## Assignment 5

https://github.com/jchryssanthacopoulos/quantum\_information/tree/main/assignment\_5

# Quantum Information and Computing AA 2022–23

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#### Time-Dependent Quantum System



■ Goal is to solve 1D time-dependent quantum harmonic oscillator

$$\hat{H}\left|\Psi(x,t)\right\rangle = i\hbar\frac{\partial}{\partial t}\left|\Psi(x,t)\right\rangle, \quad \hat{H} \equiv \hat{T} + \hat{V} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\left(\hat{x} - \frac{t}{T}\right)^2, \ t \in [0,T]$$

■ Problem was solved using **split operator method**, where state is evolved using

$$|\Psi(x,t+\Delta t)\rangle = e^{-\frac{i\hat{V}\Delta t}{2}}\mathcal{F}^{-1}e^{-i\hat{T}\Delta t}\mathcal{F}e^{-\frac{i\hat{V}\Delta t}{2}}|\Psi(x,t)\rangle$$

where  $\mathcal F$  and  $\mathcal F^{-1}$  are the Fourier transform and its inverse, respectively  $(\hbar=\omega=m=1 \text{ was used})$ 

```
! multiply by potential part of Hamiltonian
do ii = 1, N
    V(ii) = potential(x_grid(ii), time, tmax)
    final_state(ii) = cexp(cmplx(0.0, -0.5 * V(ii) * dt)) * init_state(ii)
end do
! normalize
call normalize(final_state, dx)
! call FFI to go from coordinate space to momentum space
call dfftx_plan_dft_ld(plan, Nx, final_state, state_transform, -1, 64)
call dfftx_execute_dft(plan, final_state, state_transform)
call dfftx_execute_dft(plan, final_state, state_transform)
call dfftx_execute_dft(plan, final_state, state_transform)
```

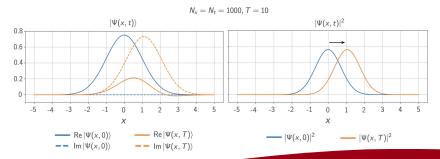
```
$ compiled/solve_time_dep_ho --xmin -15 --xmax 15 --tmax 10 --num_x_pts 100 --num_t_pts 100
xmin = 15.000
xmax = 15.000
xmax = 15.000
xmum_t_pts = 100
xmum_t_pts = 100
debug = F
output_filename = solution.txt
```

- fftw library installed from source and used to compute FFT
- System simulated using discretization parameters and wavefunction saved to file, results analyzed in Jupyter notebook

### Implementation and Solution



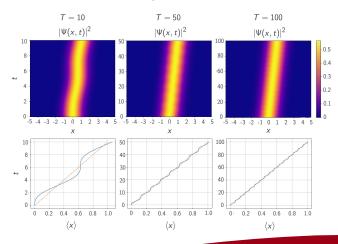
- $\blacksquare |\Psi(x,t)\rangle$  starts in ground state,  $(\frac{1}{\pi})^{1/4}e^{-\frac{x^2}{2}}$
- Position and time discretized using  $N_x$  and  $N_t$  equispaced bins of size  $\Delta x$  and  $\Delta t$
- Discretized momentum is  $\left[0, \frac{2\pi}{L}, \dots, \frac{\pi(N_x-2)}{L}, -\frac{\pi N_x}{L}, \dots, -\frac{2\pi}{L}\right]$ , where L is x range
- Wavefunction properly normalized by dividing by  $\sqrt{\sum_{i=1}^{N_x} |\Psi(x_i,t)|^2 \Delta x}$
- Wavefunction moves to right, developing complex phase, as expected



#### Wavefunction for Different T



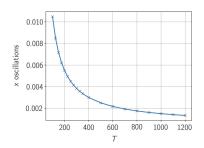
- Expected position  $\langle x(t) \rangle$  computed using  $\sum_{i=1}^{N_x} x_i |\Psi(x_i, t)|^2 \Delta x$
- Deviations from linear fit,  $\langle x(t) \rangle \frac{t}{T}$ , decrease with increasing T



#### Position Oscillations



- Position oscillations decrease when *T* increases because when *T* is large, potential changes slowly (i.e., adiabatic approximation)
- Fit of oscillations to  $\alpha T^{\beta}$  results in  $\alpha=0.60\pm0.03$  and  $\beta=-0.88\pm0.01$



lacktriangle Experiments showed that there are no oscillations with varying  $N_t$ 

## Self-Rating



- **Correctness.** Results seem sensible, but it is hard to know exactly without analytical wavefunction
- **Stability.** Code is numerically stable, but results sensitive to discretization
- Accurate Discretization. Error of split operator method is  $\mathcal{O}(\Delta t^3)$ . For  $\Delta t \sim 1$ , expected final state not reproduced
- Flexibility. Split operator method assumes  $\hat{V}$  is only function of  $\hat{x}$  and  $\hat{T}$  is only function of  $\hat{p}$ , enabling exact diagonal representation
- Efficiently. Overall time complexity is  $\mathcal{O}(N_t N_x \log N_x)$ . It took  $\sim 8$  s to solve using  $N_x = N_t = 10^4$