

Annual report research summary

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I. STATISTICAL HADRONIZATION MODELS

Heavy-ion collisions produce a hot, dense state of QCD matter commonly known as the quark-gluon plasma (QGP). The QGP then expands and freezes out into discrete particles. If we assume that the QGP is in thermal equilibrium prior to freeze-out, we can use pure statistical mechanics to calculate particle yields. Calculated yields may be fit to experimental data to extract thermal parameters of the post-QGP fireball such as temperature and chemical potential.

This approach is generally consistent with experiments at roughly the 10% level. However, more precise data have shown that particle production is not purely statistical. For example, multiplicities of particles that contain strange quarks do not always match thermal equilibrium values. Statistical models accommodate this by including non-equilibrium parameters, which must also be included in fits to data. For example the strangeness occupancy γ_s describes strangeness enhancement: $\gamma_s = 1$ implies equilibrium, $\gamma_s > 1$ enhancement, $\gamma_s < 1$ suppression.

Currently, there are a number of competing statistical models, and no consensus about the best-fit values of the various parameters, or even which parameters should be included. The Duke group plans to help resolve this dispute using SHARE, a flexible statistical model. Figure 1 shows a sample least-squares fit to hadron yield ratios from the STAR experiment at RHIC. The fit is mostly consistent with data, though the error bars are large.

Standard least-squares fits are susceptible to convergence problems and don't always provide enough information about the quality of the fit. We apply more advanced statistical techniques, namely Markov chain Monte Carlo (MCMC), which selects sets of parameters

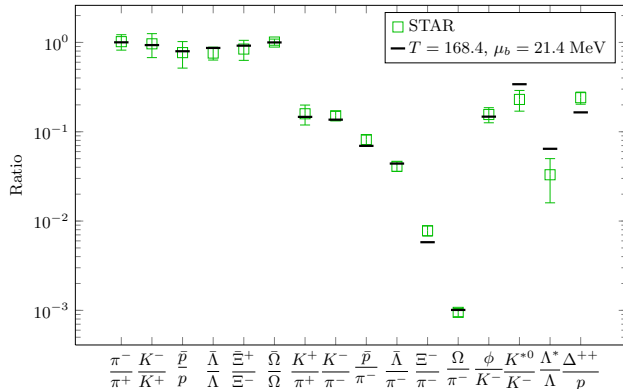


FIG. 1: Least-squares fit to identified particle ratios from STAR.

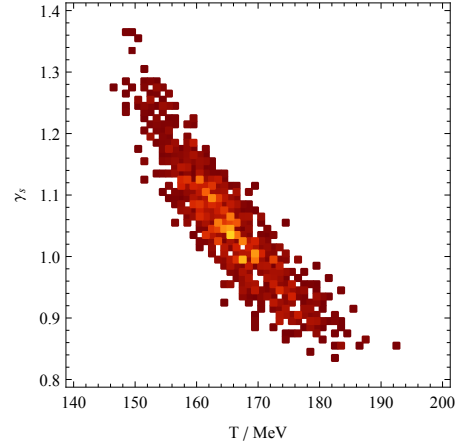


FIG. 2: MCMC analysis of T - γ_s dependence using STAR data. Colors indicate the number of entries in the Markov chain in that bin: dark red for low counts up to yellow for large counts. The preferred region is in the middle of the “island”, at $T \sim 165$ MeV and $\gamma_s \sim 1.05$.

using a Metropolis algorithm and then checks the model's agreement with data. This method should provide unbiased information about which parameters are most important and their ideal values. Figure 2 shows a sample MCMC result from varying temperature and γ_s and comparing to the same STAR data as in Fig. 1. These preliminary results are consistent with conventional statistical models, as expected. We anticipate new results this spring.

II. ONION VISUALIZATION

It is well-known that hadrons with more strange quarks have smaller scattering cross sections, and therefore cease to interact (freeze out) earlier in a post-QGP expansion. We seek to visualize the freeze-out surfaces of several baryons with different strangeness: Ω , Ξ , Λ , and the proton. The Ω baryon consists of three strange quarks ($s = 3$), therefore it should interact the least and its freeze-out surface should be smallest. The Ξ ($s = 2$) should have a somewhat larger surface, followed by the Λ ($s = 1$) and the proton ($s = 0$). When visualized, these surfaces will take on a layered appearance, hence the name “onion visualization”.

Figure 3 shows freeze-out time distributions sorted by strangeness from pure UrQMD simulations. It is clear that species with more strangeness generally freeze out earlier. Spatial distributions show similar trends. We expect to have three-dimensional visualizations this spring.

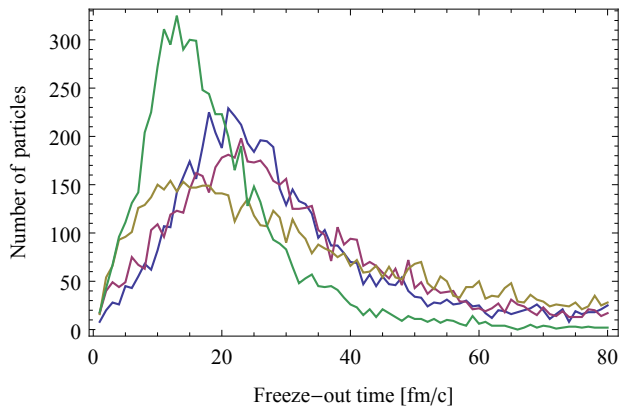


FIG. 3: Freeze-out time distributions for equal numbers of protons (blue), Λ (purple), Ξ (tan), and Ω (green).

Further work may repeat the analysis with a more realistic model including a hydrodynamic phase.

III. EVENT-BY-EVENT HEAVY-ION COLLISION SIMULATIONS

One of the major goals of heavy-ion physics is to characterize the physical properties of the QGP, perhaps most notably the shear viscosity. Since it is impossible to observe the QGP directly, the only way to accomplish this is to compare experimental data to computational models.

A complete simulation of a heavy-ion collision has many components and generally takes at least a few hours on a modern CPU. Several hours would not be a problem if the simulation only needed to be carried out once, but every collision is unique, so we cannot realistically model fluctuations with “average” events—many complete events must be simulated. Studies such as this are typically called “event-by-event”.

We must test many sets of input parameters in order to determine the best-fit to experiment. Each set of parameters requires many individual events in order to acquire adequate statistics. We expect that several million CPU hours will be required to obtain sufficient data.

In order to satisfy such computational demand, we run the simulations on the Open Science Grid (OSG), which can run hundreds or thousands of events simultaneously on computers across the USA.

Despite the capabilities of the OSG, we cannot run

events at every possible set of input parameters. We employ a Gaussian statistical emulator to provide approximate results for points in between explicitly simulated points. The emulator is similar to a regression or interpolation, but much more powerful: it supplies a statistical distribution for every point in parameter space, thus reflecting the true uncertainty of the model.

We will compare simulation data to event-by-event experimental data measured by the ATLAS experiment at the LHC. ATLAS has measured event-by-event flow coefficients, which were previously only available integrated over many events. The flow coefficients—which describe the spatial anisotropy of a given event—are a sensitive probe of the shear viscosity. We calculate the flow coefficients from simulated events and compare them to the corresponding experimental results. This method enables precise constraints on the shear viscosity and other model parameters.

The event-by-event model is currently running on the OSG. We expect to complete the study this summer.

IV. 3+1D HYDRO TO MICRO CONVERTER

Modern heavy-ion collision simulations feature a hydrodynamic phase—which models the physics of the QGP fluid—followed by freeze-out into discrete particles. A critical step in any heavy-ion simulation is the conversion from hydrodynamics to an ensemble of particles.

The hydro model evolves until a specified minimum temperature or energy density is reached. This occurs at many different points in spacetime; all these points together form a freeze-out hypersurface. The converter then statistically samples the hypersurface and calculates properties of the emitted particles. This is a very computationally intensive process—depending on the model, it may be the longest single step in the full event-by-event simulation.

In our current event-by-event model (see previous section), the hydro component runs in 2+1D (two spatial dimensions and time). The converter is designed to read a 2+1D hypersurface and produce a 3+1D particle ensemble.

Eventually, we will replace the 2+1D hydro model with a new 3+1D code, which will also require a new converter. It must run very efficiently to be viable in an event-by-event context.

This project is still in preliminary stages. We expect significant progress this summer and fall.