Simulation of Quantum Systems with Time-Evolving Block Decimation

https://github.com/jchryssanthacopoulos/quantum_information/tree/main/final_project

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Theory

Simulating Quantum Systems



■ To study a quantum system, one has to solve Schrodinger equation

$$\hat{H}|\Psi(t)\rangle = i\hbar \frac{\partial}{\partial t}|\Psi(t)\rangle$$

 One method involves direct numerical integration, where initial state is updated using time evolution operator

$$|\Psi(t+\Delta t)
angle = e^{-i\hat{H}\Delta t} |\Psi(t)
angle$$

- This requires solving system of equations at every time that scales with system size, but in many-body problems, system size is **exponential** in number of sites, *N*
- In tensor network notation, general N-body system is shown on left. Mean-field ansatz on right greatly simplifies computation, but it ignores entanglement





How does one preserve entanglement while remaining computationally tractable?

Matrix Product States



 Matrix product states generalize mean-field ansatz to allow for entanglement between sites. Graphically,



where bond between sites has fixed bond dimension $\chi.$ When $\chi=1$, mean-field approximation recovered

■ Wavefunction is given by

$$|\Psi
angle = A_1^{\mu_1} A_{\mu_1,2}^{\mu_2} \cdots A_{\mu_{N-2},N-1}^{\mu_{N-1}} A_{\mu_{N-1},N} |12 \cdots N
angle$$

where $A^{\mu_i}_{\mu_{i-1},i}$ tensors have physical dimension $i\in\{1,\ldots,d\}$ and auxiliary dimension $\mu_i\in\{1,\ldots,\chi\}$

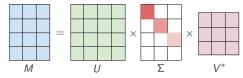
■ Number of states scales like $Nd\chi^2$, which is **polynomial** in N

How does one evolve MPS in time without breaking structure?

Factoring Quantum State into MPS



- An arbitrary quantum state can be factored as an MPS using matrix factorization technique called singular value decomposition
- SVD generalizes eigendecomposition, finding two orthonormal bases u_i , w_i and singular values σ_i such that matrix is factorized into $M = U\Sigma V^*$, with U, V unitary



■ SVD can be used to successively factor an MPS using following process:

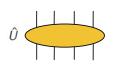


 SVD is driver behind simulating MPS quantum systems, allowing MPS structure to be preserved at each iteration. To keep bond dimension constant, singular values must be truncated

Time-Evolving Block Decimation

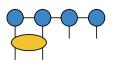


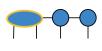
- Method for evolving quantum system while efficiently truncating large Hilbert space
- Also called t-DMRG, it evolves state using local gate operators and uses SVD to factorize back into MPS structure. The steps are:
- ① Factor time evolution operator $\hat{U}: d^N \to d^N$ into two-site gates $\hat{F}_i: d^2 \to d^2$





② Apply first gate to first two sites, contracting indices to produce new state





3 Use SVD to factor state back into MPS form, truncating to bond dimension χ





Repeat for each pair of neighbors and time step

Suzuki-Trotter Decomposition



- Approximate decomposition of Hamiltonian based on Baker-Campbell-Hausdorff formula that reduces time and storage complexity of applying time evolution operator
- Hamiltonian \hat{H} can be decomposed into odd and even operators:

$$\hat{H} = \sum_{i} \hat{h}_{i,i+1} = \sum_{i \text{ odd}} \hat{h}_{i,i+1} + \sum_{i \text{ even}} \hat{h}_{i,i+1} \equiv \hat{H}_{\text{odd}} + \hat{H}_{\text{even}}$$

■ In first-order Suzuki-Trotter, commutator is ignored, leading to error $\mathcal{O}(\Delta t^2)$:

$$\hat{\textit{U}} = e^{-i\hat{\textit{H}}\Delta t} = e^{-i\hat{\textit{H}}_{even}\Delta t} e^{-i\hat{\textit{H}}_{odd}\Delta t} e^{-i\left[\hat{\textit{H}}_{even},\hat{\textit{H}}_{odd}\right]\Delta t^2} \approx e^{-i\hat{\textit{H}}_{even}\Delta t} e^{-i\hat{\textit{H}}_{odd}\Delta t} + \mathcal{O}(\Delta t^2)$$

■ Higher orders can be built by continuing to factor Hamiltonian into finer time steps. For example, **second-order Suzuki-Trotter** comes with $\mathcal{O}(\Delta t^3)$ error:

$$\hat{U} pprox e^{-i\hat{H}_{\mathsf{even}}\Delta t/2} e^{-i\hat{H}_{\mathsf{odd}}\Delta t} e^{-i\hat{H}_{\mathsf{even}}\Delta t/2} + \mathcal{O}(\Delta t^3)$$

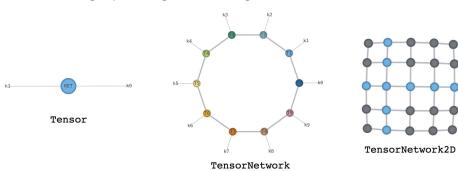
■ When applied to time steps $T/\Delta t$, errors of ST1 and ST2 are $\mathcal{O}(\Delta t)$ and $\mathcal{O}(\Delta t^2)$, respectively. ST2 has lower error but requires more operations

Code Development

QUIMB: Quantum Computing in Python



 Contains tools for working with tensors and tensor networks, including contracting, optimizing, and drawing them

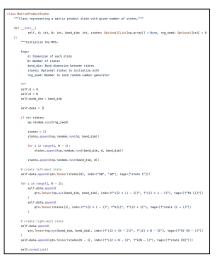


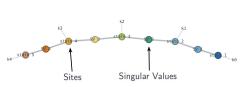
■ Although it supports more complicated geometries and algorithms, only the basic Tensor and TensorNetwork classes were used

Implementing Matrix Product States



Class used in TEBD algorithm to model basic MPS structure, initializing, contracting, and computing observables from it





Contains methods for

Generating density matrix tensor

Computing magnetization

Calculating entropy

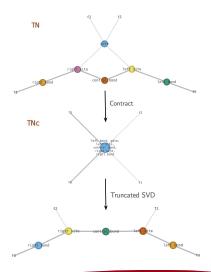
Extracting wavefunction

Implementation of TEBD Algorithm



TEBD class accepts MatrixProductState object and implements step method that applies gate on every pair of states

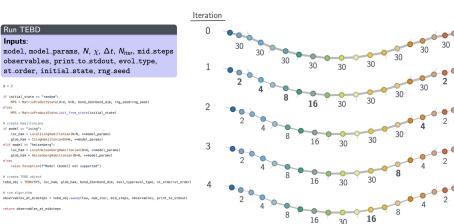
Apply Gate left_bond_T = qtn.Tensor(left_bond_data, inds=('f0', 'k1'), tags=['left bond']) left_site_T = qtn.Tensor(left_site.data, inds=('k1', 'k2', 'k3'), tags=['left site']) central bond T = gtm.Tensor(central bond.data, inds=('k3', 'k4'), tags=['central bond']) right_site_T = qtn.Tensor(right_site.data, inds=('k4', 'k5', 'k6'), tags=['right_site']) right_bond_T = qtn.Tensor(right_bond_data, inds=('k6', 'f3'), tags=('right bond')) gate_T = gtn.Tensor(gate, inds=('f1', 'f2', 'k2', 'k5'), tags=['gate']) # contract with gate TN = left_bond_T & gate_T & left_site_T & central_bond_T & right_site_T & right_bond_T TNc = TN ^ ... nshape = [self.d * left_site.data.shape[8], self.d * right_site.data.shape[2]] utemp, stemp, whtemp = LA.svd(TNc.data.reshape(nshape), full matrices=False) chitemp = min(self.bond_dim, len(stemp)) utemp = utemp[:, range(chitemp)], reshape(left site.data.shape[0], self.d * chitemp) whtemp = whtemp[range(chitemp), :],reshape(chitemp * self.d, right site.data.shape[2]) # remove environment weights to form new MPS tensors A and B left site,modify(data=(iA,iny(left bond data) @ utemp),reshape(left site,data,shape(8), self,d, chitemp)) right_site.modify(data=(whtemp @ LA.inv(right_bond_data)).reshape(chitemp, self.d, right_site.data.shape[2])) central_bond.modify(data=np.diag(stemp[range(chitemp)] / LA.norm(stemp[range(chitemp)])))



Running TEBD



- Main execution function accepts model, model parameters, and other run parameters
 - Bond dimension across chain evolves over first few iterations due to SVD



Results

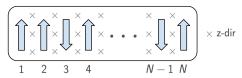
TEBD on Quantum Ising Model



■ TEBD can be applied to **quantum Ising model**, a quantum system composed of *N* spin-1/2 particles on one-dimensional lattice in presence of external magnetic field:

$$\hat{H} = J \sum_{i=1}^{N-1} \sigma_i^{\mathsf{x}} \sigma_{i+1}^{\mathsf{x}} + \lambda \sum_{i=1} \sigma_i^{\mathsf{z}}$$

where J is coupling between neighboring spins and λ is coupling to magnetic field

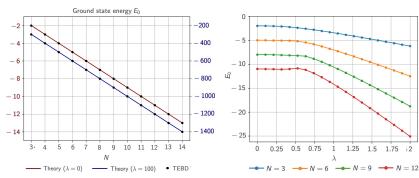


- When $\lambda=0$, all spins align when J<0 (ferromagnetic) and anti-align when J>0 (antiferromagnetic). The ground state is two-fold degenerate with energy -N+1
- When $\lambda \to \infty$, spins align to magnetic field and ground-state energy is $-\lambda N$

Ground State Energy



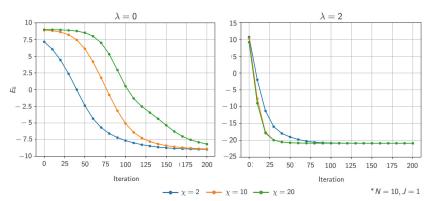
- \blacksquare Ground state energy E_0 can be determined by simulating system using TEBD
- In following experiments, TEBD was run for 500 iterations with a timestep of $\Delta t = 0.01$ and $\chi = 2$. The initial state was random, but was the same across all values of N. Results are robust to choice of χ
- TEBD reproduces E_0 reported in Assignment 7. Results match theoretical expectation very well: $E_0 = -N + 1$ for $\lambda = 0$ and $E_0 = -\lambda N$ for large λ



Effect of Bond Dimension



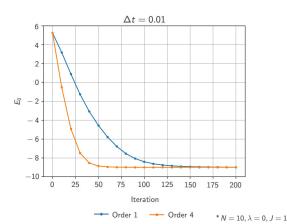
- \blacksquare Increasing bond dimension χ allows model to capture more entanglement, but it takes longer to converge when λ is small since entanglement is low
- \blacksquare When λ is beyond critical point, larger χ leads to faster convergence since degree of entanglement is higher



Effect of Suzuki-Trotter Order



- Increasing the order of Suzuki-Trotter decomposition improves convergence rate
- lacksquare Error of first order is $\mathcal{O}(\Delta t)$ while error of fourth error is $\mathcal{O}(\Delta t^4)$



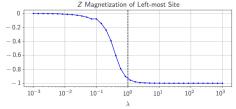
Magnetization

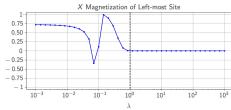


 \blacksquare Magnetization computed for site *i* along direction *j* using the magnetic operator:

$$M_{ij} = \mathbb{1}_1 \otimes \cdots \otimes \mathbb{1}_{i-1} \otimes \sigma_i^j \otimes \mathbb{1}_{i+1} \otimes \cdots \otimes \mathbb{1}_N$$

- Average magnetization is given by $\langle M_{ij} \rangle = \text{Tr}(\rho M_{ij}) = \langle \Psi | M_{ij} | \Psi \rangle$
- When λ is low, magnetization along Z is zero as spins have equal probability of being \uparrow or \downarrow . When λ increases, approaching phase transition at $\lambda=1$, spins align to magnetic field and $\langle M_{1Z} \rangle = 0$
- When Z magnetization is -1, average X magnetization is zero as Z-polarized qubit has equal probability of being \leftarrow or \rightarrow





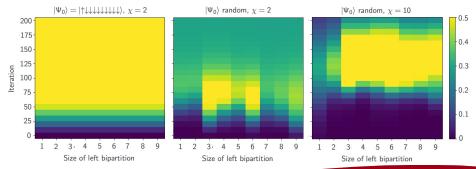
Entanglement Entropy



■ Entropy of left bipartition computed by contracting density matrix tensor:

$$S = \operatorname{Tr}(\rho_L \log \rho_L)$$
 where $\rho_L = \operatorname{Tr}_R \rho =$

- For specific initial state |↑↓↓↓↓↓↓↓⟩, entropy starts uniform and grows uniformly
- lacktriangleright For random state, entropy grows and spreads unevenly before becoming uniform. With bigger bond dimension χ , entropy across the chain persists for longer



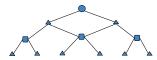
Conclusion

Conclusion





- Time-evolving block decimation is a powerful method of simulating many-body quantum systems using matrix product states
- Through efficient truncation of Hilbert space, algorithm remains computationally tractable while capturing some degree of entanglement
- TEBD can be applied to weakly-coupled one-dimensional quantum systems like Ising or Heisenberg models
- Similar techniques can be used applied using different tensor network structures, like tree tensor networks (e.g., t-MERA)



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