

# NMR Active Isotopes Exist for Nearly Every Element

<http://bouman.chem.georgetown.edu/NMRpt/NMRPerTab.html>

Select an element by clicking on it:

1 2 3 4 5 6 7 8 9 10 11 12 13 14 15 16 17 18

H																H	He
<u>Li</u>	<u>Be</u>											<u>B</u>	<u>C</u>	<u>N</u>	<u>O</u>	<u>F</u>	<u>Ne</u>
<u>Na</u>	<u>Mg</u>											<u>Al</u>	<u>Si</u>	<u>P</u>	<u>S</u>	<u>Cl</u>	<u>Ar</u>
<u>K</u>	<u>Ca</u>	<u>Sc</u>	<u>Ti</u>	<u>V</u>	<u>Cr</u>	<u>Mn</u>	<u>Fe</u>	<u>Co</u>	<u>Ni</u>	<u>Cu</u>	<u>Zn</u>	<u>Ga</u>	<u>Ge</u>	<u>As</u>	<u>Se</u>	<u>Br</u>	<u>Kr</u>
<u>Rb</u>	<u>Sr</u>	<u>Y</u>	<u>Zr</u>	<u>Nb</u>	<u>Mo</u>	<u>Tc</u>	<u>Ru</u>	<u>Rh</u>	<u>Pd</u>	<u>Ag</u>	<u>Cd</u>	<u>In</u>	<u>Sn</u>	<u>Sb</u>	<u>Te</u>	<u>I</u>	<u>Xe</u>
<u>Cs</u>	<u>Ba</u>	*	<u>Hf</u>	<u>Ta</u>	<u>W</u>	<u>Re</u>	<u>Os</u>	<u>Ir</u>	<u>Pt</u>	<u>Au</u>	<u>Hg</u>	<u>Tl</u>	<u>Pb</u>	<u>Bi</u>	Po	At	Rn
<u>Fr</u>	<u>Ra</u>	**	<u>Rf</u>	<u>Ha</u>	<u>Sg</u>	<u>Ns</u>	<u>Hs</u>	<u>Mt</u>									

* <u>La</u>	Ce	<u>Pr</u>	<u>Nd</u>	Pm	<u>Sm</u>	<u>Eu</u>	<u>Gd</u>	<u>Tb</u>	<u>Dy</u>	<u>Ho</u>	<u>Er</u>	<u>Tm</u>	<u>Yb</u>	<u>Lu</u>
** <u>Ac</u>	Th	<u>Pa</u>	<u>U</u>	Np	<u>Pu</u>	<u>Am</u>	<u>Cm</u>	<u>Bk</u>	<u>Cf</u>	<u>Es</u>	<u>Fm</u>	<u>Md</u>	<u>No</u>	<u>Lr</u>

# Review of Spin Properties

- NMR active nuclei possess an intrinsic angular momentum,  $\vec{I}$ , known as the spin angular momentum. The magnitude is

$$|\vec{I}| = \hbar [ I (I+1) ]^{1/2}$$

- Here  $I$  is the nuclear spin quantum number (integral or half-integral). If  $I = 0$ , no spin angular momentum (not NMR active)
- Associated with  $I$  is a magnetic moment,  $\vec{\mu}$ 
$$\vec{\mu} = \gamma \vec{I} = \gamma \hbar [ I (I+1) ]^{1/2}$$
- The proportionality constant is the gyromagnetic ratio,  $\gamma$
- In NMR, larger  $\vec{\mu}$  for given  $\vec{I}$  (large  $\gamma$ ), means more sensitive nucleus

# Review of Spin Properties

- In a magnetic field, otherwise degenerate (energetically equivalent) states split into nondegenerate states (known as Zeeman splitting)
- The states are quantized, with the number of states established by the spin quantum number,  $l$

$$\# \text{ levels} = 2l + 1$$

- Each of the  $2l+1$  states/levels is associated with a magnetic quantum number,  $m$

$$m = -l, -l+1, \dots, l-1, l$$

- The component of  $l$  along the  $z$  axis,  $l_z$ , is defined as follows

$$l_z = m\hbar$$

- Thus

$$\mu_z = \gamma l_z = m\gamma \hbar$$

# Review of Spin Properties

- The energies of the states resulting from the interaction of the magnetic moment with a magnetic field,  $\vec{B}$  are given by

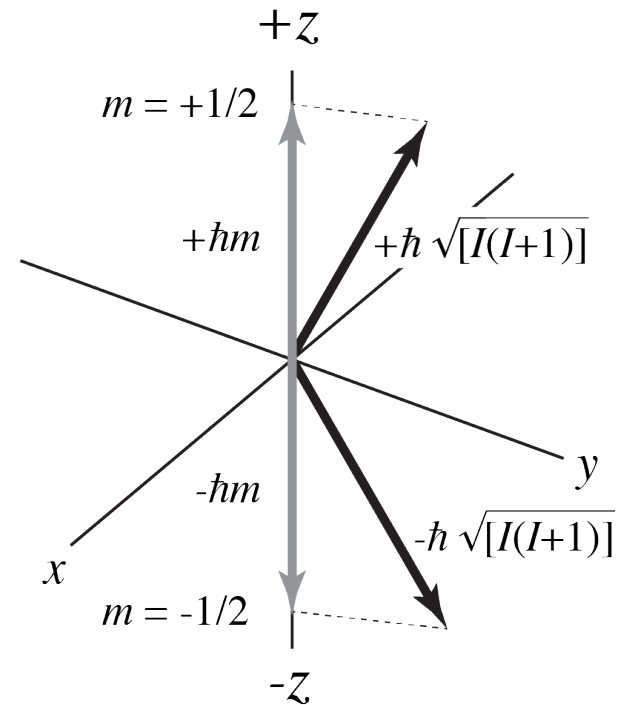
$$E = -\vec{\mu} \cdot \vec{B}$$

- The energies of the states depend on the orientations of the moments in the magnetic field, hence are proportional to the *scalar* projection of  $\vec{\mu}$  on  $\vec{B}$  (the dot product),  $\mu_z$

$$E = -\mu_z B_0 = -m\gamma \hbar B_0$$

- Here  $B_0$  is the magnetic field strength
- The 2/+1 energy levels are equally spaced. The energy difference between any two adjacent levels is

$$\Delta E = \gamma \hbar B_0$$



classical view of directional quantization for spin  $\frac{1}{2}$  nuclei

# Review of Spin Properties

- The torque exerted by  $B_0$  on the magnetic moments/dipoles promotes precession about the z-axis at a frequency given by

$$\nu_L = \gamma B_0 / (2\pi) \text{ (Larmor frequency, in Hz)}$$

$$\omega_0 = \gamma B_0 \text{ (radians/sec)}$$

- The energy difference between energy (spin) states can then be written as

$$\Delta E = h\nu_L$$

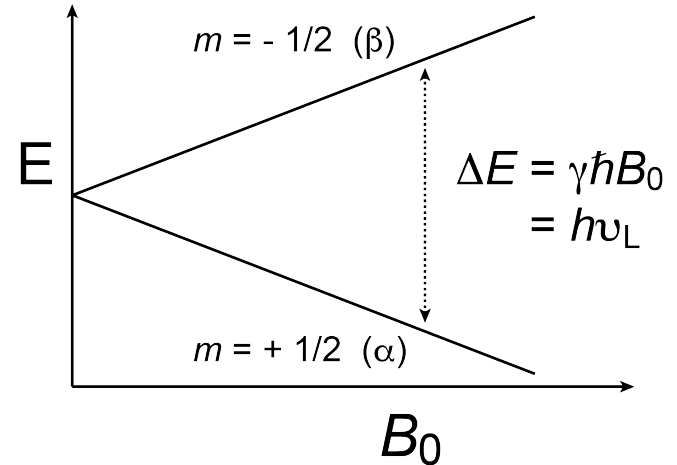
- Transitions between energy (spin) states can be effected by an electromagnetic field with an energy equal to  $\Delta E$ . This occurs when the frequency of that field,  $\nu_1$ , is equal to the Larmor frequency (*resonance condition*).

$$\nu_1 = \nu_L$$

# Review of Spin Properties

- For spin  $1/2$  ( $I = 1/2$ ), there are  $2I+1 = 2$  energy levels, with values of  $m$  equal to  $+1/2$  and  $-1/2$ , called  $\alpha$  and  $\beta$ , with energies

$$E_{\alpha} = -\frac{1}{2} \gamma \hbar B_0 \quad E_{\beta} = +\frac{1}{2} \gamma \hbar B_0$$



- From Boltzman statistics, the population ratio of these states can be estimated

$$\frac{N_{\beta}}{N_{\alpha}} = \exp\left(\frac{-\Delta E}{k_B T}\right) \approx 1 - \left(\frac{\Delta E}{k_B T}\right) \approx 1 - \left(\frac{\gamma \hbar B_0}{k_B T}\right)$$

- example:  $^1\text{H}$ , 300 °K, 5.875 Tesla (250 MHz)

$$\frac{N_{\beta}}{N_{\alpha}} = 1 - \frac{26.7519 \times 10^7 \times 1.0546 \times 10^{-27} \times 5.875}{1.3805 \times 10^{-16} \times 300} = 0.99996$$

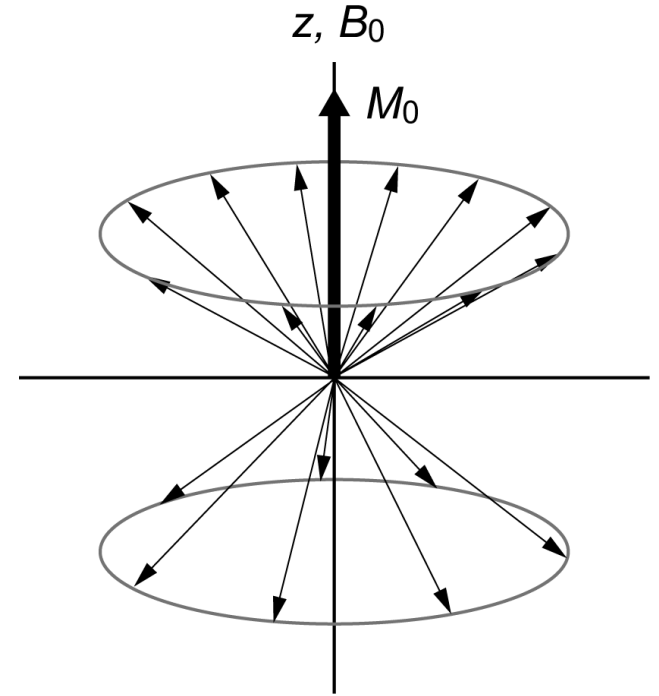
- $\Delta E$  is small, so the populations of  $\alpha$  and  $\beta$  are nearly equal, and the macroscopic magnetization is small: *NMR is insensitive*

# Review of Spin Properties

- The sum of the z-components of the nuclear dipoles in an ensemble gives the macroscopic (bulk) magnetization,  $M_0$

$$M_0 = \gamma \hbar \sum_{m=-I}^I m N_m \quad (\text{recall } \mu_z = m \gamma \hbar)$$

$$M_0 \approx \frac{N \gamma^2 \hbar^2 B_0}{k_B T (2I + 1)} \sum_{m=-I}^I m^2 \approx \frac{N \gamma^2 \hbar^2 B_0 I(I + 1)}{3 k_B T}$$



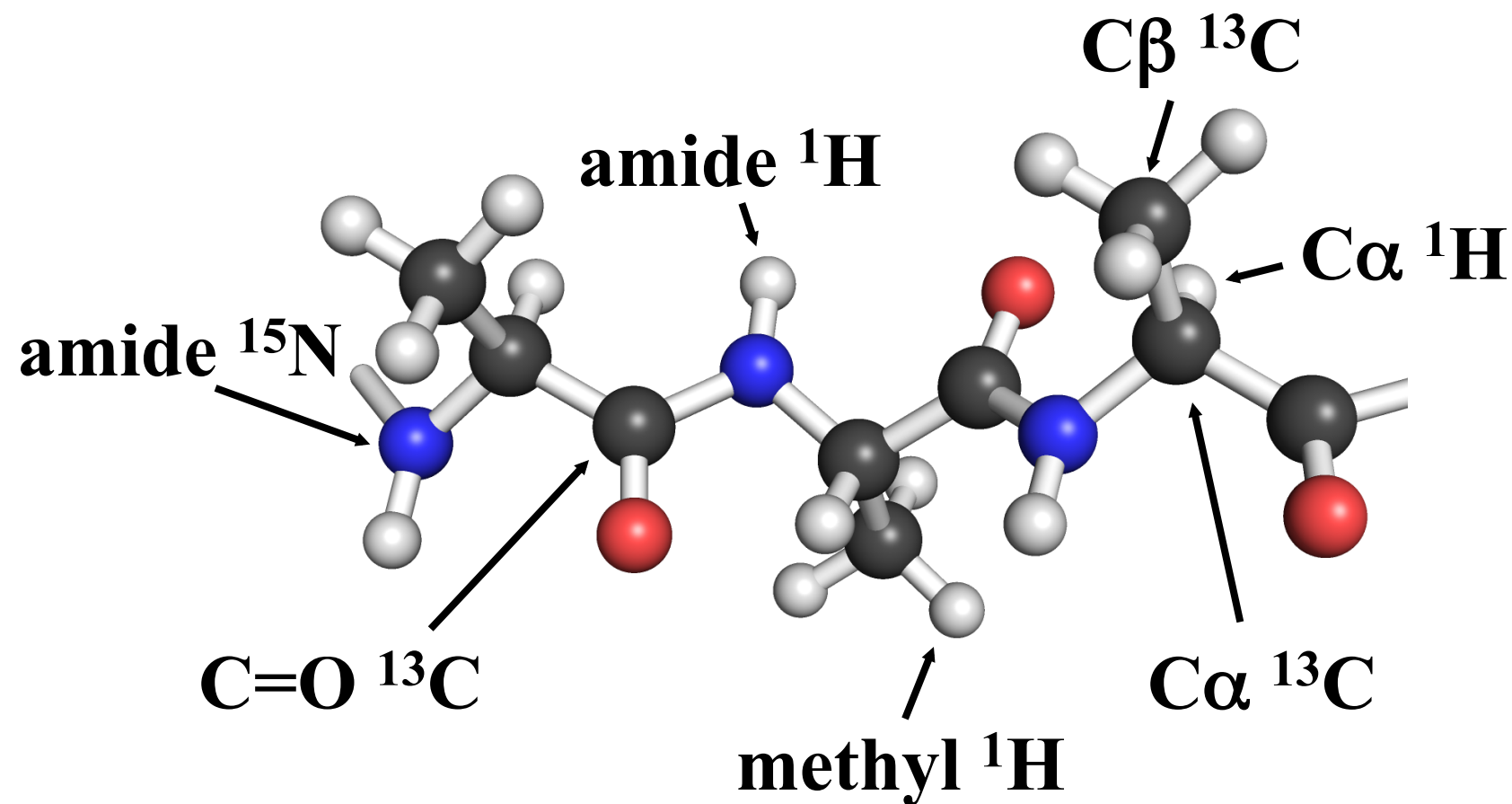
- Note: dependence on  $\gamma^2$ , linear dependence on  $B_0$ , dependence on isotopic abundance ( $N$ )

# Spin $\frac{1}{2}$ Nuclei are Most Useful in Biomolecular NMR

	$^1\text{H}$	$^{13}\text{C}$	$^{15}\text{N}$	$^{19}\text{F}$	$^{31}\text{P}$
Spin	1/2	1/2	1/2	1/2	1/2
Natural abundance	99.985%	1.108%	0.37%	100%	100%
Magnetogyric ratio ( $\gamma/10^7$ , rad T $^{-1}$ s $^{-1}$ )	26.7519	6.7283	-2.7126	25.1815	10.8394
Relative sensitivity	1.00	$1.59 \times 10^{-2}$	$1.04 \times 10^{-3}$	0.83	$6.63 \times 10^{-2}$
Relative receptivity	1.00	$1.76 \times 10^{-4}$	$3.85 \times 10^{-6}$	0.83	$6.63 \times 10^{-2}$
Magnetic moment ( $\mu/\mu_{\text{N}}$ )	4.8372	1.2166	-0.4903	4.5532	1.9601
Quadrupole moment	0	0	0	0	0
Resonance frequency (MHz)	100	25.144	10.133	94.077	40.481



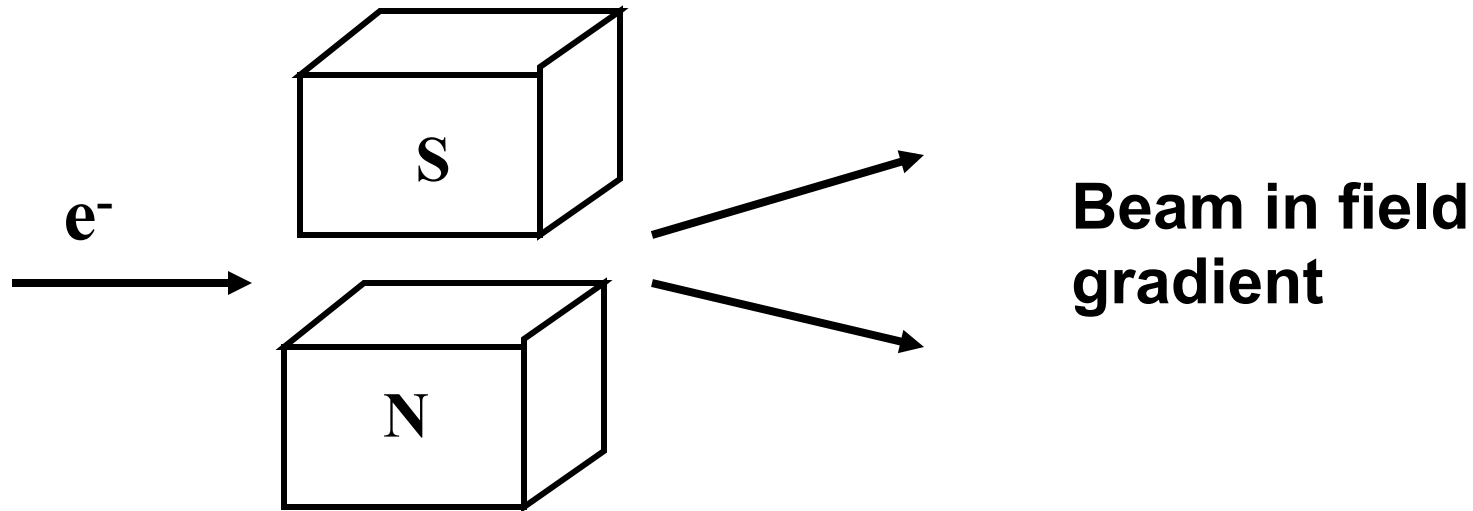
# Polypeptides are Rich in NMR Active Nuclei



# Nuclear Properties

- Not all nuclei have magnetic moments, Why?
- Not all nuclei are equally abundant, Why?
- Spins vary, Why?
- Magnetogyric ratios vary, Why?

# Fundamental Particle Properties

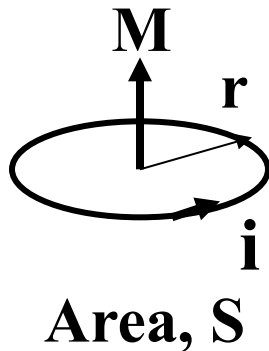


## Stern Gerlach experiment:

- demonstrated particles (electrons) possess an intrinsic angular momentum, and it is quantized
- Na atom - 1 unpaired electron  
Two spots implies quantized moments:  $\pm 1/2$   
protons and neutrons are also spin  $1/2$  particles

# Understanding Magnetic Moments

- Current Loop Model: classical analogy to connect “spin” to magnetic moment
- Can get reasonable estimate of  $\gamma$  for electron



$$\vec{M} = i \vec{S}$$

**$i$  in Coulomb  $s^{-1}$**

**$S$  in  $m^2$**

**$M$  in  $JT^{-1}$  (Tesla)**

**Estimates:  $i = -ev/(2\pi r)$ ,  $S = \pi r^2$ ,  $M$  (or  $\mu$ ) =  $-erv/2$**

$$\vec{\mu} = -e(\vec{r} \times \vec{v})/2, \quad \vec{L} = m_e \vec{r} \times \vec{v}, \quad \vec{\mu} = -e/(2m_e) \vec{L} = \gamma \vec{L} = \gamma \hbar/(2\pi) l$$

**$\gamma = -g (e/(2m_e))$ ,  $g$  = Lande  $g$  factor**

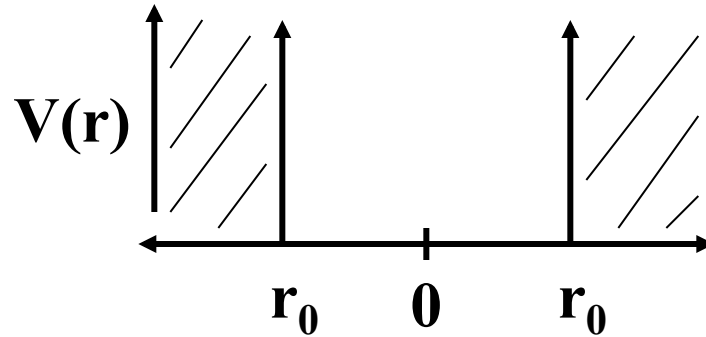
# Values of Particle Magnetogyric Ratios

**Electron:  $g \approx 2$ ,  $\gamma_e = -17.7 \times 10^{10} \text{ T}^{-1}\text{s}^{-1}$**

**Proton: expect  $1/m_p$  dependence,  $1/2000$  and positive  
 $2.7 \times 10^8 \text{ T}^{-1}\text{s}^{-1}$**

**Neutron: similar mass to proton  
 $-1.8 \times 10^8 \text{ T}^{-1}\text{s}^{-1}$**

# Heavier Nuclei: the Shell Model



**Analogous to shell model for atomic electrons**

**Some rules:**

**a) spherical particle in a box potential**

$$\psi = R_{nl}(r) Y_l^m(\theta, \phi), \quad E(n, l)$$





**ladder of energy levels like H atom, but all  $l$ s allowed  
 $l=0, 1, 2, 3$  for “s”, “p”, “d”, and “f” like atomic case**

**b) strong coupling of spin and orbit angular momentum  
quantized total:  $j = l \pm 1/2$  for spin  $1/2$  particle  
larger  $j$ , lower energy (usually)**

# Shell Model Rules Continued

- c) Treat protons and neutrons separately and fill from bottom up assuming  $2j + 1$  degeneracy**
- d) Assume particle pair strongly within levels: only unpaired spins count - total spin angular momentum given by  $j$  of level for unpaired spin**
- e) sign of moment depends on sign of moment for fundamental particle ( $+\frac{1}{2}$  for proton,  $-\frac{1}{2}$  for neutron) but changes sign when moment subtracts instead of adds to  $l$  in giving  $j$**

# Energy Level Diagram

$n+1$		$j$ ( $j = l \pm \frac{1}{2}$ )	degeneracy ( $2j + 1$ )	total
2s ( $l=0$ )		$1/2$	2	20
1d ( $l=2$ )		$3/2$	4	
		$5/2$	6	
1p ( $l=1$ )		$1/2$	2	8
		$3/2$	4	
1s ( $l=0$ )		$1/2$	2	2



**Example:**  $^{13}_6\text{C}$  ( 6 protons, 7 neutrons )

- unpaired neutron ( $-1/2$ ) in  $1p_{1/2}$  ( $j = 1 - 1/2 = 1/2$ ), so  
spin= $1/2$ , positive  $\gamma$













$n+1$		protons	neutrons	$j$ ( $j = l \pm 1/2$ )	degeneracy ( $2j + 1$ )	total
2s ( $l=0$ )				$1/2$	2	20
1d ( $l=2$ )				$3/2$	4	
				$5/2$	6	
1p ( $l=1$ )				$1/2$	2	8
				$3/2$	4	
1s ( $l=0$ )				$1/2$	2	2

**Example:**  $^{15}_7\text{N}$  ( 7 protons, 8 neutrons )

- unpaired proton (+1/2) in  $1p_{1/2}$  ( $j = 1 - 1/2 = 1/2$ ), so  
spin=1/2, negative  $\gamma$

$n+1$		protons	neutrons	$j$ ( $j = l \pm \frac{1}{2}$ )	degeneracy ( $2j + 1$ )	total
2s ( $l=0$ )				1/2	2	20
1d ( $l=2$ )				3/2	4	
				5/2	6	
1p ( $l=1$ )				1/2	2	8
				3/2	4	
1s ( $l=0$ )				1/2	2	2

**Example:**  $^{16}_8\text{O}$  ( 8 protons, 8 neutrons, two magic numbers ), spin = 0  
 - highly stable (99.76% of all oxygen on Earth)

n+1			j	degeneracy	total
	protons	neutrons	( j = l ± ½ )	( 2j + 1 )	
2s (l=0)			1/2	2	20
1d (l=2)			3/2	4	
			5/2	6	
1p (l=1)			1/2	2	8
			3/2	4	
1s (l=0)			1/2	2	2

# Particle Physics / Spin

## **Proton Spin Mystery Gains a New Clue:**

<https://www.scientificamerican.com/article/proton-spin-mystery-gains-a-new-clue1/>