

# Simulation of Quantum Systems with Time-Evolving Block Decimation

[https://github.com/jchryssanthacopoulos/quantum\\_information/tree/main/final\\_project](https://github.com/jchryssanthacopoulos/quantum_information/tree/main/final_project)

Quantum Information and Computing  
AA 2022–23

James Chryssanthacopoulos  
in collaboration with David Lange  
8 April 2023



UNIVERSITÀ  
DEGLI STUDI  
DI PADOVA

- 1** Theory
  - Simulating Quantum Systems
  - Matrix Product States
  - Time-Evolving Block Decimation
- 2** Code Development
  - Python Library for Quantum Many-Body Calculations
  - Implementation of Matrix Product States
  - Running TEBD
- 3** Results on 1D Quantum Ising Model
  - Ground State Energy
  - Magnetization
  - Entanglement Entropy

# Theory

- To study a quantum system, one has to solve Schrodinger equation

$$\hat{H} |\Psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\Psi(t)\rangle$$

- One method involves direct numerical integration, where initial state is updated using time evolution operator

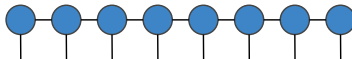
$$|\Psi(t + \Delta t)\rangle = e^{-i\hat{H}\Delta t} |\Psi(t)\rangle$$

- This requires solving system of equations at every time that scales with system size, but in many-body problems, system size is **exponential** in number of sites,  $N$
- In tensor network notation, general  $N$ -body system is shown on left. Mean-field ansatz on right greatly simplifies computation, but it ignores entanglement



How does one **preserve entanglement** while remaining **computationally tractable**?

- Matrix product states generalize mean-field ansatz to allow for entanglement between sites. Graphically,



where bond between sites has fixed **bond dimension**  $\chi$ . When  $\chi = 1$ , mean-field approximation recovered

- Wavefunction is given by

$$|\Psi\rangle = A_1^{\mu_1} A_{\mu_1,2}^{\mu_2} \cdots A_{\mu_{N-2},N-1}^{\mu_{N-1}} A_{\mu_{N-1},N}^{\mu_N} |1 2 \cdots N\rangle$$

where  $A_{\mu_{i-1},i}^{\mu_i}$  tensors have physical dimension  $i \in \{1, \dots, d\}$  and **auxiliary dimension**  $\mu_i \in \{1, \dots, \chi\}$

- Number of states scales like  $Nd\chi^2$ , which is **polynomial** in  $N$

How does one evolve MPS in time without **breaking structure**?

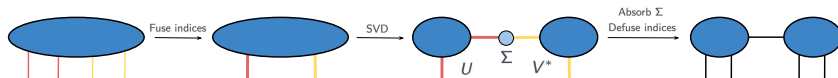
# Factoring Quantum State into MPS



- An arbitrary quantum state can be factored as an MPS using matrix factorization technique called **singular value decomposition**
- SVD generalizes eigendecomposition, finding two orthonormal bases  $u_i, w_i$  and singular values  $\sigma_i$  such that matrix is factorized into  $M = U\Sigma V^*$ , with  $U, V$  unitary

The diagram shows the SVD factorization of a matrix  $M$  into three matrices:  $U$ ,  $\Sigma$ , and  $V^*$ . Matrix  $M$  is a 4x4 grid of blue squares. Matrix  $U$  is a 4x4 grid of green squares. Matrix  $\Sigma$  is a 4x4 grid of white squares with a diagonal of red squares. Matrix  $V^*$  is a 4x4 grid of purple squares. The equation is  $M = U \Sigma V^*$ .

- SVD can be used to successively factor an MPS using following process:



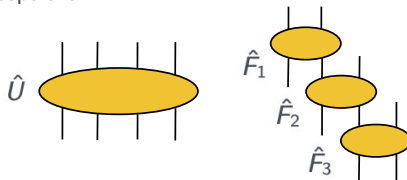
- SVD is driver behind simulating MPS quantum systems, allowing MPS structure to be preserved at each iteration. To keep bond dimension constant, singular values must be **truncated**

# Time-Evolving Block Decimation

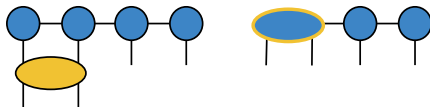


- Method for evolving quantum system while efficiently truncating large Hilbert space
- Also called **t-DMRG**, it evolves state using local gate operators and uses SVD to factorize back into MPS structure. The steps are:

- ① Factor time evolution operator  $\hat{U} : d^N \rightarrow d^N$  into two-site gates  $\hat{F}_i : d^2 \rightarrow d^2$



- ② Apply first gate to first two sites, contracting indices to produce new state



- ③ Use SVD to factor state back into MPS form, truncating to bond dimension  $\chi$



- ④ Repeat for each pair of neighbors and time step

- Approximate decomposition of Hamiltonian based on Baker-Campbell-Hausdorff formula that reduces time and storage complexity of applying time evolution operator
- Hamiltonian  $\hat{H}$  can be decomposed into odd and even operators:

$$\hat{H} = \sum_i \hat{h}_{i,i+1} = \sum_{i \text{ odd}} \hat{h}_{i,i+1} + \sum_{i \text{ even}} \hat{h}_{i,i+1} \equiv \hat{H}_{\text{odd}} + \hat{H}_{\text{even}}$$

- In **first-order Suzuki-Trotter**, commutator is ignored, leading to error  $\mathcal{O}(\Delta t^2)$ :

$$\hat{U} = e^{-i\hat{H}\Delta t} = e^{-i\hat{H}_{\text{even}}\Delta t} e^{-i\hat{H}_{\text{odd}}\Delta t} e^{-i[\hat{H}_{\text{even}}, \hat{H}_{\text{odd}}]\Delta t^2} \approx e^{-i\hat{H}_{\text{even}}\Delta t} e^{-i\hat{H}_{\text{odd}}\Delta t} + \mathcal{O}(\Delta t^2)$$

- Higher orders can be built by continuing to factor Hamiltonian into finer time steps. For example, **second-order Suzuki-Trotter** comes with  $\mathcal{O}(\Delta t^3)$  error:

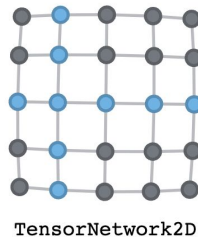
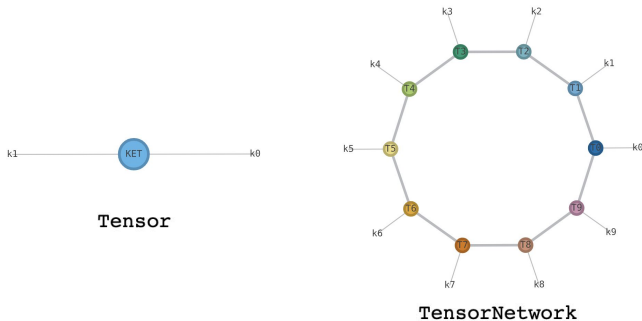
$$\hat{U} \approx e^{-i\hat{H}_{\text{even}}\Delta t/2} e^{-i\hat{H}_{\text{odd}}\Delta t} e^{-i\hat{H}_{\text{even}}\Delta t/2} + \mathcal{O}(\Delta t^3)$$

- When applied to time steps  $T/\Delta t$ , errors of ST1 and ST2 are  $\mathcal{O}(\Delta t)$  and  $\mathcal{O}(\Delta t^2)$ , respectively. ST2 has lower error but requires more operations



# Code Development

- Contains tools for working with tensors and tensor networks, including contracting, optimizing, and drawing them



- Although it supports more complicated geometries and algorithms, only the basic `Tensor` and `TensorNetwork` classes were used

# Implementing Matrix Product States



Class used in TEBD algorithm to model basic MPS structure, initializing, contracting, and computing observables from it

```
class MatrixProductState:
    """Class representing a matrix product state with given number of states."""

    def __init__(
        self, d: int, N: int, bond_dim: int, states: Optional[List[np.array]] = None, rng_seed: Optional[int] = 0
    ):
        """Initialize the MPS.

        Args:
            d: Dimension of each state
            N: Number of states
            bond_dim: Bond dimension between states
            states: Optional states to initialize with
            rng_seed: Number to seed random number generator

        """
        self.d = d
        self.N = N
        self.bond_dim = bond_dim

        self.data = []

        if not states:
            np.random.seed(rng_seed)

            states = []
            states.append(np.random.rand(d, bond_dim))

            for i in range(1, N - 1):
                states.append(np.random.rand(bond_dim, d, bond_dim))

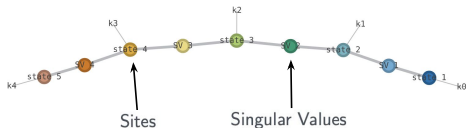
            states.append(np.random.rand(bond_dim, d))

        # create left-most state
        self.data.append(qtn.Tensor(states[0], inds=["k0", "i0"], tags=["state 1"]))

        for i in range(1, N - 1):
            self.data.append(
                qtn.Tensor(np.eye(bond_dim, bond_dim), inds=[f"i{2 + (i - 1)*3}", f"i{2 + i - 1}*3"], tags=[f"SV {i+1}"])
            )
            self.data.append(
                qtn.Tensor(states[i], inds=[f"i{2 * i - 1}", f"i{2 * i + 1}"], tags=[f"state {i + 1}"])
            )

        # create right-most state
        self.data.append(
            qtn.Tensor(np.eye(bond_dim, bond_dim), inds=[f"i{2 * (N - 2)*3}", f"i{2 * (N - 1)*3}], tags=[f"SV {N - 1}"])
        )
        self.data.append(qtn.Tensor(states[N - 1], inds=[f"i{2 * (N - 1)*3}", f"i{2 * (N - 1)*3}"], tags=[f"state {N}"]))

        self.normalize()
```



Contains methods for

Generating density matrix tensor

Computing magnetization

Calculating entropy

Extracting wavefunction

# Implementation of TEBD Algorithm



TEBD class accepts MatrixProductState object and implements step method that applies gate on every pair of states

## Apply Gate

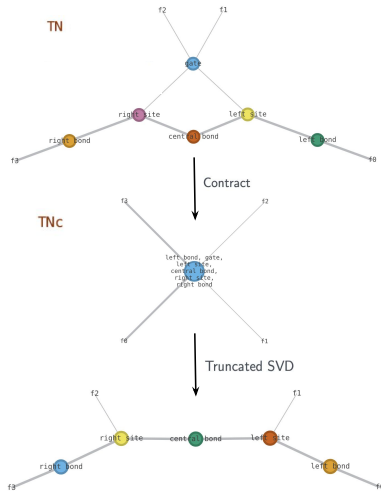
```
left_bond_T = qtn.Tensor(left_bond_data, inds=('f0', 'k1'), tags=['left bond'])
left_site_T = qtn.Tensor(left_site_data, inds=('k1', 'k2', 'k3'), tags=['left site'])
central_bond_T = qtn.Tensor(central_bond_data, inds=('k3', 'k4'), tags=['central bond'])
right_site_T = qtn.Tensor(right_site_data, inds=('k4', 'k5', 'k6'), tags=['right site'])
right_bond_T = qtn.Tensor(right_bond_data, inds=('k6', 'f3'), tags=['right bond'])
gate_T = qtn.Tensor(gate, inds=('f1', 'f2', 'k2', 'k5'), tags=['gate'])

# contract with gate
TN = left_bond_T & gate_T & left_site_T & central_bond_T & right_site_T & right_bond_T
TNc = TN ^ ...

# perform SVD
nshape = [self.d * left_site.data.shape[0], self.d * right_site.data.shape[2]]
utemp, stemp, vhtemp = LA.svd(TNc.data.reshape(nshape), full_matrices=False)

# truncate to reduced dimension
chitemp = min(self.bond_dim, len(stemp))
utemp = utemp[:, range(chitemp)].reshape(left_site.data.shape[0], self.d * chitemp)
vhtemp = vhtemp[range(chitemp), :].reshape(chitemp * self.d, right_site.data.shape[2])

# remove environment weights to form new MPS tensors A and B
left_site.modify(data=(LA.inv(left_bond_data) @ utemp).reshape(left_site.data.shape[0], self.d, chitemp))
right_site.modify(data=(vhtemp @ LA.inv(right_bond_data)).reshape(chitemp, self.d, right_site.data.shape[2]))
central_bond.modify(data=np.diag(stemp[range(chitemp)] / LA.norm(stemp[range(chitemp)]))
```

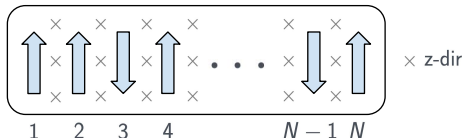


# Results

- TEBD can be applied to **quantum Ising model**, a quantum system composed of  $N$  spin-1/2 particles on one-dimensional lattice in presence of external magnetic field:

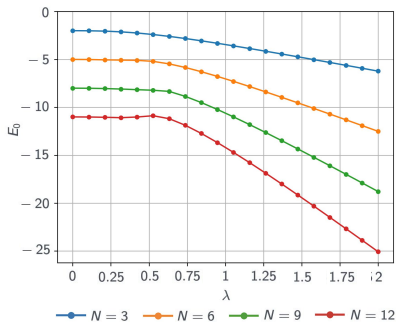
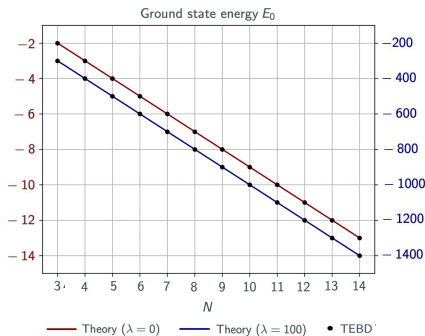
$$\hat{H} = J \sum_{i=1}^{N-1} \sigma_i^x \sigma_{i+1}^x + \lambda \sum_{i=1}^N \sigma_i^z$$

where  $J$  is coupling between neighboring spins and  $\lambda$  is coupling to magnetic field



- When  $\lambda = 0$ , all spins align when  $J < 0$  (**ferromagnetic**) and anti-align when  $J > 0$  (**antiferromagnetic**). The ground state is two-fold degenerate with energy  $-N + 1$
- When  $\lambda \rightarrow \infty$ , spins align to magnetic field and ground-state energy is  $-\lambda N$

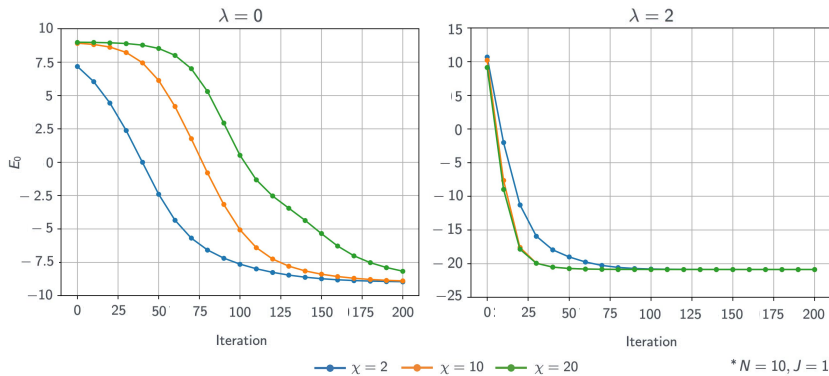
- Ground state energy  $E_0$  can be determined by simulating system using TEBD
- In following experiments, TEBD was run for 500 iterations with a timestep of  $\Delta t = 0.01$  and  $\chi = 2$ . The initial state was random, but was the same across all values of  $N$ . Results are robust to choice of  $\chi$
- TEBD reproduces  $E_0$  reported in Assignment 7. Results match theoretical expectation very well:  $E_0 = -N + 1$  for  $\lambda = 0$  and  $E_0 = -\lambda N$  for large  $\lambda$



# Effect of Bond Dimension



- Increasing bond dimension  $\chi$  allows model to capture more entanglement, but it takes longer to converge when  $\lambda$  is small since entanglement is low
- When  $\lambda$  is beyond critical point, larger  $\chi$  leads to faster convergence since degree of entanglement is higher



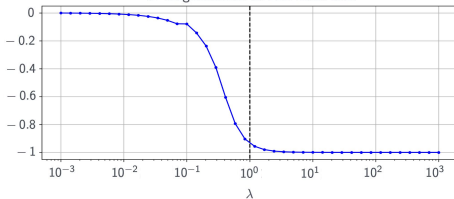


- Magnetization computed for site  $i$  along direction  $j$  using the magnetic operator:

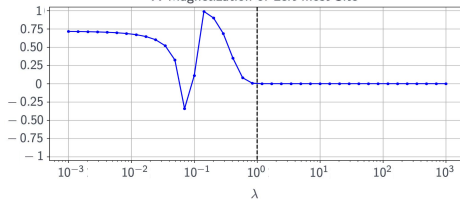
$$M_{ij} = \mathbb{1}_1 \otimes \cdots \otimes \mathbb{1}_{i-1} \otimes \sigma_i^j \otimes \mathbb{1}_{i+1} \otimes \cdots \otimes \mathbb{1}_N$$

- Average magnetization is given by  $\langle M_{ij} \rangle = \text{Tr}(\rho M_{ij}) = \langle \Psi | M_{ij} | \Psi \rangle$
- When  $\lambda$  is low, magnetization along  $Z$  is zero as spins have equal probability of being  $\uparrow$  or  $\downarrow$ . When  $\lambda$  increases, approaching phase transition at  $\lambda = 1$ , spins align to magnetic field and  $\langle M_{1Z} \rangle = 0$
- When  $Z$  magnetization is  $-1$ , average  $X$  magnetization is zero as  $Z$ -polarized qubit has equal probability of being  $\leftarrow$  or  $\rightarrow$

$Z$  Magnetization of Left-most Site



$X$  Magnetization of Left-most Site



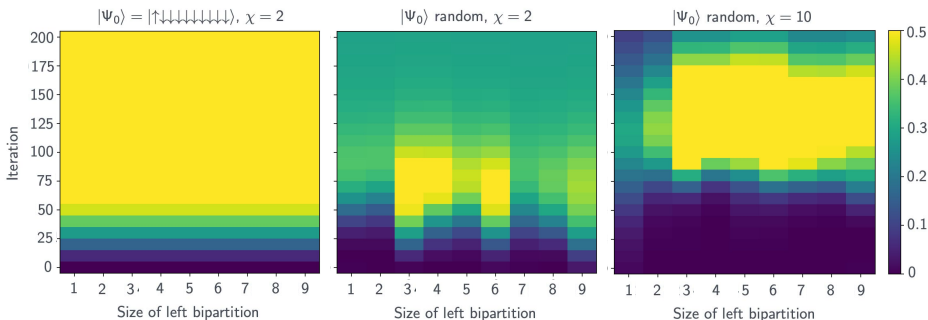
# Entanglement Entropy



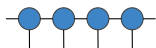
- Entropy of left bipartition computed by contracting density matrix tensor:

$$S = \text{Tr}(\rho_L \log \rho_L) \text{ where } \rho_L = \text{Tr}_R \rho =$$

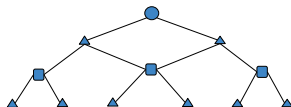
- For specific initial state  $|\uparrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\downarrow\rangle$ , entropy starts uniform and grows uniformly
- For random state, entropy grows and spreads unevenly before becoming uniform. With bigger bond dimension  $\chi$ , entropy across the chain persists for longer






# Conclusion



- Time-evolving block decimation is a powerful method of simulating many-body quantum systems using matrix product states
- Through efficient truncation of Hilbert space, algorithm remains computationally tractable while capturing some degree of entanglement
- TEBD can be applied to weakly-coupled one-dimensional quantum systems like Ising or Heisenberg models
- Similar techniques can be used applied using different tensor network structures, like tree tensor networks (e.g., t-MERA)



-  S. Paeckel, T. Köhler, A. Swoboda, S. R. Manmana, U. Schollwöck, and C. Hubig, “Time-evolution Methods for Matrix-product States,” *Annals of Physics*, vol. 411, December 2019, 167998
-  S. Montangero, “Introduction to Tensor Network Methods: Numerical Simulations of Low-Dimensional Many-Body Quantum Systems,” Springer Nature, Switzerland, 2018
-  J. Gray, “QUIMB: A Python Library for Quantum Information and Many-body Calculations,” *Journal of Open Source Software*, vol. 3, no. 29, 2018