4.2 The Neumann Condition We now consider:

We now consider either the wave equation:

$$\begin{cases} u_{tt}(x,t) = c^2 u_{xx}(x,t) & o < x < l \ t > 0 \\ u_{x}(0,t) = u_{x}(l,t) = 0 & t > 0 \\ u_{x}(0,t) = \phi(x) & u_{t}(x,0) = \psi(x) & o < x < l \end{cases}$$

or the diffusion equation:

$$\begin{cases} u_{t}(x,t) = k u_{xx}(x,t) & o < x < l & t > 0 \\ u_{t}(x,t) = u_{t}(t,t) = 0 & t \ge 0 \\ u_{t}(x,0) = \phi(x) & o < x < l \end{cases}$$

Notice that now the boundary conditions involve nx rather than n!

Using separation of variables u(x,t) = X(x) T(t) we reach the same equations for X and T as before.

X part: As before, we have $X''(x) + \beta^2 X(x) = 0$ which leads to solutions of the form:

$$X(x) = C cas(\beta x) + D siz(\beta x)$$

Let's write the derivative of this, which we will need:

$$X'(x) = -C\beta s'(\beta x) + D\beta cos(\beta x)$$
.

$$u_{\mathbf{x}}(0,t)=0 \rightarrow \mathbf{x}'(0)=0 \rightarrow -Cp \sin 0 + Dp \cos 0 = 0$$

$$u_{\chi}(\ell,t)=0 \Rightarrow \chi'(\ell)=0 \Rightarrow -Cp\sin(p\ell)=0$$

$$\beta l = n\pi \qquad \Rightarrow \qquad \beta_n = \frac{n\pi}{\ell}$$

$$\lambda_n = \left(\frac{n\pi}{\ell}\right)^2 \qquad n = 1, 2, ...$$

$$X_n(x) = \cos\left(\frac{n\pi}{\ell}x\right)$$

These are the EIGENFUNCTIONS
for the Neumann problem

Com we have eigenvalues that are not positive? I.e. com we solve $-X''(x) = \lambda X(x)$ ocxel with the boundary cond: X'(0) = X'(0) = 0 and with $\lambda \in \mathbb{C}$ which is not positive?

Try $\lambda=0$: We get X''(x)=0 so that X(x)=C+Dx, X(x)=DApply BCs: 0=X'(0)=X'(1)=D.

The can satisfy the BCs with D=0. X(x)=C (constant)

is a legitimete solution! X(x)=C an eigenvalue!

 $\lambda < 0$ or $\lambda \in \mathbb{C} \setminus \mathbb{R}$: It can be shown that such values of λ cannot be eigenvalues, but we skip that for now.

So the eigenvalues are:

$$\lambda_n = \left(\frac{n\pi}{\varrho}\right)^2 \qquad n = 0, 1, 2, \dots$$

These are the EIGENVALUES for the Neumann problem

The T(t) part will be identical to what we saw before, with the exception of the part cowing from $\lambda = 0$.

Diffusion equation: For $\lambda_n \neq 0$ we again have:

$$T'(t) = -\lambda_n k T(t)$$
 $T(t) = Ae^{-\lambda_n k t}$

For
$$\lambda=0$$
 we have $T(t)=0$ \longrightarrow $T(t)=A$.

So, for n=1,2,3,... we have as before:

$$u_n(x,t) = A_n e^{-(\frac{n\pi}{4})^2 kt} \cos(\frac{n\pi}{4}x)$$
 $n=1,2,...$

Notice that the sine is now a costne!

And we also have a up now: the spatial part is $(2\pi)^{-1} = 1$ and the temporal part is a constant which we called A above. For reasons which well become clear, we call $A_0 = 2A$, to find:

$$u(x,t) = \frac{1}{2}A_0 + \frac{2}{n}A_n e^{-\left(\frac{n\pi}{\ell}\right)^2kt} cos(\frac{n\pi}{\ell}x)$$

In addition, the initial condition will have to satisfy:

$$\phi(x) = u(x, 0) = \frac{1}{2}A_0 + \frac{2}{n}A_n \cos(\frac{n\pi}{\ell}x)$$

Wave equation: For $\lambda > 0$ we get the same behavior as we've seen before, so we have:

$$u_n(x,t) = \left[A_n \cos\left(\frac{n\pi}{\ell}et\right) + B_n \sin\left(\frac{n\pi}{\ell}et\right)\right] \cos\left(\frac{n\pi}{\ell}x\right)$$

For $\lambda = 0$ we get $X_0(x) = const$ as before. For Re T part we lowe $T''(t) = \lambda c^2 T(t) = 0$ a that $T_0(t) = A + Bt$. This To term goes with the X_0 term which is a constant. So, to conclude, the general solution has the Jorn:

$$n(x,t) = \frac{1}{2}A_0 + \frac{1}{2}B_0t + \sum_{n}[A_{n}\cos(\frac{n\pi}{\ell}ct) + B_{n}\sin(\frac{n\pi}{\ell}ct)]\cos(\frac{n\pi}{\ell}x)$$

$$\phi(x) = u(x,0) = \frac{1}{2}A_0 + \frac{2}{n}A_n\cos(\frac{n\pi}{\ell}x)$$

$$\psi(x) = u_{t}(x, 0) = \frac{1}{2}B_{0} + \frac{\pi}{n}\frac{\pi}{\ell}CB_{n}\cos(\frac{n\pi}{\ell}x)$$