

The Interaction of Light and Matter

- <http://webphysics.ph.msstate.edu/javamirror/ipmj/java/atomphoton/index.html>
- <http://webphysics.ph.msstate.edu/javamirror/ipmj/java/slitdiffr/index.html>
- <http://www.colorado.edu/physics/2000/quantumzone/index.html>
- <http://micro.magnet.fsu.edu/primer/java/doubleslit/index.html>
- <http://members.tripod.com/~vsg/interfer.htm>
- <http://micro.magnet.fsu.edu/primer/java/exciteemit/index.html>
- <http://www.colorado.edu/physics/2000/applets/a2.html>
- <http://home.a-city.de/walter.fendt/physengl/photoeffect.htm>
- <http://lectureonline.cl.msu.edu/~mmp/kap28/PhotoEffect/photo.htm>

Outline

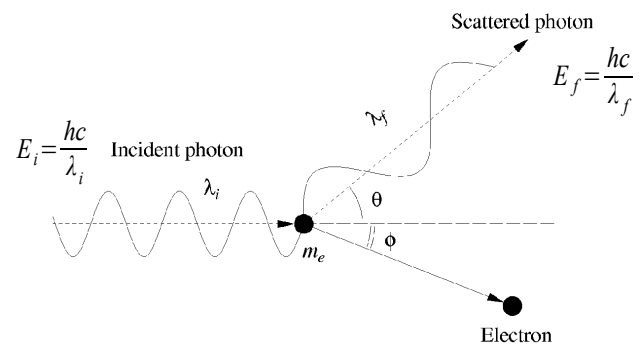
- (0) Compton scattering – light acts like a particle
- (1) Motivation: **Why spectral lines?**
 - the Birth of Spectroscopy
 - Kirchoff's Laws
- (2) **The Bohr Model of the Atom**
 - Need for a theory to describe spectral lines
- (3) **Photons – the particle nature of light**
- (4) **Quantum Mechanics and the Wave-Particle Duality**

Compton Scattering: “Particle-like” Behavior of Photon

Concept: Photon scatters off electron and loses energy, where resulting λ of scattered photon **depends on θ** .

Conservation of relativistic momentum and Energy!

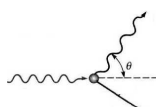
No mass for the photon but it has momentum!!!



Compton Scattering: Equation

$$\Delta\lambda = \lambda_f - \lambda_i = \frac{h}{m_e c} (1 - \cos \theta)$$

Photon OUT λ_f Scattering Angle θ
 Photon IN λ_i
 Critical $\lambda_c = 0.0024 \text{ nm}$ for e^-

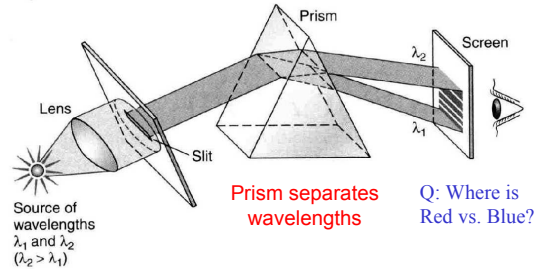


Limiting Values

- No scattering: $\theta = 0^\circ \rightarrow \cos 0^\circ = 1 \rightarrow \Delta\lambda = 0$
- "Bounce Back": $\theta = 180^\circ \rightarrow \cos 180^\circ = -1 \rightarrow \Delta\lambda = 2\lambda_c$
- Difficult to observe unless λ is small (i.e. $\Delta\lambda/\lambda > 0.01$)

Spectroscopy - history

- Trogg (50 million BC) – rainbow
- Newton (1642-1727) – decomposes light into spectrum and back again
- W. Herschel (1800) – discovers infrared
- J. W. Ritter (1801) – discovers ultraviolet
- W. Wollaston (1802) – discovers absorption lines in solar spectrum



Spectroscopy - history

- J. Herschel, Wheatstone, Alter, Talbot and Angstrom studied spectra of terrestrial things (flames, arcs and sparks) ~1810
- Joseph Fraunhofer
 - Cataloged ~475 dark lines of the solar spectrum by 1814
 - Identifies sodium in the Sun from flame spectra in the lab!
 - Looks at other stars (connects telescope to spectroscope)
- Foucault (1848) – sees absorption lines in sodium flame with bright arc behind it.



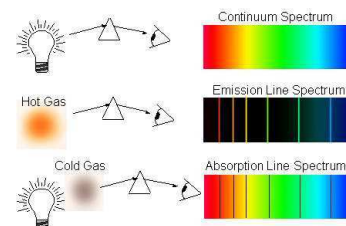
There is the need for a *new physics!*

Kirchhoff's laws (1859):

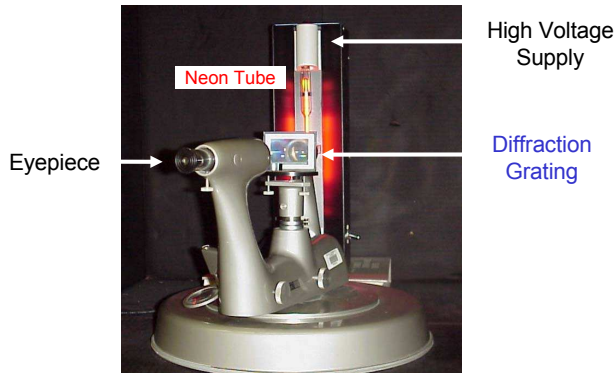
Worked with Bunsen on flame spectra

Developed a prism spectroscope

- Hot solid or dense gas, \rightarrow **blackbody radiation**
- Hot diffuse gas \rightarrow **emission lines**
- Cool diffuse gas in front of a blackbody \rightarrow **absorption lines**



Atomic Spectra Lab



$$d \sin \theta = n\lambda$$

Doppler shift

- Spectral lines allow for the measurement of radial velocities
- At low velocities, $v_r \ll c$
 - Classical Doppler effect
 - Radial velocity, v_r
 - Heliocentric correction for Earth's motion, up to 29.8 km/s, depending on direction.

$$\frac{\lambda_{obs} - \lambda_{rest}}{\lambda_{rest}} = \frac{\Delta\lambda}{\lambda_{rest}} = \frac{v_r}{c}$$

$$\Delta\lambda = \frac{v_r}{c} \lambda_{rest}$$

$$v_r = c \frac{\Delta\lambda}{\lambda_{rest}} = -14 \frac{km}{sec}$$

- Example: H_{α} is 6562.80 Å
 - Vega is measured to be 6562.50 Å
 - Coupled with the *proper motion*
 - Can determine total velocity

$$v_{\theta} = r\mu = 13 \frac{km}{s}$$

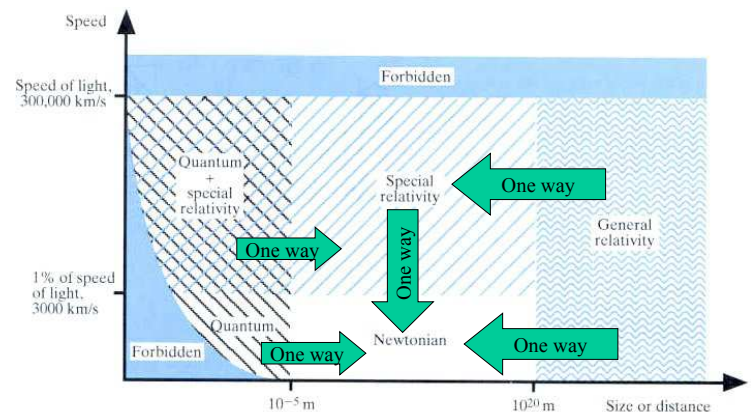
$$v = \sqrt{v_r^2 + v_{\theta}^2} = 19 \frac{km}{s}$$

What is Modern Physics?

- Modern physics only came of age in the **1900's**.
 - Physicists discovered that Newtonian mechanics does not apply when objects move **very fast** or are **very small**!
- If things move **very fast** (close to the speed of light), then **RELATIVISTIC** mechanics is necessary.
 - Cosmic particles, atomic clocks (GPS), synchrotrons.
- If things are confined to **very small** dimensions (nanometer-scale), then **QUANTUM** mechanics is necessary.
 - Atomic orbitals, quantum heterostructures.
- Need** a lot of these ideas to describe the universe

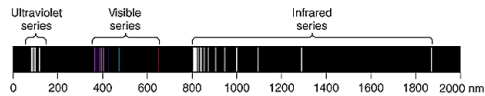
Beyond Newton

- Gravity passed every test until ~1890s
- Newton's gravity and motion is incorrect when ...



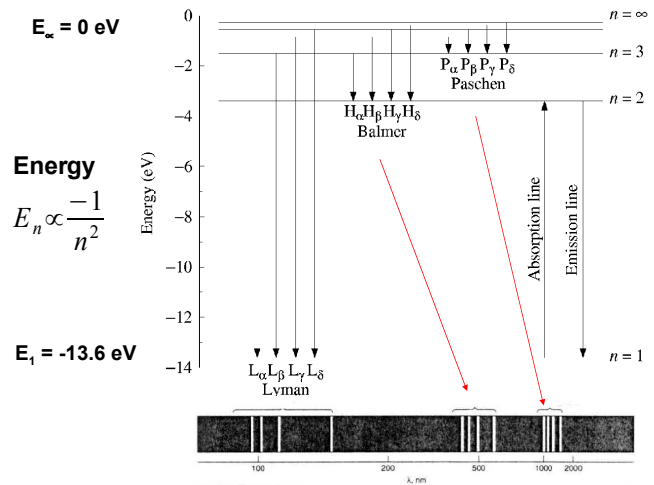
Atomic Spectra

- 1885 - **Balmer** observed Hydrogen Spectrum
 - Found empirical formula for discrete wavelengths
 - Later generalized by Rydberg for simple ionized atoms



$$\frac{1}{\lambda} = R_H \left(\frac{1}{2^2} - \frac{1}{n^2} \right) \text{ with } 2 < n$$

Atomic Spectra: Hydrogen Energy Levels



Atomic Spectra: Rydberg Formula

$$\frac{1}{\lambda} = R_H \left(\frac{1}{m^2} - \frac{1}{n^2} \right) \text{ with } m < n$$

- Rydberg constant $R_H \sim 1.097 \times 10^5 \text{ cm}^{-1}$
- $m = 1$ (Lyman), 2 (Balmer), 3 (Paschen)

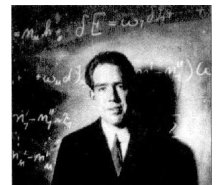
- Example for $n = 2$ to $m = 1$ transition:

$$\frac{1}{\lambda} = R_H \left(\frac{1}{1^2} - \frac{1}{2^2} \right) = \frac{3}{4} (1.097 \times 10^5 \text{ cm}^{-1})$$

$$\Rightarrow \lambda = 121.6 \text{ nm Ultraviolet}$$

Bohr Model

- 1913 – **Bohr** proposed **quantized model** of the H atom to predict the observed spectrum.
- Problem:** Classical model of the electron “orbiting” nucleus is unstable. Why unstable?
 - Electron experiences centripetal acceleration.
 - Accelerated electron emits radiation.
 - Radiation leads to energy loss.
 - Electron eventually “crashes” into nucleus.



Bohr Model: Quantization

- Solution:** Bohr proposed two "quantum" postulates
 - Electrons exist in stationary orbits (no radiation) with quantized angular momentum

$$L_n = mvr = n\hbar \quad \text{where } \hbar = \frac{h}{2\pi} = 6.58 \times 10^{-16} \text{ eV} \cdot \text{s}$$

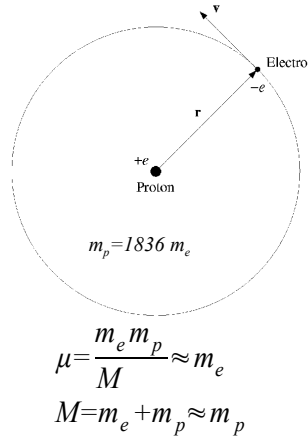
\hbar = Planck's Constant

- Atom radiates with quantized frequency ν (or energy E) only when the electron makes a transition between two stationary states.

$$h\nu = \frac{hc}{\lambda} = E_i - E_f$$

Planetary Mechanics Applied to the Atom

- Consider the attractive electrostatic force and circular motion



$$\vec{F} = \frac{q_1 q_2}{r^2} \hat{r} = \mu \frac{v^2}{r} \hat{r}$$

Note: in cgs, $e = 4.803 \times 10^{-10}$ esu

$$\frac{q_1 q_2}{r^2} = -\mu \frac{v^2}{r}$$

$$\frac{-e^2}{r^2} = -\mu \frac{v^2}{r}$$

$$\frac{1}{2} \mu v^2 = \frac{1}{2} \frac{e^2}{r} = K$$

Kinetic energy

$$U = -2K = -\frac{e^2}{r} \quad \text{Potential energy}$$

$$\mu = \frac{m_e m_p}{M} \approx m_e$$

$$M = m_e + m_p \approx m_p$$

Planetary Mechanics Applied to the Atom

- Introduce Bohr's quantized angular momentum $L = \mu vr = n\hbar$ (wrong)

$$K = \frac{1}{2} \frac{e^2}{r} = \frac{1}{2} \mu v^2 = \frac{1}{2} \frac{(\mu vr)^2}{\mu vr^2} = \frac{1}{2} \frac{(n\hbar)^2}{\mu vr^2}$$

- Solving for r $r_n = \frac{\hbar^2}{\mu e^2} n^2 = a_0 n^2$ a_0 is the Bohr radius

- Get the Total Energy in terms of n

$$E_n = -\frac{1}{2} \frac{e^2}{r} = -\frac{\mu e^4}{2\hbar^2} \frac{1}{n^2} = \frac{-13.6 \text{ eV}}{n^2} = \frac{-E_0}{n^2}$$

- Principle quantum number, $n = 1, 2, 3, \dots$

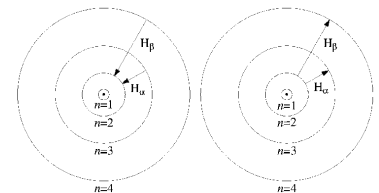
Bohr Model: Transitions

- Transitions predicted by Bohr yield general Rydberg formula

$$\delta E = -E_0 \left(\frac{1}{m^2} - \frac{1}{n^2} \right), \quad \delta E = \frac{hc}{\lambda}$$

$$\frac{1}{\lambda} = R_H \left(\frac{1}{m^2} - \frac{1}{n^2} \right) \quad \text{with } m < n$$

$$R_H = \frac{E_0}{hc} = \frac{\mu e^4}{4\pi \hbar^3 c}$$



- Applies to ionized atoms with only one electron and of nuclear charge Z .

$$\frac{1}{\lambda} = Z^2 \left(\frac{E_0}{hc} \right) \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

$$\frac{hc}{\lambda} = E_f - E_i = -Z^2 E_0 \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

Bohr Model Problem: Unknown Transition

If the wavelength of a transition in the **Balmer series** for a He^+ atom is **121 nm**, then find the corresponding transition, i.e. initial and final n values.

$$\frac{1}{\lambda} = RZ^2 \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right) = R(2)^2 \left(\frac{1}{(2)^2} - \frac{1}{n_i^2} \right)$$

where $Z = 2$ for He and $n_f = 2$ for Balmer

$$\frac{1}{4R\lambda} = \left(\frac{1}{4} - \frac{1}{n_i^2} \right)$$

$$n_i = \left(\frac{1}{4} - \frac{1}{4R\lambda} \right)^{-1/2} = \left(\frac{1}{4} - \frac{1}{4(1.1 \times 10^7 \text{ m}^{-1})(121 \times 10^{-9} \text{ m})} \right)^{-1/2} = \underline{4}$$

Bohr Model Problem: Ionization Energy

Suppose that a He atom ($Z=2$) in its ground state ($n = 1$) absorbs a photon whose wavelength is **$\lambda = 41.3 \text{ nm}$** . Will the atom be **ionized**?

➤ Find the energy of the incoming photon and compare it to the ground state ionization energy of helium, or E_0 from $n=1$ to ∞ .

$$E = \frac{hc}{\lambda} = \frac{1240 \text{ eV} \cdot \text{nm}}{41.3 \text{ nm}} = \underline{30 \text{ eV}}$$

$$E_0(\text{He}) = Z^2 \times E_0(\text{H}) = (2^2)(13.6 \text{ eV}) = 54.4 \text{ eV}$$

➤ The photon energy (30 eV) is less than the ionization energy (54 eV), so the electron will **NOT** be ionized.

Compton Scattering Problem

If a 0.511-MeV photon from a positron-electron annihilation scatters at **180°** from a free electron, then find the wavelength and energy of the Compton scattered photon.

notaiton $\lambda_1 = \lambda_i$ and $\lambda_2 = \lambda_f$

$$\lambda_2 - \lambda_1 = \lambda_C (1 - \cos \theta) = (0.00243 \text{ nm}) (1 - \cos 180^\circ) = 4.86 \times 10^{-3} \text{ nm}$$

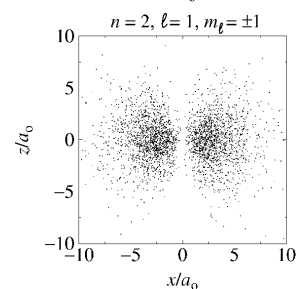
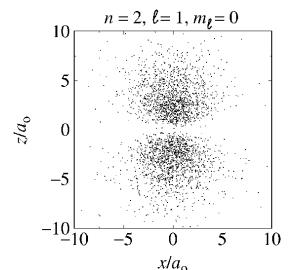
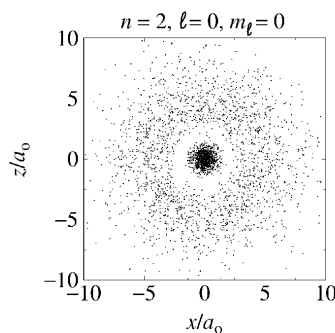
$$\lambda_1 = \frac{hc}{E_1} = \frac{1240 \text{ eV} \cdot \text{nm}}{0.511 \times 10^6 \text{ eV}} = \underline{2.43 \times 10^{-3} \text{ nm}}$$

$$\lambda_2 = \lambda_1 + \Delta\lambda = 2.43 \times 10^{-3} \text{ nm} + 4.86 \times 10^{-3} \text{ nm} = \underline{7.29 \times 10^{-3} \text{ nm}}$$

$$E_2 = \frac{hc}{\lambda_2} = \frac{1240 \text{ eV} \cdot \text{nm}}{7.29 \times 10^{-3} \text{ nm}} = 1.70 \times 10^5 \text{ eV} \text{ (or } \underline{0.17 \text{ MeV}})$$

• Probability "clouds"

– kind of the opposite to the "Plum Pudding" model



Zeeman Effect

- Measure magnetic field strengths
- Solving the Schrodinger equation yields two more quantum numbers

