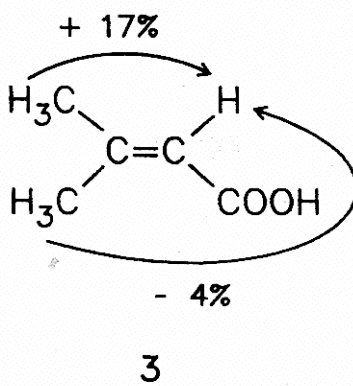
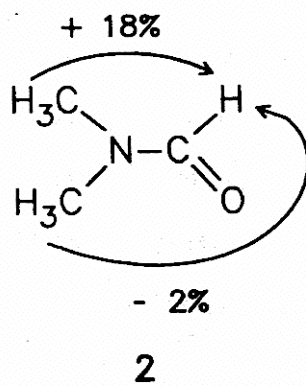
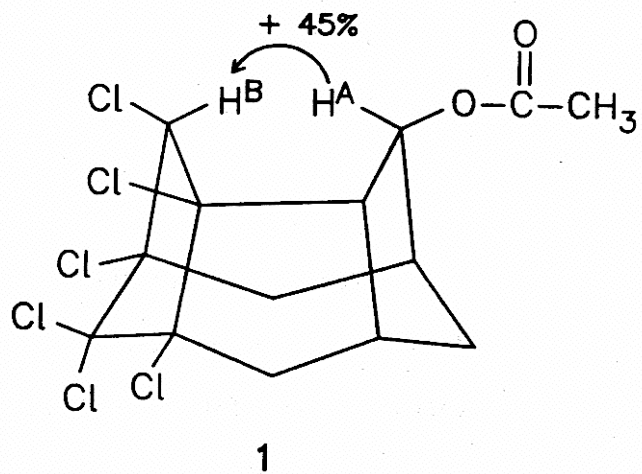


**CHEM / BCMB 4190/6190/8189**

**Introductory NMR**

Lecture 18



## The Nuclear Overhauser Effect

**NOE: Nuclear Overhauser Effect or Nuclear Overhauser Enhancement**

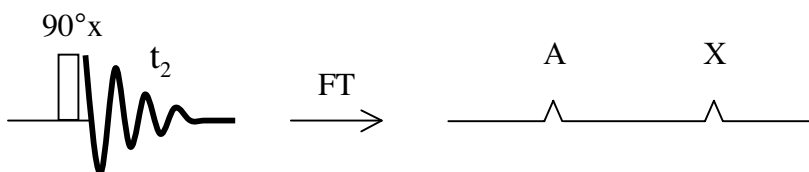
- Depends on dipole-dipole relaxation (an important mechanism of  $T_1$  relaxation)
- Is based on through-space interactions that change  $M_z$
- The NOE changes the intensities of particular signals
- The magnitude of the NOE is related to nuclear distances and molecular motions

### 1. 1D steady-state NOE experiment:

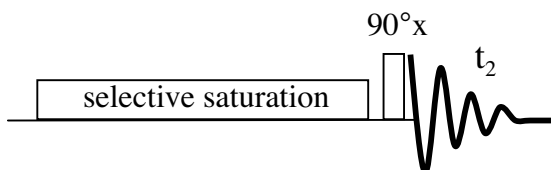
Consider a two-spin system AX:

- Assume the homonuclear case (A and X are both  $^1\text{H}$ , although they need not be)
- Assume no scalar coupling between A and X ( $J = 0$ )
- Assume the internuclear A-X distance is  $\leq \sim 5\text{\AA}$  (NOE enhancement depends on  $1/r^6$  where  $r$  is internuclear distance)

So, the simple 1D spectrum gives 2 signals with no splitting ( $J=0$ ):



To establish the NOE, we will selectively saturate the 'A' transition with an RF field to equilibrate the spin populations:



What happens to the population levels when we saturate one of the signals, and what happens afterwards?

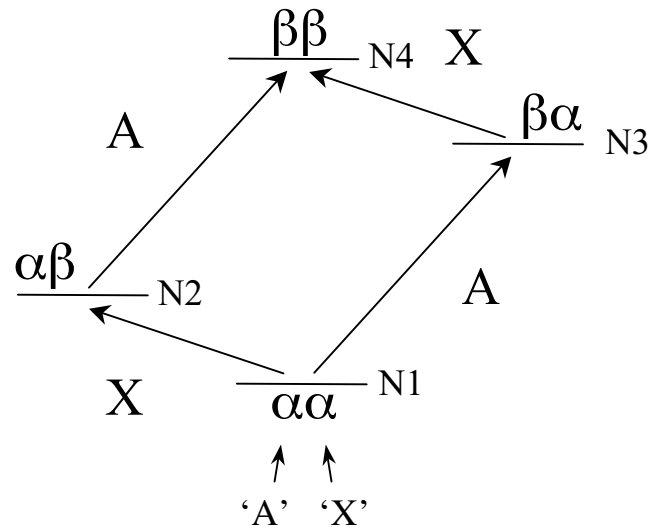
-Let's look at the energy level diagrams:

A). At equilibrium:

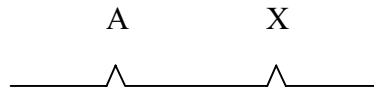
$$\begin{aligned} N4 &= N \\ N3 &= N + \Delta H \\ N2 &= N + \Delta H \\ N1 &= N + 2\Delta H \end{aligned}$$

'A' transitions:  $N1 - N3 = N2 - N4 = \Delta H$

'X' transitions:  $N1 - N2 = N3 - N4 = \Delta H$



- 'A' and 'X' transitions are allowed and observable
- 'A' and 'X' transitions give singlets in 1D spectrum

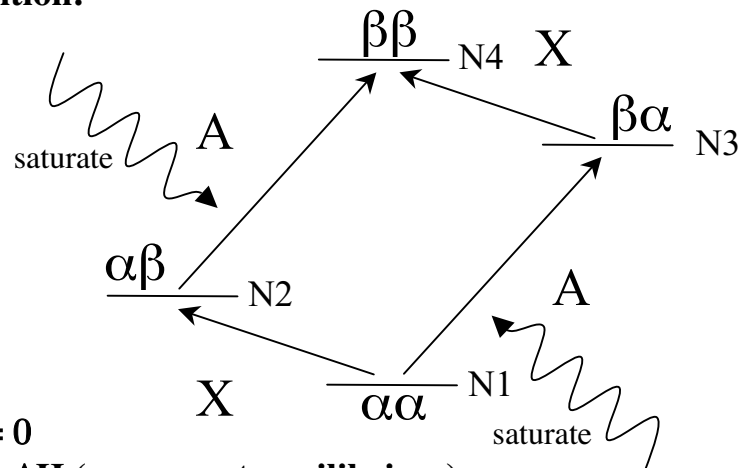


B). After saturation of the 'A' transition:

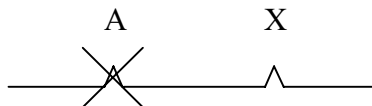
$$\begin{aligned} N4 &= N + \Delta H/2 \\ N3 &= N + 3\Delta H/2 \\ N2 &= N + \Delta H/2 \\ N1 &= N + 3\Delta H/2 \end{aligned}$$

'A' transitions:  $N1 - N3 = N2 - N4 = 0$

'X' transitions:  $N1 - N2 = N3 - N4 = \Delta H$  (same as at equilibrium)

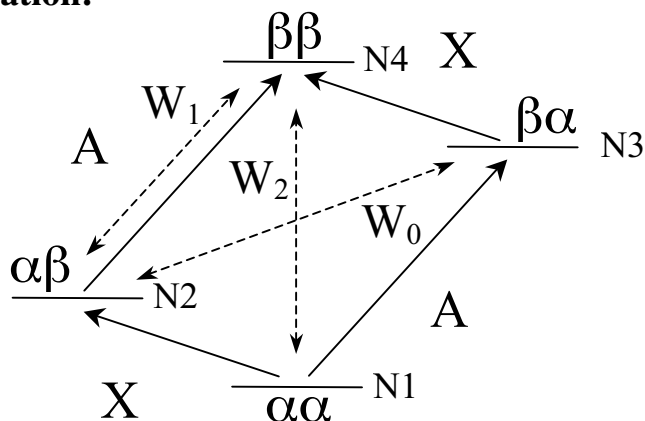


-populations of N1, N2, N3, and N4 are now different than they were at equilibrium



C). Now, allow relaxation during saturation:

	at equilibrium	after saturation
N4 =	N	$N + \Delta H/2$
N3 =	$N + \Delta H$	$N + 3\Delta H/2$
N2 =	$N + \Delta H$	$N + \Delta H/2$
N1 =	$N + 2\Delta H$	$N + 3\Delta H/2$



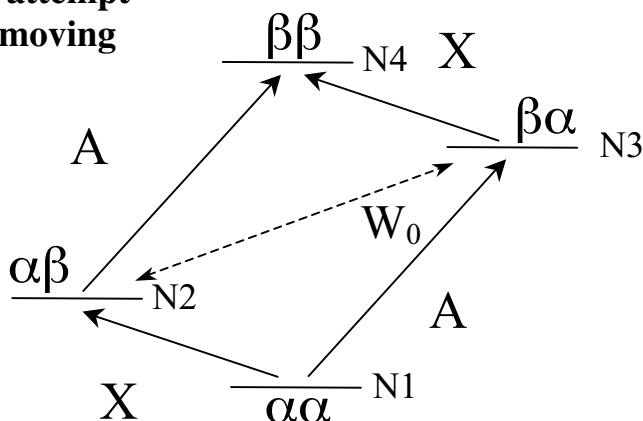
- $W_0$ ,  $W_1$ , and  $W_2$  are *rate constants or transition probabilities*
- $W_1$  is a one-spin transition or single quantum transition
- $W_0$  and  $W_2$  are two-spin transitions
- $W_0$  is a zero quantum transition
- $W_2$  is a double quantum transition
- $W_0$  and  $W_2$  are spectroscopically **FORBIDDEN** and cannot be observed
- **HOWEVER**, they are **ALLOWED** for relaxation

Let's suppose  $W_0$  is more efficient than  $W_2$  (large molecule case):

-in this case, the system will attempt to return to equilibrium by moving spins from N3 to N2

after saturation  
and  $W_0$  relaxation

N4 =  $N + \Delta H/2$   
 N3 =  $N + 3\Delta H/2 - \delta$   
 N2 =  $N + \Delta H/2 + \delta$   
 N1 =  $N + 3\Delta H/2$



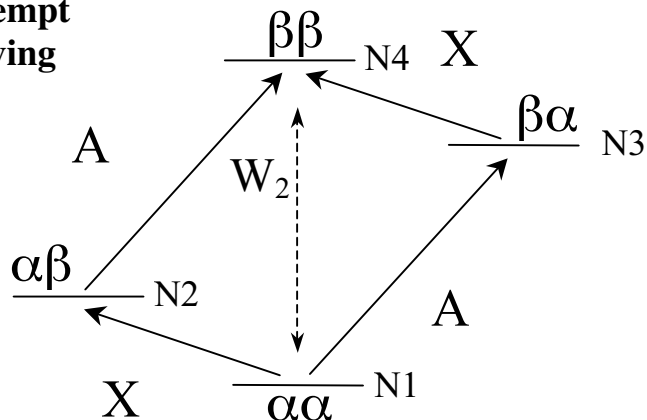
'X' transitions:  $N1 - N2 = N3 - N4 = \Delta H - \delta$ : intensity of 'X' reduced !!



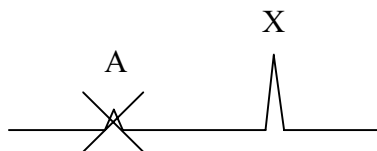
Let's suppose  $W_2$  is more efficient than  $W_0$  (small molecule case):

-in this case, the system will attempt to return to equilibrium by moving spins from N4 to N1

after saturation  
and  $W_0$  relaxation  
 $N4 = N + \Delta H/2 - \delta$   
 $N3 = N + 3\Delta H/2$   
 $N2 = N + \Delta H/2$   
 $N1 = N + 3\Delta H/2 + \delta$



'X' transitions:  $N1 - N2 = N3 - N4 = \Delta H + \delta$ : intensity of 'X' increased !!



**Summary:**

(reminder:  $\tau_c$  is the correlation time of the molecule (sec))

- for small molecules (short  $\tau_c$ )  $W_2$  dominates
- for large molecules (large  $\tau_c$ )  $W_0$  dominates
- because  $W_0$  and  $W_2$  depend on  $B_0$ , the NOE also depends on  $B_0$

## 2. Enhancement Factor for Small Molecules:

Maximal enhancement for small molecules ( $\eta$ ):

$$\eta = \gamma_a / 2\gamma_x \quad \text{A: saturated} \quad \text{X: observed}$$

( $\gamma$  is gyromagnetic ratio, of course)

Total Intensity (I):

$$I = (1 + \eta) I_0 \quad I_0 \equiv \text{original intensity}$$

Examples:

- a. Fractional increase in  $^{13}\text{C}$  NMR signal intensity with BB  $^1\text{H}$  decoupling:

$$\eta = \gamma_a / 2\gamma_x$$
$$\eta = 4 / (2 \times 1) = 2 \quad \text{i.e. signal increased by a factor of 2 (200\%)}$$

- b. Fractional increase in  $^1\text{H}$  NMR signal intensity with BB  $^{13}\text{C}$  decoupling:

$$\eta = 1 / 8 = 0.125 \quad \text{i.e. signal increased by a maximum of 12.5\%}$$
$$I = (1 + 0.125) I_0 = 1.125 \times I_0 \quad (\text{signal is slightly more intense})$$

## 3. Multi-Spin Systems:

Consider a homonuclear 3-spin system A-B-C

- assume no  $J$  coupling
- assume the molecule is a small molecule

- If A is saturated, the intensity of the signal from B is increased
- The saturation of A affects the populations of the B-C system, which causes an indirect NOE to C.
- Thus, this indirect NOE is negative

4. Dependence of Maximum  $^1\text{H}$ - $^1\text{H}$  Homonuclear NOE on  $B_0$  and molecular mass:

The magnitude and sign of the  $^1\text{H}$ - $^1\text{H}$  NOE are dependent on:

- the Larmor frequency,  $\nu$  (sometimes  $\omega$ ), which depends on  $B_0$
- the correlation time,  $\tau_c$ , which depends on molecular size (mass)
  - VERY crudely,  $\tau_c \approx 10^{-12} \text{ MW}$  (where MW is molecular mass in Daltons)
  - example:  $\tau_c \approx 10 \text{ ns}$  for 10 kDa molecule

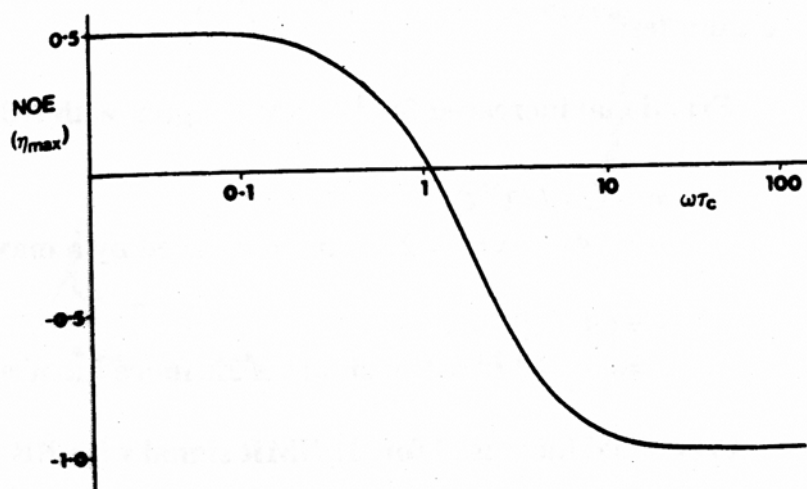


FIGURE 2.6  
Dependence of maximum homonuclear NOE enhancement on  $\omega\tau_c$ . Note the log scale of  $\omega\tau_c$ .

Note:

- NOE is *positive* for *small* molecules
- NOE is *negative* for *large* molecules
- NOE is zero at  $\omega\tau_c \approx 1$ 
  - examples: -at 200 MHz, NOE is zero for molecules with  $\tau_c \approx 5 \text{ ns}$  (approximately 5 kDa)
  - at 800 MHz, NOE is zero for molecules with  $\tau_c \approx 1.25 \text{ ns}$  (approximately 1.25 kDa)



## **5. Dependence of the magnitude of the NOE on Internuclear Distance:**

**In the absence of external relaxation:**

- the maximal enhancement is achieved
- the NOE is distance independent

**In the presence of external relaxation:**

- two-spin steady-state NOE depends on internuclear distance

$$f = \text{NOE enhancement} \propto 1/r^6$$

(r = internuclear distance)

- this effect is usually observed only for  $^1\text{H}$ - $^1\text{H}$  distances  $\leq 5 \text{ \AA}$

**Examples of external relaxation:**

**a). intermolecular relaxation caused by paramagnetic agents**

- field at a nucleus depends on the presence of unpaired electrons
- examples: dissolved  $\text{O}_2$ , manganese and other metal impurities

**b). quadrupolar relaxation**

- energy of quadrupole nuclei varies with the fluctuating electric field gradient
- dominant mechanism for nuclei with  $I > 1/2$  ( $^2\text{H}$  and  $^{14}\text{N}$ )

**c). chemical shift anisotropy (CSA)**

- anisotropy produces a different field at the nucleus depending on the orientation of the bond relative to the applied field
- significant for nuclei with large CSA ( $^{19}\text{F}$ ,  $^{31}\text{P}$ )