

# Periodicity in Exponential Function Differences

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## Abstract

The taylor series  $\sum_{m=0}^{\infty} \frac{x^{a_m}}{a_m!}$  is convergent for set of terms identified by  $a_m$  an infinite increasing sequence of positive integers. These functions all behave similar to the exponential function  $e^x$ . In particular, if the sequence  $a_m$  is an arithmetic progression with common difference  $d$  and first term  $a_0 < d$ , then the function  $f(x)$  defined by this taylor series is it's own  $d^{\text{th}}$  derivative.

## Self Derivative Series

Define the function  $\mathcal{G}_d^i$  with  $d \geq 1$  and  $0 \leq i < d$  to be the  $i - \text{th}$  offset Taylor series as

$$\mathcal{G}_d^i(x) = \sum_{n=0}^{\infty} \frac{x^{dn+i}}{(dn+i)!}.$$

Observe the following familiar functions can be written in terms of these building blocks:

$$e^x = \mathcal{G}_1^0(x)$$

$$\cosh(x) = \mathcal{G}_2^0(x) \text{ and } \sinh(x) = \mathcal{G}_2^1(x),$$

$$\cos(x) = \mathcal{G}_4^0(x) - \mathcal{G}_4^2(x) \text{ and } \sin(x) = \mathcal{G}_4^1(x) - \mathcal{G}_4^3(x).$$

By computing the  $d - \text{th}$  derivative of  $f = \mathcal{G}_d^i$  directly on the taylor series expansion we verify the claim of the abstract that  $f^{(d)} = f$ . So also we can verify the familiar claims of the above functions being their own first, second and fourth derivatives respectively.

## Other functions

If you graph the families  $\mathcal{G}_d^i$  with fixed  $d$  and  $0 \leq i < d$  we see that they cluster together by family of fixed  $d$ . In fact, it is startling looking that the graph of the 4 functions  $\mathcal{G}_4^i$  overlaid on the same graph appear to be equal for  $x \gg 0$ . The startling fact being that not only are these functions not equal but they criss cross each other as we observe writing  $\sin x$  and  $\cos x$  as above.

For what follows, it is sufficient to graph the taylor series truncated at 40 terms for  $d > 1$  and simply use  $e^x$  for  $d = 1$ . Get a feel for this by observing that

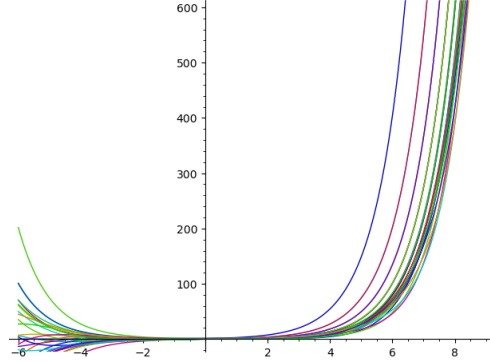


Figure 1:  $\mathcal{G}_i$  families for  $i = 1, 2, \dots, 7$

$\frac{25^{80}}{80!} < 1.0e - 7$ . It is sufficient to graph the interval  $-3 < x < 25$  to walk through the following.

The expressions for  $\sin x$  and  $\cos x$  above suggest it is reasonable to look at differences of the various  $\mathcal{G}_d^i$ . The graph below is of the 3 functions  $\mathcal{G}_3^i - \mathcal{G}_3^j$  with  $0 \leq i < j < 3$ . Observe the exponential oscillation to the left and the exponential decay to the right.

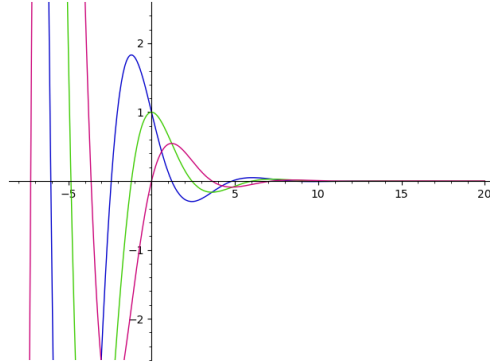


Figure 2: Unscaled graph of  $\mathcal{G}_3$  family

Now, here are these same three functions scaled by multiplying by  $e^{x/2}$ . Specifically

$$e^{x/2} \left( \mathcal{G}_3^i(x) - \mathcal{G}_3^j(x) \right).$$

Observe in this scaled version that we can easily see the periodic zeros very similar to  $\sin$  and  $\cos$ . Numerically the zeros are spaced with common difference approximately 3.628.

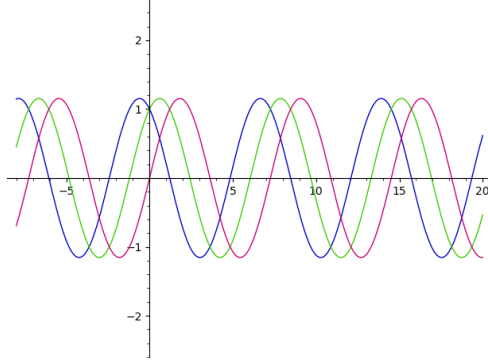


Figure 3: Scaled graph of  $\mathcal{G}_3$  family

**Conjecture 1:** There exists a periodic function  $\mathcal{H}_3^{i,j}$  such that

$$\lim_{x \rightarrow \infty} e^{x/2} \left( \mathcal{G}_3^i(x) - \mathcal{G}_3^j(x) \right) - \mathcal{H}_3^{i,j}(x) = 0.$$

However they are not mutually derivatives of each other when the scaling factor is included.

**Observation 2:** The difference  $\mathcal{G}_2^0(x) - \mathcal{G}_2^1(x)$  converges to 0 from above as  $x \rightarrow \infty$ . There is no oscillation. This is easily verified by writing  $\sinh$  and  $\cosh$  as linear combinations of  $e^x$  and  $e^{-x}$  in the usual way.

**Observation 3:** The 6 differences  $\mathcal{G}_4^i(x) - \mathcal{G}_4^j(x)$  need no scaling and are exactly periodic and mutually self-differential with a small caveat. Aside from the differences for  $\sin$  and  $\cos$ , these functions grow exponentially to the left. The difference functions are “visually periodic” for  $x > 5$ . As above in Conjecture 1 we can make this precise with a periodic limiting function  $\mathcal{H}_4^{i,j}$ .

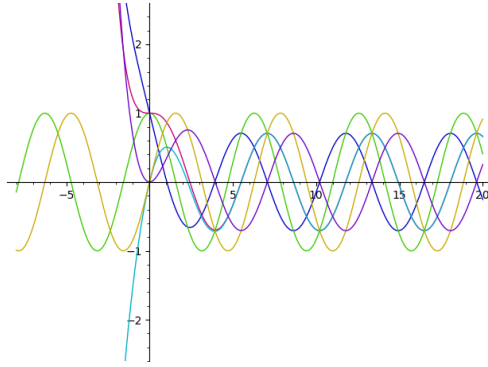


Figure 4: Unscaled graph of  $\mathcal{G}_4$  family

With correct scaling for each  $d \geq 5$ , we observe this same sort of periodic behavior for  $x \gg 0$ . The following table summarizes the numeric observation of required scaling and resulting periods and amplitudes.

d	scale	period	amplitudes			
3	0.500	7.256	1.155			
4	--	$2\pi$	0.707 1.000			
5	-0.309	6.606	0.470 0.761			
6	-0.500	7.255	0.334 0.578 0.667			
7	-0.623	8.036	0.248 0.447 0.557			
8	-0.707	8.886	0.191 0.354 0.462 0.501			
9	-0.766	9.775	0.152 0.286 0.385 0.438			
10	-0.809	10.690	0.123 0.235 0.324 0.381 0.401			
11	-0.841	11.622	0.102 0.196 0.274 0.330 0.360			
12	-0.866	12.566	0.086 0.165 0.235 0.288 0.322 0.333			

## Remaining Questions

**Question 1:** Given  $d$  can we compute the scale, period and amplitudes?

I have not been able to see any precise pattern in the scaling factor as  $d$  varies beyond the general decreasing and it possibly approaches  $-1$ .

**Question 2:** Do the functions

$$e^{\gamma x} \left( \mathcal{G}_d^i(x) - \mathcal{G}_d^j(x) \right)$$

always converge to a sinusoidal periodic function as  $x \rightarrow \infty$ ?

**Question 3:** Does

$$\left( \sum_{n=0}^{\infty} \frac{x^{nd}}{(nd)!} \right) e^{-x} \rightarrow \frac{1}{d}$$

for all  $d$  as  $x \rightarrow \infty$ ?

**Question 4 (bonus):** The abstract considered any sequence of exponents in a Taylor series. What kind of conditions on a sequence  $a_m$  with  $a_{m+1} > a_m$  enable us to make any interesting observations about  $\sum_{m=0}^{\infty} \frac{x^{a_m}}{a_m!}$ ?