# Advanced Numerical Methods for Neutron Star Interfaces

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# Acknowledgements

Motivation: Neutron Stars

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- Ian Hawke
- Southampton General Relativity Group
- STFC
- IOP Gravity Group
- Classical and Quantum Gravity



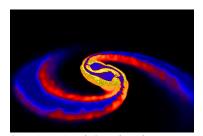
- Motivation: Neutron Stars
- Mow do we model a neutron star?
- Interfaces and Boundary Conditions
- Results
- Conclusion





# Motivation: Binary Mergers

- Neutron star binary mergers are strong candidates for the emission of detectable gravitational waves.
- Accurate gravitational wave templates are required to directly detect gravitational waves.
- Full non-linear numerical simulations are required for the plunge and merger phases.
- Therefore, accurate numerical simulations are needed.



Credit: Daniel Price (U/Exeter) and Stephan Rosswog (Int. U/Bremen)



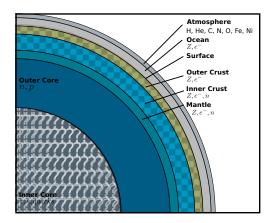
### Outline

- Mow do we model a neutron star?



## What features do we need to consider?

- Exterior Vacuum/Plasma
- Atmosphere Fluid
- Crust Elastic Matter
- Core Fluid





We want to solve the most realistic physical model possible using maximum resolution. Problem?

- Important length-scales vary over 12 orders of magnitude if including viscous boundary layers.
- Using adaptive mesh refinement cannot cover this range.

#### Solution:

Motivation: Neutron Stars

Relativistic multimodel approach developed by Millmore & Hawke. Approximate interfaces to be infinitely thin.



Simplify the most complex model by:

- Taking appropriate limits in different regions of the domain.
  - Force-free / Electro-vacuum / Multi-fluid MHD
  - Ideal MHD / Resistive MHD
  - Elasticity
- Combining the different approximations with sharp interfaces.
- Using evolution equations based on conservation laws.
  - This allows us to accurately capture the locations of shock waves
- Moving the complicated physics to the interfaces.
  - Boundary conditions to impose physics

This approach also allows a proper treatment of the surface.

Therefore, we no longer need a numerical atmosphere.



The evolution equations are a system of non-linear partial differential equations in conservation law form.

$$\frac{\partial \mathbf{q}(x,t)}{\partial x} + \frac{\partial \mathbf{F}(\mathbf{q}(x,t))}{\partial x} = 0, \tag{1}$$

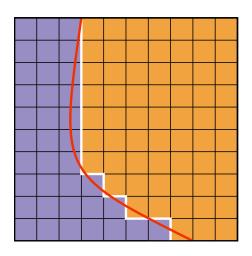
where  $\mathbf{q}$  are the conserved variables and  $\mathbf{F}$  are the fluxes. The numerical update is then given by.

$$q_i^{n+1} = q_i^n + \frac{\Delta t}{\Delta x} \left( F_{i-1/2}^n - F_{i+1/2}^n \right),$$
 (2)

where n is the time index and i is the spatial index.

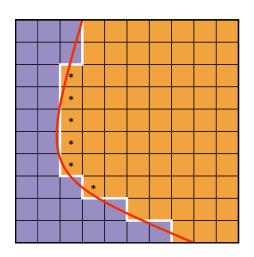


- We define a grid on which to evolve our models
- The red line indicates the true interface
- The white line indicates the numerical interface



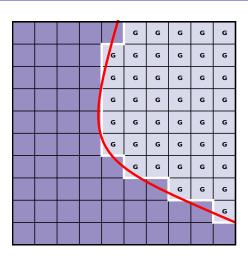


- The interface is advected with the models
- The numerical interface tracks the physical interface
- Points can change model

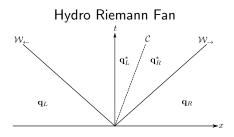




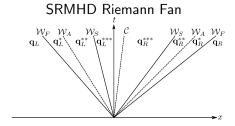
- Each model has a region of ghost zones
- As the update of each point is non-local, it is important to fill these cells correctly
- How do we impose the correct boundary conditions? By filling the ghost cells



- Each normal cell is updated by following the Godunov approch
- Each cell boundary is a Riemann problem
- A Riemann problem occurs when there is discontinuity
- We update each cell by calculating the flux through the boundaries



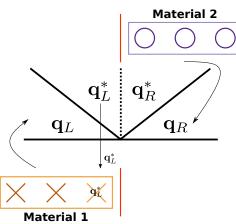
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- 2 How do we model a neutron star?
- 3 Interfaces and Boundary Conditions
- 4 Results
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## Boundary Conditions: Ghost Fluid Method

#### **Update Real Point**



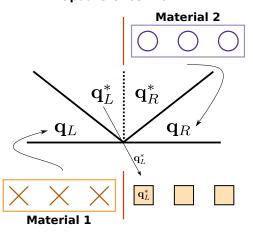
Interface

Follow the ghost fluid approach.

- Define left and right states.
- Calculate the star states using an appropriate method.
- Fill the points neighbouring the interface.
- Extrapolate into the ghost zones.



#### **Update Ghost Point**

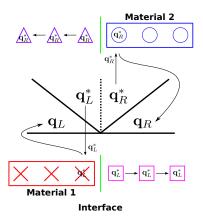


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To calculate the star states we follow the Roe approximation method and linearise about an appropriate state  $\mathbf{w}_X(\mathbf{q}_X)$  for the left and right states. Where X = L/R.

$$\partial_t \mathbf{w}_X + \hat{A}_X \partial_x \mathbf{w}_X = 0. (3)$$

The star states are then given by the following equations.

$$\mathbf{w}_L^* = \mathbf{w}_L + \sum_{i=1}^{N_L} c_i^L \mathbf{r}_L^{(i)}, \tag{4}$$

$$\mathbf{w}_{R}^{*} = \mathbf{w}_{R} - \sum_{i=1}^{N_{R}} c_{j}^{R} \mathbf{r}_{R}^{(j)}.$$
 (5)

Where  $\mathbf{r}_X$  are the right eigenvectors of the matrix  $\hat{A}_X$ .

## How do we impose the interface boundary conditions?

To calculate the coefficients  $c^X$  we need  $N = N_L + N_R$ compatibility conditions at the interface. Where  $N_X$  is the number of waves between the initial and star state.

It is through these compatibility conditions that we impose the correct physical boundary conditions at the interface.

We assume that these compatibility conditions take the form:

$$\Delta w_j = w_j^{*R} - w_j^{*L} = 0, \quad j = 1, \dots, N.$$
 (6)

This gives a linear system which we can easily solve,

$$\mathbf{w}_{R} - \mathbf{w}_{L} = \sum_{j=1}^{N_{R}} c_{j}^{R} \mathbf{r}_{R}^{(j)} + \sum_{i=1}^{N_{L}} c_{i}^{L} \mathbf{r}_{L}^{(i)}.$$
 (7)



## Outline

- Results



# Vorticity Propagation

# Toy Star

## Outline

- Conclusion



