

QUANTITATIVE STABILITY INEQUALITIES FOR VLASOV-POISSON AND HMF MODELS

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General motivation

- Get quantitative stability results of steady states f_0

$$f_0(x) = F(\sigma_0(x)), \quad \text{with monotonic } F,$$

for evolution equations that essentially **preserve the measure** (or the rearrangement) and **some functional** (energy, momentum, etc):

$$\int G(f(t, x)) dx = \int G(f(0, x)) dx, \quad \forall G. \quad f(t)^* = f(0)^*$$

$$\mathcal{H}(f(t)) = \mathcal{H}(f(0)).$$

- Quantitative means that the deviation at time t is controlled by the initial deviation from a steady state q :

$$\|f(t) - f_0\| \leq C \|f(0) - f_0\|, \quad \forall t. \quad \text{up to symmetries}$$

Examples are: **Vlasov-Poisson, HMF and 2D Euler systems.**

Most of the known stability results in these contexts are obtained via **compactness arguments.**

Two key estimates

- Goal here: Quantitative functional inequality of the generic form (up to symmetries of the system ...)

$$\|f - f_0\|_{L^1}^2 \leq C_1 \mathcal{F}(f^*, f_0^*) + C_2 (\mathcal{H}(f) - \mathcal{H}(f_0)).$$

Minimal regularity assumptions on the perturbation. For **compactly supported** steady states f_0 :

$$\mathcal{F}(f^*, f_0^*) \leq C \|f^* - f_0^*\|_{L^1}.$$

- First rearrangement inequality \implies **Stability criteria**, under which one has

$$\int \sigma_0(x) [f(x) - f_0(x)] dx \leq \mathcal{H}(f) - \mathcal{H}(f_0) + \mathcal{F}_1(f^*, f_0^*)$$

with $\mathcal{F}_1(f, f) = 0$, up to symmetries of the problem.

- Refined rearrangement inequality \implies **Stability estimate**:

$$\|f - f_0\|_{L^1}^2 \leq \int \sigma_0(x) [f(x) - f_0(x)] dx + \mathcal{F}_2(f^*, f_0^*). \quad \text{with} \quad \mathcal{F}_2(f, f) = 0 \quad \forall f.$$

- A generalized notion of rearrangement
- An extended Hardy-Littlewood type inequality: leading to a stability criteria on the steady state.
- A refined Hardy-Littlewood type inequality (RHL) for rearrangements: leading to a quantitative control of the whole distribution function.
- Applications to VP and HMF

Equimeasurability and Schwarz rearrangement

- **Equimeasurability:** two nonnegative functions $f, g \in L^1(\Omega)$ are equimeasurable if and only if:

$$\int_{\Omega} C(f(x)) dx = \int_{\Omega} C(g(x)) dx, \quad \forall C$$

or equivalently: $\forall \lambda \geq 0$,

$$\mu_f(\lambda) = \text{meas}\{x \in \Omega; f(x) > \lambda\} = \text{meas}\{x \in \Omega; g(x) > \lambda\} = \mu_g(\lambda).$$

We denote by $\text{Eq}(f)$ the set of functions equimeasurable with f .

- **The standard Schwarz symmetrization.** Let $f \in L^1(\Omega)$, then there exists a unique function $f^* \in L^1(\Omega)$ which is nonincreasing function of $|x|$ such that f^* is equimeasurable with f .
- In other terms: if $f^\#$ is the pseudo inverse of μ_f , then

$$f^*(x) = f^\#(\text{meas}(B_d(0, |x|) \cap \Omega)), \quad \forall x \in \Omega.$$

f^* is essentially the unique decreasing function of $|x|$ which is equimeasurable with f .

The Hardy-Littlewood inequality

Hardy, Littlewood, Pólya: Inequalities, 1934, Cambridge press.

Lieb and Loss: Analysis.

Burton, Annales de l'IHP (1987), Math. Ann. (1989)

Let Ω be a measurable subset of \mathbb{R}^d and let f and g be two nonnegative measurable functions on Ω . Then

$$\int_{\Omega} f(x)g(x)dx \leq \int_{\Omega} f^*(x)g^*(x)dx,$$

In particular

$$\int_{\Omega} |x|(f(x) - f^*(x))dx \geq 0.$$

A well known tool to prove symmetry properties of minimizers.

Generalized rearrangement

Let σ be a nonnegative measurable function of $\Omega \subset \mathbb{R}^d$, $d \geq 1$ such that for all $e \in [0, e_{\max})$

$$\text{meas}\{x \in \Omega, \sigma(x) = e\} = 0.$$

Let

$$a_\sigma(e) = \text{meas}\{x \in \Omega, \sigma(x) < e\}, \quad a_\sigma(e_{\max}) = |\Omega|.$$

For all $f \in L^1(\Omega)$, we define its rearrangement $f^{*\sigma}$ with respect to σ by

$$f^{*\sigma}(x) = f^\#(a_\sigma(\sigma(x))), \quad \forall x \in \Omega,$$

In particular $f^{*\sigma}$ is a decreasing function of $\sigma(x)$ and is equimeasurable with f .

Extended Hardy-Littlewood inequality

Let σ be as above. Then for any nonnegative functions $f, g \in L^1(\Omega)$ we have

$$\int_{\Omega} f(x)g(x)dx \leq \int_{\Omega} f^{*\sigma}(x)g^{*\sigma}(x)dx,$$

In particular

$$\int_{\Omega} \sigma(x)(f(x) - f^{*\sigma}(x))dx \geq 0.$$

Does this nonnegative quantity control some strong norm $\|f - f^{*\sigma}\|$?

➤ Weak answer: Saturating the inequality \implies Compactness

if $\int_{\Omega} \sigma(x)(f_n(x) - f_n^{*\sigma}(x))dx \rightarrow 0$, and if $\|f_n^{*\sigma} - f_0\|_{L^1} \rightarrow 0$ then

$$\|f_n - f_0\|_{L^1} \rightarrow 0.$$

➤ In the same spirit as in Burchard-Guo (JFA, 2004) concerning the Riez rearrangement inequality.

Theorem: refined HL inequality (ML)

Let σ be as above and b_σ the pseudo inverse of a_σ . Then for any nonnegative function $f \in L^1(\Omega)$ we have

$$\|f - f^{*\sigma}\|_{L^1}^2 \leq K(f^*, \sigma) \int_{\Omega} \sigma(x)(f(x) - f^{*\sigma}(x))dx,$$

where $K(f^*, \sigma)$ is a constant depending only on f^* and σ . More generally, for any nonnegative $f, f_0 \in L^1(\Omega)$

$$\begin{aligned} (\|f - f_0^{*\sigma}\|_{L^1} + \|f_0\|_{L^1} - \|f\|_{L^1})^2 &\leq K(f_0^*, \sigma) \left[\int_{\Omega} \sigma(x)(f(x) - f_0^{*\sigma}(x))dx \right. \\ &\quad \left. + \int_{\Omega} b_\sigma[2\mu_{f_0}(s)]\beta_{f^*, f_0^*}(s)ds \right], \end{aligned}$$

with $\beta_{f,g}(s) = \text{meas}\{x \in \Omega : f(x) \leq s < g(x)\}$.

A particular case:

Case of Schwarz symmetrization:

Corollary

For all $f \in L^1(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$, $d \geq 1$, and all $0 \leq m \leq d$, we have

$$\int_{\mathbb{R}^d} |x|^m (f(x) - f^*(x)) dx \geq K_d \|f\|_{L^\infty}^{-m/d} \|f\|_{L^1}^{-1+m/d} \|f - f^*\|_{L^1}^2,$$

$$K_d = 2^{-1+m/d} \frac{m^2}{4d^2} |B_d|.$$

This covers the Marchioro-Pulvirenti estimate used for 2D-Euler (1985): $m = 2$, and $d = 2$, and for homogeneous steady states for VP systems.

This estimate was used by [Caglioti and Rousset](#) to study long time behavior of some N particles systems (2007): homogeneous steady states to regularized VP, and Euler 2D.

First example: Vlasov-Poisson equations

$$\partial_t f + v \cdot \nabla_x f - \nabla_x \phi_f \cdot \nabla_v f = 0, \quad f(t=0, x, v) = f_0(x, v)$$

$$\phi_f(t, x) = -\frac{\gamma}{4\pi} \int_{\mathbb{R}^3} \frac{\rho_f(t, y)}{|x - y|} dy, \quad \rho_f(t, x) = \int_{\mathbb{R}^3} f(t, x, v) dv.$$

Poisson equation: $\Delta \phi_f = \gamma \rho_f$.

- Gravitational systems, $\gamma = +1$: galaxies, star clusters, etc.
- Systems of charged particles, $\gamma = -1$: Coulomb interactions. One has to bound the space domain or to confine the system by adding a given potential.

- Conservation of the **energy**: $\mathcal{H}(f) = E_{kin}(f) - \gamma E_{pot}(f)$

$$E_{kin}(f) = \frac{1}{2} \int_{\mathbb{R}^6} |v|^2 f dx dv, \quad E_{pot}(f) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla_x \phi_f|^2 dx$$

- Conservation of the **Casimir functionals** $\int_{\mathbb{R}^6} G(f) dx dv$.
- **Isotropic galactic models**:

$$f_0(x, v) = F \left(\frac{|v|^2}{2} + \phi_{f_0}(x) \right).$$

- Relative **Hamiltonian**:

$$\mathcal{H}(f) - \mathcal{H}(f_0) = \int_{\mathbb{R}^6} \left(\frac{|v|^2}{2} + \phi_0(x) \right) (f - f_0) dx dv - \frac{\gamma}{2} \int_{\Omega} |\nabla \phi_f(x) - \nabla \phi_0(x)|^2 dx$$

Known stability results: Vlasov-Poisson system

- Linear stability **Physics literature**: Gardner, Antonov, Lynden-Bell (1960'), , Kandrup-Signet (1980'), Aly-Perez (1990'), Chavanis (2000'), ..., see also Binney-Tremaine. Most rigorous result: Doremus-Baumann-Feix (1970'). Key assumption: **F decreasing**.
- Non linear stability: Essentially a **Mathematics literature** in the two last decades.
 - Stability of a subclass of steady states under general perturbations: Wolansky, Guo, Rein, Dolbeault, Lin, Hadzic, Sanchez, Soler, Lemou-Méhats-Raphaël, Rigault ...
 - Stability of the whole class of steady states $Q_0(x, v) = F \left(\frac{|v|^2}{2} + \phi_0(x) \right)$, with F decreasing: Lemou-Méhats-Raphaël 2012.
 - **A different important context**: If periodic domain in space: Homogeneous steady states:

$$f(x, v) = g_0(|v|).$$

Asymptotic stability under Penrose conditions and regularity assumptions: **Landau damping**, Mouhot-Villani 2011.

Statement of stability inequalities for VP

The energy space

$$\mathcal{E} = \{f \in L^\infty : f \geq 0, \|(1 + |\mathbf{v}|^2)f\|_{L^1} < \infty\}.$$

Theorem: Quantitative stability (ML).

We have the following

- i) There exist a constant $K_0 > 0$ depending only on f_0 such that for all $f \in \mathcal{E}$

$$\begin{aligned} \|f - f_0\|_{L^1} &\leq \|f^* - f_0^*\|_{L^1} + \\ K_0 \left[\mathcal{H}(f) - \mathcal{H}(f_0) + 2|\phi_{f_0}(0)| \|f^* - f_0^*\|_{L^1} + \|\nabla \phi_f - \nabla \phi_{f_0}\|_{L^2}^2 \right]^{1/2}. \end{aligned}$$

- ii) There exist constants $K_0, R_0 > 0$ depending only on f_0 such that, for all $f \in \mathcal{E}$ satisfying

$$\inf_{z \in \mathbb{R}^3} (\|\phi_f - \phi_{f_0}(\cdot - z)\|_{L^\infty} + \|\nabla \phi_f - \nabla \phi_{f_0}(\cdot - z)\|_{L^2}) < R_0,$$

there holds:

$$\begin{aligned} \|f - f_0(\cdot - z_{\phi_f})\|_{L^1} + \|\nabla \phi_f - \nabla \phi_{f_0}(\cdot - z_{\phi_f})\|_{L^2} &\leq \|f^* - f_0^*\|_{L^1} + \\ &K_0 [\mathcal{H}(f) - \mathcal{H}(f_0) + K_0 \|f^* - f_0^*\|_{L^1}]^{1/2}. \end{aligned}$$

Second example: HMF model

The Hamiltonian Mean Field (HMF) model is a kinetic model describing particles on a unit circle interacting via an infinite range attractive cosine potential: $f(t, \theta, v)$, $\theta \in [0, 2\pi]$, $v \in \mathbb{R}$.

$$\partial_t f + v \partial_\theta f - \partial_\theta \phi_f \partial_v f = 0, \quad f(t=0, \theta, v) = f_0(\theta, v)$$

$$\phi_f(t, \theta) = - \int_0^{2\pi} \rho_f(t, \theta') \cos(\theta - \theta') d\theta', \quad \rho_f(t, \theta) = \int_{\mathbb{R}} f(t, \theta, v) dv.$$

➤ Magnetization vector:

$$M_f = \int_0^{2\pi} \rho_f(\theta) u(\theta) d\theta, \quad u(\theta) = (\cos \theta, \sin \theta)^T.$$

$$\phi_f(\theta) = -M_f \cdot u(\theta).$$

➤ Hamiltonian:

$$\mathcal{H}(f) = \frac{1}{2} \iint v^2 f(\theta, v) d\theta dv - \frac{1}{2} |M_f|^2.$$

- Consider the following class of steady states

$$f_0(\theta, v) = F(e_0(\theta, v)), \quad \text{with} \quad e_0(\theta, v) = \frac{v^2}{2} + \phi_0(\theta),$$

$$\phi_0(\theta) = -m_0 \cos \theta, \quad m_0 > 0.$$

Note that f_0 may have an **unbounded support**.

- The relative Hamiltonian writes

$$\mathcal{H}(f) - \mathcal{H}(f_0) = \iint \left(\frac{v^2}{2} + \phi_0(x) \right) (f - f_0) d\theta dv - \frac{1}{2} |M_f - M_{f_0}|^2$$

- No stability result in this case. **Formal linear stability criteria** has been given by **Barré et al 2011** and more explicitly by **Ogawa, 2013**.

Non linear stability Criteria for HMF

- Steady states of the form

$$f_0(\theta, v) = F(e_0(\theta, v)), \quad e_0(\theta, v) = \frac{v^2}{2} - m_0 \cos \theta, \quad m_0 > 0.$$

- Assume F is **decreasing** and let

$$\kappa_0 = \int_0^{2\pi} \int_{-\infty}^{+\infty} |F'(e_0(\theta, v))| \left(\frac{\int_{\mathcal{D}} (\cos \theta - \cos \theta') (e_0(\theta, v) - \phi_0(\theta'))^{-1/2} d\theta'}{\int_{\mathcal{D}} (e_0(\theta, v) - \phi_0(\theta'))^{-1/2} d\theta'} \right)^2 d\theta dv.$$

We prove that if $\kappa_0 < 1$, then f_0 is orbitally stable.

The same criteria has been obtained formally for the linear stability by Barré et al (2011) and more explicitly by Ogawa (2013).

- f_0 may have an **unbounded support**, but in this case we assume

$$\|f_0\|_{L^p} < +\infty, \quad \forall 0 < p < 1/3.$$

Non linear stability for HMF (L, Luz, Méhats)

There exists $\delta > 0$ such that, for all $f \in L^1((1 + |v|^2)d\theta dv)$ satisfying $|M_f - M_{f_0(\cdot - \theta_f)}| < \delta$, we have

$$\begin{aligned} \|f - f_0(\cdot - \theta_f)\|_{L^1}^2 \leq & C \left(\mathcal{H}(f) - \mathcal{H}(f_0) + C(1 + \|f\|_{L^1}) \|f^* - f_0^*\|_{L^1} \right. \\ & \left. + C \int_0^{+\infty} s^2 \left(f_0^\sharp(s) - f^\sharp(s) \right)_+ ds + C \int_0^{+\infty} \mu_{f_0}(s)^2 \beta_{f^*, f_0^*}(s) ds \right), \end{aligned}$$

where

$\beta_{f^*, f_0^*}(s) = \text{meas} \{(\theta, v) \in [0, 2\pi] \times \mathbb{R} : f^*(\theta, v) \leq s < f_0^*(\theta, v)\}$, for all $s \geq 0$, and where C is a positive constant depending only on f_0 . The parameter θ_f is defined by $M_f = |M_f|(\cos \theta_f, \sin \theta_f)^T$. In particular, if f_0 is a compactly supported steady state, then

$$\|f - f_0(\cdot - \theta_f)\|_{L^1}^2 \leq C \left(\mathcal{H}(f) - \mathcal{H}(f_0) + C(1 + \|f\|_{L^1}) \|f^* - f_0^*\|_{L^1} \right).$$

Main lines of the proof: i) reduce to a functional of M_f .

We write the relative Hamiltonian in the following form ($\phi_f(\theta) = M_f \cdot u(\theta)$)

$$\mathcal{D} = \mathcal{H}(f) - \mathcal{H}(f_0) = \iint \left(\frac{v^2}{2} + \phi_f(x) \right) (f - f_0) d\theta dv + \frac{1}{2} |M_f - M_0|^2.$$

We take : $\sigma(\theta, v) = \frac{v^2}{2} + \phi + \max \phi$ and denote $f^{*\sigma} = f^{*\phi}$. Note that:
 $f_0 = f_0^{*\phi_0}$.

$$\begin{aligned} \mathcal{D} = & \iint \left(\frac{|v|^2}{2} + \phi + \max \phi \right) (f - f^{*\phi}) d\theta dv + \iint \left(\frac{v^2}{2} + \phi \right) (f^{*\phi} - f_0^{*\phi}) d\theta dv \\ & + \iint \left(\frac{v^2}{2} + \phi \right) (f_0^{*\phi} - f_0) d\theta dv + \frac{1}{2} |M_f - M_0|^2. \end{aligned}$$

We get

$$\begin{aligned} \mathcal{J}(|M_f|) - \mathcal{J}(|M_{f_0}|) & \leq \mathcal{H}(f) - \mathcal{H}(f_0) + \\ & \int_0^{+\infty} s^2 \left(f_0^\sharp(s) - f^\sharp(s) \right)_+ ds + C(1 + \|f\|_{L^1}) \|f^* - f_0^*\|_{L^1}. \end{aligned}$$

$$\mathcal{J}(|M_f|) = \iint \left(\frac{v^2}{2} + \phi \right) f_0^{*\phi} d\theta dv + \frac{1}{2} |M_f|^2.$$

ii) Stability criteria for HMF

- The stability criteria: the quadratic form

$$\mathcal{J}''(m_0) = 1 - \kappa_0 > 0, \quad \phi_0(\theta) = -m_0 \cos \theta.$$

- In case of VP, this leads to a criteria which was proved to hold true in L-Méhats-Raphaël, 2012, thanks to a new **Poincaré type inequality**.
- We have (locally)

$$\mathcal{J}(m) - \mathcal{J}(m_0) \geq C(m - m_0)^2, \quad \forall m > 0.$$

iii) Control of $\|f - f_0\|_{L^1}$

Now we first write the relative Hamiltonian in the following form

$$\mathcal{H}(f) - \mathcal{H}(f_0) = \iint \left(\frac{v^2}{2} + \phi_0(\theta) \right) (f - f_0) d\theta dv - \frac{1}{2} |M_f - M_{f_0}|^2.$$

$$\sigma(x, v) = \frac{v^2}{2} + \phi_0(\theta) + m_0, \quad f^{*\sigma} \equiv f^* \phi_0, \quad f_0^{*\phi_0} = f_0.$$

$$a_\sigma(e) = \int_0^{2\pi} (e + m_0 \cos \theta)_+^{1/2} d\theta.$$

$$\begin{aligned} \|f - f_0\|_{L^1}^2 \leq & C \left(\mathcal{H}(f) - \mathcal{H}(f_0) + C(1 + \|f\|_{L^1}) \|f^* - f_0^*\|_{L^1} + \frac{1}{2} |M_f - M_0|^2 \right) \\ & + C \int_0^{+\infty} a_\sigma^{-1}(2\mu_0(s)) \text{meas}\{f^* \leq s < f_0^*\} ds. \end{aligned}$$

There exists $\delta > 0$ such that:

If $f(t)$ is the solution the HMF model with initial data f_{init} then we have

$$\begin{aligned} \|f(t) - f_0(\cdot - \theta_{f(t)})\|_{L^1}^2 &\leq C \left[\mathcal{H}(f_{init}) - \mathcal{H}(f_0) + C(1 + \|f_{init}\|_{L^1}) \|f_{init}^* - f_0^*\|_{L^1} \right. \\ &\quad \left. + \int_0^{+\infty} s^2 \left(f_0^\#(s) - f_{init}^\#(s) \right)_+ ds + \int_0^{+\infty} \mu_0(s)^2 \text{meas}\{f_{init}^* \leq s < f_0^*\} ds. \right] \end{aligned}$$

provided that

$$|M_{f_{init}} - M_{f_0(\cdot + \theta_0)}| < \delta.$$

THANK YOU !