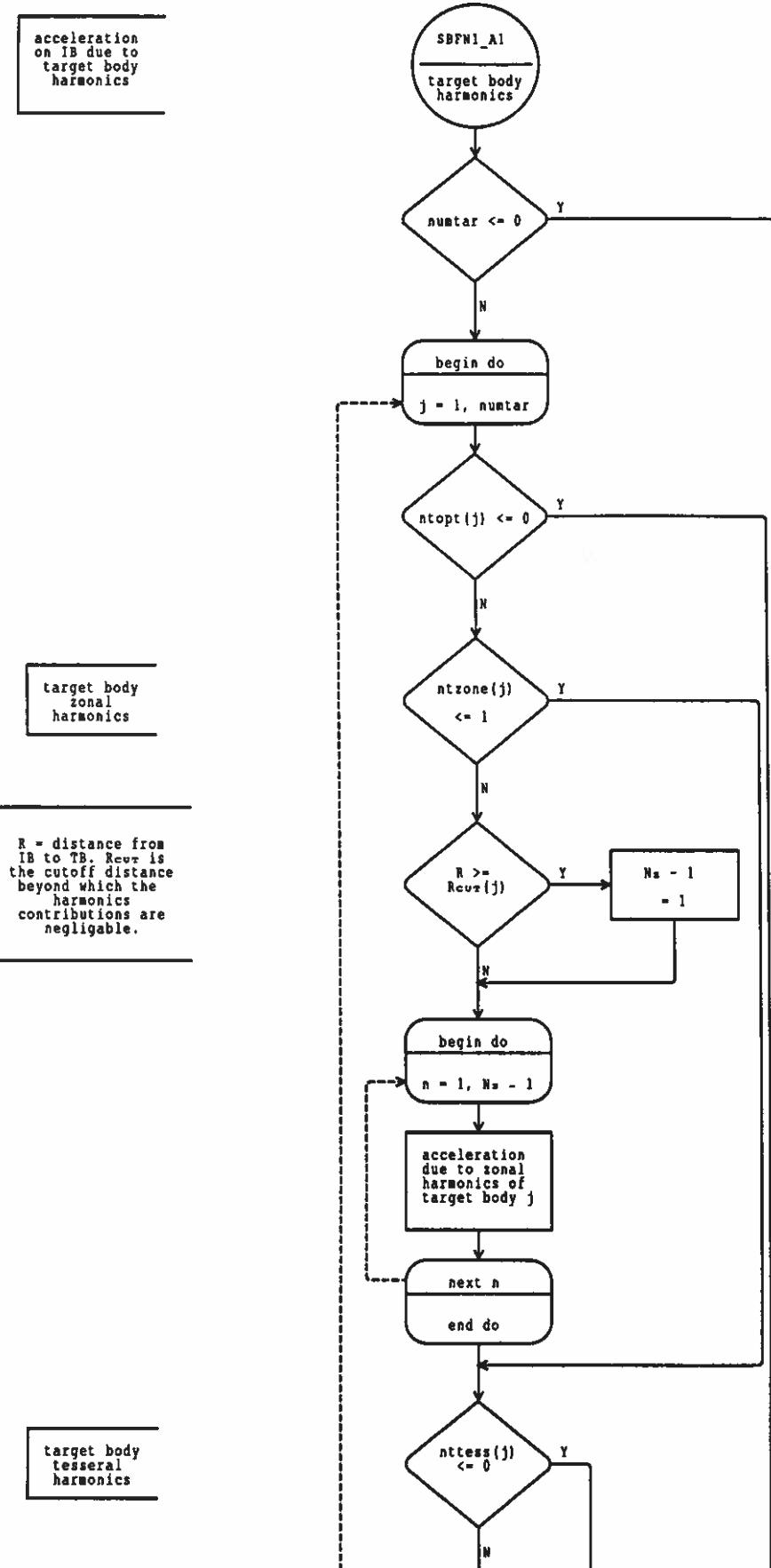
Also, for all conditionals, "Y" and "N" are equivalent to "T" and "F", respectively.

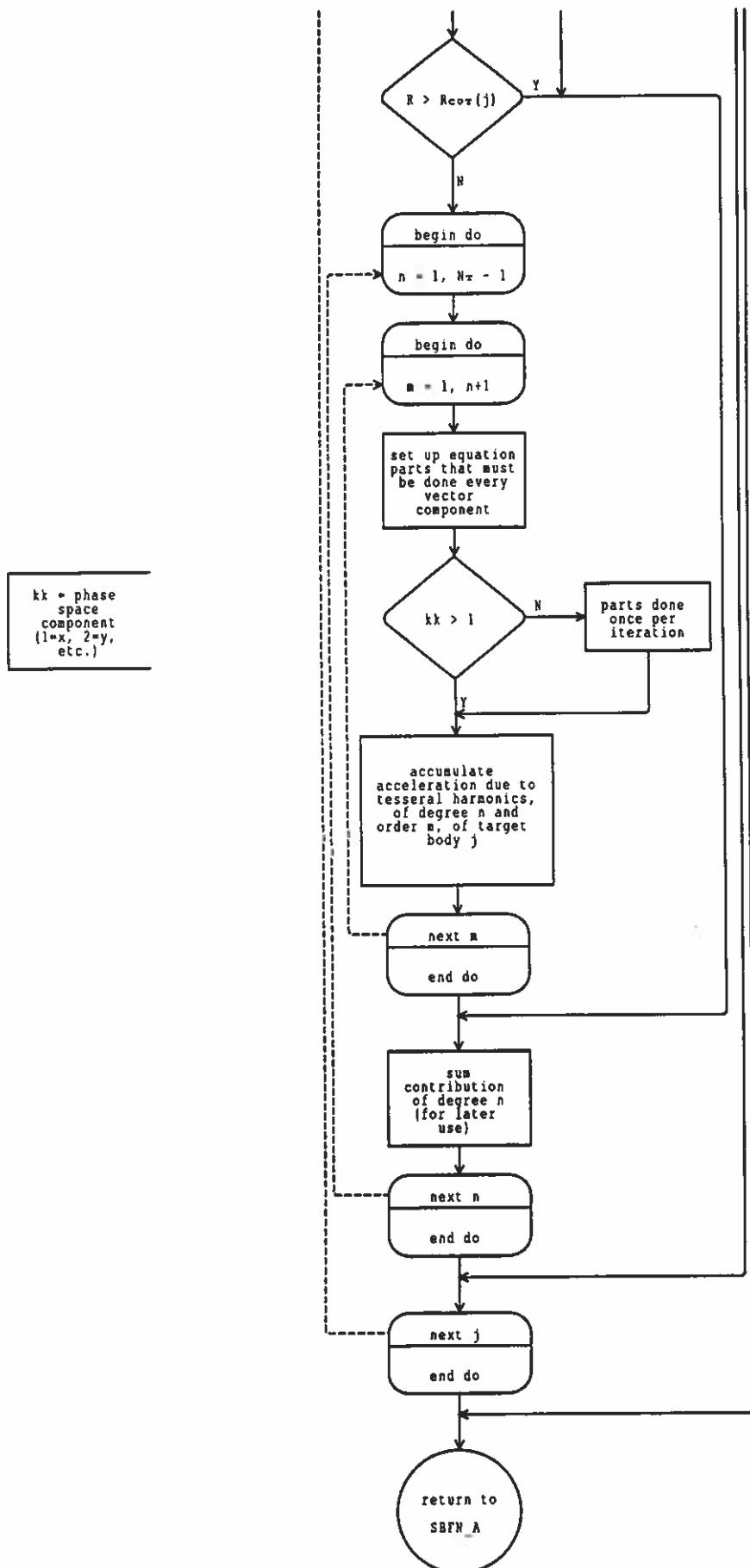


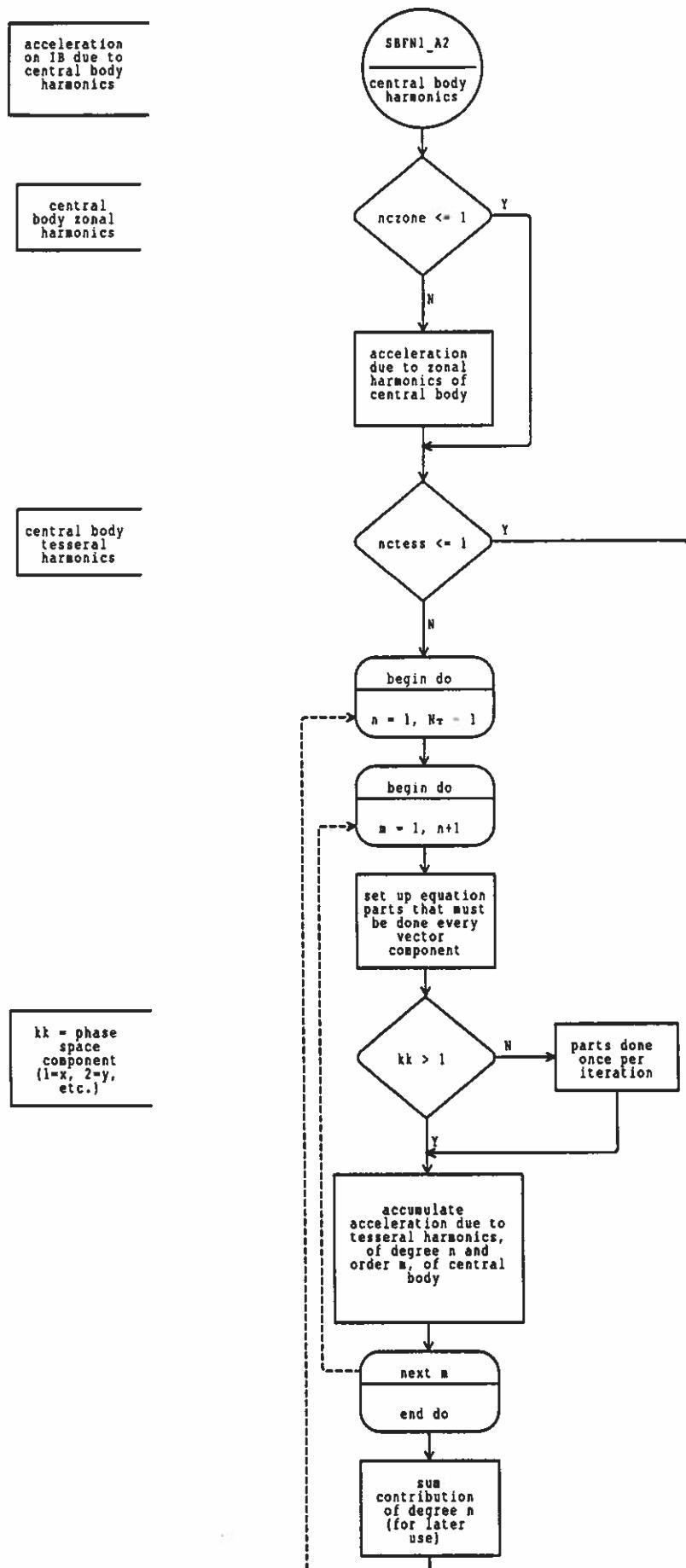


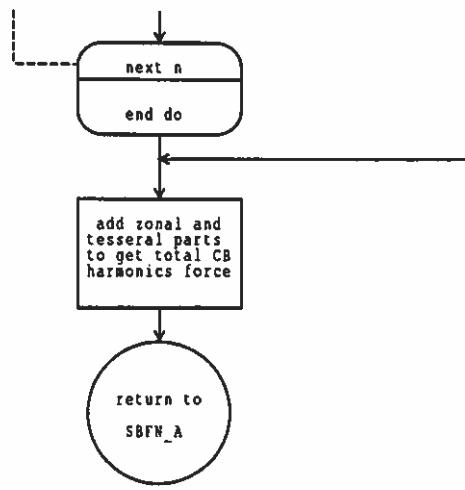


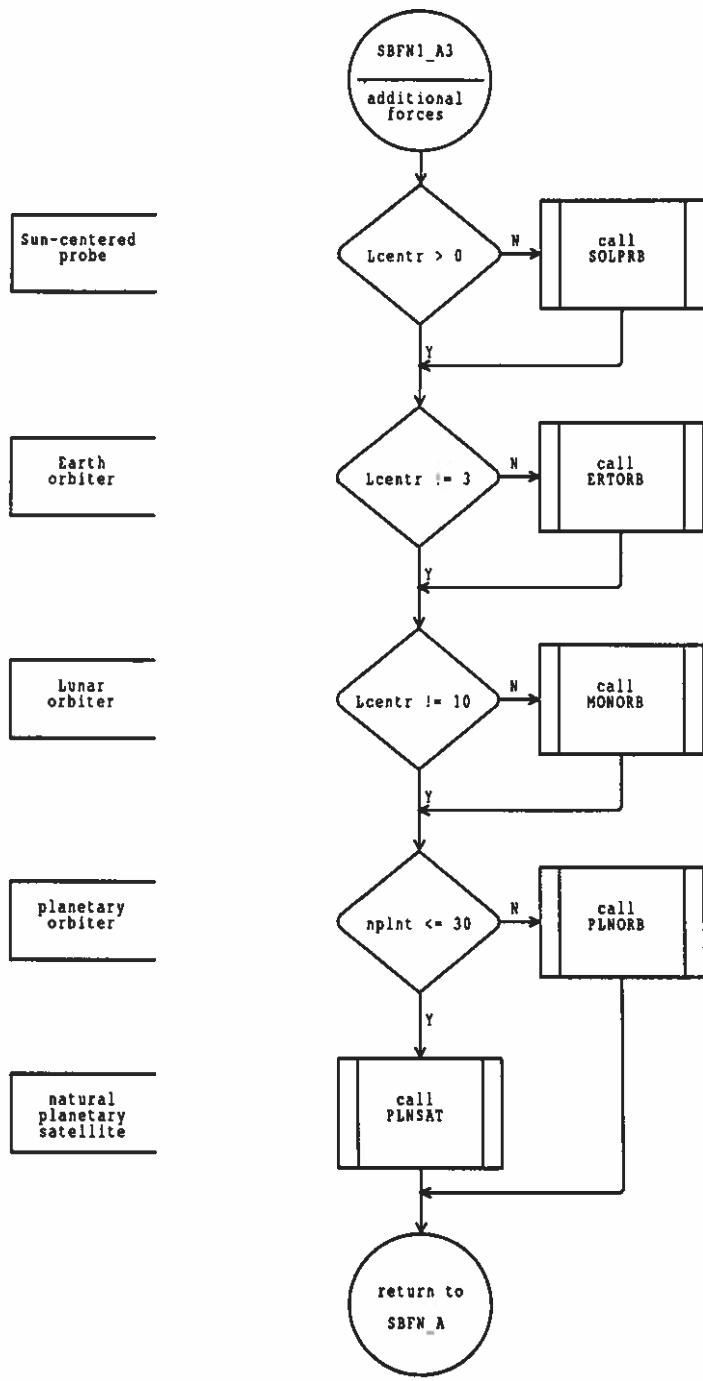


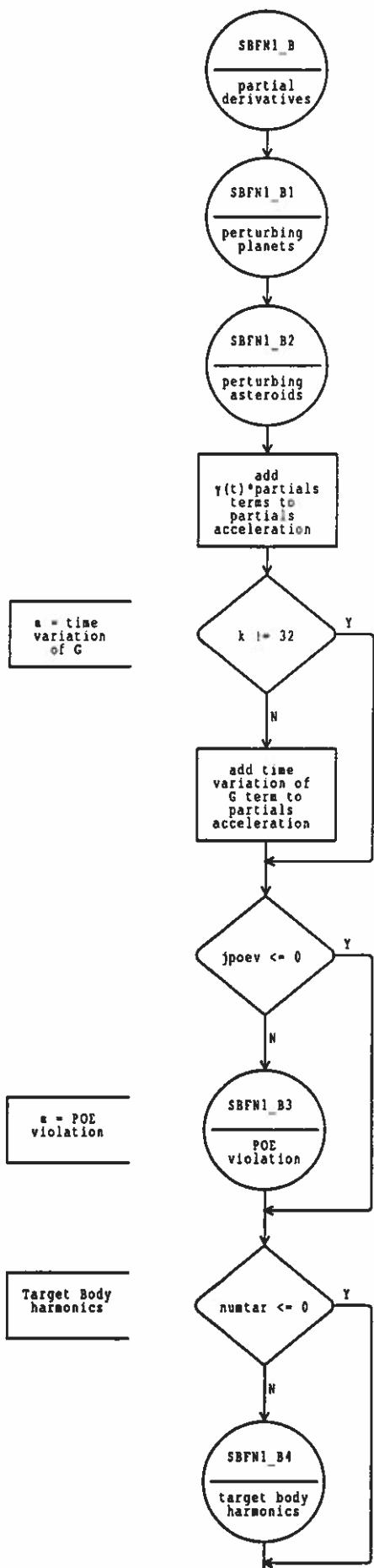


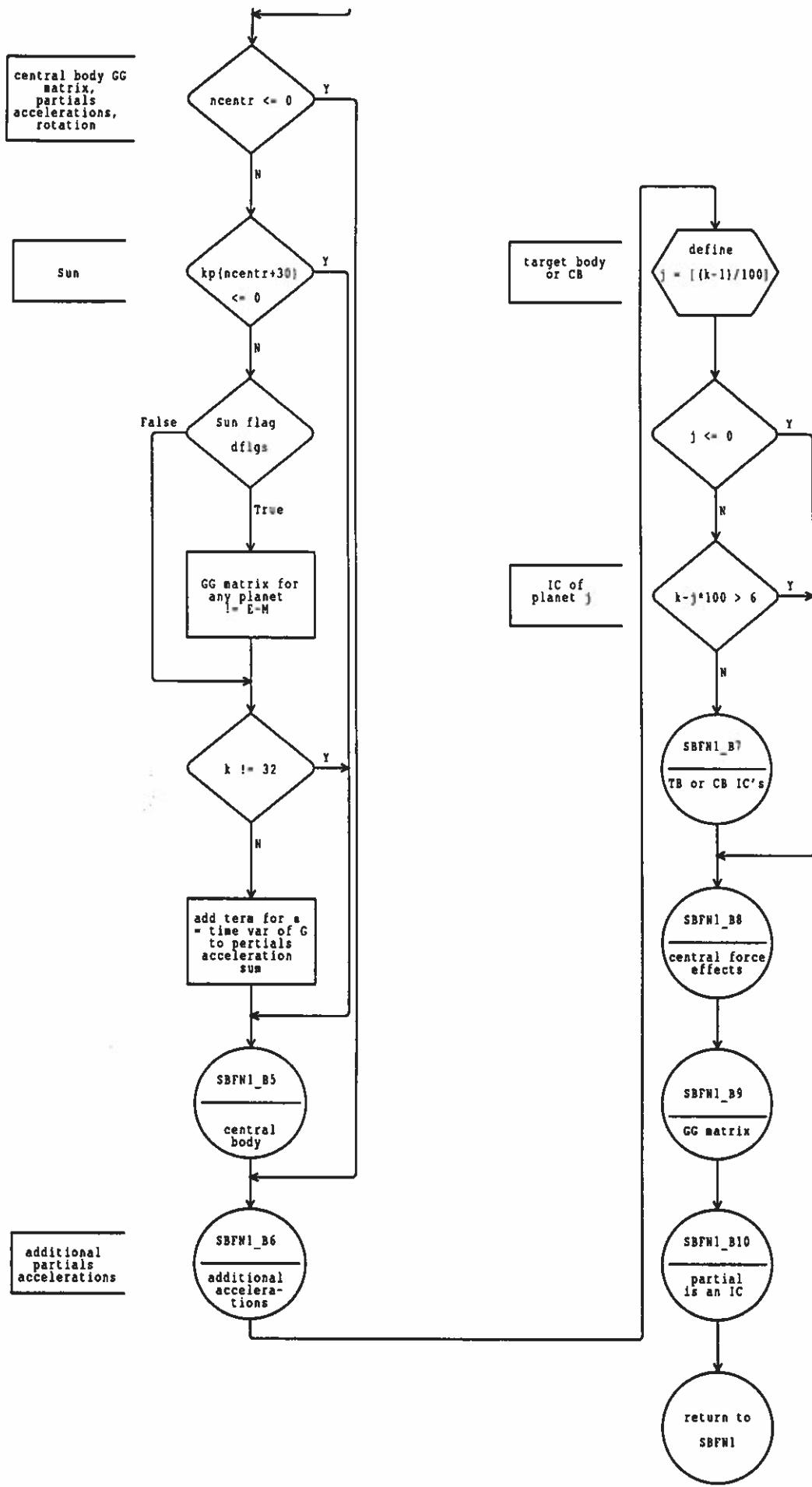


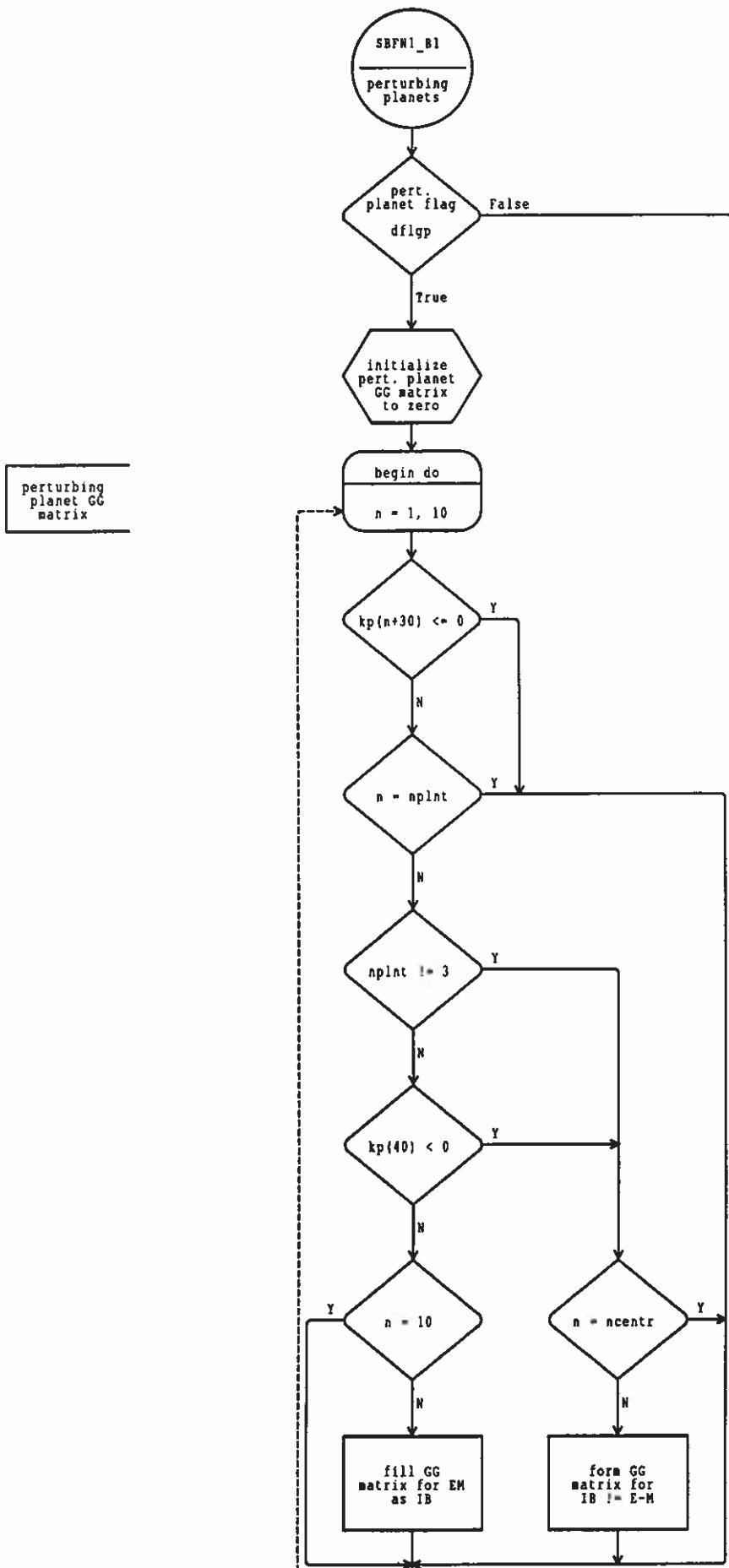


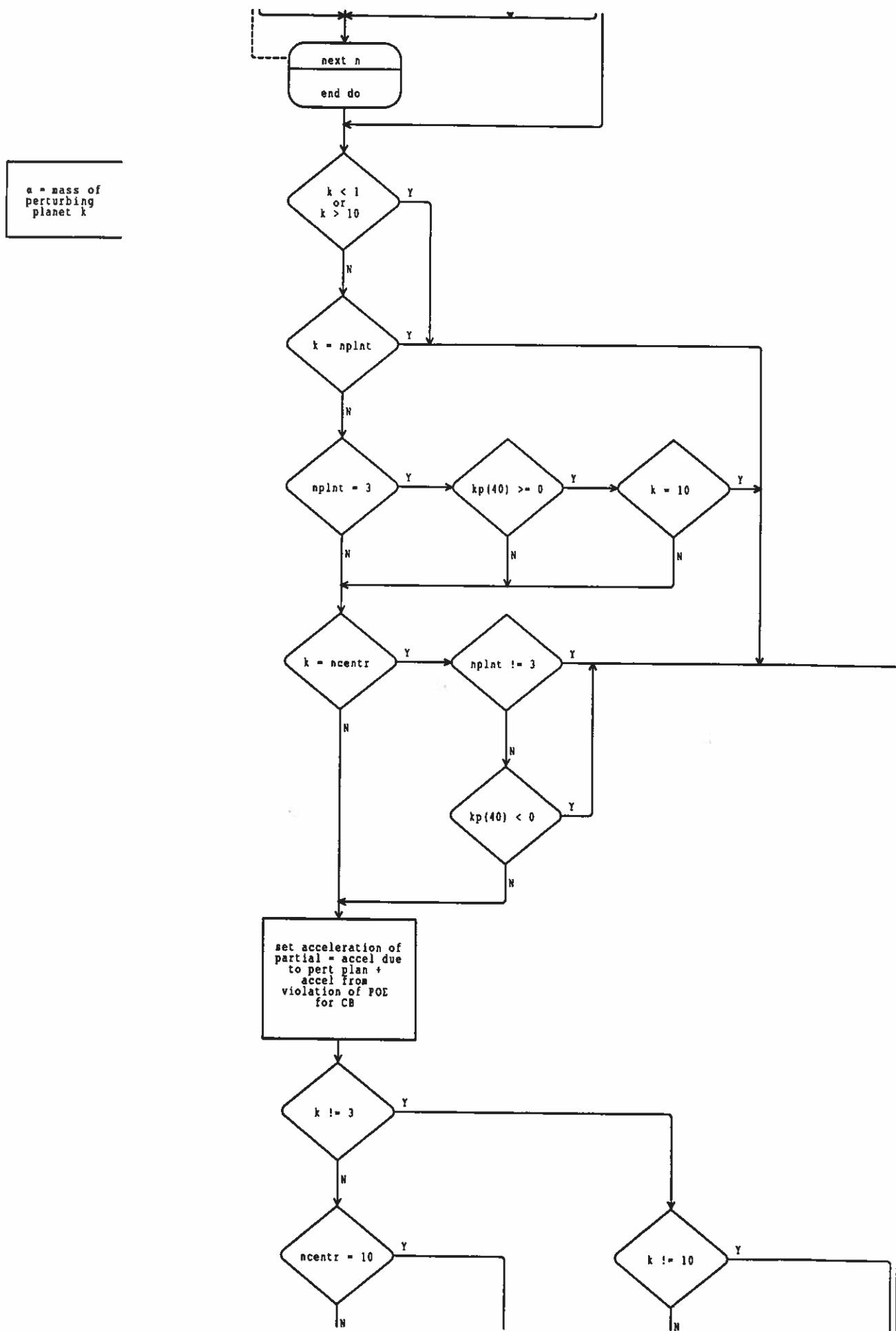


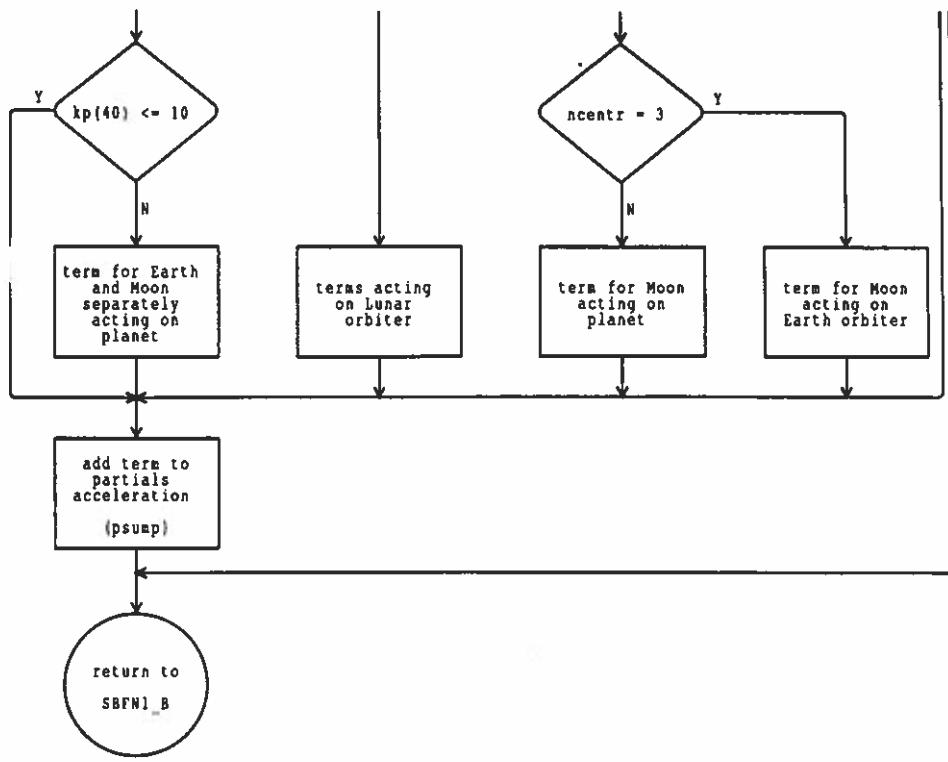


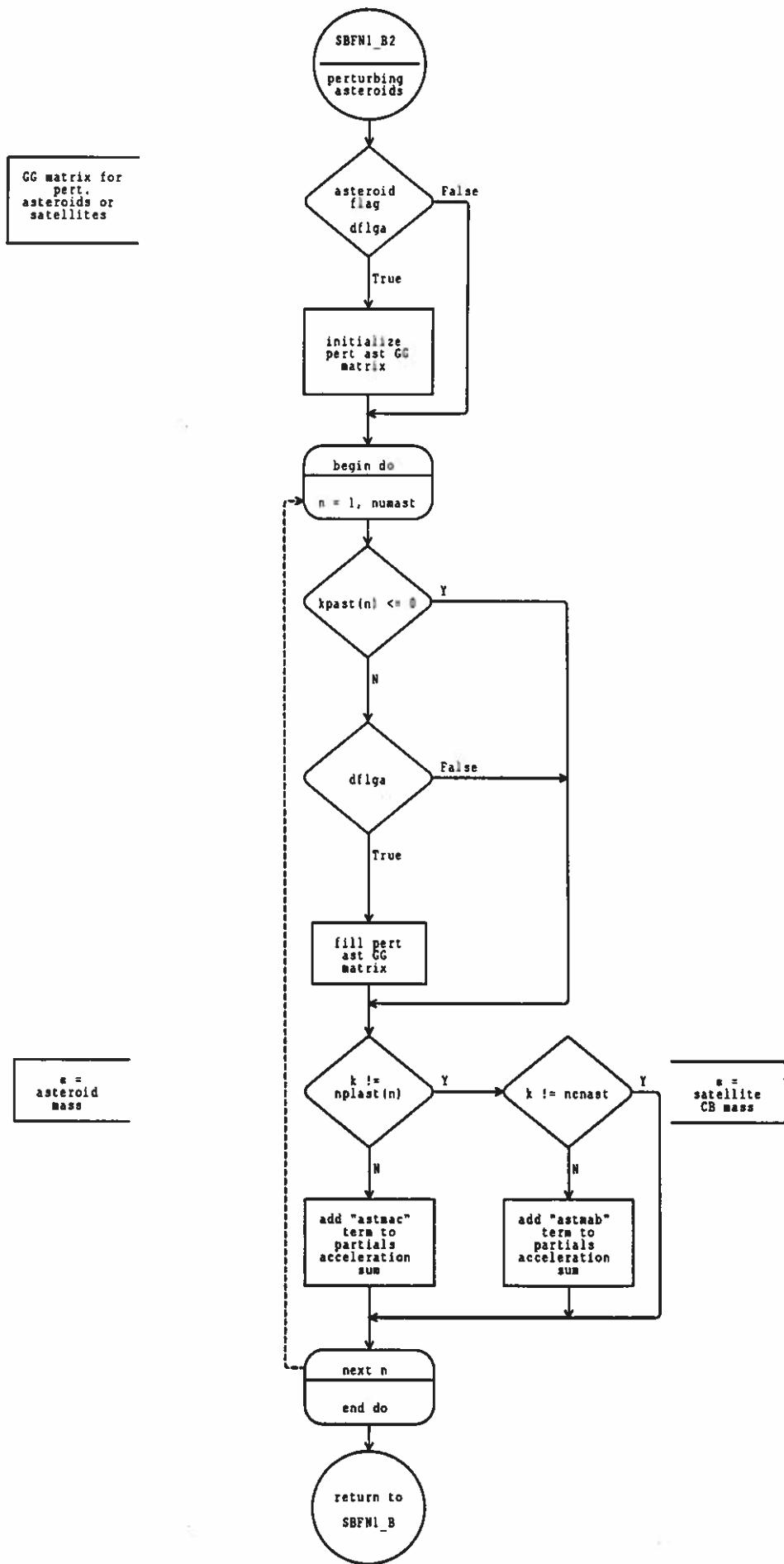




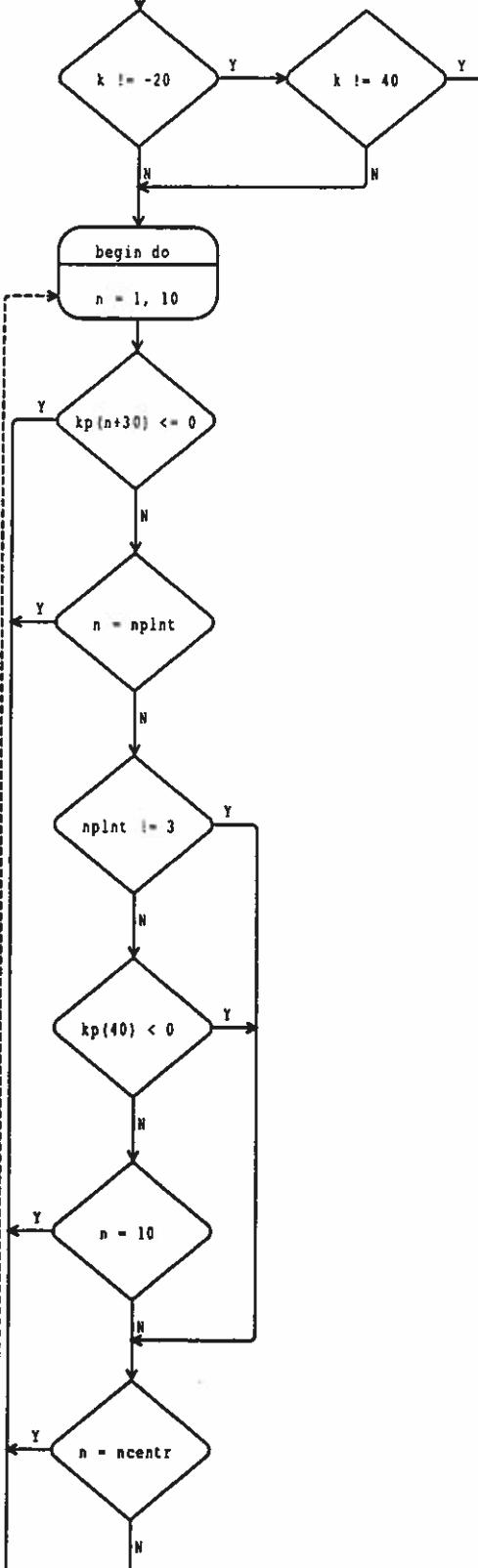


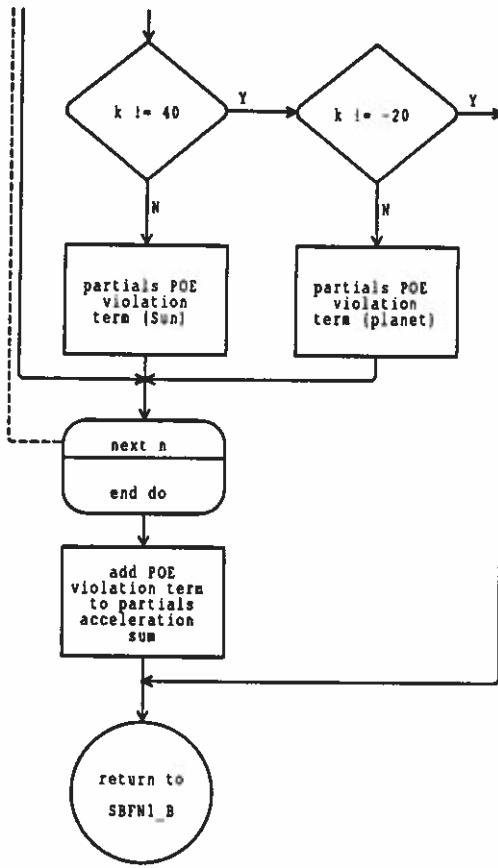


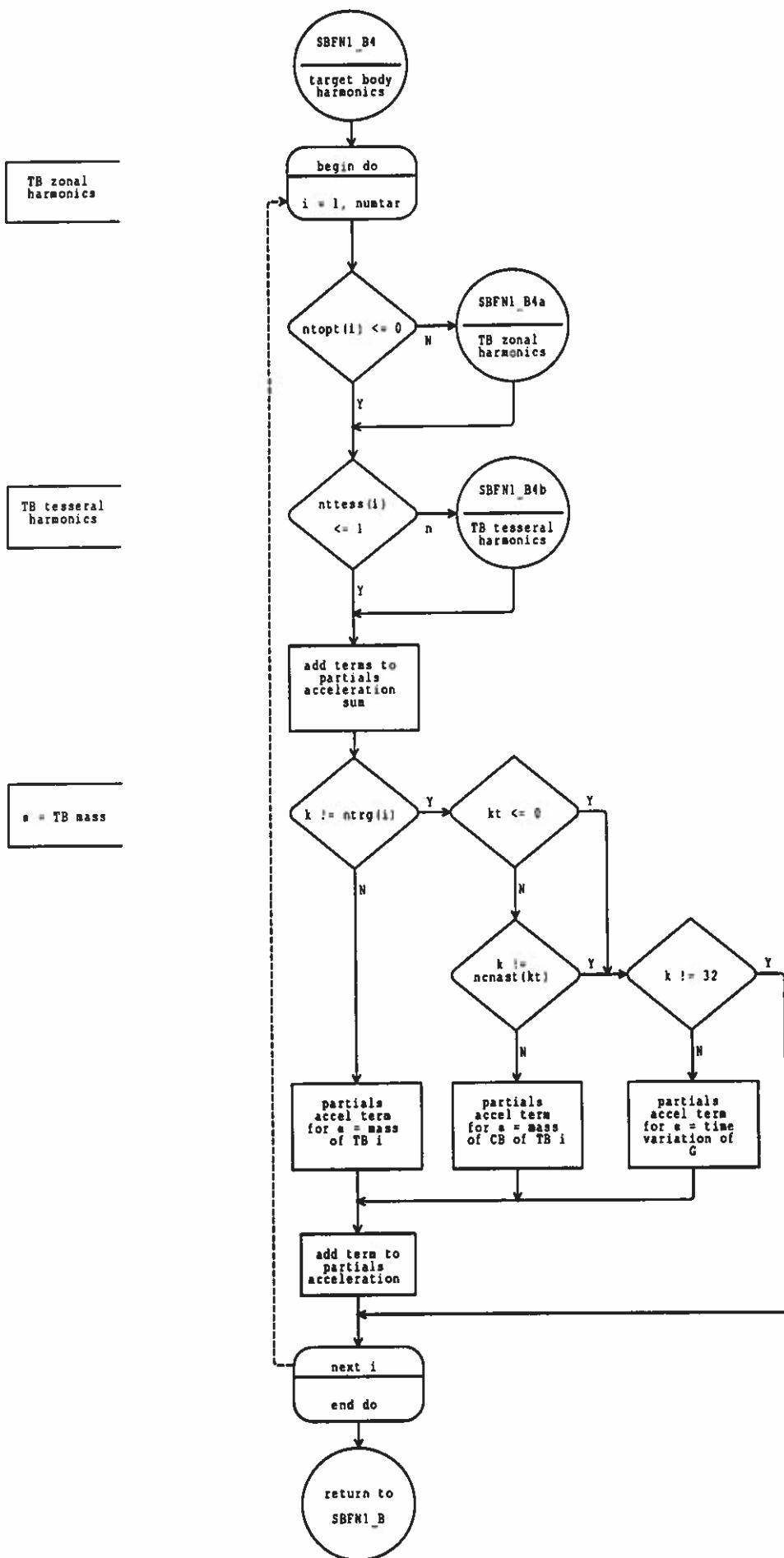


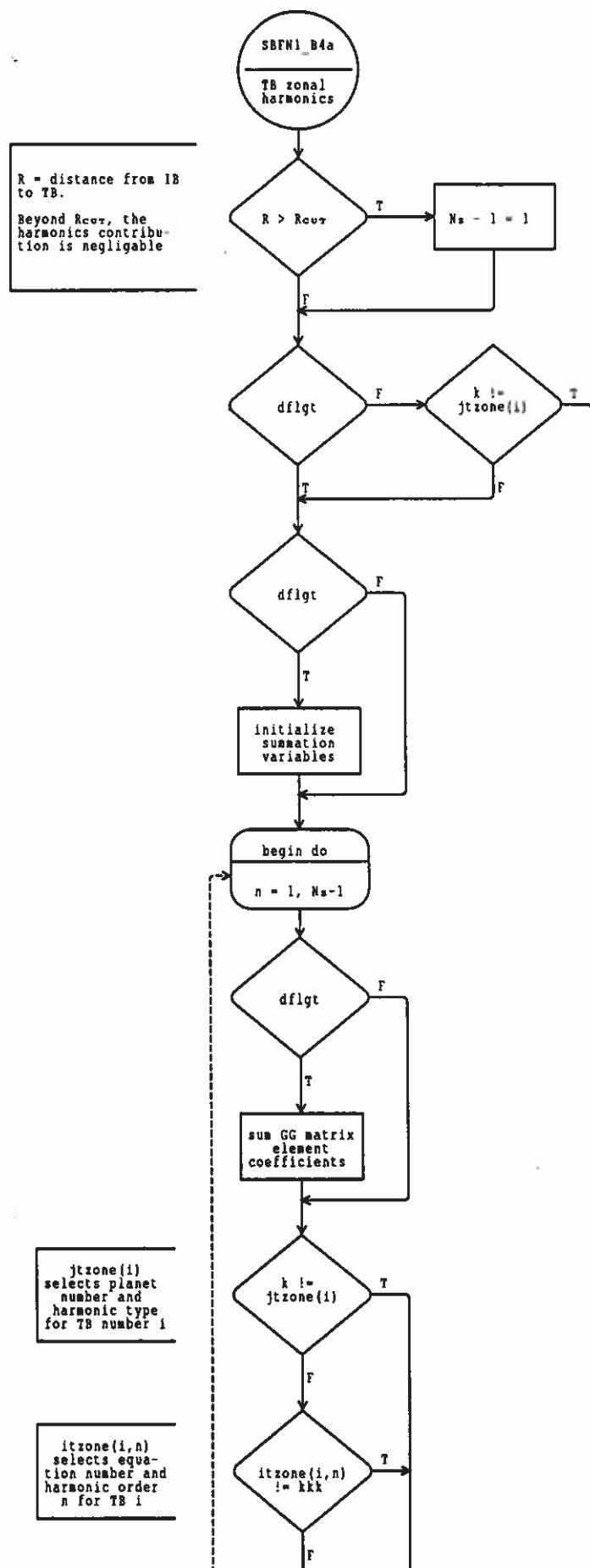


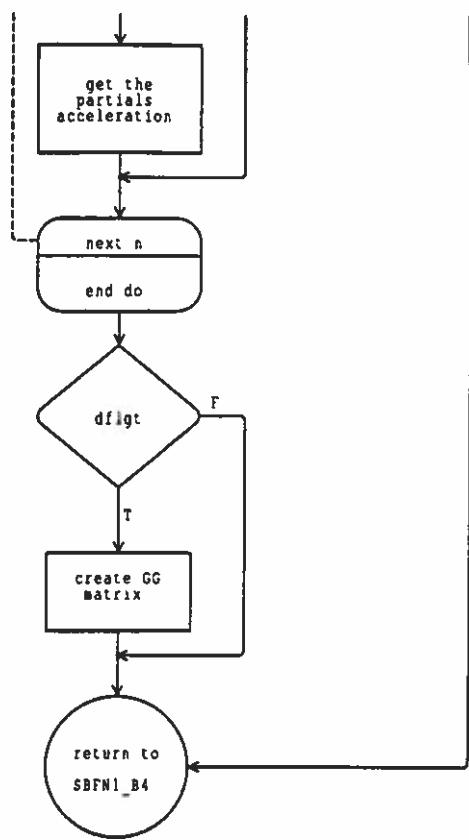
SBFW1\_B3  
POE  
violation

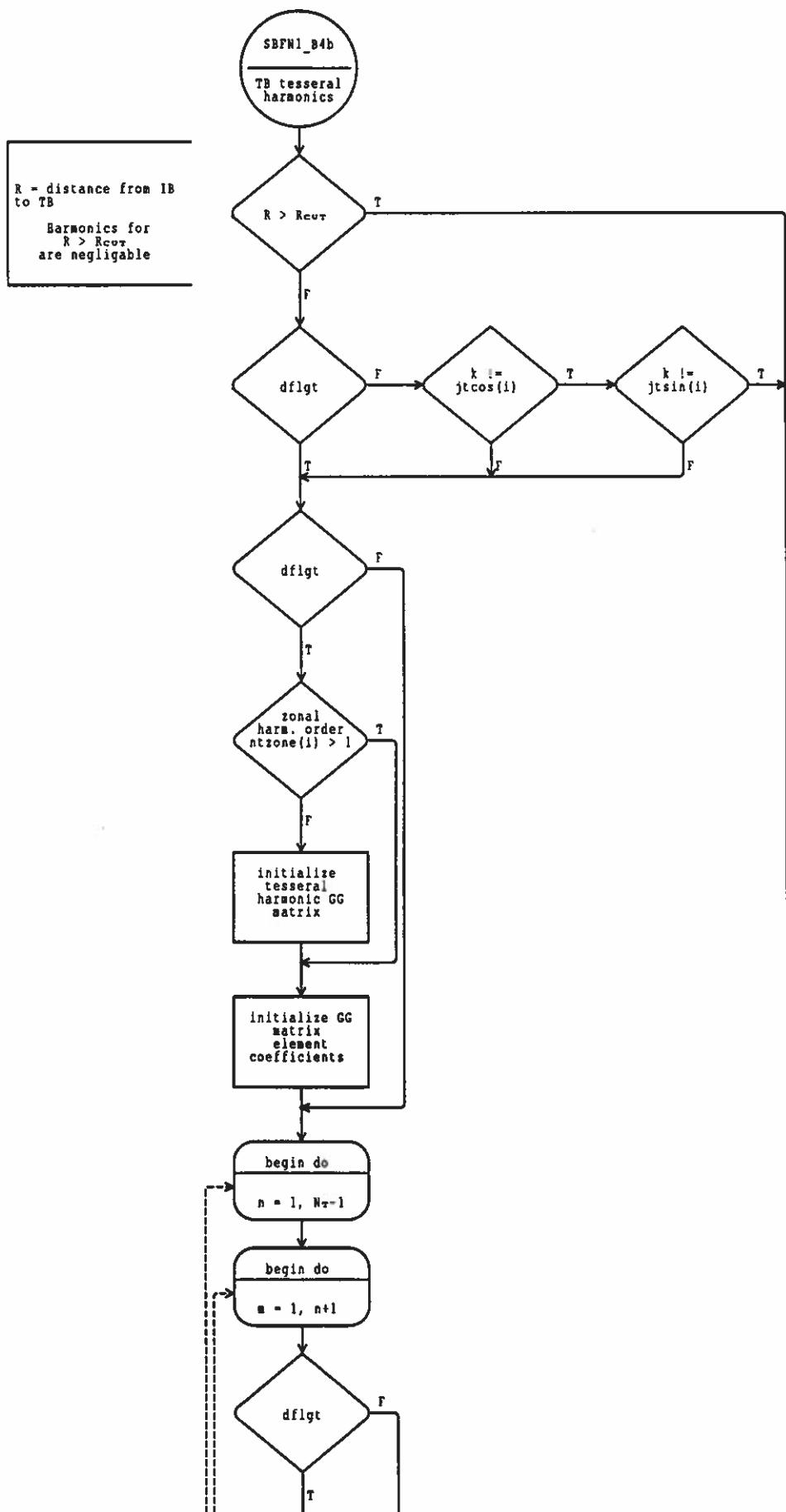


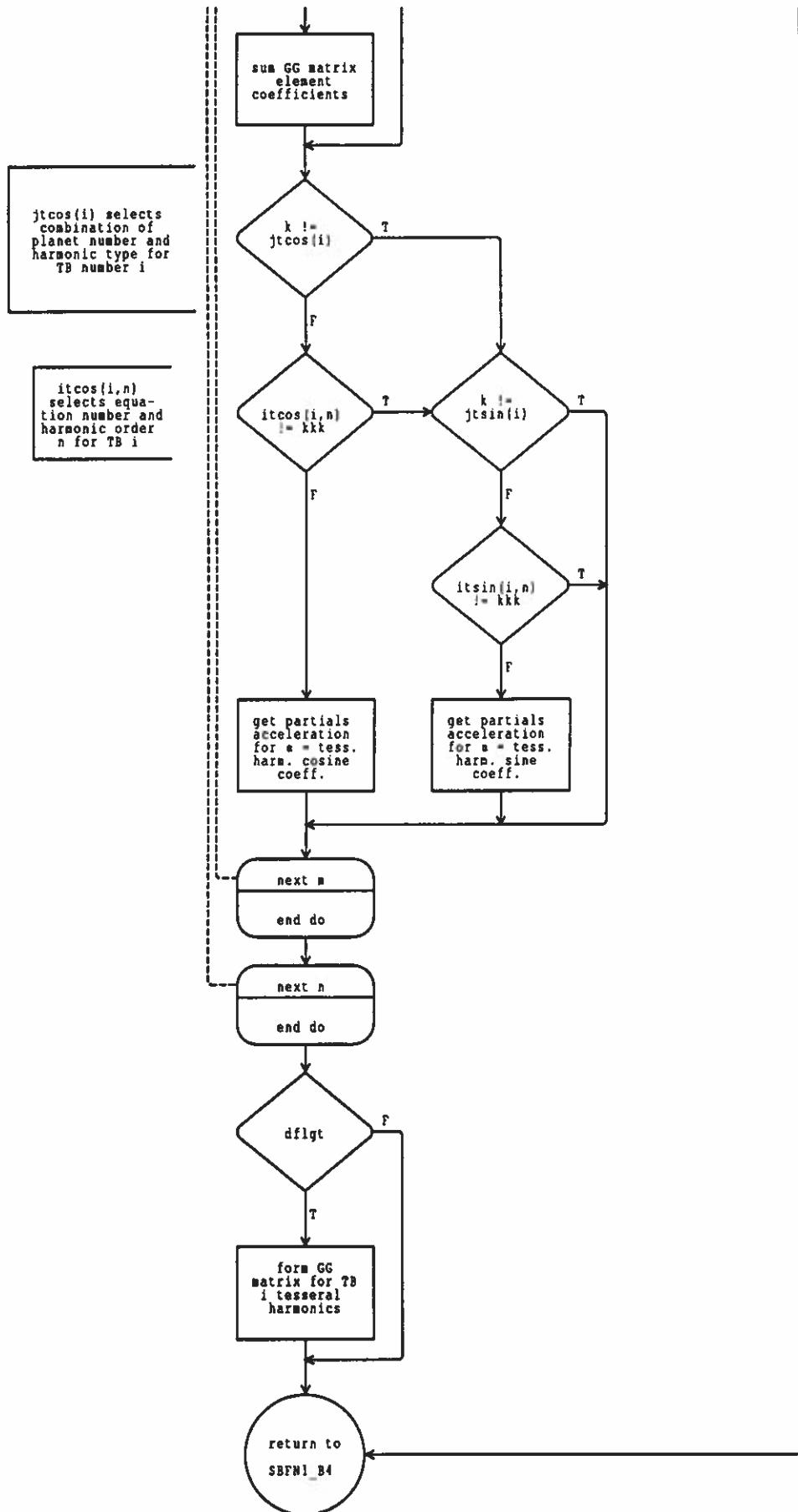


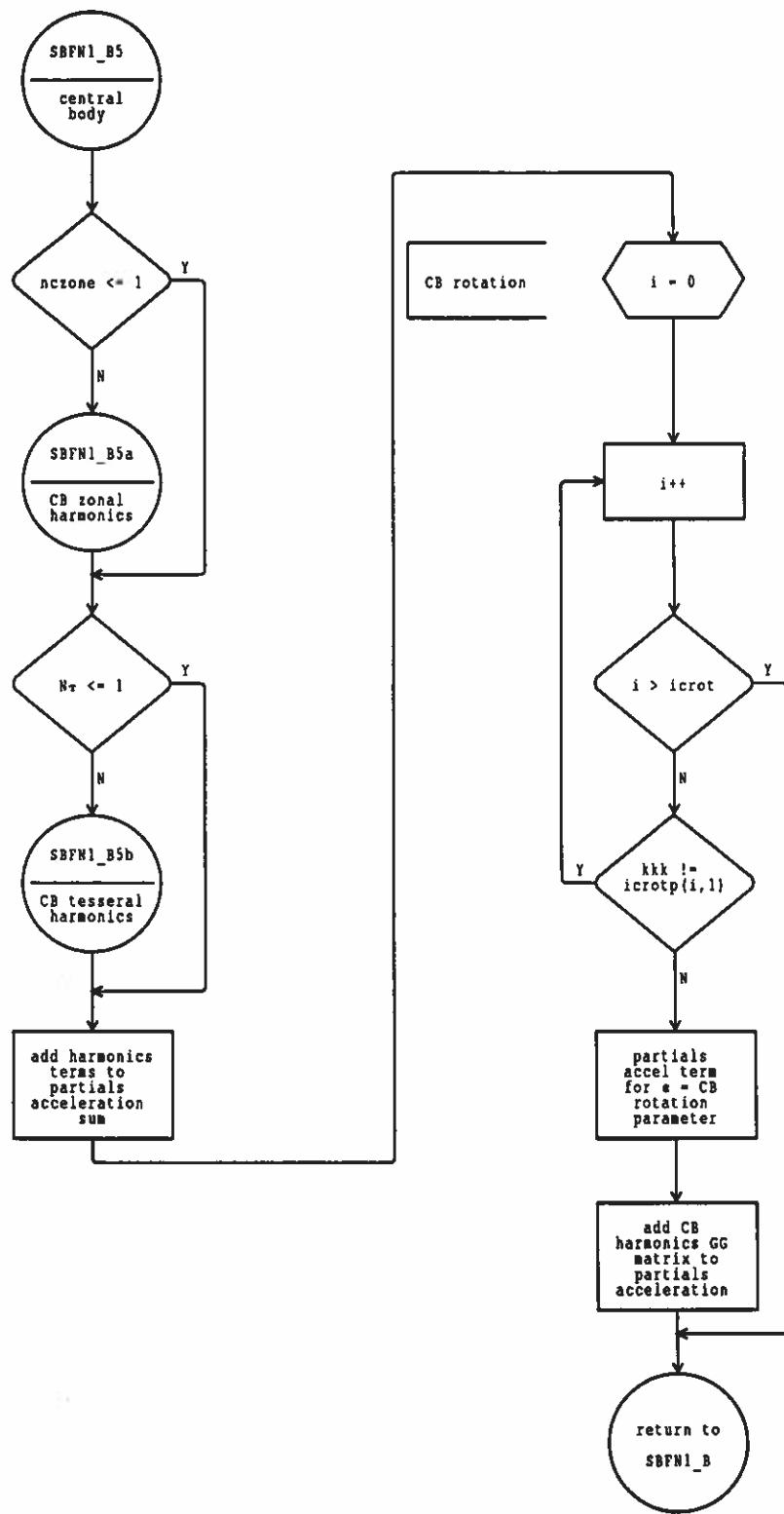


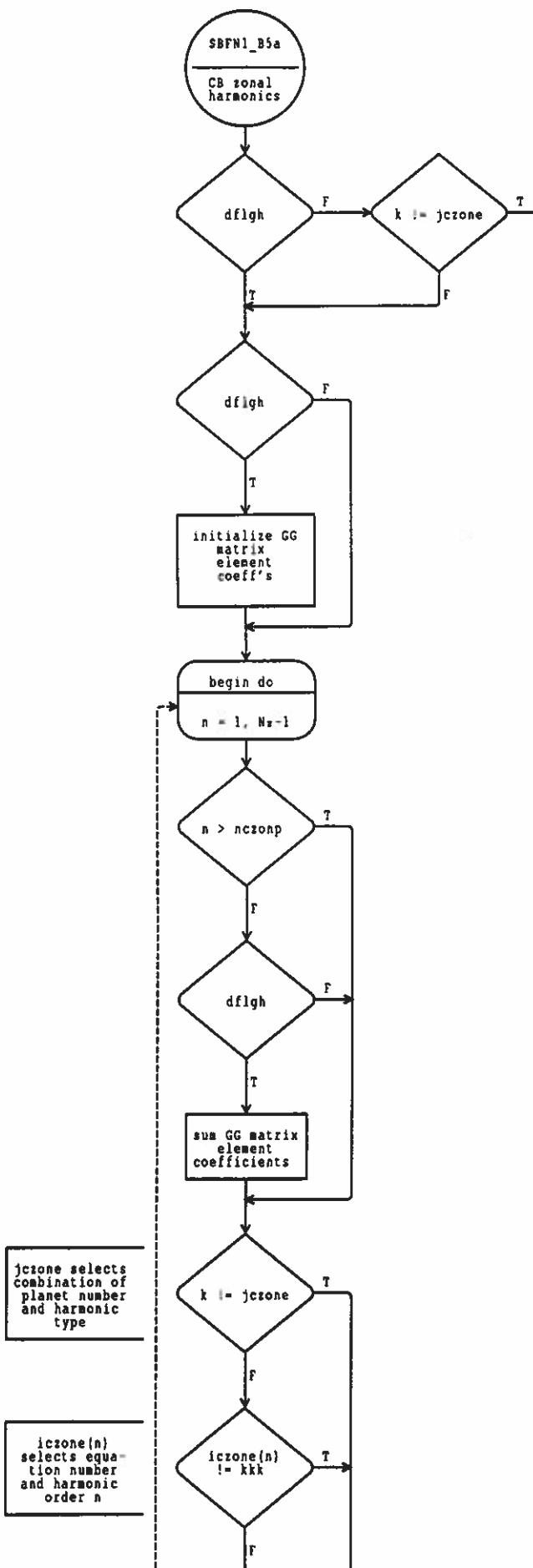


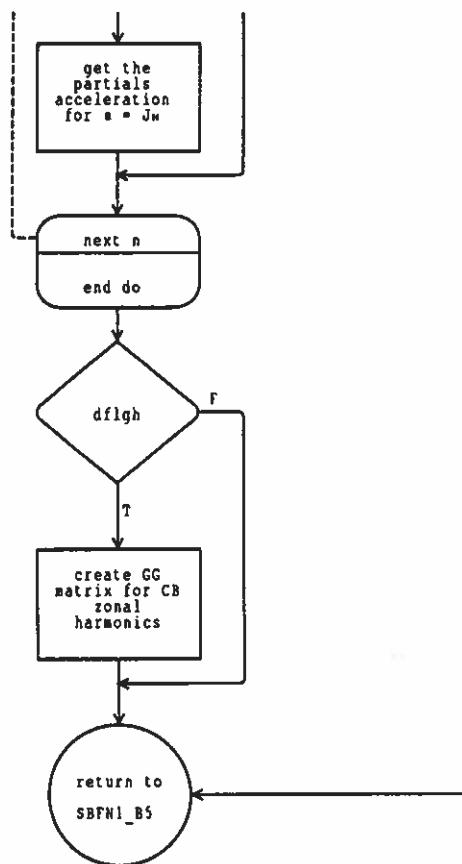


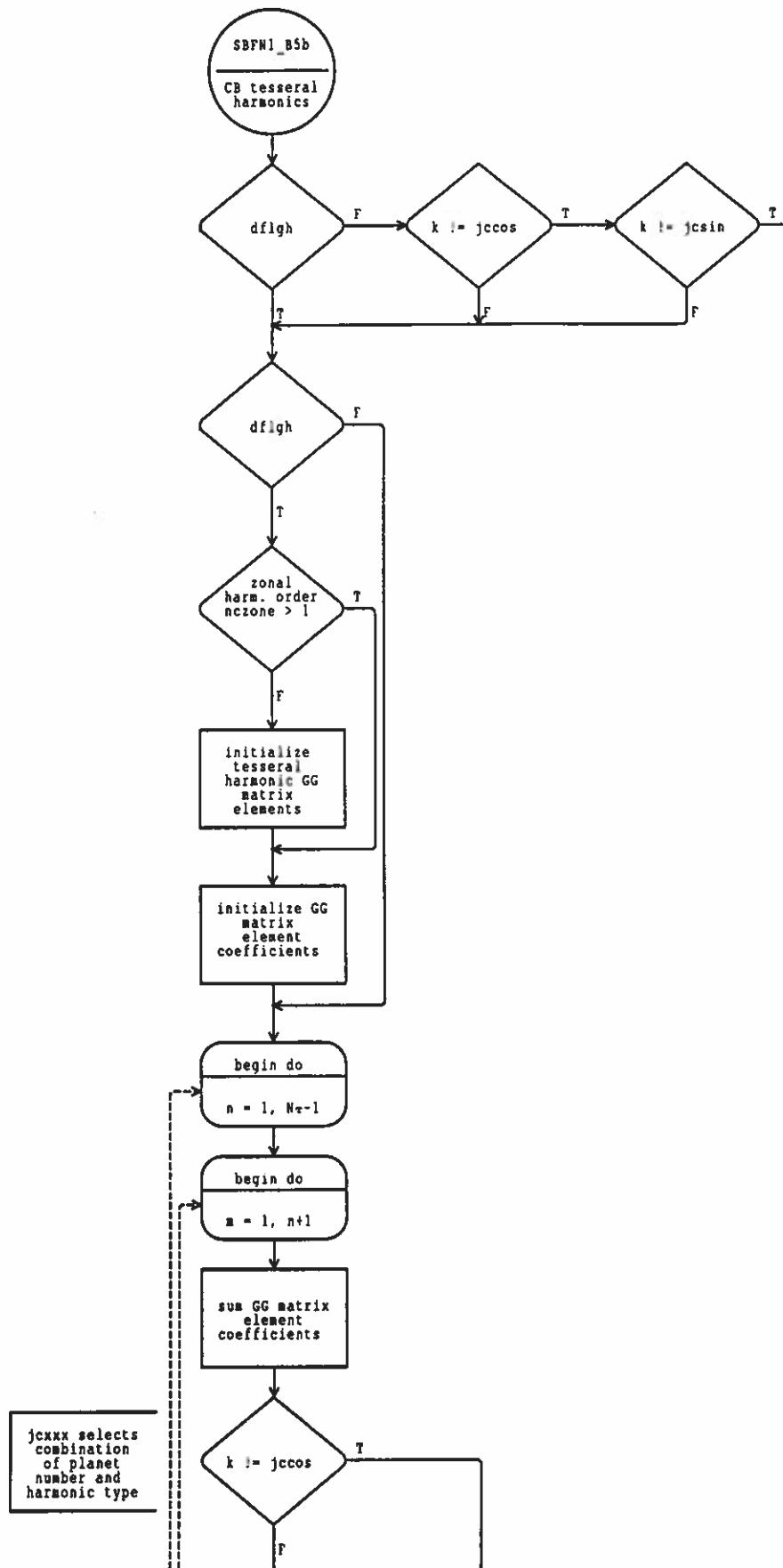


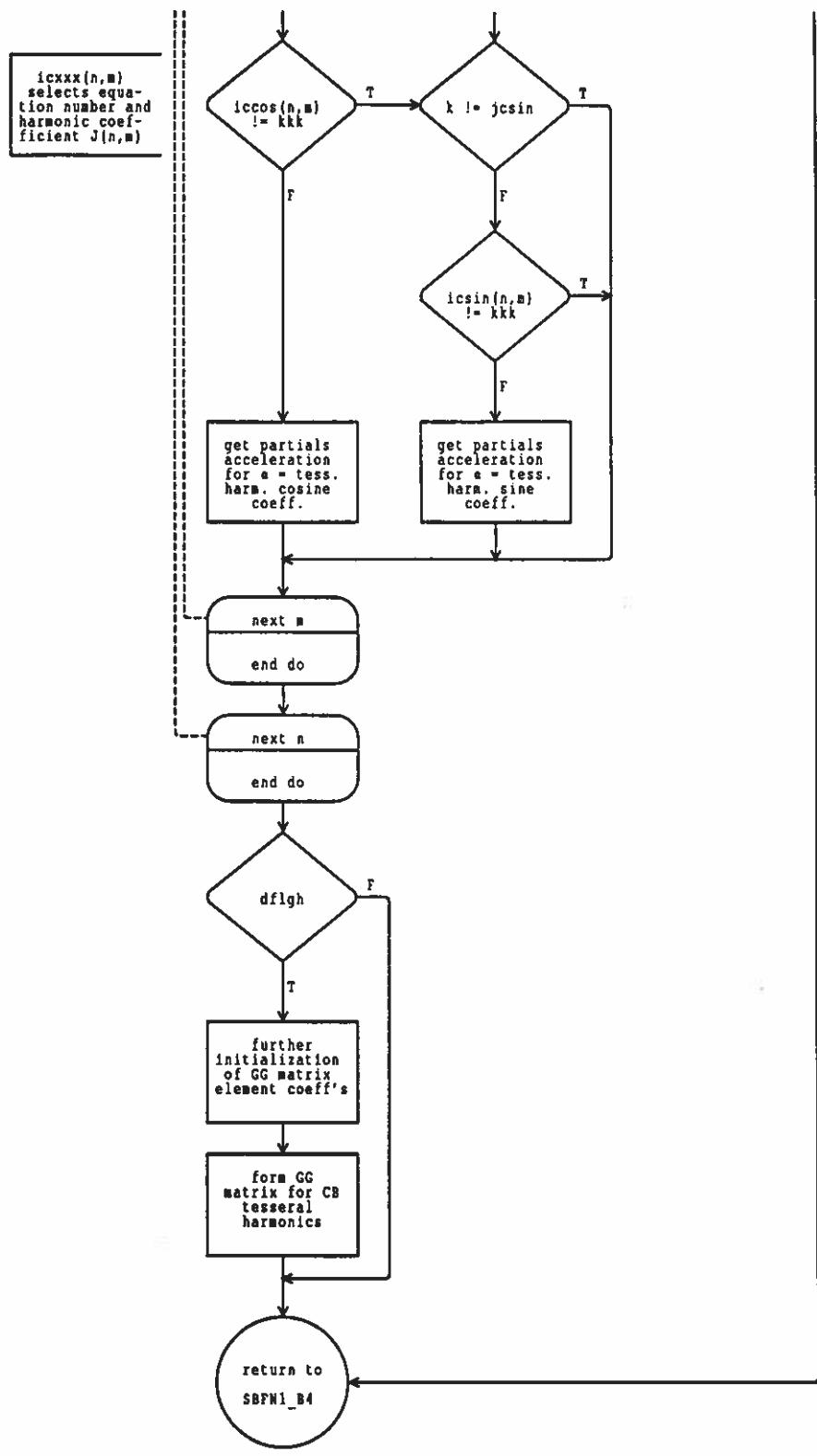


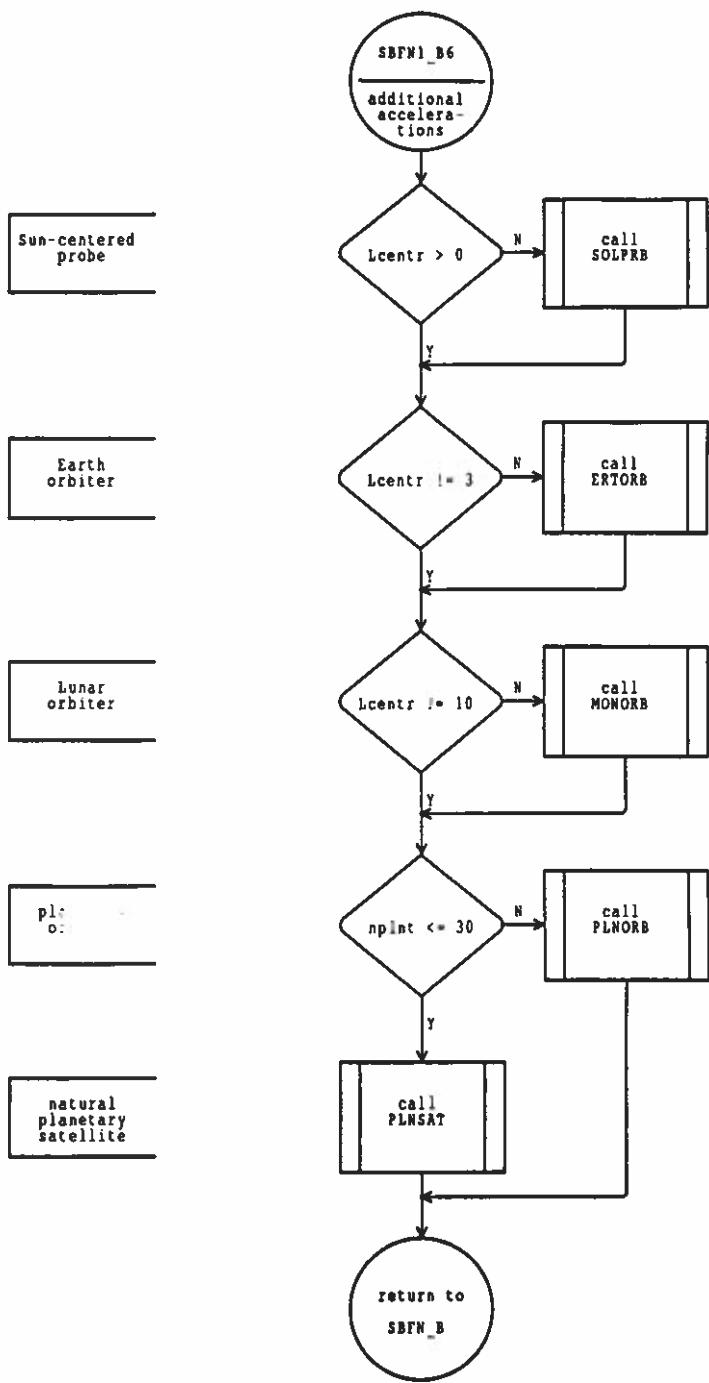












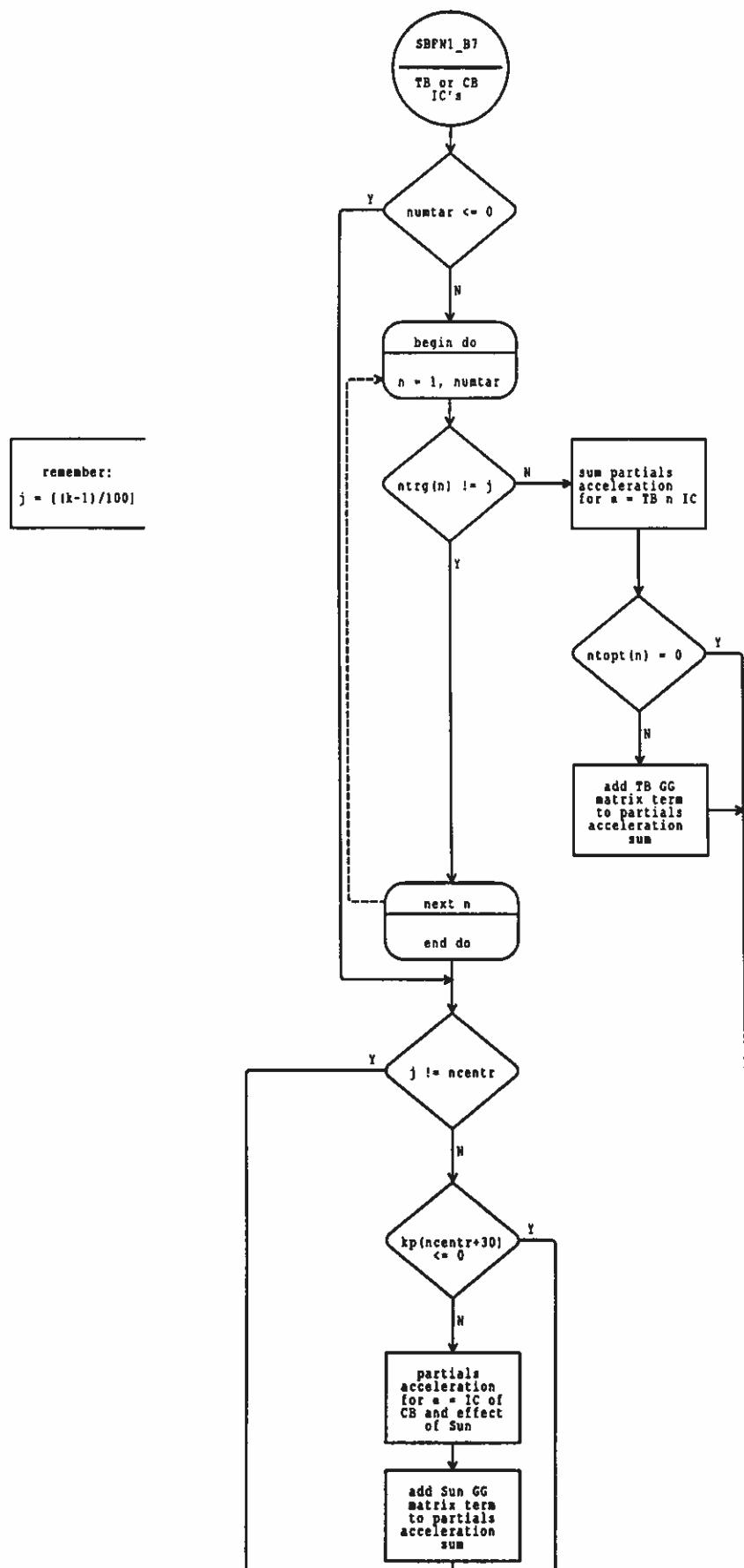
Sun-centered probe

Earth orbiter

Lunar orbiter

pl:  
o:

natural  
planetary  
satellite



planet perturbation  
depends on IB  
relative to planet,  
which depends on CB  
relative to planet

add Sun  
effects to  
planet  
partials sum

begin do

n = 1, 10

kp(n+30) <= 0

Y

N

n = ncentr

Y

N

add pert.  
planet effect  
to partials  
sum

next n

end do

add pert planet  
GG term to  
partials  
acceleration  
sum

numtar <= 0

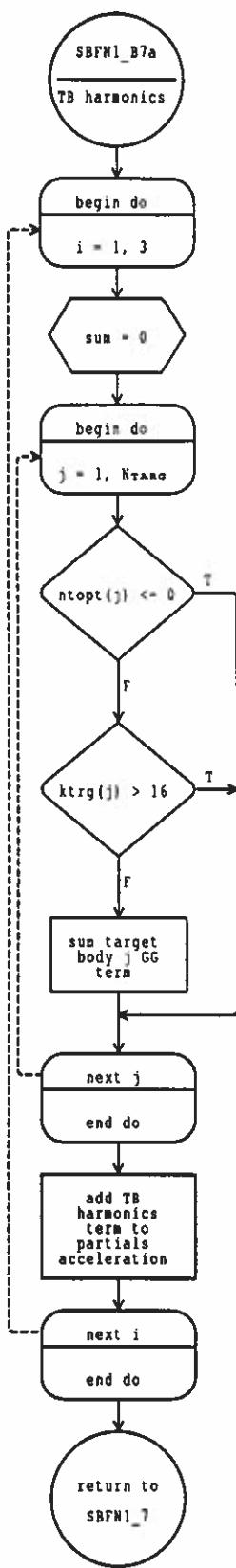
Y

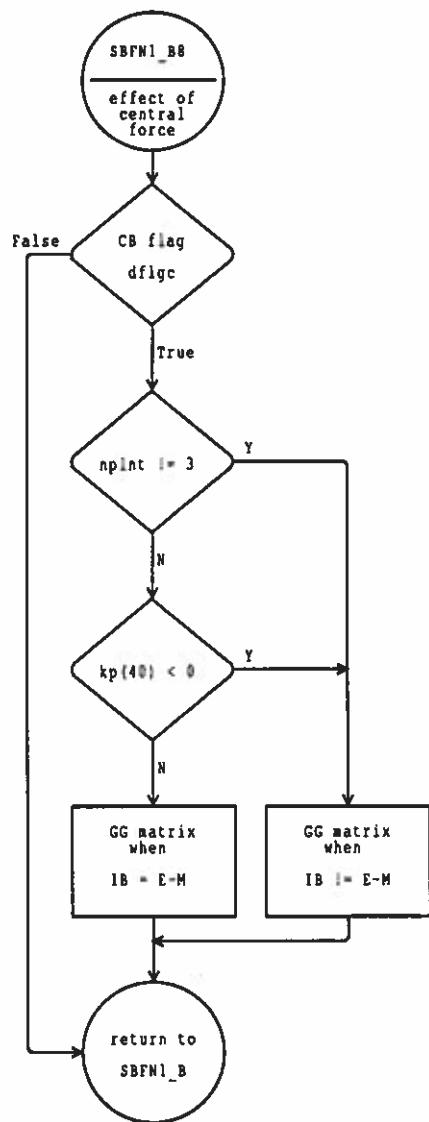
N

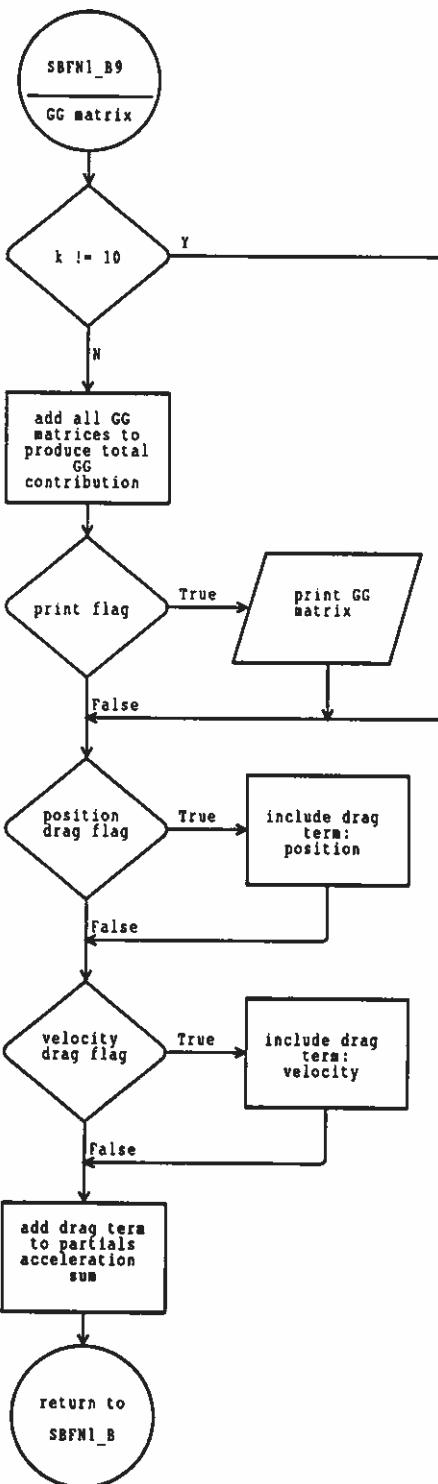
SBPN1\_B7a  
---  
TB  
harmonics

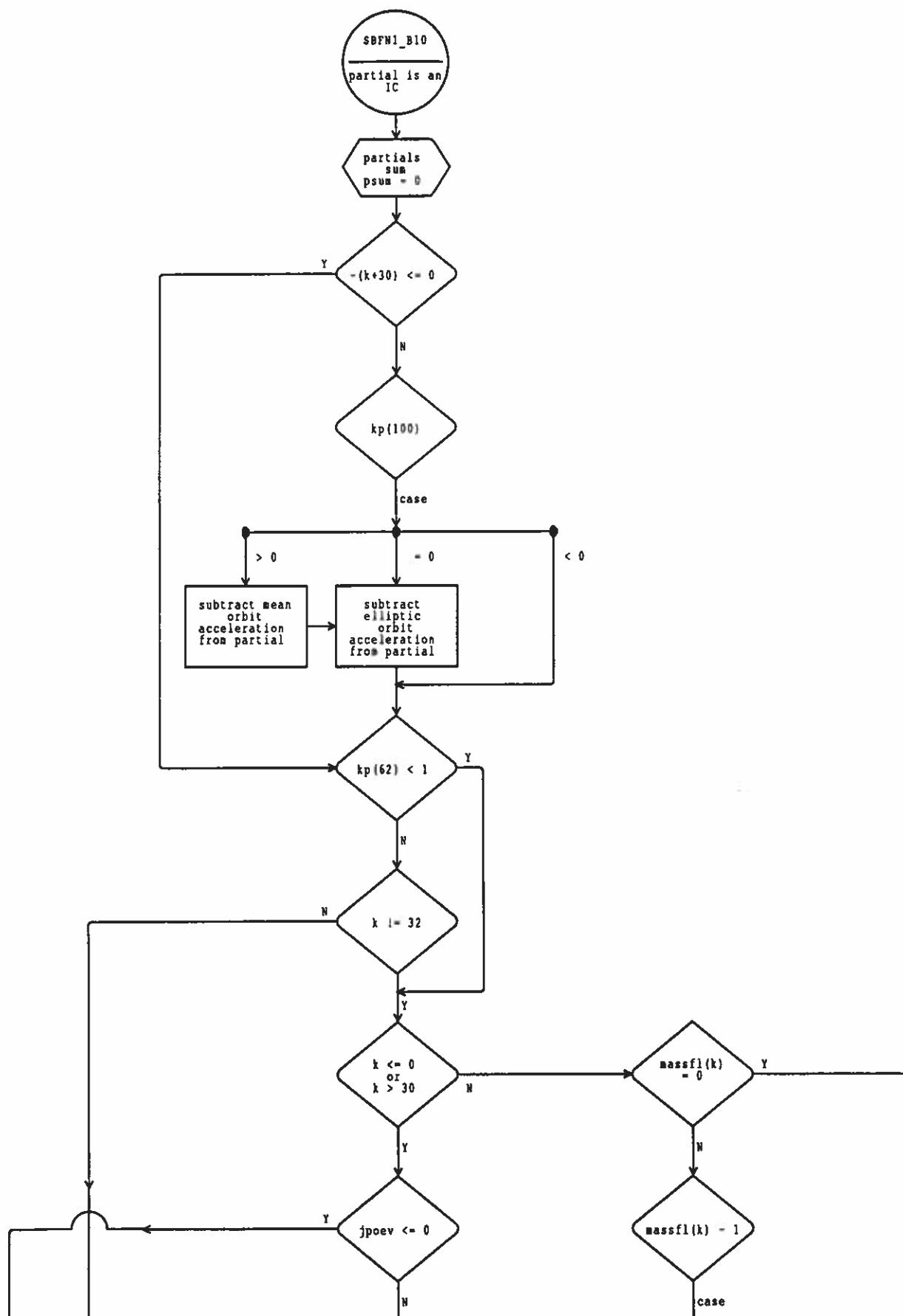
add sum to  
partials  
acceleration  
sum

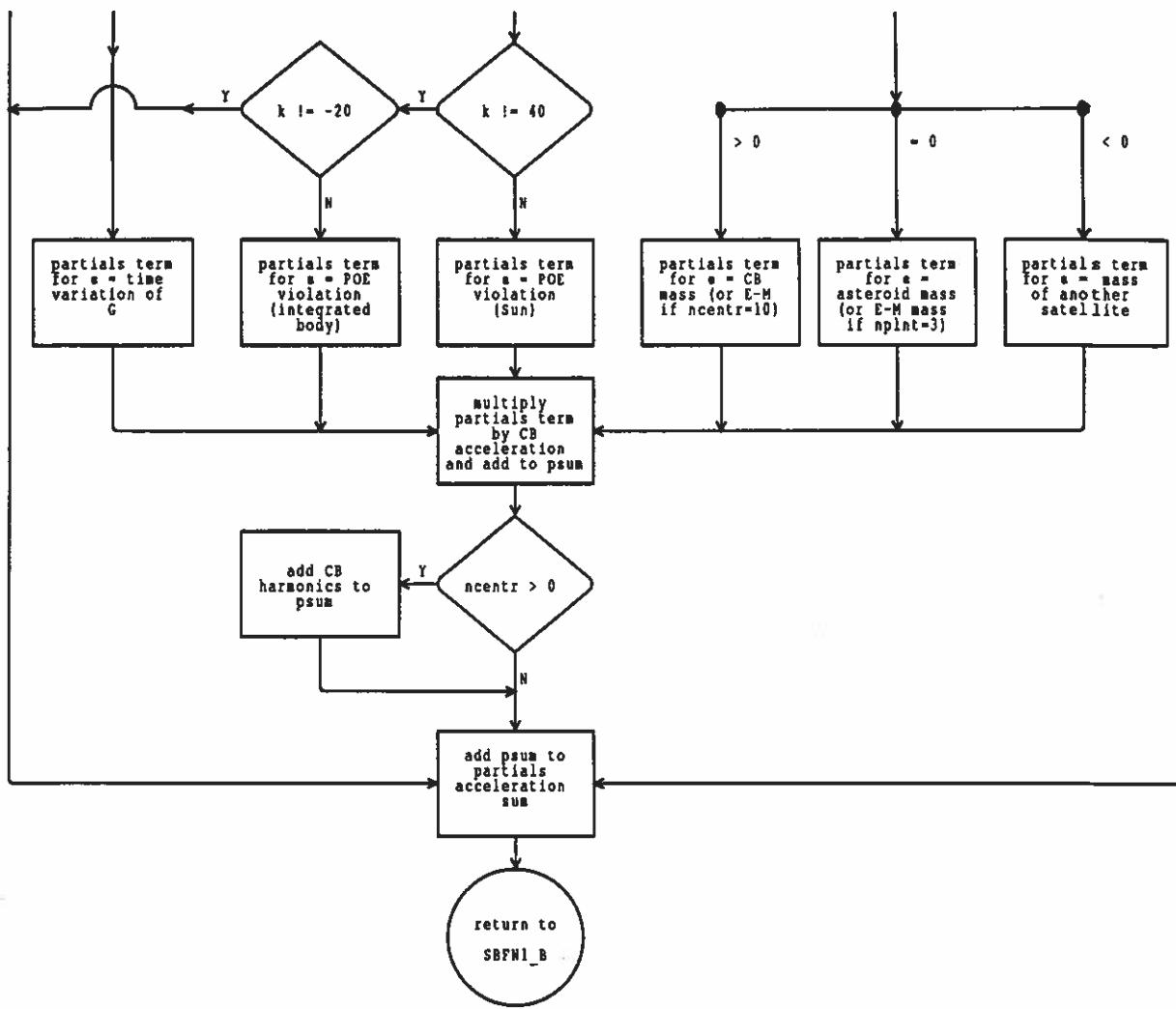
return to  
SBPN1\_B











# **Subroutine Charts**

