

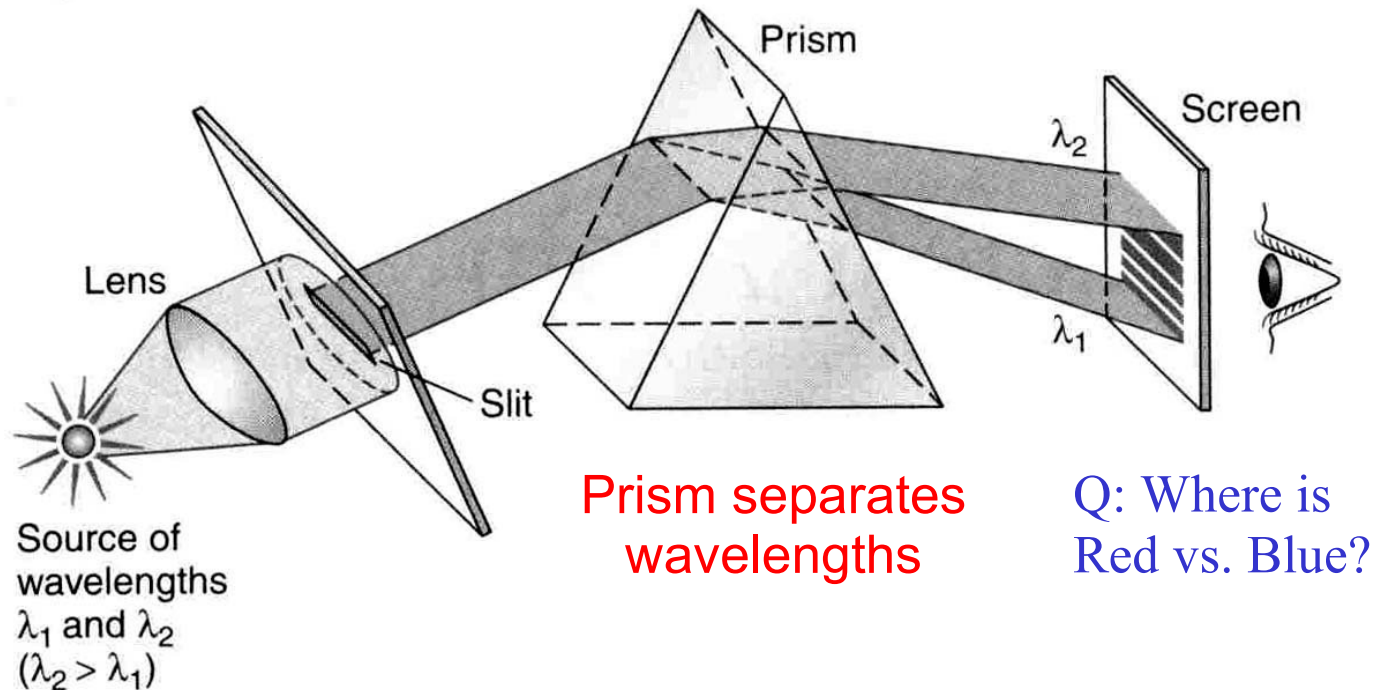
# **The Interaction of Light and Matter**

# Outline

- (1) Motivation: **Why spectral lines?**
  - the Birth of Spectroscopy
  - Kirchoff's Laws
- (2) **Photons – the particle nature of light**
  - Blackbody radiation (Planck introduces quantum of light)
  - Photoelectric Effect
  - Compton Scattering
- (3) **The Bohr Model of the Atom**
  - a theory to describe spectral lines,
- (4) **Quantum Mechanics and the Wave-Particle Duality (SKIP on ExamI)**
  - De Broglie wavelength
  - Schrodinger's probability waves.

# Spectroscopy - history

- Trogg (50 million BC) – rainbow
- Newton (1642-1727) – decomposes light into spectrum and back again
- W. Herschel (1800) – discovers infrared
- J. W. Ritter (1801) – discovers ultraviolet
- W. Wollaston (1802) – discovers absorption lines in solar spectrum



# Spectroscopy - history

- J. Herschel, Wheatstone, Alter, Talbot and Angstrom studied spectra of terrestrial things (flames, arcs and sparks) ~1810
- Joseph Fraunhofer
  - Cataloged ~475 dark lines of the solar spectrum by 1814
  - Identifies sodium in the Sun from flame spectra in the lab!
  - Looks at other stars (connects telescope to spectroscope)
- Foucault (1848) – sees absorption lines in sodium flame with bright arc behind it.

There is the need for a *new physics!*

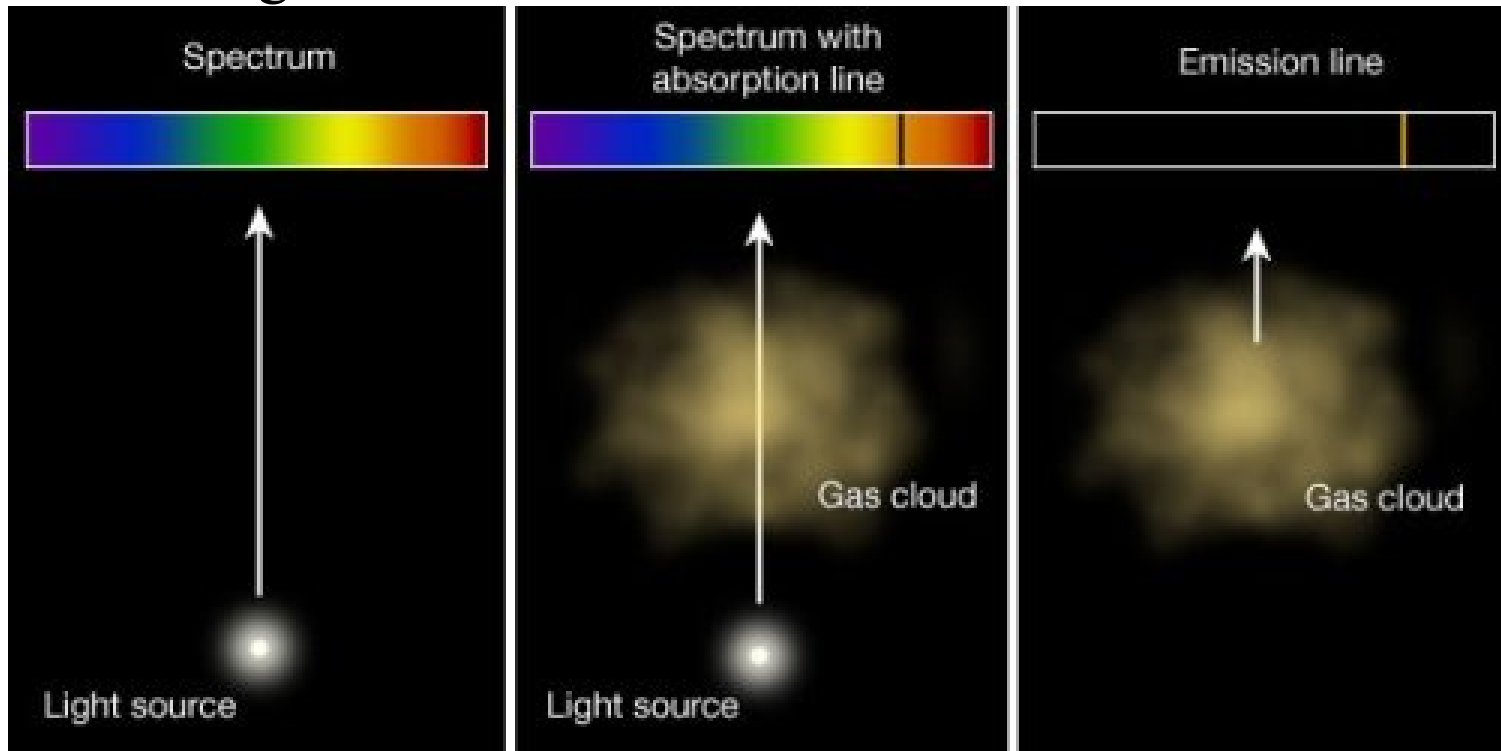


# Kirchhoff's laws (1859):

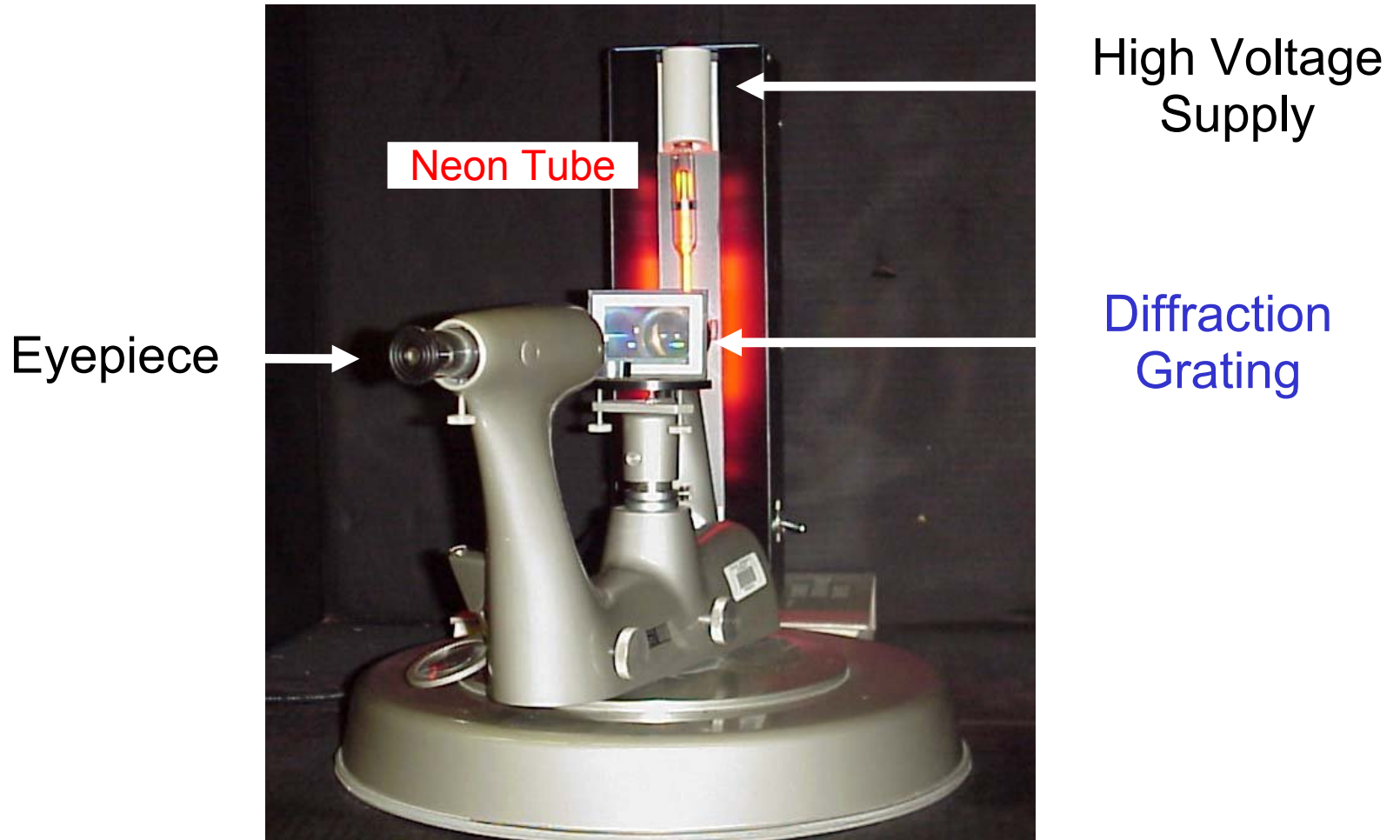
Worked with Bunsen on flame spectra

Developed a prism spectroscope

- Hot solid or dense gas, → **continuous spectrum (eg Blackbody)**
- Cool diffuse gas in front of a blackbody → **absorption lines**
- Hot diffuse gas → **emission lines**



# Atomic Spectra Lab



$$d \sin \theta = n\lambda$$

# Doppler shift

- Spectral lines allow for the measurement of radial velocities

$$z = \frac{\lambda_{obs} - \lambda_{rest}}{\lambda_{rest}} = \frac{\Delta\lambda}{\lambda_{rest}} = \frac{v_r}{c}$$

- At low velocities,  $v_r \ll c$

- Classical Doppler effect

- *Radial velocity*,  $v_r$
- *Heliocentric correction* for Earth's motion, up to  $\sim 29.8$  km/s, depending on direction.

$$\Delta\lambda = \frac{v_r}{c} \lambda_{rest}$$

- Example:  $H_\alpha$  is  $6562.80 \text{ \AA}$

- Vega is measured to be  $6562.50 \text{ \AA}$
- Coupled with the *proper motion*
  - Can determine total velocity

$$v_r = c \frac{\Delta\lambda}{\lambda_{rest}} = -14 \frac{\text{km}}{\text{sec}}$$

$$v_\theta = r\mu = 13 \frac{\text{km}}{\text{s}}$$

$$v = \sqrt{v_r^2 + v_\theta^2} = 19 \frac{\text{km}}{\text{s}}$$

# Doppler shift

- Since most galaxies are moving away, astronomers call the Doppler shift a *redshift*,  $z$ .
- At high velocities,  $v_r \lesssim c$ 
  - Relativistic redshift parameter (Ch. 4):

$$z = \frac{\Delta\lambda}{\lambda_{rest}}$$

$$z = \sqrt{\frac{1 + v_r/c}{1 - v_r/c}} - 1$$

- Example: Prob. 4.8.  
(should get:  $v_r = 0.9337c$ )



# Particle/Wave Duality - Part 1

## PART 1

- Electrons as discrete Particles
  - Measurement of  $e/m$  (CRT) and  $e$  (oil-drop expt.)
- Photons as discrete Particles
  - **Blackbody Radiation**: Planck's spectrum required quantization
  - **Photoelectric Effect**: Photons “kick out” Electrons from metals
  - **Compton Effect**: Photon *scatters* off Electron

## PART 2

- Wave Behavior: Diffraction and Interference
- Photons as Waves:  $\lambda = hc / E$ 
  - X-ray Diffraction (Bragg's Law)
- Electrons as Waves:  $\lambda = h / p$ 
  - Low-Energy Electron Diffraction (LEED)

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# Photons: Quantized Energy Particle

- Light comes in discrete energy “packets” called **photons**

Energy of  
Single Photon

$$E = h\nu = \frac{hc}{\lambda}$$

From Relativity:

$$E^2 = (pc)^2 + (mc^2)^2 \leftarrow \text{Rest mass}$$

For a Photon ( $m = 0$ ):

$$E^2 = (pc)^2 + 0 \Rightarrow E = pc$$

Momentum of  
Single Photon

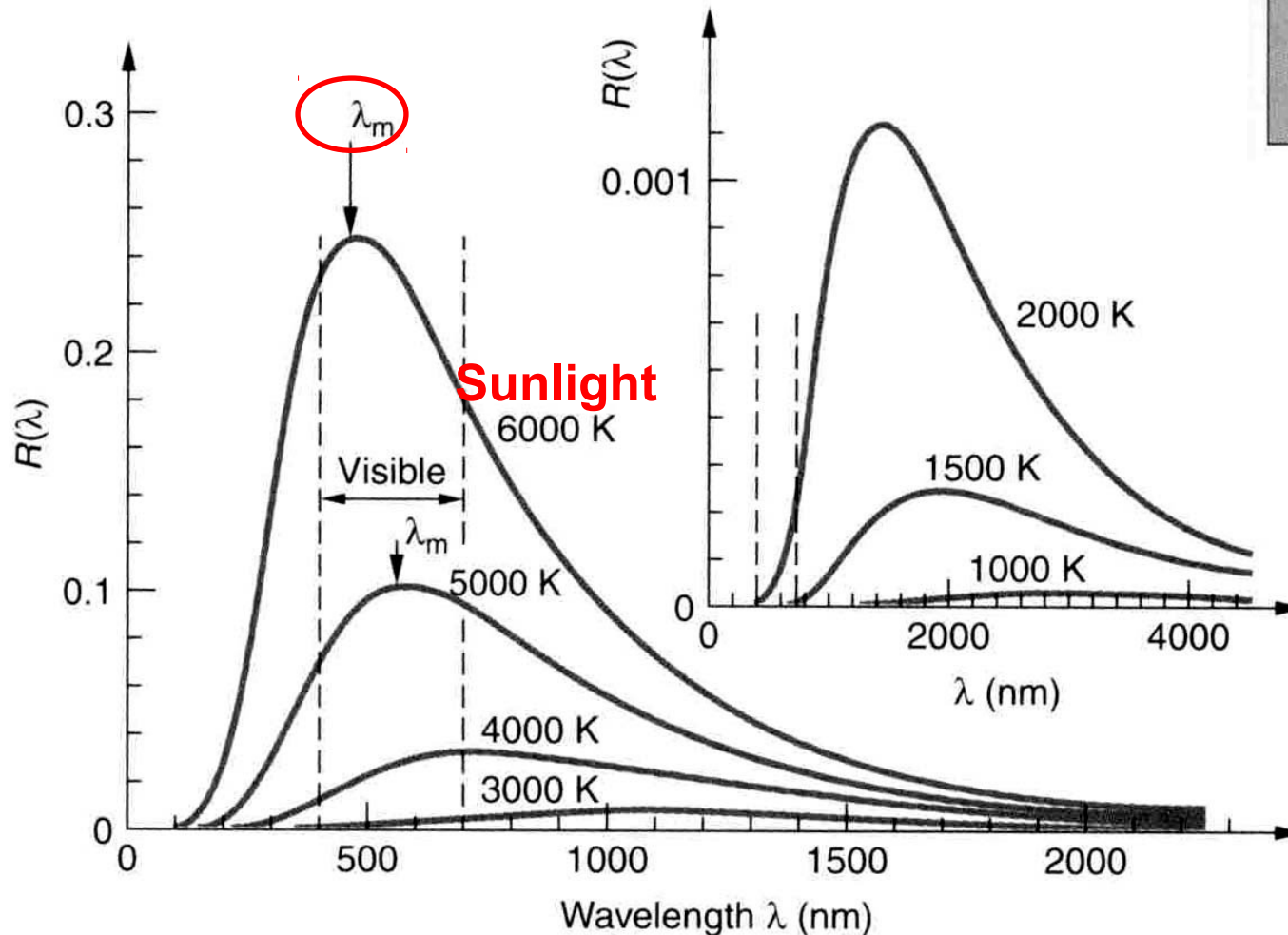
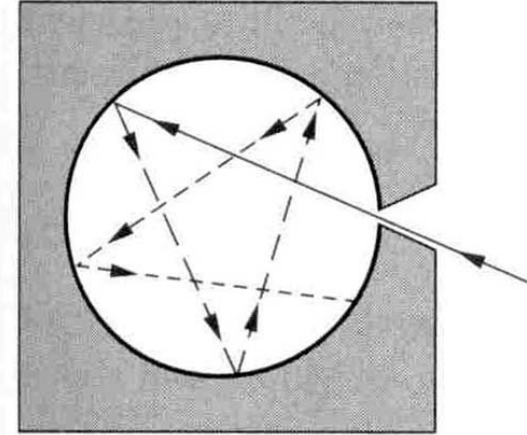
$$p = \frac{E}{c} = \frac{hc}{\lambda c} = \frac{h}{\lambda}$$

# Blackbody Radiation: First clues to quantization

Recall Wien's Law:

$$\lambda_{\max} = \frac{0.029}{T} \text{ cm} \cdot \text{K}$$

and the Stefan-Boltzmann Law:  $F = \sigma T^4$



Spectral  
Distribution  
depends only  
on Temperature

# Spectral Blackbody: Planck's Law

- Planck's Law was found empirically (trial and error!)
- Quantize the E&M radiation so that the minimum

energy for light at a given wavelength is:

$$E_{\nu} = h\nu = hc/\lambda$$

where  $h$  = Planck's Constant =  $6.626 \times 10^{-34}$  J·s.

Then  $E_{\nu} = nh\nu$ ,  $n=0, 1, 2, 3$

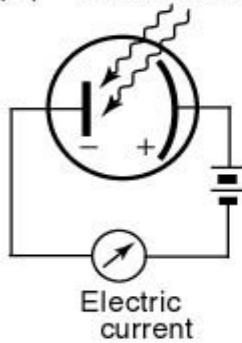
can be used in replacing the classical  $kT$  expression for the average energy in a mode.

Now the entire hot object may not have enough energy to emit one photon of light at very small wavelengths, so  $n=0$ , and the UV catastrophe can be avoided.

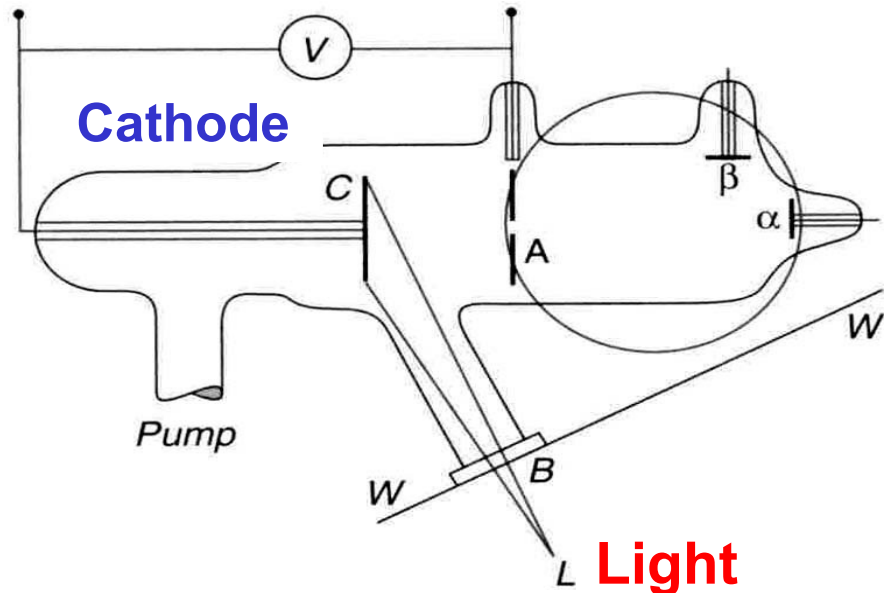
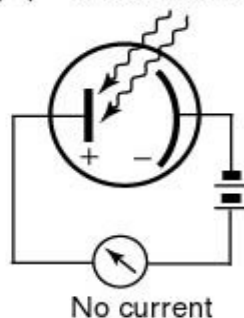
# Photoelectric Effect: “Particle Behavior” of Photon

- Shows quantum nature of light (Theory by Einstein & Expt. by Millikan).
- Photons** hit metal cathode and instantaneously eject **electrons** (requires minimum energy = work function).
- Electrons travel from cathode to anode against **retarding voltage  $V_R$**
- Electrons collected as **“photoelectric” current** at anode.
- Photocurrent becomes zero when retarding voltage  $V_R$  equals the **stopping voltage  $V_{stop}$** , i.e.  $eV_{stop} = K_e$

(A) Ultraviolet rays



(B) Ultraviolet rays



The C plate is always a source of  $e^-$ . However, a voltage can be applied that makes it positive relative to the A plate.

# Photoelectric Effect - equation

- **PHOTON IN**  $\Rightarrow$  **ELECTRON OUT**

- $e^-$  kinetic energy = Total photon energy  
–  $e^-$  ejection energy

$$K_{\max} = h\nu - \phi$$

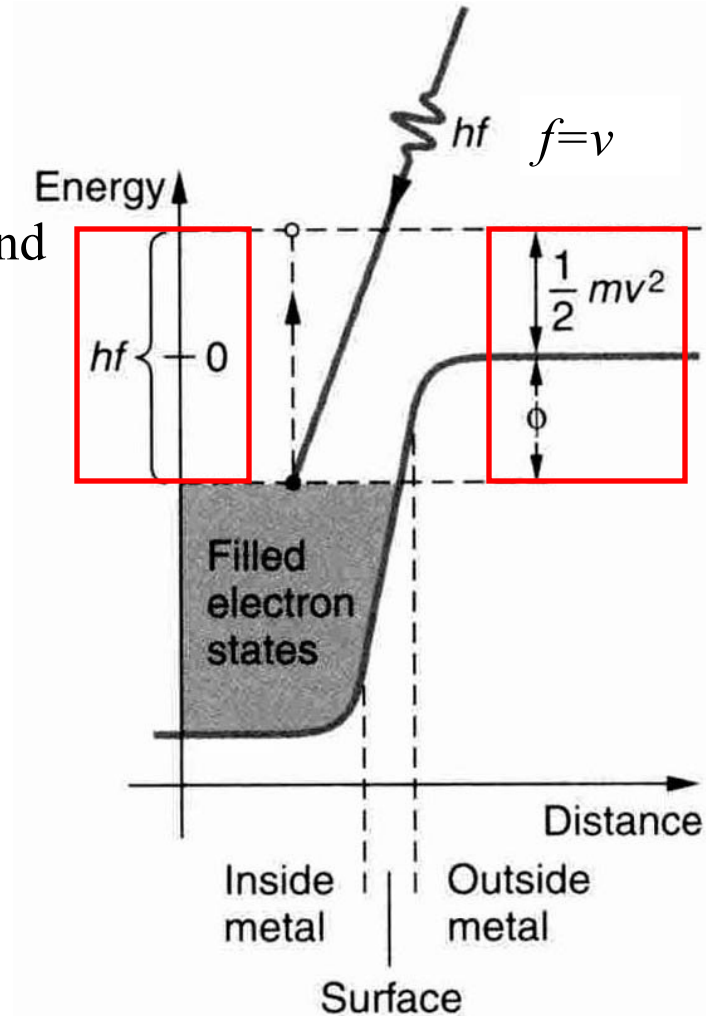
- where  $h\nu$  = photon energy,  $\phi$  = work function, and  $K_{\max}$  = kinetic energy

- $K_{\max} = eV_{\text{stop}}$  = stopping energy

- **Special Case:** No kinetic energy ( $V_0 = 0$ )

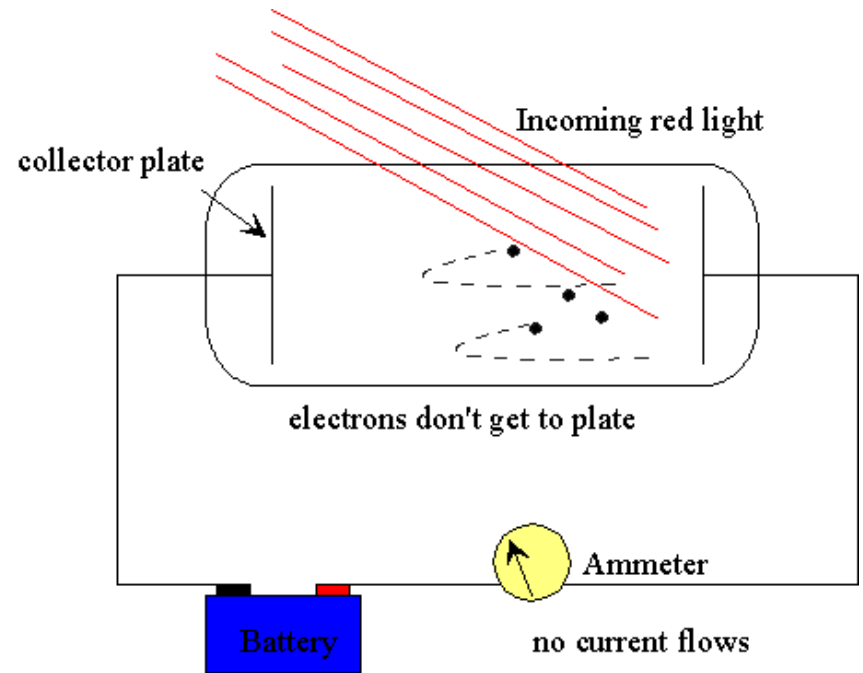
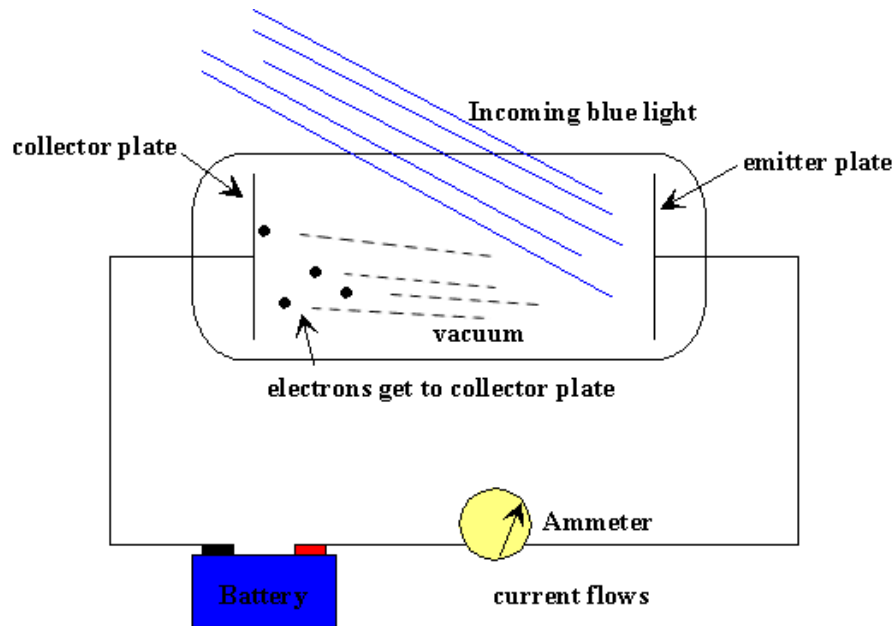
- Minimum frequency  $\nu$  to eject electron

$$h\nu_{\min} = \phi$$



# Photoelectric Effect

- In order to make electrons reach the collector plate, the light has to be “blue enough”; the intensity doesn't matter if light is red!

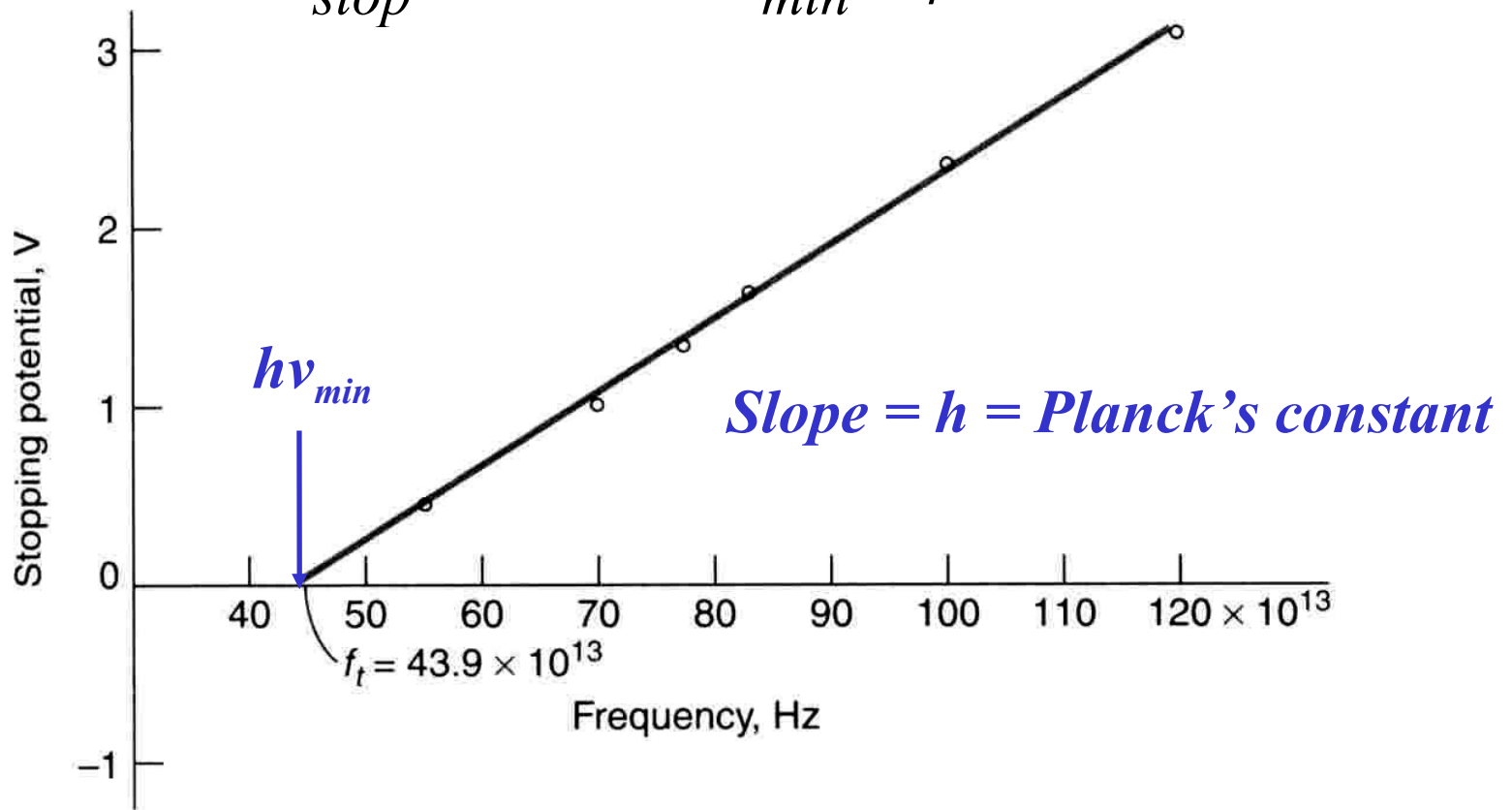




# Photoelectric Effect: $V_{\text{stop}}$ vs. Frequency

$$eV_{\text{stop}} = h\nu - \phi$$

$$V_{\text{stop}} = 0 \Rightarrow h\nu_{\text{min}} = \phi$$



# Photoelectric Effect Problem

If the work function of a metal is 2.0 eV,

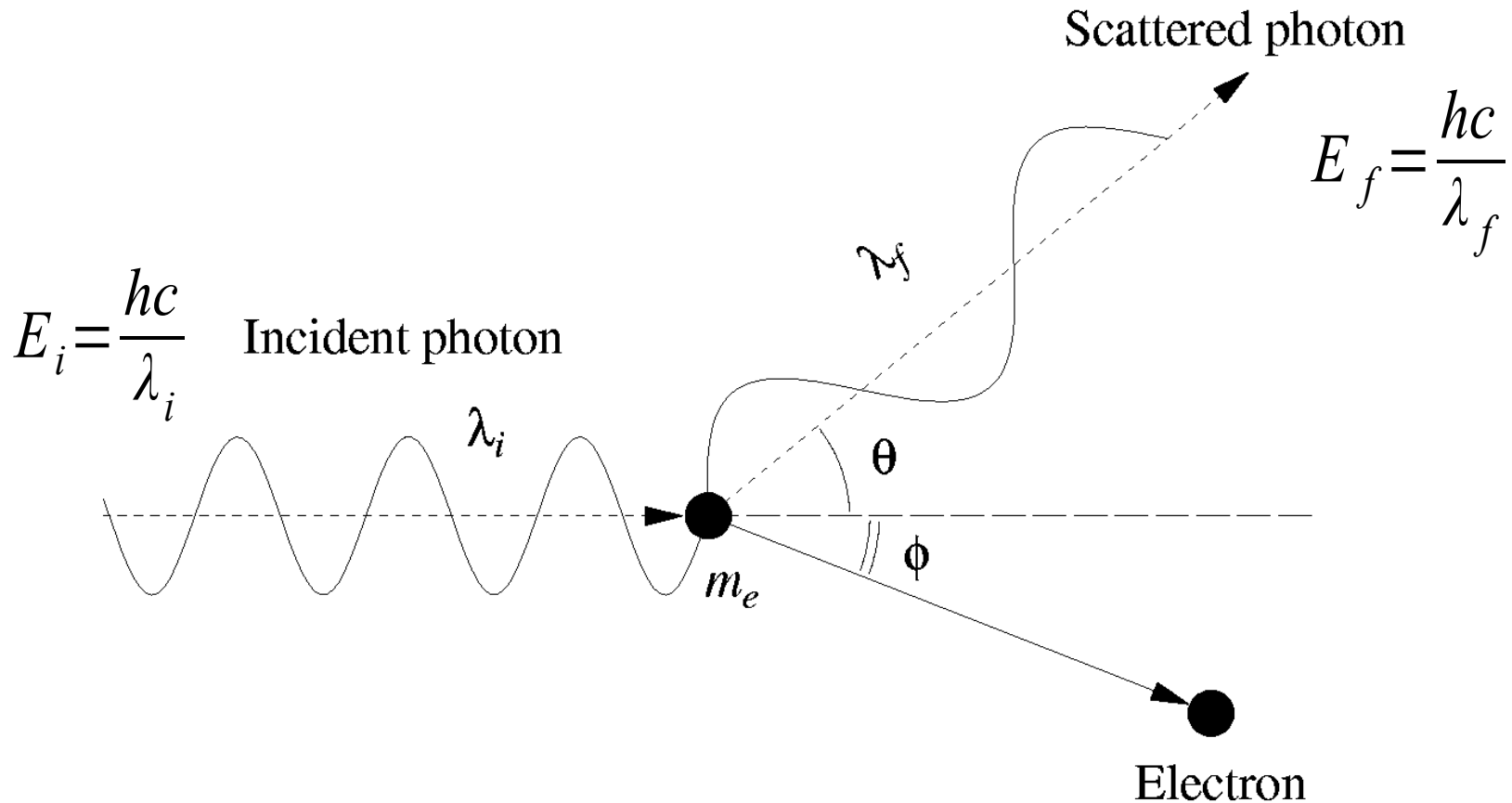
a) find the maximum wavelength  $\lambda_m$  capable of causing the photoelectric effect, and,

b) find the stopping potential if  $\lambda = \lambda_m / 2$

# Compton Scattering: “Particle-like” Behavior of Photon

**Concept:** Photon scatters off electron losing energy and momentum to the electron. The  $\lambda_f$  of scattered photon depends on  $\theta$

- Conservation of relativistic momentum and Energy!
- No mass for the photon but it has momentum!!!



# Compton Scattering: Equation

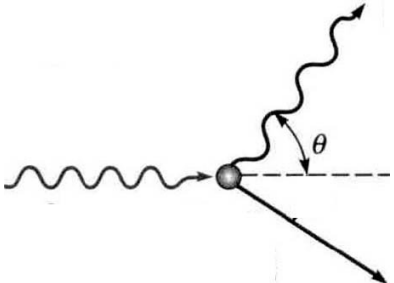
Photon OUT

Scattering Angle

$$\Delta\lambda = \lambda_f - \lambda_i = \frac{h}{m_e c} (1 - \cos \theta)$$

Photon IN

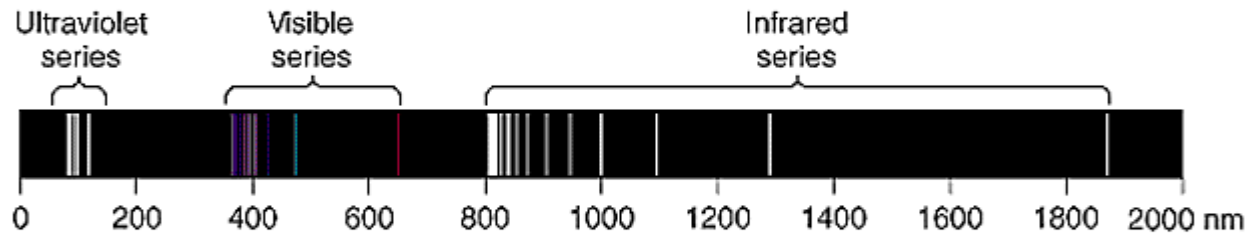
Compton wavelength  
 $\lambda_c = 0.0024 \text{ nm for } e^-$



- **Limiting Values**
  - No scattering:  $\theta = 0^\circ \rightarrow \cos 0^\circ = 1 \rightarrow \Delta\lambda = 0$
  - “Bounce Back”:  $\theta = 180^\circ \rightarrow \cos 180^\circ = -1 \rightarrow \Delta\lambda = 2\lambda_c$
- Difficult to observe unless  $\lambda$  is small (i.e.  $\Delta\lambda/\lambda > 0.01$ )

# Atomic Spectra

- 1885 - **Balmer** observed Hydrogen Spectrum
  - Found empirical formula for discrete wavelengths
  - Later generalized by Rydberg for simple ionized atoms



$$\frac{1}{\lambda} = R_H \left( \frac{1}{2^2} - \frac{1}{n^2} \right) \text{ with } 2 < n$$

# Atomic Spectra: Rydberg Formula

$$\frac{1}{\lambda} = R_H \left( \frac{1}{m^2} - \frac{1}{n^2} \right) \quad \textbf{with } m < n$$

- Gives  $\lambda$  for any lower level  $m$  and upper level  $n$  of Hydrogen.
- Rydberg constant  $R_H \sim 1.097 \times 10^7 \text{ m}^{-1}$
- $m = 1$  (Lyman),  $2$  (Balmer),  $3$  (Paschen)

- Example for  $n = 2$  to  $m = 1$  transition:

$$\frac{1}{\lambda} = R_H \left( \frac{1}{1^2} - \frac{1}{2^2} \right) = \frac{3}{4} (1.097 \times 10^7 \text{ m}^{-1})$$

$\Rightarrow \lambda = 121.6 \text{ nm}$  *Ultraviolet*

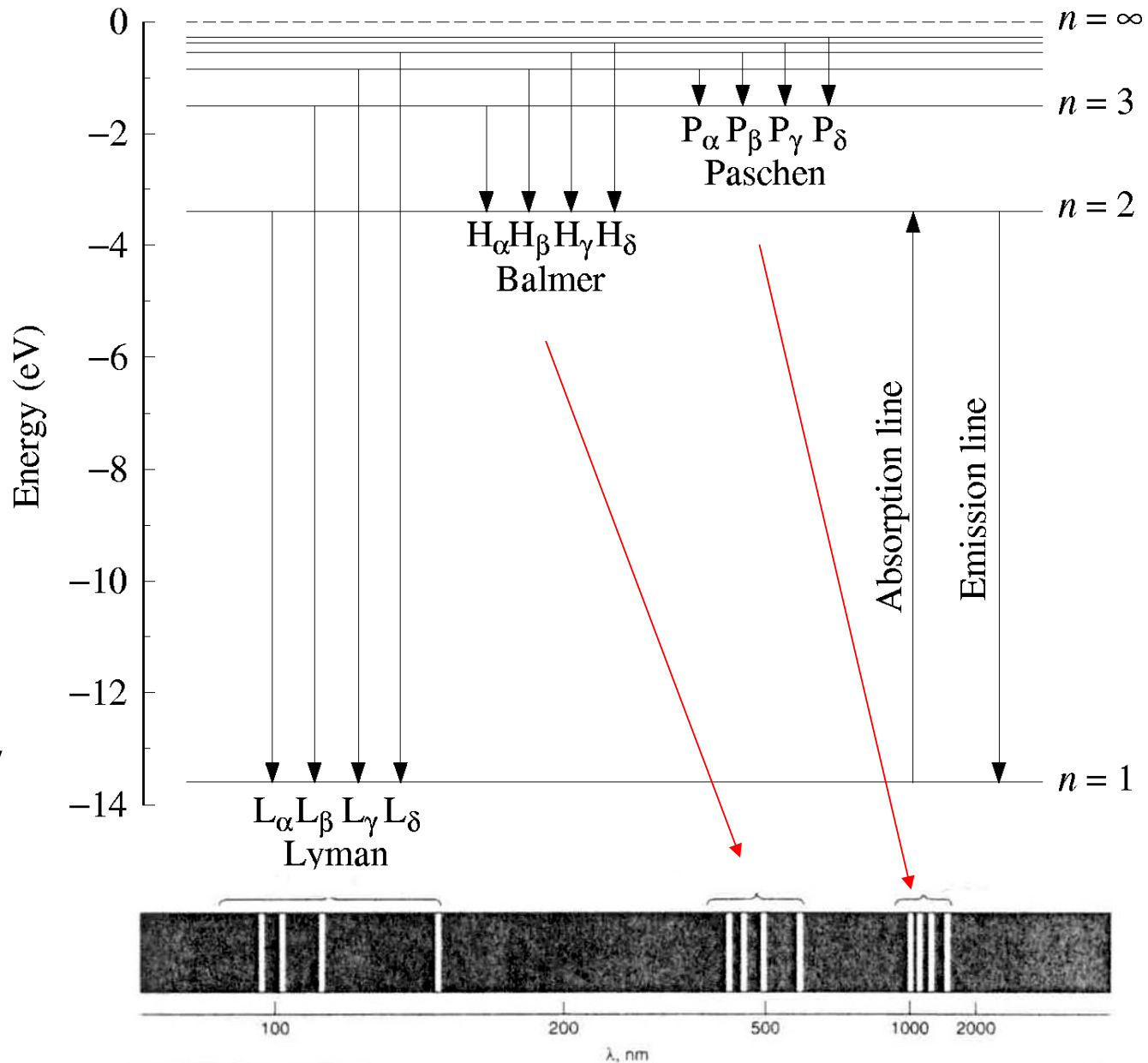
# Atomic Spectra: Hydrogen Energy Levels

$$E_{\infty} = 0 \text{ eV}$$

Energy

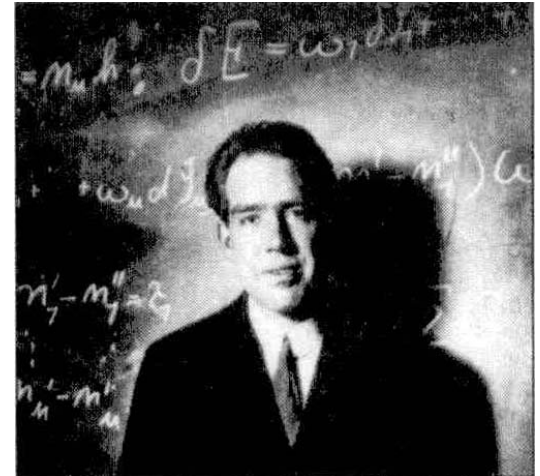
$$E_n \propto \frac{-1}{n^2}$$

$$E_1 = -13.6 \text{ eV}$$



# Bohr Model

- 1913 – **Bohr** proposed **quantized model** of the H atom to predict the observed spectrum.
- **Problem:** Classical model of the electron “orbiting” nucleus is unstable. Why unstable?
  - Electron experiences (centripetal) acceleration.
  - Accelerated electron emits radiation.
  - Radiation leads to energy loss.
  - Electron quickly “crashes” into nucleus.





# Bohr Model: Quantization

- **Solution:** Bohr proposed two “quantum” postulates
  - Electrons exist in stationary orbits (no radiation) with quantized angular momentum.

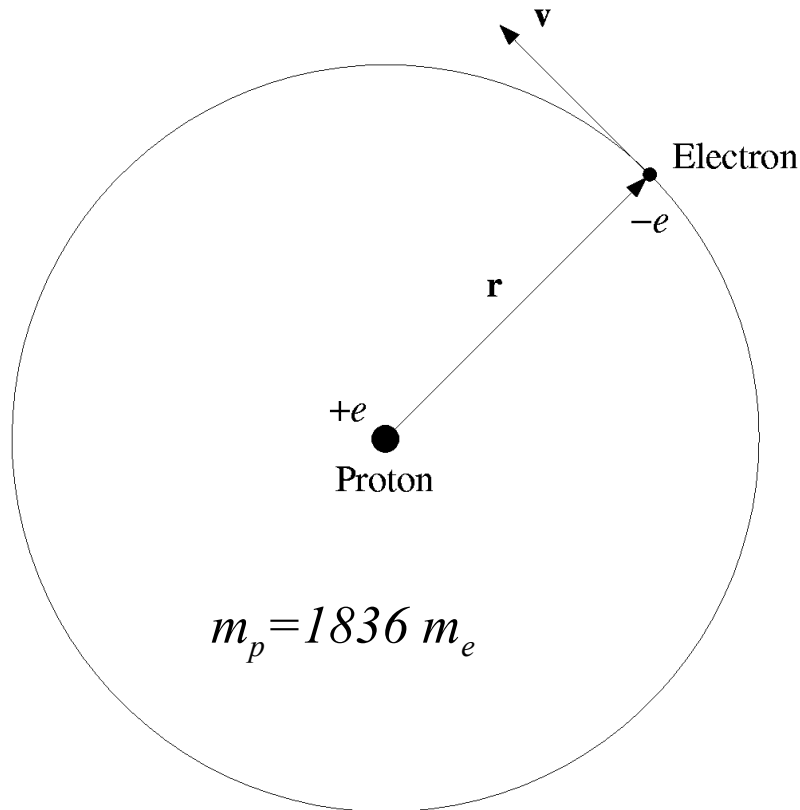
$$L_n = mvr = n \hbar \quad \left( \hbar = \frac{h}{2\pi} = 6.58 \times 10^{-16} \text{ eV} \cdot \text{s} \right)$$

- Atom radiates with quantized frequency  $\nu$  (or energy E) only when the electron makes a transition between two stationary states.

$$h\nu = \frac{hc}{\lambda} = E_i - E_f$$

# Planetary Mechanics Applied to the H Atom

- Consider the attractive electrostatic force and circular motion



$$\mu = \frac{m_e m_p}{M} \approx m_e$$

$$M = m_e + m_p \approx m_p$$

$$\vec{F} = \frac{q_1 q_2}{r^2} \hat{r} = \mu \frac{v^2}{r} \hat{r}$$

Note: in cgs,  $e = 4.803 \times 10^{-10}$  esu

$$\frac{q_1 q_2}{r^2} = -\mu \frac{v^2}{r}$$

$$\frac{-e^2}{r^2} = -\mu \frac{v^2}{r}$$

$$\frac{1}{2} \mu v^2 = \frac{1}{2} \frac{e^2}{r} = K$$

$$U = -2K = -\frac{e^2}{r}$$

Kinetic energy

Potential energy

# Planetary Mechanics Applied to the H Atom

- Introduce Bohr's quantized angular momentum  $L = \mu v r = n \hbar$  (wrong)

$$K = \frac{1}{2} \frac{e^2}{r} = \frac{1}{2} \mu v^2 = \frac{1}{2} \frac{(\mu v r)^2}{\mu r^2} = \frac{1}{2} \frac{(n \hbar)^2}{\mu r^2}$$

- Solving for  $r$

$$r_n = \frac{\hbar^2}{\mu e^2} n^2 = a_0 n^2$$

$a_0$  is the Bohr radius

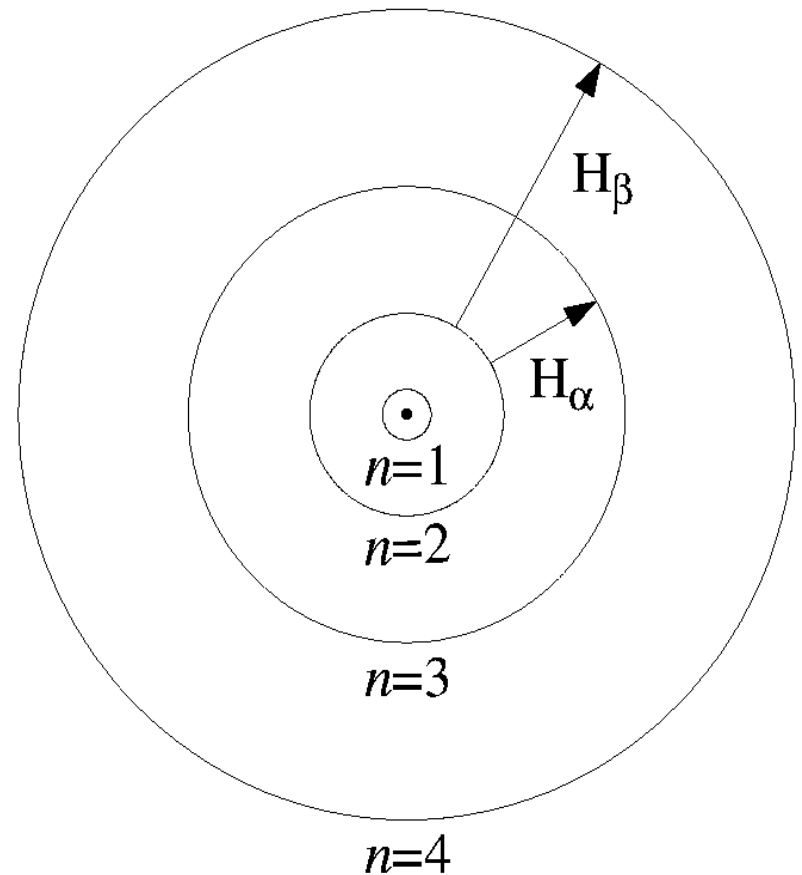
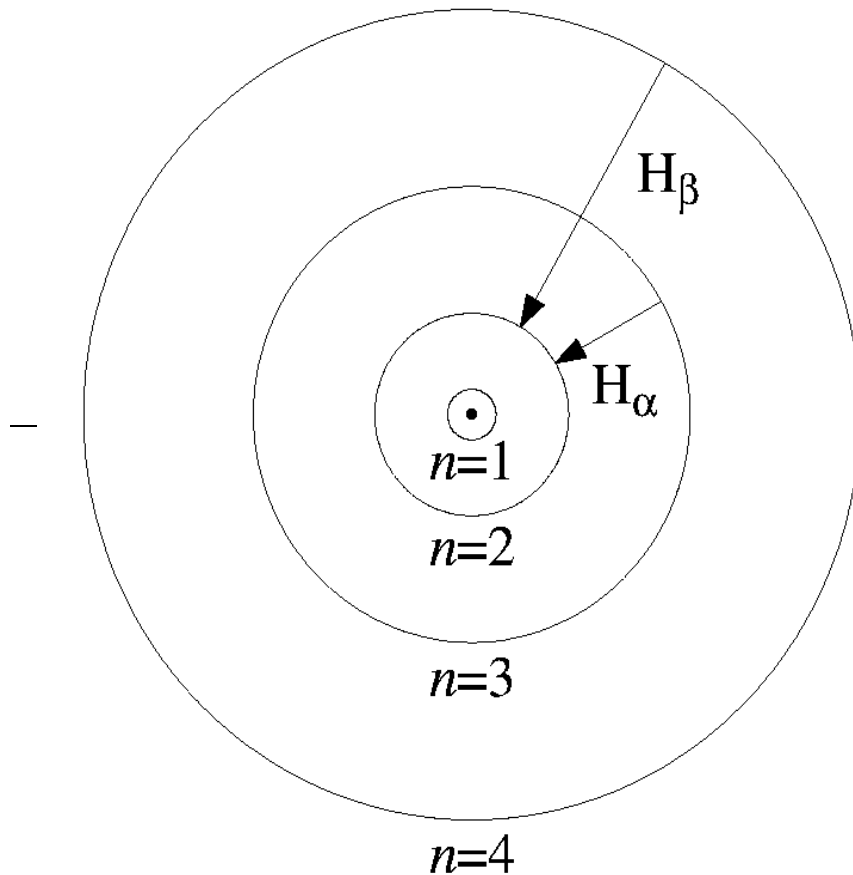
- Get the Total Energy in terms of  $n$ . (Recall  $E_{tot} = \langle U \rangle / 2$ )

$$E_n = -\frac{1}{2} \frac{e^2}{r} = -\frac{\mu e^4}{2 \hbar^2} \frac{1}{n^2} = \frac{-13.6 \text{ eV}}{n^2} = \frac{-E_0}{n^2}$$

- Principle quantum number,  $n = 1, 2, 3, \dots$

# Bohr Model: Transitions

- Transitions predicted by Bohr yield **general Rydberg formula**



# Bohr Model Problem: Unknown Transition

If the wavelength of a transition in the **Balmer series** for a **He<sup>+</sup>** atom is **121 nm**, then find the corresponding transition, i.e. initial and final n values.

$$\frac{1}{\lambda} = RZ^2 \left( \frac{1}{n_f^2} - \frac{1}{n_i^2} \right) = R(2)^2 \left( \frac{1}{(2)^2} - \frac{1}{n_i^2} \right)$$

where  $Z = 2$  for He and  $n_f = 2$  for Balmer

$$\frac{1}{4R\lambda} = \left( \frac{1}{4} - \frac{1}{n_i^2} \right)$$

$$n_i = \left( \frac{1}{4} - \frac{1}{4R\lambda} \right)^{-1/2} = \left( \frac{1}{4} - \frac{1}{4(1.1 \times 10^7 \text{ m}^{-1})(121 \times 10^{-9} \text{ m})} \right)^{-1/2} = \underline{4}$$

# Bohr Model Problem: Ionization Energy

Suppose that a He atom ( $Z=2$ ) in its ground state ( $n = 1$ ) absorbs a photon whose wavelength is  $\lambda = 41.3 \text{ nm}$ . Will the atom be **ionized**?

➤ *Find the energy of the incoming photon and compare it to the ground state ionization energy of helium, or  $E_0$  from  $n=1$  to  $\infty$ .*

$$E = \frac{hc}{\lambda} = \frac{1240 \text{ eV nm}}{41.3 \text{ nm}} = \underline{30 \text{ eV}}$$

$$E_0(\text{He}) = Z^2 \times E_0(\text{H}) = (2^2)(13.6 \text{ eV}) = 54.4 \text{ eV}$$

➤ *The photon energy (30 eV) is less than the ionization energy (54 eV), so the electron will NOT be ionized.*

# Bohr Model Problem: Series Limit (book)

Find the **shortest wavelength** that can be emitted by the **Li<sup>++</sup> ion**.

➤ *The shortest  $\lambda$  (or highest energy) transition occurs for the highest initial state ( $n_i = \infty$ ) to the lowest final state ( $n_f = 1$ ).*

$$\frac{1}{\lambda} = RZ^2 \left( \frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

where  $Z = 3$  for Li,  $n_i = \infty$ , and  $n_f = 1$  for shortest  $\lambda$

$$\frac{1}{\lambda} = (1.1 \times 10^7 \text{ m}^{-1}) (3)^2 \left( \frac{1}{(1)^2} - \frac{1}{(\infty)^2} \right) = 10.1 \text{ nm}$$

# Particle/Wave Duality - Part 2

## PART 1

- Electrons as discrete Particles
  - Measurement of  $e/m$  (CRT) and  $e$  (oil-drop expt.)
- Photons as discrete Particles
  - **Blackbody Radiation**: Temp. Relations & Spectral Distribution
  - **Photoelectric Effect**: Photon “kicks out” Electron
  - **Compton Effect**: Photon “scatters” off Electron

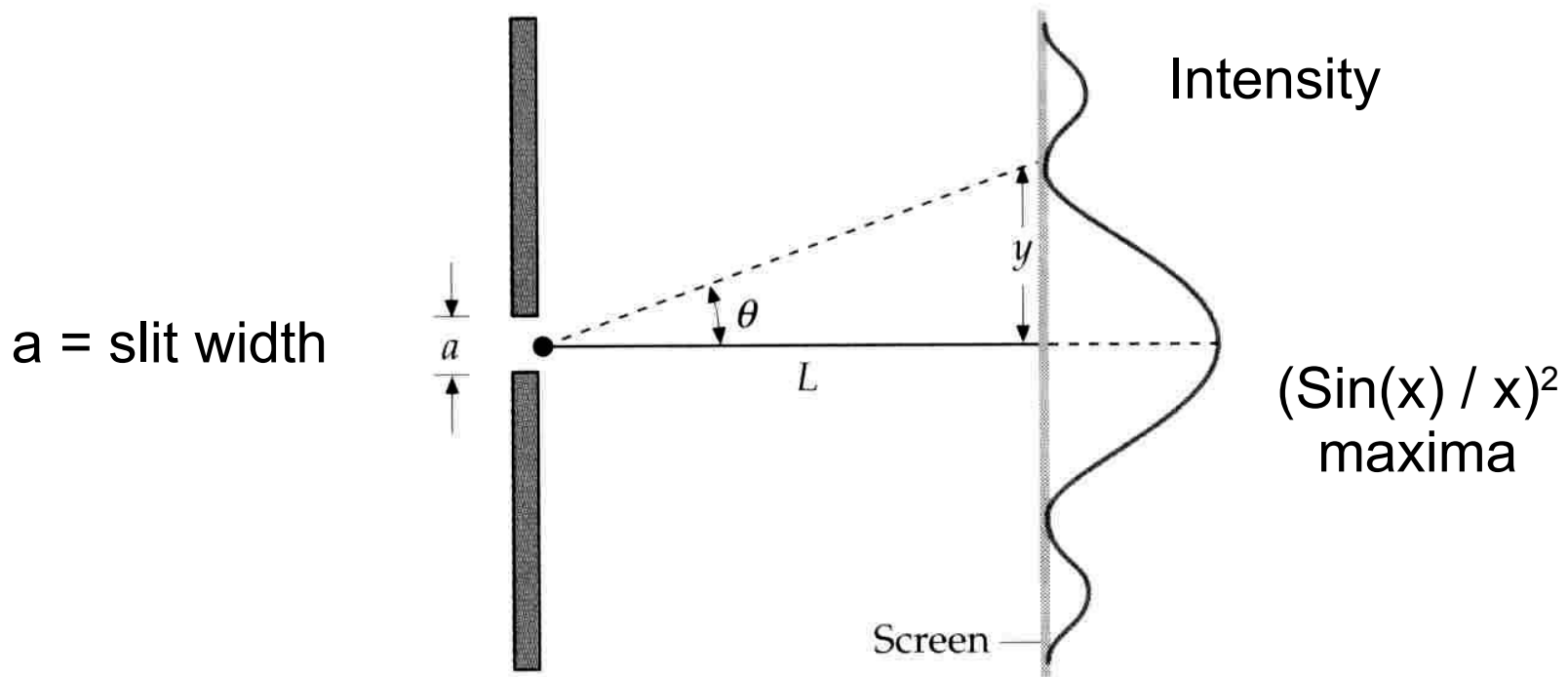
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- **Wave Behavior**: Diffraction and Interference
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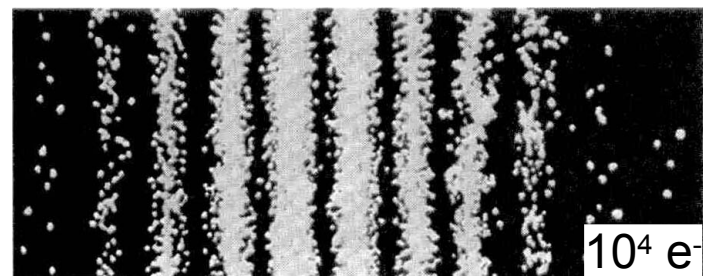
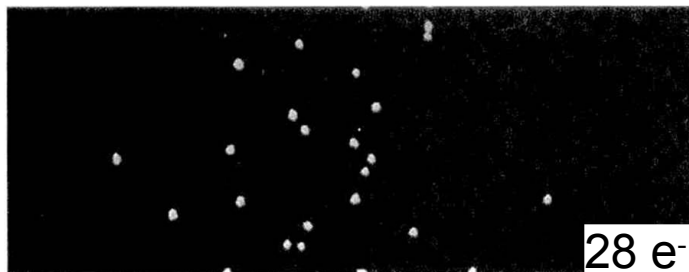


# Wave Property: Single-Slit Diffraction

$$\text{Minima: } n\lambda = a \sin \theta$$



Diffraction Pattern of **Electron Waves**



# Electrons: Wave-like Behavior

$$\lambda = \frac{h}{p}$$

- Every particle has a wavelength given by:
- Question: Why don't we **observe effects** of particle waves (i.e., diffraction and interference) in day-to-day life?
- Answer: Wavelengths of most macroscopic objects are **too small** to interact with slits, BUT atomic-sized objects DO behave like waves!

**Macroscopic** – ping pong ball

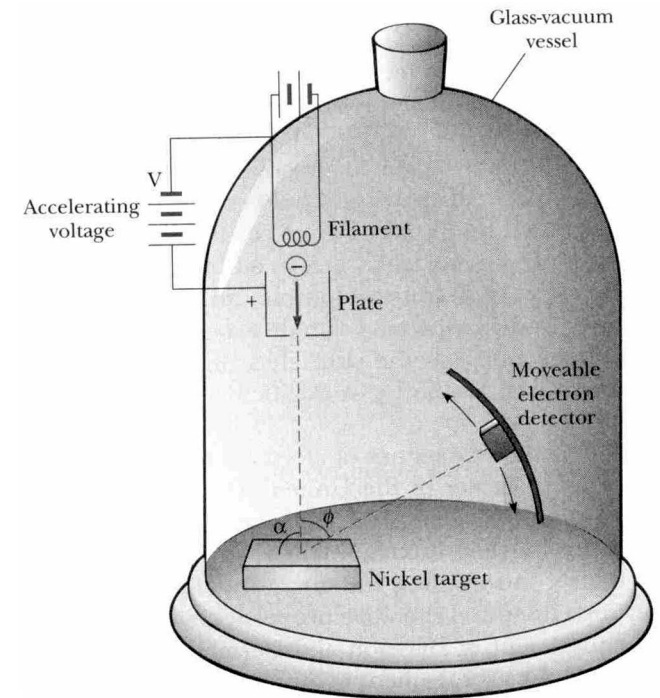
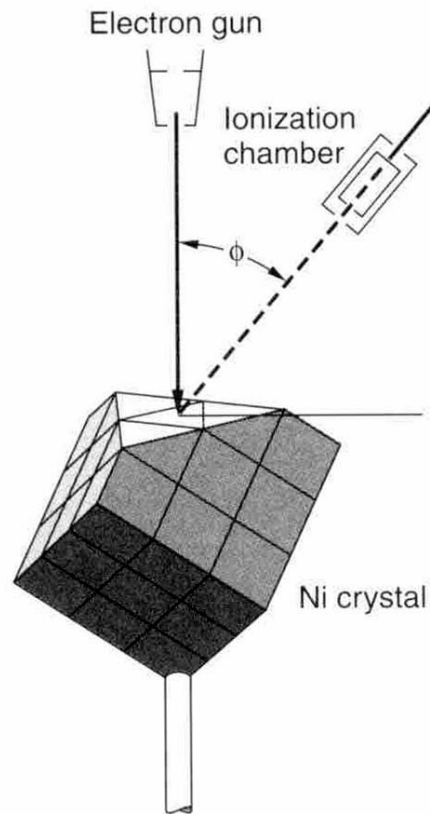
$$\lambda = \frac{h}{mv} = \frac{6.63 \times 10^{-34} \text{ J} \cdot \text{s}}{(2 \times 10^{-3} \text{ kg})(5 \text{ m/s})} = 6.6 \times 10^{-32} \text{ m} \text{ (immeasurably small! )}$$

**Microscopic** – “slow electron” (1% speed of light)

$$\lambda = \frac{h}{mv} = \frac{6.63 \times 10^{-34} \text{ J} \cdot \text{s}}{(9.1 \times 10^{-31} \text{ kg})(10^6 \text{ m/s})} = 7.3 \times 10^{-10} \text{ m (atomic dimension )}$$

# Electron Diffraction: Wave-like Behavior

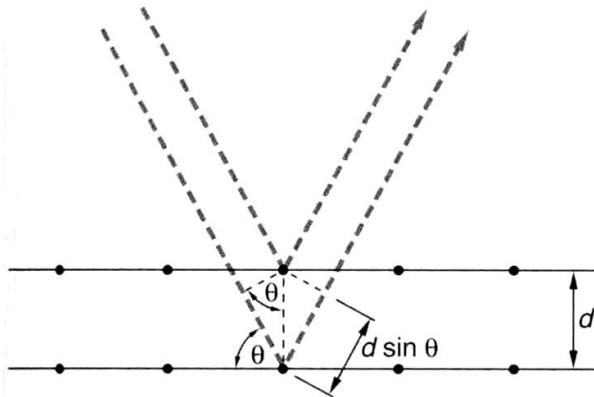
- 1927 – Davisson and Germer studied the diffraction of an electron beam from a nickel crystal surface and observed discrete spots (maxima).
- Modern day technique now: Low Energy Electron Diffraction (LEED).



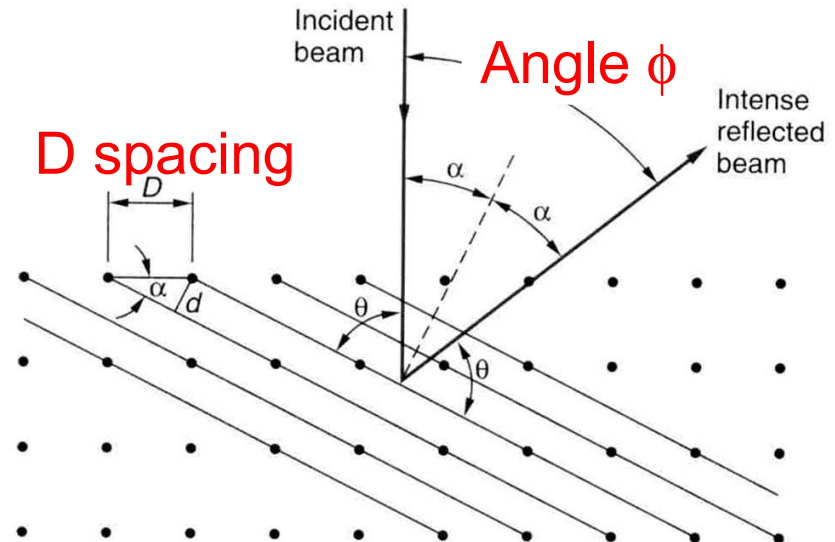
# Electron Diffraction: LEED Equation

**Concept:** Use [Bragg's Law](#) for X-ray scattering and then substitute appropriate angles, where  $\lambda$  is now the [electron](#) wavelength.

## X-ray Diffraction



## Electron Diffraction

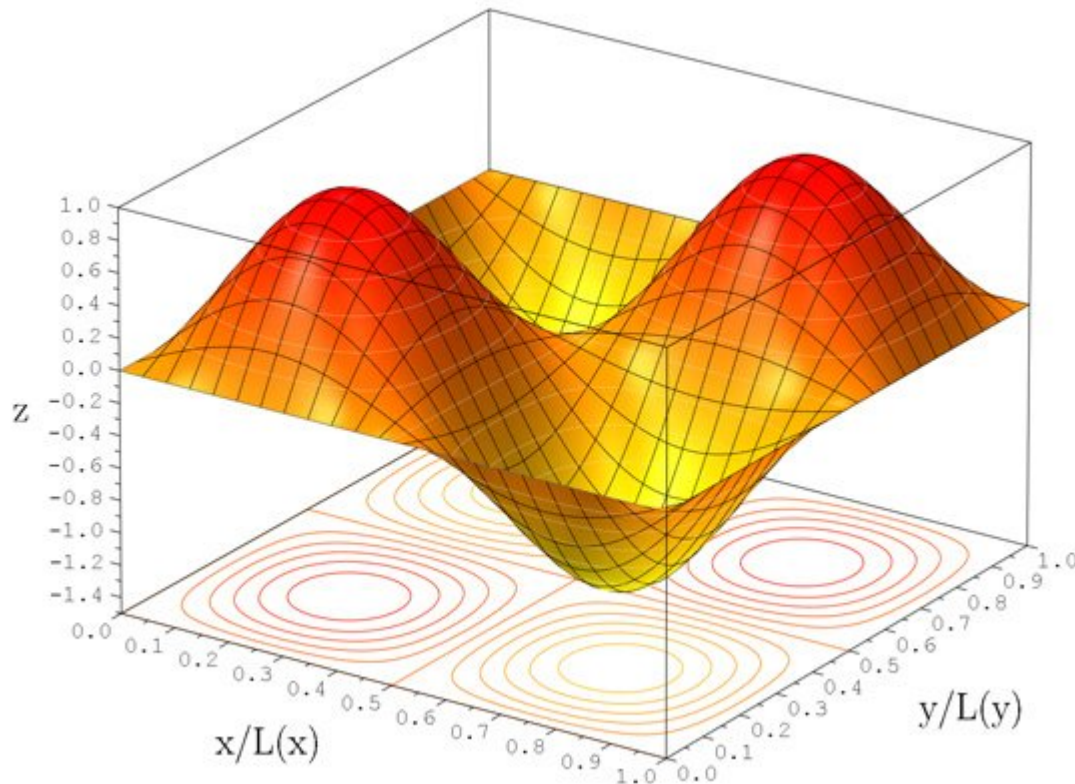


$$n\lambda = \underbrace{2d}_{D \sin \alpha} \underbrace{\sin \theta}_{\cos \alpha} = 2D \underbrace{\sin \alpha \cos \alpha}_{\frac{1}{2} \sin 2\alpha \text{ by trig}} = D \sin 2\alpha$$

$$n\lambda = D \sin 2\alpha = D \sin \phi$$

# Wave/Particle Duality

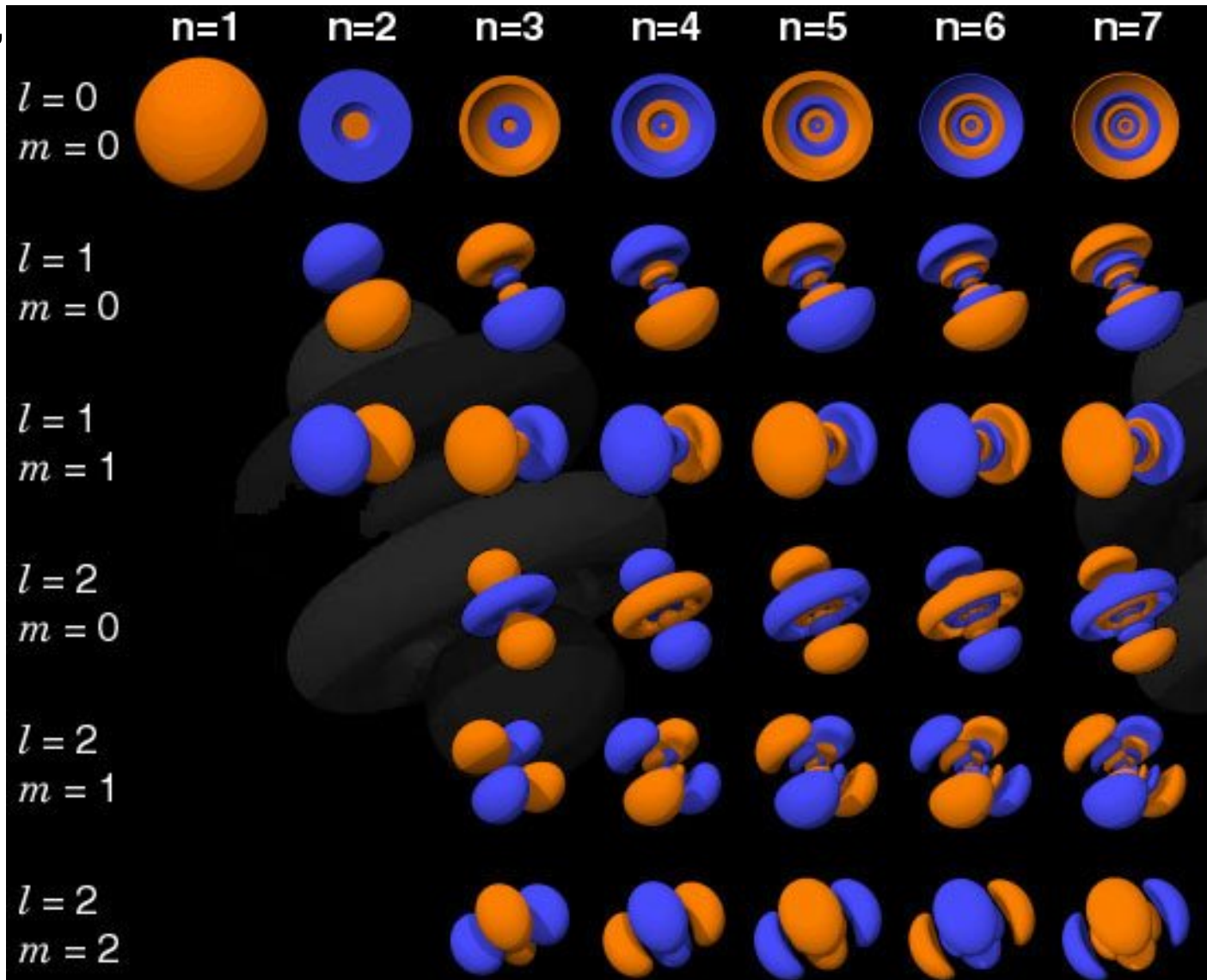
- The particle wavefunction,  $\psi$ , is the “probability amplitude” (see figure “Z”), a complex number.
- Probability density =  $|\psi|^2$  gives the probability of where we might find the particle. ( this must be positive)
- Can have destructive and constructive interference





# Wave/Particle Duality

- This picture shows some of the possible electron probability densities for different quantum states of the H atom.
- Electron “clouds”



- Probability “clouds”
  - kind of the opposite of the “Plum Pudding” model

