

3.2

Suppose A and B are operators on a finite-dimensional Hilbert space and suppose that $AB - BA = cI$ for some constant c . Show that $c = 0$.

Let $\dim H = n$.

So $\text{tr}(AB - BA) = \text{tr}(cI) = nc$.

$$(AB)_{ij} = \sum_{k=1}^n a_{ik} b_{kj}$$

$$(BA)_{ij} = \sum_{k=1}^n b_{ik} a_{kj}$$

$$(AB - BA)_{ij} = \sum_{k=1}^n (a_{ik} b_{kj} - b_{ik} a_{kj})$$

$$\begin{aligned} \text{tr}(AB - BA) &= \sum_{i=1}^n (AB - BA)_{ii} \\ &= \sum_{i=1}^n \sum_{k=1}^n (a_{ik} b_{ki} - b_{ik} a_{ki}) \\ &= \sum_{i=1}^n \sum_{k=1}^n a_{ik} b_{ki} - \sum_{i=1}^n \sum_{k=1}^n b_{ik} a_{ki} \\ &= \sum_{i=1}^n \sum_{k=1}^n a_{ki} b_{ik} - \sum_{i=1}^n \sum_{k=1}^n b_{ik} a_{ki} \\ &= \sum_{i=1}^n \sum_{k=1}^n a_{ki} b_{ik} - \sum_{i=1}^n \sum_{k=1}^n a_{ki} b_{ik} \\ &= 0 \end{aligned}$$

But $I \neq 0$.

$\therefore c = 0$

3.3

If A is a bounded operator on a Hilbert space H , then there exists a unique bounded operator A^* on H satisfying $\langle A\phi, \psi \rangle = \langle \phi, A^*\psi \rangle$ for all ϕ and ψ in H . The operator A^* is called the *adjoint* of A , and A is called *self-adjoint* if $A^* = A$.

- (a) Show that for any bounded operator A and constant $c \in \mathbb{C}$, we have $(cA)^* = \bar{c}A^*$, where \bar{c} is the complex conjugate of c .

Note: $\langle \phi, A\psi \rangle = \langle \phi, B\psi \rangle \implies A = B$

Assume $\phi, \psi \in H$.

$$\begin{aligned} \langle \phi, (cA)^*\psi \rangle &= \langle (cA)\phi, \psi \rangle \\ &= \langle A(c\phi), \psi \rangle \\ &= \langle c\phi, A^*\psi \rangle \\ &= c \langle \phi, A^*\psi \rangle \\ &= \langle \phi, \bar{c}(A^*\psi) \rangle \\ &= \langle \phi, (\bar{c}A^*)\psi \rangle \end{aligned}$$

$$\therefore (cA)^* = \bar{c}A^*$$

- (b) Show that if A and B are self-adjoint, then the operator $\frac{1}{i\hbar}[A, B]$ is also self-adjoint.

Note: $(A + B)^* = A^* + B^*$ and $(AB)^* = B^*A^*$

$$\begin{aligned} (i\hbar[A, B])^* &= [i\hbar(AB - BA)]^* \\ &= \overline{i\hbar}(AB - BA)^* \\ &= -i\hbar[(AB)^* - (BA)^*] \\ &= -i\hbar(B^*A^* - A^*B^*) \\ &= -i\hbar(BA - AB) \\ &= i\hbar(AB - BA) \\ &= i\hbar[A, B] \end{aligned}$$

Therefore $i\hbar[A, B]$ is self-adjoint.

3.5

Suppose that ψ is a unit vector in $L^2(\mathbb{R})$ such that the functions $x\psi(x)$ and $x^2\psi(x)$ also belong to $L^2(\mathbb{R})$. Show that:

$$\langle X^2 \rangle_\psi > \left(\langle X \rangle_\psi \right)^2$$

Let $a = \langle X \rangle_\psi$ and consider the integral:

$$\int_{-\infty}^{\infty} (x - a)^2 |\psi(x)|^2 dx$$

Note that the integrand is positive (for non-constant x) and thus so is the integral.

$$\begin{aligned} \int_{-\infty}^{\infty} (x - a)^2 |\psi(x)|^2 dx &= \int_{-\infty}^{\infty} (x^2 - 2ax + a^2) |\psi(x)|^2 dx \\ &= \int_{-\infty}^{\infty} x^2 |\psi(x)|^2 dx - 2a \int_{-\infty}^{\infty} x |\psi(x)|^2 dx + a^2 \int_{-\infty}^{\infty} |\psi(x)|^2 dx \\ &= \langle X^2 \rangle_\psi - 2 \langle X \rangle_\psi \langle X \rangle_\psi + \left(\langle X \rangle_\psi \right)^2 \cdot 1 \\ &= \langle X^2 \rangle_\psi - 2 \left(\langle X \rangle_\psi \right)^2 + \left(\langle X \rangle_\psi \right)^2 \\ &= \langle X^2 \rangle_\psi - \left(\langle X \rangle_\psi \right)^2 \\ &> 0 \end{aligned}$$

$$\therefore \langle X^2 \rangle_\psi > \left(\langle X \rangle_\psi \right)^2$$

3.6

Consider the Hamiltonian \hat{H} for a quantum harmonic oscillator, given by:

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{k}{2} x^2$$

where j is the spring constant of the oscillator. Show that the function

$$\psi_0(x) = e^{-\frac{\sqrt{km}}{2\hbar} x^2}$$

is an eigenvector for \hat{H} with eigenvalue $\frac{\hbar\omega}{2}$ where $\omega = \sqrt{\frac{k}{m}}$ is the classical frequency of the oscillator.

$$\frac{d}{dx} \psi_0(x) = \frac{d}{dx} e^{-\frac{\sqrt{km}}{2\hbar} x^2} = -\frac{\sqrt{km}}{h} x e^{-\frac{\sqrt{km}}{2\hbar} x^2}$$

$$\frac{d^2}{dx^2} \psi_0(x) = -\frac{\sqrt{km}}{h} \left(e^{-\frac{\sqrt{km}}{2\hbar} x^2} - \frac{\sqrt{km}}{h} x^2 e^{-\frac{\sqrt{km}}{2\hbar} x^2} \right) = -\frac{\sqrt{km}}{h} \psi_0(x) + \frac{km}{\hbar^2} x^2 \psi_0(x)$$

$$\begin{aligned}
\hat{H}\psi_0(x) &= \left(-\frac{\hbar^2}{2m}\frac{d^2}{dx^2} + \frac{k}{2}x^2\right)\psi_0(x) \\
&= -\frac{\hbar^2}{2m}\frac{d^2}{dx^2}\psi_0(x) + \frac{k}{2}x^2\psi_0(x) \\
&= -\frac{\hbar^2}{2m}\left(-\frac{\sqrt{km}}{h}\psi_0(x) + \frac{km}{\hbar^2}x^2\psi_0(x)\right) + \frac{k}{2}x^2\psi_0(x) \\
&= \frac{\hbar}{2}\sqrt{\frac{k}{m}}\psi_0(x) - \frac{k}{2}x^2\psi_0(x) + \frac{k}{2}x^2\psi_0(x) \\
&= \frac{\hbar\omega}{2}\psi_0(x)
\end{aligned}$$