

Phy 981 Project

Josh Bradt

March 9, 2015

Part (a)

The Hamiltonian used in the problem is given as

$$\hat{H} = \hat{H}_0 + \hat{V}$$

where

$$\hat{H}_0 = \xi \sum_{p\sigma} (p-1) a_{p\sigma}^\dagger a_{p\sigma} \quad (1)$$

and the matrix element of the potential is

$$\langle q+q- | V | s+s- \rangle = -g. \quad (2)$$

The form of this matrix element implies that the two-body part of the Hamiltonian can be written in an anti-symmetric form as

$$\hat{V} = -g \sum_{pq} a_{p+}^\dagger a_{p-}^\dagger a_{q-} a_{q+}. \quad (3)$$

Both parts of the Hamiltonian commute with the spin projection operator

$$\hat{S}_z = \frac{1}{2} \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma}.$$

For the unperturbed part, the commutator can be written down as

$$[\hat{H}_0, \hat{S}_z] = \frac{\xi}{2} \left[\sum_{p\sigma} (p-1) a_{p\sigma}^\dagger a_{p\sigma}, \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} \right]. \quad (4)$$

Recall, however, that the creation and annihilation operators anticommute. Thus, this commutator boils down to whether

$$a_{p\sigma}^\dagger a_{p\sigma} a_{p'\sigma'}^\dagger a_{p'\sigma'} \stackrel{?}{=} a_{p'\sigma'}^\dagger a_{p'\sigma'} a_{p\sigma}^\dagger a_{p\sigma}. \quad (5)$$

Since the operators are paired together, it will take an even number of permutations to make the right-hand side of (5) look like the left-hand side. Therefore, the two sides are equal, and the commutator (4) is zero.

The commutator between the potential and the spin projection operator is

$$[\hat{V}, \hat{S}_z] = -\frac{g}{2} \left[\sum_{pq} a_{p+}^\dagger a_{p-}^\dagger a_{q-} a_{q+}, \sum_{p\sigma} \sigma a_{p\sigma}^\dagger a_{p\sigma} \right].$$

As above, this reduces to commuting strings of paired creation and annihilation operators, and as before, it equals zero.

The total spin operator is

$$\hat{S}^2 \equiv \hat{S}_z^2 + \frac{1}{2} (\hat{S}_+ \hat{S}_- + \hat{S}_- \hat{S}_+) \quad (6)$$

where

$$\hat{S}_\pm \equiv \sum_p a_{p\pm}^\dagger a_{p\mp}.$$

We showed above that \hat{S}_z commutes with both parts of the Hamiltonian, so \hat{S}^2 commutes with the Hamiltonian if \hat{S}_\pm commutes with \hat{H} . Fortunately, this is the same problem as we solved before:

$$[\hat{H}_0, \hat{S}_\pm] = \xi \left[\sum_{p\sigma} (p-1) a_{p\sigma}^\dagger a_{p\sigma}, \sum_p a_{p\pm}^\dagger a_{p\mp} \right] = 0$$

since this is, once again, the problem of permuting pairs of operators, and

$$[\hat{V}, \hat{S}_\pm] = -g \left[\sum_{pq} a_{p+}^\dagger a_{p-}^\dagger a_{q-} a_{q+}, \sum_p a_{p\pm}^\dagger a_{p\mp} \right] = 0$$

for the same reason.

Now we can use the pair creation and annihilation operators

$$\hat{P}_p^+ \equiv a_{p+}^\dagger a_{p-}^\dagger, \quad \hat{P}_p^- \equiv a_{p-} a_{p+}$$

to rewrite the Hamiltonian as

$$\hat{H} = \sum_{pq} (p-1) a_{p\sigma}^\dagger a_{p\sigma} - g \sum_{pq} \hat{P}_p^+ \hat{P}_q^-. \quad (7)$$

This Hamiltonian commutes with the product of the two operators.

$$[\hat{H}, \hat{P}_p^+ \hat{P}_q^-] = \left[\sum_{pq} (p-1) a_{p\sigma}^\dagger a_{p\sigma}, a_{p+}^\dagger a_{p-}^\dagger a_{q-} a_{q+} \right]$$

Once again, this is the permutation of pairs of operators, so the result is zero.

Label	State
$ 1\rangle$	$ 1 + 1-\rangle$
$ 2\rangle$	$ 2 + 2-\rangle$

Table 1: Basis states for two particles and two levels

Part (b)

For the effective Hilbert space of the two lowest single-particle states with two particles, there are only two available configurations (assuming that pairs may not be broken): $|1 + 1-\rangle$ and $|2 + 2-\rangle$. We can compute the matrix elements for these configurations using Wick's theorem.

$$\begin{aligned}\langle q + q - | \hat{H}_0 | s + s - \rangle &= \xi \sum_{p\sigma} (p-1) \langle q + q - | a_{p\sigma}^\dagger a_{p\sigma} | s + s - \rangle \\ \langle q + q - | \hat{H}_0 | s + s - \rangle &= \xi \sum_{p\sigma} (p-1) \langle 0 | a_{q+} a_{q-} a_{p\sigma}^\dagger a_{p\sigma} a_{s-}^\dagger a_{s+}^\dagger | 0 \rangle\end{aligned}$$

There are two possible contractions here:

$$\begin{array}{c} \overbrace{a_{q+} a_{q-} a_{p\sigma}^\dagger a_{p\sigma} a_{s-}^\dagger a_{s+}^\dagger} \\ \overbrace{a_{q+} a_{q-} a_{p\sigma}^\dagger a_{p\sigma} a_{s-}^\dagger a_{s+}^\dagger} \end{array} \left| \begin{array}{l} \delta_{qs} \delta_{sp} \delta_{\sigma+} \\ \delta_{qs} \delta_{qp} \delta_{\sigma-} \end{array} \right.$$

Therefore, each particle contributes its own single-particle energy:

$$\langle q + q - | \hat{H}_0 | s + s - \rangle = 2\xi(s-1)\delta_{qs} \quad (8)$$

The matrix element for the potential was given in the problem statement, and is written in (2).

Now, it's easiest to compute each column of the Hamiltonian matrix separately. For brevity, I'll label the basis states as shown in Table 1. Then,

$$\begin{aligned}\hat{H} |1\rangle &= 2\xi(1-1) |1\rangle - g(|1\rangle + |2\rangle) = -g |1\rangle - g |2\rangle \\ \hat{H} |2\rangle &= 2\xi(2-1) |2\rangle - g(|1\rangle + |2\rangle) = -g |1\rangle + (2\xi - g) |2\rangle\end{aligned}$$

This produces the Hamiltonian matrix

$$\mathbf{H} = \begin{pmatrix} -g & -g \\ -g & 2\xi \end{pmatrix} \quad (9)$$

I used Mathematica to diagonalize this matrix and find the eigenvalues:

$$E_1 = -g + \xi - \sqrt{g^2 + \xi^2}, \quad E_2 = -g + \xi + \sqrt{g^2 + \xi^2} \quad (10)$$

Notice that if $g \rightarrow 0$, $E_1 \rightarrow 0$ and $E_2 \rightarrow 2\xi$ as expected. This result is plotted in Fig. 1.

Label	State
$ 1234\rangle$	$ 1 + 1 - 2 + 2-\rangle$
$ 1256\rangle$	$ 1 + 1 - 3 + 3-\rangle$
$ 1278\rangle$	$ 1 + 1 - 4 + 4-\rangle$
$ 3456\rangle$	$ 2 + 2 - 3 + 3-\rangle$
$ 3478\rangle$	$ 2 + 2 - 4 + 4-\rangle$
$ 5678\rangle$	$ 3 + 3 - 4 + 4-\rangle$

Table 2: Basis states for four particles and four levels

Part (c)

For four particles and four holes (and no broken pairs), there are six available basis states. These are listed and labeled in Table 2.

Acting on the states one-by-one, we find

$$\begin{aligned}
\hat{H} |1234\rangle &= 2\xi |1234\rangle - g(2 |1234\rangle + |1256\rangle + |1278\rangle + |5634\rangle + |7843\rangle) \\
&= 2\xi |1234\rangle - g(2 |1234\rangle + |1256\rangle + |1278\rangle + |3456\rangle + |3478\rangle) \\
\hat{H} |1256\rangle &= 4\xi |1256\rangle - g(2 |1256\rangle + |1234\rangle + |1278\rangle + |3456\rangle + |7856\rangle) \\
&= 4\xi |1256\rangle - g(2 |1256\rangle + |1234\rangle + |1278\rangle + |3456\rangle + |5678\rangle)
\end{aligned}$$

and so on. In other words, the diagonal elements (the expectation values) get a contribution from the single-particle energies, and any matrix element whose bra differs from its ket by at most one pair of single particle states gets a contribution of $-g$. The expectation values also get a contribution of $-2g$ from the two-body operator since they can be formed by destroying and then re-creating either pair. Using this information, the Hamiltonian matrix can be written as:

$$\mathbf{H} = \begin{pmatrix} 2\xi - 2g & -g & -g & -g & -g & 0 \\ -g & 4\xi - 2g & -g & -g & 0 & -g \\ -g & -g & 6\xi - 2g & 0 & -g & -g \\ -g & -g & 0 & 6\xi - 2g & -g & -g \\ -g & 0 & -g & -g & 8\xi - 2g & -g \\ 0 & -g & -g & -g & -g & 10\xi - 2g \end{pmatrix}. \quad (11)$$

There is no analytical solution for the eigenvalues of this matrix, but if we let $\xi = 1$ and $g = 0$, then the eigenvalues are $\{2, 4, 6, 6, 8, 10\}$, the single-particle energies. Fig. 2 shows a plot of the eigenvalues as a function of g . Based on this, it appears that the ground state becomes bound for g greater than about 0.65.

Part (d)

I wrote a small Python module to calculate the eigenvalues. It can be found in the GitHub repository https://github.com/jbradt/structure_work/ in the

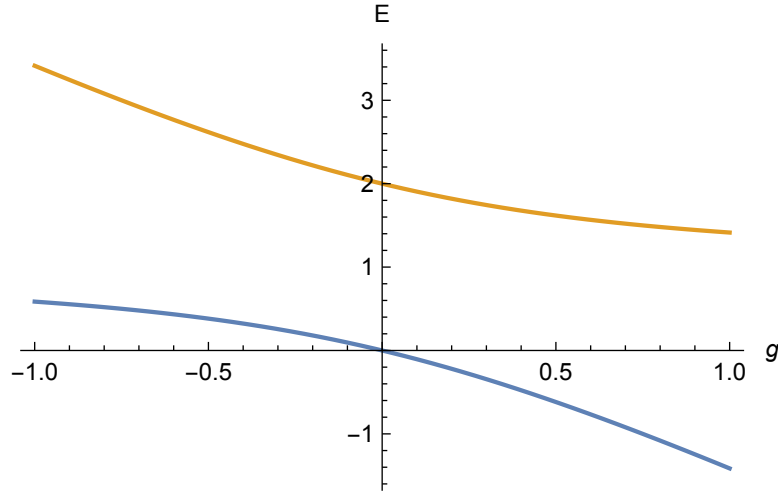


Figure 1: Analytical eigenvalues for two particles and two levels ($\xi = 1$)

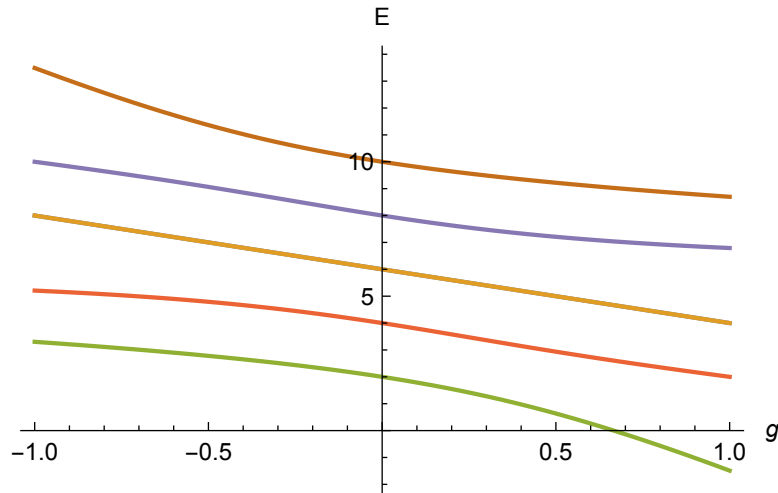


Figure 2: Analytical eigenvalues for four particles and four levels ($\xi = 1$)

directory `project/shellcode`. This program is perhaps not the fastest thing in the world, but it was written with simplicity in mind, and not efficiency.

The program generates the single-particle states for the pairing model internally and then finds the possible Slater determinants. These are then used to find the Hamiltonian matrix, which can then be diagonalized. The interactions for the Hamiltonian are written directly into the code, which is not ideal, but is perhaps less inconvenient in Python than it would be in a compiled language. The usage of the code is documented in the docstrings of the source file itself.

After running the code for the cases in the previous two parts, I created two graphs of the results for varying values of g . Fig. 3 shows the results for two particles and two levels, and Fig. 4 shows the results for four particles and four levels. These results look just like the analytical results from the previous parts.

Part (e)

One way to test the code is to remove the pairing interaction. This is equivalent to letting $g \rightarrow 0$. If I generate the Hamiltonian matrix for four particles and four levels, leaving the pairs-only restriction in place to make comparisons easy, I get the following matrix:

$$\mathbf{H} = \begin{pmatrix} 2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 4 & 0 & 0 & 0 & 0 \\ 0 & 0 & 6 & 0 & 0 & 0 \\ 0 & 0 & 0 & 6 & 0 & 0 \\ 0 & 0 & 0 & 0 & 8 & 0 \\ 0 & 0 & 0 & 0 & 0 & 10 \end{pmatrix}.$$

This matrix is diagonal, so the eigenvalues are $\{2, 4, 6, 6, 8, 10\}$, which are just the single-particle energies, as expected.

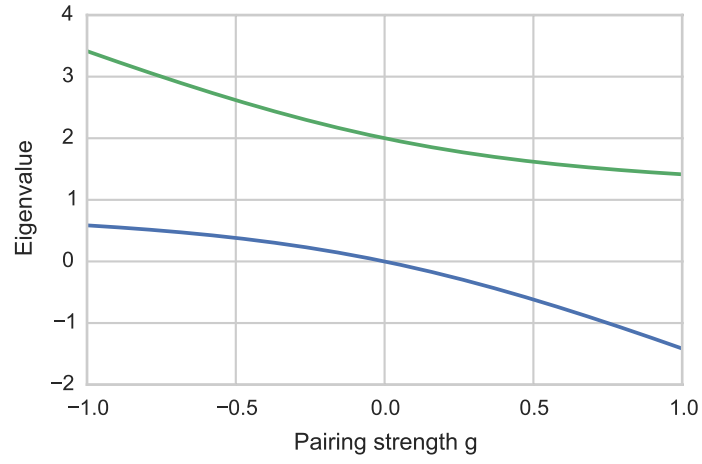


Figure 3: Calculated eigenvalues for two particles and two levels ($\xi = 1$)

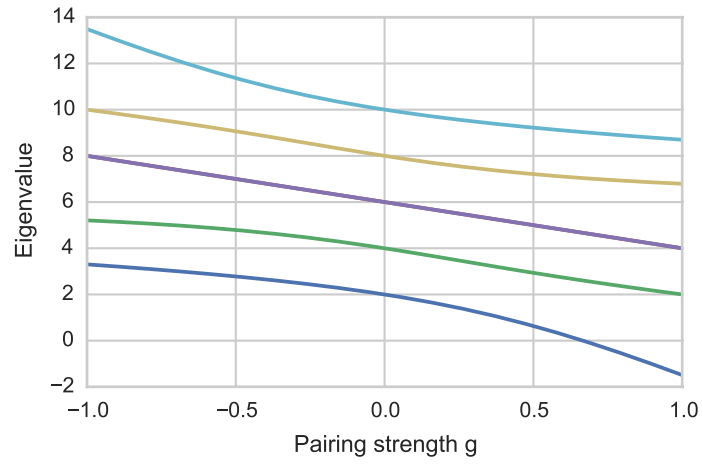


Figure 4: Calculated eigenvalues for four particles and four levels ($\xi = 1$)