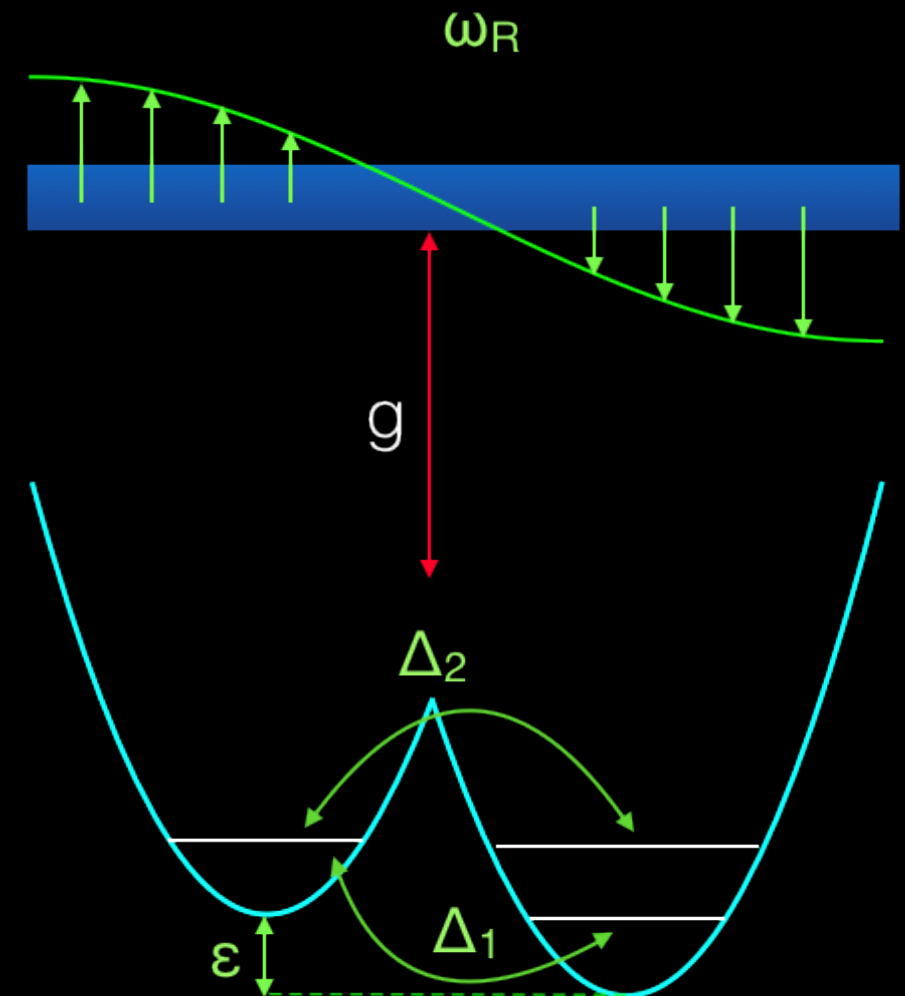


Long-range two-qubit gates: hybrid qubit & quadrupolar

Phys. Rev. A 104, 032612 (2021)
Phys. Rev. Research 3 (1), 013171 (2021)

Outline

- Background and basics
- Hybrid qubit
 - HQ basics
 - HQ sweet spots
 - Resonant gate
 - Cavity frequency problem
 - SB gate
- Charge-quadrupole qubit
 - CQ basics
 - Experiment
 - Noise characterization
- Conclusions



Background

Single-qubit gates with 99% fidelities are already achievable

See: Veldhorst et al. Nat. Nano **9**, 981-985 (2014)

Kawakami et al. PNAS October **18**, 2016 113 (42) 11738-11743

Takeda et al. Science Advances **2**, 8, e1600694

Yoneda et al. Nat. Nano **13**, 102-106 (2018)

Two-qubit gates based on exchange have been implemented. >90% achieved

See: Veldhorst et al. Nature **526**, 410-414 (2015)

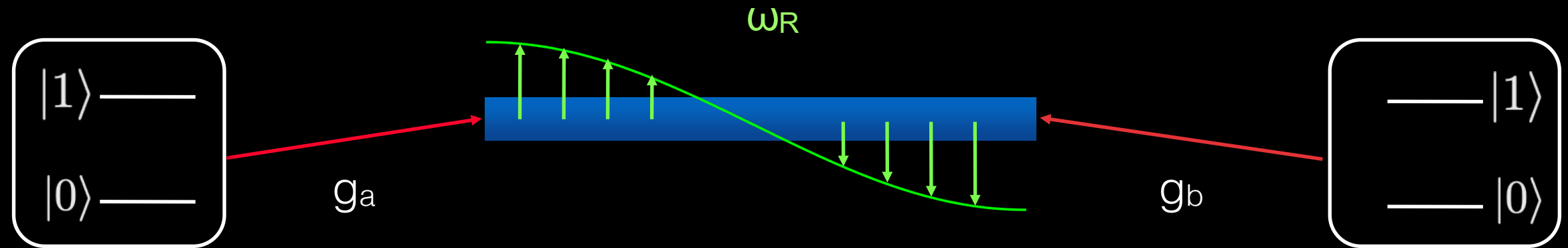
Zajac et al. Science **359**, 6374, 439-442 (2018)

Watson et al. Nature **555**, 633-637 (2018)

Huang et al. Nature **569**, 532-536 (2019)

Xue et al. PRX **9**, 021011 (2019)

On the way towards long-range two-qubit gates



Spin-cQED architectures allow long-range coupling between qubits

Not limited by nearest neighbor interactions

Potential for hybrid quantum systems

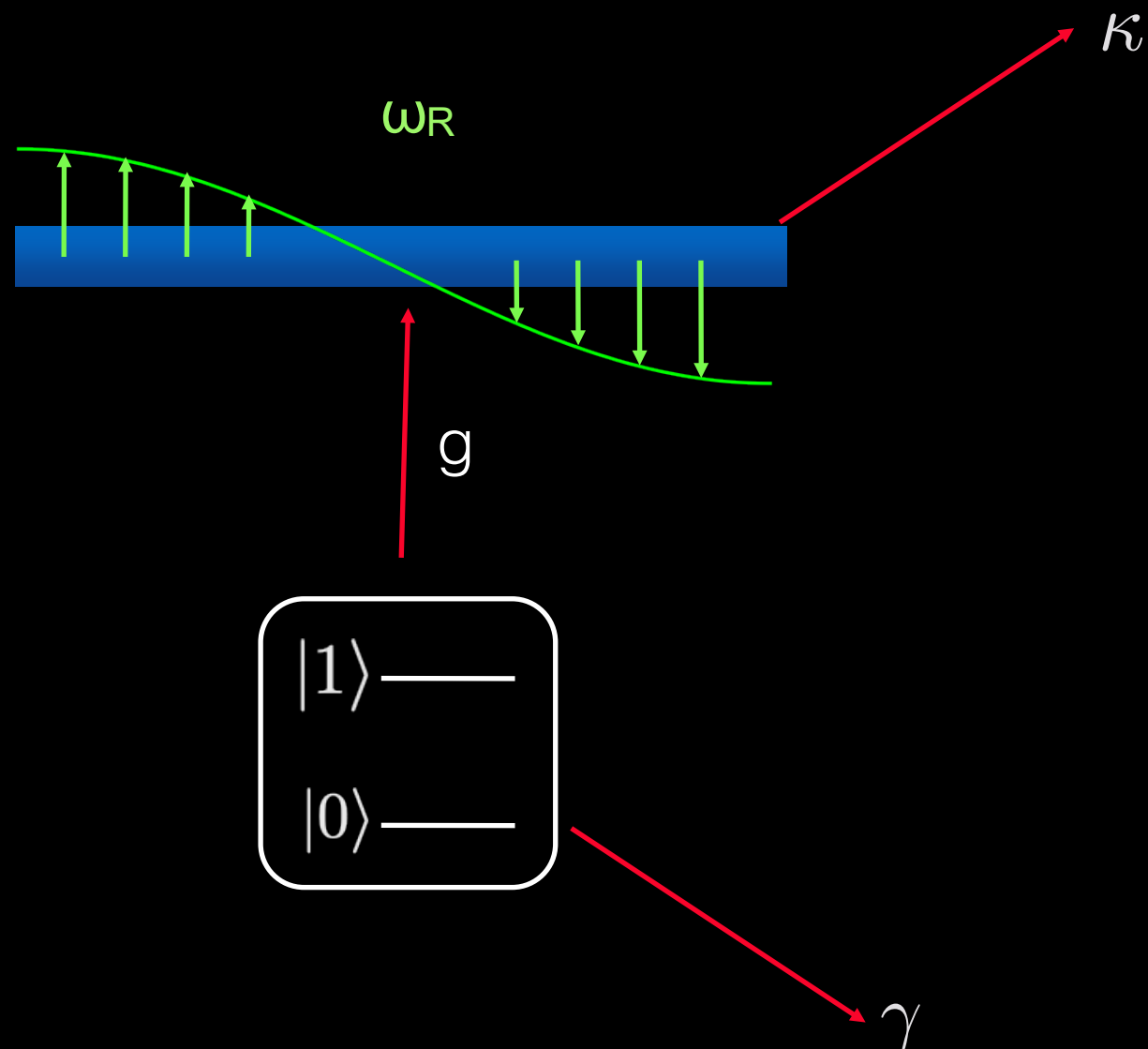
Coupling spin qubits with superconducting resonators

Strong coupling regime achieved with semiconductor qubits

See: Landig et al. Nature **560**, 179-184 (2018)

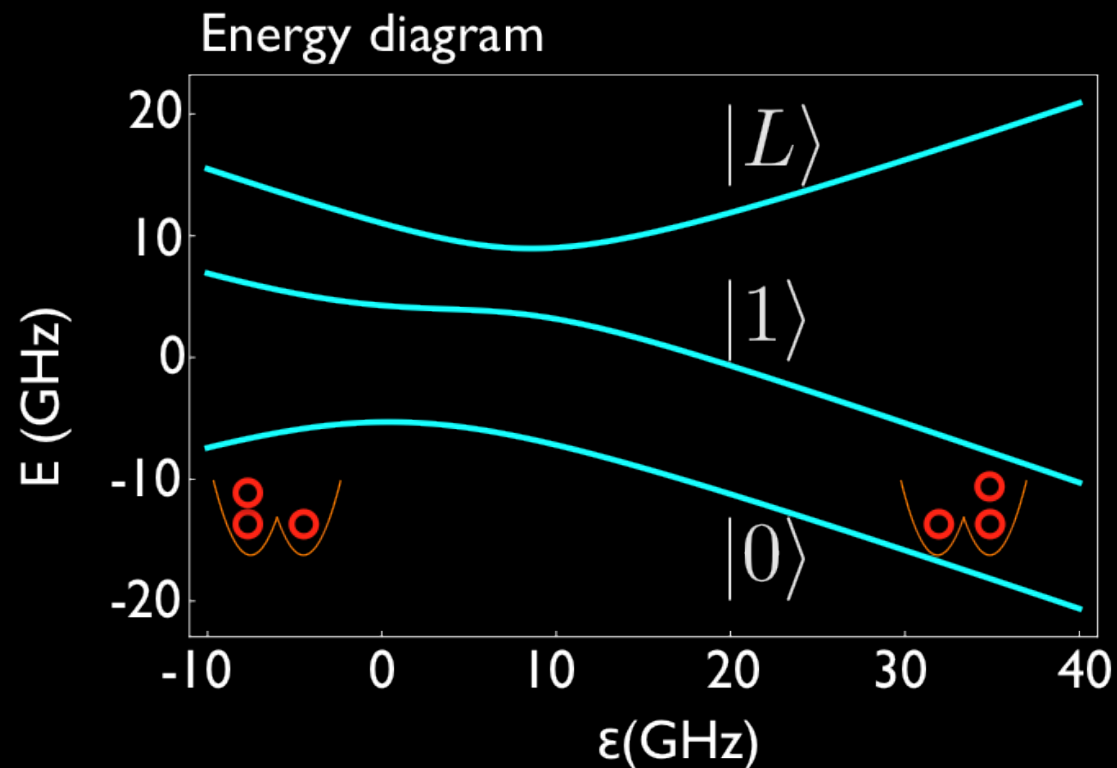
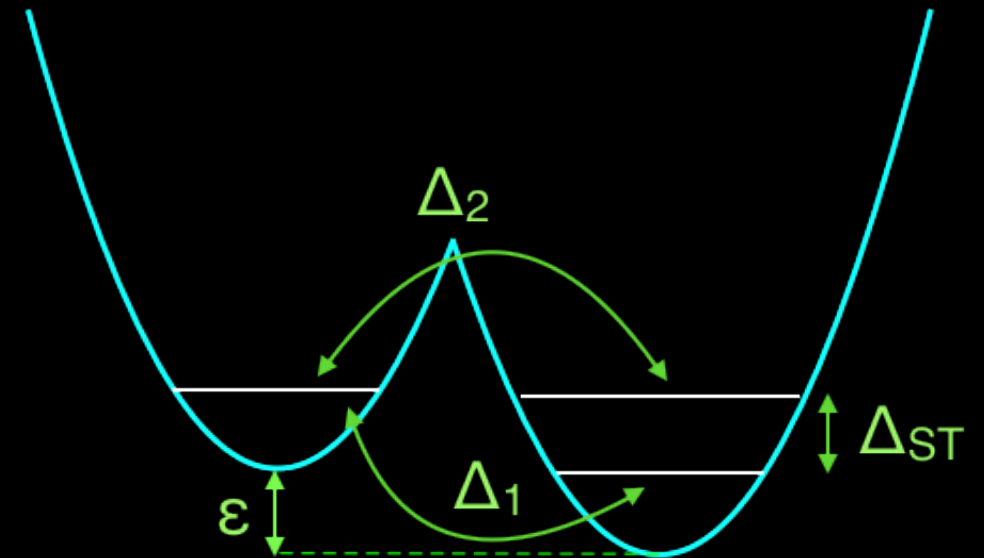
X. Mi et al. Nature **555**, 599-603 (2018)

Samkharadze et al. Science 10.1126/science (2018)



Hybrid qubit

$$H = \begin{pmatrix} -\varepsilon/2 & 0 & \Delta_1 \\ 0 & -\varepsilon/2 + \Delta_{ST} & \Delta_2 \\ \Delta_1 & \Delta_2 & \varepsilon/2 \end{pmatrix}$$

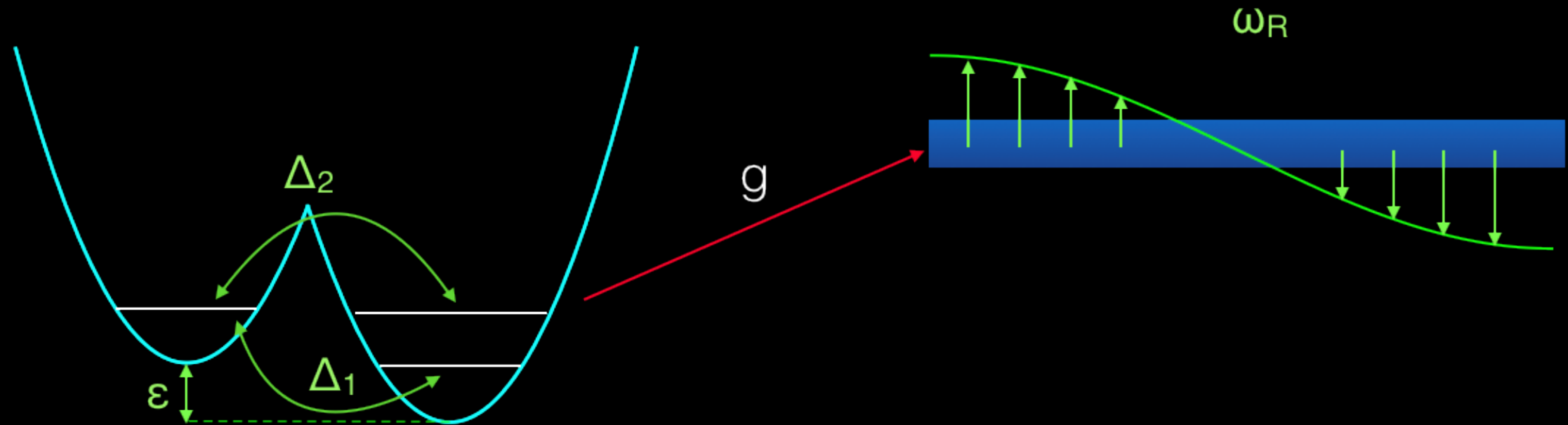


Two anticrossings related to the qubit parameters

Far-detuned regime

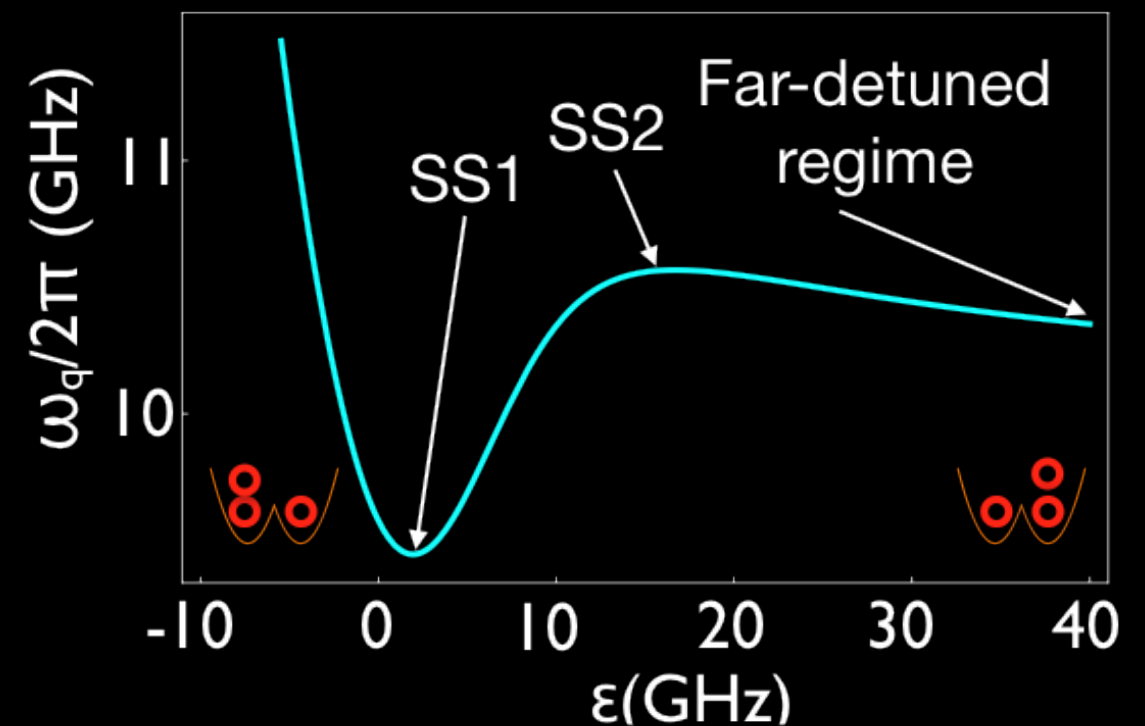
Dipolar coupling maximized near the 1st anticrossing

Strategy: Maximize g/γ

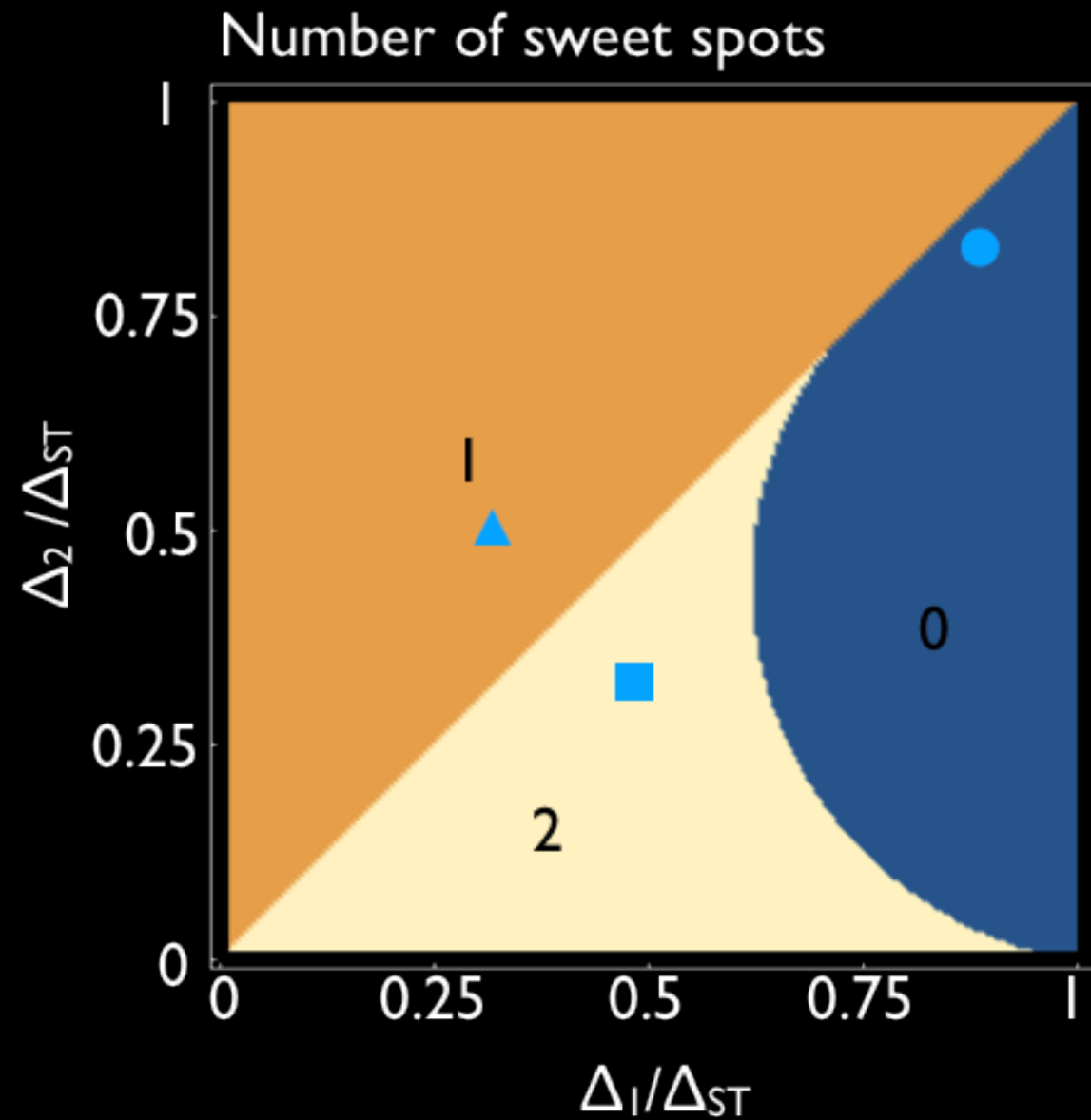


Achieving strong coupling requires $g/\gamma > 1$

Strong coupling candidates: SSs, far-detuned regime



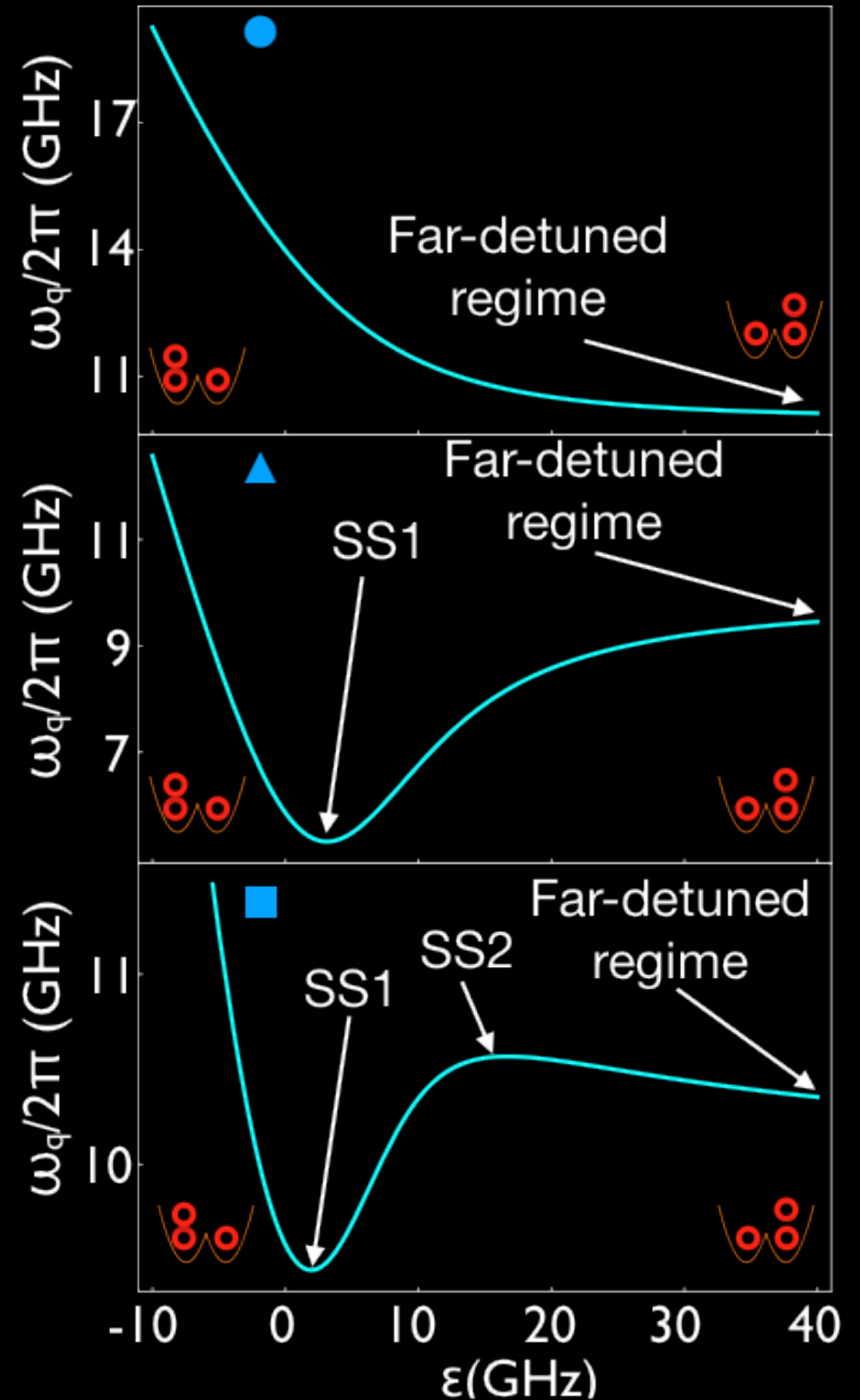
Basic properties of the HQ



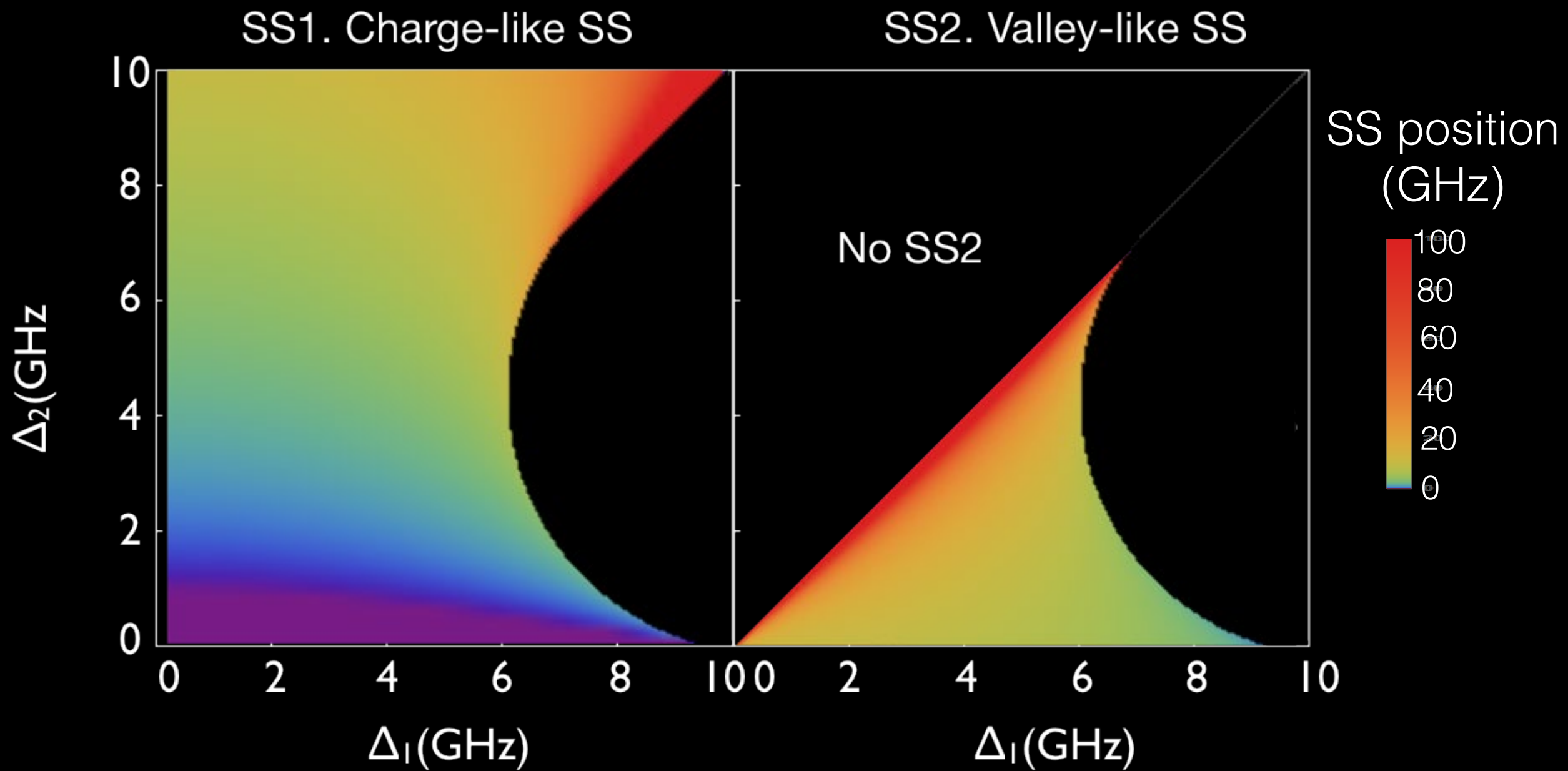
SS1: higher coupling, bad coherence

SS2: Smaller coupling with improved coherence

Far-detuned regime: Small coupling with good coherence

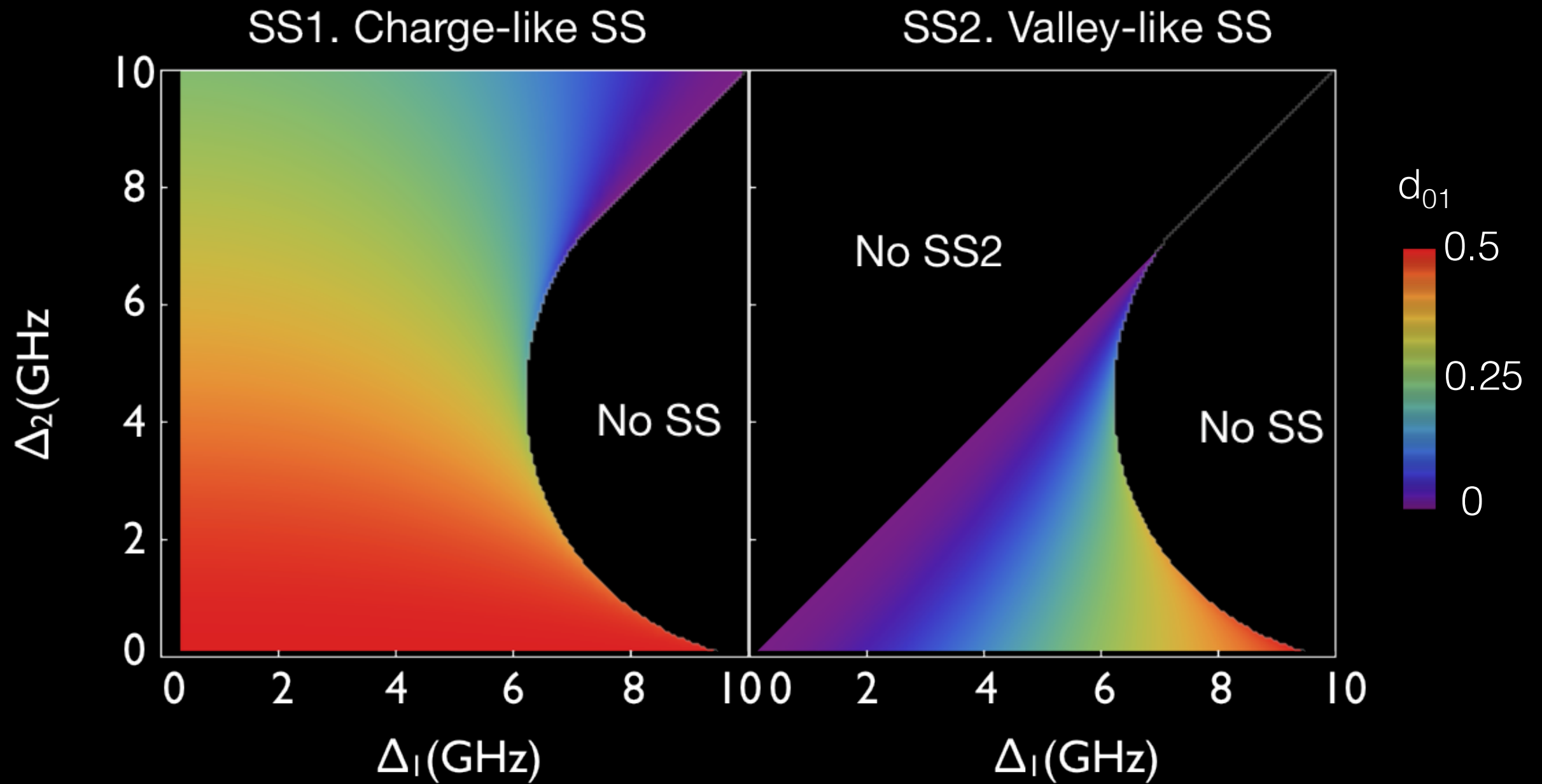


Sweet spot position.



Sweet spot properties. $V_S=10$ GHz

Dipolar coupling

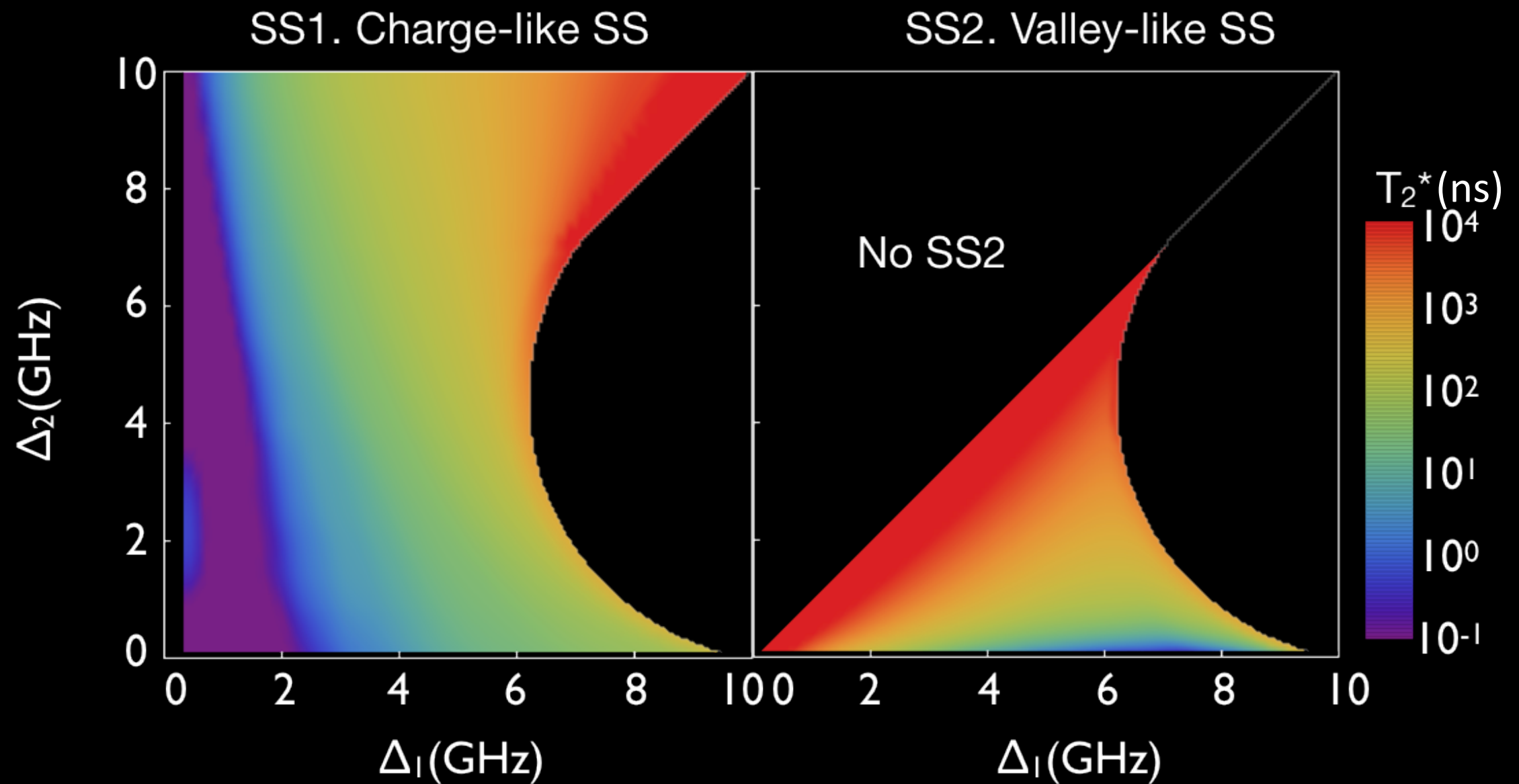


Sweet spot properties. $V_S=10$ GHz

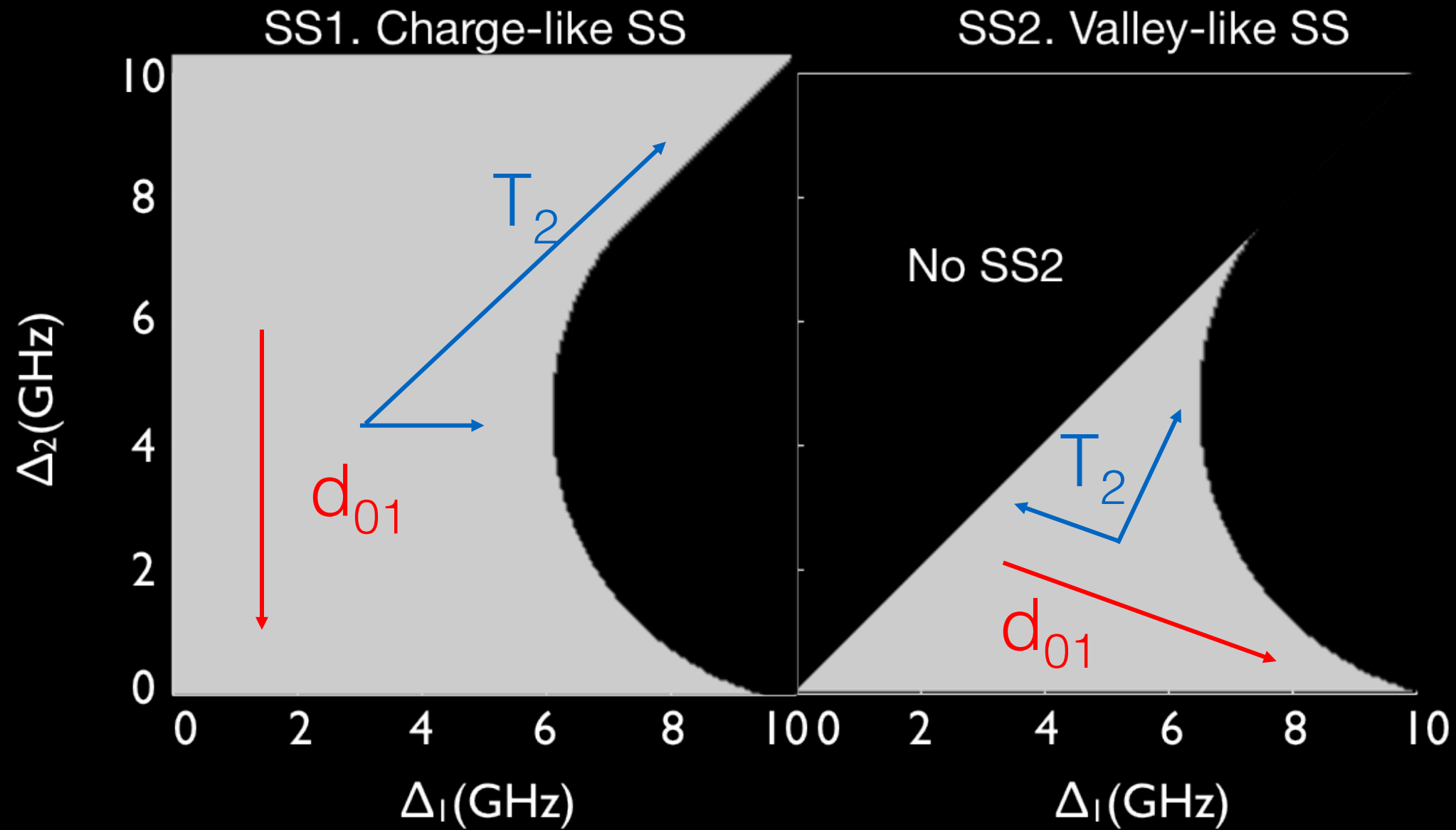
We estimate coherence times by simulating $1/f$ noise

$$\sigma_\varepsilon = 2\mu eV$$

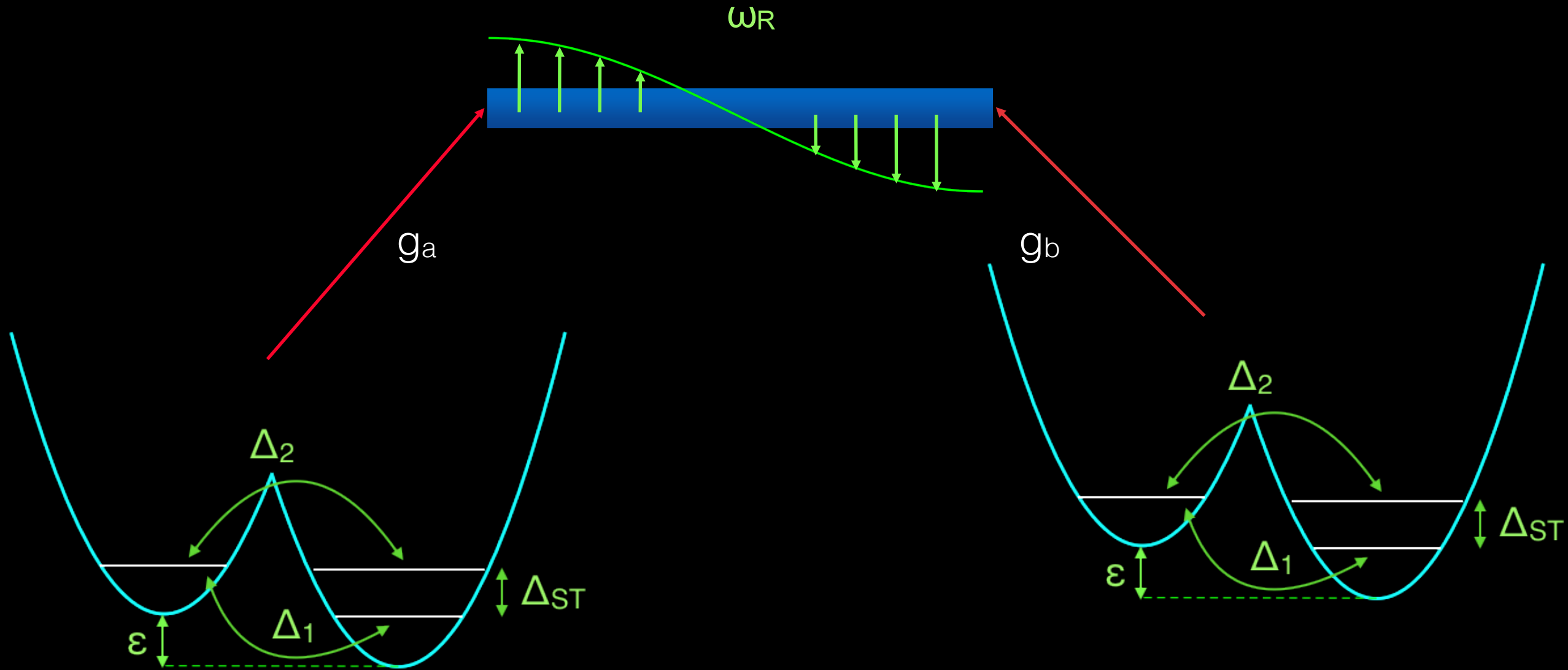
Coherence times



Overview of the sweet spot properties



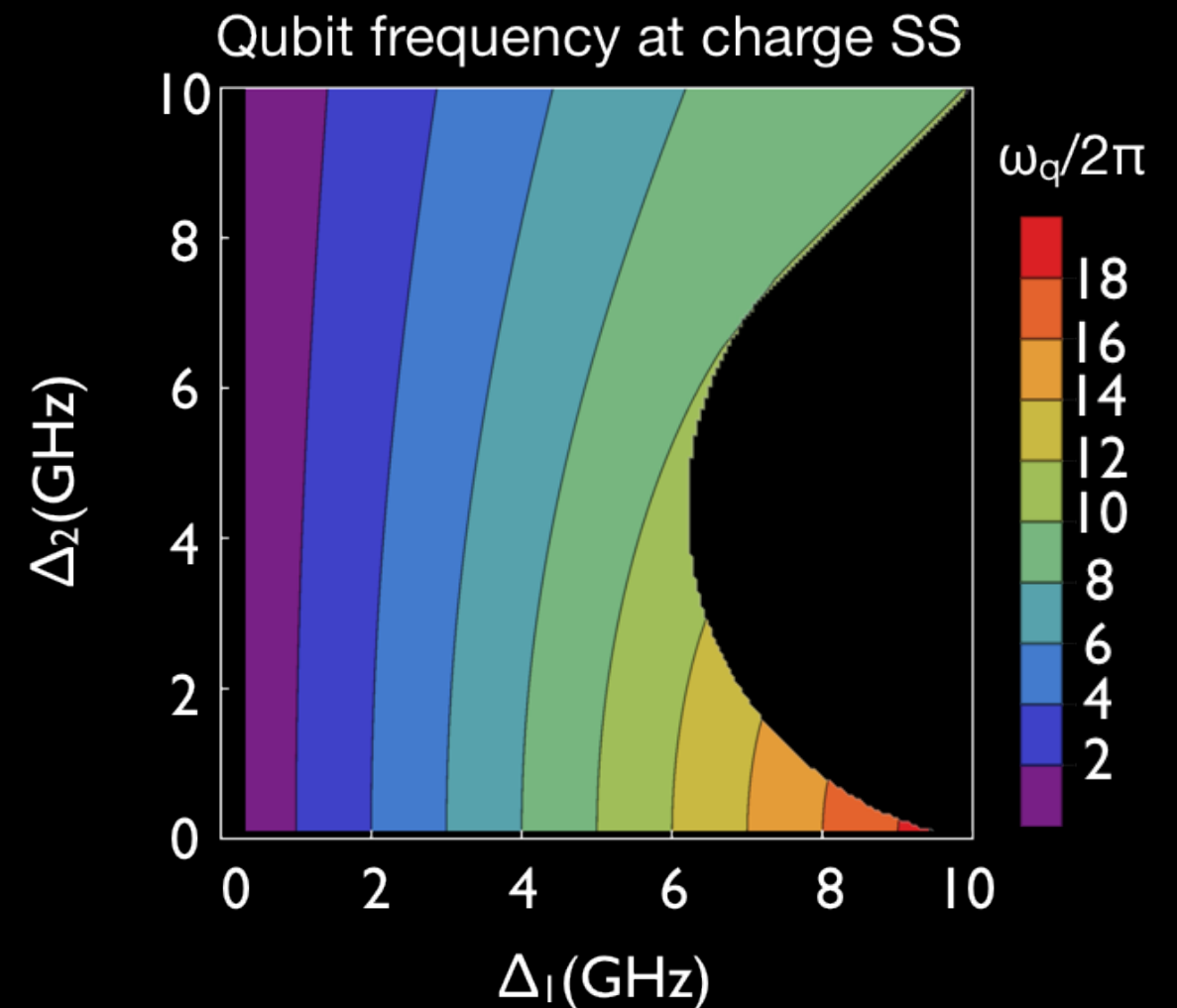
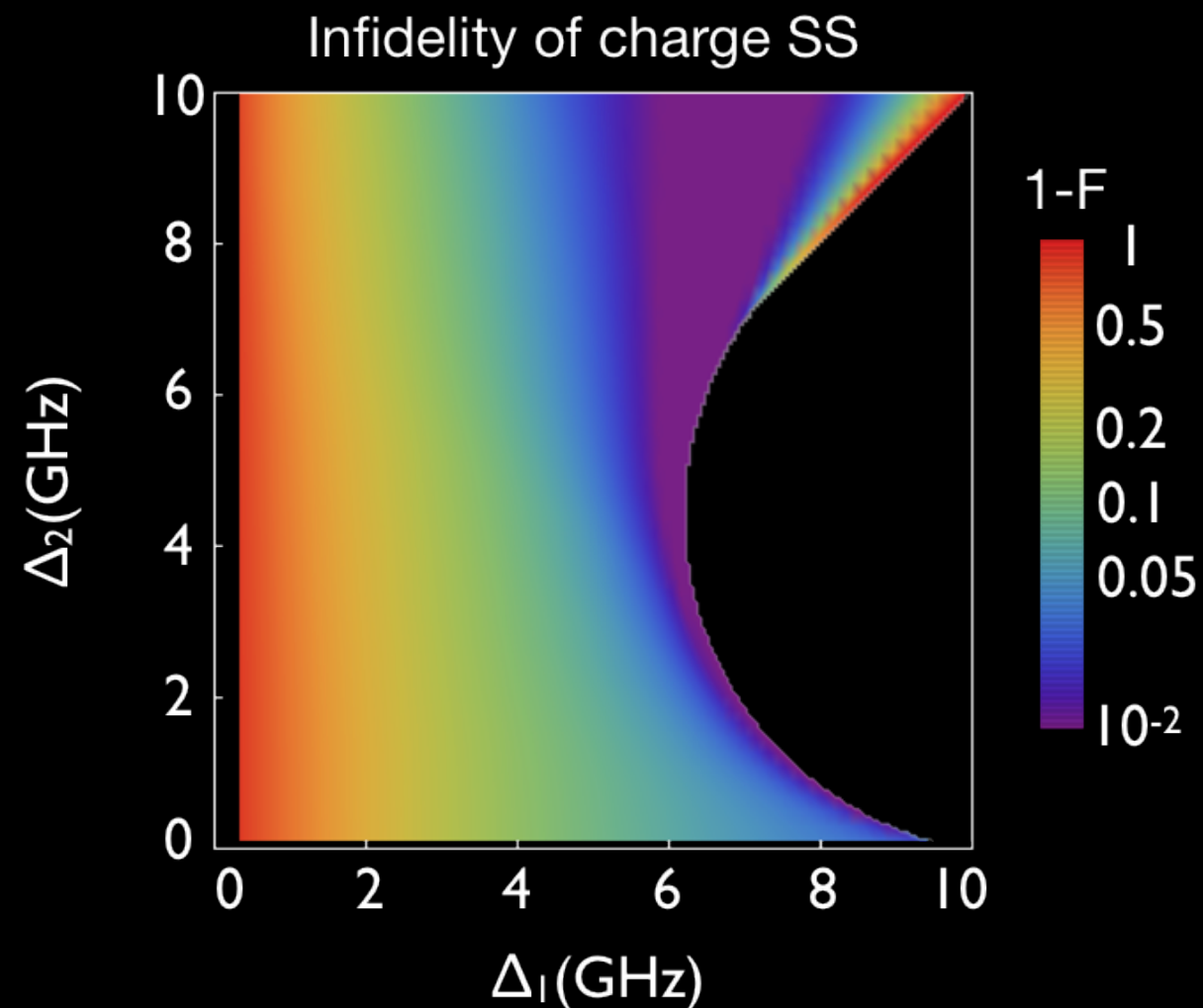
Coupling two hybrid-qubits through a resonator



Upper bound two-qubit gate fidelities. Charge SS

Resonant operation with $\kappa=0$ gives the upper bound gate fidelities for a $01 \leftrightarrow 10$ operation

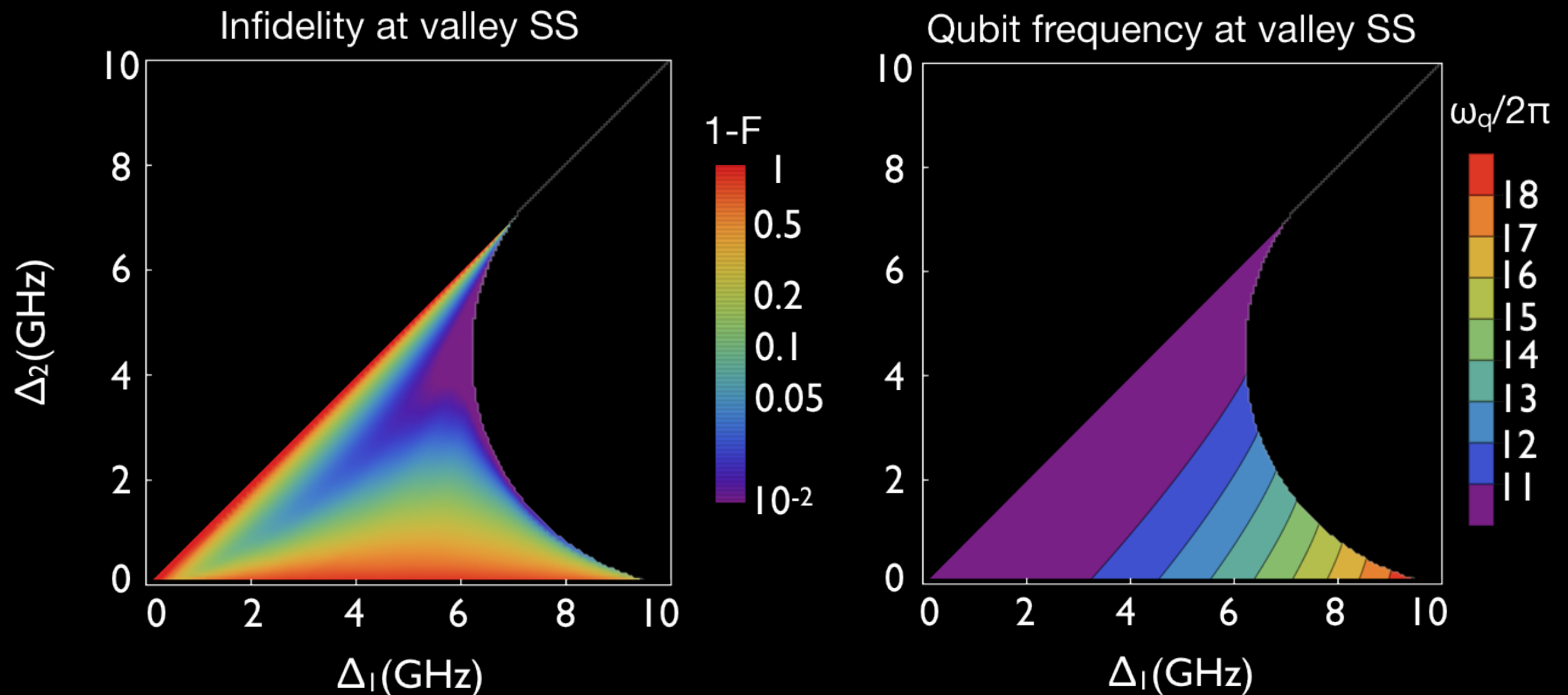
$$H = \omega_c a^\dagger a + \sum_{I=1,2} \frac{\omega_{q,i}}{2} \sigma_z^i + g_0 d_{01}^i (a + a^\dagger) \sigma_x^i$$



Upper bound two-qubit gate fidelities. Valley SS

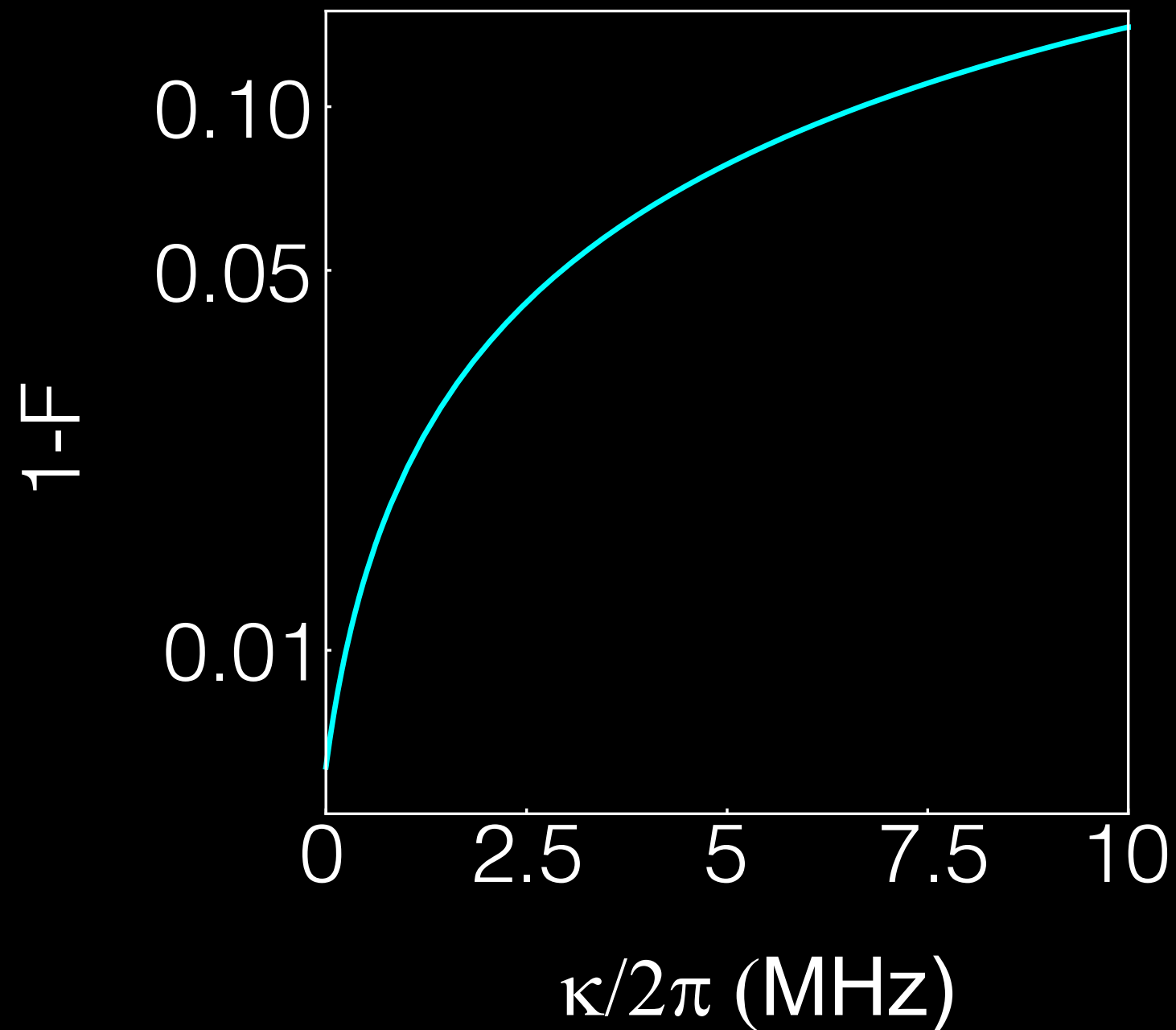
Resonant operation with $\kappa=0$ gives the upper bound gate fidelities for a $01 \leftrightarrow 10$ operation

$$H = \omega_c a^\dagger a + \sum_{I=1,2} \frac{\omega_{q,i}}{2} \sigma_z^i + g_0 d_{01}^i (a + a^\dagger) \sigma_x^i$$



Infidelity due to finite Q

Realistic resonators have finite decay rates (1-10 MHz)



SB transitions mediated by driving detuning and tunnel coupling

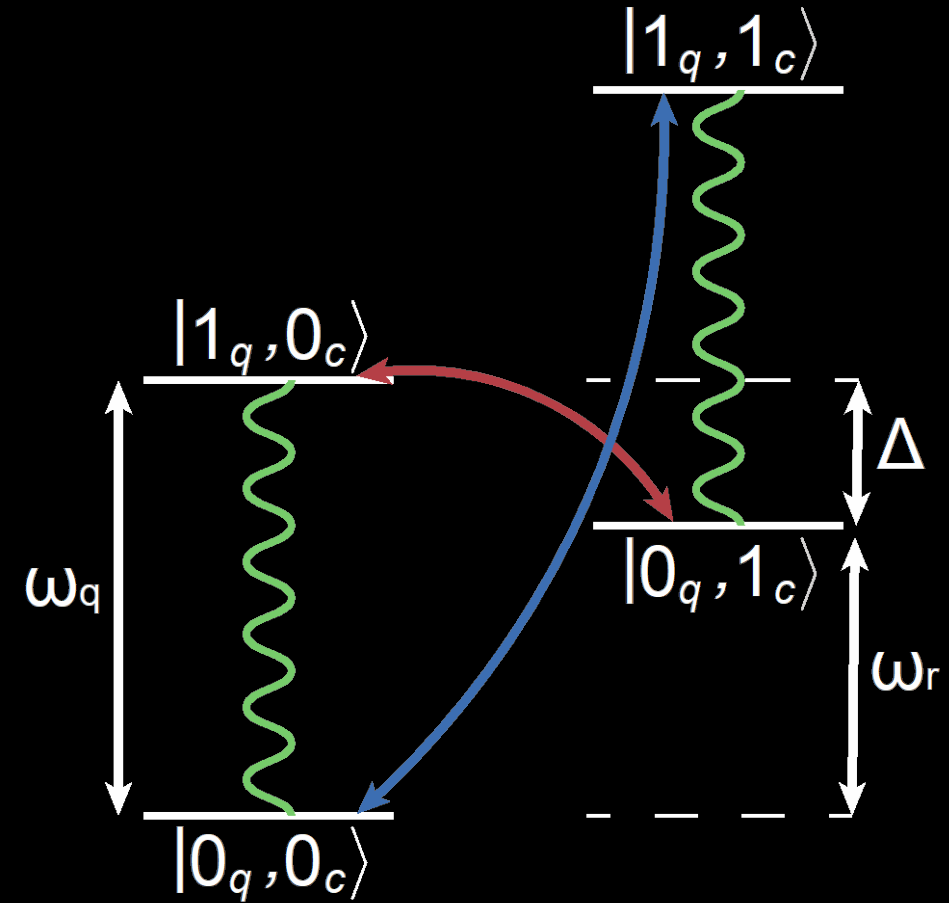
$$H = \omega_c a^\dagger a + \frac{\omega_{q,i}}{2} \sigma_z^i + g_0 d_{01} \sigma_x^i (a + a^\dagger) + d_{01} \delta \varepsilon \cos(\omega_{q,i} t + \phi) \sigma_x^i$$

$$\omega_c - \omega_q = \pm d_{01} \delta \varepsilon$$

$$H_{S+} = \frac{g}{2} (e^{i\phi} \sigma^- a + e^{i\phi} \sigma^+ a^\dagger)$$

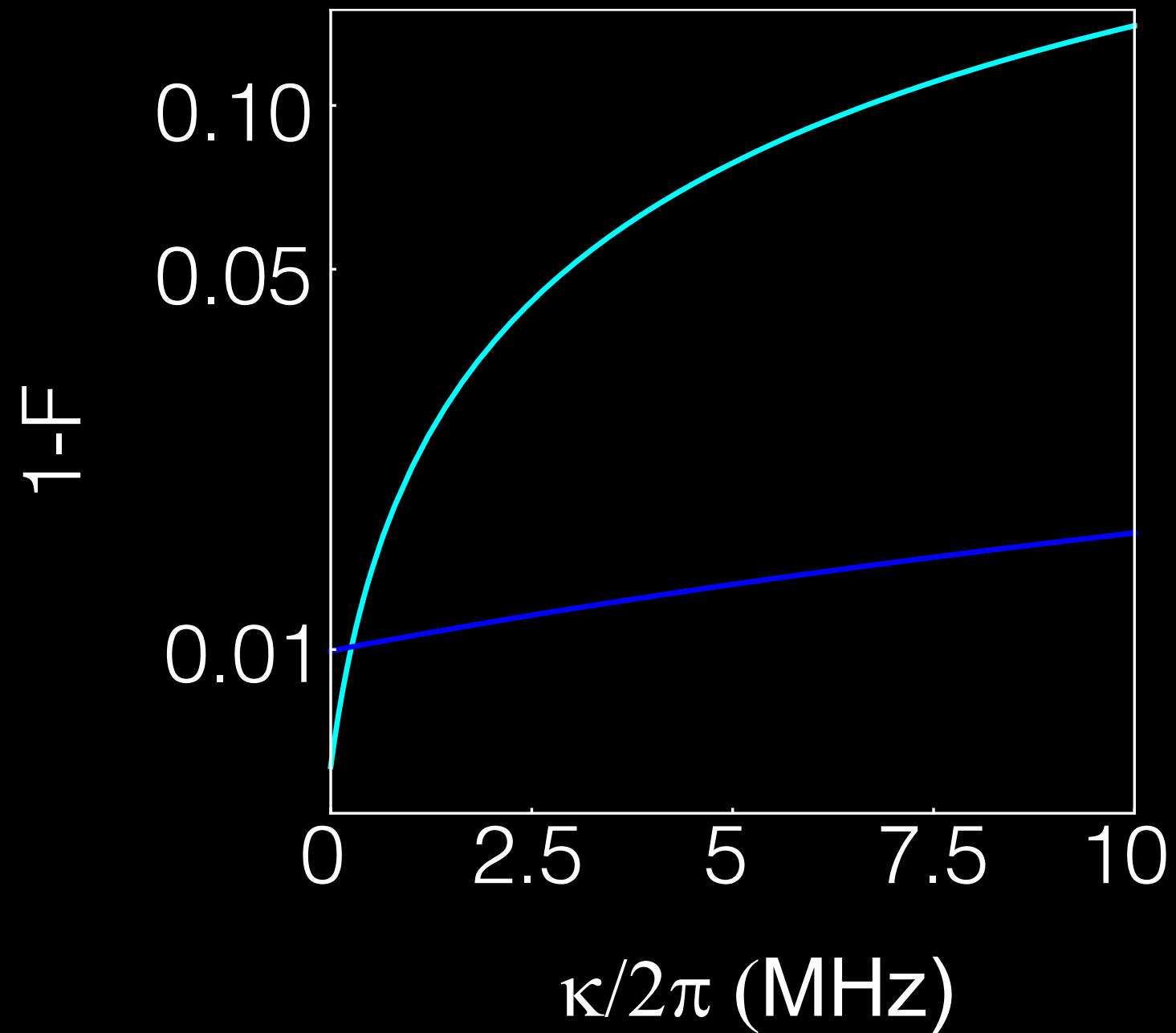
$$H_{S-} = \frac{g}{2} (e^{i\phi} \sigma^+ a + e^{i\phi} \sigma^- a^\dagger)$$

This method allows to work detuned away from resonance while keeping fast gates



$$U_{CZ} = R_{z,a}\left(\frac{\pi}{\sqrt{2}}\right) S_a^+(\pi, \phi) S_a^-(\pi/2, 0) S_a^-(\pi\sqrt{2}, \pi/2) S_a^-(\pi/2, 0) R_{z,b}(\pi/\sqrt{2}) S_a^+(\pi, \phi)$$

Sideband transition fidelities decay slower



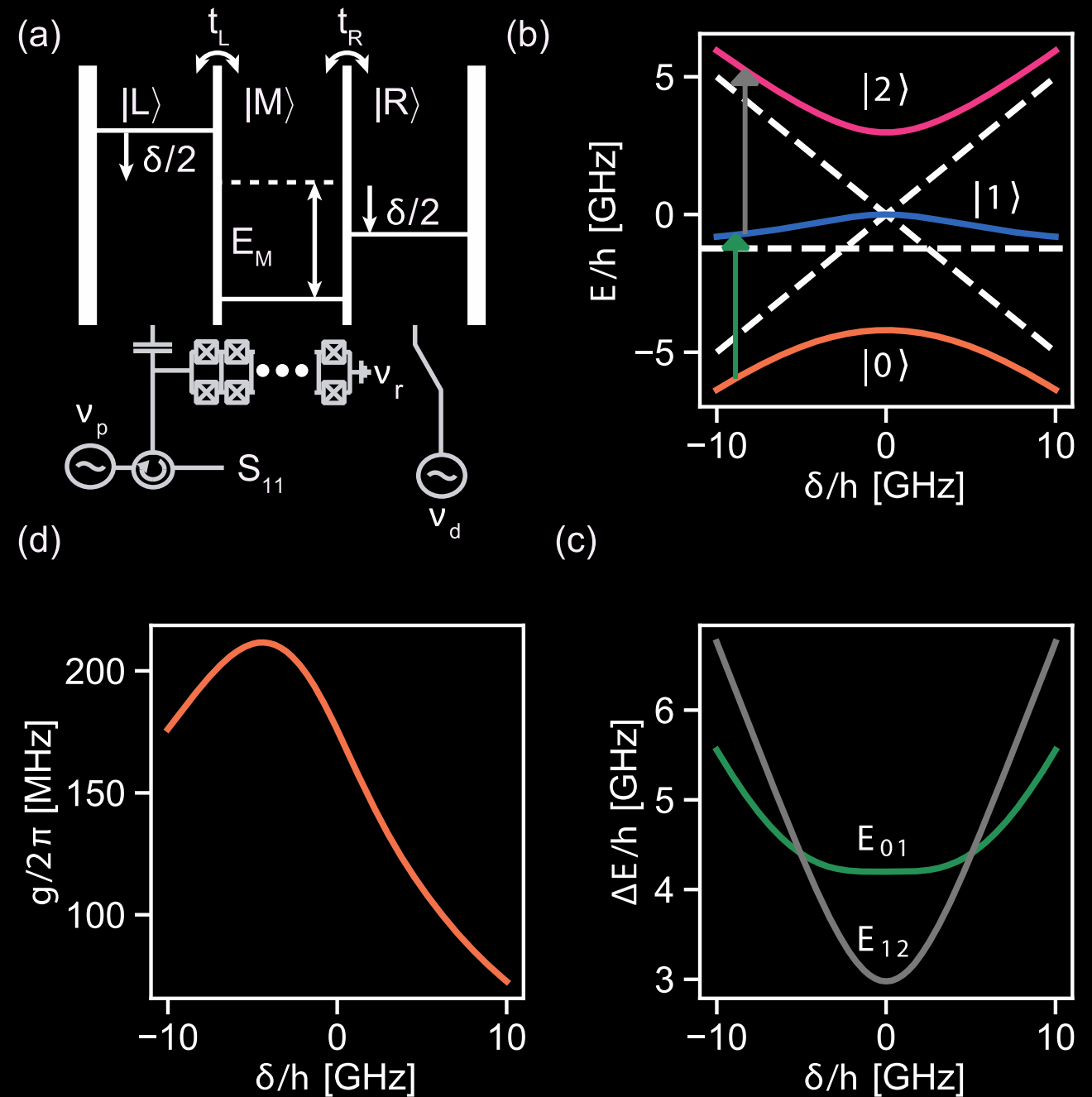
Coupling a Charge quadrupole qubit to a SC resonator

Phys. Rev. Research 3 (1), 013171 (2021)

$$H = \begin{pmatrix} \delta/2 & t_L & 0 \\ t_L & E_M & t_R \\ 0 & t_R & -\delta/2 \end{pmatrix}$$

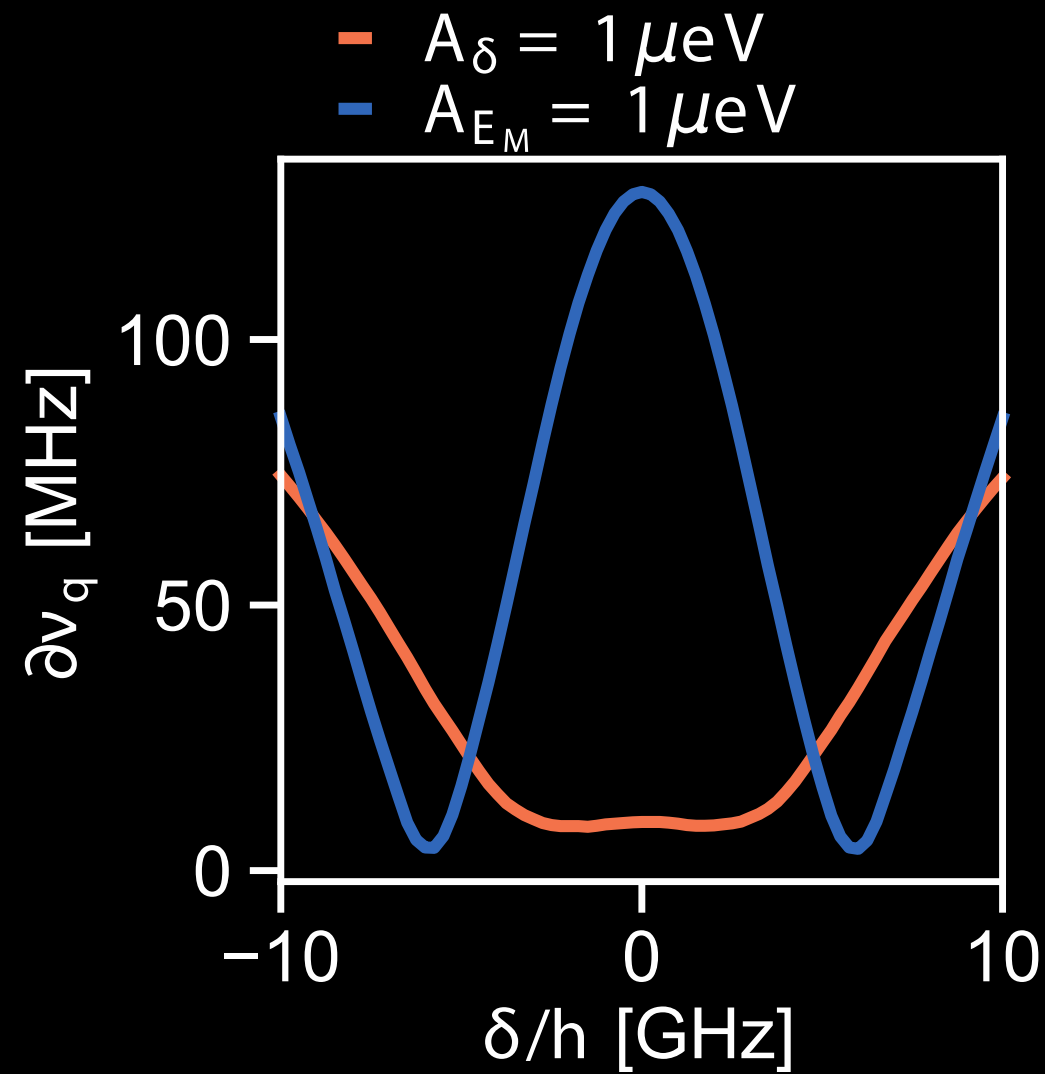
$$E_M^{\text{Opt}} \approx -0.493t$$

At this optimal point, the qubit is in a 2nd order SS in dipolar detuning



Operating in the quad detuning vs dipolar detuning

If noise is dominated by long-range charged sources
dipolar noise dominates



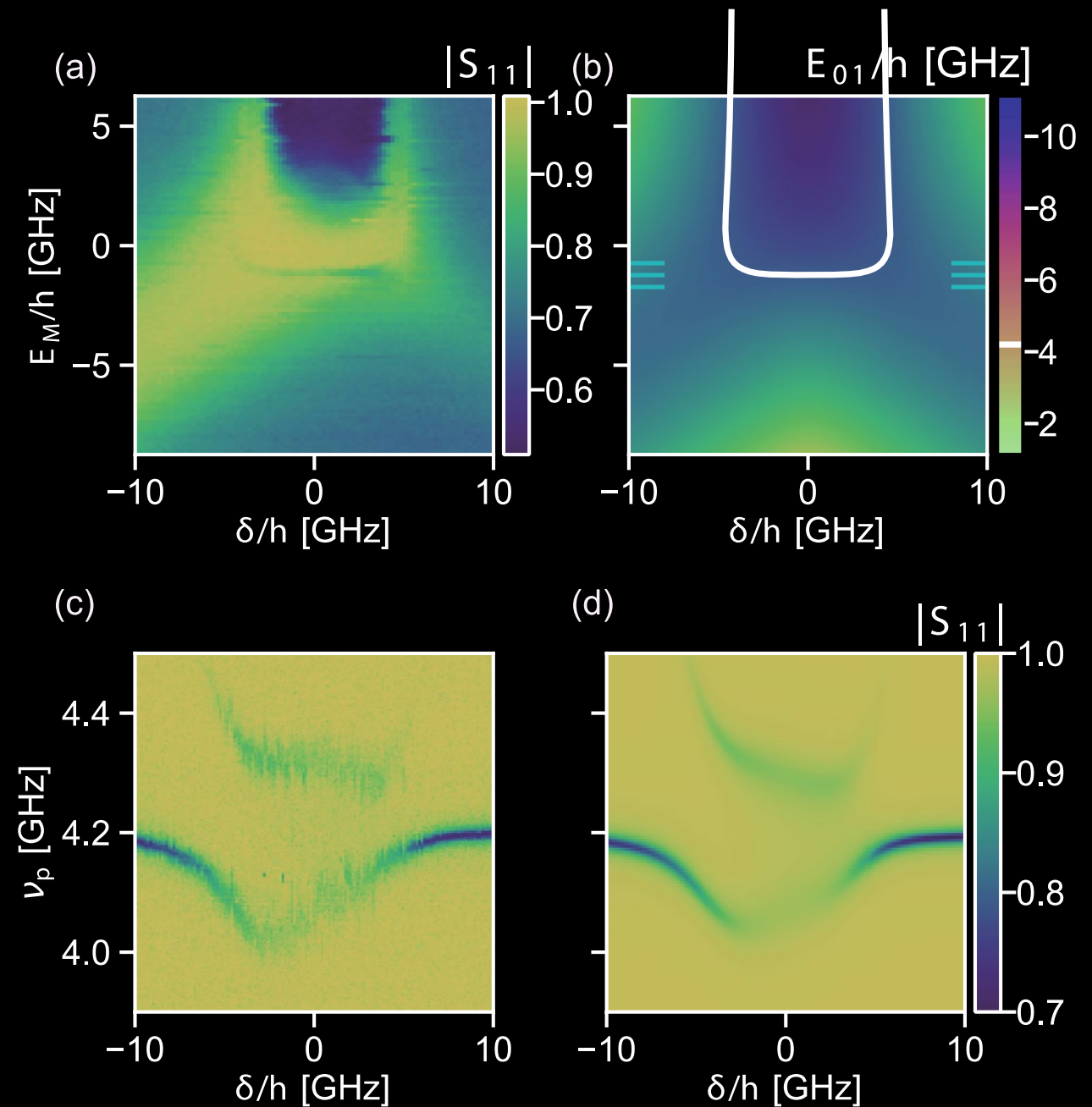
Experimental results and fitting

Fitting to IO theory yields

$$g/2\pi = 203 \text{ MHz}$$

$$\kappa_{\text{int}}/2\pi = 14 \text{ MHz}$$

$$\kappa_{\text{ext}}/2\pi = 1.3 \text{ MHz}$$



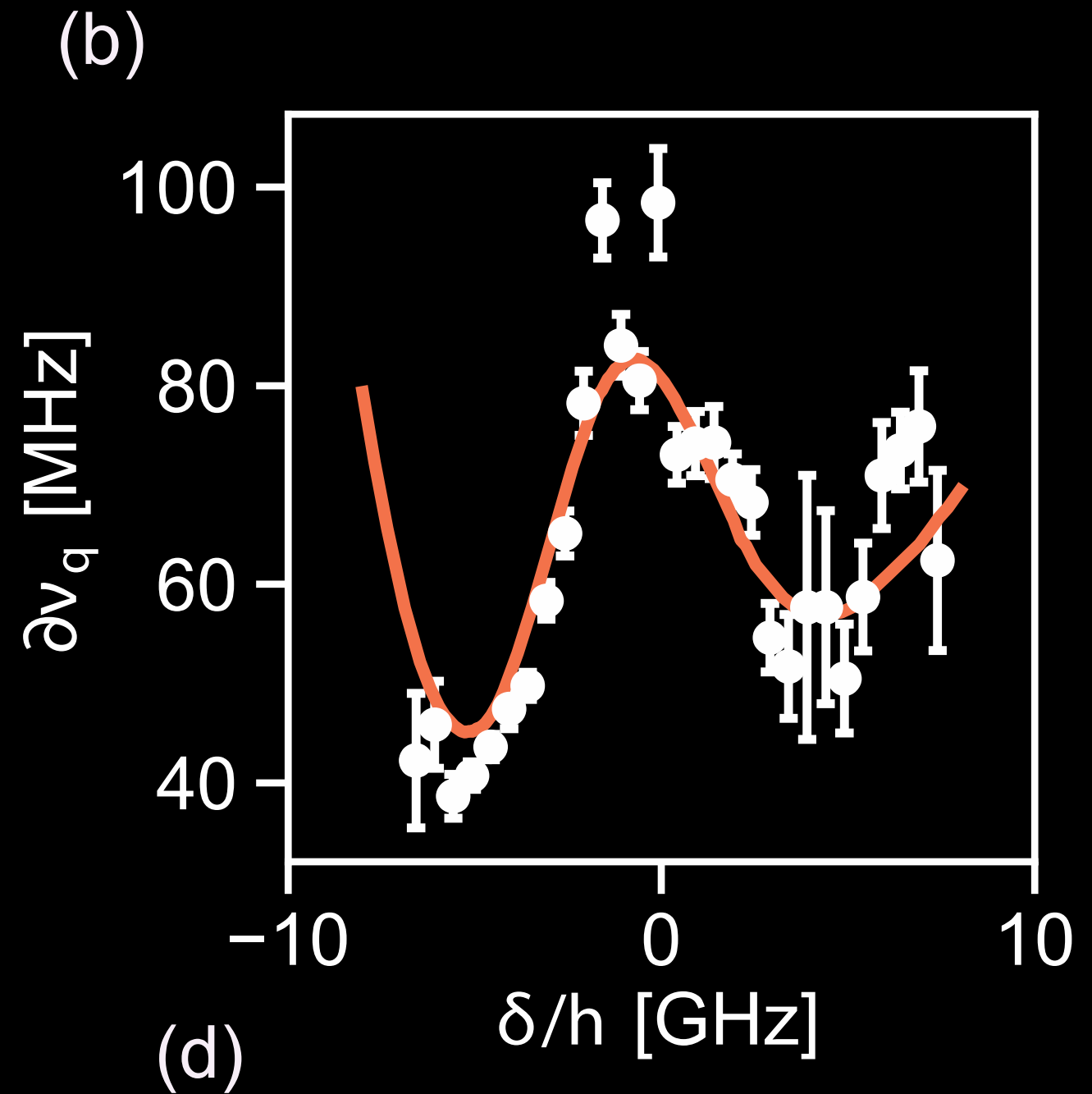
Noise characterization

1/f noise simulations allow to fit the decoherence rate

$$A_{\delta}=1.95 \text{ } \mu\text{eV}$$

$$A_{\text{EM}}=0.935 \text{ } \mu\text{eV}$$

$$\rho=-0.084$$



Conclusion

Superconducting resonators are a promising technology for long-range two-qubit gates

Coupling a HQ to a SCR imposes a frequency constraint.
Can be reduced by operating off-resonance

The CQ qubit allows to spatially characterize charge noise.
The experimental results seem to point towards close-range noise