

# **Data Analysis Techniques for Continuous Gravitational Wave Searches**

Joseph Bayley

Submitted in fulfilment of the requirements for the  
Degree of Doctor of Philosophy

School of Physics and Astronomy  
College of Science and Engineering  
University of Glasgow



University  
of Glasgow

February 2017

# **Abstract**

# Contents

<b>Abstract</b>	<b>i</b>
<b>Acknowledgements</b>	<b>x</b>
<b>Declaration</b>	<b>xi</b>
<b>1 Introduction</b>	<b>1</b>
1.1 Gravitational waves . . . . .	2
1.2 Sources and signals . . . . .	5
1.2.1 Transient . . . . .	5
1.2.2 Stochastic . . . . .	7
1.2.3 Continuous waves . . . . .	7
1.3 Detectors . . . . .	11
1.3.1 Laser Interferometers . . . . .	11
1.4 Bayes Theorem . . . . .	14
<b>2 Searching for continuous gravitational waves</b>	<b>15</b>
2.1 Continuous signal model . . . . .	15
2.1.1 Basic probability . . . . .	17
2.1.2 Bayesian Inference . . . . .	18
2.2 Continuous wave searches . . . . .	20
2.2.1 Fully coherent . . . . .	20
2.2.2 Semi coherent . . . . .	21
2.3 Motivation . . . . .	23
<b>3 SOAP: A generalised application of the Viterbi algorithm to search for continuous gravitational wave sources.</b>	<b>24</b>
3.1 Introduction . . . . .	25
3.2 Viterbi algorithm . . . . .	27
3.3 The transition matrix . . . . .	28
3.4 Single detector . . . . .	29

3.5	Multiple detectors . . . . .	33
3.6	Memory . . . . .	35
3.7	Summed input data . . . . .	36
3.8	Line-aware statistic . . . . .	37
3.9	Line aware statistic for consistent amplitude . . . . .	41
3.10	Testing the algorithm . . . . .	44
3.10.1	S6 injections into gapless Gaussian noise . . . . .	47
3.10.2	S6 injections into Gaussian noise with gaps . . . . .	49
3.10.3	Tests on the S6 MDC . . . . .	52
3.11	Optimisation of Line-aware statistic. . . . .	53
3.11.1	S6 injections into gapless Gaussian noise and Gaussian noise with gaps . . . . .	53
3.11.2	S6 MDC injections . . . . .	53
3.11.3	Sensitivity with frequency . . . . .	56
3.11.4	Searching for non CW sources . . . . .	57
3.11.5	Computational cost . . . . .	57
3.12	Discussion . . . . .	59
<b>4</b>	<b>Machine learning for continuous wave searches</b>	<b>61</b>
4.1	Introduction . . . . .	61
4.2	Neural networks . . . . .	62
4.2.1	Neurons . . . . .	63
4.2.2	Network structure . . . . .	63
4.2.3	Activation functions . . . . .	64
4.2.4	Training . . . . .	65
4.2.5	Loss functions . . . . .	65
4.3	Convolutional Neural Networks . . . . .	66
4.3.1	Convolutional layers . . . . .	66
4.3.2	Max pooling layers . . . . .	67
4.3.3	CNN structure . . . . .	67
4.4	CW search . . . . .	68
4.4.1	Network structure . . . . .	68
4.5	Network Visualisation . . . . .	68
4.6	Data and simulations . . . . .	68
<b>5</b>	<b>Detector Characterisation with SOAP</b>	<b>72</b>
5.1	Instrumental lines . . . . .	72
5.2	Identifying and monitoring instrumental lines . . . . .	75
5.3	Identifying and cleaning lines with SOAP . . . . .	76
5.4	Summary pages . . . . .	78

<b>6</b>	<b>Summary</b>	<b>80</b>
<b>A</b>	<b>Continuous gravitational wave injections</b>	<b>81</b>
A.0.1	CW Signal . . . . .	81
A.0.2	Time series and complex fast Fourier transform (FFT) injections . . . .	81
A.0.3	Spectrogram injections . . . . .	81
A.0.4	Signals in data . . . . .	82

# List of Tables

2.1	From [48], shows the computational cost for the first 4 months of advanced Laser Interferometer Gravitational-wave Observatory (LIGO) for each search. One million standard units (MSU), where one standard unit is one core-hour on a standard core. ‘ Expected computational costs of searches using the first four months of advanced LIGO data with each search pipeline. These estimates are for a different data observing time from that of the mock data challenge (MDC), and do not cover the same parameter space as each other or the MDC. The Einstein@Home searches uses the computing resources of the Einstein@Home project and is designed to run for 6 - 10 months in the Einstein@Home grid.’ .	23
3.1	Table shows the ranges of the search parameters and their optimised values for injections into gapless Gaussian noise, Gaussian noise with gaps and the S6 MDC. For gapless Gaussian noise and Gaussian noise with gaps, there are 10 parameter values spaced linearly between the limits. For the S6 MDC the parameters, $\tau$ , $w_L$ and $w_S$ were distributed in log space between the limits and $p(M_L)/p(M_N)$ is distributed uniformly. . . . .	50

# List of Figures

1.1	Shows how the plus and cross polarisation's affect a ring of test particles. This assumes the wave is travelling out of the page and the effects have been greatly exaggerated. This also shows an example of how this effects the test masses of an interferometer. This will be described in more detail in Sec. 1.3. . . . . .	4
1.2	Each gravitational-wave (GW) signal type can be categorised based on its signal length and how well the signal is modelled. Transient signals which are short duration, include both well modelled compact binary coalescence (CBC) signals and unknown Burst signals. Long duration signals include well modelled continuous wave (CW) signals and unknown Stochastic signals. . . . .	6
1.3	The neutron star has some ellipticity which mean that the star deforms such that it is not symmetric around its rotation axis. This is thought of as a 'mountain' on the surface of the neutron star. The diagram above shows an extremely exaggerated view of what this could look like. . . . .	9
1.4	The r-modes can travel in either direction in the star. in the case when the r-mode is moving clockwise and the neutron star is moving anticlockwise, if the neutron star is rotating fast enough it can cause unstable emission of GW. This image was reconstructed from [30]. . . . .	10
1.5	This figure shows a basic interferometer.ETMY and ETMX refer to the external test masses, which are just mirrors at the end of the interferometer arms. ITMY and ITMX refer to the internal test masses, these create a cavity in the interferometers arms which can build up laser power. BS is the beam splitter which splits the Laser beam equally to each arm, this them recombines the beams back to the detector. . . . .	12
1.6	The antenna response is shown as in [23] for the plus, cross and average polarisation's. The detectors arms lie on the x and y axis in the above plots. . . . .	13
1.7	The different noise sources affect the sensitivity of the LIGO detectors at different frequencies. This shows the various sources how the affect the noise curve [10]. . . . .	14

3.1	This figure shows an example of a time-frequency spectrogram which is typical of LIGO data which is searched through. Here an instrumental line has been injected at 100.006 Hz. The white track shows a random walk track through this spectrogram whereas the red line shows the track which gives the highest sum of detector power. The second panel shows a histogram of a subset of all paths which can be found through the given spectrogram from start to finish. This is a subset as the total number of paths is too large to calculate. The value of the statistic which comes from the optimal path is $\sim 1700$ . This is much larger than any of the random tracks in our subset and much larger than the mean of all tracks. . . . .	25
3.2	Fig. 3.4c shows the observed data, i.e the log-likelihood values $C_{j,k}$ . Fig. 3.2b shows the calculated log-probabilities $V_{j,k}$ . $A_{j,k}$ is shown in parentheses, where the up, centre or down (UCD) components correspond to $i = [-1, 0, 1]$ respectively. The red line shows the path that gives the maximum probability. The transition matrix for the UCD jumps is $[0, 1, 0]$ and corresponds to the un-normalised prior log-probabilities of these jumps occurring. . . . .	31
3.3	Lookup tables using the line aware statistic in Eq. 3.29. . . . .	42
3.4	Lookup tables using the line aware statistic for consistent amplitude as in Sec. 3.9. Each of these use the parameters $p_s(\lambda) = 4, p_l(\lambda) = 5$ and $p(M_L)/p(M_G) = 0.03$ . . . . .	45
3.5	The results that the SOAP algorithm returns from an injection with an optimal signal-to-noise-ratio (SNR) of 90, i.e., the SNR in H1 is 64 and the SNR in L1 is 62. The signal is injected into Gaussian noise, where the 1800 s short Fourier transforms (SFTs) have been summed over 1 day. The top panel shows a simulation of summed SFTs from H1, the second panel shows the same for L1, the third panel shows the values proportional to the log-odds ratios in Eq. 3.24. The log-odds have been normalised such that the sum of all the odds ratios in every time bin are equal to 1. The bottom panel shows the injected signal track (black dotted) and the track found in the ‘imaginary’ detector by the two-detector SOAP search with the line-aware statistic (red), both of these tracks are at the geo-centre. In this case the root median square (RMS) of the difference between the Viterbi track and injected signal track was $\sim 1$ bin, where 1 bin is 0.00056 Hz wide. . . . .	48



- 3.6 Panels 3.6a and 3.6c show the detection efficiency as a function of SNR and depth respectively. Here SNR is the the integrated SNR which we would expect to recover from the available data. The four curves refer to injections into gapless Gaussian noise (red), Gaussian noise with gaps in data, where the noise floor is either fixed (blue-dashed) or it is moving with time (orange) in the same way as the S6 MDC and injections into real data i.e., the S6 MDC. In the gapless Gaussian noise case, the recovered integrated SNR refers to the SNR the injection would have if it had the same amount of data as in the cases with gaps. The curves are made by fitting a sigmoid Eq. 3.39 to binomial detection data with a 1% false alarm rate, as explained in Sec. 3.10.1, the error bounds are the 5% and 95% intervals. At 95% efficiency and a 1% false alarm rate, this shows we can detect to an SNR of  $\sim 60$  and a sensitivity depth of  $\sim 34 \text{ Hz}^{-1/2}$  for gapless Gaussian noise and an SNR of  $\sim 69$  and  $72$  and a sensitivity depth of  $\sim 13 \text{ Hz}^{-1/2}$  and  $\sim 10 \text{ Hz}^{-1/2}$  for the Gaussian with gaps case with fixed noise floor and moving noise floor respectively. For the S6 MDC we can detect an SNR of  $\sim 74$  and a sensitivity depth of  $\sim 13 \text{ Hz}^{-1/2}$ . Panels 3.6b and 3.6d show the RMS of the difference between the injected signal track and the track found by SOAP as a function of SNR and sensitivity depth respectively. This is shown in units of bins where each bin is  $0.00056 \text{ Hz}$  wide. . . . . 51
- 3.7 In Gaussian noise the transition matrix parameter  $\tau$  and the width of the prior on the signal case  $w_S$  were optimised. The key part to remember when reading this plot is that the lower the value of SNR the better the search has performed. Therefore darker blue areas are when the search performed better. This map shows that there is a line in parameter space where the search performed best. Also in Gaussian noise, the search is not that sensitive to the choice of parameter. The red lines on here shows the parameters used in the searches in this section. . . . . 54
- 3.8 When using real S6 data, all four parameters of the search were optimised over on simulations in real data. The plot above shows the SNR at 80% efficiency for each of the parameters where the ranges are in Tab. ???. Lower values of SNR mean the search is performing better. The red lines show the parameters used in the searches in this section and the sections that follow. Whilst this does not seem optimal, the search does not underperform much when a poor choice of parameters is used. . . . . 55
- 3.9 This figure shows an example of when the line aware statistic is used compared to a version when it is not. . . . . 56

3.10	The sensitivity of the SOAP search in this configuration decreases as the frequency of the pulsar increases. This setup of data for the search however, was optimised for the 100-200 frequency band and can be changed for different frequencies. 3.10a shows the efficiency curves with 1% false alarm rate for each frequency. 3.10b shows the values from the efficiency curves at 90% efficiency.	58
4.1	Basic neuron . . . . .	64
4.2	A neural network is structured with layers. This consists of the input layer which is the data which you would like to analyse. Then this passes to a number of ‘hidden’ layers, in the above diagram there are two. These are just all layers between the input and the output. The output layer is then the desired output, above there is a single neuron as output, so this could perhaps classify the input to a value between 0 and 1. Every neuron in a layer is connected to the output of all neurons in the previous layer. . . . .	65
4.3	There are many different activations function which are used, and essentially any function can be defined if necessary. Above is shown a subset of the more commonly used functions. The linear function is not used, however, is there to compare to common non linear function. . . . .	69
4.5	Max pooling layer . . . . .	71
4.6	This shows a visualisation of the convolution neural network used for vtebri maps above . . . . .	71
5.1	There are many features in the LIGO detectors averaged spectral density. This image was taken from [83] where I have annotated some of the stronger lines. The power line from the mains in the USA is at 60 Hz. Some of the calibration lines are around 30 Hz, 331 Hz and 1083 Hz. The various violin modes of the suspensions are at 300, 500, 600 and 900 Hz where I have only marked the 500 Hz mirror suspension modes [83]. . . . .	73
5.2	Example of output plot from SOAP line search. . . . .	77
5.3	different types of instrumental line which appear in SOAP spectrograms . . . .	77
5.4	example plot of the summary pages . . . . .	79

# Acknowledgements

# Declaration

# Chapter 1

## Introduction

Gravitational waves were first predicted in 1915 as a consequence of Einstein's theory of general relativity [1]. They are theorised as ripples in the fabric of space-time. The first observational evidence that GW exist came from observations of the Hulse-Taylor binary [2, 3]. This observation showed a binary pulsar system whose separation was decreasing with time. If the separation of two orbiting objects is decreasing then the system is losing energy somewhere. The loss in energy matched the general relativity (GR) prediction which assumed the energy was lost to GW. This gave hope of GW existence and helped lead the way to designing instruments which could directly detect them. The first direct detection of gravitational waves was made in 2015 when the two LIGO detectors in the US [4] identified a signal from a binary black hole (BBH) system. This was not only the first observation of GW but gave information on a yet unobserved astrophysical system. This has since been followed by many more detections of BBH signals involving LIGO and Virgo including [5, 6]. In 2017 the LIGO detectors observed the first binary neutron star (BNS) system [7] which had a corresponding electromagnetic counterpart. This allowed verification of the source from optical counterparts and extended the era of multi-messenger astronomy. These detections opened up the field of gravitational wave astronomy, where many more detections are expected to give more information on the universe and objects within it.

As well as searching for BBH and BNS signals, there are many efforts to detect other types of GW signals. This thesis focuses on efforts to search for a particular type of GW which are thought to originate from rapidly rotating neutron stars. In this chapter 1 and 2 I will review introductory material. This includes a general introduction to the generation of gravitational waves in Sec. 1.1 and their sources in Sec. 1.2. I will then introduce instruments used to detect GW in Sec. 1.3. In Chapter 2 I will introduce the general model for CWs and current methods used to detect them. Chapters 3, 4 and 5 will go into detail about techniques developed by the author to search for CW signals. Finally I will summarise this work and discuss future developments in chapter 6.

## 1.1 Gravitational waves

In general relativity, gravity is thought of as the curvature of space-time and matter moves according to this curvature. The matter in the universe also has an effect on the curvature of the space-time. The larger the mass of matter the more the space-time is distorted. Space-time can generally be described by Einstein's field equations,

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}. \quad (1.1)$$

where  $G_{\mu\nu}$  is the Einstein tensor and  $T_{\mu\nu}$  is the stress-energy tensor. The stress energy tensor essentially describes the matter and energy in the universe. Its components contain information on the density of energy and momentum. The Einstein tensor contains information on the curvature of the universe. This can be derived directly from the metric tensor  $g_{\mu\nu}$  which describes the geometry of the universe. Einstein's equations then explain how the curvature of space-time changes with the energy and matter within it. In empty space one can assume that the geometry of space-time is flat, i.e. there is no curvature to space-time. The metric tensor for this can then be defined as,

$$g_{\mu\nu} = \eta_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}. \quad (1.2)$$

Each index of this matrix refers to a space-time dimension, i.e.  $x^0 = t$ ,  $x^1 = x$ ,  $x^2 = y$  and  $x^3 = z$ . Measuring a distance  $dx$  in space-time can be different for different observers, therefore, one needs a measure which is invariant for every observer. This is the space-time interval  $ds$ , also known as the line element, between two 'events' in space-time. This is defined as,

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu. \quad (1.3)$$

As in Einstein's notation this is a sum over the indices  $\mu$  and  $\nu$ . Eq. 1.3 is essentially Pythagoras's theorem, therefore, can be thought to describe the space-time 'distance' between the two events. For flat space-time,  $\eta_{\mu\nu}$ , this can then be written as,

$$ds^2 = -c^2 dt^2 + dx^2 + dy^2 + dz^2. \quad (1.4)$$

The Einstein equations Eq. 1.1 then demonstrate how the curvature of space-time  $G_{\mu\nu}$  depends on the matter and energy distribution  $T_{\mu\nu}$  within it.

A gravitational wave can be described as a ripple in this space time. The simplest way to visualise this is just a small time dependent change to the flat space-time metric  $\eta_{\mu\nu}$ . In

linearised theory of gravity, the space-time metric  $g_{\mu\nu}$  can be defined as,

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad (1.5)$$

where  $\eta_{\mu\nu}$  is the metric for flat space-time and  $h_{\mu\nu}$  is some perturbation, where  $|h_{\mu\nu}| \ll 1$  [8]. In this linearised theory the perturbations to the metric tensor are assumed to be small, therefore, Einstein's field equations can be solved such that the solution is a plane wave. I will skip over lots of the calculation here, however, more information on this derivation can be found in [8] **JOE: more references** By using  $g_{\mu\nu}$  from Eq. 1.5, we can write the linearised Einstein equations as,

$$\square h_{\mu\nu} = -16\pi T_{\mu\nu}, \quad (1.6)$$

where  $\square$  is the d'Alembert operator which in flat space is defined by,

$$\square = -\frac{1}{c^2} \frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad (1.7)$$

In empty space there is no matter, therefore, all the components of the stress energy tensor are zero, i.e.  $T_{\mu\nu} = 0$ . This allows Eq. 1.6 to be reduced to,

$$\square h_{\mu\nu} = 0. \quad (1.8)$$

This then takes the form of a common physics problem known as a wave equation. This follows the same form as in electrodynamics, therefore, the general solutions can be written down as,

$$h_{\mu\nu} = A_{\mu\nu} e^{ik_\alpha x^\alpha}, \quad (1.9)$$

where each component of  $h_{\mu\nu}$  is a sinusoid travelling along vector  $k_\alpha$  with amplitude  $A_{\mu\nu}$  [9]. At this point the set of equations are not simple, the symmetric tensor  $A_{\mu\nu}$  has 10 independent components. This can be greatly simplified by choosing a different gauge (coordinate system) where the metric perturbation is both transverse and traceless (TT) [8]. This is just a choice of coordinate system which does not change any current assumptions. A traceless metric is one where the sum of the diagonal elements are 0 and a transverse metric is when the oscillations are perpendicular to the direction of travel. This gauge imposes two conditions: one is that  $h_{\mu\nu}$  is traceless, i.e. that the sum of the diagonal elements are 0 and the other is that  $h_{\mu\nu}$  is transverse. The transverse element means that the oscillations of the wave happen perpendicular to the direction of travel. At this point we can choose that the wave is travelling in the  $z$  direction which means that  $k = (\omega, 0, 0, k)$ . By then adopting the TT gauge there are only two unique

components to the metric such that the perturbation is defined as,

$$h_{\mu\nu} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & h_+ & h_\times & 0 \\ 0 & h_\times & -h_+ & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} e^{i(kt-wt)}. \quad (1.10)$$

The two unique components are then the two polarisations of gravitational waves,  $h_+$  and  $h_\times$ . The affect of each of the polarisations on a ring of test particles can be seen in Fig. 1.1 where the gravitational wave is travelling out of the page.

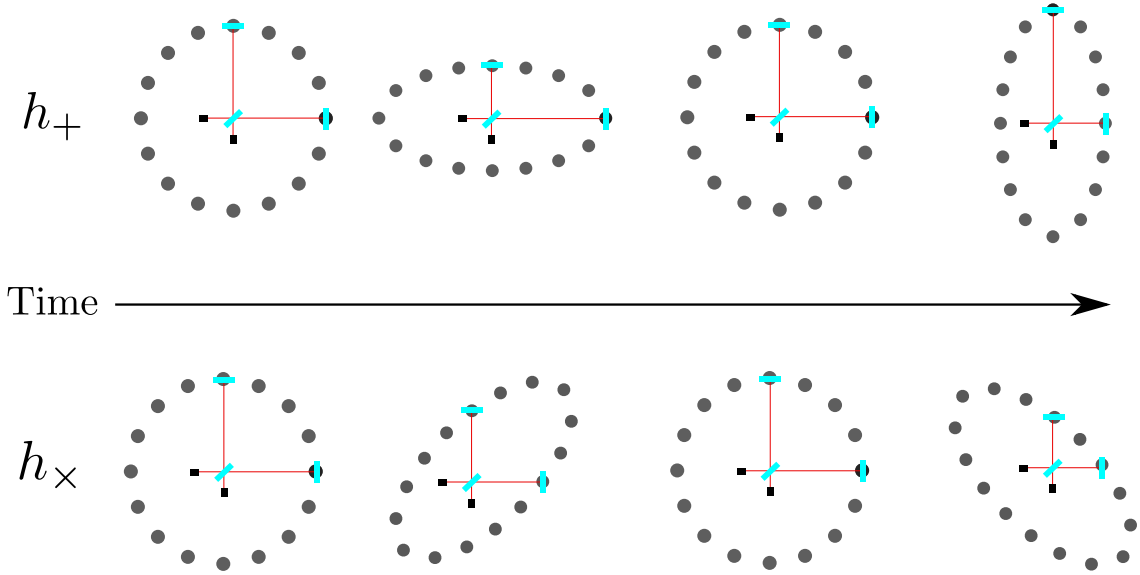


Figure 1.1: Shows how the plus and cross polarisation's affect a ring of test particles. This assumes the wave is travelling out of the page and the effects have been greatly exaggerated. This also shows an example of how this effects the test masses of an interferometer. This will be described in more detail in Sec. 1.3.

### Generating gravitational waves

To generate gravitational waves we go back to Eq. 1.6 where we include the stress-energy term on the right hand side. Following the derivation in [8], one can find that the gravitational wave amplitude is related to the second moment of the mass distribution. The second moment of the mass distribution  $I_{\mu\nu}$  is defined as,

$$I_{\mu\nu}(t) = \int \rho(t, \mathbf{x}) x^\mu x^\nu d^3x, \quad (1.11)$$



where  $\rho$  is the mass density, and  $x_i$  and  $x_j$  are the coordinates [8]. This is the quadrupole moment tensor without the trace subtracted. The gravitational wave amplitude is then defined as,

$$h_{\mu\nu} = \frac{2}{r} \frac{d^2 I_{\mu\nu}(t-r)}{dt^2}. \quad (1.12)$$

This has a slight modification in the TT gauge, see [8], however, has the same relationship between the mass quadrupole and the GW amplitude. This shows that for a GW to be generated, the second derivative of the mass quadrupole moment is needed. A mass quadrupole moment only exists when the mass distribution is not spherically symmetric. Therefore, a mass which is asymmetric and accelerating will produce a GW.

Systems which will produce detectable GWs are generally rapidly rotating high mass systems which have some asymmetry around their rotation axis. The sources of these GW will be described in the following section.

## 1.2 Sources and signals

There are many potential sources for GW. The expected sources can be split into 3 general categories based on their signal type: Transient, Stochastic and CWs. These categories are chosen based on the length of the signal and how well modelled the signal is. Fig. 1.2 shows an example of each of the signals and their category. In the sections that follow, I will give an overview of the potential sources of each of these signal categories and their wave-forms.

### 1.2.1 Transient

Transient sources of gravitational waves give short duration signals which are observable from milliseconds to tens of seconds depending on the source. Some of these sources will emit signals for a much longer time, however these are at a lower frequency and lower amplitude and not observable by current ground based detectors. Transient signals can be further split into two categories based on how well they are modelled. CBCs have well modelled wave-forms and bursts are generally from un-modelled or unknown sources.

#### Compact Binary Coalescence

CBCs originate from the slow in-spiral and merge of two compact objects which are gravitationally bound. The objects in-spiral as they lost energy to the radiation of gravitational waves. Dependent on the masses of the two objects, the gravitational waves generated by the system can be detected by ground based detector such as LIGO [10] and Virgo [11]. In fact, the only detections to date have been originated from this source which include [4, 5, 7].

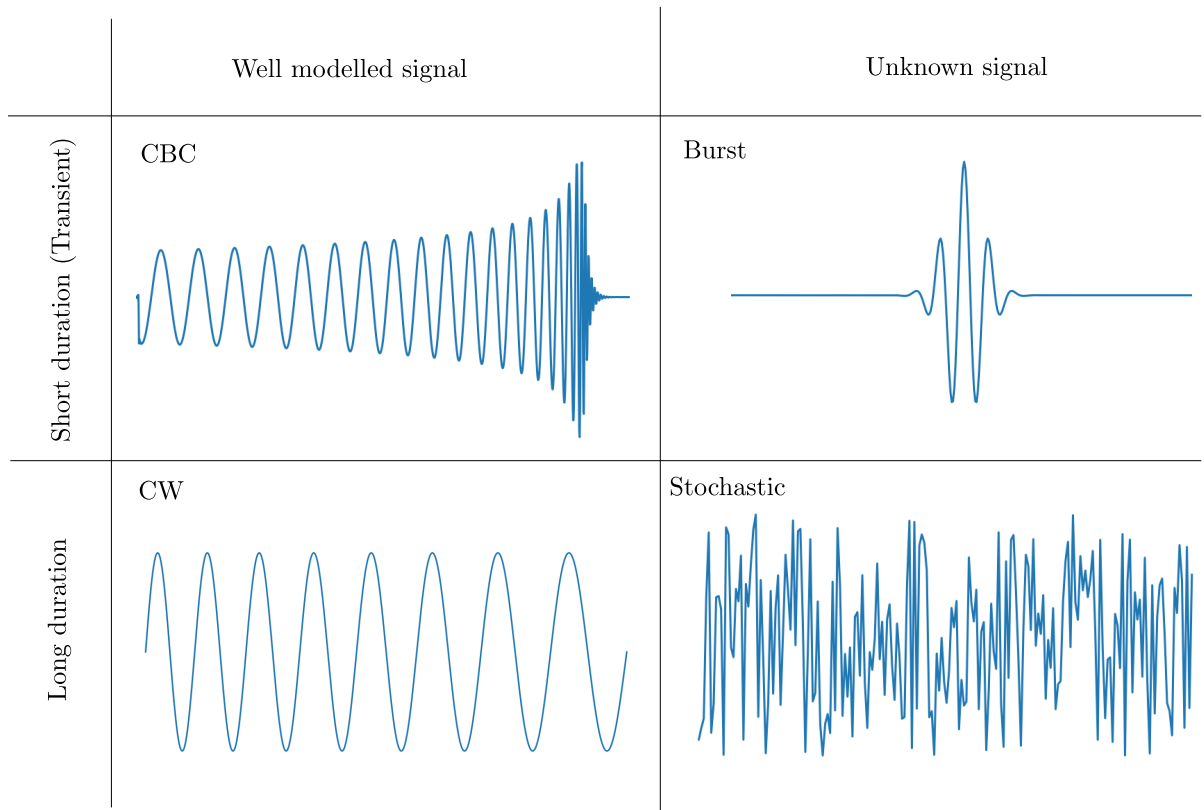


Figure 1.2: Each GW signal type can be categorised based on its signal length and how well the signal is modelled. Transient signals which are short duration, include both well modelled CBC signals and unknown Burst signals. Long duration signals include well modelled CW signals and unknown Stochastic signals.

The compact objects referred to here are either black holes or neutron stars. There are generally 3 types of CBC source: BBH, BNS and neutron star black hole (NSBH). The general structure of the wave-form is the same for each of these and follows a ‘chirp’ where the GW frequency increases with time until merger. An example of this is shown in Fig. 1.2. For higher mass systems such as BBH these signals are detectable by ground based detectors for  $< 1$ s. For lower mass systems such as BNS they can be detected for longer periods  $\mathcal{O}(10)$ s.

The wave-form of a CBC signal is generally split into three separate components: the in-spiral, the merger and ring-down []. **JOE: rewrite this paragraph** The inspiral is when the two compact objects are orbiting each other. As they lose energy to gravitational waves, the radius of the orbit decreases and therefore the frequency increases. The merger is the period when the two objects begin to join to become a single object. The ring-down is the GW emitted of the merged object. The joint object can oscillate whilst it settles into its final state.

In systems which have a neutron star, during the in-spiral when the objects are close, the neutron star can deform due tidal interactions between the objects [12, 13]. This becomes useful as it will affect the generated waveform and can help place limits on and determine the equation of state (EOS) for the dense matter in a neutron star [12, 13]. BNS systems also offer a way to observe objects in multiple different channels, or what is known as multi-messenger astronomy.

This is where the object can be viewed in the electromagnetic (EM) spectrum as well as in gravitational waves. This offers much in the field of astronomy as it can aid in the measurement of the Hubble constant [14]. Observations of BBH systems can also give information on how black holes and BBHs form, more details on this can be found in [15, 16].

## **Burst**

Burst sources are also short duration however, are un-modelled or difficult to model. This means that the exact wave-form of the signal is unknown. There are a few possible reasons for the lack in knowledge of the waveform: the physics of the system is too complicated to model in a reasonable amount of time or there is no model of the source. As there is no model to generate wave-forms, burst searches cannot use matched filtering as in CBC searches. Rather burst searches look for short bursts in power which is coherent between multiple detectors [17, 18]. There are a number of systems which could potentially emit a short duration burst signals. These include core collapse supernovae [19], gamma ray bursts (GRBs) [20], cosmic strings [21] and other unknown sources. Detecting GW from one of these sources could offer more insight into the processes inside hostile environments.

As burst searches are un-modelled, they are sensitive to almost any signal which is coherent between detectors. This allows them to also search for signals from CBC as well as any GW signal from an unknown source.

### **1.2.2 Stochastic**

The stochastic background has no signal model, however, is expected to be a persistent source of GW in the background of the detector. The stochastic background is the incoherent sum of many unresolved GW signals. The source of these signals can be anything from cosmological sources such as cosmic strings to CBC signals. These signals can be thought of as the GW analogue of the cosmic microwave background (CMB). The signal is assumed to be isotropic such that it can be observed at any point on the sky [22]. As the stochastic background is essentially just noise, it is not possible to detect with a single detector [22]. Rather, searches for the stochastic background correlate signals between multiple detectors [23, 22]. When detected, these signals may be able to offer insights into the early universe and its formation.

### **1.2.3 Continuous waves**

CWs are long duration signals which which can be well modelled. The signals last for times greater than the observation runs of ground based detectors and in general have a fixed or slowly varying frequency. There are a number of potential sources of CWs including CBCs long before their merger. Long before a CBC merger, the two compact objects are orbiting each other whilst slowly in-spiralling. This will emit a long duration, almost fixed frequency sinusoid as

its signal. This signal however, is at lower frequency than ground based detectors can detect, therefore space based detectors such as laser interferometer space antenna (LISA) [danzman] are expected to observe this type of CW.

The primary source for many CWs searches is rapidly rotating neutron stars. Neutron stars originate when a massive star collapses and are the remnant of this collapse, they are objects with incredibly high density and are highly magnetised. They have masses around  $1.4 - 2 M_{\odot}$  contained in a star with radius of  $\sim 10$  km with magnetic fields ranging from  $10^8 - 10^{15}$  G [24]. Despite many observations in the electromagnetic spectrum and a large amount of research, these objects are not well understood. A key part of neutron stars which is not understood is the EOS. A review of the current understanding can be found in [89]. The EOS relates quantities such as the pressure, temperature and volume of a neutron star and dictates how the neutron star matter behaves. Observations of GWs from neutron stars can place limits on the EOS of this type of matter. These observations have already been made in the form of BNS mergers [7]. However, independent observations of rapidly rotating neutron stars can add to this understanding.

For a neutron star to emit a gravitational wave it needs to have some asymmetry in its mass distribution around its rotation axis, this follows from Eq. 1.12. There are a number of different mechanisms which could cause this and emit GWs, some of these are reviewed in [25, 26, 27, 28]. Here I will summarise two main theories: Neutron star mountains and neutron star oscillations.

## Mountains

One of more likely mechanisms for detectable GW emission from neutron stars is from ‘mountains’ on the surface of the star. These are permanent deformations on the surface which are non axisymmetric, i.e. the deformation is not symmetric around the rotation axis. Fig. 1.3 shows an exaggerated example of a deformation.

This deformation or asymmetry can be quantified by the ellipticity  $\varepsilon$  of the neutron star. This is defined using the principal moments of inertia,

$$\varepsilon = \frac{I_{xx} - I_{yy}}{I_{zz}}, \quad (1.13)$$

where  $I_{zz}, I_{xx}, I_{yy}$  are the principal moment of inertia. This is when the star is rotating around the  $z$  axis so  $I_{zz}$  is along the rotation axis.

There are a number of theories which describe the origin of this axisymmetry. If the pulsar is in a binary system and accreting material from its companion star, the material can be funnelled towards the magnetic poles by the magnetic field, thereby causing a hot spot [27]. This ‘hot spot’ could cause a deformation on the surface of the star which is not axi-symmetric. The magnetic stresses from strong magnetic fields within the star, could potentially also cause non axi-symmetric deformations to the star. Finally the spin down of the pulsar itself could cause

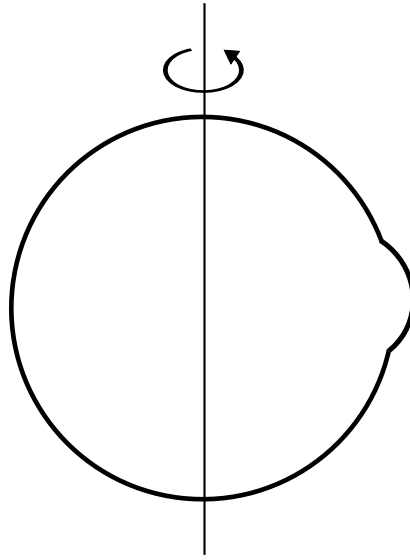


Figure 1.3: The neutron star has some ellipticity which means that the star deforms such that it is not symmetric around its rotation axis. This is thought of as a ‘mountain’ on the surface of the neutron star. The diagram above shows an extremely exaggerated view of what this could look like.

stresses in the crust of the star until the point of breaking, it then after this break which could leave a distortion in the crust [29]. More details on the signal waveform of this type of GW and methods to search for it will be explained in Sec. 2.

### Neutron star oscillations

There are a number of oscillation modes within a star such as f-modes, p-modes and r-modes [29]. Each of these waves are oscillations in the star similar to oscillations in the earth which are used for seismology. The difference between the different modes are the restoring force bringing the perturbed state back to equilibrium. For example, f-modes use gravity as the restoring force where the oscillations happen in the crust of the star. The more promising of these for gravitational wave emission and detection is the r-mode [28]. These are oscillations in the neutron superfluid part of the star, where the restoring force is the Coriolis effect from the rotation of the star. Fig. 1.4 shows an highly exaggerated view of a neutron star with an oscillation mode travelling in each direction. If these modes are excited in a non-rotating star, then they will emit GW where the GW carries away angular momentum [30]. For the mode travelling in a clockwise direction, this angular momentum is positive and the mode travelling anti-clockwise the angular momentum is negative. Therefore, the GW is taking away either positive or negative angular momentum depending on the direction of rotation. The emission of GW damps the modes and the magnitude of the perturbation decreases making them extremely difficult to detect. Now if the neutron star is rotating, this can lead to an effect called the Chandrasekhar, Friedman and

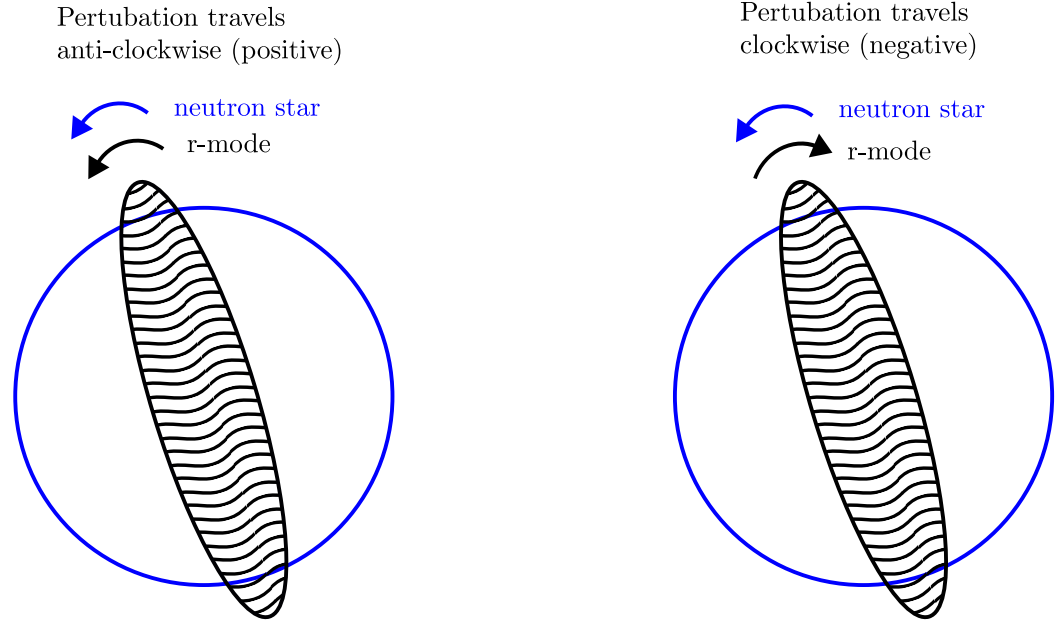


Figure 1.4: The r-modes can travel in either direction in the star. in the case when the r-mode is moving clockwise and the neutron star is moving anticlockwise, if the neutron star is rotating fast enough it can cause unstable emission of GW. This image was reconstructed from [30].

Schutz (CFS) instability [31, 32]. As the rotation speed of the neutron star increases, there are two different effects on the modes travelling in opposite directions. For the mode travelling anti-clockwise with the stars rotation, the mode will appear to be travelling faster, therefore, will emit more GW taking away more angular momentum. This means that this mode will be damped more rapidly. The interesting affect is for the mode travelling clockwise, opposite to the neutron stars rotation. At a certain rotation rate, the mode will be ‘frozen’ from the observers perspective and no GW will be emitted. As the rotation rate increases further, the mode will appear to travel anti-clockwise to an observer, i.e. the mode is dragged in the opposite direction by the stars rotation. Here it is key to remember that this mode had negative angular momentum as in the neutron stars frame it is still travelling clockwise. As the mode rotates its emits positive angular momentum, which is then subtracted from the modes negative angular momentum. The magnitude of the angular momentum then increases such that more GW are released. This effect causes the amplitude of the oscillation to grow and therefore become unstable. Therefore, a neutron star is unstable to GW emission if it is rotating sufficiently fast [28]. For a more detailed view on how r-modes generate GW see [33, 30]

## 1.3 Detectors

The theory mentioned above and the indirect detection of gravitational waves from the Hulse-Taylor binary pulsar system left little doubt as to whether GW existed. The real challenge was to design an instrument which could directly detect gravitational waves. There were a number of different proposed methods for the design of the instrument which includes: resonant bar detectors, both ground based and space based interferometers, pulsar timing arrays and cosmic microwave background detectors. Resonant bar detectors were initially designed and built by Joseph Weber [34]. These are large cylinders of metal which should resonate as a gravitational wave passes by. There are a few different designs of this type of detector, including an omnidirectional design [35]. The majority of these detectors are no longer operational. Pulsar timing arrays aim to use the accurate arrival time of pulses from millisecond pulsars to measure GW [36]. As a GW passes between the pulsar and the observer, the arrival time of the pulses should change. Whilst a detection has not been made using pulsar timing arrays, these methods are still in use. Cosmic microwave background detectors aimed to look for evidence of gravitational waves in the polarisation's of the CMB [37]. These use the a range of detectors to look at the CMB however, are yet to confirm a detection of a GW signal. The most commonly known design of a GW detector is the ground based interferometer, these made the first detection of GW in 2015 [4]. These are the focus of this section as the analysis that will follow uses data from the LIGO detectors in the USA [38, 10] and Virgo detector in Italy [11, 39].

### 1.3.1 Laser Interferometers

Laser interferometers use the interference of light to measure a length with high precision. The majority of this section will focus on ground based interferometers such as LIGO and Virgo [10, 11]. A simple design of an interferometer is shown in Fig. 1.5. Here it shows how the laser is equally split into two, each of these beams is sent down perpendicular arms of the detector. The light then returns to the beam splitter where the two beams are combined and sent to a photo-detector. At the output, there is an interference pattern between the two beams. If the length of one of the arms is changed then the interference pattern will change as the phase of one beam changes with respect to the other. Each of the arms is a Fabry-Perot cavity [], this increases the amount of time light spends within an arm and therefore increases the detectors effective arm length. This helps to increase the sensitivity of the detector **JOE: remove this sentence or go into more detail on extra features.**

This can be used in gravitational wave detection as the mirrors at the end of each arm of the interferometer can be treated as ‘free’ test masses. Fig. 1.5 shows the effect of a GW in free test masses. In the interferometer, this affect essentially changes the relative lengths of the two arms. The change in the interference pattern with time is then related to the GW. Actual ground based GW detectors such as LIGO [38] and Virgo [11] are much more complicated than described

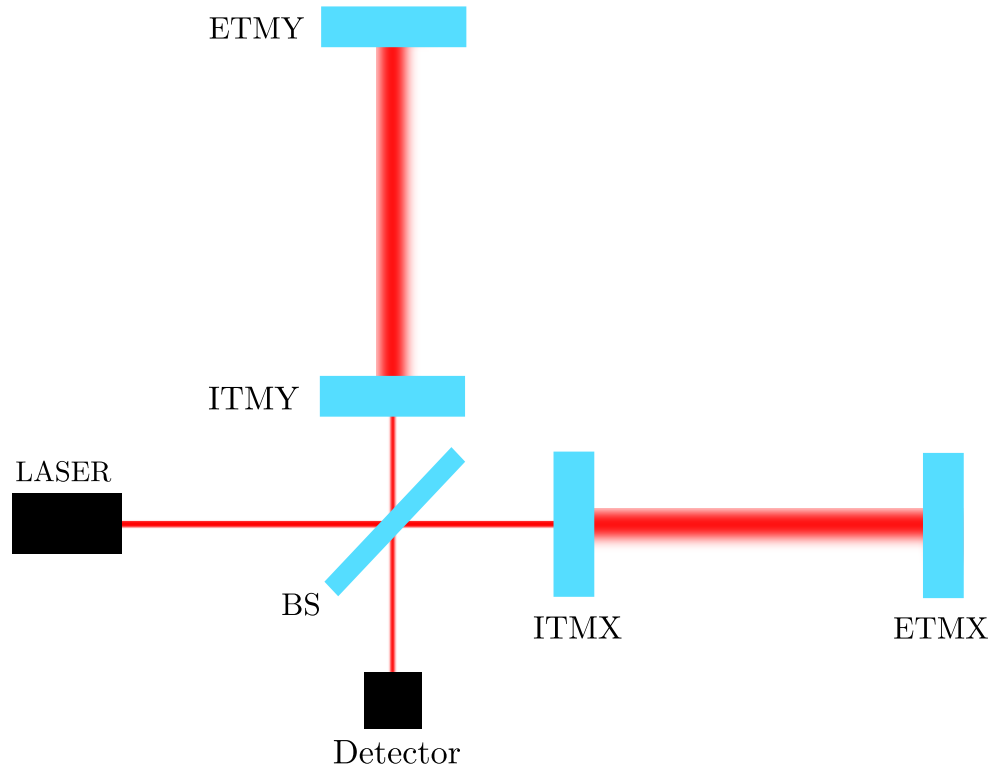


Figure 1.5: This figure shows a basic interferometer. ETMY and ETMX refer to the external test masses, which are just mirrors at the end of the interferometer arms. ITMY and ITMX refer to the internal test masses, these create a cavity in the interferometers arms which can build up laser power. BS is the beam splitter which splits the Laser beam equally to each arm, this then recombines the beams back to the detector.

above. They use many techniques to reduce effects on the detector which are not astrophysical. Some of these effects are listed in Sec. 1.3.1

### Detector response

An important factor to know when using detector data to search for astrophysical signals is the detector's response. This measures how sensitive the detector is to different locations on the sky. An example of the antenna response for LIGO is in Fig. 1.6 where the detector's arms lie on the  $x$  and  $y$  axis of the image. This is clear when thinking about how a gravitational wave affects the test masses. As the GW is transverse to its propagation, when the detector is face on to the source, there will be a maximum change in the arm lengths and therefore a maximum sensitivity. In the same way the sensitivity will be at a minimum when edge to the source.



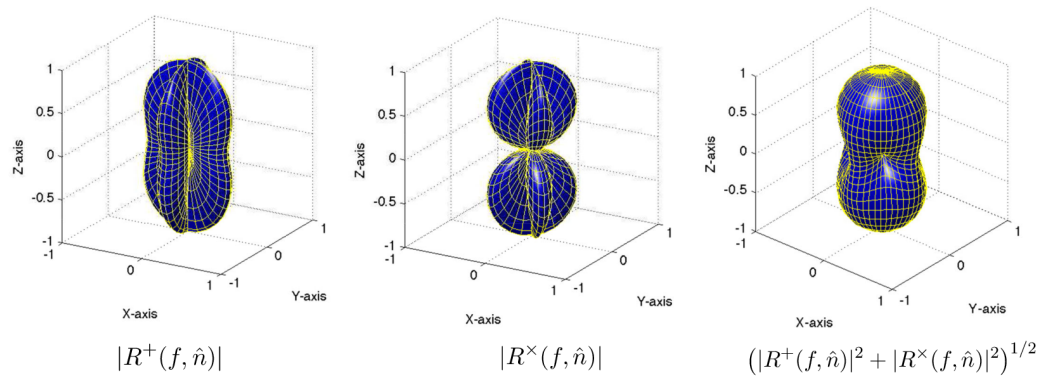


Figure 1.6: The antenna response is shown as in [23] for the plus, cross and average polarisation's. The detectors arms lie on the x and y axis in the above plots.

### Noise sources

To increase the sensitivity of the LIGO detectors, any effect on the output of the interferometer which is not astrophysical needs to be reduced. This involves understanding what causes certain noise features in the detector, and how the affect of these can be reduced. Within the detector, there are many sources of noise. Some of these noise sources and their affect on the detectors strain sensitivity are all shown in Fig. 1.7 from [10]. Here I will summarise some of the limiting sources and also sources which become useful for understanding later sections.

**Seismic noise** This originates from the seismic activity of the earth and effects the lower frequencies. This can be earthquakes or ocean waves. Seismic waves cause the mirrors to oscillate and induce a change in the length of the arm. This is reduced by having multi stage suspensions in the detectors.

**Coating noise** This is in general due to two main factors, the thermal noise of the coating and brownian noise. The Brownian noise is from the mechanical dissipation in the coating and the thermal noise is due to thermal dissipation. The Brownian noise is the dominant factor as shown in Fig. 1.7. This is reduced by using different coatings on the mirrors.

**Quantum noise** Quantum noise is a fundamental limit due to the statistical uncertainty of counting photons. This limits the sensitivity at many frequencies. There are methods to reduce this include squeezing of the light [1].

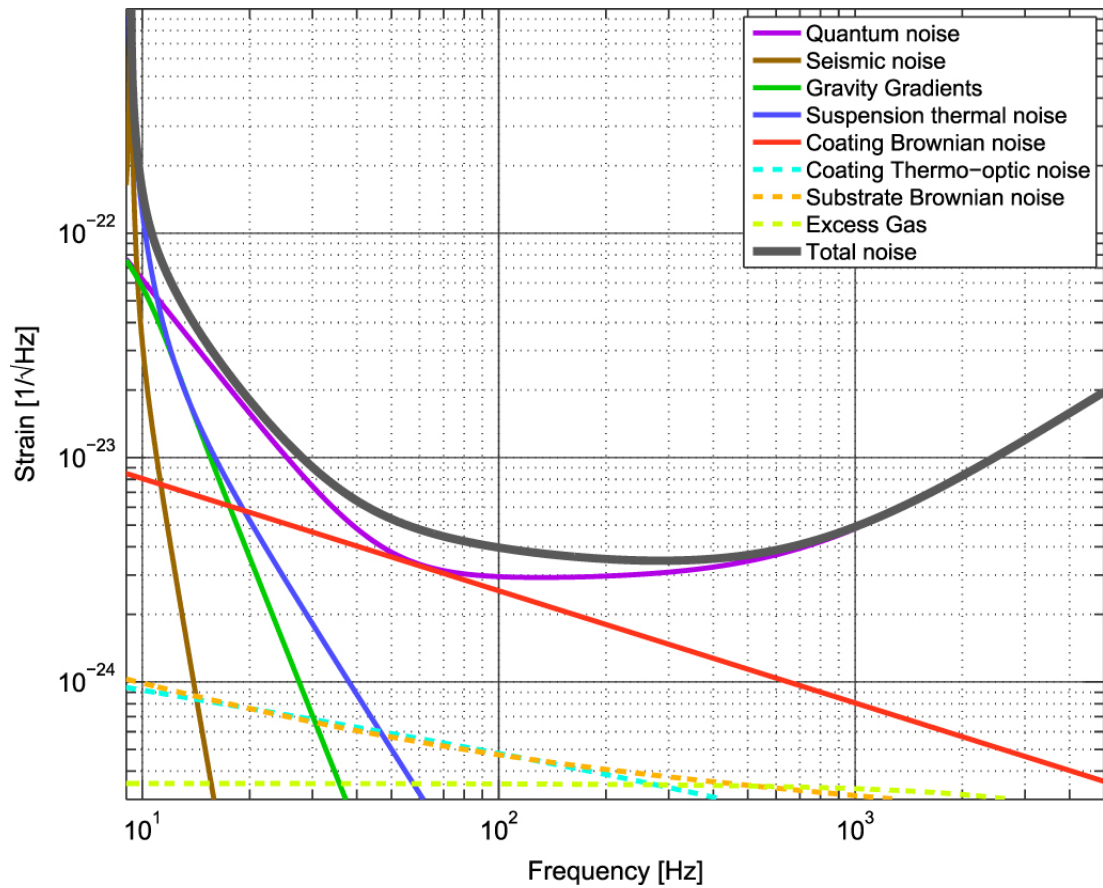


Figure 1.7: The different noise sources affect the sensitivity of the LIGO detectors at different frequencies. This shows the various sources how they affect the noise curve [10].

**Electronics noise** Whilst this is not shown in Fig. 1.7, this becomes important to searches described later. These have a different effect on the detector which is more narrowband frequency lines. This is generated by the digital and analogue electronics that are used to measure the signal.

There are also many other sources of noise in the detector which I have not listed. However, these are often not the limiting cases of noise or are not relevant to this thesis. In Sec. 5 I will go into more detail about specific noise sources in the detector known as instrumental lines and how they can be monitored and potentially removed.

## Chapter 2

# Searching for continuous gravitational waves

Continuous gravitational waves have particular challenges when it comes to their detection. Whilst I have described the potential sources of the signal and its approximate signal type in Sec. 1.2.3, to perform a search for this signal more detail is needed. In this section I will go into more detail on the ‘mountain’ model in Sec. 1.3 and its wave-form description. This model is then used in various search methods for CW signals. In Sec. 2.2 I will overview a subset of current searches for CW signals. Sec. 2.3 explains the motivation for the majority of the work in this thesis.

### 2.1 Continuous signal model

The model of a GW signal from a pulsar is relatively simple, it is a quasi-sinusoidal signal. This means that the signal is a sinusoid with a slowly varying frequency. One reason for the slow variance in the frequency is due to the energy loss to GW as the pulsar spins down. Here the signal is modelled to originate from an isolated triaxial neutron star rotating around a principal axis. The parameters of each pulsar can be split into two sections: the Doppler components  $(\alpha, \delta, f)$  and its amplitude components  $(\psi, \phi_0, \iota, h_0, \theta)$ . This ignores any orbital parameters which would be present if the star was in a binary system and higher order frequency derivatives. They are defined as follows: the sky positions  $\alpha$  and  $\delta$  refer to the right ascension and declination.  $f$  refers to the source frequency and its derivatives.  $\psi$  and  $\phi_0$  and  $h_0$  are the GW polarisation, initial phase and amplitude respectively.  $\iota$  is the inclination angle which is how much the source is tilted relative to the observer.  $\theta$  is the ‘wobble angle’ or the angle between the rotation axis and the symmetry axis of the neutron star.

The definition of the GW from a neutron star here follows that in [26, 45, 46]. The amplitude

of the GW can be defined as,

$$h(t) = F_+(t)h_+(t) + F_\times(t)h_\times(t), \quad (2.1)$$

where  $h_+, h_\times$  are the plus and cross polarisations functions as in Eq. 1.10 and  $F_+, F_\times$  are the antenna pattern functions to the two polarisations. These are defined by,

$$\begin{aligned} h_+(t) &= h_0 \frac{1 + \cos^2(\iota)}{2} \cos(\Phi(t)) \\ h_\times(t) &= h_0 \cos(\iota) \sin(\Phi(t)) \end{aligned} \quad (2.2)$$

The plus and cross polarised components then depend on the GW amplitude  $h_0$ , the inclination angle of the source  $\iota$  and the phase evolution of the GW. Here I have chosen to assume a small wobble angle  $\theta$ , however, this is included in [45]. The phase of the wave  $\Phi(t_{\text{SSB}})$  at the solar system barycenter (SSB) can be defined as,

$$\Phi(t_{\text{SSB}}) = \phi_0 + 2\pi \left[ f(t_{\text{SSB}} - t_0) + \frac{1}{2} \dot{f}(t_{\text{SSB}} - t_0)^2 + \dots \right]. \quad (2.3)$$

This consists of an initial phase  $\phi_0$  which is the phase at time  $t_0$ , the frequency of the signal  $f$  and its derivative  $\dot{f}$  at time  $t_0$ . Here we show the phase to second order, however, this can be easily extended if necessary. The time at the SSB  $t_{\text{SSB}}$  can be transformed to the time  $t$  at the detector by,

$$t_{\text{SSB}} = t - \frac{\mathbf{r}_d \cdot \mathbf{k}}{c} + \delta t. \quad (2.4)$$

Here  $\mathbf{r}_d$  is the position of the detector with reference to the SSB,  $\mathbf{k}$  is a unit vector in the direction of the source. This essentially takes into account the Doppler shift of the signal due to the movement of the detector, i.e. as the earth rotates and orbits the sun.  $c$  is the speed of light and  $\delta t$  is extra corrections from the Einstein, Binary and Shapiro delay []. The amplitudes  $h_0$  in Eq. 2.2 are defined by,

$$h_0 = \frac{16\pi^2 G \epsilon I f^2}{c^4 r}, \quad (2.5)$$

where  $G$  is the gravitational constant,  $c$  is the speed of light,  $\epsilon$  is the ellipticity of the star,  $f$  is the sum of the frequency of rotation of the star and the frequency of precession,  $r$  is the distance to the star and  $I_{zz}$  is the moment of inertia with respect to the rotation axis  $z$ . The ellipticity of the star  $\epsilon$  is a measure of the distortion of the star around its rotation axis and is defined by,

$$\epsilon = \frac{I_{xx} - I_{yy}}{I_{zz}}, \quad (2.6)$$

where  $I_{xx}, I_{yy}$  and  $I_{zz}$  are the moments of inertia for each axis.

In Eq. 2.1,  $F_+(t)$  and  $F_\times(t)$  are the antenna pattern functions of the detector. These describe how sensitive a detector is to a particular location on the sky at any given time. The amplitude

of the signal will vary dependent on the orientation and location of the detector relative to the source. This is described in Sec. 1.3 and the response to sky location is shown in Fig. 1.6. These components are defined in [45] as,

$$\begin{aligned} F_+(t) &= \sin \zeta [a(t) \cos(2\psi) + b(t) \sin(2\psi)], \\ F_\times(t) &= \sin \zeta [b(t) \cos(2\psi) - a(t) \sin(2\psi)], \end{aligned} \quad (2.7)$$

where  $\zeta$  is the angle between the arms of the detectors,  $\psi$  is the polarisation angle of the GW and  $a(t)$  and  $b(t)$  are defined in [45] and relate the sky location to the orientation of the detector at a given time. A full derivation of this can be found in [45] where each of these terms are expanded.

Eqs. 2.1 - 2.7 then describe the amplitude and phase evolution of a signal at a given detector location and orientation.

## 2.2 Bayes Theorem

A key part in understanding the different methods to search for GW or any data analysis, is understanding probability and statistics. This gives understanding of the random processes underlying all measured quantities. Whilst there are generally two approaches to statistics: Frequentist and Bayesian, here I will focus on the Bayesian approach.

### 2.2.1 Basic probability

Initially I will define some basic concepts of probability. We can define the probability of some event  $A$  as  $p(A)$  where probabilities follow  $0 \leq p(A) \leq 1$  and some other event  $B$  which has a probability  $p(B)$  and follows  $0 \leq p(B) \leq 1$ .

**Union** A union is the probability of either and event  $A$  happening or event  $B$  happening. This is written as,  $p(A \cup B)$ .

**Intersection** An intersection is then the probability that both and event  $A$  and an event  $B$  happens. This is written as  $p(A \cap B)$ .

**Independent and dependent Events** If the events  $A$  and  $B$  are independent, i.e. the event  $A$  does not affect the outcome of event  $B$ , then,

$$p(A \cap B) = p(A)p(B). \quad (2.8)$$

However, if the event  $A$  is dependent on event  $B$ , i.e. the event  $A$  affects event  $B$  or vice

versa, then the joint probability of both events is,

$$p(A \cap B) = p(A)p(B | A) = p(B)p(A | B). \quad (2.9)$$

Here  $p(B | A)$  means the probability of event  $B$  happening given that event  $A$  has happened.

**Conditional probability** Conditional probability arises from situations where the outcome of one event will affect the outcome of future events. The definition of this arises from the the dependent events defined above in Eq. 2.9,

$$p(A | B) = \frac{p(A \cap B)}{p(B)}. \quad (2.10)$$

**Bayes Theorem** Bayes theorem can then be defined using conditional probabilities. i.e we can use

$$p(A | B) = \frac{p(A \cap B)}{p(B)} \quad \text{and} \quad p(B | A) = \frac{p(A \cap B)}{p(A)} \quad (2.11)$$

such that then,

$$p(B)p(A | B) = p(A)p(B | A) \quad (2.12)$$

and this is rearranged to Bayes theorem,

$$p(A | B) = \frac{p(A)p(B | A)}{p(B)} \quad (2.13)$$

### 2.2.2 Bayesian Inference

We can take Bayes theorem from Sec. 2.1.1 and apply it to a problem which involves inferring some parameters from some model. Here we can relabel the events  $A$  and  $B$  with the data  $\mathbf{d}$  and the parameters  $\boldsymbol{\theta}$  of some model  $I$ . Eq. 2.1.2 then becomes,

$$p(\boldsymbol{\theta} | \mathbf{d}, I) = \frac{p(\boldsymbol{\theta}, I)p(\mathbf{d} | \boldsymbol{\theta}, I)}{p(\mathbf{d} | I)} \quad (2.14)$$

where each of the components are assigned names:  $p(\boldsymbol{\theta} | \mathbf{d})$  is the posterior distribution,  $p(\boldsymbol{\theta})$  is the prior distribution,  $p(\mathbf{d} | \boldsymbol{\theta})$  is the likelihood and  $p(\mathbf{d})$  is the evidence.

**Posterior** The posterior distribution describes the probability of a parameter  $\boldsymbol{\theta}$  in some model  $I$  given some data  $\mathbf{d}$ . For many problems this is the distribution which is most useful as it informs you the most likely set of parameters of your model given some observation.

**Prior** The Prior distribution is a key part of Bayesian statistics. This distribution describes any information which you have prior to the observation. This is a distribution defined by the user, where you define a distribution of the parameters based on what you expect to be true.

**Likelihood** The likelihood is where the observation is included in the calculation. This tells you how probable it is to get the observed data  $d$  given the model  $I$  with the set of parameters  $\theta$ .

**Evidence** The evidence is the probability of the data itself given the choice of model. This is found by integrating the likelihood over all possible values of  $\theta$  weighting them by our prior belief of that value of  $\theta$ . This known as a marginal distribution and is defined by,

$$p(d | I) = \int p(\theta, I) p(d | \theta, I) d\theta. \quad (2.15)$$

Bayes theorem then gives a description of the probability distribution of some parameters in a model given some observation. Often when using Bayesian statistics the aim is to find posterior distribution of parameters. There are very few cases where this can be calculated analytically, therefore, numerical methods are almost always used to find the posterior. This can be difficult to calculate numerically especially in problems where the parameters space has many dimensions. The most difficult part to calculate is the evidence in Eq. 2.15, this involves calculating an integral over all possible parameters. There is however, a way around having to calculate this. For any given mode  $I$ , the evidence  $p(d | I)$  is the same for any set of parameters  $\theta$  in Eq. 2.14. The evidence is then just a normalisation factor for the posterior distribution. When different models are not being compared, and we assume the model  $I$  to be true, we no longer need to calculate the evidence. The posterior distribution can then be found by sampling,

$$p(\theta | d, I) \propto p(\theta, I) p(d | \theta, I). \quad (2.16)$$

To find this posterior you could then calculate (sample) the value for every point in parameter space. This however, is very computationally expensive and often the posterior distribution is located in a small area in parameter space. Therefore, the majority of the time is sampling a area of parameter space where the posterior is close to zero, and this is not particularly useful. A method titled Markov-Chain Monte Carlo (MCMC) was proposed [40] to deal with this issue, more information on this can be found in [41, 42]. This builds up the posterior distribution by randomly jumping around in the parameter space. It starts by calculating the posterior value for a particular point in parameter space. Then it will randomly jump to another parameter space point. The posterior can then be calculated again, if the posterior value is higher then the jump is ‘accepted’. This just means that the parameter values of this point are stored. If the posterior value is lower than the previous step then the jump is accepted with some probability. This means that there is a random chance that a value lower than the current is accepted. As the accepted positions aim for areas where the posterior is higher, MCMC does not waste time calculating areas in parameter space of low posterior values. The samples which were accepted then build the posterior distribution.

In certain situations it can be useful to calculate the evidence in Eq. 2.15. For example, if there are two different models which could represent the data, the evidence can be used to determine which of the two models is more likely. This is known as a Bayes factor where two models  $I_1$  and  $I_2$  are compared and is defined as,

$$B = \frac{p(\mathbf{d} | I_1)}{p(\mathbf{d} | I_2)} \quad (2.17)$$

This then requires the calculation of the evidence. To estimate the evidence efficiently a method known as Nested sampling can be used, this is explained in detail in [43, 44]. By calculating the Bayes factor, which is similar to a likelihood ratio, one can find the posterior odds of a particular model by using,

$$\frac{p(I_1 | \mathbf{d})}{p(I_2 | \mathbf{d})} = \frac{p(I_1)}{p(I_2)} \frac{p(\mathbf{d} | I_1)}{p(\mathbf{d} | I_2)}. \quad (2.18)$$

This can be written as *posterior odds* = *prior odds*  $\times$  *bayes factor*. This is then a comparison of how likely different models are given some observation.

The methods described above then provide a way to estimate parameters of a model given some data. Also this provides a way to compare different models given some observation. In following sections the methods described above are used to estimate various parameters.

## 2.3 Continuous wave searches

There are many different methods to search for continuous gravitational waves. They can be split into three general categories: Targeted searches, Directed searches and All-sky searches. The main difference between these different categories is the amount which is known about the source prior to the search. For targeted searches the sky position  $(\alpha, \delta)$  and rotation frequency are known from electromagnetic observations, i.e. X-ray, radio or  $\gamma$ -ray. Directed searches have information on the sky position  $(\alpha, \delta)$  but not the rotation frequency. For all-sky searches, there is no prior knowledge of the pulsar, therefore, is a search for unknown pulsars. In general the searches in each of these categories use two distinct techniques: Fully coherent searches and Semi-coherent searches

### 2.3.1 Fully coherent

A Fully coherent search generally uses a pre generated waveform which follows the model described in Sec. 2.1. This contains all the phase information of the signal. The set of parameters which generated the waveform which ‘matches’ the best can then be considered as the optimum set of parameters given the data. This is known as a matched filter [1] and is used in CW searches in [46].

The matched filter maximises the signal to noise ratio for a given filter, in this case the filter



is our CW model. The matched filter used for CW models is defined in [47] and it titled the  $\mathcal{F}$ -statistic. This essentially maximises a likelihood with respect to the parameters. If one assumes that the noise  $n$  is Gaussian and zero mean, the data  $x$  can be written as,

$$x(t) = n(t) + h(t). \quad (2.19)$$

The likelihood can then be written as,

$$\log \Lambda = (x | h) - \frac{1}{2} (h | h) \quad (2.20)$$

where the product  $(x | y)$  is defined as,

$$(x | y) = 4\mathcal{R} \int_{-\infty}^{\infty} \frac{\tilde{x}^X(f) \tilde{y}^{X*}(f)}{S^X(f)} df. \quad (2.21)$$

This is fully expanded into the  $\mathcal{F}$ -statistic in [45], however, it is essentially this likelihood function which is maximised.

Targeted searches look for a specific pulsar which has been observed in the electromagnetic spectrum. These observations give information such as the sky position and the frequency evolution of the source. Using knowledge of the earth's position around the sun, which is well known, one can use the accurate sky position and frequency of a known source to find its phase evolution in Eq. 2.3. This means that for this type of search one can maximise the likelihood with respect to the parameters  $h_0, \phi_0, \iota$  and  $\psi$ . Another method which uses templates is described in [46], this uses a Bayesian approach.

This type of search can take long periods of time. This is due both to the size of the parameter space and the amount of data which needs to be searched. CW searches need long observation times in order to accumulate the required SNR for detection. Therefore, most searches use data from an entire LIGO observing run which can last for  $\mathcal{O}(1)$  years. Given that the sampling rate for the GW channel is 16 kHz usually downsampled to  $\sim 4$  kHz, the quantity of data is large.

Whilst the fully coherent matched filter searches have methods to reduce the computational time for known sources, in all-sky and directed searches, this type of search is not feasible. This is because all-sky and directed searches have a wider parameter space, therefore, enough templates need to be made to sufficiently cover the large parameter space. This task quickly becomes impossible for coherent matched filtering for an entire observing run due to the amount of time needed. This problem led to the development of semi-coherent searches which will be introduced in the next section.

### 2.3.2 Semi coherent

Semi-coherent searches offered a solution to searching over large parameters spaces and large amounts of data. As is directed and all-sky searches the phase evolution of the source is not known, one cannot use a coherent search for the entire observing run. It may however, be possible to approximately describe the phase for a shorter length of time known as the coherence time,  $T_{coh}$ . The general idea of a semi-coherent search is to break the data-set into smaller section which each can be analysed coherently. The coherent analysis can use the matched filter as described in Sec. 2.2.1 or another method such as a Fourier transform. The results from each of these individual sections can the be combined incoherently using various methods which will be summarised later. This method can greatly reduce the time taken for the analysis depending on the coherence length, however, will always come with some loss in sensitivity.

There are many different types of semi-coherent search which use various methods to incoherently combine the coherently analysed results. I will summarise some of these searches below, some of these searches were summarised and compared in [48]. Many of these searches use a set of 1800s long Fourier transforms as the input data, known as SFTs. This is a default for many all-sky CW searches, where it assumes that the signal remains within one frequency bin during that 1800s.

**Stack-slide** Stack uses a set of Fourier transforms of the data known as SFTs, specifically it uses the power spectrum of these. Each of the separate Fourier transforms (segments) is shifted up or down relative to the others to account for the Doppler modulation of the source. The power from each can then be stacked. More explanation of this can be found in [49, 50]

**Hough** The Hough transform is based on the stack-slide algorithm. The main difference is that the detection statistic for each segment is assigned a weight of 0 or 1 depending if it crossed a detection threshold. The Hough transform can the create a ‘Hough map’ which gives a view of the data in parameter space. This approach is explained in greater detail in [51, 52]. This method has been applied in two main ways known as Sky Hough [51] and Frequency Hough [52, 53].

**Einstein@Home** Einstein at home uses the  $\mathcal{F}$ -statistic mentioned above in various stages. It has a hierarchical structure where it starts with a coarse parameter space with shorter coherence times. This search then provides a list of candidates from this run in coarse parameter space. The parameter space is then more finely sampled around the parameters of the candidates and this process is repeated. The search can also increases the coherence length when searching around given candidates to improve the sensitivity of the search. This algorithm has many additions which are explained in more detail in [54, 55, 48] This provides the most sensitive all-sky CW search, however, uses a large amount of computing

Table 2.1: From [48], shows the computational cost for the first 4 months of advanced LIGO for each search. One MSU, where one standard unit is one core-hour on a standard core. ‘Expected computational costs of searches using the first four months of advanced LIGO data with each search pipeline. These estimates are for a different data observing time from that of the MDC, and do not cover the same parameter space as each other or the MDC. The Einstein@Home searches uses the computing resources of the Einstein@Home project and is designed to run for 6 - 10 months in the Einstein@Home grid.’

Pipeline	Expected runtime of O1 search
Powerflux	6.8 MSU
Time domain $\mathcal{F}$ -statistic	1.6 MSU
Frequency Hough	0.9 MSU
Sky Hough	0.9 MSU
Einstein@Home	100-170 MSU

power. This is achieved by using a distributed computing projects, more details can be found at [56].

**Time domain  $\mathcal{F}$ -statistic** The time domain  $\mathcal{F}$ -statistic splits the data into narrowband segments of length  $\sim 2$  days [48]. Then a coherent search using the  $\mathcal{F}$ -statistic is applied to each of these segments. Values of this statistic above a threshold are stored. Coincidences are then found in each segment, where candidates are selected best on a given threshold. This is explained in greater detail in [57, 48].

**Powerflux** Powerflux uses a standard set of 1800s SFTs. For each point in parameter space, the power in this set of SFTs along the frequency track is recorded. This power is then weighted depending on the antenna pattern and noise of the detector. In longer stretches of  $\sim 1$  month, the weighted power is summed. Any point in parameter space which produces high power in each of these stretches is identified as a potential signal. This search can then be repeated around each candidate with a finer resolution in parameter space. This is explained in more detail and tested in [58, 48, 59]

**Viterbi** **JOE: should probably mention other Viterbi searches**

Each of these searches uses a large computational cost. In [48] a MDC was conducted to compare the sensitivity of some of the searches, where an expected runtime for and O1 search was presented. Result in O1 for some of these searches can now be found in [60]. The results of this are shown in Tab.2.1. Even the fastest of these searches takes close to 1 million core-hours to search through four months of data. Running computing clusters can be costly, therefore, the focus of much of this thesis is to develop techniques to reduce the computational cost of an all-sky search for CW.

## 2.4 Motivation

The searches described above are computationally intensive, where the fastest takes  $\sim 1$  million core-hours. The motivation for much of the work that follows is to develop search algorithms which can reduce this computational time without a large loss in sensitivity to CW signals. **JOE:**  
**more**

## Chapter 3

# SOAP: A generalised application of the Viterbi algorithm to search for continuous gravitational wave sources.

The SOAP search is a semi-coherent CW search algorithm that aims to reduce the computational time needed to find a potential signal. The algorithm looks through narrow-banded time-frequency spectrograms of data to find the ‘most probable track’ in frequency through it. This ‘most probable track’ is then the most likely track which a pulsars frequency would follow. The motivation of the search is simple, if we looked at a frequency band in a spectrogram as in Fig. 3.1, we could find every possible randomly wandering track from a starting frequency bin to an end frequency bin. For each of these tracks the sum of the spectrogram power along the track can be found such that for each track there is a single value. Fig. 3.1 shows a histogram of a subset of these values. If there is a signal present in this band then it could be assumed that the signal which gives the highest sum of spectrogram power is the track which is most likely to follow the signals frequency track. In Fig. 3.1 the optimum track in red shows gives a statistic value of  $\sim 1700$  which is far outside the main distribution of summed powers. The red track follows that an injected signal. This demonstrates that the sum of the spectrogram power along a track which follows a signal is outside the distribution of tracks randomly walk through noise. Therefore, it can be assumed that if the frequency track with the highest sum of spectrogram power is found, then the corresponding track is most likely for follow a signal. Now as calculating all possible tracks is computationally expensive. Given that in the example in Fig. 3.1, the spectrogram is frequency 180 bins wide  $M$  and 400 time segments long  $N$  and after each segment the track has  $T$  possible options to jump to (in this case it is 180 options), The total number of possible tracks is  $T^{M \times N}$ . For this spectrogram this value is  $180^{72000}$ , this is an unreasonable number of tracks to possibly calculate. This is where the Viterbi algorithm [61] is useful as it can efficiently find the track which gives the maximum sum of power. For an equivalent search the Viterbi algorithm would have to do  $T * M * N$  calculations for find the

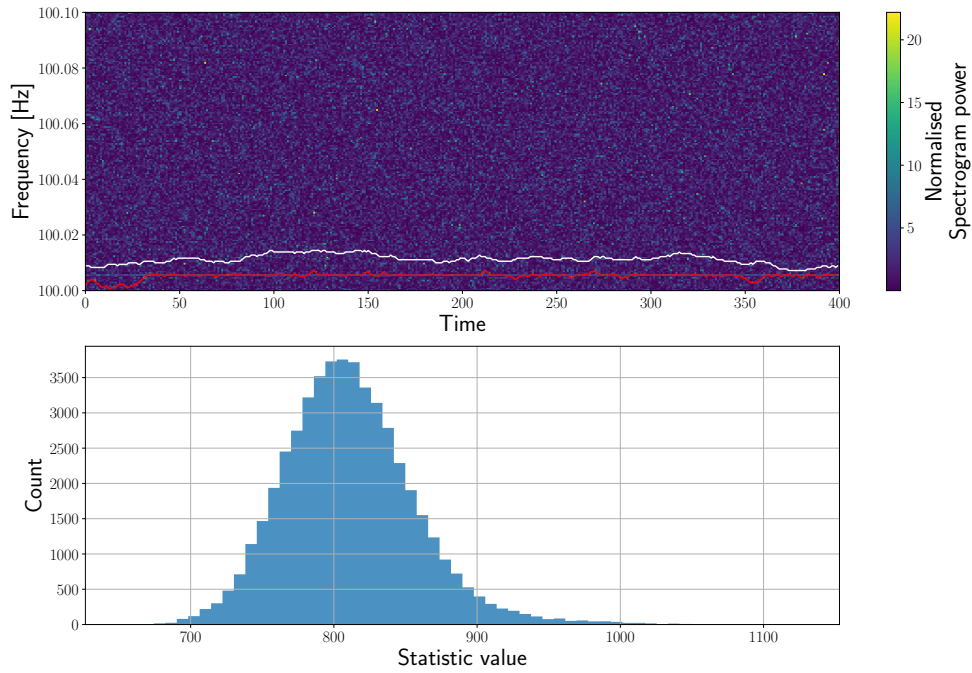


Figure 3.1: This figure shows an example of a time-frequency spectrogram which is typical of LIGO data which is searched through. Here an instrumental line has been injected at 100.006 Hz. The white track shows a random walk track through this spectrogram whereas the red line shows the track which gives the highest sum of detector power. The second panel shows a histogram of a subset of all paths which can be found through the given spectrogram from start to finish. This is a subset as the total number of paths is too large to calculate. The value of the statistic which comes from the optimal path is  $\sim 1700$ . This is much larger than any of the random tracks in our subset and much larger than the mean of all tracks.

optimum track. The method of how this is done is explained in the following sections.

The majority of this chapter that follows has been reviewed and published as in [77]. This work was done by the author under the supervision of Prof. Graham Woan and Dr. Chris Messenger. The work in Sec. 3.9, Sec. 3.11, Sec. 3.11.3 and Sec. 3.11.4 which is supplementary material.

### 3.1 Introduction

One of the main targets for current ground based GW detectors, including LIGO [38, 10] and Virgo [39, 11], are sources of continuous gravitational waves. These are long-duration, quasi-monochromatic sinusoidal signals that are well-modelled by a Taylor series expansion in the signal phase. A likely source of such signals are rapidly spinning non axisymmetric neutron stars. A number of possible emission mechanisms are outlined in [62, 63].

These types of GWs are expected to give strain amplitudes that are significantly below the detector’s noise spectral density, and need sensitive search algorithms for detection. The most sensitive method is to use a coherent matched filter which requires knowledge of the waveform beforehand such that it can be coherently correlated with the data. This approach is used in searches for gravitational signals from known pulsars such as [46, 64, 45, 90, 65]. For broad parameter space searches, where the parameters of the signal are unknown, a large number of template waveforms must be used to sufficiently cover the parameter space. This approach rapidly becomes computationally impractical as the search space grows, so semi-coherent search methods have been developed to deliver the maximum overall sensitivity for a given computational cost. Semi-coherent searches break the data up into sections of either time or frequency and perform a coherent analysis on these sections separately. These intermediate results can then be recombined incoherently in a number of different ways to form the final search result outlined in [66, 67] and references therein.

The analysis that we present here is known as SOAP [68] and is based on the Viterbi algorithm [61]. The algorithm models a process that has a discrete number of states at discrete time steps, and computes the set of states which gives the highest probability (suitably defined) given the data. Our implementation of SOAP is intended as a stand-alone search which is naturally non-parametric and has broad applications to both searches for known signal types and signals which have an unknown frequency evolution. The algorithm works in time-frequency plane, where our ‘states’ are represented by the time and frequency coordinates of a potential signal. We can then find the most probable set of frequencies a possible signal could have, i.e., we can find the most probable track in frequency as a function of time. This is not the first application of the Viterbi algorithm to GW data. Another variant of the algorithm [91] has recently been used, amongst other applications, as part of a CW search to track a pulsar with randomly wandering spin frequency [92, 69, 70, 71, 72]. We develop an alternative version which is aimed to be applied more generally to search for any long duration signals using just SFTs.

In the next section we will describe the Viterbi algorithm and the basic SOAP implementation to GW time-frequency data. We then describe additional features to the algorithm, including the use of data from multiple detectors. As well as this we describe methods used to ignore instrumental effects in the data, such as incoherently summing data and a ‘line aware’ statistic. In the final section as well as a test of the computational cost of the search, we show results of a search performed on datasets of increasing complexity: Gaussian noise with no gaps (i.e., contiguous in time), Gaussian noise with gaps simulating real data more accurately, and finally real LIGO data taken during the sixth science run.

### 3.2 Viterbi algorithm

The Viterbi algorithm is an efficient method for determining the most probable set of states (a single ‘track’ of steps on the time-frequency plane) in a Markov model dependent on data, where the model has a discrete number of states at each step. Rather than computing the probability of every possible track and selecting the most probable, the algorithm maximises this probability after every discrete step. As a result, a partial track which cannot ultimately be the most probable is rejected before the next step is calculated, and only a fraction of all possible tracks need to be computed to find the one that is most probable.

In this work we apply the Viterbi algorithm to a GW strain time-series to find the most probable track of a single variable-frequency signal in the noisy data. We divide the time series into  $N$  equal-length and contiguous segments  $\mathbf{x}_j$ , defining the set  $D \equiv \{\mathbf{x}_j\}$ . The ‘states’ in the model correspond to the frequencies a signal could have in each segment. A ‘track’ is a list of such frequencies  $\mathbf{v} \equiv \{v_j\}$ , where  $v_j$  is the frequency in the segment  $\mathbf{x}_j$ .

Our objective is to calculate the most probable track given the data, i.e., the track that maximises  $p(\mathbf{v} | D)$ . Using Bayes theorem, this posterior probability can be written as

$$p(\mathbf{v} | D) = \frac{p(\mathbf{v})p(D | \mathbf{v})}{p(D)}, \quad (3.1)$$

where  $p(\mathbf{v})$  is the prior probability of the track,  $p(D | \mathbf{v})$  is the likelihood of the track (i.e., the probability of the data given the track) and  $p(D)$  is the model evidence (or marginalised likelihood).

The Viterbi algorithm treats the track as the result of a Markovian process, such that the current state depends only on the previous state. It is therefore useful to split the track’s prior into a set of transition probabilities such that

$$\begin{aligned} p(\mathbf{v}) &= p(v_{N-1}, \dots, v_1, v_0) \\ &= p(v_{N-1} | v_{N-2})p(v_{N-2} | v_{N-3}) \dots p(v_1 | v_0)p(v_0) \\ &= p(v_0) \prod_{j=1}^{N-1} p(v_j | v_{j-1}), \end{aligned} \quad (3.2)$$

where  $p(v_0)$  is the prior probability that the signal in the first time step has a frequency  $v_0$  and  $p(v_j | v_{j-1})$  is the prior ‘transition’ probability for  $v_j$  given the frequency at the last step was  $v_{j-1}$ .

The noise in each of the segments can be treated as independent, so the likelihood component in Eq. 3.1 can be factorised as

$$p(D | \mathbf{v}) = \prod_{j=0}^{N-1} p(\mathbf{x}_j | v_j), \quad (3.3)$$



where  $p(\mathbf{x}_j | \mathbf{v}_j)$  is the likelihood of our signal having a frequency  $\mathbf{v}_j$  in the  $j$ th segment.

Using Eqs. 3.1, 3.2 and 3.3, the posterior probability is then

$$p(\mathbf{v}|D) = \frac{p(\mathbf{v}_0)p(\mathbf{x}_0|\mathbf{v}_0) \prod_{j=1}^{N-1} p(\mathbf{v}_j|\mathbf{v}_{j-1})p(\mathbf{x}_j|\mathbf{v}_j)}{\sum_S \left\{ p(\mathbf{v}_0)p(\mathbf{x}_0|\mathbf{v}_0) \prod_{j=1}^{N-1} p(\mathbf{v}_j|\mathbf{v}_{j-1})p(\mathbf{x}_j|\mathbf{v}_j) \right\}}, \quad (3.4)$$

where in the denominator we must sum over all possible tracks  $S$ . We require the specific track, or set of frequencies,  $\hat{\mathbf{v}}$  that maximises the posterior probability. Therefore, as the denominator in Eq. 3.4 is a sum over all possible tracks, the track which maximises the posterior is the same track which maximises the numerator on the right-hand side of Eq. 3.4, i.e.,

$$p(\hat{\mathbf{v}}|D) \propto \max_{\mathbf{v}} \left[ p(\mathbf{v}_0)p(\mathbf{x}_0|\mathbf{v}_0) \prod_{j=1}^{N-1} p(\mathbf{v}_j|\mathbf{v}_{j-1})p(\mathbf{x}_j|\mathbf{v}_j) \right]. \quad (3.5)$$

This track also maximises the log of the probability and can be written as,

$$\begin{aligned} \log p(\hat{\mathbf{v}}|D) = \max_{\mathbf{v}} & \left\{ \log p(\mathbf{v}_0) + \log p(\mathbf{x}_0|\mathbf{v}_0) \right. \\ & \left. \sum_{j=1}^{N-1} \left[ \log p(\mathbf{v}_j|\mathbf{v}_{j-1}) + \log p(\mathbf{x}_j|\mathbf{v}_j) \right] \right\} + \text{const.} \end{aligned} \quad (3.6)$$

The Viterbi algorithm finds the most probable track  $\hat{\mathbf{v}}$  by calculating the quantities in Eq. 3.6 for each frequency at each time step. In the following sections we explain how this is achieved in practice.

### 3.3 The transition matrix

An important concept when using the Viterbi algorithm is the ‘transition matrix’  $T$ , which is defined as the matrix that stores the prior log-probabilities  $\log p(\mathbf{v}_j | \mathbf{v}_{j-1})$ . These transition probabilities depend only on the size and direction of the transition, and in our case correspond to a jump in frequency when moving from the  $(j-1)$ th to the  $j$ th state. It is within the transition matrix that we impose some loose model constraints. For example it is usual in the time-frequency plane for frequencies to only have discrete values (frequency bins) and a track might only be allowed to move by one bin in each time step, restricting it to a UCD transition or ‘jump’ or equivalently setting the size of the first dimension of the transition matrix  $n_1 = 3$ . We can also impose that the transition probabilities are independent of the current track location in frequency, i.e.  $p(\mathbf{v}_j | \mathbf{v}_{j-1}) = p(\mathbf{v}_{j+k} | \mathbf{v}_{j+k-1})$ . This leads to the transition matrix containing only three numbers, corresponding to the three prior log-probabilities that the track was in the

corresponding UCD frequency bin at the previous time step. These numbers are chosen to reflect the prior probability of a frequency deviation in the track and depend on the class of signals that one wishes to detect. For the majority of examples that follow, a symmetric transition matrix is used, i.e. the probability of a transition up a frequency bin is equal to the probability of a transition down a frequency bin. This allows us to parameterise the one dimensional transition matrix with a single value, this value is the ratio of the probability of a transition to the same frequency bin, to either up or down a frequency bin.

In later sections we will consider more complex situations in which the transition matrix describes the prior probability associated with sequences of even earlier transitions (‘memory’) and the case where there are multiple detectors. In these cases the number of dimensions of the transition matrix can grow substantially to account for the extra complexity of the problem.

### 3.4 Single detector

We will first consider the simple case of a single dataset  $D$ , generated by a single gravitational wave detector, and consider only a one-dimensional transition matrix. We will make use of discrete Fourier transforms so that frequencies, and hence the track frequencies, are also discrete. These frequencies will be indexed by  $k$  and therefore  $\nu_j \rightarrow \nu_{j,k} = k(j)\Delta f$  where  $\Delta f = 1/T$  is the frequency bin width for a segment of duration  $T$ .

The Viterbi algorithm determines the most probable track on the time-frequency plane by calculating the value of Eq. 3.6 for every discrete Fourier frequency, incrementally in time. In other words, at each time segment it finds the most probable earlier track which ends at each particular frequency. On reaching the final segment it can look back to identify the most probable track connecting segment 1 to segment  $N$ .

There are two main components to Eq. 3.6: the transition probabilities  $p(\nu_j | \nu_{j-1})$  and the likelihoods  $p(\mathbf{x}_j | \nu_j)$ . The transition probabilities are pre-calculated and stored in a transition matrix according to Sec. 3.3 above. To calculate the likelihood we follow the approach of [93] which gives, under the assumption of a single sinusoidal signal in additive Gaussian noise in data segment  $\mathbf{x}_j$ ,

$$p(\mathbf{x}_j | \nu_{j,k}, \sigma_{j,k}, I) \propto \exp [C(\nu_{j,k})]. \quad (3.7)$$

where  $C_{j,k}(\nu_{j,k})$  is the Schuster periodogram normalised to the noise variance at frequency  $\nu_{j,k}$  of segment  $j$ . This is equivalent to the log-likelihood, and is defined as

$$C(\nu_{j,k}) \equiv C_{j,k} = \frac{1}{\sigma_{j,k}^2} \frac{1}{N_s} \left| \sum_{r=0}^{N_s-1} x_{j,r} e^{i\nu_{j,k} t_r} \right|^2, \quad (3.8)$$

where  $N_s$  is the number of data points in each segment and  $t_r$  is the time corresponding to  $x_{j,r}$ ,

the  $r$ th sample in the  $j$ th data segment.  $\sigma_{j,k}^2$  is the noise variance and is calculated as an estimate of the noise power spectral density (PSD) in the  $k$ th sample and the  $j$ th data segment. It is worth noting at this point that it is also possible to write this as a likelihood ratio, and therefore write out detection statistic as a log-odds ratio, however, we will discuss this in more depth in Sec. 3.8. The log-likelihoods of each segment can be calculated at discrete frequencies before running the algorithm by computing the power spectra for each segment from discrete Fourier transforms of the data. In the GW field these standard data forms are known as SFTs.

The Viterbi algorithm records two quantities for each frequency and time bin: The first,  $V_{j,k}$ , contains the value defined by Eq. 3.6, which is the log-probability of the most probable path ending in position  $j,k$ . The second,  $A_{j,k}$ , is the transition, or ‘jump’, used to achieve the most probable path. The algorithm can be divided into three main sections: initialisation, iteration and identification. These three sections are described in pseudo-code in Alg. 3.1 and a simple demonstration of the algorithm at work is shown in Fig. 3.2.

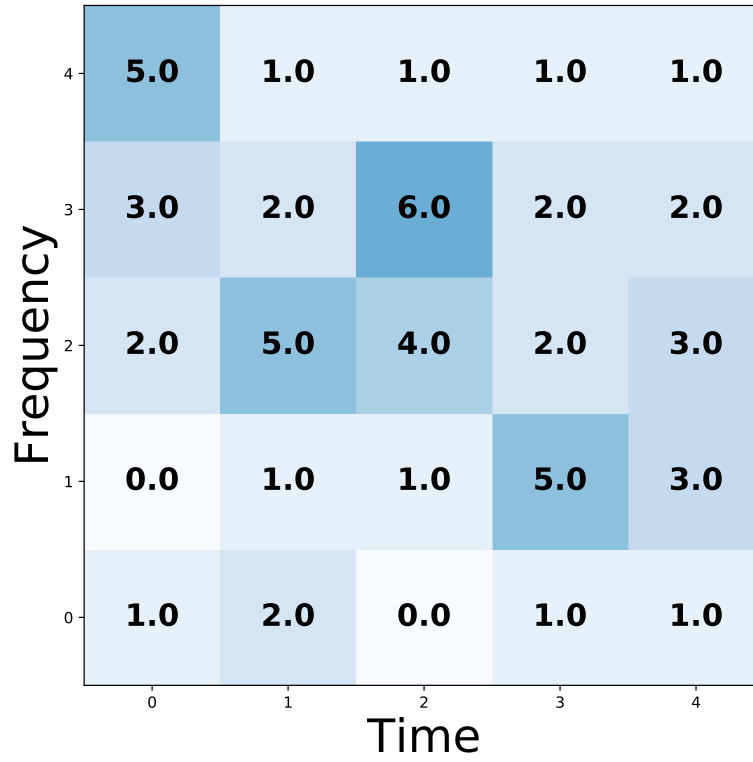
```

1: Input:  $C, T$  {log-likelihood, transition matrix}
2: Output:  $\hat{\mathbf{v}}, V, A$  {most probable track, track probabilities, jumps}
3:
4: Initialisation
5: for Frequency ( $v_{0,k}$ ),  $k = 0 \rightarrow M - 1$  do
6:    $V_{0,k} = C_{0k}$ 
7:    $A_{0,k} = 0$ 
8: end for
9:
10: Iteration
11: for Segment,  $j = 0 \rightarrow N - 1$  do
12:   for Frequency ( $v_{j,k}$ ),  $k = 0 \rightarrow M - 1$  do
13:      $V_{j,k} = \max_i (C_{j,k} + T_i + V_{j-1,j+i})$ 
14:      $A_{j,k} = \operatorname{argmax}_i (C_{j,k} + T_i + V_{j-1,j+i})$ 
15:   end for
16: end for
17:
18: Identification
19:  $\hat{v}_{N-1} = \operatorname{argmax}_k (V_{N-1,k})$ 
20: for Segment,  $j = N - 1 \rightarrow 0$  do
21:    $\hat{v}_j = \hat{v}_{j+1} + A_{j, \hat{v}_{j+1}}$ 
22: end for

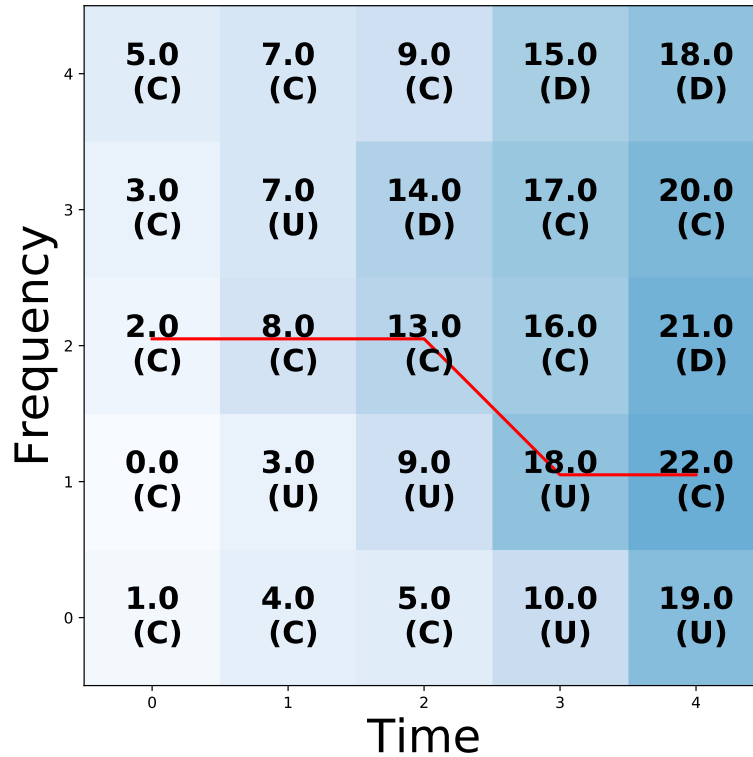
```

ALGORITHM 3.1: The Viterbi algorithm in pseudo-code.  $N$  is the number of segments,  $M$  is the number of frequency bins in each segment. Here the maximisations over  $i$  run between  $\pm(n_1 - 1)/2$  where  $n_1$  is the size of the transition matrix. The values from Eq. 3.6 are stored in  $V$ , and the jumps are stored in  $A$ . The most probable track is denoted by  $\hat{\mathbf{v}}$ .

**Initialisation** The two parts of Eq. 3.6,  $\log p(v_0)$  and  $\log p(\mathbf{x}_0 | v_0)$ , must be computed before the main recursive part of the algorithm can start. Therefore, the initialisation section



(a) The input data



(b) The log-probabilities, jumps, and most probable path

Figure 3.2: Fig. 3.4c shows the observed data, i.e the log-likelihood values  $C_{j,k}$ . Fig. 3.2b shows the calculated log-probabilities  $V_{j,k}$ .  $A_{j,k}$  is shown in parentheses, where the UCD components correspond to  $i = [-1, 0, 1]$  respectively. The red line shows the path that gives the maximum probability. The transition matrix for the UCD jumps is  $[0, 1, 0]$  and corresponds to the unnormalised prior log-probabilities of these jumps occurring.

(lines 5–8) in Alg. 3.1 calculates the first column in the lower panel of Fig. 3.2. A priori, there is no preferred initial frequency, so we take the log-prior  $\log p(v_{0,k})$  to be uniform over the complete frequency range. As a result, this does not affect the maximisation for any jump, therefore, can be omitted from the calculation. We then use the pre-calculated log-likelihood values  $C_{0,k}$  to fill the track probabilities  $V_{0,k}$ . There is no previous position to jump from in this case, so the transition probabilities are irrelevant and  $A_{0,k}$  are set to zero.

**Iteration** The main part of the calculation is the sum in Eq. 3.6. Lines 11–16 in Alg. 3.1 calculate the most probable tracks that end at each frequency bin for each segment by using

$$V_{j,k} = \max_i (C_{j,k} + T_i + V_{j-1,k+i}), \quad (3.9)$$

where  $i$  is the size and direction of the jump. For example, in Fig. 3.2 columns 1–4 are calculated in order using Eq. 3.9, where it maximises over three possible previous positions in frequency. These positions are the frequency bins UCD of the current position. The size and direction of the jump,  $i$ , which gives the maximum probability is then saved to  $A_{j,k}$ . These are shown in parentheses below the log-probabilities in Fig. 3.2 where UCD correspond to values of  $i = [-1, 0, 1]$  respectively.

**Identification** The final stage of the algorithm identifies the most probable track. This is done by initially finding the highest log-probability values in the final time segment,  $\max_k (V_{N-1,k})$  (line 19 in Alg. 3.1). In the lower panel of Fig. 3.2 this is located at position  $j, k = 4, 1$  with  $V_{4,1} = 22$ . To find the track which corresponds to this, the values in  $A_{j,k}$  are followed backwards from this position (lines 20–21). For example, in Fig. 3.2 the final position is  $j, k = 4, 1$  and  $A_{j,k} = \text{Center} = 0$ , this means that at the previous segment the most probable track was at position  $j, k = 4 - 1, 1 + 0 = 3, 1$ . At this time  $A_{3,1} = R = 1$ , therefore, the next track element is at  $j, k = 3 - 1, 1 + 1 = 2, 2$ . This then continues until  $j = 0$  whereupon these retraced positions constitute the most probable track, highlighted in red in Fig. 3.2.

The most probable track is the one traced backwards from the highest probability final segment frequency position. However, tracks can also be traced back from any of the end-frequency positions, returning the most probable track conditional on a given final position. Such tracks should not be confused with the being equal to the second, third, fourth, etc. most probable tracks. Information regarding the rankings and properties of all possible tracks (excluding the most probable and conditionally most probable tracks) is lost during the maximisation procedures computed at each stage in the algorithm – a necessary consequence of the algorithm’s speed and efficiency.

### 3.5 Multiple detectors

If there are  $Q$  detectors operating simultaneously we have  $Q$  sets of data which can be combined appropriately to provide input to the Viterbi search described above. We must also modify the allowed transitions encoded within the transition matrix to take account of the extra prior constraints that are now available.

The received instantaneous frequency of a given astrophysical signal will be nearly the same for all ground-based GW detectors, and our algorithm should be sensitive to tracks that show this consistency in frequency. However there *will* be small differences between the frequencies measured at detectors that are not co-located, due to differential Doppler shifts caused by Earth rotation. As a result the signal could fall in different frequency bins at each detector.

To account for these small differences in signal tracks in each detector, we reference the observed tracks to a third (pseudo) detector located at the centre of the Earth which would be insensitive to Earth spin. The signal frequencies in each real detector are then allowed to vary within a certain number of frequency bins from the track in the reference detector. In the examples that follow, we only consider the possibilities that the track in each real detector is no more than one frequency bin away from the reference track. We can tune the length of the SFTs to ensure this is a valid assumption. As well as differences in signal frequency, due to antenna patterns and other effects, the measured signal amplitude may differ between the detectors. In the following example we assume that the signal has the same amplitude in each detector, however, in Sec. 3.8 we discuss the case where they differ.

We will now show how the algorithm in Sec. 3.4 can be modified to handle a two-detector network (i.e.,  $Q = 2$ ), however any number of detectors can easily be accommodated. In the two detector case the joint probability of two (real) tracks,  $\mathbf{v}^{(1)}$  and  $\mathbf{v}^{(2)}$ , and the geocentric track  $\mathbf{v}$ , given the data, is

$$p(\mathbf{v}, \mathbf{v}^{(1)}, \mathbf{v}^{(2)} | D^{(1)}, D^{(2)}) \propto p(\mathbf{v}) p(\mathbf{v}^{(1)}, \mathbf{v}^{(2)} | \mathbf{v}) p(D^{(1)} | \mathbf{v}^{(1)}) p(D^{(2)} | \mathbf{v}^{(2)}), \quad (3.10)$$

where  $D^{(1)}$  and  $D^{(2)}$  represent the data from the two detectors. The main difference between this and that described in Sec. 3.4 is that the track probabilities  $V_{j,k}$  are stored for the geocentric pseudo-detector. The main iterative calculation (defined for the single detector case in Eq. 3.9) now becomes

$$V_{j,k} = \max_{i,l,m} (C_{j,k+l}^{(1)} + C_{j,k+m}^{(2)} + T_{i,l,m} + V_{j-1,k+i}), \quad (3.11)$$

where  $C^{(1)}$  and  $C^{(2)}$  refer to the log-likelihoods in detectors 1 and 2 respectively and the transition matrix  $T$  is an  $n_1 \times n_2 \times n_3$  matrix, where  $n_1$  dimension refers to the jump from the previous time step,  $n_2$  and  $n_3$  refer to the relative frequency positions in each real detector. The transition matrix is now three-dimensional and holds the prior log-probabilities of  $p(\mathbf{v})$  and

$p(v^{(1)}, v^{(2)} | v)$ . We now need to maximise over three indices:  $i, l$  and  $m$ . The index  $i$  refers to the size and direction of the jump at the geocentre (as before). The indices  $l$  and  $m$  refer to the number of frequency bins by which the two real tracks deviate from the geocentre track. For example, if the most probable track in the geocentred detector is in bin  $j, k = 5, 12$  and the values of  $i, l, m = 0, -1, 1$ , then detector 1 is in position  $j, k = 5, 11$  and detector 2 is in position  $j, k = 5, 13$  and the geocentred track was in the position  $j, k = 4, 12$  at the previous time step. As a result, the track at the geocentre is only affected by Doppler modulations from the Earth's orbit whereas the tracks in the real detectors include Doppler modulations from the Earth's spin.

At every time step the frequency bin position for each real detector is forced to be within  $n_l$  or  $n_m$  bins of the track in the geocentred detector, where  $n_l$  and  $n_m$  depend on how much each detector could possibly be Doppler shifted. As mentioned previously, we only consider the case where  $n_l = 1$  and  $n_m = 1$ , allowing the track from each real detector to be at most one frequency bin away from the geocentred track position. While we tune the SFT length to keep this condition for different frequencies, it is also possible to tune the values of  $n_l$  and  $n_m$  to get a similar effect. The implementation of the multi-detector algorithm is similar to the single detector case described in Sec. 3.4. However in the single detector case there is only a single variable to be maximised over for each time-frequency bin. This variable is the frequency jump from the position in the previous segment. For the multi-detector case there are at least three variables to be maximised over: the probability of the jump,  $i$ , at the geo-centre and the probability of the signal being in the surrounding positions in each of  $Q$  real detectors,  $l, m, \dots$ . The values of  $i, l, m, \dots$  are then saved to  $A_{j,k}$  and are ultimately used to reconstruct the most probable consistent tracks in each real detector.

As in Sec. 3.4, there are three main sections: Initialisation, iteration, and the identification. For the multi-detector case each element is modified as follows.

**Initialisation** The first-row calculation (lines 5–8) in Alg. 3.1, are now modified to additionally maximise over the real detector track positions  $l$  and  $m$ . For each time-frequency bin the maximum sum of the log-likelihoods is saved together with the frequency locations of the corresponding tracks in the real detectors. The index  $i = 0$  is kept constant as there is no previous position.

**Iteration** To process the subsequent time segments, lines 13–14 in Alg. 3.1 are modified to account for two (or more) detectors. Line 13 of Alg. 3.1 is changed to calculate Eq. 3.11, the log-probability of a track at the geocentre ending in bin  $j, k$  given that signal is in the real detector positions of  $j, k + l$  and  $j, k + m$ . Line 14 is then modified so that  $A_{j,k}$  stores the jump values,  $i$ , and the real detector positions,  $l$  and  $m$ , which returned the highest probability.

**Identification** The most probable track is identified in the same way as for the single detector case, first by finding the maximum value in the final time step of  $V_{j,k}$  (line 19 in Alg. 3.1).



The track at the geocentre can then be found by iteratively following the jump values stored in  $A_{j,k}$  back from this position. The track in each of the real detectors is determined by using the values of  $l$  and  $m$  indices also stored in  $A_{j,k}$  to find the relative position of the track in each real detector compared to the geocentre.

This method can be extended to more than two detectors by including additional datasets and expanding the corresponding number dimensions of the maximisation procedures in the iterative steps.

### 3.6 Memory

In this section we extend the basic Viterbi algorithm to improve its sensitivity to non-stochastic signals where there is some knowledge of its frequency evolution. We do this by including a form of ‘memory’ and this extension applies to both the single and multiple-detector cases. Rather than considering only the previous step in our decision-making process, we now include the previous  $m + 1$  steps and expand the transition matrix to include these values. A memory of  $m = 0$  therefore corresponds to the methods described in previous sections. With a non-zero memory the transition matrix can a-priori make certain sequences of jumps more probable and assign different prior probabilities for these jump sequences e.g., ‘up then centre’ may be less preferable to ‘centre then centre’. As a result we can increase the chance of the most probable track matching an expected astrophysical signal. In a single detector search with a memory of  $m = 1$ , if we only allow UCD transitions, then for every frequency bin we save 3 values. These are proportional to the log-probabilities of a track coming from a UCD bin in the previous time step, where the maximisation is over the corresponding UCD bins two time steps back. Eq. 3.11 then is then modified to,

$$V_{j,k,s} = \max_h (C_{j,k} + T_{s,h} + V_{j-1,k+s,k+s+h}), \quad (3.12)$$

where  $s$  and  $h$  refer to the UCD jumps at the time step  $j - 1$  and  $j - 2$  respectively. Similar to the previous two sections, the algorithm is split into three parts: initialisation, iteration, and the track identification:

**Initialisation** The initialisation process needs to populate the first  $m + 1$  steps before the main iteration can start. At the first time step, the elements  $V_{0,k,s}$  are set to the log-likelihoods  $C_{0,k}$  as in Sec. 3.4. There is no previous time step, so the element  $s$  is not relevant. At the second time step,  $V_{1,k,s}$  is calculated using Eq. 3.12, where there is no maximisation over  $h$ , it is assumed to be 0, or a center jump. As there is no data before  $j = 0$ , the maximisation at this point will always return the jump which has the largest prior probability, which in this case is a center jump. Therefore, the maximisation returns the same value for all frequency bins and can be set to a center jump.



**Iteration** For all following time steps the values for each element of  $V_{j,k,s}$  in Eq. 3.12 are calculated. This quantity is proportional to the log-probability of the track ending in time-frequency bin  $j,k$ , which was in the previous position of  $j-1, k+s$ . The corresponding value of  $h$  that maximised the log-probability of the track is recorded in  $A_{j,k,s}$ .

**Identification** The most probable track is identified in a similar way to the non-memory cases, by finding the highest-valued last element,  $V_{N-1,k,s}$ . The values of  $s$  and  $h$  are then followed back to find the most probable track. As an example, let us assume the most probable track finishes in bin  $j,k,s = 10,5,0$ , where the value of  $m$  is  $A_{10,5,0} = 1 = \text{up}$ . The previous position is then  $j,k,s = 10-1, 5+s, m = 10-1, 5+0, 1 = 9,5,1$  with a value  $A_{9,5,1} = 0 = \text{Center}$ , and the next track position is  $j,k,s = 9-1, 5+1, 0 = 8,6,0$  etc. The values of  $j,k$  along this track describes most probable path.

The number of elements over which one must search increases rapidly with memory length, and has a strong impact on the computational cost of the analysis. For the single detector Viterbi approach the number of calculations made is  $3 \times N \times M$  if we only allow UCD jumps, where  $N$  and  $M$  are the number of time and frequency bins respectively. When memory is included this increases to  $3^{m+1} \times N \times M$ .

## 3.7 Summed input data

In this section a method of incoherently-summing a set of SFTs to increase the SNR of a signal in a segment is outlined. To be more precise, it is actually the log-likelihoods which are summed, i.e. the quantity in Eq. 3.8. We can write the new summed set of data  $F_j$  as,

$$F_j = \sum_i^{N_s} C_{i,k} \quad (3.13)$$

where  $N_s$  is the number of SFTs to sum together and the log-likelihood  $C(v_{i,k})$  is defined in Eq. 3.8. We can see this is possible by looking at Eq. 3.7, where we can use the product of likelihoods,

$$\begin{aligned} p(D | v) &\propto p(x_1, x_2 \dots x_n | v) \\ &\propto p(x_1 | v) \dots p(x_n | v) \\ &\propto \exp \left( \sum_i C_{j,k} \right). \end{aligned} \quad (3.14)$$

If the data contains gaps where the detector was not observing, then we fill the gaps in the power spectrum with a constant value which is the expectation value of the log-likelihood. The procedure of filling in the gaps of the data is completed before any summing. Therefore, the data

should have the same mean regardless of how much real data is in each sum. In the examples that follow, we sum the SFTs over the length of one day.

The main motivation for summing the data is to increase the SNR of a signal in the segments. The risk is that a signal can move between adjacent frequency bins during a day. To reduce this risk, we choose the frequency bin width such that it is more likely that a signal will be contained within a single frequency bin that cross a bin edge. In practise, to ensure that this is true, the segment or SFT length and the number of segments which are summed can be tuned for each search. As well as increasing the SNR, summing over one day should average out the antenna pattern. This means that the log-likelihood value in any bin should be more similar between detectors, however, there is still some variation due to the sky localisation and polarisation.

This also has two main effects on the transition matrix, the first is that as each segment of data is now one day long, a jump between frequency bins is far more likely, therefore, the transition matrix elements are modified to account for this. The second is that as the data is averaged over one day, the signal should remain in the same frequency bin between detectors, therefore, there is no longer a need for the multi-dimensional transition matrix described in Sec. 3.5.

The volume of the data is also reduced by a factor of  $1/N_s$ , therefore, the time taken for the algorithm to run is also reduced by the same factor.

### 3.8 Line-aware statistic

The multiple-detector algorithm described in Sec. 3.5 returns the most probable track of a common signal assumed to be in Gaussian noise. As a consequence the algorithm will return large values of the log-likelihood even if there are inconsistent values of SFT power between the detectors, either from non-Gaussian noise or because the signal is not equally strong in the two detectors. However a signal with unequal power in the two detectors is more likely to be a non-Gaussian instrumental line than an astrophysical signal. The line-aware statistic described in this section is designed to make the search more robust to such instrumental artefacts within realistic non-Gaussian data whilst maintaining sensitivity to astrophysical signals.

For most of the analysis examples presented here we use data which is the incoherent sum of 30-minute normalised SFTs over a day (described in more detail in Sec. 3.7). As a result the effects of the detector antenna patterns and of differential Doppler shifts are significantly reduced, and any signal should have a broadly similar summed log-likelihood in the same frequency bin in each detector. The statistic can then be modified such that we expect a similar log-likelihood in each detector.

We first consider the model of Gaussian noise with no signal present. Within a single summed segment, the likelihood of Gaussian noise at frequency  $\nu$  is given by a  $\chi^2$  distribution,

$$p(F_j|\nu_j, M_N, I) = \frac{1}{2^{d/2}\Gamma(d/2)} F_j^{d/2-1} \exp\left\{-\frac{F_j}{2}\right\} \quad (3.15)$$

where  $F_j$  is the frequency domain power summed over sub-segments within a single day, as described in Sec. 3.7 and  $d$  is the number of degrees of freedom, equal to twice the total number of summed SFTs.  $M_N$  represents the model that the data is simply Gaussian noise. In the presence of a signal (model  $M_S$ ), the power should follow a non central  $\chi^2$  distribution in which the non-centrality parameter  $\lambda$  is the square of the SNR, ( $\lambda = \rho_{\text{opt}}^2$ ), i.e.

$$p(F_j | \mathbf{v}_j, \lambda, M_S, I) = \frac{1}{2} \exp \left\{ -\frac{F_j + \lambda}{2} \right\} \left( \frac{F_j}{\lambda} \right)^{d/4-1/2} I_{d/2-1} \left( \sqrt{\lambda F_j} \right). \quad (3.16)$$

If a signal is present we therefore expect the SFT powers in both detectors to follow Eq. 3.16. Assuming for the moment that the noise variance is the same in both, we can determine the evidence for model  $M_S$  by marginalising over  $\lambda$ ,

$$p(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j, M_S, I) = \int_0^\infty p(\lambda, w_s) p(F_j^{(1)} | \mathbf{v}_j, \lambda, M_S, I) p(F_j^{(2)} | \mathbf{v}_j, \lambda, M_S, I) d\lambda. \quad (3.17)$$

We set the prior on  $\lambda$  to be an exponential distribution of width  $w$ , this is done somewhat arbitrarily as we expect the majority of signals to have a low SNR. This distribution follows,

$$p(\lambda, w) = \exp \left( \frac{-\lambda}{w} \right). \quad (3.18)$$

On the other hand, if an instrumental line is present in one of the detectors we expect to see signal-like power in that detector and noise-like power in the other. The evidence for this ‘line’ model ( $M_L$ ) is therefore

$$p(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j, M_L, I) = \int_0^\infty p(\lambda, w_L) \left[ p(F_j^{(1)} | \mathbf{v}_j, M_N, I) p(F_j^{(2)} | \mathbf{v}_j, \lambda, M_S, I) + p(F_j^{(1)} | \mathbf{v}_j, \lambda, M_S, I) p(F_j^{(2)} | \mathbf{v}_j, M_N, I) \right] d\lambda, \quad (3.19)$$

The third option to consider is the simple case of approximately Gaussian noise in both of the detectors,

$$p(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j, \lambda, M_G, I) = p(F_j^{(1)} | \mathbf{v}_j, M_G, I) p(F_j^{(2)} | \mathbf{v}_j, M_G, I). \quad (3.20)$$

The posterior probability of model  $M_{GL}$ , which contains the probability of Gaussian noise or

Gaussian noise with a line in one detector, (taken as mutually exclusive) is

$$p(M_{\text{GL}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j, I) = p(M_{\text{G}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j, I) + p(M_{\text{L}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j, I), \quad (3.21)$$

where we assume that  $M_{\text{G}}$  and  $M_{\text{L}}$  are mutually exclusive.

We can now find the posterior odds ratio for the presence of a signal over noise or a line,

$$\begin{aligned} O_{\text{SGL}}(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j) &= \frac{p(M_{\text{S}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j)}{p(M_{\text{GL}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j)} = \frac{p(M_{\text{S}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j)}{p(M_{\text{G}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j) + p(M_{\text{L}} | F_j^{(1)}, F_j^{(2)}, \mathbf{v}_j)} \\ &= \frac{p(M_{\text{S}})p(F_j^{(1)}, F_j^{(2)} | M_{\text{S}}, \mathbf{v}_j)}{p(M_{\text{G}})p(F_j^{(1)}, F_j^{(2)} | M_{\text{G}}, \mathbf{v}_j) + p(M_{\text{L}})p(F_j^{(1)}, F_j^{(2)} | M_{\text{L}}, \mathbf{v}_j)} \\ &= \frac{p(F_j^{(1)}, F_j^{(2)} | M_{\text{S}}, \mathbf{v}_j)p(M_{\text{S}})/p(M_{\text{G}})}{p(F_j^{(1)}, F_j^{(2)} | M_{\text{G}}, \mathbf{v}_j) + p(F_j^{(1)}, F_j^{(2)} | M_{\text{L}}, \mathbf{v}_j)p(M_{\text{L}})/p(M_{\text{N}})} \end{aligned} \quad (3.22)$$

In practice it is convenient to use the log odds ratio,

$$\begin{aligned} \log [O_{\text{SGL}}(F_j^{(1)}, F_j^{(2)})] &= \log [p(F_j^{(1)}, F_j^{(2)} | M_{\text{S}})] \\ &\quad - \left[ \log (p(F_j^{(1)}, F_j^{(2)} | M_{\text{G}}) \right. \\ &\quad \left. + p(F_j^{(1)}, F_j^{(2)} | M_{\text{L}})p(M_{\text{L}})/p(M_{\text{G}})) \right] \end{aligned} \quad (3.23)$$

As we are only interested in the maximum of  $\log [O_{\text{SGL}}(F_j^{(1)}, F_j^{(2)})]$ , the factor  $\log [p(M_{\text{S}})/p(M_{\text{G}})]$  can be dropped from the expression.

In this version of the Viterbi algorithm, rather than storing a value proportional to the log-probabilities as in Sec. 3.5, here we store a value proportional to the log-odds ratio. Here we take the log-odds ratio defined in Eq. 3.23, which is the log-odds of a signal having a similar power in each detector, and add the log-prior odds  $p(\mathbf{v} | M_{\text{S}})/(p(\mathbf{v} | M_{\text{N}}) + p(\mathbf{v} | M_{\text{L}}))$  which is the log-prior or any particular track. By assuming that the track transitions for the line and noise model are equally probable for any jump, we set the denominator of the prior-odds is a constant  $b$ . This then means Eq. 3.11 is modified to,

$$\begin{aligned} \hat{V}_{i,j} &= \max_{k,l,m} (T_{k,l,m} + b + V_{i-1,j+k} \\ &\quad + \log [O_{\text{SGL}}(F_j^{(1)}, F_j^{(2)})]) \end{aligned} \quad (3.24)$$

where  $\hat{V}$  refers to a log-odds ratio. The maximised statistic now has three tuneable parameters: the width,  $w_{\text{S}}$  in Eq. 3.18, on the prior for a signal SNR squared,  $p_{\text{S}}(\lambda)$ , the width,  $w_{\text{L}}$  of

the prior in the case of a line,  $p_L(\lambda)$ , and the ratio of the prior on the line and noise models,  $p(M_L)/p(M_G)$ . These parameters are optimised for each search, where we initially estimate the SNR of a signal we hope to be sensitive to in each time slice, then use this as a guide for the width of the signal prior. This is then repeated for an expected line SNR and this is used for the width of the line prior. The ratio of line and noise models runs in the range 0 to 1, we set this limit as we do not expect an instrumental line to be as likely as Gaussian noise in any particular frequency bin.

This line-aware statistic can be applied in a more powerful way when we use multiple detectors and is similar to the approach in [87]. The multiple-detector algorithm described in Sec. 3.5 returns the most probable track of a common signal assumed to be in Gaussian noise. As a consequence the algorithm will return large values of the log-likelihood even if there are inconsistent values of SFT power between the detectors, either from non-Gaussian noise or because the signal is not equally strong in the two detectors. However a signal with unequal power in the two detectors is more likely to be a non-Gaussian instrumental line than an astrophysical signal. The line-aware statistic described in this section is designed to make the search more robust to such instrumental artefacts within realistic non-Gaussian data whilst maintaining sensitivity to astrophysical signals.

For most of the analysis examples presented here we use data which is the incoherent sum of 30-minute normalised SFTs over a day (described in more detail in Sec. 3.7). As a result the effects of the detector antenna patterns and of differential Doppler shifts are significantly reduced, and any signal should have a broadly similar summed log-likelihood in the same frequency bin in each detector. The statistic can then be modified such that we expect a similar log-likelihood in each detector.

In a similar way to the single-detector case, we can write out the evidence for each of the three models as follows. If a signal is present we therefore expect the SFT powers in both detectors to follow Eq. 3.16. Assuming for the moment that the noise variance is the same in both, we can determine the evidence for model  $M_S$  by marginalising over  $\lambda$ ,

$$p(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j, M_S, I) = \int_0^\infty p(\lambda, w_s) p(F_j^{(1)} | \mathbf{v}_j, \lambda, M_S, I) p(F_j^{(2)} | \mathbf{v}_j, \lambda, M_S, I) d\lambda. \quad (3.25)$$

We set the prior on  $\lambda$  the same as in the single detector case in Eq. 3.18. In this case, if an instrumental line is present in one of the detectors we expect to see signal-like power in that

detector and noise-like power in the other. The evidence for this ‘line’ model ( $M_L$ ) is therefore

$$\begin{aligned} p(F_j^{(1)}, F_j^{(2)} | v_j, M_L, I) &= \int_0^\infty p(\lambda, w_L) \\ &\left[ p(F_j^{(1)} | v_j, M_N, I) p(F_j^{(2)} | v_j, \lambda, M_S, I) \right. \\ &\left. + p(F_j^{(1)} | v_j, \lambda, M_S, I) p(F_j^{(2)} | v_j, M_N, I) \right] d\lambda, \end{aligned} \quad (3.26)$$

The third option is the simple case of approximately Gaussian noise in both of the detectors,

$$\begin{aligned} p(F_j^{(1)}, F_j^{(2)} | v_j, \lambda, M_G, I) &= p(F_j^{(1)} | v_j, M_G, I) \\ &p(F_j^{(2)} | v_j, M_G, I). \end{aligned} \quad (3.27)$$

We can now find the posterior odds ratio for the presence of a signal over noise or a line by following the same steps as in Eq. ???. Once again we write this as a log-odds ratio,

$$\begin{aligned} \log \left[ O_{S/GL}^{(2)}(F_j^{(1)}, F_j^{(2)}) \right] &= \log \left[ p(F_j^{(1)}, F_j^{(2)} | M_S) \right] \\ &- \left[ \log \left( p(F_j^{(1)}, F_j^{(2)} | M_G) \right) \right. \\ &\left. + p(F_j^{(1)}, F_j^{(2)} | M_L) p(M_L) / p(M_G) \right] \end{aligned} \quad (3.28)$$

The factor  $\log [p(M_S)/p(M_G)]$  can again be dropped from the expression.

For the multi-detector case we then modify Eq. 3.11 to,

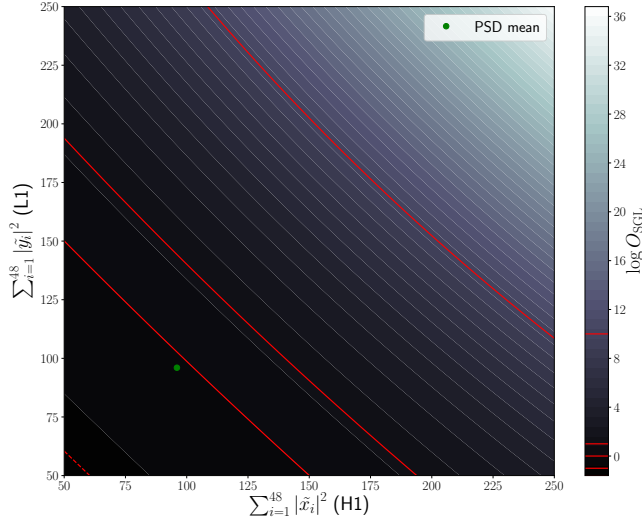
$$\begin{aligned} \hat{V}_{i,j} &= \max_{k,l,m} (T_{k,l,m} + b + V_{i-1,j+k} \\ &+ \log \left[ O_{S/GL}^{(2)}(F_j^{(1)}, F_j^{(2)}) \right]), \end{aligned} \quad (3.29)$$

where  $\hat{V}$  refers to a log-odds ratio. This is then optimised over the same three parameters as the single detector case.

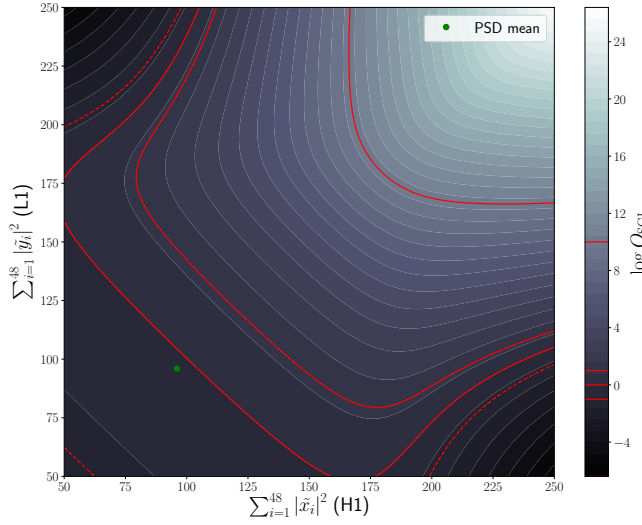
Fig. 3.3 shows an example of the output of the statistic in Eq. 3.28 for different FFT powers  $F$ .

### 3.9 Line aware statistic for consistent amplitude

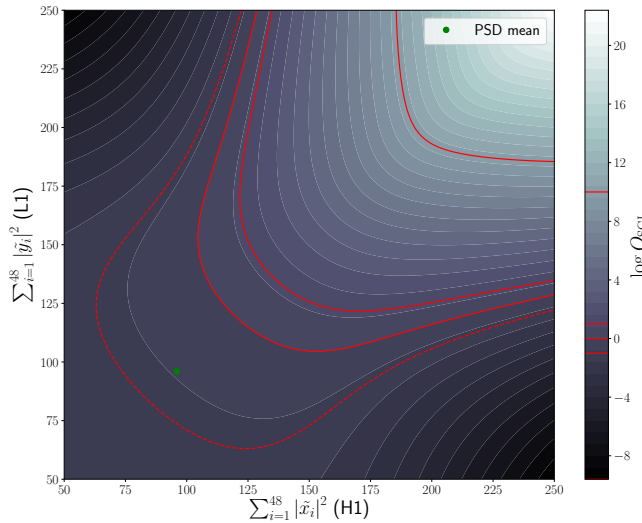
In Sec. 3.8 the ‘line aware’ statistic was designed to penalise high SFT powers in a single detector and reward powers which have a similar SNR. This is often a useful statistic to use when the detectors have similar sensitivities, however, this is not always the case. During an observing run of a gravitational wave detector, their sensitivity will vary with time due to fluctuating or new noise sources or potentially upgrades which can increase the sensitivity. A change in the sensitivity, or noise floor, affects the SNR of a possible signal in the data, i.e. a lower noise floor



(a) This shows the distribution of the lines aware statistic plotted against the FFT power in each detector. This example is for parameters  $p_s(\lambda) = 4, p_l(\lambda) = 0$  and  $p(M_L)/p(M_G) = 0$ . So the line part of the statistic is not operating.



(b) This shows the distribution of the lines aware statistic plotted against the FFT power in each detector. This example is for parameters  $p_s(\lambda) = 4, p_l(\lambda) = 5$  and  $p(M_L)/p(M_G) = 0.03$ . Here we include the line part of the statistic.



(c) This shows the distribution of the lines aware statistic plotted against the FFT power in each detector. This example is for parameters  $p_s(\lambda) = 4, p_l(\lambda) = 5$  and  $p(M_L)/p(M_G) = 1$ . Here the effect of lines is expected to be larger than the previous panel on the search. Therefore, the statistic forces the two detectors to have more similar power.

Figure 3.3: Lookup tables using the line aware statistic in Eq. 3.29.

results in a higher SNR. In this section the above ‘line aware’ statistic is modified to account for the difference in sensitivities of the detectors, and therefore search for a consistent amplitude between detectors as opposed to SNR.

There are two main factors which are taken into account when determining how sensitive a detector is in a particular time interval: the PSD of detector and the duty cycle. The PSD of the detector is essentially how sensitive the detector is at that time and the duty cycle is the fraction of time in a given interval that the detector was collecting data. A decrease in the duty cycle and an increase in the PSD will decrease the SNR and vice-versa. To search for consistent amplitude Eq. 3.24 is modified by weighting each detector by its PSD and duty cycle.

The definition of SNR is taken from [ ] as,

$$\rho_0^2 = \frac{h_0^2 T}{S} (\alpha_1 A + \alpha_2 B + \alpha_3 C), \quad (3.30)$$

where  $\rho_0$  is the optimal SNR,  $h_0$  is the signal amplitude,  $T$  is the time of observation,  $S$  is the noise PSD and the term in brackets include the antenna pattern of the detector. The signal with amplitude  $h_0$  will be the same amplitude at both detectors (H and L), therefore we can relate the SNR in each detector by,

$$\rho_L^2 = \frac{\rho_H^2 S_H T_L}{S_L T_H} \frac{(\alpha_1 A_L + \alpha_2 B_L + \alpha_3 C_L)}{(\alpha_1 A_H + \alpha_2 B_H + \alpha_3 C_H)}. \quad (3.31)$$

For the majority of the analysis that follows the SFTs are summed over one day, this will be explained in greater detail in Sec. 3.7. The components in the above equation which have the form  $(\alpha_1 A + \alpha_2 B + \alpha_3 C)$ , account for the antenna pattern of the earth as it rotates. These can be approximated to be the same for the two detectors H1 and L1 as we will essentially be averaging out the daily modulation by summing SFTs. Therefore we can simplify the above Eq. 3.31 to,

$$\rho_L^2 = \frac{\rho_H^2 S_H T_L}{S_L T_H} = l \rho_H^2. \quad (3.32)$$

This then gives a factor  $l = S_H T_L / S_L T_H$  which relates the SNR of each detector, where  $S$  and  $T$  are values that are known for a given data-set prior to running the search.

This ratio of SNRs can be included in the integral over SNR for the signal model in Eq. 3.25 as follows,

$$p(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j, M_S, I) = \int_0^\infty p(\lambda, w_s) p(F_j^{(1)} | \mathbf{v}_j, \lambda, M_S, I) p(F_j^{(2)} | \mathbf{v}_j, l\lambda, M_S, I) d\lambda. \quad (3.33)$$

Similarly, the line model in Eq. 3.26 can be modified as,



$$p(F_j^{(1)}, F_j^{(2)} | \mathbf{v}_j, M_L, I) = \int_0^\infty p(\lambda, w_L) \left[ p(F_j^{(1)} | \mathbf{v}_j, M_N, I) p(F_j^{(2)} | \mathbf{v}_j, l\lambda, M_S, I) + p(F_j^{(1)} | \mathbf{v}_j, \lambda, M_S, I) p(F_j^{(2)} | \mathbf{v}_j, M_N, I) \right] d\lambda. \quad (3.34)$$

Fig. 3.4 shows an example of the values of the statistic described in Eq. 3.34 plotted against a range of FFT powers from each detector. This demonstrated how the statistic accounts for a difference in sensitivity on detectors by allowing the FFT power or effectively SNR to vary more.

In Fig. 3.4 we show an example of two detectors with large differences in sensitivity, and how the statistic which takes this into account can improve the search sensitivity in this case.

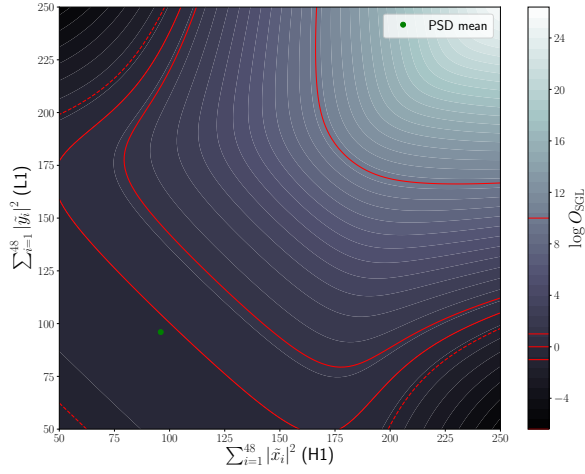
### 3.10 Testing the algorithm

The sensitivity of the algorithm was tested by searching for artificial signals from isolated pulsars added to three types of noise-like data: continuous Gaussian noise, Gaussian noise but with periods of missing data, and real detector data (the S6 MDC [48]). The S6 MDC refers to a standardised set of simulated signals which are injected into real data, this set is also what is used for the injections into the two Gaussian noise cases. We describe each of the tests in more detail in Sec. 3.10.1, 3.10.2 and 3.10.3, but several common pre-processing steps are performed before running these datasets through the Viterbi algorithm:

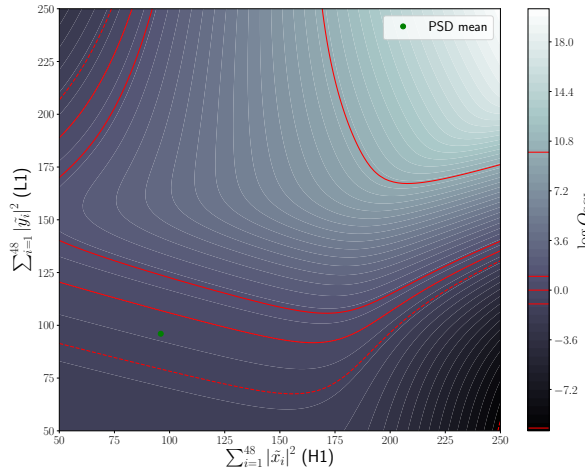
1. We read SFTs generated from 1800 s stretches of data in 2 Hz bands between 100 and 200 Hz. The SFTs length is chosen to ensure that any signal is likely to be contained within the width of a single frequency bin during the length of one day, rather than being split across the bin edges (see below).
2. We estimate the noise PSD for each SFT by calculating a running median over frequency using LALSuite code `XLALSFTtoRnmed` [73], this includes a bias factor to convert this to the mean and has a width of 100 bins. We then normalise the SFT by dividing it by its running median, giving the noise-like parts of the spectrum a mean power of approximately one.
3. The SFTs are then summed over one day, as described in Sec. 3.7. The signal parameters are chosen so that within the frequencies of the search, the signal will not fall in more than two frequency bins over this period.

The differential Doppler shift of a signal seen at two detector sites due to the Earth's rotation  $\Delta f_{\text{rot}}^{(1,2)}$  is simply

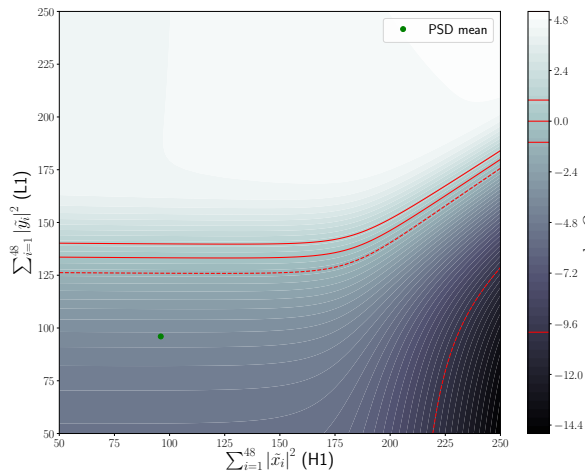
$$\Delta f_{\text{rot}}^{(1,2)} = \frac{(\mathbf{v}^{(1)} - \mathbf{v}^{(2)}) \cdot \hat{\mathbf{s}}}{c} f_0, \quad (3.35)$$



(a) This shows the distribution of the lines aware statistic plotted against the FFT power in each detector. This example is for parameters  $p_s(\lambda) = 2, p_l(\lambda) = 0$  and  $p(M_L)/p(M_G) = 0$ . So the line part of the statistic is not operating.



(b) This shows the distribution of the lines aware statistic plotted against the FFT power in each detector. This example is for parameters  $p_s(\lambda) = 2, p_l(\lambda) = 2$  and  $p(M_L)/p(M_G) = 1$ . Here the line part of the statistic has the same SNR and the signal part, i.e. we expect the SNR of a signal to be similar to that of a line.



(c) This shows the distribution of the lines aware statistic plotted against the FFT power in each detector. This example is for parameters  $p_s(\lambda) = 2, p_l(\lambda) = 10$  and  $p(M_L)/p(M_G) = 1$ . Here we expect the SNR of a line to be larger than a signal.

Figure 3.4: Lookup tables using the line aware statistic for consistent amplitude as in Sec. 3.9. Each of these use the parameters  $p_s(\lambda) = 4, p_l(\lambda) = 5$  and  $p(M_L)/p(M_G) = 0.03$ .

where  $\mathbf{v}^{(1,2)}$  is the velocity of detector 1, 2 in an inertial reference frame,  $f_0$  is the instantaneous signal frequency in the frame,  $\hat{\mathbf{s}}$  is the unit vector in the direction of the source and  $c$  is the speed of light. The maximum difference in frequency seen by the two LIGO detectors is

$$\Delta f_{\text{rot}} \approx 6.5 \times 10^{-7} f_0, \quad (3.36)$$

so the frequency measured from a source in the equatorial plane with  $f_0 = 200$  Hz will differ by up to  $1.3 \times 10^{-4}$  Hz in the two detectors. This is  $\sim 4$  times smaller than the frequency bin width of 1800 s SFTs ( $5.6 \times 10^{-4}$  Hz), so signals at frequencies lower than this are likely to appear in the same frequency bin in the two detectors. Therefore, whilst at higher frequencies we still allow the signal to be in different frequency bins between the detectors, in the following searches, we do not allow this.

4. The data is then split into 0.1 Hz sub-bands which are overlapping by 0.05 Hz. These were chosen to ensure that signals are contained within a sub-band over the year. On these timescales the important contributions to the frequency evolution are the spin-down rate of the pulsar and the Doppler shift due to the earth orbit. To investigate the doppler shift, we can look at a signal at 200 Hz, using Eq. 3.35 we can calculate the maximum shift in frequency due to the earths orbit as,

$$\Delta f_{\text{orbit}} = \frac{2\pi R_o}{T_o} \frac{1}{c} f_0 \approx 9.9 \times 10^{-5} f_0, \quad (3.37)$$

where  $T_o$  and  $R_o$  are the earth orbit time and radius. This gives a maximum doppler shift of 0.019 Hz, this is a  $\sim 1/5$  of the width of a sub-band, therefore, is more likely to be totally contained within a sub-band than crossing over the edge. To account for the cases where the signal frequency crosses over the edge of a sub-band, the sub-bands overlap by 0.05 Hz so that the majority of the signals should be completely contained within at least one of the sub-bands. To investigate the spin-down of the pulsar, we look at the length of data,  $T = 4.05 \times 10^7$  s and we choose a sub-band width of 0.1 Hz. For a signal to drift over the width of a whole sub-band we would need  $\dot{f}$  of,

$$\frac{df}{dt} > \left| \frac{-0.1}{4.05 \times 10^7} \right| = 2.4 \times 10^{-9} \text{ Hz/s}. \quad (3.38)$$

The majority of the injections that follow satisfy this condition, signals which are greater than this, and therefore drift over multiple bands, are vetoed from the search.

5. The two detector Viterbi algorithm is then run using the line aware statistic (see Sec. 3.8). There are 4 parameters which we optimise in this search. The transition probabilities, where we have one parameter  $\tau$  which is the ratio of the probability of going straight to the probability of going either up or down. Due to the averaging procedure, the sig-

nals received at each detector are forced to follow a common track which is equal to the ‘imaginary’ detectors track. The other three parameters,  $w_S, w_L$  and  $p(M_L)/p(M_N)$ , are described in Sec. 3.8.

6. The algorithm then returns the most probable track through the data, and the value  $\propto$  the log-odds in the final time step, i.e., the maximum final value,  $\max_j(V_{N,j})$ , in Eq. 3.24, which is then our detection statistic.

As an example of what the algorithm returns, Fig. 3.5 shows the tracks in the two detectors, H1 and L1. This also shows the log-odds ratio of ending in any frequency bin, i.e., all the elements in Eq. 3.24. In this figure, each time segment of the odds ratios have been normalised such that the sum of the odds ratios is 1.

In the following tests there are two main quantities which we use to determine the sensitivity. These are sensitivity depth  $\mathcal{D}$  and the optimal SNR  $\rho$ . The sensitivity depth,  $\mathcal{D}$ , is defined in [74] as,

$$\mathcal{D}(f) = \frac{\sqrt{S_h(f)}}{h_0}, \quad (3.39)$$

where  $S_h(f)$  is the single-sided noise PSD and  $h_0$  is the GW amplitude. The optimal SNR is defined as,

$$\rho^2 = \sum_X 4\Re \int_0^\infty \frac{\tilde{h}^X(f) \tilde{h}^{X*}(f)}{S^X(f)} df, \quad (3.40)$$

where  $X$  indexes the detectors and  $\tilde{h}(f)$  is the Fourier transform of the time series of the signal  $h(t)$ . This expression is defined in [47] for a double-sided PSD and we have defined it for the more common single-sided case.

### 3.10.1 S6 injections into gapless Gaussian noise

The first test involves injecting signals into Gaussian noise. The power spectrum of a Gaussian noise time-series follows a  $\chi^2$  distribution with two degrees of freedom, therefore, as we search through the power spectrum, we generate spectrograms which follow a  $\chi^2$  distribution. These spectrograms are 0.1 Hz wide and are set at 0.05 Hz intervals between 100 Hz and 200 Hz. The bins are 1./1800 Hz wide and 1800s long, where the total length of data is the same as S6, i.e.,  $\sim 1.3$  years. We then generate the signals, where the pulsars parameters are fixed to the same values as the injections in the S6 MDC in this band, these values are outlined in [48].

The values of  $f_0$  for the injections were not always centred in a sub-band, therefore a number of sub-bands contained only part of the injected signal. These sub-bands were ignored as they contaminated the signal statistics and only the sub-band which contained the whole signal was accepted. This reduced the number of sub-bands from 2000 to 1762 with the removal of 238 sub-bands containing only part of a signal. This set also includes signals that drift across multiple sub-bands due to their high spin-down rate. Only two signals were removed due to their spin-

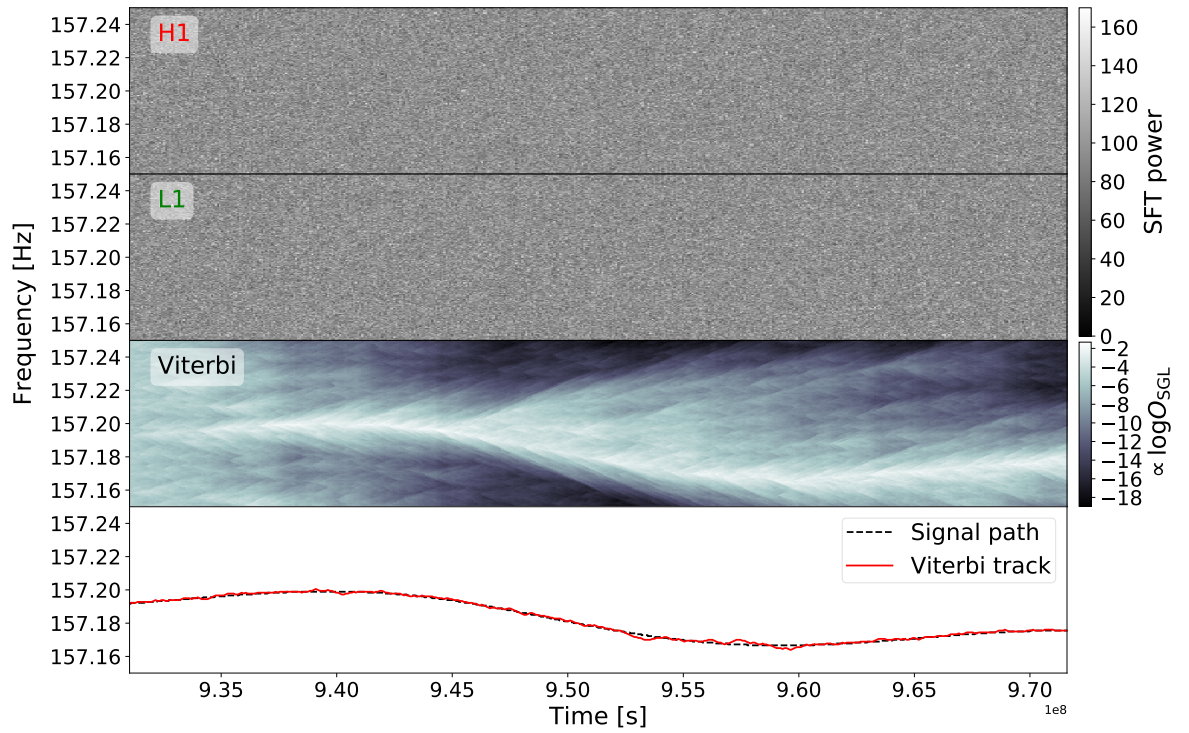


Figure 3.5: The results that the SOAP algorithm returns from an injection with an optimal SNR of 90, i.e., the SNR in H1 is 64 and the SNR in L1 is 62. The signal is injected into Gaussian noise, where the 1800 s SFTs have been summed over 1 day. The top panel shows a simulation of summed SFTs from H1, the second panel shows the same for L1, the third panel shows the values proportional to the log-odds ratios in Eq. 3.24. The log-odds have been normalised such that the sum of all the odds ratios in every time bin are equal to 1. The bottom panel shows the injected signal track (black dotted) and the track found in the ‘imaginary’ detector by the two-detector SOAP search with the line-aware statistic (red), both of these tracks are at the geo-centre. In this case the RMS of the difference between the Viterbi track and injected signal track was  $\sim 1$  bin, where 1 bin is 0.00056 Hz wide.

down values, which were  $> 5 \times 10^{-9}$  Hz/s, these were the two hardware injections in the 100-200 Hz band.

For each injection the SOAP algorithm returns the detection statistic described in Sec. 3.8 and 3.10. We calculate a false alarm rate, which is the fraction of bands that have no injection that do exceed a given threshold. This is set to 1% and is used as a detection threshold. We then take all of the bands and if they pass the threshold we set them as detected, i.e., 1, and if they do not they are set as not detected, i.e., 0. This then leaves us with a set of binomial data, where the efficiency curves later in the paper are sigmoids which have been fitted to this. The sigmoid follows,

$$s(x; x_0, k) = \frac{1}{1 - \exp(-k(x - x_0))}. \quad (3.41)$$

The fit is done by sampling the posterior, i.e.,

$$p(x_0, k | b) \propto p(x_0, k) p(x | x_0, k), \quad (3.42)$$

where  $p(x_0, k)$  is the prior and we set to a flat prior and  $p(x | x_0, k)$  is the likelihood function which is defined by,

$$p(\bar{x} | x_0, k) = \prod_{j=0}^n \frac{n!}{k!(n-k)!} s(x_j | x_0, k)^k (1 - s(x_j | x_0, k))^{n-k}. \quad (3.43)$$

To plot the efficiency curves and lower and upper error bounds, we sample Eq. 3.42 using MCMC and then take the mean and the 5th and 95th percentiles respectively for each point in SNR or depth and plot these. Fig. 3.6a and 3.6c then show the efficiency curves for the analyses plotted against the signals optimal SNR and depth respectively. The parameters of the search and their optimised values are shown in Tab. 3.1. Where we set the prior on the line model to 0 as this part is irrelevant to this search due to the lack of lines in the data.

From this we can determine that in Gaussian noise without gaps, the Viterbi algorithm can detect to an SNR of  $\sim 60$  and a depth of  $\sim 33 \text{ Hz}^{-1/2}$  with 95% efficiency at a 1% false alarm.

Fig. 3.6b and 3.6d, show the RMS of the difference between the injected signal track and the track found by Viterbi for SNR and sensitivity depth respectively. This shows that at SNR of 60, where we are detecting signals with a 95% efficiency, the signals have a mean RMS of  $\sim 2$  frequency bins. Here one bin width is 0.00056 Hz therefore, we have and RMS of  $\sim 0.0012$  Hz.

### 3.10.2 S6 injections into Gaussian noise with gaps

In the second test, we attempt to more closely mirror the S6 MDC [48] in two stages. The first uses the same injection method as Sec. 3.10.1 however, removes the SFTs where there are



Table 3.1: Table shows the ranges of the search parameters and their optimised values for injections into gapless Gaussian noise, Gaussian noise with gaps and the S6 MDC. For gapless Gaussian noise and Gaussian noise with gaps, there are 10 parameter values spaced linearly between the limits. For the S6 MDC the parameters,  $\tau$ ,  $w_L$  and  $w_S$  were distributed in log space between the limits and  $p(M_L)/p(M_N)$  is distributed uniformly.

	$\tau$	$w_S$	$w_L$	$p(M_L)/p(M_N)$
<b>Gapless Gaussian</b>				
limits	[1.0,1.3]	[0.1,5.0]	None	0.0
optimised	1.1	2.06	None	0.0
<b>Gaussian with gaps</b>				
limits	[1.0,1.3]	[0.1,5.0]	None	0.0
optimised	1.1	2.06	None	0.0
<b>S6 MDC</b>				
limits	[1.0,1.1]	[0.1,5.0]	[0.1,6.0]	[0.0,1.0]
optimised	1.00000001	4.0	5.0	0.0387

gaps in S6. The second uses the same injection method again including gaps, however, uses a different value for the noise floor for each SFT, this is calculated for each band and SFT from S6 data.

Both detectors in S6 had a duty cycle of  $\sim 50\%$  [75], which means that there are sections of time where there is no data in either one or both detectors. In the sections where one detector is observing but the other is not, the multi detector statistic will not behave correctly as it only has access to data from a single detector. In these sections we switch from using the multi-detector statistic to the single-detector statistic using the same parameters, these are both defined in defined in Sec. 3.8.

The process of removing sub-bands and generating efficiency curves is the same as in Sec. 3.10.1.

We set a 1% false alarm rate and generate an efficiency curve for SNR and depth in Fig. 3.6a and Fig. 3.6c respectively. From these efficiency plots we can see to an SNR of  $\sim 72$  or a depth of  $\sim 13 \text{ Hz}^{-1/2}$  at a 95% confidence with a false alarm of 1%.

The parameters of the search which were optimised and their optimised values are shown in Tab. 3.1.

In Fig. 3.6b and 3.6d show the RMS of the difference between the injected signal track and the track found by Viterbi for SNR and sensitivity depth respectively. This shows that at SNR of 72, where we are detecting signals with a 95% efficiency, the signals have a mean RMS of  $\sim 10$  frequency bins (0.0056 Hz).

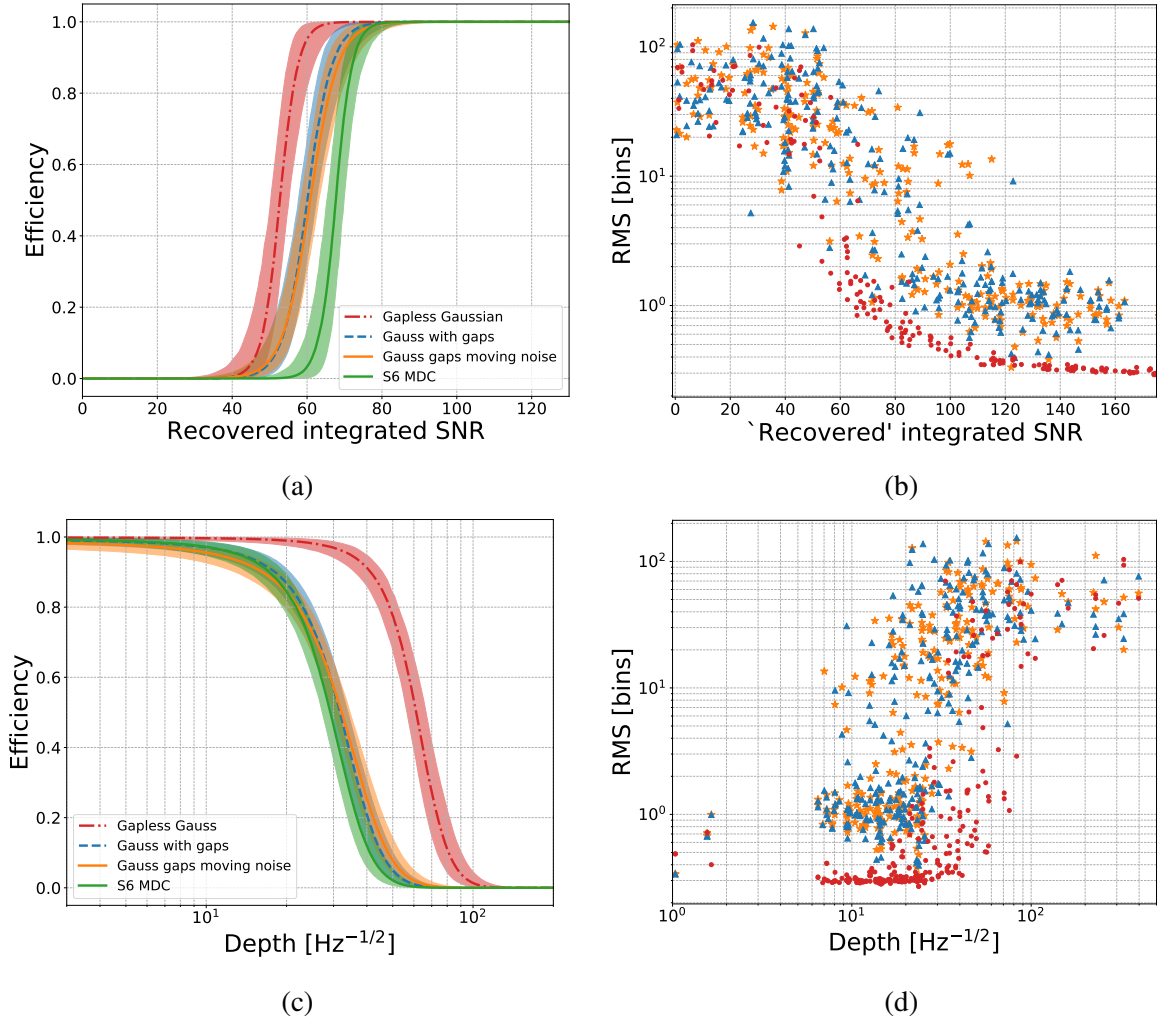


Figure 3.6: Panels 3.6a and 3.6c show the detection efficiency as a function of SNR and depth respectively. Here SNR is the the integrated SNR which we would expect to recover from the available data. The four curves refer to injections into gapless Gaussian noise (red), Gaussian noise with gaps in data, where the noise floor is either fixed (blue-dashed) or it is moving with time (orange) in the same way as the S6 MDC and injections into real data i.e., the S6 MDC. In the gapless Gaussian noise case, the recovered integrated SNR refers to the SNR the injection would have if it had the same amount of data as in the cases with gaps. The curves are made by fitting a sigmoid Eq. 3.39 to binomial detection data with a 1% false alarm rate, as explained in Sec. 3.10.1, the error bounds are the 5% and 95% intervals. At 95% efficiency and a 1% false alarm rate, this shows we can detect to an SNR of  $\sim 60$  and a sensitivity depth of  $\sim 34 \text{ Hz}^{-1/2}$  for gapless Gaussian noise and an SNR of  $\sim 69$  and  $72$  and a sensitivity depth of  $\sim 13 \text{ Hz}^{-1/2}$  and  $\sim 10 \text{ Hz}^{-1/2}$  for the Gaussian with gaps case with fixed noise floor and moving noise floor respectively. For the S6 MDC we can detect an SNR of  $\sim 74$  and a sensitivity depth of  $\sim 13 \text{ Hz}^{-1/2}$ . Panels 3.6b and 3.6d show the RMS of the difference between the injected signal track and the track found by SOAP as a function of SNR and sensitivity depth respectively. This is shown in units of bins where each bin is  $0.00056 \text{ Hz}$  wide.



### 3.10.3 Tests on the S6 MDC

For a more direct comparison to other CW searches and to see how the algorithm performs with real data, we test the two detector SOAP algorithm using the S6 MDC. We focus this search on the 100-200 Hz band, there are two main reasons for this, one being that this is LIGO's most sensitive band and the other is that for much higher frequencies the signal will drift over larger frequency ranges, therefore, our SFT length will have to be changed. Here the 1800 s SFTs are split as in Sec. 3.10, whereafter normalisation, the data is split into 0.1 Hz wide sub-bands overlapping by 0.05 Hz.

The two detector SOAP algorithm using the line-aware statistic in Sec. 3.8 is then run on each sub-band under the assumption that the detectors have the same sensitivity. For this search we have four parameters which we optimise, the ranges and optimised values are shown in Tab. 3.1.

As in Sec. 3.10.2, only the sub-bands which contained the entire frequency evolution of the signal were selected. Out of the 2000 sub-bands, 238 were removed due to the sub-band only containing part of the signal's frequency evolution. The main difference between the analysis for Gaussian noise and real data is that the real data is contaminated with instrumental lines. This means that whilst the techniques described in Sec. 3.8 reduce the number of contaminated bands with a high statistic value, there are still instrumental lines which are coincident between the detectors and which could not be removed with these techniques. Within the data there are a large number of lines at integer Hertz, which are seen in coincidence between the two detectors, these are thought to originate from digital electronics [76]. Therefore the frequency bins  $\pm 1$  bin of each integer frequency in Hertz were removed and filled with the expectation value of the noise. To remove instrumental effects at other frequencies, the sub-bands which gave values of our statistic above a chosen threshold were investigated by eye. In this case 344 sub-bands were investigated, and any which were contaminated were vetoed. From these 344 sub-bands, 193 were removed from the analysis. The predominant feature in the bands which were removed were broad spectral features which lasted the whole run. Therefore, out of the 2000 sub-bands which are searched over, a total number of 431 sub-bands were removed.

The process to calculate the efficiency curves is the same as in Sec. 3.10.2 and 3.10.1.

Fig. 3.6c and Fig. 3.6a show the efficiency curves for SNR and depth respectively. These show that we can detect at SNR of  $\sim 74$  and a sensitivity depth of  $\sim 13 \text{ Hz}^{-1/2}$  with an efficiency of 95% at a false alarm of 1%. These results can then be compared to other searches in the S6 MDC comparison paper [48]. Whilst we only search in the 100 - 200 Hz range, the closest comparison in [48] is the test in the 40 - 500 Hz range, such as in Fig. 4 in [48]. Here our algorithm sits roughly in the middle of all other searches in terms of sensitivity.

### 3.11 Optimisation of Line-aware statistic.

For the above searches we used optimised versions on the line aware statistic, however, we have yet to explain how this was optimised. The aim is to find the best parameters for any given search; the four parameters are  $\tau, p_S(\lambda), p_L(\lambda)$  and  $p(M_L)/p(M_G)$ . We find the optimum values empirically by running the entire search for each parameter value that need to be tested. This is possible as the search is relatively fast, this will be explained in Sec .3.11.5. The line aware statistic is time consuming to calculate, therefore, to reduce the computational time, it is pre-calculated and places into lookup tables such that it is calculated once and called many times. These lookup tables were calculated for values of FFT power  $F$  in the range 1 to 400 in each of the detectors as shown in Fig. 3.3. For each parameter the ranges in which they were optimised were chosen based on expected SNRs of each of the injections.

We can use a measure of the sensitivity of that search and pick the lookup table which gives the highest sensitivity. We measure the sensitivity by taking the value of SNR which is at 90% efficiency at 1% false alarm.

#### 3.11.1 S6 injections into gapless Gaussian noise and Gaussian noise with gaps

For injections into Gaussian noise, we know that there are no instrumental lines, therefore, we do not need to optimise over the ‘lines’ part of the statistic and can set the parameter  $p(M_L)/p(M_G)$  to zero which renders the parameter  $p_L(\lambda)$  redundant. This then reduces the complexity of the problem by leaving us with only two parameters to optimise over,  $\tau$  and  $p_S(\lambda)$ . Whilst this optimisation was partially done in Sec. 3.10, with the result in Tab. 3.1, this is repeated more completely here. The parameters were optimised in the range shown in Tab. ??.

Fig. 3.7 shows the values for the SNR at 80% efficiency for each of the parameter values which were tested, the optimum which is ..... **JOE: run test and fill in**

#### 3.11.2 S6 MDC injections

As the S6 MDC data-set is real detector data, there are many examples of instrumental lines. This is where we expect the line-aware statistic to have the greatest effect in rejecting lines. Here the all of the four parameters are optimised over in the ranges described in Tab. 3.1. This greatly increases the number of lookup tables which need to be generated and therefore the number of time the search needs to be run. This optimisation is then carried over to all future searches. Fig .3.8 shows the slices through parameters space for each of the parameters where the values are the SNRs at 90% efficiency.

Now the parameters have been optimised for both the Gaussian noise case and the real S6 data case, we can test the effect of the statistic when used on a band which has an instrumental

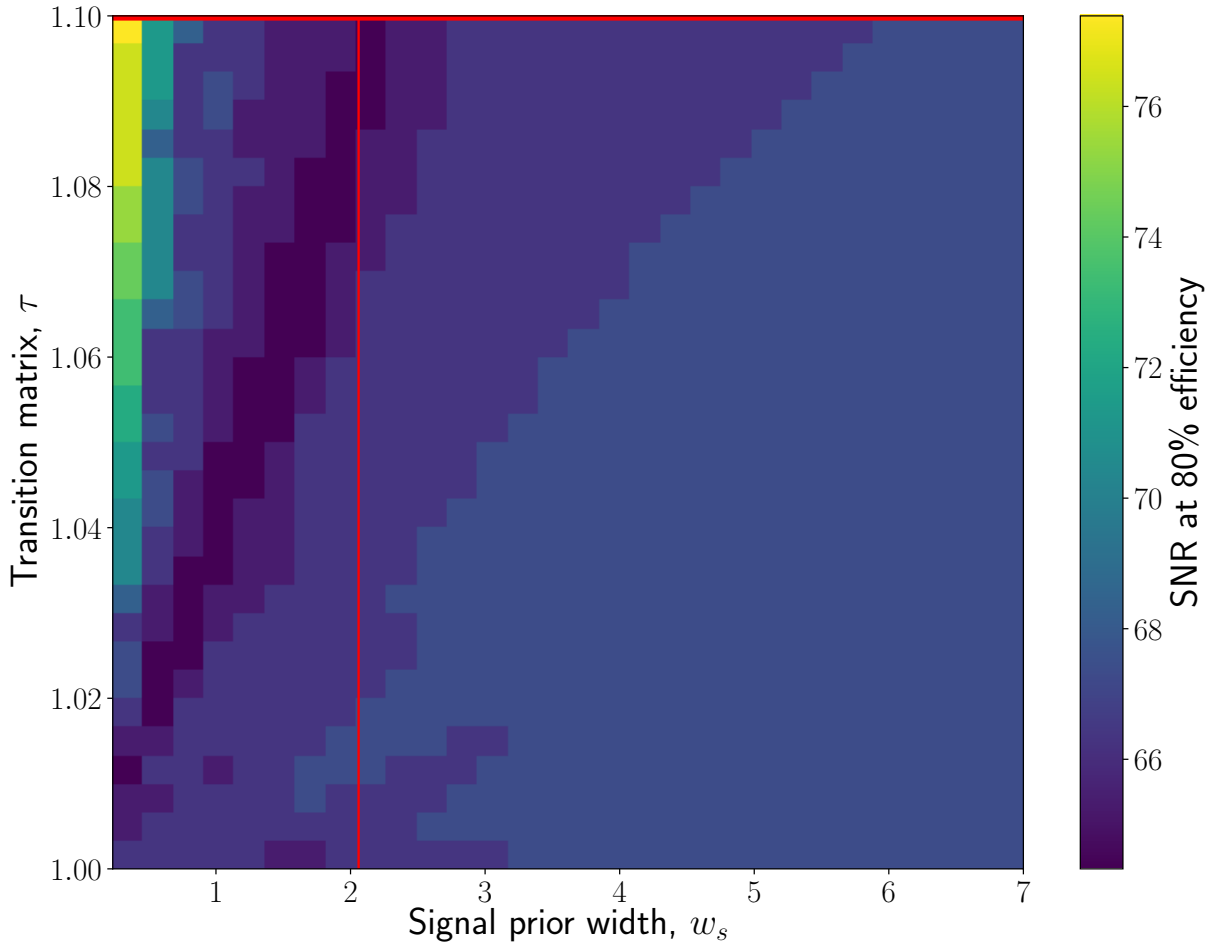


Figure 3.7: In Gaussian noise the transition matrix parameter  $\tau$  and the width of the prior on the signal case  $w_s$  were optimised. The key part to remember when reading this plot is that the lower the value of SNR the better the search has performed. Therefore darker blue areas are when the search performed better. This map shows that there is a line in parameter space where the search performed best. Also in Gaussian noise, the search is not that sensitive to the choice of parameter. The red lines on here shows the parameters used in the searches in this section.

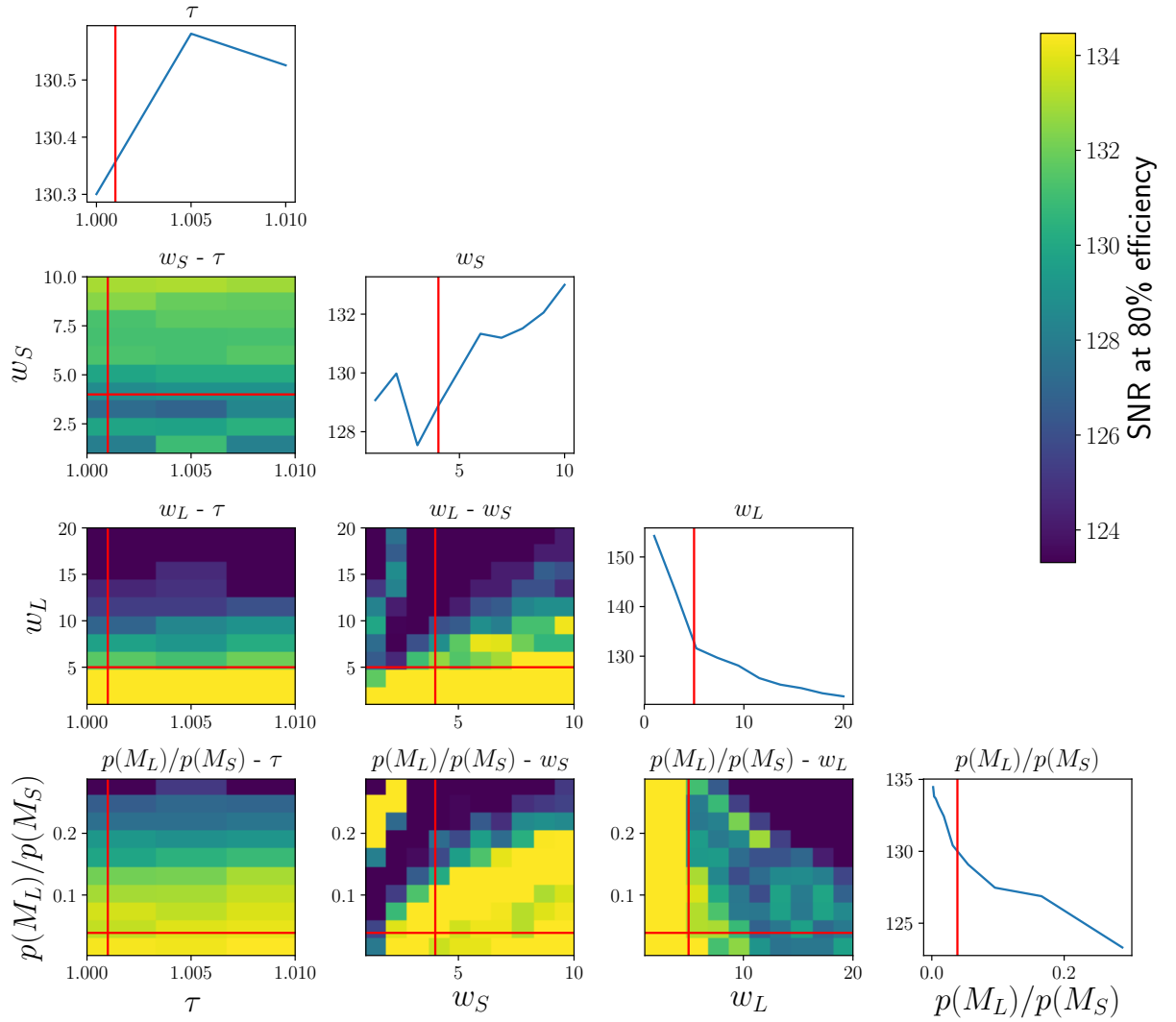


Figure 3.8: When using real S6 data, all four parameters of the search were optimised over on simulations in real data. The plot above shows the SNR at 80% efficiency for each of the parameters where the ranges are in Tab. ???. Lower values of SNR mean the search is performing better. The red lines show the parameters used in the searches in this section and the sections that follow. Whilst this does not seem optimal, the search does not underperform much when a poor choice of parameters is used.

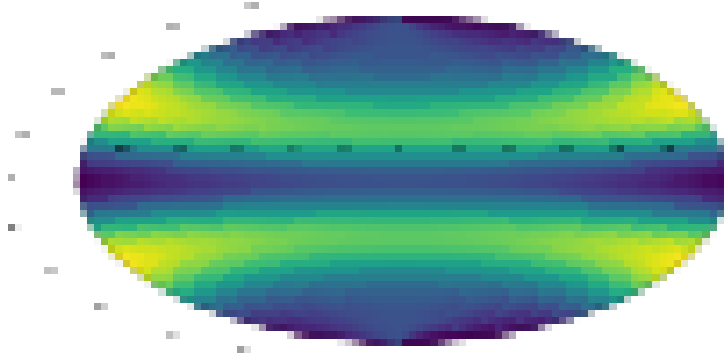


Figure 3.9: This figure shows an example of when the line aware statistic is used compared to a version when it is not.

line injected. In Fig. 3.9 an example a CW signal injected into Gaussian noise in both detectors. This also shows a simple example of an instrumental line injected into one detector. In the Gaussian noise optimised case the search is looking for high power in both detectors which the strong instrumental line satisfied when the astrophysical signal is weak. The S6 optimised statistic looks for more consistent SNR in each of the detectors. The two different lines show the statistic optimised on the Gaussian noise injections and the S6 MDC injections. This demonstrates how the ‘line aware’ statistic improves the robustness of the algorithm against non astrophysical signals.

### 3.11.3 Sensitivity with frequency

All of the above tests, the search was conducted in the range from 100-200 Hz. This was chosen to be within the most sensitive band of LIGO as this is where a signal is most likely to be discovered. However, signals can appear at much higher frequencies also. Therefore, it is important to see how the sensitivity of the search varies with the frequency.

For this test we simulated CW signals in Gaussian noise with no gaps in data. The injections used the same source parameters as in the a6 MDC [48] and the tests above. This has the exception that the integrated recovered SNR of the signal is sampled uniformly between 50 and 500. These injections were then made at frequencies of 100,250,500,750,1000,1500 and 2000 Hz, where the band width is 2 Hz. i.e. the simulations were in frequency bands 100-102 Hz, 250-252 Hz etc. The setup of the search was the same as in the above sections. Here each sub-band is 0.1 Hz wide, and the parameters of the SOAP search were as in Tab. 3.1.

Fig.?? the resulting efficiency curves from each of these tests. This is for a 1% false alarm rate, which means that 1% of sub-bands which contained no injection crossed the detection threshold. This plot shows how the sensitivity of the search drops as the frequency increases. This is perhaps unfair to the algorithm as we used the setup of the search which has been op-

timised for the range 100-200 Hz. Optimising the search means, choosing the parameters of SOAP, the key parameter which will affect this is the transition matrix. As the simulated signals frequency is increased, the scale of the Doppler modulation will also increase. This means that at higher frequencies the signal is more likely to jump more than a single frequency bin. The current setup of the search does not allow this size of jump, therefore, would struggle to identify this type of track. The other main factor which will decrease this sensitivity is the sub-band width of 0.1 Hz. Similarly to before, as the frequency increases the scale of the Doppler modulation will increase as in Eq. 3.37. For example at 1000 Hz, the Doppler shift is  $\sim 0.1$  Hz, the signal is then more likely to not be fully contained within a frequency band. Therefore, the search can not accumulate all of the injected SNR

Whilst these parameters of the search can be changed to depend on frequency, this was meant as a small test to see how sensitivity varies with frequency.

### 3.11.4 Searching for non CW sources

### 3.11.5 Computational cost

One of the main strengths of this search is the drastically reduced computational cost when compared to other current CW searches. The scaling of the computing cost can be estimated for a single detector by looking at the number of calculations that need to be made. The number of calculations for a single detector search,  $N_{\text{calcs}}^{(1)}$  is,

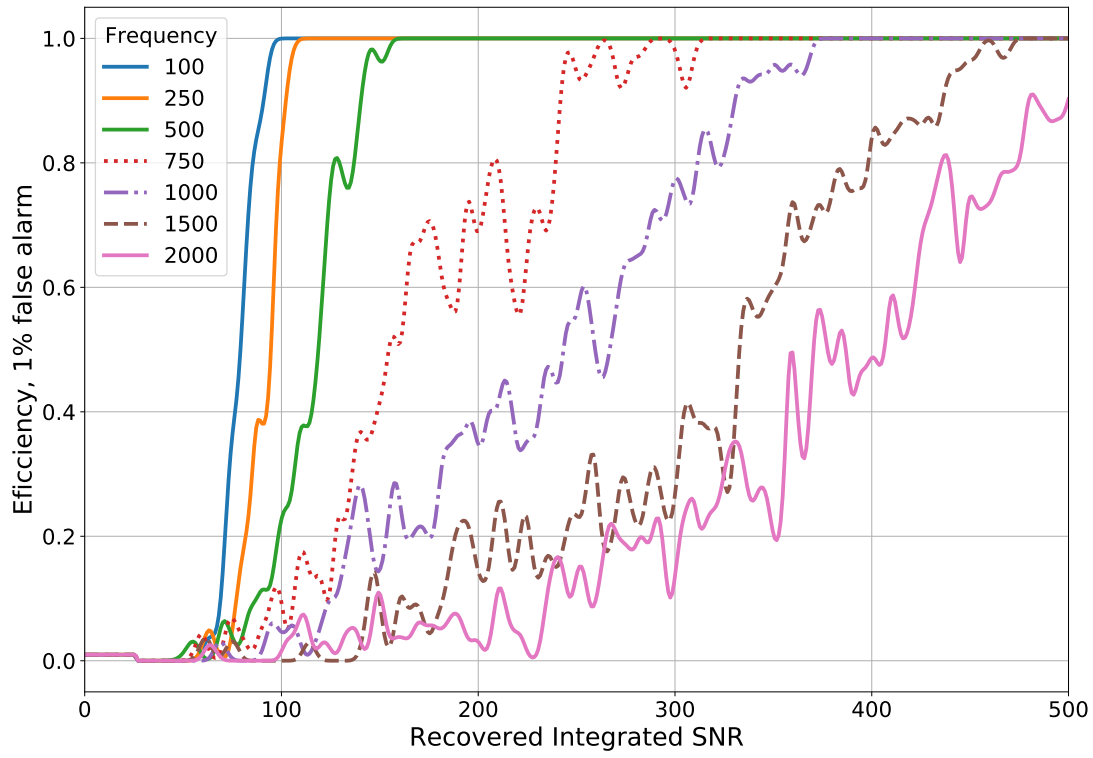
$$N_{\text{calcs}}^{(1)} = n_1^m NM, \quad (3.44)$$

where  $n_1$  is the size of the transition matrix,  $N$  is the number of SFTs,  $M$  is the number of frequency bins and  $m$  is the amount of memory described in Sec. 3.6. Where the computing cost scales linearly with the number of frequency bins and SFTs. In the following test we ignore ‘memory’ and look at the time taken for the single detector search where the time taken to read and save data is ignored. Here the data is the same size as the S6 MDC for a single detector search and the search is over a 0.1 Hz band, where we set  $n_1 = 3$ . This test, and the following test, was run locally on a MacBook Air with a 1.3 GHz Intel Core i5 processor. We can then write the time taken,  $T$ , as,

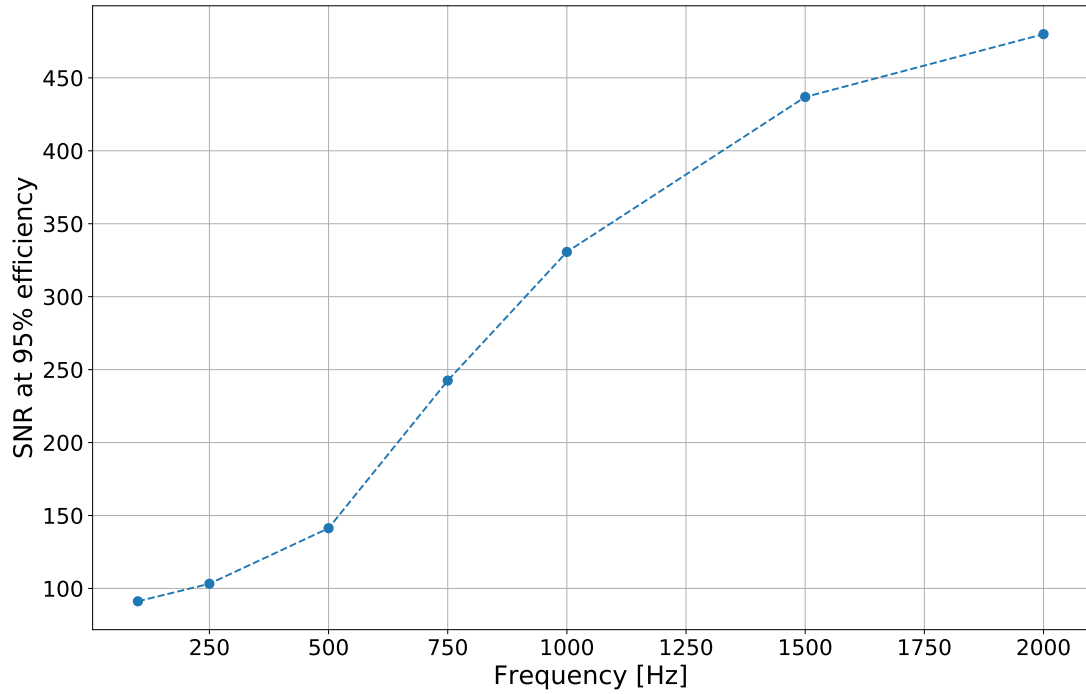
$$T = 0.56 \text{ sec} \left( \frac{N}{22538} \right) \left( \frac{M}{180} \right) \left( \frac{N_{\text{bands}}}{1} \right), \quad (3.45)$$

where  $N_{\text{bands}}$  is the number of different frequency bands. For the multiple,  $Q$ , detector case, we can then generalise Eq. 3.44 and write the number of calculations  $N_{\text{calcs}}^{(Q)}$  as,

$$N_{\text{calcs}}^{(Q)} = NM n_1^m \prod_{q=1}^Q n_{q+1}, \quad (3.46)$$



(a)



(b)

Figure 3.10: The sensitivity of the SOAP search in this configuration decreases as the frequency of the pulsar increases. This setup of data for the search however, was optimised for the 100-200 frequency band and can be changed for different frequencies. 3.10a shows the efficiency curves with 1% false alarm rate for each frequency. 3.10b shows the values from the efficiency curves at 90% efficiency.

where  $n_1$  is the first dimension of the transition matrix,  $Q$  is the number of detectors and  $n_{q+1}$  is the size of the transition matrix element which refers to detector  $q$ . For our tests we set  $n_1 = n_{q+1} = 3$  and use 2 detectors i.e.,  $Q = 2$  which each have the same size data as the previous test. The actual time taken to run however, depends on the version of the algorithm which is run. For example, including the line aware statistic slows the search slightly. For the two detector where two SFT powers are summed,

$$T = 1.35\text{s} \left( \frac{N}{22538} \right) \left( \frac{M}{180} \right) \left( \frac{N_{\text{bands}}}{1} \right). \quad (3.47)$$

The same search now including the line aware statistic, which is implemented using a lookup table, changes this to,

$$T = 25.7\text{s} \left( \frac{N}{22538} \right) \left( \frac{M}{180} \right) \left( \frac{N_{\text{bands}}}{1} \right). \quad (3.48)$$

Other searches, excluding Einstein@home which takes on the order of months to run ( $> 100$  million core-hours [48]), take 1 – 10 million core-hours [48]. This search should take  $\mathcal{O}(10^3)$  core-hours.

## 3.12 Discussion

In this paper we describe an application of the Viterbi algorithm, called SOAP, to search for continuous sources of gravitational waves. This paper outlines the method and derives the statistics behind the method in a consistent Bayesian formalism. It then presents the results from the first set of tests of sensitivity for the SOAP algorithm on three separate datasets.

We tested SOAP on a set of fake isolated pulsar signals in the 100 – 200 Hz range, based on 1800s SFTs summed over one day. The three datasets that included these signals comprised continuous Gaussian noise, Gaussian noise but with temporal gaps corresponding to LIGO dead times in the S6 data run, and real data, i.e., the S6 MDC. Although a major attraction of SOAP is its sensitivity to a wide range of signal types, in the tests above it was optimised to detect isolated pulsar signals below 100 Hz with low spin-down to offer a comparison with other CW searches. From these tests, by setting a 95% efficiency and a false alarm of 1%, we found that in the case of continuous Gaussian data we could detect a signal with an optimal SNR of  $\sim 60$  and a depth of  $\sim 33 \text{ Hz}^{-1/2}$  with an RMS of the difference between the injected and Viterbi track being  $\sim 2$  frequency bins (0.0012 Hz). When gaps were introduced into the data to simulate S6 we could detect a signal with an SNR  $\sim 72$  and a depth of  $\sim 10 \text{ Hz}^{-1/2}$ , with an RMS of  $\sim 10$  bins (0.0056 Hz). The drop in sensitivity here is simply because there is  $\sim 50\%$  less data compared to the previous case. Finally, in the S6 MDC we could detect a signal with an SNR  $\sim 74$  and a depth of  $\sim 13 \text{ Hz}^{-1/2}$ . These real data contain non-Gaussian artefacts such as instrumental lines and this causes a further drop in sensitivity. Whilst not a full comparison to other searches in the



S6 MDC [48], as we only tested on a subset of the bands, this search has a sensitivity which is comparable to some other CW searches, however offers a massive increase in speed.

We chose the specific frequency band to search over as the data which we used, i.e., the summed data, becomes less effective at frequencies much higher than 200 Hz, and using the parameters of our simulations, signals can spread over many frequency bins in a day, reducing sensitivity further, however this can be mitigated by using shorter SFTs or performing their summation over 12 (rather than 24) hours.

The methods described in this paper present a basic approach for gravitational-wave signal searches using SOAP. However there are several further developments that could increase its sensitivity. Some of these are outlined below:

One of the main features which reduces the sensitivity of the search is non-gaussianities within the data, namely instrumental lines. Although we have a statistic which penalises these features, in some cases it will also penalise a strong signal. For example, when the amplitude of the noise floor is high for one detector or the duty factor is lower, the signal will appear more like an instrumental line to this statistic. We hope to improve the search statistic in the future by searching for consistent amplitudes as opposed to consistent SNR, i.e, the statistic will take the amplitude of the noise floor and the duty factor into account.

One variation of this method which has been described in this paper is ‘memory’, which is where the tracks jump in frequency is determined by the previous  $n$  jumps. This has yet to be fully tested, however, we expect that this will increase our sensitivity to signals where have a better idea of their frequency evolution. This however, comes at a cost in computational time which we can estimate given Eq. 3.46 in Sec. 3.11.5.

Further additions to the search include using the Fourier transform of the SFT power along the Viterbi track as a detection statistic. If the Viterbi track follows that from an astrophysical signal, then we should see the effects of the antenna pattern in this Fourier transform as a peak at half a sidereal day. If the track follows something which is not astrophysical then this should not be seen this peak in this Fourier transform. This only applies to the search directly on the SFTs not the summed data, as the antenna pattern variations will have been averaged out in the summing.

As well as searching for astrophysical signals, SOAP can also be used to search for and identify instrumental lines. Here we use single detector data, or multiple channels from a single detector, to identify quasi-monochromatic features on the data for further study.

Whilst this paper presents initial tests on sensitivity, further tests will be needed for a full comparison to other CW search methods. This search, however, aims to look for signals which may not follow the standard frequency evolution and is intended to return potentially interesting candidates for a more sensitive followup.

# Chapter 4

## Machine learning for continuous wave searches

Machine learning is a term which was used by Arthur Samuel in 1959. He described it as a "Field of study that gives computers the ability to learn without being explicitly programmed" [1]. This can be thought of as a subset of artificial intelligence.

With the development in computing in recent years, including GPUs and the languages used to program with them, machine learning has become more accessible.

### 4.1 Introduction

Gravitational wave detectors such as LIGO [38, 10] and VIRGO [11, 39] search for a number of different targets. Some targets such as CBCs have been observed [7, 5, 4], however, other primary sources such as CWs are yet to be observed. CWs are well modelled quasi-sinusoidal signals with a duration much longer than observing times of detectors. The source of these signals is thought to be rapidly rotating neutron stars which can emit GW if there is some asymmetry around its rotation axis. This can be caused by various mechanisms as described in [62]. These signals have small amplitude, which if detected will be below the noise PSD of the detector. Therefore, sensitive search algorithms are needed to find the signals. These algorithms generally fall into three categories: Targeted, directed, and all-sky searches, listed in order of how much is known a priori about the source from EM observations.

In targeted searches the sky position, frequency, and its derivatives are assumed to be well known, in directed searches only the sky position is known and in all-sky searches the sky position and frequency of the source is unknown. The most sensitive of these are targeted searches which use coherent matched filtering [46, 45]. These use template waveforms which are generated using the information already known about the source, then correlated this with the data. Directed and all-sky searches have a much broader parameter space to search, therefore, many templates are needed to sufficiently cover the parameter space. Using the coherent matched filter

for broader parameter space searches becomes unfeasible due to the amount of computing time that is needed. This led to the development of semi-coherent searches where the data is divided up into smaller segments which can be analysed separately and then the results can be recombined incoherently using various methods [67, 66]. Semi-coherent searches result in a trade off between sensitivity and computing time.

The analysis here is presented mainly as an addition to an existing semi-coherent search algorithm titled SOAP [77]. This is a fast and largely un-modelled search which finds tracks of high FFT power in time-frequency spectrograms. When applied to multiple detectors using a line-aware statistic, SOAP looks for frequency bins which have both a high power and are similar in each detector. This means that at a given frequency at a given time, SOAP will penalise frequencies where the FFT power is largely different in each detector. The algorithmic details summarised in Sec. 3.

One effect which limits the sensitivity of SOAP and many other GW searches is the effect of ‘instrumental lines’. These can be anything from long duration fixed frequency or wandering lines to shorter duration fixed frequency transients. There are certain types of instrumental line which the SOAP search can struggle to distinguish from an astrophysical signal even with the development of a ‘line aware’ statistic in [77]. Currently the method used to reduce the effect of these lines is to manually look at the SOAP output and the spectrograms for each sub-band to determine whether the sub-band is contaminated by instrumental effects. This process is slow, requires a lot of human input and is subject to human error. When the search runs over a larger bandwidth, it will no longer be practical to look through all bands.

We aim to automate how the search deals with instrumental lines by using convolutional neural networks (CNNs). These have been used extensively in image classification and we explain this in more detail in Sec. 4. CNNs have already been shown to detect gravitational wave signals from CBCs in [78, 79, 80] and other deep learning techniques have been used in searching for CW signals in [81].

In Sec. 3 we will summarise the basics of how the SOAP search works. In Sec. 4 we explain how CNNs operate followed by how we generate data to train the CNN in Sec. ???. We then describe the entire search from raw data to results in Sec. ??? and finally in Sec. 3.10 we show the results from this search and compare to similar analyses.

## 4.2 Neural networks

Throughout this section I will summarise one machine learning technique which are known as Neural networks. Neural networks, as the name may suggest, was developed as a way for a computer to mimic the neurons in the brain. To understand why this would be useful, I will give short example which will be followed through this section. A good example used extensively in explanations of neural networks is the ability to identify hand written digits. This is a simple

task which a brain can complete with ease. However, writing a traditional algorithm to perform this same task is very difficult. The algorithm would have to identify a particular shape, which has a huge amount of variation. Neural networks offer a way to deal with this problem as they can be trained on large datasets. This is similar to how a human brain is ‘trained’. In a lifetime or a brain many examples of different symbols are seen and each time a new one is seen the brain ‘updates’ itself based on what is observed. This process is essentially replicated for a neural network, where the algorithm can be updated such that it can correctly identify each digit.

Neural networks have many parameters which can be modified or ‘trained’. These parameters and their application are grouped into objects called neurons, many of these neurons are used to build a neural network.—

### 4.2.1 Neurons

Neurons are the building blocks of any neural network. They perform simple operations on any number of input values and then output a single value. The output  $o$  of a neuron is defined by the equation,

$$o = f \left( b + \sum_{i=1}^N w_i x_i \right), \quad (4.1)$$

where  $b$  is the bias,  $x$  is the input,  $w$  is the weights,  $f$  is the activation function and  $o$  is the output. Here the input  $x$  represents either the data which is input, i.e the pixels of the image which contains the digit in the example above, or the output of another neuron. The weights  $w$  then represents how important each of those data points are to this problem, or specifically this neuron. The bias  $b$  is then just an extra factor which can shift the data by a fixed value. The activation function  $f$  is then a function which can have many forms, in the simplest case in a neuron known as a ‘perceptron’, it provides a cut where any value above a given threshold is 1 and any below is 0. However, there are many different types of activation functions which can be applied to different situations. This will be explained in more detail in Sec. 4.2.3.

In the example in Fig. 4.1 I have shows a neuron which has 4 input variables, or 4 input data points. When a network is trained, or when it learns, the weights applied to each of the inputs and the bias are updated to better represent the input data. This training procedure is explained in more detail in Sec. 4.2.4 Many neurons are then used in combination with each other to develop a neural network which can be applied to more complex problems.

### 4.2.2 Network structure

The structure of a neural network is defined by the user and there is no set way to design a network. However, the general layout of a neural network is defined by structures called layers, sometimes known as fully connected layers. These are rows of  $N$  neurons which all take the

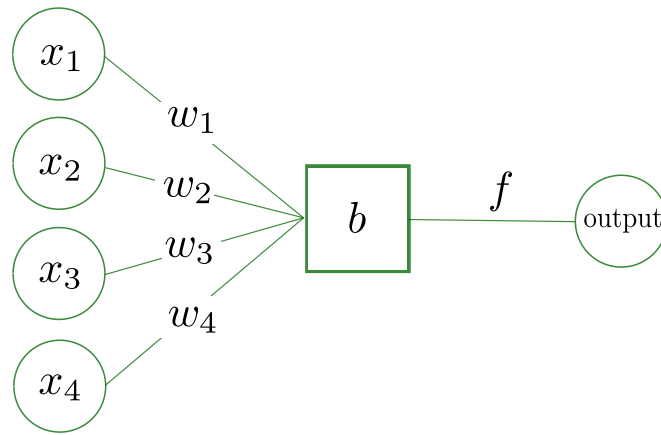


Figure 4.1: Basic neuron

same input such that there is  $N$  output values. An example of a simple neural network is shown in Fig. 4.2. The first layer is the input layer, this is just the data points from an input example. In the example of hand drawn digits, this would be the pixels from the image of the digit. The final layer represents the information that you intend the network to extract from the input data. In the hand drawn digit example, this could have 10 output values corresponding to each digit 0-9. Each of these outputs is then a value which is related to the probability of that digit being present in the image.

When designing a network, the user will have a defined input layer size from the data and a set number for the output layer which represents, for a classification example, the number of output classes. The number of hidden layers and the number of neurons in those hidden layers can be arbitrarily changed. In general if the data contains more complex information the size or complexity of the network will need to be increased for it to be able to extract the information. If there is a small number of training examples and a large and complex network, it may be able to learn the input data set as opposed to the general information that they represent.

### 4.2.3 Activation functions

The activation function is how the output of a neuron is transformed. The most simple activation function is a cut as described in Sec. 4.2.1, however, this type of activation does not perform well. Activation functions are generally based on a few properties. The activation function is generally non-linear, this allows networks with multiple layers to be used to approximate a function. A linear activation function means that any number of layers in a network is equivalent to a single layer network. Another property which is desired in activation function is that it is continuously differentiable. This is to allow algorithms such as gradient descent to optimise the network. The functions are found to perform better if they are monotonic and smooth. There are many choices when defining this in the network, some of the available options are shown in Fig. 4.3. One of

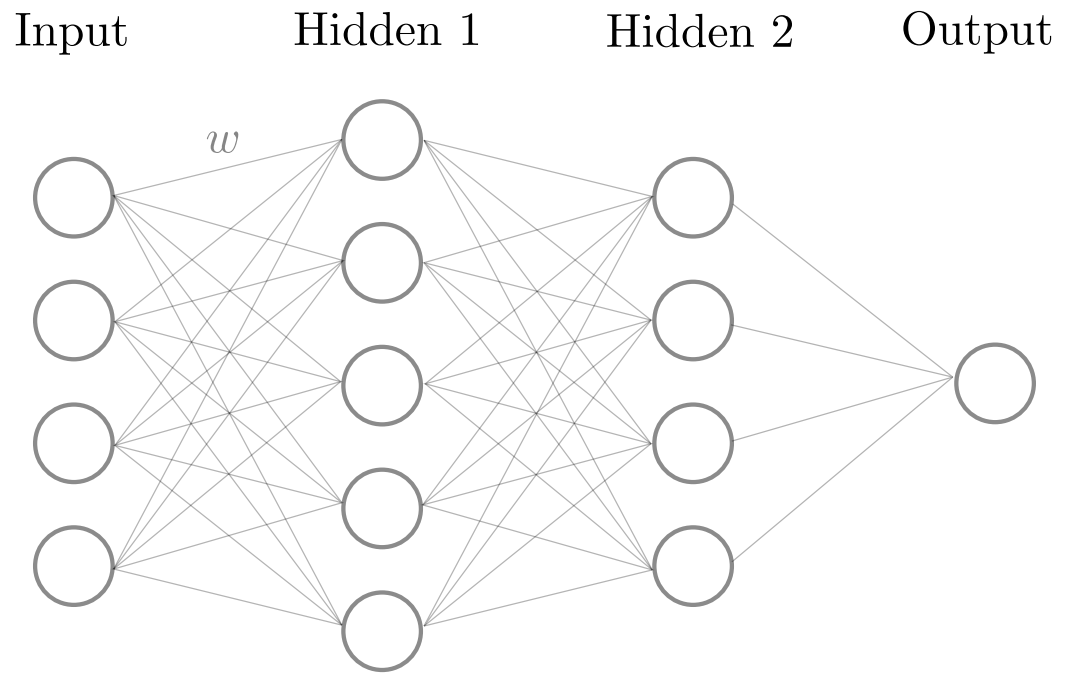


Figure 4.2: A neural network is structured with layers. This consists of the input layer which is the data which you would like to analyse. Then this passes to a number of ‘hidden’ layers, in the above diagram there are two. These are just all layers between the input and the output. The output layer is then the desired output, above there is a single neuron as output, so this could perhaps classify the input to a value between 0 and 1. Every neuron in a layer is connected to the output of all neurons in the previous layer.

the more commonly used activation function is the LeakyRELU function, this is explained in more detail in [82]. In the work that follows we use the Leaky RELU function and the sigmoid function.

#### 4.2.4 Training

For example, imagine you are playing darts and have the aim of trying to hit bullseye. After each dart that misses you would adjust where you aim the dart based on how far away from the bullseye it landed.

#### 4.2.5 Loss functions

When training a network as described in Sec. 4.2.4 one needs to find how far from the truth the result is. Depending on the problem, there are a number of ways to do this.

## 4.3 Convolutional Neural Networks

Convolutional neural networks are based on the same principles as the neural networks described above. The difference being that CNNs use additional types of layers which help to contain information of pixels relative to each other. In a standard neural network, it does not have any information on the location of any input value. For example, in an input image every pixel is flattened into a single dimensional array before being input to the network. CNNs try to retain some location information for each of the pixels. Whilst CNNs can be used for single dimensional inputs, they were generally developed for images. As well as retaining some location information, they reduce the size of the networks parameters.

### 4.3.1 Convolutional layers

Convolutional layers have some similarities to standard fully connected layers as described in Sec. 4.2.2. The main difference being how the weights are applied to the inputs. To show the difference I have an example of a simple input image which has 9 pixels, this is shown in Fig. ??.

A fully connected neural network would flatten this image and apply Eq. 4.1 to the inputs, such that the output is,

$$\begin{aligned}
 O_k &= f \left( \sum_i \sum_j I_{i,j} w_{i,j,k} + b_k \right) \\
 O_1 &= f (I_1 w_{1,1} + I_2 w_{2,1} + I_3 w_{3,1} + I_4 w_{4,1} + I_5 w_{5,1} + I_6 w_{6,1} + I_7 w_{7,1} + I_8 w_{8,1} + I_9 w_{9,1} + b_1) \\
 &\dots
 \end{aligned} \tag{4.2}$$

If the size of the convolutional filter was  $2 \times 2$ , then the convolutional layer only has 5 parameters to vary as opposed to 10 in this case. This would then use the equation,

$$O_{i,j} = \sum_m \sum_n K_{m,n} I_{i-m,j-n} \tag{4.3}$$

$$\begin{aligned}
 O_1 &= f (I_1 w_1 + I_2 w_2 + I_4 w_3 + I_5 w_4 + b) \\
 O_2 &= f (I_2 w_1 + I_3 w_2 + I_5 w_3 + I_6 w_4 + b) \\
 O_3 &= f (I_4 w_1 + I_5 w_2 + I_7 w_3 + I_8 w_4 + b) \\
 O_4 &= f (I_5 w_1 + I_6 w_2 + I_8 w_3 + I_9 w_4 + b).
 \end{aligned} \tag{4.4}$$

Whilst this is hard to picture mathematically, it can be easier seen as a  $2 \times 2$  filter which is convolved with the input image. Fig. ?? shows how the convolutional filter is applied. The output of a convolutional layer is then an image which is half the filter length less length in each

dimension. In practice, the input images are often padded with zeros such that the output image is the same size as the input image.

A convolutional layer has a number of different parameters which can be varied when setting up a model. Below I list each of the adaptable parameters and what they do.

**Filter size** The filter size is the size and shape of the convolutional filter. The filter does not have to be square, however must be less than the dimensions of the image.

**Number of filters** The number of filters can be any value. If you have  $K$  filter kernels, then the convolutional layer will output  $K$  filtered images.

**Activation function** The activation function is generally kept the same for each of the layers, however this can be set here. The different types have been explained in Sec. 4.2.3.

**Stride** The stride is when the convolution is applied to every other pixel. This method reduces the size of the output by the same factor. This is not used in the rest of this work.

The convolutional layers will reduce the number of parameters used in each network or model, which can speed up the training procedure. In a normal neural network the image is flattened, therefore, any information which related the location of a pixel to another is lost. Convolutions can keep hold of this information and look for similar patterns within an image.

### 4.3.2 Max pooling layers

Max pooling layers are designed to reduce the size of the problem whilst holding on to as much important information as possible. The idea of this layer is relatively simple, it reduces the image size by taking the maximum value in a region of a given size. For example, in Fig. ??, the image is reduced by a 2x2 max pooling layer.

### 4.3.3 CNN structure

CNNs are usually structured such that they can extract larger features from an input image, then the outputs from this are passed on to be classified. To extract the larger features any number of different convolutional layers and max-pooling layers used. Once this then outputs  $K$  images, the values are flattened into a one dimensional array and passed to a set of fully connected layers. This is essentially the same type of network described in Sec. 4.2. This can then classify the network into the chosen output.

CNNs are then trained in the same way as neural networks as described in Sec. 4.2.4.



## **4.4 CW search**

### **4.4.1 Network structure**

**Viterbi statistic**

**Viterbi Maps**

**Spectrograms**

**Viterbi map and spectrograms**

**Viterbi map and Viterbi statistic**

**Viterbi map, Viterbi statistic and spectrograms**

## **4.5 Network Visualisation**

Neural network are generally hard to visualise due to the large number of parameters in the network that have to be varied. However, there are methods which can be used to see how input data is affected by the network. This can be useful to see how the network performs when given certain types of data and gives some insight into how the networks work

In the examples above we use CNNs, the first few layers of this are build using convolutional filters. The filters weights should ideally correspond to the shape of the feature which one wants to extract from the image. Generally this is only useful to picture at the first layer as the network can make subsequent layer and representations quite abstract.

## **4.6 Data and simulations**

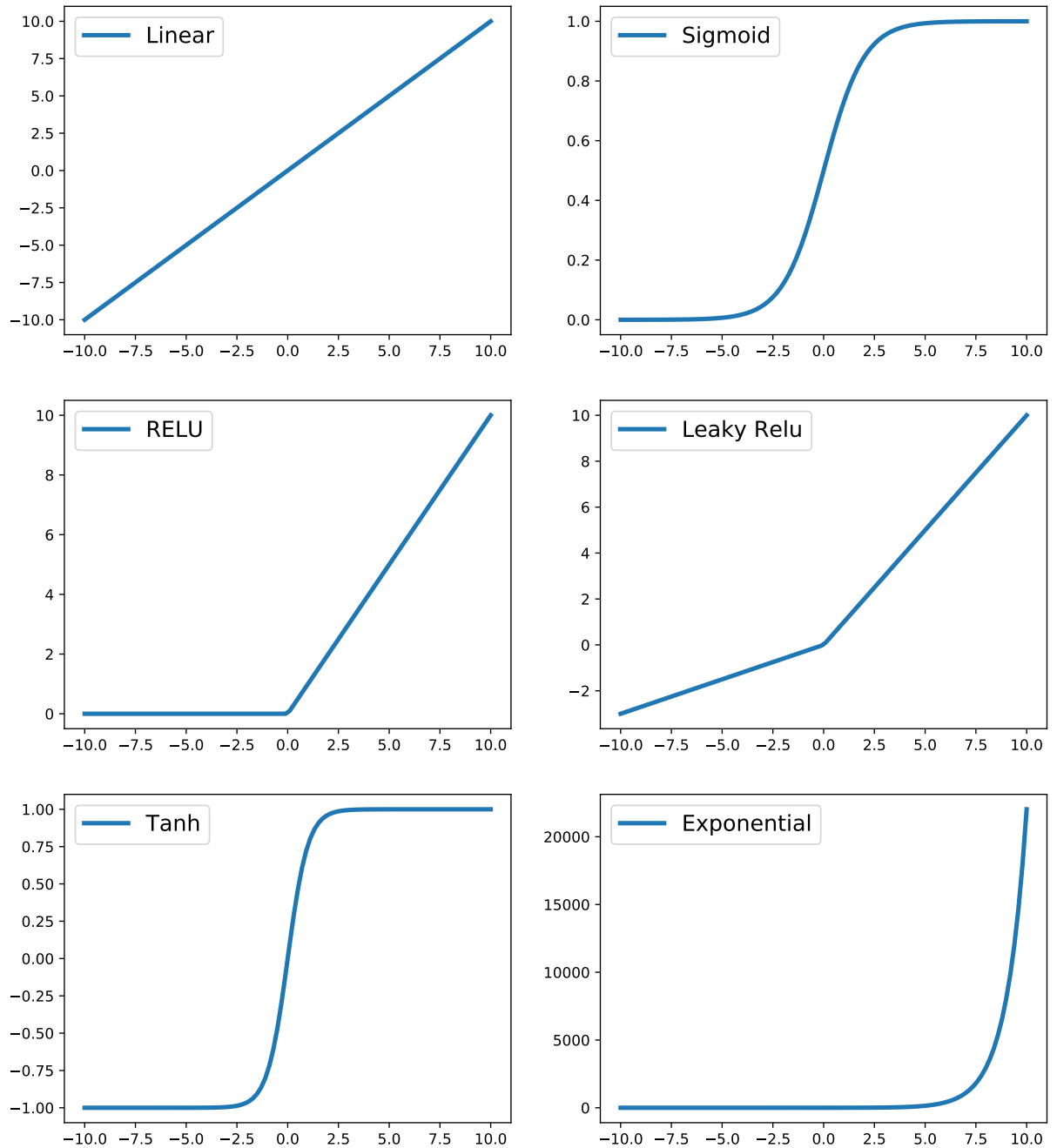
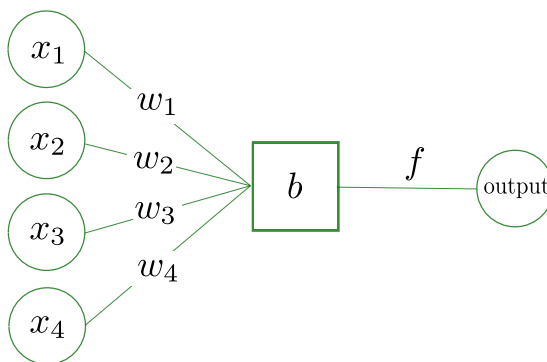
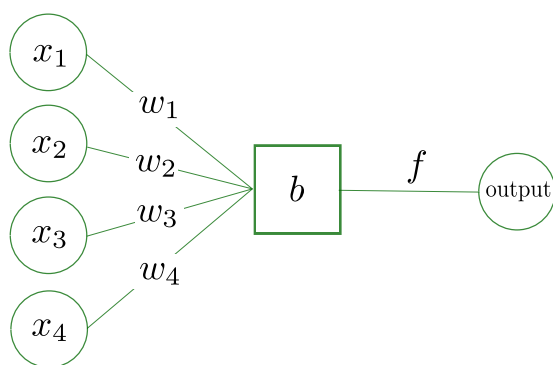


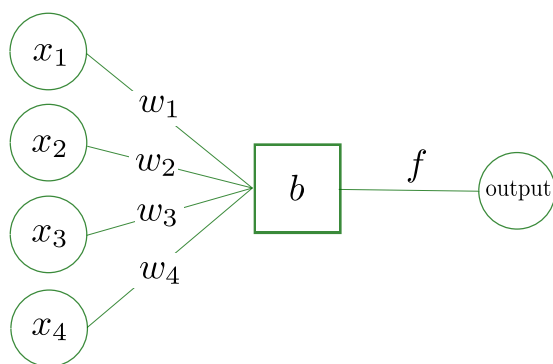
Figure 4.3: There are many different activations function which are used, and essentially any function can be defined if necessary. Above is shown a subset of the more commonly used functions. The linear function is not used, however, is there to compare to common non linear function.



(a) Simple neural network



(b) Simple neural network



(c) Simple neural network

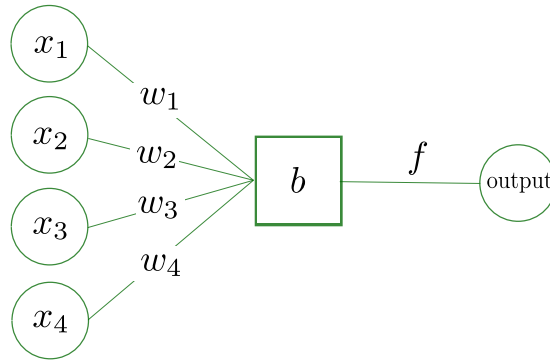


Figure 4.5: Max pooling layer

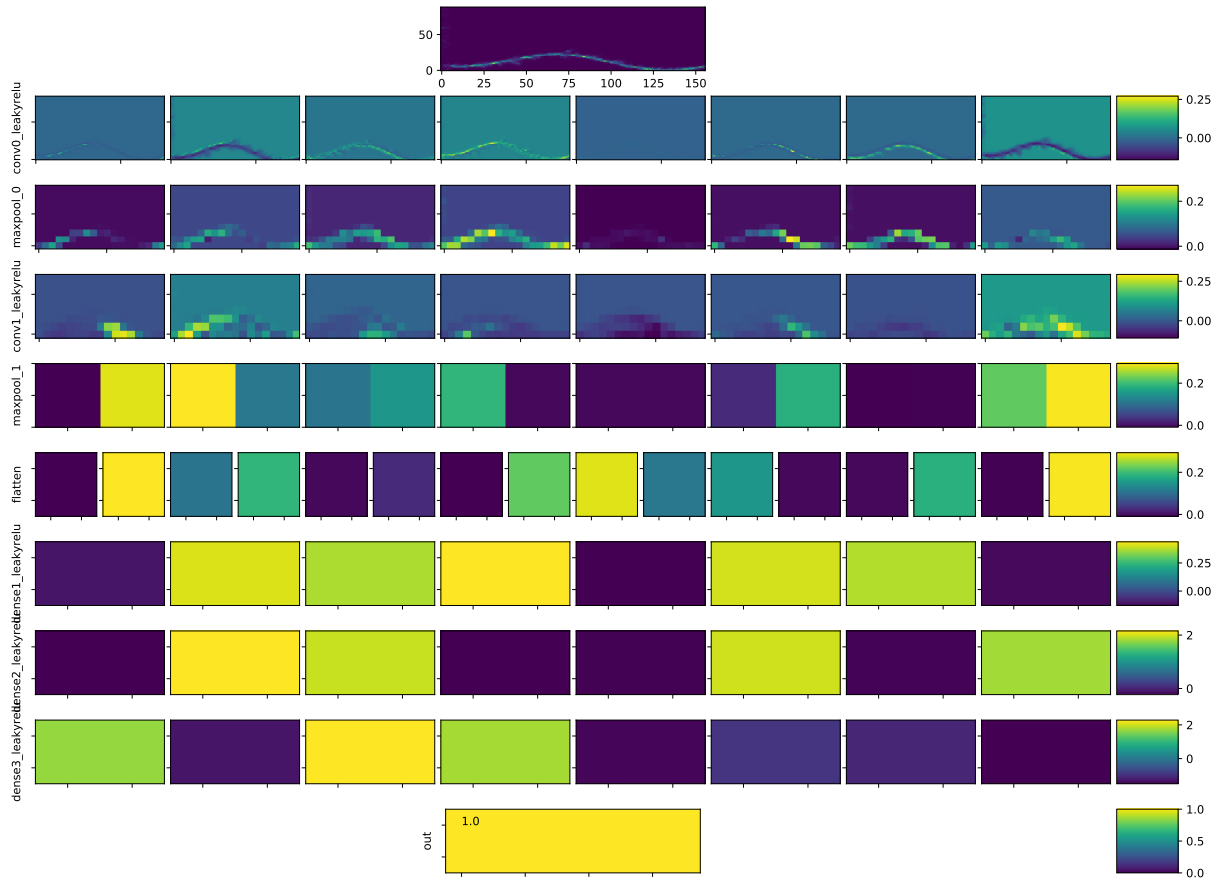


Figure 4.6: This shows a visualisation of the convolution neural network used for vitebri maps above

# Chapter 5

## Detector Characterisation with SOAP

When searching for GW signals, it is important to understand the origins of noise in the detectors data which does not originate from an astrophysical source. A large fraction of GW search algorithms, including SOAP above, assume that the detectors noise follows a Gaussian distribution. However, the detectors contain artefacts which are not distributed as a Gaussian. These artefacts can negatively affect many searches for GW as they can be easily mistaken for a real GW signal. Some of the potential sources of these artefacts have been mentioned in Sec. 1.3.1. There are many different classes of artefact, some key types include: glitches, these are short duration signals and instrumental lines which are long duration. To conduct a reliable search there are two main tasks which are necessary for detector characterisation. The first is identifying the artefact such that any search knows that part of data is contaminated. The search can then remove that section of data, or use more sophisticated techniques to deal with the artefact []. The second task is to find the source of the artefact. If the source of the artefact is found, it can potentially be removed or limited for future data runs.

The focus of this chapter is on a particular class of artefact called instrumental lines and how they affect CW searches. Sec. 5.1 will introduce different classes of instrumental line and how it affects a CW search. Sec. 5.2 will outline how these artefacts are detected and monitored, and describe current tools used for this task. Sec. 5.3 will describe how the CW search algorithm introduced in Sec. 3 can be used to search for instrumental lines. Finally Sec. 5.4 will show the outputs of the search and this is displayed for ease of use.

### 5.1 Instrumental lines

Instrumental lines have the general structure that they are persistent noise artefacts. There are many classes of instrumental line including: narrow, fixed frequency spectral artefacts or broad features which have a time varying frequency known as wandering lines. Many of these lines can make it difficult to distinguish them from an astrophysical signal. For many search methods, lines affect results in two main ways. They can cause the search to produce outliers which are

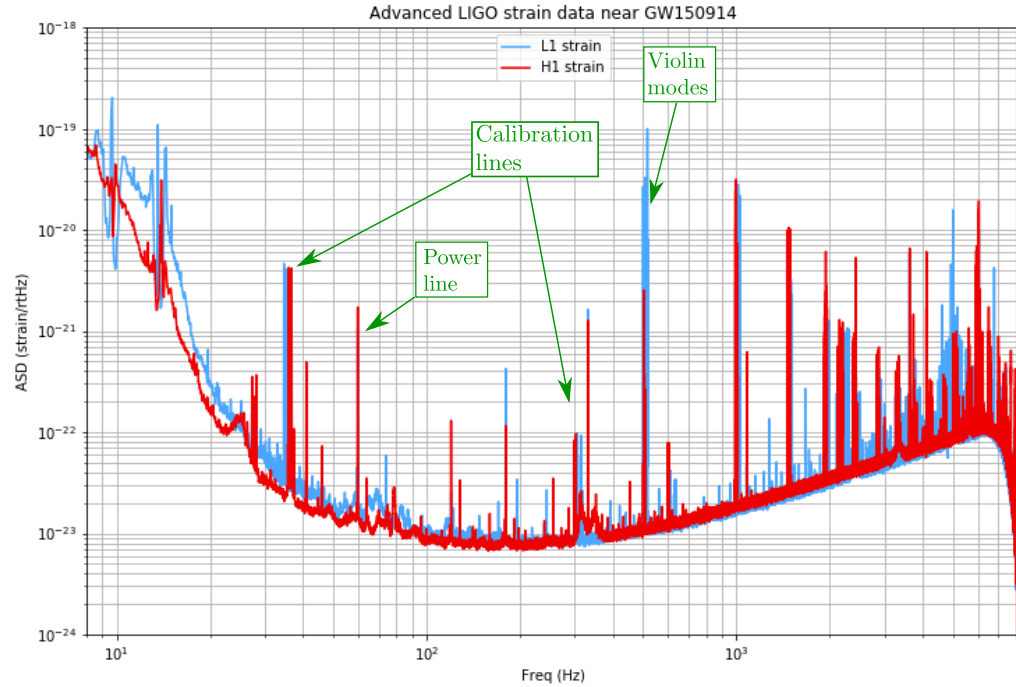


Figure 5.1: There are many features in the LIGO detectors averaged spectral density. This image was taken from [83] where I have annotated some of the stronger lines. The power line from the mains in the USA is at 60 Hz. Some of the calibration lines are around 30 Hz, 331 Hz and 1083 Hz. The various violin modes of the suspensions are at 300, 500, 600 and 900 Hz where I have only marked the 500 Hz mirror suspension modes [83].

considered as GW candidates. If the line is close to the GW signal in frequency, then it can conceal the power of the GW, or if it is overlapping, the search can be almost impossible. It is therefore crucial to understand the structure and origin of these lines when performing a search for GW specifically CW and stochastic searches.

Some instrumental lines are clearly visible when looking at one of the LIGO detectors frequency spectrum is shown in Fig. 5.3. This figure demonstrates some lines which have larger amplitudes. There are however, many more weaker lines which become visible when spectra are averaged over longer times. The frequency spectrum in Fig. 5.3 shows the time averaged spectra of the GW channel of the LIGO detectors. The lines seen in the spectrum are not from any GW and are usually from ground based sources. To see the lines in the GW channel, they must be coupled in via some mechanism. There are a number of ways in which this happens which are outlined in [84]. This includes coupling via shared power and grounds. When different components share the same power supplies, if a component draws power with a given period, then the voltage will decrease at this frequency. Another component which shares this same power supply can then also see this drop in voltage and this can potentially affect a stored output. Another

mechanism is coupling through magnetic fields, this is common when cables are nearby and can couple the magnetic field to different systems. Coupling can also occur through a physical connection, known as mechanical coupling, for example the resonances of the suspension fibres which couple directly to the mirrors.

Many of the spectral lines seen in the frequency spectrum, Fig. 5.3 are fundamental to the design of the detector. These cannot be eliminated, therefore need to be understood such that their effect on searches is minimised. Some of these lines are listed below.

**Power line** The power line harmonics are fundamental to the detector and originate from the US mains supply. These lines exist at 60 Hz which is the frequency of the mains alternating current oscillates at [75].

**Violin modes** The violin modes are associated with the suspensions of the mirrors and the beam splitter in the detector. These are designed to have a narrow frequency spectrum such that they contaminate as smaller part of the spectrum as possible. These are the lines around 500 Hz for the mirrors and 300, 600 and 900 Hz for the beam-splitter [83] in Fig. 5.3.

**Calibration lines** The mirrors of the LIGO detectors, are held at a resonance of the cavity in the arms of the interferometer. This then requires a feedback loop to hold them at resonance as a GW passes the detector. Calibration lines are used to calibrate this feedback loop such that the arm length changes are accurate [85, 76].

In [86], techniques to counter the affect of power lines and calibration lines.

Along with the fundamental lines of the detector which cannot be removed at the source, there are a large number of other lines whose source has been found and can be removed. Many of these are from mechanisms described earlier such as shared power supplies or grounds, these can be removed by, for example, using a different supply for different systems.

These lines have large effect on all searches for GWs both if the astrophysical signals frequencies overlap with the frequency of the line and can cause outliers. Long duration searches for CWs are particularly sensitive to this type of artefact. As described in Sec. 2, CW are long duration signals with a slowly varying frequency. In the case of an isolated neutron star, the signal which is searched for is narrow-band and a fixed frequency which is Doppler modulated by the earth's rotation and orbit. For certain areas of parameter space, the astrophysical signal of an isolated neutron star can appear very similar to an instrumental line. Many of these lines can be mitigated by using multiple detectors data. If a signal appears in a one detector and not the others, then it is likely that the signal is from an instrumental lines. These contaminated frequency bands can be removed of a statistic similar to that described in Sec. 3.8 or [87] can be used to limit their effect. However, there are many examples of instrumental line which appear at the same or similar frequencies in multiple detectors. These pose a real challenge to some CW searches, and require a lot of investigation to limit their affect.

## 5.2 Identifying and monitoring instrumental lines

When a detector is running, it is very important to identify instrumental lines and monitor them. This can then lead to either locating the origin of the line such that it can be removed, or allowing it to be flagged for other search algorithms. The lines are generally identified in the GW channel, this is the output of the detector which GW are observed and is the data using is previous chapters. This is the channel where the affect of lines is intended to be minimised.

As well and the GW channel the detector records many different channels known as auxiliary channels. These channels monitor many components of the detector. For example, the seismometers which are located near the corner stations and end stations are all channels which are monitored. There are magnetometers located around the electronics on the site to measure how these may be affecting the detector. These channels can be very useful in identifying the source of an instrumental line. These channels which monitor the environment are known as physical environment monitor (PEM). The main goal is to reduce the number of artefacts in the GW such that it is as close to Gaussian noise as possible. If an artefact shows up in the GW channel in coincidence with one of the PEM or other channels, then this is an indicator that the artefact originates from something related to that PEM. For example, say a magnetometer located near some electronics has a long duration line in its spectrum at the same frequency as the main GW channel. This indicates that the magnetic field of that piece of electronics is somehow coupling into the detector.

There are a number of tools which are used to monitor these spectral lines, along with a team of people which regularly look though the results, a summary of these for the first two observing runs of LIGO can be found in [84]. Some of the tools used are described below.

**Fscan** Fscan [76] takes FFTs of the raw detector data, typically these are 1800s long. This is done for all of the auxiliary channels as well as the GW channel. The FFTs are then averaged over a day and spectrograms are generated. After known lines such as Violin modes and power lines are subtracted, noise lines can be identified. A threshold can be set and anything in the spectrograms which exceed this threshold are flagged as a line. These can then be compared across multiple different channels. More detail on how the lines are identified can be found in [76]

**Coherence** Coherence searches for the coherence between different channels and different detectors. This is similar to searches for stochastic gravitational waves. More detail of how this works can be found in [84]

**Finetooth** Many of the instrumental lines found are part of combs. These are repeating structures where the harmonics (teeth) of a line make up the comb which is defined by the start frequency and tooth spacing. Finetooth is a tool which identifies and monitors these combs [].



**NoEMi** NoEMi uses FFTs and identifies peaks within the spectrum. I can then find coincidences with auxiliary channels and label if there are overlapping peaks with the GW channel. This can then track the lines with time. All of the information is stored in a database where more detail on its operation can be found in [88].

These tools offer many ways for the detector characterisation team to identify instrumental lines and hunt for their source. A summary of these efforts for the advanced LIGO data can be found in [84]. The following sections describe how the SOAP search described in Sec. 3 can be used as an extra tool to aid in the identification and monitoring of instrumental lines.

### 5.3 Identifying and cleaning lines with SOAP

The SOAP search has been run on a number of observing runs to search for CW. One of the major factors which limited the sensitivity of the search is instrumental lines. Whilst for the majority of this thesis I have worked on developing techniques to limit the affect instrumental lines, the search also has the ability to identify these lines well. In this section, I will explain the setup of the search to identify instrumental lines.

Sec. 3 described how multiple detectors can be used to increase the sensitivity of a search for CWs. However, when searching for instrumental lines, the aim is to have the reverse effect, therefore, a simple search is run separately on each detector. This then removed any of the statistics developed in Sec. 3.8 and uses the FFT power as the statistic in the SOAP search. The single detector soap search then has one parameter to vary, the transition matrix parameter. This governs how probable the frequency track is to transition up straight or down a frequency bin. In this search we are aiming to find any non Gaussian artefacts, therefore, we allow an equal probability for the track to jump in any direction. The Fscan search described above generates FFTs every day of varying lengths, this include 1800 s long FFTs. We use this to our advantage as the search is already set up to search through 1800 s long time-frequency spectrograms. For this line search, we split the 1800 s long Fscan FFTs into 0.2 Hz wide sub-bands and run the single detector search on each sub-band. This then returns the same outputs as described in Sec. 3 and Sec. 4: the frequency track (Viterbi track), a Viterbi map and a Viterbi statistic. Here the Viterbi statistic is just the sum of the FFT power along the frequency track. This outputs a plot as shown in Fig. 5.2.

The Viterbi statistics are then ranked in order, where the largest value is most likely to be from an instrumental artefact. The top results can be compared to the known line lists from the other line tools mentioned in Sec. 5.2.

- overview of soap search and what type is used for line searches
- how this is applied and how to interpret the output for line searches

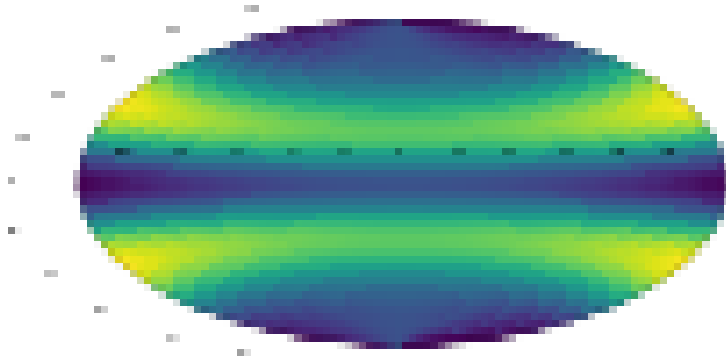


Figure 5.2: Example of output plot from SOAP line search.

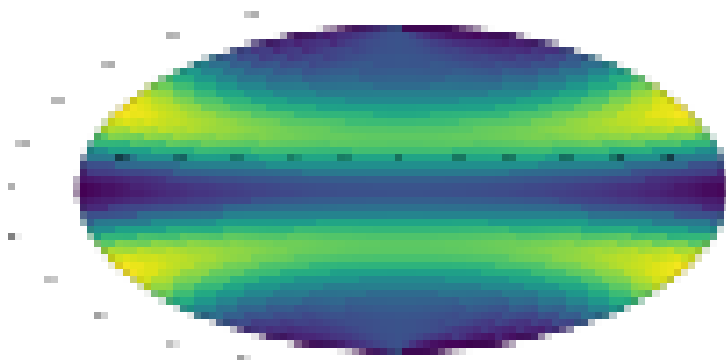


Figure 5.3: different types of instrumental line which appear in SOAP spectrograms

- SOAP can identify weak lines when they are wandering or fixed frequency
- 

## 5.4 Summary pages

Summary pages are an important tool when searching for instrumental lines. There is such a large amount of data both in frequency and time space and channels to search though when looking for instrumental lines. Summary pages should provide an easy way to view the important information to identifying the line as well as be easy to navigate to find a particular area of the spectrum which is of interest. These summary pages exist for the above searches in [line summary pages].

For the SOAP search summary pages were generated for each observing run and for the two LIGO detectors. This was done for various timescales: for the entire observing run and separately for each month. This allows the variation of a line to be observed for the entire length and also artefacts on shorter timescales to be observed. An example of the output plot is shown in Fig. 5.4 for the entire O1 observing run. The plots three main panels which are explained below.

**Spectrograms** The time-frequency spectrograms show the power spectrum of the 1800 s FFTs. The optimal Viterbi track is then overlaid on top.

**Viterbi maps** The Viterbi maps show the ‘log-probability’ of a signal being present in any particular frequency bin at a given time. This can be thought of as a distilled view of the spectrograms above.

**Sensitivity and power** The final panel in the example plots show an estimate of the sensitivity with time. The plot also shows the FFT power along the optimal tracks path. This allows the user to see how the sensitivity varies with the spectrogram power along the signal track. This can be useful when trying to identify some spectral artefacts.

- why summary pages are useful
- how to read them and what they mean
- how to use the information with other searches

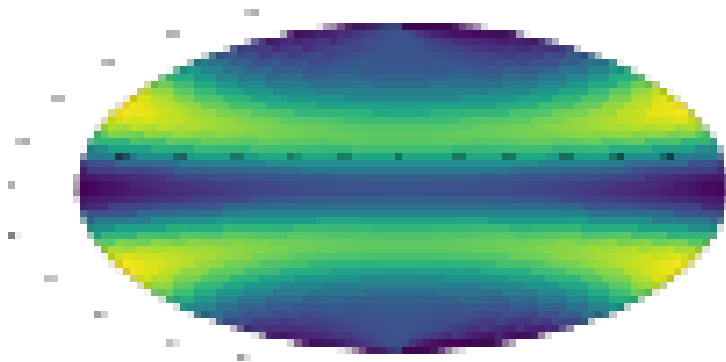


Figure 5.4: example plot of the summary pages

# **Chapter 6**

## **Summary**

# Appendix A

## Continuous gravitational wave injections

In this section I outline how we inject a CW signal into data. This can generally be done in two different ways: simulating a signal in the time domain and injecting into time domain noise or simulating the power spectrum of a signal and injecting the signal into a PSD.

### A.0.1 CW Signal

This section has been covered in Sec. ??, The CW signal is generated by using a lalsuite package where ....

### A.0.2 Time series and complex FFT injections

Injections into time-series data is relatively simple. Given a set of parameters for the source the signal can be generated in the time-series, this is then just summed with the time-series which it is injected into. Similarly with the FFT, the time-series of the signal at the correct time and for the correct duration is generated, the complex FFTs are then summed.

### A.0.3 Spectrogram injections

To inject into a spectrogram the power spectrum of the signal will need to be simulated. In our injection we do not have access to a time-series, therefore, we do not simulate the signal in the same way, rather we use the signals estimated SNR

It can be shown that the PSD of Gaussian noise with zero mean and unit variance is a  $\chi^2$  distribution with 2 degrees of freedom. Therefore, if we want to generate a spectrogram for Gaussian noise, we just generate a two dimensional array of values distributed as  $\chi^2$  with two degrees of freedom. Assuming that there is some sinusoidal signal with a given SNR within a Gaussian noise time-series with zero mean and unit variance, the FFT power in a particular frequency bin can be estimated using a non-central  $\chi^2$  distribution with 2 degrees of freedom, where the non centrality parameter is the square of the SNR. To calculate the SNR in a given

frequency bin the equation in [47] for optimal SNR was used,

$$\rho(0)^2 = \frac{1}{2} h_0^2 T S^{-1} [\alpha_1 A + \alpha_2 B + \alpha_3 C], \quad (\text{A.1})$$

where  $h_0$  is the GW amplitude,  $T$  is the total observing time in seconds,  $S^{-1}$  is the mean PSD noise floor. The values of  $\alpha$  are then defined by,

$$\begin{aligned} \alpha_1 &= (\mathcal{A}^1)^2 + (\mathcal{A}^3)^2 \\ \alpha_2 &= (\mathcal{A}^2)^2 + (\mathcal{A}^4)^2 \\ \alpha_3 &= \mathcal{A}^1 \mathcal{A}^2 + \mathcal{A}^3 \mathcal{A}^4 \end{aligned} \quad (\text{A.2})$$

where,

$$\begin{aligned} \mathcal{A}^1 &= A_+ \cos(2\psi_0) \cos(2\phi) - A_\times \sin(2\psi_0) \sin(2\phi) \\ \mathcal{A}^2 &= A_+ \cos(2\psi_0) \sin(2\phi) + A_\times \sin(2\psi_0) \cos(2\phi) \\ \mathcal{A}^3 &= A_+ \sin(2\psi_0) \cos(2\phi) - A_\times \cos(2\psi_0) \sin(2\phi) \\ \mathcal{A}^4 &= A_+ \sin(2\psi_0) \sin(2\phi) + A_\times \cos(2\psi_0) \cos(2\phi) \end{aligned} \quad (\text{A.3})$$

The signals frequency varies with time and will not always be located at the center of a frequency bin, therefore, when taking the FFT some of the power is spread over multiple frequency bins. In our injections into the power spectrum we need to account for this effect. For a given frequency bin width

#### A.0.4 Signals in data

The data recorded from a detector,  $x(t)$ , will include the signal model described in Eq. ?? above, but it will be buried in the noise of the detector. If we assume the noise is Gaussian distributed and the noise and signal add linearly, then,

$$x(t) = n(t) + h(t; \mathcal{A}, \lambda), \quad (\text{A.4})$$

where  $n(t)$  is the noise,  $h(t)$  is the signal and  $\mathcal{A}$  and  $\lambda$  refer to the amplitude and Doppler parameters respectively. The optimal signal to noise ratio (SNR) squared of this signal is defined as the scalar product of the signal with itself,

$$\rho^2(0) = (\mathbf{h} | \mathbf{h}) = \sum_X (h^X | h^X), \quad (\text{A.5})$$

where if there is more than one detector the SNR squared for each detector  $X$  can be summed [47]. The scalar product of two time series,  $x(t)$  and  $y(t)$ , is defined by,

$$(x | y) = 4\Re \int_0^\infty \frac{\tilde{x}(f)\tilde{y}^*(f)}{S_n(f)} df, \quad (\text{A.6})$$

where  $\tilde{x}(f)$  is the Fourier transform of  $x(t)$ ,  $\tilde{y}^*(f)$  is the complex conjugate of the Fourier transform of  $y(t)$  and  $S_n(f)$  is the single sided noise power spectral density [47].



# Bibliography

- [1] A. Einstein. “Die Grundlage Der Allgemeinen Relativitätstheorie [AdP 49, 769 (1916)]”. en. In: *Annalen der Physik* 14.S1 (2005), pp. 517–571. ISSN: 1521-3889. DOI: [10 . 1002 / andp . 200590044](https://doi.org/10.1002/andp.200590044). URL: [https : // onlinelibrary . wiley . com / doi / abs / 10 . 1002 / andp . 200590044](https://onlinelibrary.wiley.com/doi/abs/10.1002/andp.200590044) (visited on 02/26/2020).
- [2] Joel M. Weisberg, Joseph H. Taylor, and Lee A. Fowler. “Gravitational Waves from an Orbiting Pulsar”. In: *Sci. Am.* 245.4 (Oct. 1981), pp. 74–82. ISSN: 0036-8733. DOI: [10 . 1038 / scientificamerican1081-74](https://doi.org/10.1038/scientificamerican1081-74).
- [3] J. M. Weisberg and J. H. Taylor. “Relativistic Binary Pulsar B1913+16: Thirty Years of Observations and Analysis”. In: (July 2004). URL: [http : // arxiv . org / abs / astro-ph / 0407149](http://arxiv.org/abs/astro-ph/0407149).
- [4] B. P. Abbott et al. “Observation of Gravitational Waves from a Binary Black Hole Merger”. In: *Phys. Rev. Lett.* 116.6 (Feb. 2016), p. 061102. ISSN: 10797114. DOI: [10 . 1103 / PhysRevLett . 116 . 061102](https://doi.org/10.1103/PhysRevLett.116.061102). URL: [https : // link . aps . org / doi / 10 . 1103 / PhysRevLett . 116 . 061102](https://link.aps.org/doi/10.1103/PhysRevLett.116.061102).
- [5] B. P. P. Abbott et al. “GW170814: A Three-Detector Observation of Gravitational Waves from a Binary Black Hole Coalescence”. In: *Phys. Rev. Lett.* 119.14 (Oct. 2017), p. 141101. ISSN: 10797114. DOI: [10 . 1103 / PhysRevLett . 119 . 141101](https://doi.org/10.1103/PhysRevLett.119.141101). URL: [https : // link . aps . org / doi / 10 . 1103 / PhysRevLett . 119 . 141101](https://link.aps.org/doi/10.1103/PhysRevLett.119.141101).
- [6] The LIGO Scientific Collaboration et al. “GW190425: Observation of a Compact Binary Coalescence with Total Mass  $3.4 M_{\odot}$ ”. In: (Jan. 2020). URL: [http : // arxiv . org / abs / 2001 . 01761](http://arxiv.org/abs/2001.01761).
- [7] B. P. P. Abbott et al. “GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral”. In: *Phys. Rev. Lett.* 119.16 (Oct. 2017), p. 161101. ISSN: 10797114. DOI: [10 . 1103 / PhysRevLett . 119 . 161101](https://doi.org/10.1103/PhysRevLett.119.161101). URL: [https : // link . aps . org / doi / 10 . 1103 / PhysRevLett . 119 . 161101](https://link.aps.org/doi/10.1103/PhysRevLett.119.161101).
- [8] Eanna E. Flanagan and Scott A. Hughes. “The Basics of Gravitational Wave Theory”. In: *New J. Phys.* 7 (Jan. 2005). DOI: [10 . 1088 / 1367-2630 / 7 / 1 / 204](https://doi.org/10.1088/1367-2630/7/1/204). URL: [http : // arxiv . org / abs / gr-qc / 0501041](http://arxiv.org/abs/gr-qc/0501041) [http : // dx . doi . org / 10 . 1088 / 1367-2630 / 7 / 1 / 204](http://dx.doi.org/10.1088/1367-2630/7/1/204).

- [9] Collin Capano. “Searching for Gravitational Waves from Compact Binary Coalescence Using LIGO and Virgo Data”. In: *Phys. - Diss.* (Dec. 2011). URL: [https://surface.syr.edu/phy\\_etd/114](https://surface.syr.edu/phy_etd/114).
- [10] J. Aasi et al. “Advanced LIGO”. In: *Class. Quantum Gravity* 32.7 (Apr. 2015), p. 074001. ISSN: 13616382. DOI: [10.1088/0264-9381/32/7/074001](https://doi.org/10.1088/0264-9381/32/7/074001). URL: <http://stacks.iop.org/0264-9381/32/i=7/a=074001?key=crossref.20895763c84bce3f8929251031b2475c>.
- [11] F. Acernese et al. “Advanced Virgo: A Second-Generation Interferometric Gravitational Wave Detector”. In: *Class. Quantum Gravity* 32.2 (Jan. 2015), p. 024001. ISSN: 13616382. DOI: [10.1088/0264-9381/32/2/024001](https://doi.org/10.1088/0264-9381/32/2/024001). URL: <http://stacks.iop.org/0264-9381/32/i=2/a=024001?key=crossref.f3779cf21eab5f6631c6650a20>.
- [12] Francisco Hernandez Vivanco et al. “Measuring the Neutron Star Equation of State with Gravitational Waves: The First Forty Binary Neutron Star Merger Observations”. In: *Phys. Rev. D* 100.10 (Nov. 2019), p. 103009. DOI: [10.1103/PhysRevD.100.103009](https://doi.org/10.1103/PhysRevD.100.103009). URL: <https://link.aps.org/doi/10.1103/PhysRevD.100.103009> (visited on 02/26/2020).
- [13] Ian Harry and Tanja Hinderer. “Observing and Measuring the Neutron-Star Equation-of-State in Spinning Binary Neutron Star Systems”. en. In: *Class. Quantum Grav.* 35.14 (June 2018), p. 145010. ISSN: 0264-9381. DOI: [10.1088/1361-6382/aac7e3](https://doi.org/10.1088/1361-6382/aac7e3). URL: <https://doi.org/10.1088/1361-6382/aac7e3> (visited on 02/26/2020).
- [14] The LIGO Scientific Collaboration and The Virgo Collaboration et al. “A Gravitational-Wave Standard Siren Measurement of the Hubble Constant”. en. In: *Nature* 551.7678 (Nov. 2017), pp. 85–88. ISSN: 0028-0836, 1476-4687. DOI: [10.1038/nature24471](https://doi.org/10.1038/nature24471). URL: <http://www.nature.com/articles/nature24471> (visited on 02/26/2020).
- [15] Michael Zevin et al. “Constraining Formation Models of Binary Black Holes with Gravitational-Wave Observations”. en. In: *ApJ* 846.1 (Sept. 2017), p. 82. ISSN: 0004-637X. DOI: [10.3847/1538-4357/aa8408](https://doi.org/10.3847/1538-4357/aa8408). URL: <https://doi.org/10.3847/1538-4357/aa8408> (visited on 02/26/2020).
- [16] Ilya Mandel and Alison Farmer. “Merging Stellar-Mass Binary Black Holes”. In: *arXiv:1806.05820 [astro-ph, physics:gr-qc]* (June 2018). URL: <http://arxiv.org/abs/1806.05820> (visited on 02/26/2020).
- [17] Neil J. Cornish and Tyson B. Littenberg. “Bayeswave: Bayesian Inference for Gravitational Wave Bursts and Instrument Glitches”. In: *Class. Quantum Gravity* 32.13 (July 2015), p. 135012. ISSN: 13616382. DOI: [10.1088/0264-9381/32/13/135012](https://doi.org/10.1088/0264-9381/32/13/135012).

- [18] S. Klimenko et al. “A Coherent Method for Detection of Gravitational Wave Bursts”. In: *Class. Quantum Gravity*. Vol. 25. IOP Publishing, June 2008, p. 114029. DOI: [10.1088/0264-9381/25/11/114029](https://doi.org/10.1088/0264-9381/25/11/114029).
- [19] Christian D. Ott. “The Gravitational Wave Signature of Core-Collapse Supernovae”. In: *Class. Quantum Gravity* 26.6 (Sept. 2008). DOI: [10.1088/0264-9381/26/6/063001](https://doi.org/10.1088/0264-9381/26/6/063001). URL: <http://arxiv.org/abs/0809.0695><http://dx.doi.org/10.1088/0264-9381/26/6/063001>.
- [20] J. Aasi et al. “Search for Gravitational Waves Associated with  $\gamma$ -Ray Bursts Detected by the Interplanetary Network”. In: *Phys. Rev. Lett.* 113.1 (June 2014), p. 011102. ISSN: 10797114. DOI: [10.1103/PhysRevLett.113.011102](https://doi.org/10.1103/PhysRevLett.113.011102).
- [21] Thibault Damour and Alexander Vilenkin. “Gravitational Radiation from Cosmic (Super)Strings: Bursts, Stochastic Background, and Observational Windows”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 71.6 (Mar. 2005), pp. 1–13. ISSN: 15502368. DOI: [10.1103/PhysRevD.71.063510](https://doi.org/10.1103/PhysRevD.71.063510).
- [22] Nelson Christensen. “Stochastic Gravitational Wave Backgrounds”. In: *Reports Prog. Phys.* 82.1 (Nov. 2018). DOI: [10.1088/1361-6633/aae6b5](https://doi.org/10.1088/1361-6633/aae6b5). URL: <http://arxiv.org/abs/1811.08797><http://dx.doi.org/10.1088/1361-6633/aae6b5>.
- [23] Joseph D. Romano. “Searches for Stochastic Gravitational-Wave Backgrounds”. In: (Aug. 2019). URL: <http://arxiv.org/abs/1909.00269>.
- [24] Sushan Konar. “Magnetic Fields of Neutron Stars”. In: *J. Astrophys. Astron.* 38.3 (Sept. 2017). DOI: [10.1007/s12036-017-9467-4](https://doi.org/10.1007/s12036-017-9467-4). URL: <http://arxiv.org/abs/1709.07106><http://dx.doi.org/10.1007/s12036-017-9467-4>.
- [25] Kostas Glampedakis and Leonardo Gualtieri. “Gravitational Waves from Single Neutron Stars: An Advanced Detector Era Survey”. In: (Sept. 2017), pp. 673–736. DOI: [10.1007/978-3-319-97616-7\\_12](https://doi.org/10.1007/978-3-319-97616-7_12). URL: <http://arxiv.org/abs/1709.07049>[http://dx.doi.org/10.1007/978-3-319-97616-7\\_12](http://dx.doi.org/10.1007/978-3-319-97616-7_12).
- [26] Keith Riles. “Recent Searches for Continuous Gravitational Waves”. In: *Mod. Phys. Lett. A* 32.39 (Dec. 2017), p. 1730035. ISSN: 0217-7323, 1793-6632. DOI: [10.1142/S021773231730035](https://doi.org/10.1142/S021773231730035). URL: <http://arxiv.org/abs/1712.05897> (visited on 02/20/2020).
- [27] B. Haskell et al. “Detecting Gravitational Waves from Mountains on Neutron Stars in the Advanced Detector Era”. In: *Mon. Not. R. Astron. Soc.* 450.3 (July 2015), pp. 2393–2403. ISSN: 0035-8711. DOI: [10.1093/mnras/stv726](https://doi.org/10.1093/mnras/stv726). URL: <http://academic.oup.com/mnras/article/450/3/2393/1056611/Detecting-gravitational-waves-from-mountains-on>.

- [28] Paul D. Lasky. “Gravitational Waves from Neutron Stars: A Review”. In: *Publ. Astron. Soc. Aust.* 32 (Aug. 2015). DOI: [10.1017/pasa.2015.35](https://doi.org/10.1017/pasa.2015.35). URL: <http://arxiv.org/abs/1508.06643><http://dx.doi.org/10.1017/pasa.2015.35>.
- [29] Werner Becker, ed. *Neutron Stars and Pulsars*. Vol. 357. Astrophysics and Space Science Library. Berlin, Heidelberg: Springer Berlin Heidelberg, 2009. ISBN: 978-3-540-76964-4. DOI: [10.1007/978-3-540-76965-1](https://doi.org/10.1007/978-3-540-76965-1). URL: <http://link.springer.com/10.1007/978-3-540-76965-1>.
- [30] Ian Jones. *The CFS Instability*. URL: <http://www.personal.soton.ac.uk/dij/cfs.html>.
- [31] S. Chandrasekhar. “Solutions of Two Problems in the Theory of Gravitational Radiation”. In: *Phys. Rev. Lett.* 24.11 (Mar. 1970), pp. 611–615. DOI: [10.1103/PhysRevLett.24.611](https://doi.org/10.1103/PhysRevLett.24.611). URL: <https://link.aps.org/doi/10.1103/PhysRevLett.24.611> (visited on 02/29/2020).
- [32] J. L. Friedman and B. F. Schutz. “Secular Instability of Rotating Newtonian Stars”. In: *The Astrophysical Journal* 222 (May 1978), pp. 281–296. ISSN: 0004-637X. DOI: [10.1086/156143](https://doi.org/10.1086/156143). URL: <http://adsabs.harvard.edu/abs/1978ApJ...222..281F> (visited on 02/29/2020).
- [33] Benjamin J. Owen et al. “Gravitational Waves from Hot Young Rapidly Rotating Neutron Stars”. In: 58 (1998), pp. 1–15. DOI: [10.1103/PhysRevD.58.084020](https://doi.org/10.1103/PhysRevD.58.084020). URL: <http://arxiv.org/abs/gr-qc/9804044><http://dx.doi.org/10.1103/PhysRevD.58.084020>.
- [34] J. Weber. “Observation of the Thermal Fluctuations of a Gravitational-Wave Detector”. In: *Phys. Rev. Lett.* 17.24 (Dec. 1966), pp. 1228–1230. ISSN: 00319007. DOI: [10.1103/PhysRevLett.17.1228](https://doi.org/10.1103/PhysRevLett.17.1228).
- [35] A. De Waard et al. “MiniGRAIL, the First Spherical Detector”. In: *Class. Quantum Gravity*. Vol. 20. IOP Publishing, May 2003, S143. DOI: [10.1088/0264-9381/20/10/317](https://doi.org/10.1088/0264-9381/20/10/317).
- [36] George Hobbs and Shi Dai. “Gravitational Wave Research Using Pulsar Timing Arrays”. In: *Natl. Sci. Rev.* 4.5 (Sept. 2017), pp. 707–717. ISSN: 2095-5138. DOI: [10.1093/nsr/nwx126](https://doi.org/10.1093/nsr/nwx126). URL: <https://academic.oup.com/nsr/article/4/5/707/4566535>.
- [37] P. A.R. Ade et al. “Constraints on Primordial Gravitational Waves Using Planck, WMAP, and New BICEP2/ Keck Observations through the 2015 Season”. In: *Phys. Rev. Lett.* 121.22 (Nov. 2018), p. 221301. ISSN: 10797114. DOI: [10.1103/PhysRevLett.121.221301](https://doi.org/10.1103/PhysRevLett.121.221301).

- [38] B. P. Abbott et al. “LIGO: The Laser Interferometer Gravitational-Wave Observatory”. In: *Reports Prog. Phys.* 72.7 (July 2009), p. 076901. ISSN: 00344885. DOI: [10.1088/0034-4885/72/7/076901](https://doi.org/10.1088/0034-4885/72/7/076901). URL: <http://stacks.iop.org/0034-4885/72/i=7/a=076901?key=crossref.669f7c34d7d28503fd7f84b53db6048f>.
- [39] F. Acernese et al. “Status of Virgo”. In: *Class. Quantum Gravity* 25.11 (June 2008), p. 114045. ISSN: 02649381. DOI: [10.1088/0264-9381/25/11/114045](https://doi.org/10.1088/0264-9381/25/11/114045). URL: <http://stacks.iop.org/0264-9381/25/i=11/a=114045?key=crossref.62c81a5e3c6d5978d66477c5fb155f0c>.
- [40] Nicholas Metropolis et al. “Equation of State Calculations by Fast Computing Machines”. In: *J. Chem. Phys.* 21.6 (June 1953), pp. 1087–1092. ISSN: 00219606. DOI: [10.1063/1.1699114](https://doi.org/10.1063/1.1699114). URL: <http://aip.scitation.org/doi/10.1063/1.1699114>.
- [41] Don van Ravenzwaaij, Pete Cassey, and Scott D. Brown. “A Simple Introduction to Markov Chain Monte–Carlo Sampling”. In: *Psychon. Bull. Rev.* 25.1 (Feb. 2018), pp. 143–154. ISSN: 15315320. DOI: [10.3758/s13423-016-1015-8](https://doi.org/10.3758/s13423-016-1015-8).
- [42] Sanjib Sharma. “Markov Chain Monte Carlo Methods for Bayesian Data Analysis in Astronomy”. In: *Annu. Rev. Astron. Astrophys.* 55.1 (Aug. 2017), pp. 213–259. ISSN: 0066-4146. DOI: [10.1146/annurev-astro-082214-122339](https://doi.org/10.1146/annurev-astro-082214-122339).
- [43] John Skilling. “Nested Sampling for General Bayesian Computation”. In: *Bayesian Anal.* 1.4 (2006), pp. 833–860. ISSN: 19360975. DOI: [10.1214/06-BA127](https://doi.org/10.1214/06-BA127).
- [44] Joshua S Speagle. “Dynesty: A Dynamic Nested Sampling Package for Estimating Bayesian Posteriors and Evidences”. In: *Mon. Not. R. Astron. Soc.* (Apr. 2019). DOI: [10.1093/mnras/staa278](https://doi.org/10.1093/mnras/staa278). URL: <http://arxiv.org/abs/1904.02180%20http://dx.doi.org/10.1093/mnras/staa278>.
- [45] B. Schutz. “Data Analysis of Gravitational-Wave Signals from Spinning Neutron Stars: The Signal and Its Detection”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 58.6 (Aug. 1998), p. 063001. ISSN: 15502368. DOI: [10.1103/PhysRevD.58.063001](https://doi.org/10.1103/PhysRevD.58.063001).
- [46] Réjean J. Dupuis and Graham Woan. “Bayesian Estimation of Pulsar Parameters from Gravitational Wave Data”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 72.10 (Nov. 2005), p. 102002. ISSN: 15507998. DOI: [10.1103/PhysRevD.72.102002](https://doi.org/10.1103/PhysRevD.72.102002). URL: <https://link.aps.org/doi/10.1103/PhysRevD.72.102002>.
- [47] Reinhard Prix. “Search for Continuous Gravitational Waves: Metric of the Multidetector F-Statistic”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 75.2 (Jan. 2007), p. 023004. ISSN: 15507998. DOI: [10.1103/PhysRevD.75.023004](https://doi.org/10.1103/PhysRevD.75.023004). URL: <http://arxiv.org/abs/gr-qc/0606088%20http://dx.doi.org/10.1103/PhysRevD.75.023004>.

- 75.023004%2010.1103/PhysRevD.75.069901%20<https://link.aps.org/doi/10.1103/PhysRevD.75.023004>.
- [48] S. Walsh et al. “Comparison of Methods for the Detection of Gravitational Waves from Unknown Neutron Stars”. In: *Phys. Rev. D* 94.12 (Dec. 2016), p. 124010. ISSN: 24700029. DOI: [10.1103/PhysRevD.94.124010](https://doi.org/10.1103/PhysRevD.94.124010). URL: <https://link.aps.org/doi/10.1103/PhysRevD.94.124010>.
- [49] Patrick R. Brady and Teviet Creighton. “Searching for Periodic Sources with LIGO. II. Hierarchical Searches”. In: *Phys. Rev. D* 61.8 (Feb. 2000), p. 082001. DOI: [10.1103/PhysRevD.61.082001](https://doi.org/10.1103/PhysRevD.61.082001). URL: <https://link.aps.org/doi/10.1103/PhysRevD.61.082001> (visited on 02/25/2020).
- [50] Curt Cutler, Iraj Gholami, and Badri Krishnan. “Improved Stack-Slide Searches for Gravitational-Wave Pulsars”. In: *Phys. Rev. D* 72.4 (Aug. 2005), p. 042004. DOI: [10.1103/PhysRevD.72.042004](https://doi.org/10.1103/PhysRevD.72.042004). URL: <https://link.aps.org/doi/10.1103/PhysRevD.72.042004> (visited on 02/25/2020).
- [51] Badri Krishnan et al. “Hough Transform Search for Continuous Gravitational Waves”. In: *Phys. Rev. D* 70.8 (Oct. 2004), p. 082001. DOI: [10.1103/PhysRevD.70.082001](https://doi.org/10.1103/PhysRevD.70.082001). URL: <https://link.aps.org/doi/10.1103/PhysRevD.70.082001> (visited on 02/25/2020).
- [52] F. Antonucci et al. “Detection of Periodic Gravitational Wave Sources by Hough Transform in The  $\theta$ - $\phi$  Plane”. In: *Class. Quantum Grav.* 25.18 (Sept. 2008), p. 184015. ISSN: 0264-9381. DOI: [10.1088/0264-9381/25/18/184015](https://doi.org/10.1088/0264-9381/25/18/184015). URL: <https://doi.org/10.1088/0264-9381/25/18/184015> (visited on 02/25/2020).
- [53] Pia Astone et al. “Method for All-Sky Searches of Continuous Gravitational Wave Signals Using the Frequency-Hough Transform”. In: *Phys. Rev. D* 90.4 (Aug. 2014), p. 042002. DOI: [10.1103/PhysRevD.90.042002](https://doi.org/10.1103/PhysRevD.90.042002). URL: <https://link.aps.org/doi/10.1103/PhysRevD.90.042002> (visited on 02/25/2020).
- [54] Avneet Singh et al. “Results of an All-Sky High-Frequency Einstein@Home Search for Continuous Gravitational Waves in LIGO’s Fifth Science Run”. In: *Phys. Rev. D* 94.6 (Sept. 2016), p. 064061. DOI: [10.1103/PhysRevD.94.064061](https://doi.org/10.1103/PhysRevD.94.064061). URL: <https://link.aps.org/doi/10.1103/PhysRevD.94.064061> (visited on 02/26/2020).
- [55] Maria Alessandra Papa et al. “Hierarchical Follow-up of Subthreshold Candidates of an All-Sky Einstein@Home Search for Continuous Gravitational Waves on LIGO Sixth Science Run Data”. In: *Phys. Rev. D* 94.12 (Dec. 2016), p. 122006. DOI: [10.1103/PhysRevD.94.122006](https://doi.org/10.1103/PhysRevD.94.122006). URL: <https://link.aps.org/doi/10.1103/PhysRevD.94.122006> (visited on 02/26/2020).



- [56] *Einstein@Home*. URL: <https://einsteinathome.org/> (visited on 02/26/2020).
- [57] J. Aasi et al. “Implementation of an  $\ell$ -Statistic All-Sky Search for Continuous Gravitational Waves in Virgo VSR1 Data”. en. In: *Class. Quantum Grav.* 31.16 (Aug. 2014), p. 165014. ISSN: 0264-9381. DOI: [10.1088/0264-9381/31/16/165014](https://doi.org/10.1088/0264-9381/31/16/165014). URL: <https://doi.org/10.1088/0264-9381/31/16/165014> (visited on 02/26/2020).
- [58] J. Abadie et al. “All-Sky Search for Periodic Gravitational Waves in the Full S5 LIGO Data”. In: *Phys. Rev. D* 85.2 (Jan. 2012), p. 022001. DOI: [10.1103/PhysRevD.85.022001](https://link.aps.org/doi/10.1103/PhysRevD.85.022001). URL: <https://link.aps.org/doi/10.1103/PhysRevD.85.022001>.
- [59] LIGO Scientific Collaboration and Virgo Collaboration et al. “Comprehensive All-Sky Search for Periodic Gravitational Waves in the Sixth Science Run LIGO Data”. In: *Phys. Rev. D* 94.4 (Aug. 2016), p. 042002. DOI: [10.1103/PhysRevD.94.042002](https://link.aps.org/doi/10.1103/PhysRevD.94.042002). URL: <https://link.aps.org/doi/10.1103/PhysRevD.94.042002> (visited on 02/26/2020).
- [60] LIGO Scientific Collaboration and Virgo Collaboration et al. “All-Sky Search for Periodic Gravitational Waves in the O1 LIGO Data”. In: *Phys. Rev. D* 96.6 (Sept. 2017), p. 062002. DOI: [10.1103/PhysRevD.96.062002](https://link.aps.org/doi/10.1103/PhysRevD.96.062002). URL: <https://link.aps.org/doi/10.1103/PhysRevD.96.062002> (visited on 02/28/2020).
- [61] Andrew J. Viterbi. “Error Bounds for Convolutional Codes and an Asymptotically Optimum Decoding Algorithm”. In: *IEEE Trans. Inf. Theory* 13.2 (Apr. 1967), pp. 260–269. ISSN: 15579654. DOI: [10.1109/TIT.1967.1054010](http://ieeexplore.ieee.org/document/1054010/). URL: <http://ieeexplore.ieee.org/document/1054010/>.
- [62] Reinhard Prix. “Gravitational Waves from Spinning Neutron Stars”. In: *Neutron Stars Pulsars*. Berlin, Heidelberg: Springer Berlin Heidelberg, 2009, pp. 651–685. DOI: [10.1007/978-3-540-76965-1\\_24](http://link.springer.com/10.1007/978-3-540-76965-1_24). URL: [http://link.springer.com/10.1007/978-3-540-76965-1\\_24](http://link.springer.com/10.1007/978-3-540-76965-1_24).
- [63] Benjamin J. Owen. “Probing Neutron Stars with Gravitational Waves”. In: (2009).
- [64] P. Astone et al. “A Method for Detection of Known Sources of Continuous Gravitational Wave Signals in Non-Stationary Data”. In: *Class. Quantum Gravity* 27.19 (Oct. 2010). ISSN: 02649381. DOI: [10.1088/0264-9381/27/19/194016](https://doi.org/10.1088/0264-9381/27/19/194016).
- [65] B. P. Abbott et al. “Searches for Gravitational Waves from Known Pulsars at Two Harmonics in 2015-2017 LIGO Data”. In: (2019).

- [66] Teviet Creighton. “Searching for Periodic Sources with LIGO. II. Hierarchical Searches”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 61.8 (Feb. 2000), p. 082001. ISSN: 15502368. DOI: [10.1103/PhysRevD.61.082001](https://doi.org/10.1103/PhysRevD.61.082001). URL: <https://link.aps.org/doi/10.1103/PhysRevD.61.082001>.
- [67] B. P. Abbott et al. “All-Sky Search for Continuous Gravitational Waves from Isolated Neutron Stars Using Advanced LIGO O2 Data”. In: *Phys. Rev. D* 100.2 (July 2019), p. 024004. ISSN: 2470-0010. DOI: [10.1103/physrevd.100.024004](https://doi.org/10.1103/physrevd.100.024004). URL: <https://link.aps.org/doi/10.1103/PhysRevD.100.024004>.
- [68] David R. Ellis. *Snakes on a Plane*. Action, Crime, Thriller. IMDb ID: tt0417148 event-location: USA. Aug. 2006.
- [69] S. Suvorova et al. “Hidden Markov Model Tracking of Continuous Gravitational Waves from a Binary Neutron Star with Wandering Spin. II. Binary Orbital Phase Tracking”. In: *Phys. Rev. D* 96.10 (Nov. 2017), p. 102006. DOI: [10.1103/PhysRevD.96.102006](https://doi.org/10.1103/PhysRevD.96.102006). URL: <https://link.aps.org/doi/10.1103/PhysRevD.96.102006>.
- [70] B. P. Abbott et al. “Search for Gravitational Waves from Scorpius X-1 in the First Advanced LIGO Observing Run with a Hidden Markov Model”. In: *Phys. Rev. D* 95.12 (June 2017), p. 122003. DOI: [10.1103/PhysRevD.95.122003](https://doi.org/10.1103/PhysRevD.95.122003). URL: <https://link.aps.org/doi/10.1103/PhysRevD.95.122003>.
- [71] B. P. Abbott et al. “Search for Gravitational Waves from a Long-Lived Remnant of the Binary Neutron Star Merger GW170817”. In: (2018).
- [72] Ling Sun and Andrew Melatos. “Application of Hidden Markov Model Tracking to the Search for Long-Duration Transient Gravitational Waves from the Remnant of the Binary Neutron Star Merger GW170817”. In: *arXiv e-prints*, arXiv:1810.03577 (Oct. 2018), arXiv:1810.03577.
- [73] LIGO Scientific Collaboration. “LIGO Algorithm Library - LALSuite”. In: (2018). DOI: [10.7935/GT1W-FZ16](https://doi.org/10.7935/GT1W-FZ16).
- [74] Berit Behnke, Maria Alessandra Papa, and Reinhard Prix. “Postprocessing Methods Used in the Search for Continuous Gravitational-Wave Signals from the Galactic Center”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 91.6 (Mar. 2015), p. 064007. ISSN: 15502368. DOI: [10.1103/PhysRevD.91.064007](https://doi.org/10.1103/PhysRevD.91.064007). URL: <https://link.aps.org/doi/10.1103/PhysRevD.91.064007>.
- [75] J. Aasi et al. “Characterization of the LIGO Detectors during Their Sixth Science Run”. In: *Class. Quantum Gravity* 32.11 (2015). ISSN: 13616382. DOI: [10.1088/0264-9381/32/11/115012](https://doi.org/10.1088/0264-9381/32/11/115012).



- [76] Michael Coughlin, the Ligo Scientific Collaboration, and the Virgo Collaboration. “Noise Line Identification in LIGO S6 and Virgo VSR2”. In: *Journal of Physics: Conference Series* 243.1 (2010), p. 012010. URL: <http://stacks.iop.org/1742-6596/243/i=1/a=012010>.
- [77] Joe Bayley, Graham Woan, and Chris Messenger. “SOAP: A Generalised Application of the Viterbi Algorithm to Searches for Continuous Gravitational-Wave Signals”. In: (Mar. 2019). DOI: [10.1103/PhysRevD.100.023006](https://doi.org/10.1103/PhysRevD.100.023006). URL: <http://arxiv.org/abs/1903.12614> <http://dx.doi.org/10.1103/PhysRevD.100.023006>.
- [78] Hunter Gabbard et al. “Matching Matched Filtering with Deep Networks for Gravitational-Wave Astronomy”. In: *Phys. Rev. Lett.* 120.14 (Apr. 2018), p. 141103. ISSN: 10797114. DOI: [10.1103/PhysRevLett.120.141103](https://doi.org/10.1103/PhysRevLett.120.141103). URL: <https://link.aps.org/doi/10.1103/PhysRevLett.120.141103>.
- [79] Daniel George and E. A. Huerta. “Deep Learning for Real-Time Gravitational Wave Detection and Parameter Estimation: Results with Advanced LIGO Data”. In: *Phys. Lett. Sect. B Nucl. Elem. Part. High-Energy Phys.* 778 (Mar. 2018), pp. 64–70. ISSN: 03702693. DOI: [10.1016/j.physletb.2017.12.053](https://doi.org/10.1016/j.physletb.2017.12.053). URL: <https://www.sciencedirect.com/science/article/pii/S0370269317310390?via%5C%3Dihub>.
- [80] Timothy D. Gebhard et al. “Convolutional Neural Networks: A Magic Bullet for Gravitational-Wave Detection?” In: *Phys. Rev. D* 100.6 (Sept. 2019), p. 063015. ISSN: 24700029. DOI: [10.1103/PhysRevD.100.063015](https://doi.org/10.1103/PhysRevD.100.063015). URL: <https://link.aps.org/doi/10.1103/PhysRevD.100.063015>.
- [81] Christoph Dreissigacker et al. “Deep-Learning Continuous Gravitational Waves”. In: *Phys. Rev. D* 100.4 (Aug. 2019), p. 044009. ISSN: 24700029. DOI: [10.1103/PhysRevD.100.044009](https://doi.org/10.1103/PhysRevD.100.044009). URL: <https://link.aps.org/doi/10.1103/PhysRevD.100.044009>.
- [82] Andrew L. Maas. “Rectifier Nonlinearities Improve Neural Network Acoustic Models”. In: 2013.
- [83] *GW Open Science Center*. URL: <https://www.gw-openscience.org/olspeclines/> (visited on 02/29/2020).
- [84] P. B. Covas et al. “Identification and Mitigation of Narrow Spectral Artifacts That Degrade Searches for Persistent Gravitational Waves in the First Two Observing Runs of Advanced LIGO”. In: *Phys. Rev. D* 97.8 (Apr. 2018), p. 082002. ISSN: 24700029. DOI: [10.1103/PhysRevD.97.082002](https://doi.org/10.1103/PhysRevD.97.082002).

- [85] D. Tuyenbayev et al. “Improving LIGO Calibration Accuracy by Tracking and Compensating for Slow Temporal Variations”. en. In: *Class. Quantum Grav.* 34.1 (Dec. 2016), p. 015002. ISSN: 0264-9381. DOI: [10.1088/0264-9381/34/1/015002](https://doi.org/10.1088/0264-9381/34/1/015002). URL: <https://doi.org/10.1088%5C%2F0264-9381%5C%2F34%5C%2F1%5C%2F015002> (visited on 02/28/2020).
- [86] Derek Davis et al. “Improving the Sensitivity of Advanced LIGO Using Noise Subtraction”. en. In: *Class. Quantum Grav.* 36.5 (Feb. 2019), p. 055011. ISSN: 0264-9381. DOI: [10.1088/1361-6382/ab01c5](https://doi.org/10.1088/1361-6382/ab01c5). URL: <https://doi.org/10.1088%5C%2F1361-6382%5C%2Fab01c5> (visited on 02/28/2020).
- [87] David Keitel et al. “Search for Continuous Gravitational Waves: Improving Robustness versus Instrumental Artifacts”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 89.6 (2014), pp. 1–19. ISSN: 15502368. DOI: [10.1103/PhysRevD.89.064023](https://doi.org/10.1103/PhysRevD.89.064023).
- [88] T Accadia et al. “The NoEMi (Noise Frequency Event Miner) Framework”. In: *J. Phys.: Conf. Ser.* 363 (June 2012), p. 012037. ISSN: 1742-6596. DOI: [10.1088/1742-6596/363/1/012037](https://doi.org/10.1088/1742-6596/363/1/012037). URL: <http://stacks.iop.org/1742-6596/363/i=1/a=012037?key=crossref.4b8bfd9320b240a51529da2d97549ab6> (visited on 02/28/2020).
- [89] James M. Lattimer and M. Prakash. “The Equation of State of Hot, Dense Matter and Neutron Stars”. In: *Phys. Rep.* 621 (Dec. 2015), pp. 127–164. DOI: [10.1016/j.physrep.2015.12.005](https://doi.org/10.1016/j.physrep.2015.12.005). URL: <http://arxiv.org/abs/1512.07820> <http://dx.doi.org/10.1016/j.physrep.2015.12.005>.
- [90] The LIGO Scientific Collaboration et al. “First Search for Gravitational Waves from Known Pulsars with Advanced LIGO”. In: 12 (2017). ISSN: 1538-4357. DOI: [10.3847/1538-4357/aa677f](https://doi.org/10.3847/1538-4357/aa677f). URL: <http://arxiv.org/abs/1701.07709> <http://dx.doi.org/10.3847/1538-4357/aa677f>.
- [91] S Suvorova et al. “Hidden Markov Model Tracking of Continuous Gravitational Waves from a Neutron Star with Wandering Spin”. In: *Phys. Rev. D - Part. Fields, Gravit. Cosmol.* 93.12 (2016), pp. 1–17. ISSN: 15502368. DOI: [10.1103/PhysRevD.93.123009](https://doi.org/10.1103/PhysRevD.93.123009).
- [92] L. Sun et al. “Hidden Markov Model Tracking of Continuous Gravitational Waves from Young Supernova Remnants”. In: (Oct. 2017). DOI: [10.1103/physrevd.97.043013](https://doi.org/10.1103/physrevd.97.043013). URL: <https://arxiv.org/abs/1710.00460>.
- [93] G Larry Bretthorst. “Bayesian Spectrum Analysis and Parameter Estimation”. In: *Springer-Verlag*. 1988, p. 220. ISBN: 0-387-96871-7.