# CHEM / BCMB 4190/6190/8189

# **Introductory NMR**

Lecture 8

#### **Relaxation:**

#### Applying and rf pulse:

- alters the population ratios
- produces transverse magnetic field components  $(M_x$  and  $M_v)$

The system returns to equilibrium via two relaxation processes:

T<sub>1</sub> Relaxation: in the static field direction, spin-lattice or longitudinal.

T<sub>2</sub> Relaxation: perpendicular to the static field, spin-spin or transverse.

## **Spin-Lattice Relaxation:**

Energy absorbed from the rf pulse is transferred to the surroundings, lattice and neighboring molecules.

<sup>1</sup>H: order of seconds.

<sup>13</sup>C: vary from milliseconds(large molecules) to hundreds of seconds.

# Factors contributing to $T_1$ relaxation:

- dipole-dipole relaxation (ie NOE).
- spin-rotation relaxation
- chemical shift anisotropy.
- scalar coupling.
- electric quadrupolar relaxation
- unpaired electrons in paramagnetic compounds.

$$\tau_c$$
= Correlation time  $\propto 1/T_1$ 

- interval between two successive rotations.
- the faster a molecule moves the greater the  $T_1$

# Relaxation:

T<sub>1</sub>: spin-lattice relaxation time (longitudinal)

T<sub>2</sub>: spin-spin relation time (transverse)

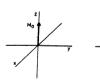
After we remove the pulse, the magnetization vector returns to equilibrium.

Bloch assumed that the processes were first-order and defined them by  $T_1$  and  $T_2$ .

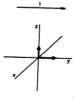
$$dM_z/dt = -M_z - M_o/T_1$$
so
$$M_z = M_0(1-e^{-t/T_1})$$

$$dM_{x'}/dt = -M_{x'}/T_2$$

$$dM_{y'}/dt = -M_{y'}/T_2$$



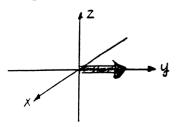


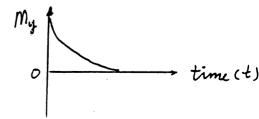






# On-resonance





# Experimental Determination of T<sub>1</sub>: **( 3** ¹H 180% $\tau = 0.70_{s}$ $\tau = 0.23$ s

3.8

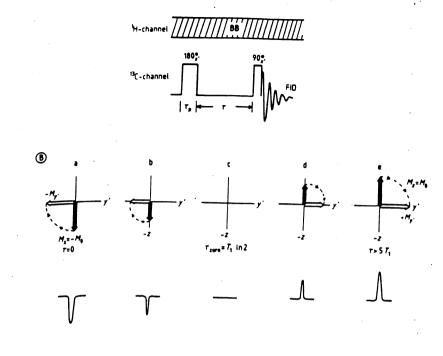
3.6

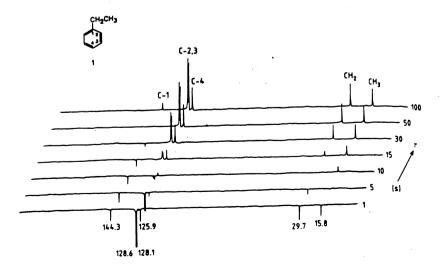
3.4

3.2

2.8

2.6





How do we determine the  $T_1$  value?

$$\mathbf{M_0 - M_z} = \mathbf{A} \mathbf{e}^{-\tau/\mathrm{T}1}$$

After a  $180^{\circ}$  pulse during the inversion recovery experiment  $M_0 = -\dot{M_z}$  so  $A = 2M_0$ :

$$M_0 - M_z = 2M_0 e^{-\tau/T_1}$$

$$\ln (M_0 - M_z) = \ln 2M_0 - \tau / T_1$$

Substituting I for M where I is intensity of signal:

$$\ln (I_0 - I_z) = \ln 2I_0 - \tau / T_1$$

$$\mathbf{T}_{\mathrm{zero}} = \mathbf{T}_1 \ln 2$$

 $T_1$  values can vary greatly from one nuclei to the next in the same molecule.

Need to wait  $5T_1$  periods for system to return to equilibrium so:

$$(5 T_1 - 180^{\circ}_{x'} - \tau - 90^{\circ}_{x'} - FID)_n$$

Why measure  $T_1$ ?

- $-\pi/2$  pulse elicits maximal NMR signal at thermal equilibrium.
- for repetitive signal averaging this is not always the best pulse angle.
- need to wait about 5T<sub>1</sub> for z-magnetization to return to equilibrium.
- more efficient to repeat sooner.

#### **Ernst Angle:**

$$\cos \alpha_{\rm E} = e^{-Tr/T1}$$

 $T_r$  = the delay between pulses.

- best value of  $\alpha$  actually depends on offsets of signals in spectrum but this is an average value.

If stick with  $\pi/2$  pulses then:

the optimum repetition rate is  $T_r = 1.27 T_1$ 

- in this case you are about 80% efficient.

# <u>Influence of protons on <sup>13</sup>C T<sub>1</sub> values:</u>

- main contribution is via dipole-dipole interactions, so directly bonded protons can have large effects.
- increasing the number of protons shortens the T1 although the correlation is not always direct.

# Influence of molecular size on T<sub>1</sub>:

- T1 values decrease as the size of the molecules

## **Suppression of Water Signal:**

$$(180^{\circ}_{x'} - \tau - 90^{\circ}_{x'} - FID)_{n}$$

 $T_1$  for  $H_2O$  is ~ 3sec and thus is greater than  $T_1$  of most organic molecules so one can chose  $\tau$  to minimize solvent signal.

## Spin-Spin Relaxation: (T2)

After  $90^\circ_{x'}$  pulse  $M_z=0$  and net magnetization is along the y-axis. The population of  $N_\alpha$  and  $N_\beta$  are equal .

Phase Coherence: a small proportion of nuclear dipoles are bunched along the y-axis.

 $T_2$  - determines how fast  $M_{x'}$  and  $M_{y'}$  return to zero.

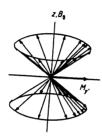
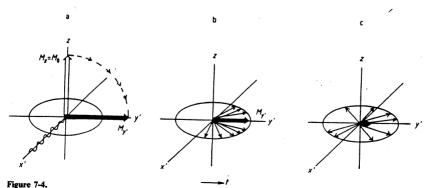


Figure 1-12.
Pictorial representation of phase coherence: after a 90°, pulse a fraction of the nuclear spins (not all!) are bunched together in phase as they precess about the field direction 2.



Decay of the transverse magnetization  $M_y$  (in the rotating coordinate system x', y', z) after a  $90_x^o$  pulse, as the precessing bunched spins fan out as a result of magnetic field inhomogeneities.

#### How do we view this?:

Nuclei change from one energy state to the other and we lose the phase coherence.

- this is not the main contribution.

Magnetic Field inhomogeneities: there is a small  $\Delta B_0$  throughout the sample which leads to fanning out of signals.

- -fanning out process causes signals that are chemically equivalent to precess with slightly different Larmor frequencies.
- the net result is an increase in the line width of the various signals.

By definition though  $T_1 \ge T_2$ 

#### Experimental Determination of T<sub>2</sub>:

Want to measure a true value which is independent of the  $\,\Delta B_0$  fanning out process.

#### Spin-echo experiment:

After initial  $90^{\circ}_{x'}$  pulse you apply a  $180^{\circ}_{x'}$  pulse at set interval  $\tau$ .

At  $2\tau$  - resultant transverse magnetization is -y' direction.

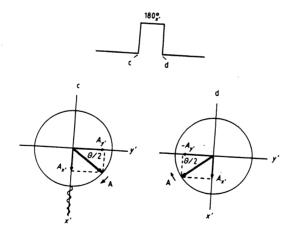
 $4\tau$  - magnetization is y' direction.

6τ - magnetization is -y' direction.

Etc.....we get an echo with alternating phase at intervals of  $2\tau$ .

However, the net magnetization is decreasing with each echo due to true spin-spin relaxation.

 $90_{x'}^{\circ} - \tau - 180_{x'}^{\circ} - \tau (1st echo) - \tau - 180_{x'}^{\circ} - \tau (2nd echo) \dots$   $90_{x'}^{\circ} \cdot 180_{x'}^{\circ} \cdot 180_{x'}^{\circ$ 



- the  $180^\circ_{x'}$  pulse turns the  $A_{y'}$  component into the (-y') direction but has not effect on the  $A_{x'}$ . The angle between vector A and the y'-axis is unchanged if we disregard sign
- the effect of a  $180^{\circ}_{\,x'}$  pulse is that vectors A and B undergo reflections in the x', z' plane.

Quantitatively:

$$\mathbf{M}_{\mathbf{v}'} = \mathbf{A} \mathbf{e}^{-t/T2}$$

At t=0,  $A=M_0$ 

$$\ln M_v = \ln M_0 - t/T_2$$

since I  $\alpha$   $M_{y}$ .

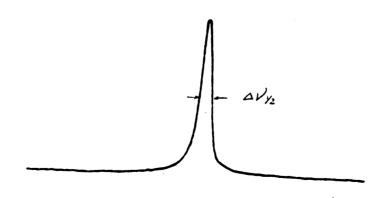
$$\ln I(t) = \ln I_0 - t/T_2$$

plot I(t) vs t where I(t) is intensities at each echo. Slope of line is -1/ $T_2$ .

For the most part we do not concern ourselves with T<sub>2</sub> measurement.

Of more consequence is  $T_2^*$  which takes into account the inhomogeneity of the magnetic field.

Line Widths of NMR Signals



$$\Delta v_{V2} = \frac{1}{\pi T_2^*}$$

$$\frac{1}{T_2^*} = \frac{\gamma \, \Delta B_0}{2} + \frac{1}{T_2}$$