

MATH292 NOTES

Contents

1	Unit 1	2
1.1	Lecture 1	2
1.1.1	What is a Differential Equation?	2
1.1.2	First Order ODEs	2
1.1.3	Vector Fields and Integral Curves	3
1.1.4	First Order Linear Equations	4
1.1.5	Extras	5
1.2	Lecture 2	6
1.2.1	Solution Curves	6
1.2.2	Flow Transformations	6
1.2.3	Bernoulli Equations	7
1.2.4	Riccati Equations	8
1.2.5	Reduction of Order	10
2	Unit 2	12
2.1	Lecture 3	12
2.1.1	Monotonicity on Maximal Intervals	12
2.1.2	Lipschitz Continuity	15
2.2	Lecture 4	16
2.2.1	Barrows Formula	16
2.2.2	Existance and Uniqueness for First Order Autonomous Equations	17
2.2.3	Existance and Uniqueness within Maximal Intervals	19

1 Unit 1

1.1 Lecture 1

1.1.1 What is a Differential Equation?

A **differential equation** relates a function to its derivatives. An **ordinary differential equation** (ODE) is a differential equation in which the function only has one independent variable. Differential equations often arise from physical laws such as Newton's second law:

$$x''(t) = \frac{1}{m}F(x(t))$$

1.1.2 First Order ODEs

Order

The order of a differential equation refers to highest derivative found in the differential equation.

It follows that a first order ODE refers to an ODE that only contains first derivatives. The general form for a first order ODE is

$$G(x, y, y') = 0 \tag{1}$$

for some real-valued function G . We can apply the *Implicit Function Theorem* on G if G is continuously differentiable and

$$G(x_0, y_0, z_0) = 0 \quad \text{and} \quad \frac{\partial}{\partial z}G(x_0, y_0, z_0) \neq 0$$

to give a much 'nicer' form of

$$y' = g(x, y)$$

The solution to a first order ODE is the function $y(x)$ that satisfies (1). The **standard form** of a first order ODE is an equation that gives the derivative y' as a function of y . There are two main types of standard forms.

Standard Forms

An **autonomous** first order ODE is one in which the derivative only depends on y . The standard form of an autonomous first order ODE is

$$y' = f(y).$$

A **non-autonomous** first order ODE is one in which the derivative depends on x as well as the value of the function y itself. The standard form of a non-autonomous first order ODE is

$$y' = g(x, y).$$

Seperable ODEs

Seperable first order ODEs are ODEs in the form

$$f(y)y' = g(x).$$

These can be solved with (relative) ease by integrating both sides and then solving algebraically.

Example 1.1.1: A Seperable First Order ODE

Solve the differential equation

$$y' = e^{x+y}$$

Solution:

$$\begin{aligned} y' = e^{x+y} &\Rightarrow e^{-y} y' = e^x \\ &\Rightarrow \int e^{-y} dy = \int e^x dx \\ &\Rightarrow -e^{-y} = e^x + C \\ &\Rightarrow e^{-y} = -e^x + C \\ &\Rightarrow y = -\ln(e^x + C) \end{aligned}$$

1.1.3 Vector Fields and Integral Curves

Consider a function $\vec{x}'(t) = v(\vec{x}(t))$ that is continuous on some open subset $U \subseteq \mathbb{R}^n$ and has values in \mathbb{R}^n . Then, $v(x)$ is a **vector field**. To solve this equation means to find all continuously differentiable functions $\vec{x}(t)$ defined on some time interval (a, b) with values in U for which the equation is satisfied $\forall t \in (a, b)$.

Vector Function and Integral Curves

An **autonomous first order system of ODEs** in standard form is an equation in the form

$$\vec{x}'(t) = v(\vec{x}(t)).$$

A **non-autonomous first order system of ODEs** in standard form is an equation in the form

$$\vec{x}'(t) = v(t, \vec{x}(t)).$$

The solutions to either an autonomous or non-autonomous system, $\vec{x}(t)$ are called **integral curves** of the vector field.

As a basic example consider the vector field on \mathbb{R}^2 given by $v(x, y) = (-y, x)$. Then, the differential equation $\vec{x}'(t) = v(\vec{x}(t))$ is simply a more efficient way to write

$$\begin{aligned} x'(t) &= -y(t) \\ y'(t) &= x(t) \end{aligned}$$

We might solve this by considering the function $f(\vec{x}) = \|\vec{x}\|^2 = x^2 + y^2$ (don't worry about where this function came from for now!) and apply the multivariable chain rule,

$$\frac{d}{dt} [f(\vec{x}(t))] = \nabla f(\vec{x}(t)) \cdot \vec{x}'(t) = (2x, 2y) \cdot (-y, x) = 0.$$

Recall that $\vec{x}'(t) = v(\vec{x}(t)) = (-y, x)$. Since

$$\frac{df(\vec{x}(t))}{dt} = 0$$

we know that $f(\vec{x}(t))$ must be a constant function. Thus, $f(\vec{x}(t)) = f(x, y) = C = x^2 + y^2$ for some constant $C \in \mathbb{R}$. It should be clear to see that $f(\vec{x}(t))$ simply represents a circle centered at the origin with radius \sqrt{C} .

Slope Function

let $v(x, y) = (f(x, y), g(x, y))$ be any continuous vector field on an open set U of \mathbb{R}^2 . On the subset open V on which $f(x, y) \neq 0$, define the slope function for the vector field to be

$$s(x, y) = \frac{g(x, y)}{f(x, y)}.$$

We define the slope function because an integral curve of V is the function $y(x)$ that satisfies the differential equation $y' = s(x, y)$. This is important because there can be occasions in which you may be able to solve the slope equation $y' = s(x, y)$ despite not being able to explicitly solve for the integral curves of a vector field.

1.1.4 First Order Linear Equations

Differential equations in the form

$$x'(t) = p(t)x(t) + q(t) \quad (2)$$

are called **first order linear equations**. By letting p and q be continuous on some interval (a, b) (with $a = -\infty$ and $b = \infty$ allowed) we consider the primitive (or anti-derivative) of p , $P(t) = \int p(t) dt$. Now, notice that

$$\frac{d}{dt} [x(t)e^{-P(t)}] = [x'(t) - p(t)x(t)] e^{-P(t)}.$$

But,

$$x'(t) - p(t)x(t) = q(t).$$

So,

$$\begin{aligned} \frac{d}{dt} [x(t)e^{-P(t)}] &= [x'(t) - p(t)x(t)] e^{-P(t)} \\ \Rightarrow \frac{d}{dt} [x(t)e^{-P(t)}] &= q(t)e^{-P(t)} \\ \Rightarrow \int \frac{d}{dt} [x(t)e^{-P(t)}] dt &= \int q(t)e^{-P(t)} dt \\ \Rightarrow x(t)e^{-P(t)} + C &= \int q(t)e^{-P(t)} dt \\ \Rightarrow x(t) &= e^{P(t)} \left(\int q(t)e^{-P(t)} dt + C \right). \end{aligned} \quad (3)$$

And thus (3) is the general solution to a first order linear differential equation.

You may be asked to solve an **initial value problem** (IVP) which is simply when an initial condition is provided (i.e. $x(t_0) = x_0$). In order to solve an IVP for a first order linear ODE all you must do is solve for the general solution and then plug in the initial condition and solve for the constant of integration. Usually when solving an IVP we have two main questions:

1. Is there a solution? (existence)
2. If there is a solution, is it unique? (uniqueness)

Example 1.1.2: First Order Linear Differential Equation

Solve the differential equation

$$x'(t) = tx(t) + 1$$

Solution: We'll start by noticing that $p(t) = t$ so $P(t) = \frac{1}{2}t^2$. Thus,

$$\begin{aligned} x'(t) = tx(t) + 1 &\Rightarrow \frac{d}{dt} \left[x(t)e^{-\frac{1}{2}t^2} \right] = 1e^{-\frac{1}{2}t^2} \\ &\Rightarrow \int \frac{d}{dt} \left[x(t)e^{-\frac{1}{2}t^2} \right] dt = \int e^{-\frac{1}{2}t^2} dt \\ &\Rightarrow x(t)e^{-\frac{1}{2}t^2} + C = \int e^{-\frac{1}{2}t^2} dt \\ &\Rightarrow x(t) = \boxed{e^{\frac{1}{2}t^2} \left(\int e^{-\frac{1}{2}t^2} dt + C \right)} \end{aligned}$$

For those curious, the integral on the right handside is not an easy integral to evaluate (it requires the power series for e^x) so for simplicity sake we leave it as is and thus the boxed answer is our general solution in simplest form.

1.1.5 Extras**Flow Transformation**

For a first order linear equation in the form $x'(t) = p(t)x(t) + q(t)$ with p and q continuous on (a, b) , define

$$\Phi_{t_1, t_0}(x) = x(t_1)$$

for all $t_0 \neq t_1 \in (a, b)$ and where $x(t)$ is the unique solution to the IVP $x(t_0) = x_0$.

Matrix Exponentials

Consider e^A where A is an $n \times n$ matrix over \mathbb{C} defined as

$$e^A = \sum_{n=0}^{\infty} \frac{A^n}{n!}.$$

If $A = PDP^{-1}$ (i.e. if A is diagonalized) then,

$$e^A = Pe^D P^{-1}$$

where

$$e^D = \begin{pmatrix} e^{\lambda_1} & 0 \\ 0 & e^{\lambda_2} \end{pmatrix} \text{ if } D = \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix}.$$

1.2 Lecture 2

1.2.1 Solution Curves

Solution Curves

Consider the first order ODE given by

$$x'(t) = p(t)x(t) + q(t)$$

defined on an open interval (a, b) . A **solution curve** is the graph of any specific solution.

Note that the general solution of any first order ODE is of the form

$$x(t) = cf(t) + g(t)$$

with $c \in \mathbb{R}$ and $f(t) \neq 0$.

Any two different solutions curves of a given first order ODE will never intersect. This follows from the fact that any two different solutions of the ODE correspond to different values of the constant, c , and since $f(t) \neq 0$, the values of the two solutions at any t must be different. Hence, through each point in the (t, x) plane given by $(a, b) \times \mathbb{R}$, there is exactly one solution curve. Conversely, every family of functions of the form

$$x(t) = cf(t) + g(t)$$

with g and f continuously differentiable and such that $f(t) \neq 0, \forall t \in (a, b)$ is the solution set of a first order linear equation on (a, b) . To see this combine

$$x(t) = cf(t) + g(t)$$

and

$$c = \frac{1}{f(t)}(x(t) - g(t))$$

to deduce

$$\begin{aligned} x'(t) &= \frac{f'(t)}{f(t)}(x(t) - g(t)) + g'(t) \\ &= \frac{f'(t)}{f(t)}x(t) + \left(g'(t) - \frac{f'(t)}{f(t)}g(t)\right) \end{aligned}$$

1.2.2 Flow Transformations

Consider the ODE

$$x'(t) = \frac{-x(t)}{t} + t$$

on the interval $t > 0$. The field of solution curves extends over the entire right half of the (t, x) plane. If we think of the equation as describing the motion of a point on the line, such that the point is at x_0 at time t_0 , then by following the unique solution curves through (t_0, x_0) , we can see where the particle is at every other $t > 0$. Through each point in the right half of the plane there is exactly one solution curve. We define a function $\Phi_{t_1, t_0} : \mathbb{R} \rightarrow \mathbb{R}$ as follows: $\Phi_{t_1, t_0}(x_0)$ is the intersection of the solution curve passing through (t_0, x_0) with the line $t = t_1$. In our graph, locate the point at height x_0 in the vertical line $t = t_0$. Follow the solution curve $x(t)$ through this point until it intersects the line $t = t_1$. The height of the intersection, which is $x(t_1)$, is $\Phi_{t_1, t_0}(x_0)$.

Flow Transformations

Consider the IVP

$$x'(t) = p(t)x(t) + q(t), \quad x(t_0) = x_0$$

with p, q continuous on (a, b) . For every $t_0 \neq t_1 \in (a, b)$, define

$$\Phi_{t_1, t_0}(x_0) = x(t_1)$$

where x is the unique solution to the IVP.

One important fact about flow transformations is that for any distinct $t_0, t_1, t_2 \in (a, b)$ and $x_0 \in \mathbb{R}$, we have

$$\Phi_{t_2, t_1}(\Phi_{t_1, t_0}(x_0)) = \Phi_{t_2, t_0}(x_0).$$

Proof. The LHS is obtained by starting from the point (t_0, x_0) and moving along the solution curve of the ODE to meet the vertical line $t = t_1$, then moving along the solution curve passing $(t_1, x(t_1))$ until it intersects the vertical line $t = t_2$ where x is the unique solution to the IVP $x(t_0) = x_0$. The RHS is obtained by starting from the point (t_0, x_0) and moving along the solution curve of the ODE to meet the vertical line $t = t_2$. Both sides are now the point $(t_2, x(t_2))$. \square

Example 1.2.1

Compute the flow transformation for the IVP

$$x'(t) = \frac{-x(t)}{t} + t, \quad x(t_0) = x_0.$$

Solution: This is a basic linear differential equation and be solved using an integrating factor of $\mu(t) = e^{\int \frac{1}{t} dt} = t$. Upon solving you should arrive at a final solution of

$$x(t) = \frac{t^2}{3} + \frac{c}{t}$$

and so, solving for c ,

$$c = t_0 \left(x_0 - \frac{t_0^2}{3} \right) \Rightarrow x(t) = \frac{t^2}{3} + \frac{t_0 x_0 - \frac{t_0^3}{3}}{t}.$$

Thus,

$$\Phi_{t_1, t_0}(x_0) = x(t_1) = \frac{t_1^2}{3} + \frac{3t_0 x_0 - t_0^3}{3t_1}$$

1.2.3 Bernoulli Equations

Bernoulli Equations

Any ODE of the form

$$x'(t) = p(t)x(t) + q(t)x^n(t), \quad n \neq 0, 1 \in \mathbb{R}$$

is called a **Bernoulli Equation**.

To solve a Bernoulli equation we use the change of variables

$$v = x^{1-n} \Rightarrow v' = (1-n)x^{-n}x'$$

under the assumption that $x(t) \neq 0$ for all t . We can then manipulate our Bernoulli Equation in the following manner:

$$\begin{aligned} x'(t) = p(t)x(t) + q(t)x^n(t) &\Rightarrow x'(t) - p(t)x(t) = q(t)x^n(t) \\ &\Rightarrow x^{-n}(t)x'(t) - p(t)x^{1-n}(t) = q(t) \\ &\Rightarrow \frac{1}{1-n}v' - vp(t) = q(t). \end{aligned}$$

Thus resulting in a first order linear differential equation that we know how to solve. You will then solve for v , followed by using your solution for v to solve for $x(t)$.

Example 1.2.2

Solve the following ODE:

$$x'(t) = x^2(t)$$

Solution:

$$\begin{aligned} x'(t) = x^2(t) &\Rightarrow x^{-2}x' = 1 \\ &\Rightarrow v = x^{-1} \Rightarrow v' = -x^{-2}x' \\ &\Rightarrow -v' = 1 \\ &\Rightarrow v = -t + c \\ &\Rightarrow x^{-1} = -t + c \\ &\Rightarrow x(t) = \frac{1}{-t + c} \end{aligned}$$

for $c \in \mathbb{R}$.

1.2.4 Riccati Equations

Riccati Equations

Any ODE of the form

$$x'(t) = p(t) + q(t)x(t) + r(t)x^2(t)$$

is called a **Riccati Equation**.

There is no general method to solve a Riccati equation unless you're given a particular solution, $x_1(t)$, in which case the general solution will be of the form

$$x(t) = x_1(t) + u(t)$$

where $u(t)$ is the general solution of the Bernoulli equation

$$u'(t) = [q(t) + 2r(t)x_1(t)]u(t) + r(t)u^2(t). \quad (4)$$

Now, substituting in for x into the Riccati equation we see that

$$\begin{aligned} (x_1 + u)' - p - q(x_1 + u) - r(x_1 + u)^2 &= 0 \\ \Rightarrow (x_1' - p - qx_1 - rx_1^2) + (u' - (q + 2rx_1)u - ru^2) &= 0 \\ \Rightarrow u' - (q + 2rx_1)u - ru^2 &= 0. \end{aligned}$$

With the last implication following from the fact that the first term is simply just the Riccati equation rewritten.

Example 1.2.3

Solve the following ODE:

$$x'(t) = -\frac{1}{t}x(t) + \frac{1}{t^3}x^2(t) + 2t$$

Solution: Since all the coefficients are powers of gt , it is natural to see if there is a solution of the form

$$x_1(t) = ct^\alpha$$

for $c, \alpha \in \mathbb{R}$. Substituting this in the ODE, we get

$$c\alpha t^{\alpha-1} = -ct^{\alpha-1} + c^2t^{2\alpha-3} + 2t.$$

Since this is a functional equivalence, all the powers of t on the LHS must be equal to all the powers of t on the RHS. This leads us to the fact that

$$\alpha - 1 = \alpha - 1 = 2\alpha - 3 = 1 \Rightarrow \alpha = 2.$$

Now, since all the powers of t are equal, we can cancel them which leaves us with

$$2c = -c + c^2 + 2 \Rightarrow c^2 + c + 2 = 0 \Rightarrow (c - 1)(c - 2) = 0.$$

Thus, either $x_1 = t^2$ or $x_1 = 2t^2$. With this in mind, we can now solve for u using (4) and choosing $x_1 = t^2$.

$$\begin{aligned} u'(t) &= \frac{1}{t}u + \frac{1}{t^3}u^2 \Rightarrow u^{-2}u' - \frac{1}{t}u^{-1} = \frac{1}{t^3} \\ \Rightarrow v &= u^{-1} \Rightarrow v' = -u^{-2}u' \\ \Rightarrow -v' - \frac{1}{t}v &= \frac{1}{t^3} \\ \Rightarrow v't + v &= -\frac{1}{t^2} & (\mu(t) = t) \\ \Rightarrow vt &= \int -\frac{1}{t^2}dt = \frac{1}{t} + c \\ \Rightarrow v &= \frac{1}{t^2} + \frac{c}{t} = \frac{1+ct}{t^2} \\ \Rightarrow u &= \frac{1}{v} = \frac{t^2}{1+ct} \\ \Rightarrow x(t) &= x_1(t) + u(t) = t^2 + \frac{t^2}{1+ct} \end{aligned}$$

1.2.5 Reduction of Order

Some second order ODEs can be reduced to first order equations that we know how to solve. General second order ODEs have the form

$$F(t, x, x', x'') = 0$$

for some real-valued function F defined on a subset of \mathbb{R}^4 . There are two cases when dealing with reduction of order problems.

Case 1: F does not depend on x . In this case $F(t, x', x'') = 0$ so we let $y = x'$ which turns F into $F(t, y, y')$ which is first order.

Example 1.2.4

Solve the following ODE:

$$tx'' + x' = t^3$$

Solution: Let $y = x'$ and substitute,

$$\begin{aligned} ty' + y &= t^3 \Rightarrow y' + t^{-1}y = t^2 \\ \Rightarrow \mu(t) &= e^{\int t^{-1} dt} = t \\ \Rightarrow ty' + y &= t^3 \\ \Rightarrow ty &= \int t^3 dt = \frac{t^4}{4} + c \\ \Rightarrow y &= \frac{t^3}{4} + \frac{c}{t} \\ \Rightarrow x &= \int y(t) dt = \int \frac{t^3}{4} + \frac{c}{t} dt = \frac{t^4}{16} + c \ln(t) + d \end{aligned}$$

for $c, d \in \mathbb{R}$.

Case 2: F does not depend on t . Once again we let $y = x'$, but also define $X(t) = (x(t), y(t))$ and the vector field $v(x, y) = (y, f(x, y))$. Then $x(t)$ solves $x'' = f(x, x')$ if, and only if, $X(t)$ solves $X'(t) = v(X(t))$. Our strategy is to first find a first order ODE for y . Solving this will tell us the curve traced out by $X(t)$ for t near t_0 . Once we know the explicit function, $y(x)$, recall that $y = x'$, and so $x'(t) = y(x(t))$ is a first order equation for $x(t)$. To actually find an equation for $y(x)$ note that the slope of the graph of this function, namely $y'(x)$, is given by the slope of the vector $v(x, y) = (y, f(x, y))$ which is

$$y' = \frac{f(x, y)}{y}. \quad (5)$$

And so, (5) is a separable first order equation we can solve.

Example 1.2.5: Simple Harmonic Motion

Solve the following ODE:

$$x'' + \omega^2 x = 0, \quad \omega \in \mathbb{R}^+$$

Solution: Notice that

$$\begin{aligned}\frac{dy}{dx} &= \frac{f(x, y)}{y} = \frac{-\omega^2 x}{y} \\ \Rightarrow \int y \, dy &= \int -\omega^2 x \, dx \\ \Rightarrow \frac{y^2}{2} &= \frac{-\omega^2 x^2}{2} + c \\ \Rightarrow \omega^2 x^2 + y^2 &= 2c \\ \Rightarrow \omega^2 x^2 + (x')^2 &= 2c\end{aligned}$$

for $c \in \mathbb{R}$. The case $c = 0$ is trivial so let $2c = r^2 \geq 0$. Then $\omega^2 x^2 + (x')^2 = r^2$ so the point $(\omega x(t), x'(t))$ is on the circle centered at the origin with radius r . Using this to convert to polar we can see that

$$\begin{cases} \omega x(t) &= r \cos(\theta(t)) \\ x'(t) &= r \sin(\theta(t)) \end{cases}$$

But,

$$\begin{aligned}(\omega x(t))' &= \omega x'(t) \\ \Rightarrow \frac{d}{dt}(r \cos(\theta(t))) &= \omega r \sin(\theta(t)) \\ \Rightarrow -r \sin(\theta(t)) \theta'(t) &= \omega r \sin(\theta(t)) \\ \Rightarrow -x'(t) \theta'(t) &= \omega x'(t) \\ \Rightarrow -\theta'(t) &= \omega \\ \Rightarrow \theta(t) &= -\omega t + \theta_0\end{aligned}$$

for $\theta_0 \in \mathbb{R}$. Hence,

$$x(t) = \frac{r \cos(-\omega t + \theta_0)}{\omega}$$

where $r \geq 0$ and $\theta_0 \in [0, 2\pi)$

2 Unit 2

2.1 Lecture 3

2.1.1 Monotonicity on Maximal Intervals

Equilibrium Points and Steady State Solutions

If v is a vector field on \mathbb{R} and $v(x_0) = 0$, then x_0 is called an **equilibrium point** of v . For any $t_0 \in \mathbb{R}$, the function $x(t) = x_0$ for all t is a solution of the IVP

$$x'(t) = v(x(t)), \quad x(t_0) = x_0.$$

Such a constant solution is called a **steady state solution** for $x'(t) = v(x(t))$.

Maximal Intervals

An interval (a, b) is a maximal interval for v if $v(x) \neq 0$ for all $x \in (a, b)$ and if either $a = -\infty$ or $v(a) = 0$ and either $b = \infty$ or $v(b) = 0$.

For example, consider $v : \mathbb{R} \rightarrow \mathbb{R}$ given by

$$v(x) = x(1 - x).$$

Then, v will have three maximal intervals. Namely $(-\infty, 0)$, $(0, 1)$ and $(1, \infty)$. Note that these maximal intervals were found by setting $v(x) = 0$ and solving for x as the defining characteristic of a maximal interval is $v(x) \neq 0$.

Monotone functions

Let $f : A \rightarrow B$ be a function. Then, f is said to be **monotone increasing** if, and only if,

$$x_1 < x_2 \Rightarrow f(x_1) \leq f(x_2).$$

Similarly, f is **monotone decreasing** if, and only if,

$$x_1 < x_2 \Rightarrow f(x_1) \geq f(x_2).$$

As well, f is **strictly** monotone increasing if, and only if,

$$x_1 < x_2 \Rightarrow f(x_1) < f(x_2),$$

and **strictly** monotone decreasing if, and only if,

$$x_1 < x_2 \Rightarrow f(x_1) > f(x_2).$$

Fact

Let $f : A \rightarrow B$ be a function. If f is strictly monotone increasing or decreasing on an interval $(a, b) \subseteq A$, then the inverse function $f^{-1} : B \rightarrow A$ exists.

Suppose (a, b) is a maximal interval for v and $x_0 \in (a, b)$. Then $v(x_0) \neq 0$ by definition of a maximal interval. Now, assume that $v(x_0) > 0$. Since v is continuous and $v(x) \neq 0$ through (a, b) , we can apply the intermediate value theorem to see that $v(x) > 0$ on (a, b) . It follows that $x(t)$ must be strictly increasing on (a, b) since

$$x'(t) = v(x(t)).$$

A similar statement can be made for $v(x_0) < 0$ and $x(t)$ being strictly decreasing. Thus, $x(t)$ is strictly monotone (either increasing or decreasing) over (a, b) and hence is an invertible function onto its range. We denote the inverse function $t(x)$.

We define

$$T_a = \lim_{x \rightarrow a} t(x) \quad \text{and} \quad T_b = \lim_{x \rightarrow b} t(x)$$

under the assumption that v is positive on (a, b) and that $x(t)$ and $t(x)$ both exist and are strictly increasing functions. Then, $x(t)$ is invertible from (T_a, T_b) onto (a, b) with $t(x)$ going the opposite way.

If v is negative, and thus $x(t)$ and $t(x)$ are decreasing functions, then simply swap T_a and T_b such that $x(t)$ is invertible from (T_b, T_a) onto (a, b) and $t(x)$ invertible from (a, b) onto (T_a, T_b) .

Example 2.1.1: Solution on a Maximal Interval

Consider the differential equation

$$v(x) = x'(t) = x(1 - x), \quad x(t_0) = x_0 \in (0, 1).$$

As previously mentioned, we know that $(0, 1)$ is a maximal interval for v . Observe that

$$\frac{1}{x(1-x)} = \frac{1}{x} + \frac{1}{1-x} = \frac{d}{dx} \left[\ln \left(\frac{x}{1-x} \right) \right].$$

If we define

$$y = F(x) = \ln \left(\frac{x}{1-x} \right)$$

then,

$$\frac{d}{dt} F(x(t)) = 1 \tag{6}$$

because

$$\begin{aligned} \frac{d}{dt} F(x(t)) &= F'(x(t)) x'(t) \\ &= \frac{d}{dx} \left[\ln \left(\frac{x}{1-x} \right) \right] x'(t) \\ &= \frac{1}{x(1-x)} \cdot x(1-x) = 1. \end{aligned}$$

Now, integrating both sides of (6),

$$\int_{t_0}^t \frac{d}{dt} F(x(t)) dt = \int_{t_0}^t 1 dt \Rightarrow F(x(t)) - F(x(t_0)) = t - t_0 \Rightarrow F(x(t)) = F(x_0) + t - t_0, \tag{7}$$

recalling that $x(t_0) = x_0$. In order to solve this for $x(t)$ start by noticing that

$$\begin{aligned} F(x) = y &\Rightarrow y = \ln \left(\frac{x}{1-x} \right) \\ &\Rightarrow e^y = \frac{x}{1-x} \\ &\Rightarrow e^y(1-x) = x \\ &\Rightarrow e^y = x + e^y x \\ &\Rightarrow \frac{e^y}{1+e^y} = x. \end{aligned}$$

Since $y = F(x) = F(x_0) + t - t_0$,

$$\begin{aligned} x &= \frac{e^y}{1 + e^y} \\ &= \frac{e^{F(x_0) + t - t_0}}{1 + e^{F(x_0) + t - t_0}} \\ &= \frac{e^{F(x_0)} e^{t - t_0}}{1 + e^{F(x_0)} e^{t - t_0}}. \end{aligned} \tag{8}$$

But

$$e^{F(x_0)} = e^{\ln\left(\frac{x_0}{1-x_0}\right)} = \frac{x_0}{1-x_0}.$$

Applying this to (8) and simplifying we can see that

$$\begin{aligned} \frac{e^{F(x_0)} e^{t - t_0}}{1 + e^{F(x_0)} e^{t - t_0}} &= \frac{\frac{x_0}{1-x_0} e^{t - t_0}}{1 + \frac{x_0}{1-x_0} e^{t - t_0}} \\ &= \frac{x_0 e^{t - t_0}}{(1-x_0) + x_0 e^{t - t_0}}. \end{aligned}$$

It should be clear that $x(t_0) = x_0$. Since the denominator will never be 0, it follows that $x(t)$ is defined for all t . Because

$$\lim_{t \rightarrow -\infty} x(t) = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} x(t) = 1,$$

we can see that $x(t)$ is a strictly monotone increasing function from $(-\infty, \infty)$ onto $(0, 1)$. Thus, the inverse function $t(x)$ must exist (from the earlier fact). Solving (7) for t we can see that

$$t(x) = t_0 + F(x) - F(x_0),$$

but,

$$F(x) = \frac{1}{v} \Rightarrow F(x) - F(x_0) = \int_{x_0}^x \frac{1}{v(z)} dz.$$

So,

$$t(x) = t_0 + \int_{x_0}^x \frac{1}{v(z)} dz$$

with

$$\int_{x_0}^1 \frac{1}{v(z)} dz = \infty \quad \text{and} \quad \int_{x_0}^0 \frac{1}{v(z)} dz = -\infty$$

for any $x_0 \in (0, 1)$. This implies that $T_0 = -\infty$ and $T_1 = \infty$.

2.1.2 Lipschitz Continuity

Metric Spaces

A metric space is a double $M = (M, d)$ that consists of a set M and a distance *metric*, d . In order for M to be a metric space, the distance metric must satisfy the following axioms:

1. Non-negativity: For any $x, y \in M$, $d(x, y) \geq 0$
2. Identity: For any $x, y \in M$, $d(x, y) = 0 \Leftrightarrow x = y$
3. Symmetry: For any $x, y \in M$ $d(x, y) = d(y, x)$
4. Triangle Inequality: For any $x, y, z \in M$

$$d(x, z) \leq d(x, y) + d(y, z)$$

The prototypical metric space is \mathbb{R}^n with the euclidean metric (the standard distance formula).

Lipschitz Continuity

Let $f : X \rightarrow Y$ be a function between metric spaces. It is L -Lipschitz if there is a constant $L > 0$ such that for any $x_1, x_2 \in X$

$$d_y(f(x_1), f(x_2)) \leq L \cdot d_x(x_1, x_2)$$

where d_x and d_y are the metrics for X and Y respectively.

Fact

If $f : (a, b) \rightarrow \mathbb{R}^n$ is differentiable and

$$|f'(x)| \leq L, \quad \forall x \in (a, b)$$

then, f is L -Lipschitz continuous. Formally, if f is L -Lipschitz continuous on (a, b) then

$$|f(x) - f(y)| \leq L|x - y|, \quad \forall x, y \in (a, b).$$

We can prove this fact as follows.

Proof. Let $x, y \in (a, b)$ such that $x < y$ and assume that $|f'(x)| \leq L$. By the Mean Value Theorem,

$$f(y) - f(x) = f'(c)(y - x)$$

for $c \in (x, y)$. So,

$$|f(y) - f(x)| \leq L|y - x|$$

□

2.2 Lecture 4

2.2.1 Barrows Formula

Theorem 2.1. Let v be continuous and let (a, b) be a maximal interval for v and $x_0 \in (a, b)$. Fix any $t_0 \in \mathbb{R}$ and define $t(x)$ on (a, b) by

$$t(x) = t_0 + \int_{x_0}^x \frac{1}{v(z)} dz. \quad (9)$$

Then $t(x)$ is a strictly monotone function on (a, b) so that

$$T_a = \lim_{x \downarrow a} t(x) \quad \text{and} \quad T_b = \lim_{x \uparrow b} t(x) \quad (10)$$

both exist.

If v is positive on (a, b) then $T_a < t_0 < T_b$, and if $x(t)$ is the inverse function to $t(x)$, then $x(t)$ is a solution of

$$x'(t) = v(x(t)), \quad x(t_0) = x_0. \quad (11)$$

Moreover, every solution to (11) that is defined on any subinterval of (T_a, T_b) containing t_0 equals the restriction of $x(t)$ to this subinterval. In particular, there is a unique solution of (11) defined on (T_a, T_b) . If v is negative on (a, b) the same conclusion is valid provided we interchange T_a and T_b .

Proof. We first suppose that v is positive on (a, b) and define the function $t(x)$ on (a, b) by Barrows formula (9). For any $x \in (a, b)$ that isn't x_0 , let J be the closed interval with endpoint x and x_0 , which is either $[x_0, x]$ in the case $x_0 < x < b$, or else $[x, x_0]$ in the case $a < x < x_0$.

Since v is continuous on J , it has a minimum value attained somewhere in J , and since v is positive everywhere on (a, b) , and hence everywhere on J , there is a $c > 0$ so that $v(x) \geq c$ for all $x \in J$. Therefore $1/v(x)$ is continuous and bounded on J . It follows that

$$\int_{x_0}^x \frac{1}{v(z)} dz$$

is a proper integral for each $x \in (a, b)$. Thus,

$$t(x) = t_0 + \int_{x_0}^x \frac{1}{v(z)} dz$$

does define a function on (a, b) . By the Fundamental Theorem of Calculus, this function is differentiable, and

$$\frac{d}{dx} t(x) = \frac{1}{v(x)}. \quad (12)$$

Since the right hand side is continuous, $t(x)$ is continuously differentiable on (a, b) . Since $v(x) > 0$ on (a, b) , $t(x)$ is strictly monotone increasing on (a, b) . Therefore, the limit defining T_a and T_b in (10) both exist with $T_a = -\infty$ and $T_b = \infty$ allowed, and $t(x)$ is an invertible function from (a, b) onto (T_a, T_b) . Since $t(x)$ is differentiable, by the Inverse Function Theorem, $x(t)$ is differentiable and

$$\frac{d}{dt} x(t) = \left(\frac{d}{dx} t(x) \right)^{-1} \Big|_{x=x(t)}.$$

By (12)

$$\left(\frac{d}{dx} t(x) \right)^{-1} \Big|_{x=x(t)} = v(x(t))$$

which shows that $x'(t) = v(x(t))$, and clearly $t(x_0) = t_0$, so $x(t_0) = x_0$.

Now suppose that $y(t)$ is a continuously differentiable function defined on some interval (S, T) with $t_0 \in (S, T)$. Suppose also that

$$y'(t) = v(y(t)) \quad \text{and} \quad y(t_0) = x_0,$$

and that $y(t) \in (a, b)$ for all $t \in (S, T)$.

Then, for all $t \in (S, T)$, $y'(t) = v(y(t)) > 0$ and so $y(t)$ is strictly increasing monotone and hence it is an invertible function from (S, T) onto its range. Let $t(y)$ be the inverse function. By the Inverse Function Theorem, $t(y)$ is differentiable and

$$\frac{d}{dy}t(y) = \left(\frac{d}{dt}y(t) \right)^{-1} \Big|_{t=y(t)} = \frac{1}{v(y)}.$$

Then, by the Fundamental Theorem of Calculus,

$$t(y) - t(x_0) = \int_{x_0}^y \frac{1}{v(z)} dz$$

for all y such that $y = y(t)$ for some $t \in (S, T)$. That is, $t(y)$ is given by Barrows formula and hence $y(t)$ is the inverse of the function defined by Barrow's formula. In other words, $y(t) = x(t)$ on its domain of definition (S, T) . This proves the uniqueness of the solution - for as long as it stays inside (a, b) . \square

2.2.2 Existence and Uniqueness for First Order Autonomous Equations

Let v be continuous and strictly positive on (a, b) with $v(a) = v(b) = 0$. We can apply Barrows formula to compute T_a and T_b subject to $x(t_0) = x_0$ for any $x_0 \in (a, b)$. For example,

$$T_b = \int_{x_0}^b \frac{1}{v(z)} dz.$$

However, since $v(b) = 0$,

$$T_b = t_0 + \lim_{x \rightarrow b} \int_{x_0}^x \frac{1}{v(z)} dz$$

which is either convergent or divergent. If it is divergent,

$$\lim_{x \rightarrow b} \int_{x_0}^x \frac{1}{v(z)} dz = \infty$$

in which case it takes an infinitely long time for the solutions $x(t)$ to reach the endpoint b of the maximal interval. Thus, the solution exists and remains in (a, b) for all $t > t_0$ if, and only if,

$$\int_{x_0}^b \frac{1}{v(z)} dz < \infty.$$

A similar analysis shows that $T_a > -\infty$, and thus the solution exists and remains in (a, b) for all $t < t_0$ if, and only if,

$$\int_{x_0}^a \frac{1}{v(z)} dz = - \int_a^{x_0} \frac{1}{v(z)} dz > -\infty.$$

Theorem 2.2. Let v be continuous and strictly positive on a maximal interval (a, b) . Then for any t_0 and any $x_0 \in (a, b)$, $x'(t) = v(x(t))$ with $x(t_0) = x_0$ has a unique solution that is defined for all t , and remains in (a, b) for all $t \in \mathbb{R}$ if, and only if,

$$\int_a^{x_0} \frac{1}{|v(z)|} dz = \int_{x_0}^b \frac{1}{|v(z)|} dz = \infty$$

Theorem 2.3. Let $v(x)$ be continuous on $(a - \delta, a + \delta)$, $\delta > 0$ and suppose that the only solution of $v(x) = 0$ in $(a - \delta, a + \delta)$ is $x = a$. Then, for any t_0 , there is a solution of $x'(t) = v(x(t))$ for all t if, and only if, either $v(x) > 0$ on $(a, a + \delta)$ and

$$\int_a^{a+\delta} \frac{1}{v(z)} dz < \infty$$

is satisfied, in which case there is a solution moving right or $v(x) < 0$ on $(a - \delta, a)$ and

$$\int_a^{a-\delta} \frac{1}{v(z)} dz < \infty$$

is satisfied, in which case there's a solution moving to the left. In particular, if

$$\int_{a-\delta}^a \frac{1}{|v(z)|} dz = \int_a^{a+\delta} \frac{1}{|v(z)|} dz = \infty,$$

then the steady state solution is the unique solution.

Consider the vector field $v : \mathbb{R} \rightarrow \mathbb{R}$ defined by

$$v(x) = \begin{cases} \sqrt{x^2 - 1}, & \text{if } x < -1 \\ \sqrt{1 - x^2}, & \text{if } -1 \leq x \leq 1 \\ \sqrt{x^2 - 1}, & \text{if } x > 1 \end{cases}$$

It is easy to check that v is continuous and that $(-\infty, -1)$, $(-1, 1)$, $(1, \infty)$ are max intervals for v . Consider the IVP $x'(t) = v(x(t))$, $x(0) = 0$. By Barrows formula,

$$\begin{aligned} t(x) &= t_0 + \int_{x_0}^x \frac{1}{v(z)} dz \\ &= \int_0^x \frac{1}{\sqrt{1 - z^2}} dz = \arcsin(x) \quad \forall x \in (-1, 1). \end{aligned}$$

So, $x(t) = \sin t$ for $t \in (-\frac{\pi}{2}, \frac{\pi}{2})$. Then

$$\lim_{x \rightarrow -\frac{\pi}{2}} x(t) = -1 \quad \text{and} \quad \lim_{x \rightarrow \frac{\pi}{2}} x(t) = 1$$

thus, $T_{-1} = -\frac{\pi}{2}$ and $T_1 = \frac{\pi}{2}$. So, on (T_{-1}, T_1) the solution is unique. However, there are infinitely many choices to continue the solution (T_{-1}, T_1) . Since $v(\pm 1) = 0$, the solution is instantaneously at rest at $x = \pm 1$. One solution is to let it stay at rest:

$$x(t) = \begin{cases} -1, & \text{if } t < -\frac{\pi}{2} \\ \sin t, & \text{if } -\frac{\pi}{2} \leq t \leq \frac{\pi}{2} \\ 1, & \text{if } t > \frac{\pi}{2} \end{cases}$$

For $x > 1$, we can integrate $\frac{1}{v(z)} = \frac{1}{\sqrt{x^2-1}}$ using Barrows Formula. Consider the IVP $x'(t) = v(x(t))$, $x(\frac{\pi}{2}) = 1$, for $t > \frac{\pi}{2}$. By Barrow's formula

$$t(x) = \frac{\pi}{2} + \int_1^x \frac{1}{\sqrt{z^2-1}} dz = \frac{\pi}{2} + \ln(x + \sqrt{x^2-1}).$$

Hence,

$$x + \sqrt{x^2-1} = e^{t-\pi/2}.$$

Let $a = e^{t-\pi/2}$.

$$\begin{aligned} \Rightarrow x + \sqrt{x^2-1} &= a \\ \Rightarrow \sqrt{x^2-1} &= a - x \\ \Rightarrow x^2-1 &= (a-x)^2 \\ \Rightarrow x &= \frac{a^2+1}{2a} = \cosh\left(t - \frac{\pi}{2}\right). \end{aligned}$$

We have another solution for the IVP:

$$x(t) = \begin{cases} -1, & \text{if } t < -\frac{\pi}{2} \\ \sin t, & \text{if } -\frac{\pi}{2} \leq t \leq \frac{\pi}{2} \\ \cosh t, & \text{if } t > \frac{\pi}{2} \end{cases}$$

2.2.3 Existence and Uniqueness within Maximal Intervals

Let (a, b) be a maximal interval for an L -Lipschitz vector field v on \mathbb{R} . For any $x_0 \in (a, b)$ the solution of the IVP $x'(t) = v(x(t))$, $x(t_0) = x_0$ exists and is well-defined for all $t \in \mathbb{R}$. In particular for every $t_0 \in \mathbb{R}$. In particular for every $t_0 \in \mathbb{R}$, $x_0 \in (a, b)$, there is a unique solution curve of the IVP passing through (t_0, x_0) .

Proof. For simplicity assume $v(x) > 0$ on (a, b) , so the solution of the IVP is increasing. We now show that $x(t) < b$ for all $t \in \mathbb{R}$. Since $x(t)$ is increasing, we only need to consider $t > t_0$, since otherwise $x(t) \leq x_0 < b$. There are two cases. Suppose $b = \infty$. Since v is L -Lipschitz for all $z > x_0$,

$$v(z) - v(x_0) \leq |v(z) - v(x_0)| \leq L|z - x_0| = L(z - x_0).$$

So, $v(z) \leq v(x_0) + L(z - x_0)$. By Barrows formula we have that

$$\begin{aligned} t(x) &= t_0 + \int_{x_0}^x \frac{1}{v(z)} dz \\ &\geq t_0 + \int_{x_0}^x \frac{1}{v(x_0) + L(z - x_0)} dz \\ &= t_0 + \frac{1}{L} \ln\left(\frac{1 + L(z - x_0)}{v(x_0)}\right) \end{aligned}$$

for all $x > x_0$. Since,

$$\lim_{x \rightarrow \infty} \frac{1}{L} \ln\left(\frac{1 + L(z - x_0)}{v(x_0)}\right) = \infty \Rightarrow \lim_{x \rightarrow \infty} t(x) = \infty$$

therefore, $x(t) < b = \infty$ for all t .

Now assume $b < \infty$. We know $v(b) = 0$ so by Thm (2.2) we need to show

$$\int_{x_0}^b \frac{1}{v(z)} dz = \infty.$$

The L -Lipschitz condition gives

$$\begin{aligned}v(z) &= v(z) - v(b) \\&= |v(z) - v(b)| \\&\leq L|z - b| = L(z - b)\end{aligned}$$

So, $v(z) \leq L(b - z)$ and

$$\int_{x_0}^b \frac{1}{v(z)} dz \geq \int_{x_0}^b \frac{1}{L(b - z)} dz = \infty.$$

Similar, $x(t) > a$ for all t (regardless if $a = -\infty$ or $a > -\infty$).

□