# Overlapping Domain Decomposition with Non-matching Grids

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#### 1 Introduction

In this paper we consider two topics. In Section 2 we introduce a new macro-hybrid formulation based on overlapping domain decomposition for linear elliptic equations with symmetric positive definite operators. The problem is discretized by the mortar element method using non-matching grids on the interfaces between subdomains. An iterative method of an optimal order of arithmetical complexity is proposed for solving the arising algebraic systems in the case of regular quasiuniform hierarchical grids. An example of such a formulation was originally given in [Kuz95b]. The approach proposed here has many common points with the decentralization methods studied more than twenty years ago in [BLT74, Lem74]. In these papers the authors used splittings of bilinear forms between different subdomains to decompose a variational problem.

The second important topic is presented in Section 4 where we consider an extension of results from [Kuz95a, Kuz95b] to the case of overlapping subdomains. Here we present several results which mainly concern the construction of the interface preconditioner.

In Section 5 results of numerical experiments with 2D- and 3D-overlapping subdomains are given.

## 2 Macro-hybrid based on Overlapping Domain Decomposition

Let us consider a model elliptic problem

$$-\Delta u + cu = f \quad \text{in } \Omega$$

$$\frac{\partial u}{\partial \mathbf{r}} = 0 \quad \text{on } \partial \Omega$$
(2.1)

where  $f \in L_2(\Omega)$  is a given function,  $c \equiv \text{const} \in (0; 1]$ ,  $\partial \Omega$  is the boundary of a domain  $\Omega$  and  $\mathbf{n}$  is the outer unit normal vector to  $\partial \Omega$ . For the sake of simplicity we assume that  $\Omega$  is a polygon in  $\mathbf{R}^p$ , p = 2, 3 with diam  $(\Omega) \sim O(1)$ , and all further subdomains of  $\Omega$  are also polygons with diameters O(1).

The classical weak formulation of (2.1) is: find  $u \in H^1(\Omega)$  such that

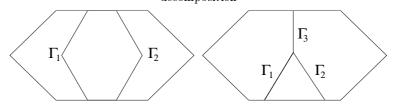
$$\Phi(u) = \min_{v \in H^1(\Omega)} \Phi(v), \tag{2.2}$$

where

$$\Phi(v) = \int_{\Omega} \left[ |\nabla v|^2 + cv^2 - 2fv \right] d\Omega. \tag{2.3}$$

Let  $\Omega_1$  and  $\Omega_2$  be two overlapping subdomains of  $\Omega$  ( $\Omega_1 \cap \Omega_2 \neq \emptyset$ ) such that  $\overline{\Omega_1 \cup \Omega_2} = \overline{\Omega}$ . We assume that subdomains  $\Omega_1$  and  $\Omega_2$  are regularly shaped. Examples of such a partitioning of  $\Omega$  into two subdomains are given in Fig. 1 ( $\Omega_1$  is located on the left, i.e.  $\Gamma_2 \subset \partial \Omega_1 \cap \Omega$ ).

Figure 1 a) Overlapping and b) overlapping / nonoverlapping domain decomposition



We denote the intersection of  $\Omega_1$  and  $\Omega_2$  by  $\Omega_{12}$  and define two bilinear forms

$$a_k(u, v) = \int_{\Omega_k} \left[ a_k \nabla v \cdot \nabla u + c_k u v \right] d\Omega, \qquad k = 1, 2, \tag{2.4}$$

two linear forms

$$l_k(v) = \int_{\Omega_k} f_k v d\Omega, \qquad k = 1, 2, \tag{2.5}$$

and two quadratic functionals

$$\psi_k(v) = a_k(v, v) - 2l_k(v), \qquad k = 1, 2.$$
 (2.6)

The coefficients  $a_k$ ,  $c_k$  and functions  $f_k$  are defined by

$$a_{k} = \begin{cases} 1 & \text{in } \Omega_{k} \setminus \Omega_{12} \\ q_{k} & \text{in } \Omega_{12} \end{cases} \qquad c_{k} = \begin{cases} c & \text{in } \Omega_{k} \setminus \Omega_{12} \\ q_{k}c & \text{in } \Omega_{12} \end{cases}$$

$$f_{k} = \begin{cases} f & \text{in } \Omega_{k} \setminus \Omega_{12} \\ q_{k}f & \text{in } \Omega_{12} \end{cases}$$

$$(2.7)$$

where  $q_k$  are positive constants, k = 1, 2 such that  $q_1 + q_2 = 1$ . It is important that

$$\psi_k(v) = q_k \Phi(v), \quad \forall v \in H^1(\Omega), \text{ supp } v \in \Omega_{12}, \ k = 1, 2.$$
 (2.8)

To introduce and to analyze macro-hybrid formulations of elliptic problems we have to deal with interfaces between subdomains. To this end we introduce the following notation:

$$\gamma_g = (\partial \Omega_1 \cap \Omega) \cup (\partial \Omega_2 \cap \Omega), 
\gamma_{in} = \partial \Omega_1 \cap \partial \Omega_2, 
\partial \gamma_{in} = \bar{\gamma}_{in} \setminus \mathring{\gamma}_{in}.$$
(2.9)

Here  $\mathring{\gamma}_{in}$  is the interior part of  $\gamma_{in}$  with respect to (p-1)-dimensional topology,  $\gamma_g$  is said to be the "global" interface (with respect to the macro-hybrid formulation to be presented), and  $\partial \gamma_{in}$  is the set of cross points in the case p=2 and the set of interedges in the case p=3.

The set  $\gamma_g \setminus \partial \gamma_{in}$  can be presented as the union of nonoverlapping open subsets  $\Gamma_s$ ,  $s = 1, \ldots, s_g$  such that each  $\Gamma_s$  is a piece-wise linear curve in the case of p = 2 and a piece-wise linear surface in the case of p = 3. It is obvious that this partitioning  $\gamma_g$  into nonoverlapping open subsets is unique.

For examples given in Fig. 1 we have  $s_g = 2$  in the case a) and  $s_g = 3$  in the case b). Respectively,  $\partial \gamma_{in}$  consists of four points both in the cases a) and b).

Now we introduce the space  $V = H^1(\Omega_1) \times H^1(\Omega_2)$ , the space

$$W = \left\{ \bar{v} = (v_1, v_2) : \ \bar{v} \in V, \int_{\Gamma_s} (v_1 - v_2) \mu \, ds = 0, \, \forall \mu \in H^{-1/2}(\Gamma_s), \, s = 1, \dots, s_g \right\}$$
(2.10)

and the quadratic functional

$$\psi(\bar{v}) = \psi_1(v_1) + \psi_2(v_2), \quad \bar{v} \in V. \tag{2.11}$$

It can be shown (see, for instance [BF91]) that under the assumptions made the following macro-hybrid formulation of problem (2.1):

$$\bar{u} \in V : \ \psi(\bar{u}) = \min_{\bar{v} \in W} \psi(\bar{v})$$
 (2.12)

has a unique solution and is equivalent to problem (2.2). We understand the equivalence in the sense that

$$u(x) = u_k(x) \quad \forall x \in \Omega_k, \tag{2.13}$$

where u is the solution function to (2.2).

Problem (2.12) has also an equivalent formulation in terms of Lagrange multipliers. For instance, in the case of example in Fig. 1b it can be presented in the following

form: find  $(\bar{u}, \bar{\lambda}) \in V \times \Lambda$  such that

$$a_{1}(u_{1}, v_{1}) + \int_{\Gamma_{1}} \lambda_{1}v_{1} ds + \int_{\Gamma_{2}} \lambda_{2}v_{1} ds + \int_{\Gamma_{3}} \lambda_{3}v_{1} ds = l_{1}(v_{1}),$$

$$a_{2}(u_{2}, v_{2}) - \int_{\Gamma_{1}} \lambda_{1}v_{2} ds - \int_{\Gamma_{2}} \lambda_{2}v_{2} ds - \int_{\Gamma_{3}} \lambda_{3}v_{2} ds = l_{2}(v_{2}),$$

$$\int_{\Gamma_{1}} (u_{1} - u_{2})\mu_{1} ds = 0,$$

$$\int_{\Gamma_{2}} (u_{1} - u_{2})\mu_{2} ds = 0,$$

$$\int_{\Gamma_{2}} (u_{1} - u_{2})\mu_{3} ds = 0,$$

$$(2.14)$$

 $\forall (\bar{v}, \bar{\mu}) \in V \times \Lambda$ . Here  $\Lambda = \prod_{s=1}^{3} H^{-1/2}(\Gamma_s)$ . It can be easily shown that

$$\lambda_1 = q_2 \frac{\partial u_2}{\partial \mathbf{n}_2} \text{ on } \Gamma_1, \quad \lambda_2 = -q_1 \frac{\partial u_1}{\partial \mathbf{n}_1} \text{ on } \Gamma_2, \quad \lambda_3 = -\frac{\partial u_1}{\partial \mathbf{n}_1} \text{ on } \Gamma_3,$$

$$(2.15)$$

where  $\mathbf{n_1}$  and  $\mathbf{n_2}$  are the outer normal vectors to  $\partial\Omega_1$  and  $\partial\Omega_2$ , respectively. Recall that  $u_1 \equiv u$  in  $\Omega_1$  and  $u_2 \equiv u$  in  $\Omega_2$ .

In a compact form (2.14) can be presented [GW88, BF91] by: find  $(\bar{u}, \bar{\lambda}) \in V \times \Lambda$  such that

$$\begin{array}{lcl} \hat{a}(\bar{u},\,\bar{v}) + b(\bar{\lambda},\,\bar{v}) & = & \hat{l}(\bar{v}), \\ b(\bar{\mu},\,\bar{u}) & = & 0, & \forall (\bar{v},\,\bar{\mu}) \in V \times \Lambda. \end{array} \tag{2.16}$$

Here

$$V = \prod_{k=1}^{m} V_k, \qquad V_k = H^1(\Omega_k), \ k = 1, \dots, m,$$

$$\Lambda = \prod_{s=1}^{s_g} \Lambda_s, \qquad \Lambda_s = H^{-1/2}(\Gamma_s), \ s = 1, \dots, s_g,$$

$$\hat{a}(\bar{u}, \bar{v}) = \sum_{k=1}^{m} a_k(u, v), \quad \hat{l}(\bar{v}) = \sum_{k=1}^{m} l_k(v),$$

$$b(\bar{\mu}, \bar{u}) = \sum_{s=1}^{3} \int (u_1 - u_2) \mu_s \, ds,$$

$$(2.17)$$

where m=2 and  $s_g=3$ .

**Remark 1** Generalization to a larger number of subdomains is straightforward. For instance, if we use the formulation (2.16)-(2.17) we have to assume that for any simply connected subdomain  $G \subset \Omega \setminus \gamma_g$  a positive constant  $q_G$  exists such that

$$a_k(u, v) = q_G a(u, v), \quad l_k(v) = q_G l(v), \quad \forall u, v \in H^1(\Omega), \quad \text{supp } v \subset G, \quad k = \overline{1, m}.$$

$$(2.18)$$

**Remark 2** If  $\int_{\Omega} f d\Omega = 0$  and  $c \ll 1$  then problem (2.1) can be considered as a singular perturbation of the Neumann problem

$$\begin{aligned}
-\Delta u &= f & \text{in } \Omega \\
\frac{\partial u}{\partial \mathbf{n}} &= 0 & \text{on } \partial \Omega.
\end{aligned} (2.19)$$

## 3 The Mortar Element Method and Algebraic Systems

We consider the only case when  $\Omega_{kh}$  are conforming triangular, p=2, or tetrahedral, p=3, partitions of  $\Omega_k$ ,  $k=1,\ldots,m$ , and  $\gamma_g$  does not intersect the interiors of the grid cells. Then  $V_{kh}$  are the standard piecewise linear finite element subspaces of  $V_k \equiv H^1(\Omega_k)$ ,  $k=1,\ldots,m$ . The finite element subspaces  $\Lambda_{sh} \subset \Lambda \equiv H^{-\frac{1}{2}}(\Gamma_s)$ ,  $s=1,\ldots,s_g$  are chosen using the mortar element technique from [BMP89, BM94, Kuz95b].

The mortar finite element discretization of (2.16)–(2.17) is defined by: find  $(\bar{u}_h, \bar{\lambda}_h) \subset V_h \times \Lambda_h$  such that

$$\begin{array}{rcl} \hat{a}(\bar{u}_h,\bar{v}) + b(\bar{\lambda}_h,\bar{u}_h) & = & \hat{l}(\bar{v}), \\ b(\bar{\mu},\bar{u}_h) & = & 0, \end{array}$$

 $\forall (\bar{v}, \bar{\mu}) \in V_h \times \Lambda_h$  where  $V_h = \prod_{k=1}^m V_{kh}$  and  $\Lambda_h = \prod_{s=1}^{s_g} \Lambda_{sh}$ . Problem (3.20) leads to an algebraic system

$$\mathcal{A}x = y \tag{3.20}$$

with a saddle-point matrix

$$\mathcal{A} = \begin{pmatrix} A & B^T \\ B & O \end{pmatrix} \tag{3.21}$$

and vectors

$$x = \begin{pmatrix} u \\ \lambda \end{pmatrix}, \qquad y = \begin{pmatrix} f \\ 0 \end{pmatrix}.$$
 (3.22)

Here A is a symmetric positive definite matrix and Ker  $B^T = 0$ . It follows immediately that det  $A \neq 0$ .

For further analysis we need a more detailed description of A and B in block forms. The simplest block representations of A and  $B^T$  are:

$$A = \begin{pmatrix} A_1 & O & O \\ O & \ddots & O \\ O & O & A_m \end{pmatrix}, \qquad B^T = \begin{pmatrix} B_1^T \\ \vdots \\ B_m^T \end{pmatrix}$$
(3.23)

Here the kth block corresponds to the degrees of freedom of the finite element space  $V_{kh}$ ,  $k = 1, \ldots, m$ .

For each subdomain  $\Omega_k$  we partition degrees of freedom (grid nodes) into two groups. In the second group denoted by  $\gamma$  we collect the degrees of freedom which correspond to the grid nodes belonging to  $\gamma_k = \gamma_g \cap \bar{\Omega}_k$ . All other degrees of freedom we collect in the first group denoted by I. These partitionings induce the following block representations:

$$A_k = \begin{pmatrix} A_{kI} & A_{kI\gamma} \\ A_{k\gamma I} & A_{k\gamma} \end{pmatrix}, \qquad B_k^T = \begin{pmatrix} O \\ B_{k\gamma}^T \end{pmatrix}. \tag{3.24}$$

Let  $\mathcal{B}$  be a symmetric positive definite matrix and  $\mathcal{H} = \mathcal{B}^{-1}$ . Since  $\mathcal{A} = \mathcal{A}^T$  the preconditioned Lanczos [MK74, Kuz95b] can be used to solve system (3.20). In this paper we also recommend the preconditioned conjugate method based on the  $\mathcal{B}$ -norm of minimal errors [MK74]:

$$\hat{p}_{l} = \begin{cases}
\mathcal{H}\xi^{0}, & l = 1, \\
\mathcal{H}\xi^{l-1} - \alpha_{l}\hat{p}_{l-1}, & l > 1,
\end{cases}$$

$$p_{l} = \mathcal{H}A\hat{p}_{l}, \qquad (3.25)$$

$$x^{l} = x^{l-1} - \beta_{l}p_{l},$$

$$\alpha_{l} = \frac{(\xi^{l-1}, \, \mathcal{A}\hat{p}_{l-1})_{\mathcal{H}}}{(\mathcal{A}\hat{p}_{l-1}, \, \mathcal{A}\hat{p}_{l-1})_{\mathcal{H}}}, \quad \beta_{l} = \frac{(\xi^{l-1}, \, \hat{p}_{l})}{(\mathcal{A}\hat{p}_{l}, \, \mathcal{A}\hat{p}_{l})_{\mathcal{H}}},$$

where  $\xi^l = \mathcal{A}x^l - y$  are the residual vectors,  $l = 1, 2, \ldots$  Assume that the eigenvalues of  $\mathcal{H}\mathcal{A}$  belong to the union of segments  $[d_1; d_2]$  and  $[d_3; d_4]$  with  $d_1 \leq d_2 < 0 < d_3 \leq d_4$ . Then the convergence estimate

$$||x^{l} - x||_{\mathcal{H}} \le 2q^{l}||x^{0} - x||_{\mathcal{H}}, \quad l \ge 1,$$
 (3.26)

holds [MK74] where  $q = \frac{\hat{d} - \check{d}}{\hat{d} + \check{d}}$ ,  $\hat{d} = \max\{d_4; |d_1|\}$ , and  $\check{d} = \min\{d_3; |d_2|\}$ .

## 4 Block Diagonal Preconditioning

We propose a preconditioner  $\mathcal{H}$  as a block diagonal matrix:

$$\mathcal{H} = \begin{pmatrix} H_A & O \\ O & H_\lambda \end{pmatrix} \tag{4.27}$$

where  $H_A$  is also a block diagonal matrix:

$$H_A = \begin{pmatrix} H_1 & O & O \\ O & \ddots & O \\ O & O & H_m \end{pmatrix}. \tag{4.28}$$

All blocks are symmetric positive definite matrices.  $H_k$  are said to be the subdomain preconditioners, and  $H_{\lambda}$  is said to be the interface preconditioner.

If matrices  $H_k$  are spectrally equivalent to the matrices  $A_k^{-1}$  with constants independent of the value of the coefficient c, and if a matrix  $H_{\lambda}$  is spectrally equivalent to the matrix  $S_{\lambda}^{-1}$  with  $S_{\lambda}$  given by

$$S_{\lambda} = BA^{-1}B^{T} \equiv \sum_{k=1}^{m} B_{k\gamma} S_{k\gamma}^{-1} B_{k\gamma}^{T}$$
(4.29)

with the constants independent of the value of c then the values of  $\hat{d}, \check{d}$  in (3.26) are positive constants [Kuz95a] also independent of c. Here

$$S_{k\gamma} = A_{k\gamma} - A_{k\gamma I} A_{kI}^{-1} A_{kI\gamma} \tag{4.30}$$

are the Schur complements. Our aim is to construct a preconditioner  $\mathcal{H}$  spectrally equivalent [Kuz95a] to the matrix  $\mathcal{A}^{-1}$  with constants independent of c.

#### Subdomain Preconditioners

Let us define matrices  $\overset{\circ}{A}_k$  and  $M_k$  by:

$$(\overset{\circ}{A}_{k} v, w) = \int_{\Omega_{k}} \nabla v_{h} \cdot \nabla w_{h} d\Omega, \qquad (4.31)$$

$$(M_{k} v, w) = \int_{\Omega_{k}} v_{h} w_{h} d\Omega$$

 $\forall v_h, w_h \in V_{kh}, k = 1, \ldots, m$ . Thus, matrices  $\overset{\circ}{A}_k$  are the stiffness matrices for the operator  $-\Delta$  with the Neumann boundary conditions, and  $M_k$  are the corresponding mass matrices. It can be easily shown [Kuz95b] that

$$A_k^{-1} \sim \left( \stackrel{\circ}{A}_k + M_k \right)^{-1} + \frac{1}{c} P_k$$
 (4.32)

where  $P_k M_k$  is the  $M_k$ -orthogonal projector onto Ker  $\stackrel{\circ}{A}_k$  and the sign " $\sim$ " denotes the spectral equivalence. Moreover, the constants of the spectral equivalence in (4.32) are independent of the value of c.

Suppose that a matrix  $\overset{\circ}{H}_k$  is spectrally equivalent to the matrix  $(\overset{\circ}{A}_k + M_k)^{-1}$ . Then the matrix

$$H_k = \overset{\circ}{H}_k + \frac{1}{c} P_k \tag{4.33}$$

is spectrally equivalent to matrix  $A_k^{-1}$  with constants independent of the value of c.

We have plenty of choices for  $\overset{\circ}{H}_k$ ,  $k = 1, \ldots, m$ .

#### Interface Preconditioner

We can easily show [Kuz95b] that

$$S_{k\gamma}^{-1} \sim \tilde{S}_{k\gamma}^{-1} + \frac{1}{c} P_{k\gamma} \tag{4.34}$$

where  $\tilde{S}_{k\gamma}^{-1}$  is the Schur complement for the matrix  $\overset{\circ}{A}_k + M_k$  and  $P_{k\gamma} M_{k\gamma}$  is the  $M_{k\gamma}$  orthogonal projector onto Ker  $S_{k\gamma}$  in the case c=0. Moreover, the constants of equivalence in (4.34) are independent of the value of c. Here  $M_{k\gamma}$  is the interface mass matrix defined by:

$$(M_{k\gamma}v, w) = \int_{\gamma_h} v_h w_h ds \quad \forall v_h, w_h \in V_{k\gamma h}$$

$$(4.35)$$

where  $V_{k\gamma h}$  is the trace of  $V_{kh}$  into  $\gamma_k = \partial \Omega_k \cap \Omega$ ,  $k = 1, \ldots, m$ . Let the matrices

$$\overset{\circ}{H}_{k} = \begin{pmatrix} \overset{\circ}{H}_{kI} & \overset{\circ}{H}_{kI\gamma} \\ \overset{\circ}{H}_{k\gamma I} & \overset{\circ}{H}_{k\gamma} \end{pmatrix}$$
(4.36)

be spectrally equivalent to the matrices  $(\mathring{A}_k + M_k)^{-1}$ ,  $k = 1, \ldots, m$ . We can also prove that the matrix

$$\hat{S}_{\lambda} = \sum_{k=1}^{m} B_{k\gamma} (\mathring{H}_{k\gamma} + \frac{1}{c} P_{k\gamma}) B_{k\gamma}^{T}$$

$$\tag{4.37}$$

is spectrally equivalent to  $S_{\lambda}$  with constants independent of the value of c.

To construct the interface preconditioner  $H_{\lambda}$  we shall use the preconditioned Chebyshev iterative procedure [BPS86, Kuz95a]. Let  $\hat{H}_{\lambda}$  be a symmetric positive defined matrix and  $\nu_{\lambda} = \lambda_{\max}/\lambda_{\min}$  where  $\lambda_{\max}$  and  $\lambda_{\min}$  are the maximal and minimal eigenvalues of  $\hat{H}_{\lambda}\hat{S}_{\lambda}$ , respectively. Then for any  $t_{\lambda} \sim \sqrt{\nu_{\lambda}}$  the matrix

$$H_{\lambda} = \left[ I_{\lambda} - \prod_{t=1}^{t_{\lambda}} \left( I_{\lambda} - \alpha_{t} \hat{H}_{\lambda} \hat{S}_{\lambda} \right) \right] \hat{S}_{\lambda}^{-1}$$
 (4.38)

is spectrally equivalent to the matrix  $S_{\lambda}^{-1}$  where  $\{\alpha_t\}$  is a set of the corresponding Chebyshev parameters.

Let  $\hat{B}_{\lambda}$  be a symmetric positive definite matrix such that  $1 \in [\mu_{\min}; \mu_{\max}]$  where  $\mu_{\min}$  and  $\mu_{\max}$  are the minimal and maximal eigenvalues of the matrix  $\hat{B}_{\lambda}^{-1} \sum_{k=1}^{m} B_{k\gamma} H_{k\gamma}^{\circ} B_{k\gamma}^{T}$ , respectively. Then for the choice  $\hat{H}_{\lambda} = \hat{R}_{\lambda}^{-1}$  where

$$\hat{R}_{\lambda} = \hat{B}_{\lambda} + \frac{1}{c} \sum_{k=1}^{m} B_{k\gamma} P_{k\gamma} B_{k\gamma}^{T}, \tag{4.39}$$

the estimate

$$\nu_{\lambda} \le \hat{\nu}_{\lambda} \equiv \mu_{\text{max}} / \mu_{\text{min}} \tag{4.40}$$

holds.

A solution algorithm for a system

$$\hat{R}_{\lambda}z = g$$

is presented in [Kuz95b, KW95]. It includes a so called "coarse grid" problem based on the projectors  $P_{k\gamma}$ ,  $k=1,\ldots,m$ .

Arithmetical Complexity for Hierarchical Grids

Assume that grids  $\Omega_{kh}$  are regular, quasiuniform and hierarchical with the average grid step size  $h \sim \sqrt[p]{N}$  where N is the dimension of matrix  $\mathcal{A}$ .

In this case we can use various V-cycle multilevel preconditioners to define matrix  $\overset{\circ}{H_k}$  in (4.33). These preconditioners are spectrally equivalent to the matrices  $(\overset{\circ}{A_k} + M_k)^{-1}$ ,  $k = 1, \ldots, m$  and have the optimal order of arithmetical complexity [Osw94, Xu92], i. e. the multiplication with such a preconditioner by a vector costs O(N) arithmetical operations.

Our choice  $H_{k\gamma}$  in (4.37) as the corresponding blocks of V-cycle multilevel preconditioner (BPX or MDS-type) is based on two observations. The first one is obvious: spectral equivalence of  $H_{k\gamma}$  and  $\tilde{S}_{\gamma k}^{-1}$  follows directly from the spectral equivalence of  $H_k$  and  $(\hat{A}_k + M_k)^{-1}$ ,  $k = 1, \ldots, m$ . The second observation is rather technical and concerns implementation algorithms for V-cycle multilevel preconditioners: multiplication of  $H_{k\gamma}$  by a vector can be implemented with  $O(h^{1-p})$  arithmetical operations. The latter observation has at least one very important consequence: the corresponding matrix  $\hat{S}_{\lambda}$  can be multiplied by a vector with  $O(h^{1-p})$  arithmetical operations, i.e. multiplication with  $\hat{S}_{\lambda}$  has the optimal order of arithmetical complexity.

It remains to choose preconditioner  $\hat{R}_{\lambda}$ , and we do not need an optimal preconditioner because the dimension of  $S_{\lambda}$  is much smaller than the dimension of A.

In paper [Kuz95a] we proposed to choose  $\hat{B}_{\lambda}$  being equal to a scalar matrix which is a spectrally equivalent to the matrix  $\sum_{k=1}^{m} B_{k\gamma} M_{k\gamma}^{-1} B_{k\gamma}^{T}$ . With this choice, obviously

$$\nu_{\lambda} < \text{const} \cdot h^{-2}$$

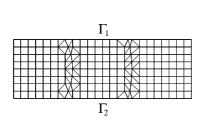
where the constant is independent of h and c, and the multiplication  $B_{\lambda}^{-1}$  by a vector can be implemented with  $O(h^{1-p})$  arithmetical operations.

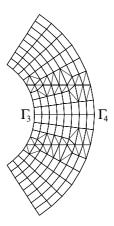
On the basis of the latter facts we conclude that  $t_{\nu}$  should be proportional to  $h^{-1}$ , and arithmetical complexity of the corresponding preconditioner  $H_{\lambda}$  in (4.38) is of the order  $O(h^{1-p})$ . In some particular cases we can prove [BPS86, Kuz95a] that  $t_{\nu} \sim h^{-1/2}$  and consequently the arithmetical complexity of  $H_{\lambda}$  is of the order  $O(h^{1/2-p})$ .

# 5 Numerical Experiment

The numerical experiments have been performed for two test cases.

Figure 2 Cartesian and polar locally fitted grids





The first test case is presented by the union of a rectangle and a segment of a ring. In the rectangular subdomain we have a rectangular cartesian grid and in the segment we have an orthogonal polar grid. Both grids are fitted to the interface boundary which consists of two straight segments  $\Gamma_1$  and  $\Gamma_2$ , and two circle's segments  $\Gamma_3$  and  $\Gamma_4$ . These grids are given in Fig. 2. Here  $\partial \Omega_1 \cap \Omega = \Gamma_1 \cup \Gamma_2$  and  $\partial \Omega_2 \cap \Omega = \Gamma_3 \cup \Gamma_4$ .

Table 1 Cartesian / Polar grids

Cartesian grids in $\Omega_1$	Polar grids in $\Omega_2$	Number of Chebyshev iterations	Number of Lanczos iterations
24  imes 8	$16 \times 8$	13	68
$48 \times 16$	32  imes 16	22	72
96  imes 32	64  imes 32	32	75
$192\times 64$	$128\times 64$	45	77

The second test case is presented by the union of two parallelepipeds. The intersection of them is also a parallelepiped. Both grids are uniform and rectangular ones.

The results of numerical experiments are given in Tables 1 and 2. Two first columns contain information about the grids: product of numbers of nodes for each of the coordinates. The third column contains information about number of iterations  $t_{\lambda}$  used in construction of the interface preconditioner (4.38). The last column contains the number of iterations which were needed to reduce the  $\mathcal{H}$ -norm of the initial residual  $\xi^0$  of  $10^6$  times by the preconditioned Lanczos method.

Grids in $\Omega_1$	Grids in $\Omega_2$	Number of Chebyshev iterations	Number of Lanczos iterations
$16 \times 16 \times 16$	$16 \times 8 \times 8$	24	54
32  imes 32  imes 16	32  imes 16  imes 16	38	51
64  imes 64  imes 32	64  imes 16  imes 16	52	51

 Table 2
 Intersecting parallelepipeds: uniform grids

One can see that  $t_{\lambda}$  grows up proportionally to  $h^{-1/2}$  and that the number of Lanczos iterations is almost constant for both test cases.

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