

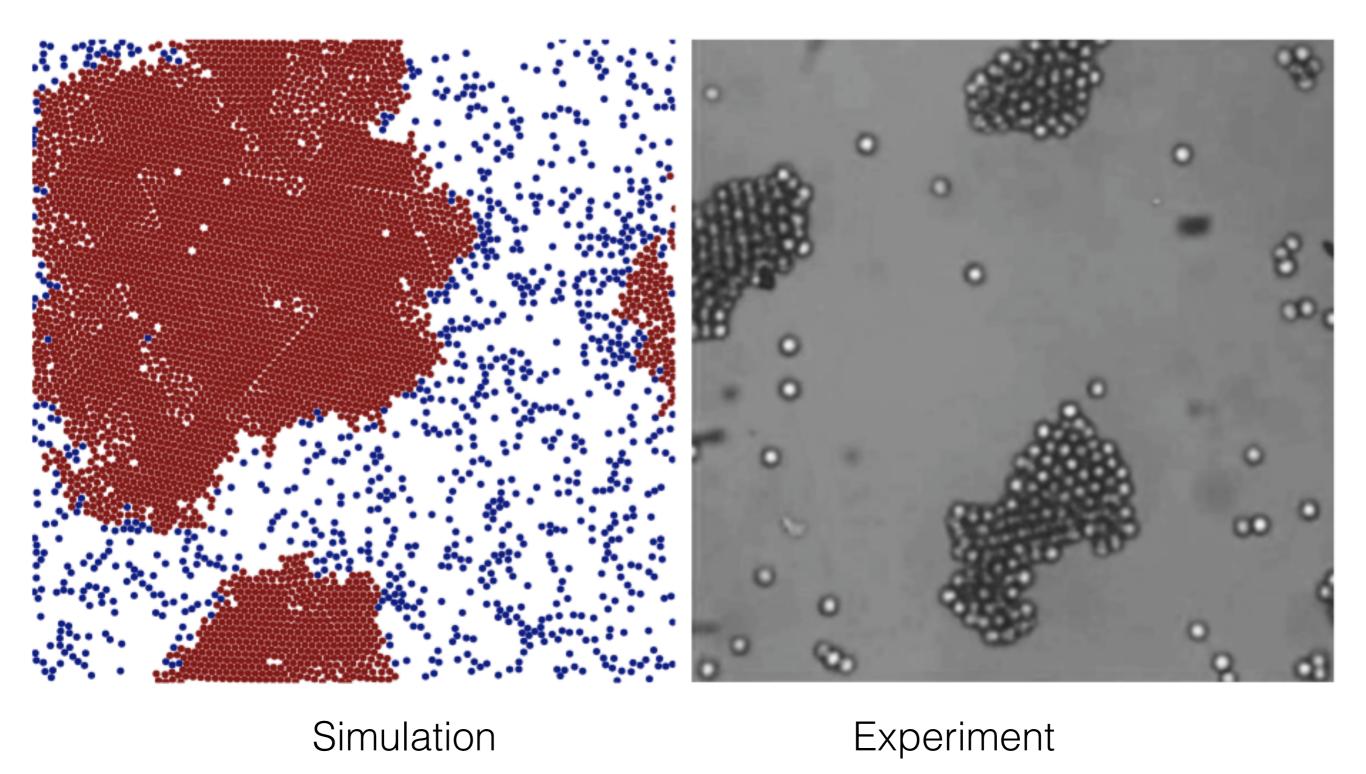
Microscopic Degrees of Freedom

Particle Positions \mathbf{r}_k Particle Directors \mathbf{e}_k

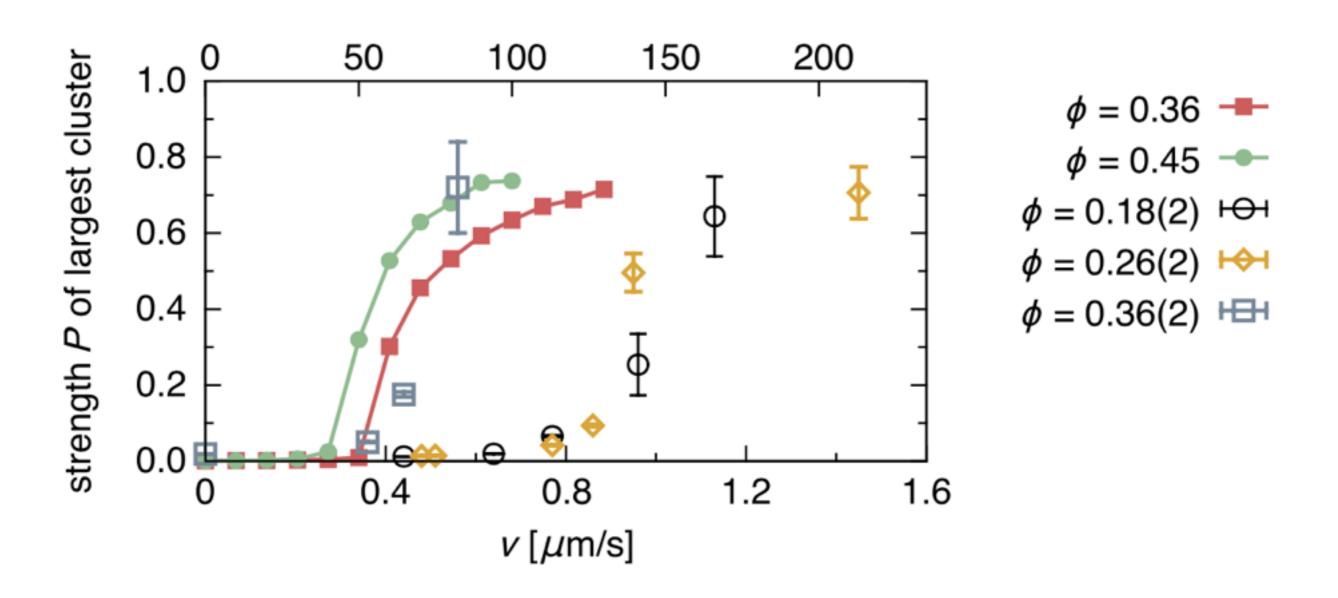
Microscopic Equations of Motion

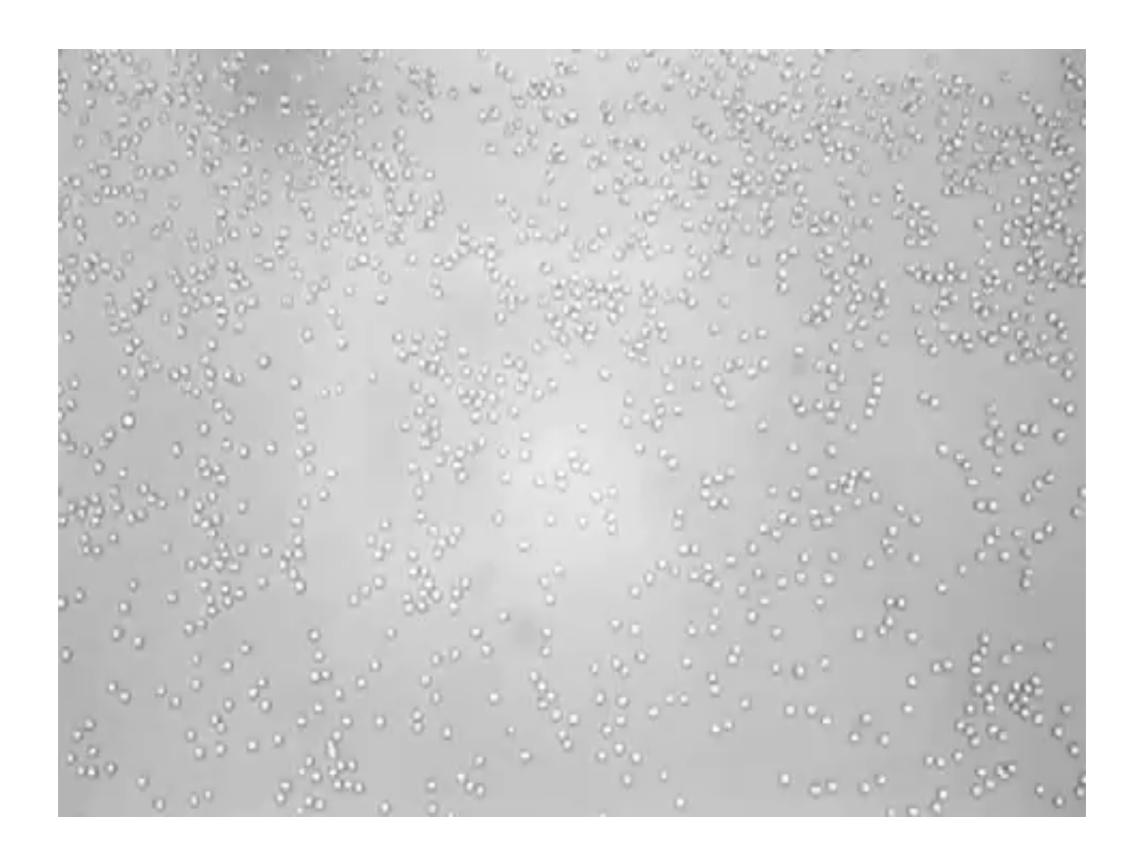
$$\dot{\mathbf{r}}_k = -\nabla_k U + V \mathbf{e}_k + f_k$$

$$\dot{\mathbf{e}}_k = g_k$$



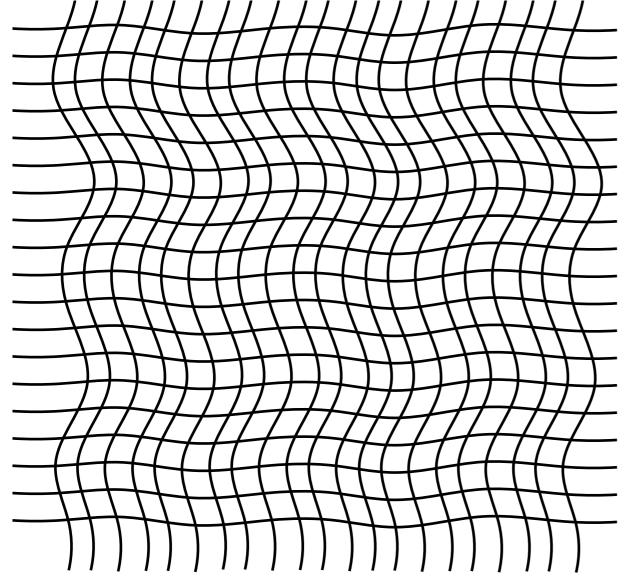
Buttinoni et al. PRL 2013





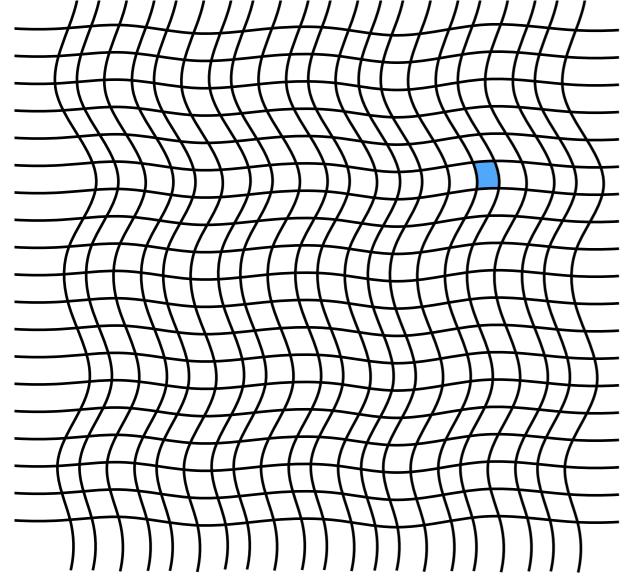
Continuum Mechanics

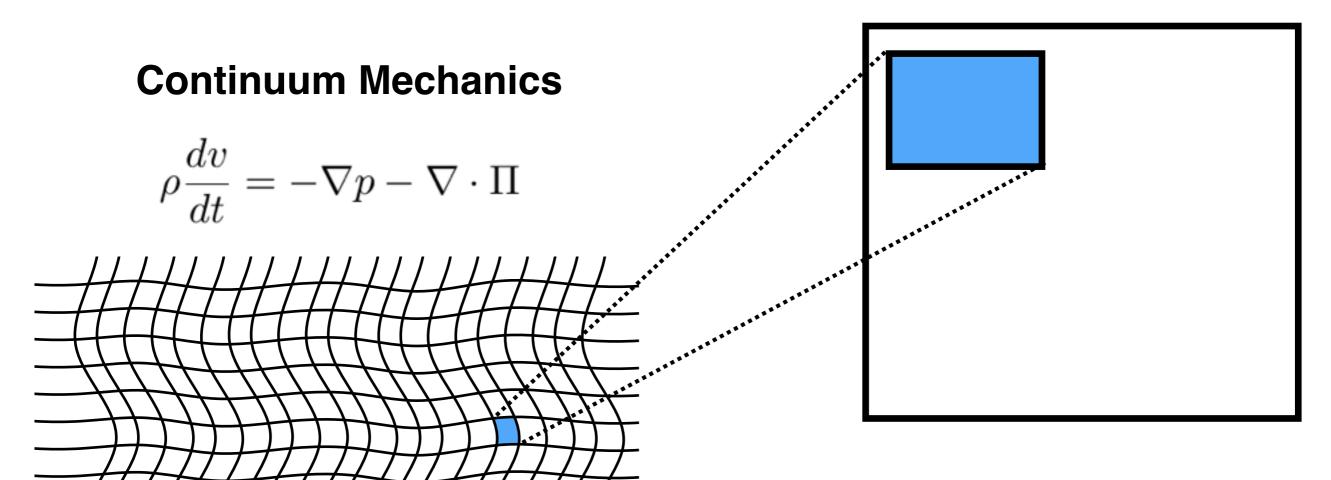
$$\rho \frac{dv}{dt} = -\nabla p - \nabla \cdot \Pi$$



Continuum Mechanics

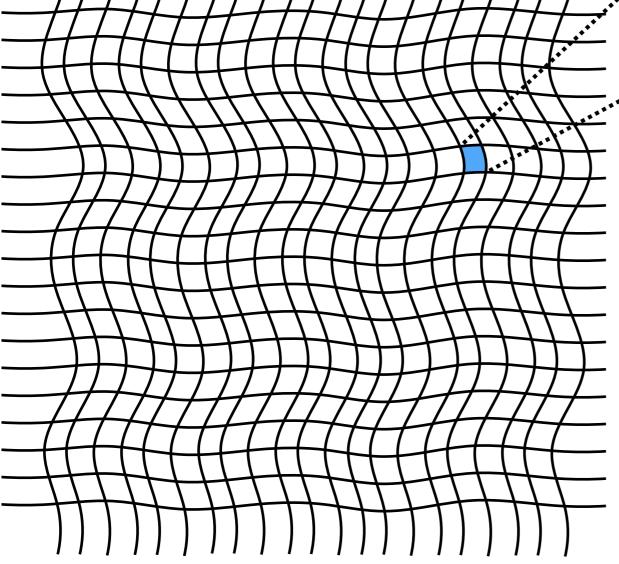
$$\rho \frac{dv}{dt} = -\nabla p - \nabla \cdot \Pi$$





Rational Thermodynamics, C. Truesdell Non-Equilibrium Thermodynamics, de Groot, Mazur





work rate W(t)

Work

heat rate $\mathcal{D}(t)$

internal energy $\mathcal{F}(t)$

heat bound $\mathcal{B}(t)$

Thermodynamics

Rational Thermodynamics, C. Truesdell Non-Equilibrium Thermodynamics, de Groot, Mazur

Heat

Heat is absorbed through boundaries

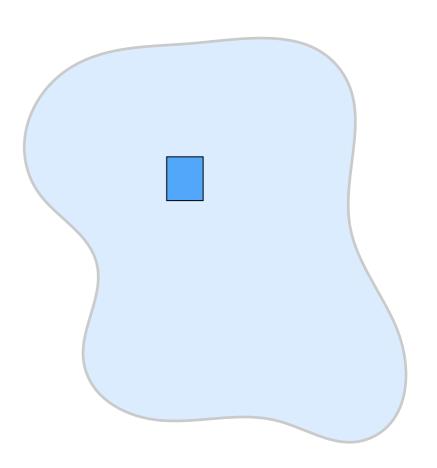
$$\mathcal{D}(t) = -\rho^{-1} \nabla \cdot J_q$$

Entropy has both a current and a source

$$\mathcal{B}(t) = \rho^{-1}T\left(-\nabla \cdot J_s + \sigma\right)$$

Work is defined within the mechanical theory

$$\mathcal{W}(t) = -p\frac{dJ}{dt} - \rho^{-1}\Pi : \nabla v$$



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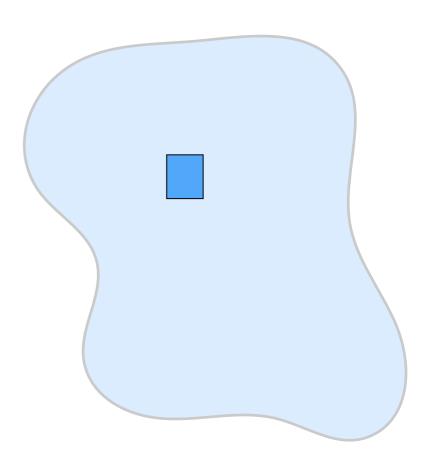
$$\mathcal{W}(t) = -p\frac{dJ}{dt} - \rho^{-1}\Pi : \nabla v$$

Assume Local equilibrium $\mathcal{B}(t) = \dot{\mathcal{F}}(t) + p\dot{J}$

$$\mathcal{B}(t) = \dot{\mathcal{F}}(t) + pJ$$

Assume entropy production takes the form

$$\sigma = \sum_{r} J_r :: X_r \qquad J_r = \lim(\{X_q\})$$



After some algebra...

$$\rho \frac{dv}{dt} = -\nabla p + \eta \triangle v + \left(\frac{1}{3}\eta + \eta_v\right) \nabla \left(\nabla \cdot v\right)$$

This is Navier-Stokes!

It's simple to add in temperature gradients and multiple reacting species.

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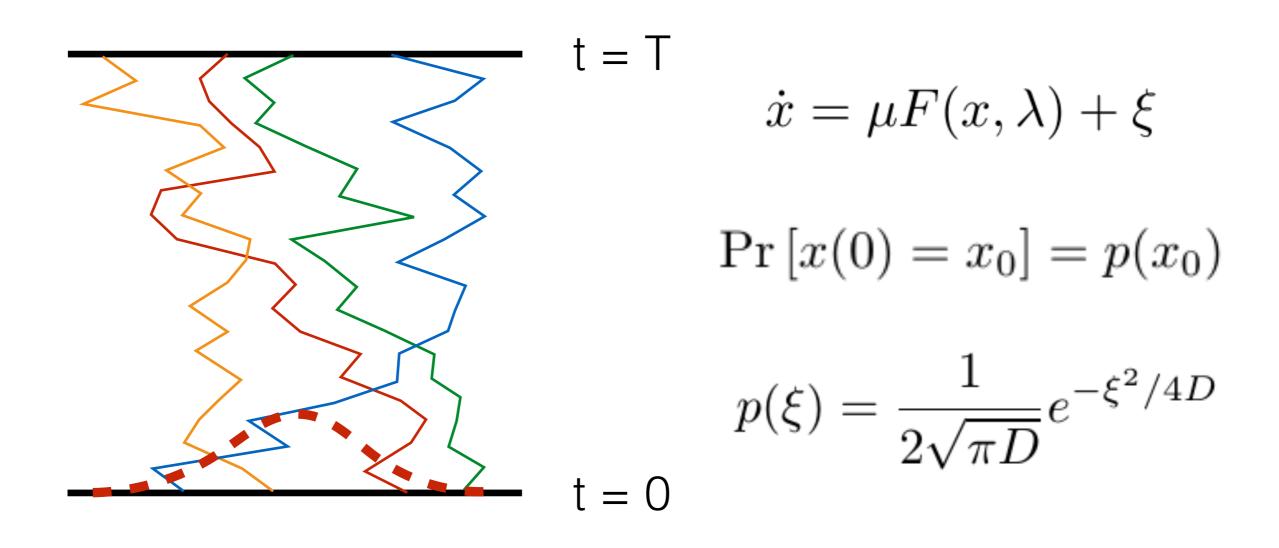
It's simple to add in temperature gradients and multiple reacting species.

Local Equilibrium

→ LIT
Form of EP

→ Hydrodynamics

Goal: Associate thermodynamical quantities to individual trajectories.



Can we quantify the "irreversibility" of a process?

Does a reversed trajectory generated by a reversed protocol look like a forwards trajectory?

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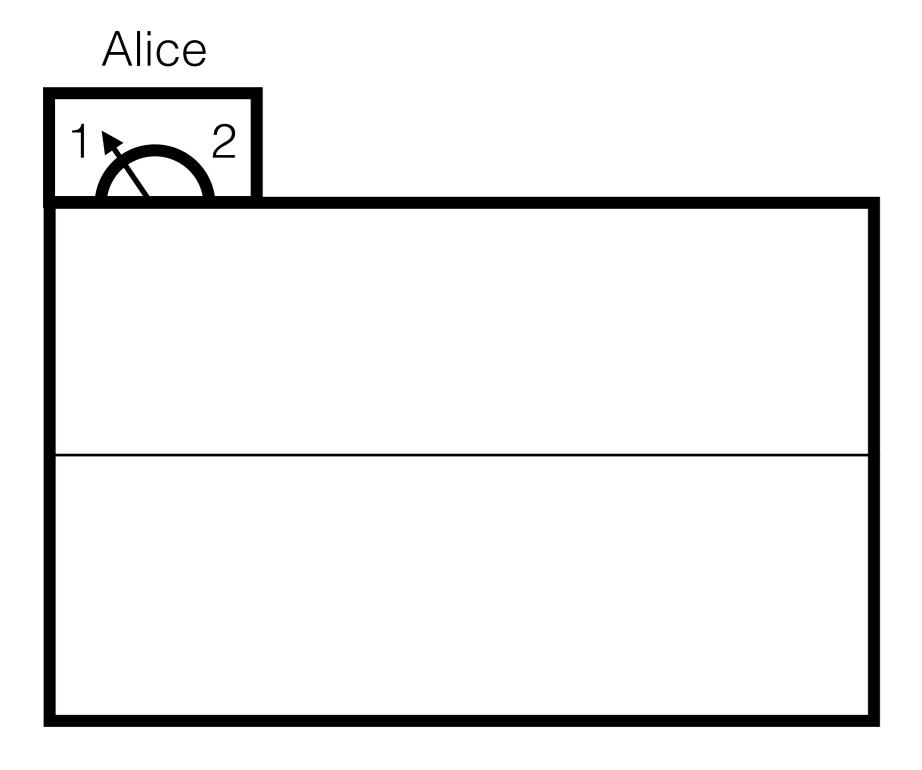
Define a protocol F(x,t) and two probability distributions p and q on \mathbb{R} . Consider the following two samplers from the set of paths $x:[0,T]\to\mathbb{R}$:

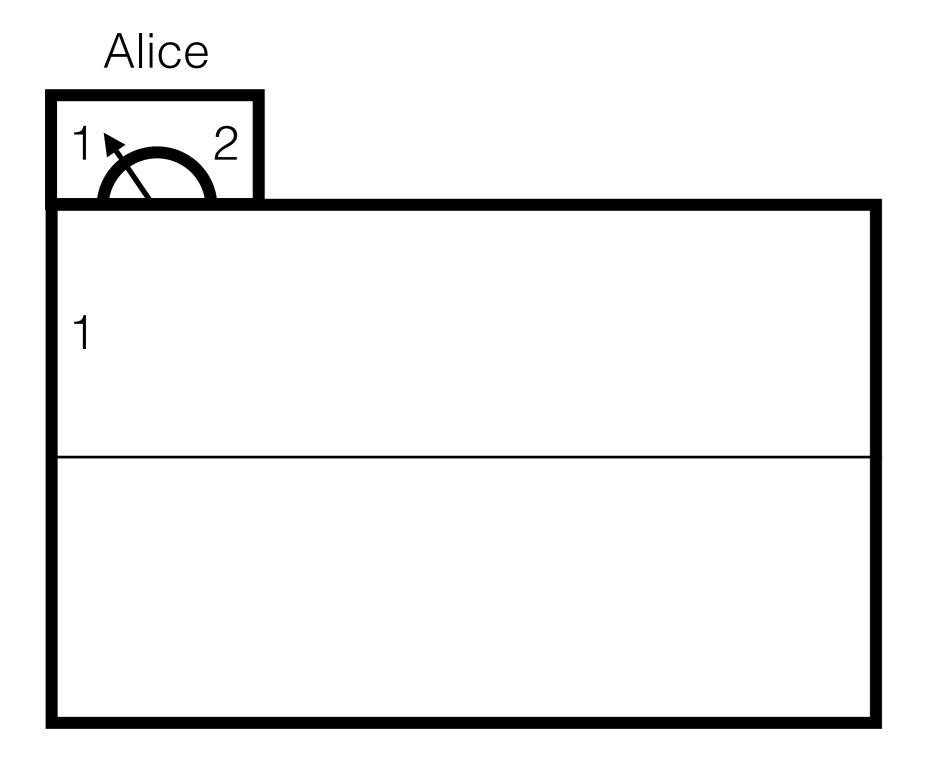
Sampler 1

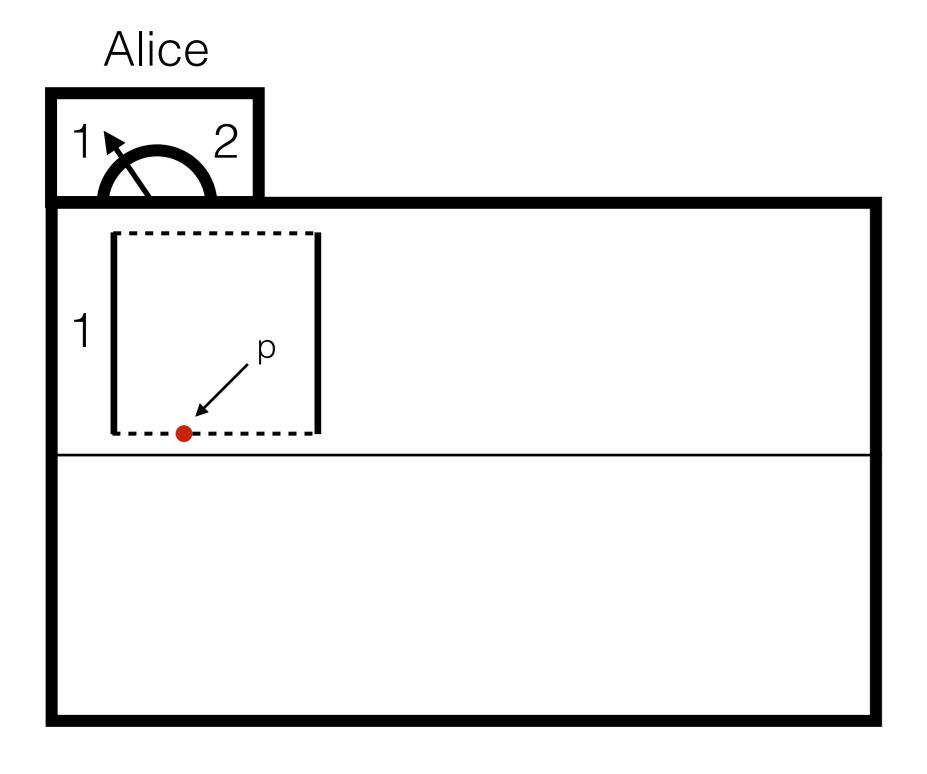
- 1. Generate a point x_0 from the distribution p.
- 2. Generate a path $x:[0,T]\to\mathbb{R}$ with SDE $\dot{x}(t)=\mu F(x(t),t)+\xi(t)$ with initial condition $x(0)=x_0$.
- 3. Return x.

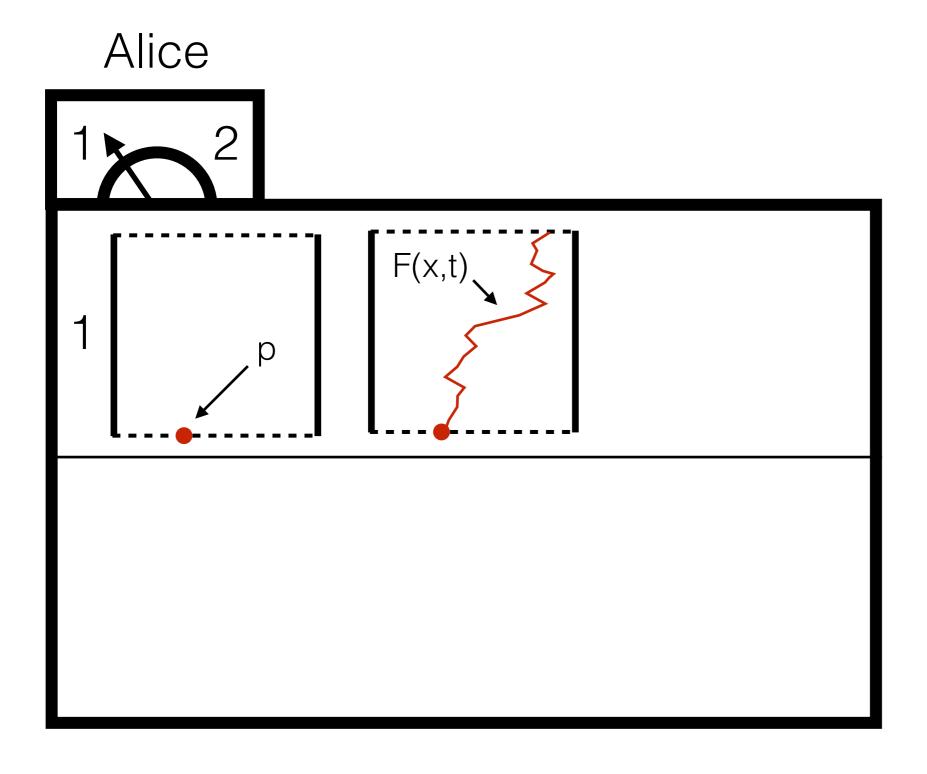
Sampler 2

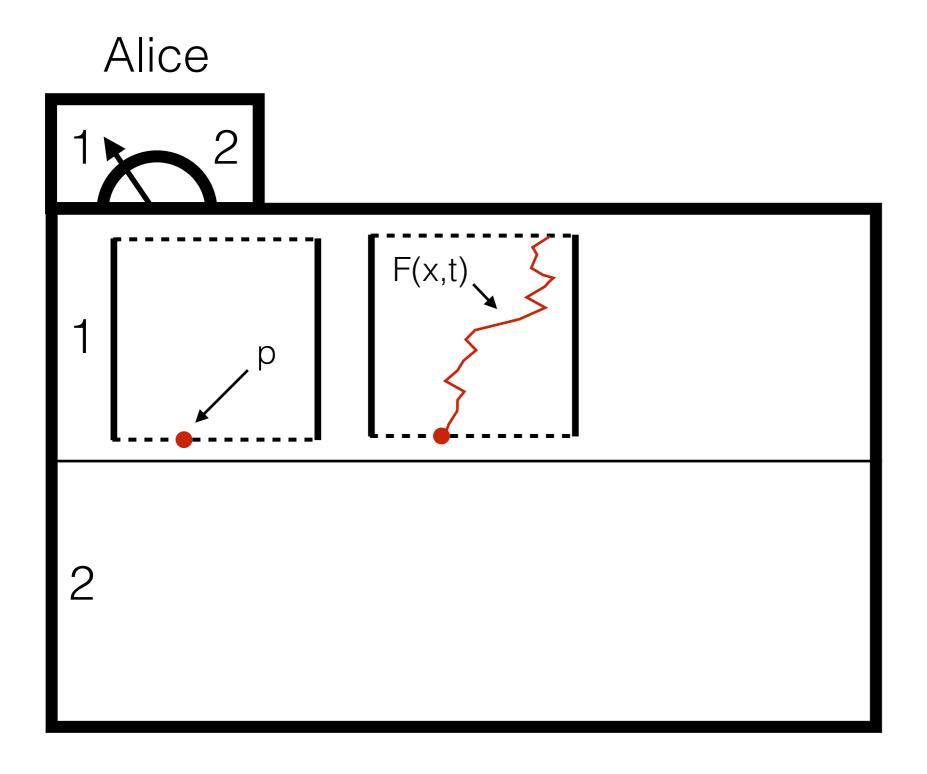
- 1. Generate a point y_0 from the distribution q.
- 2. Generate path $y:[0,T] \to \mathbb{R}$ with SDE $\dot{y}(t) = \mu F(y(t),T-t) + \xi(t)$ with initial condition $y(0) = y_0$.
- 3. Define x(t) = y(T t).
- 4. Return x.

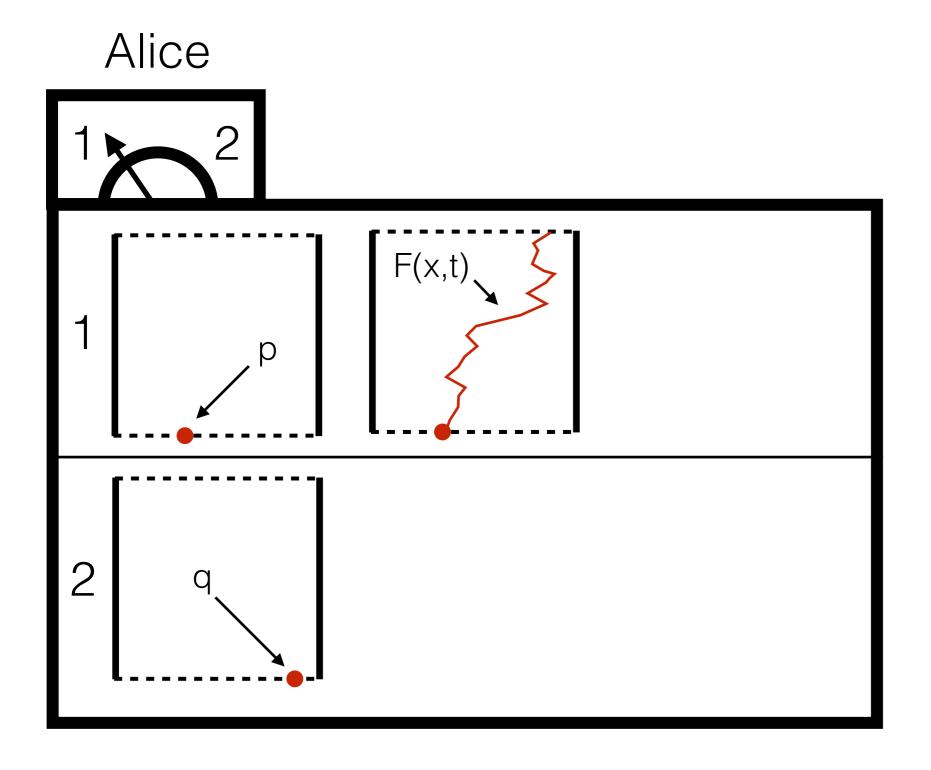


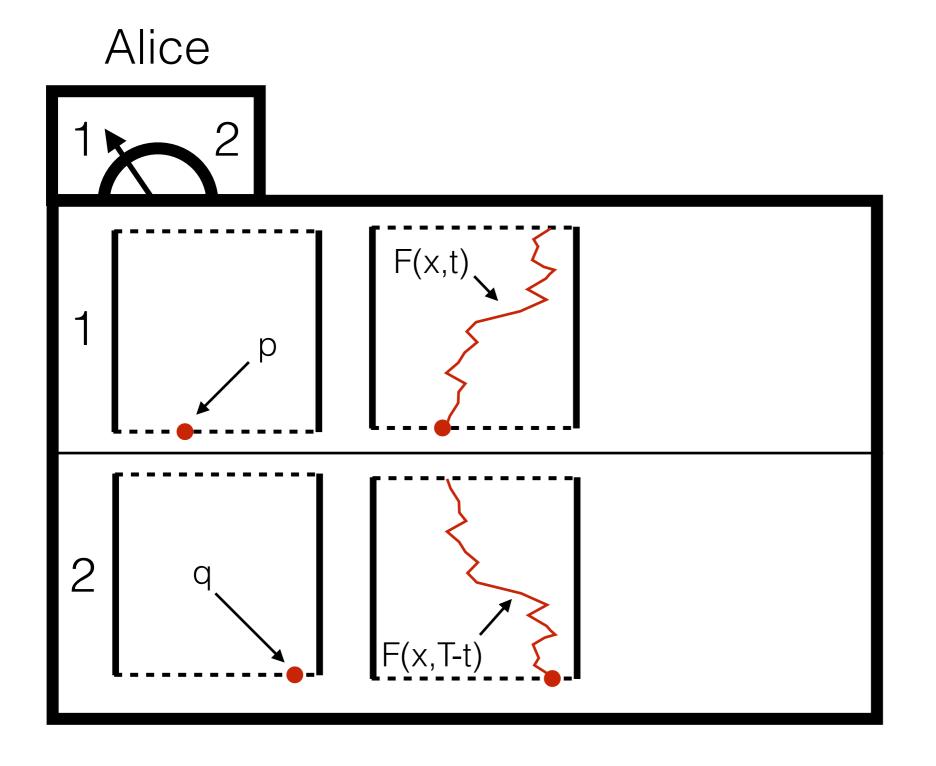


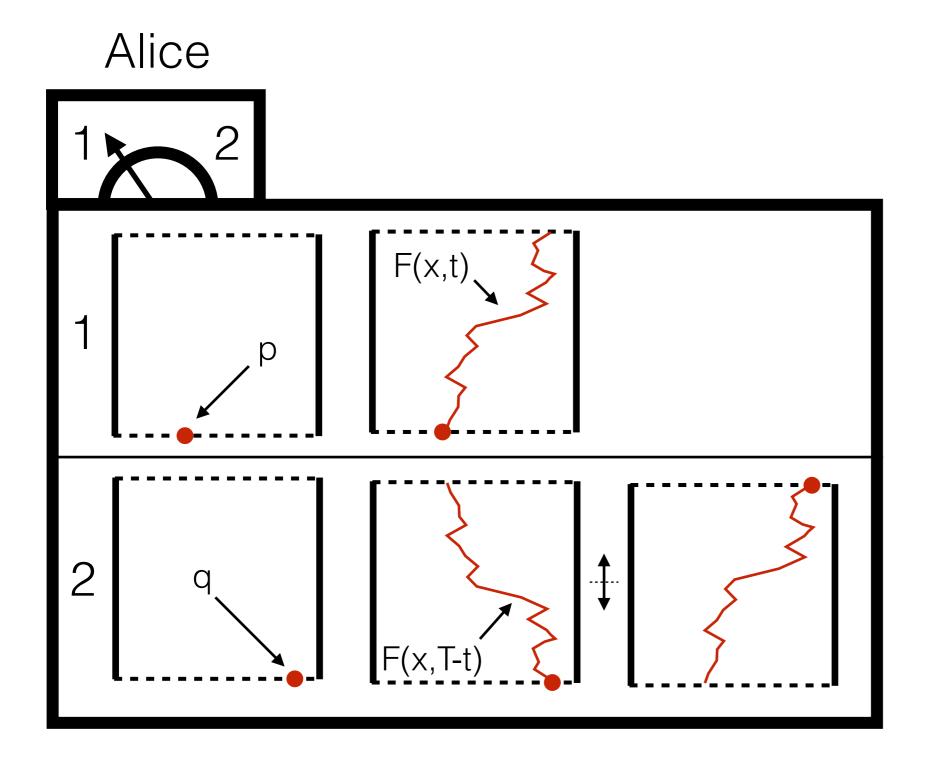


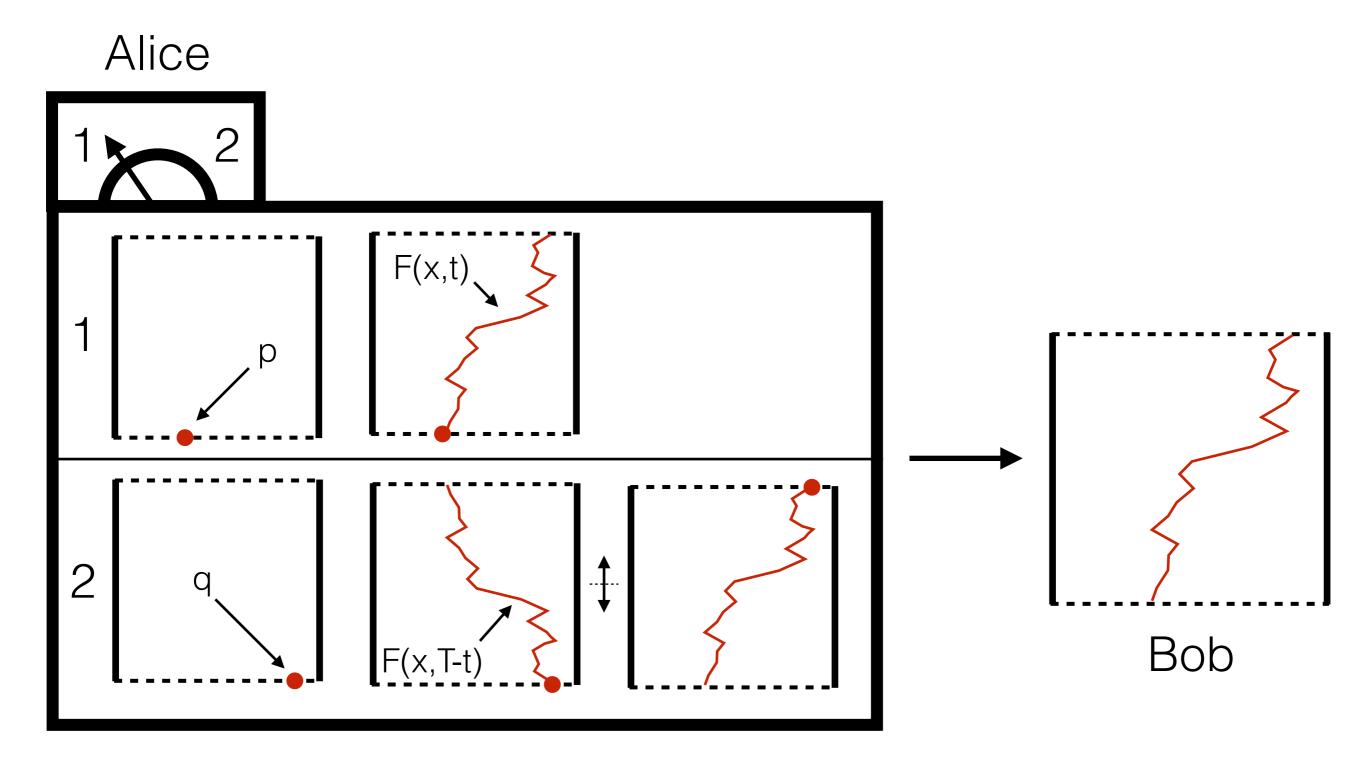


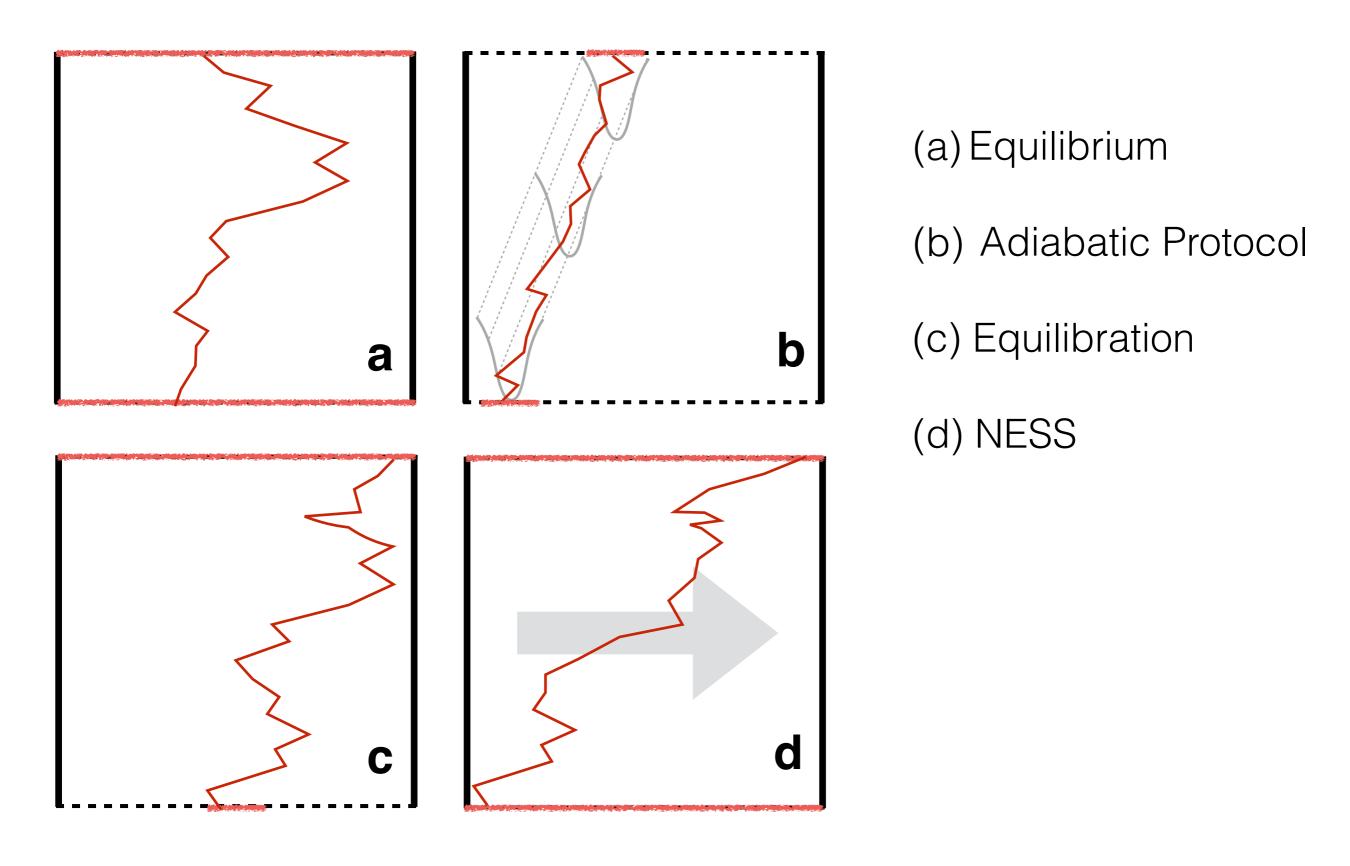


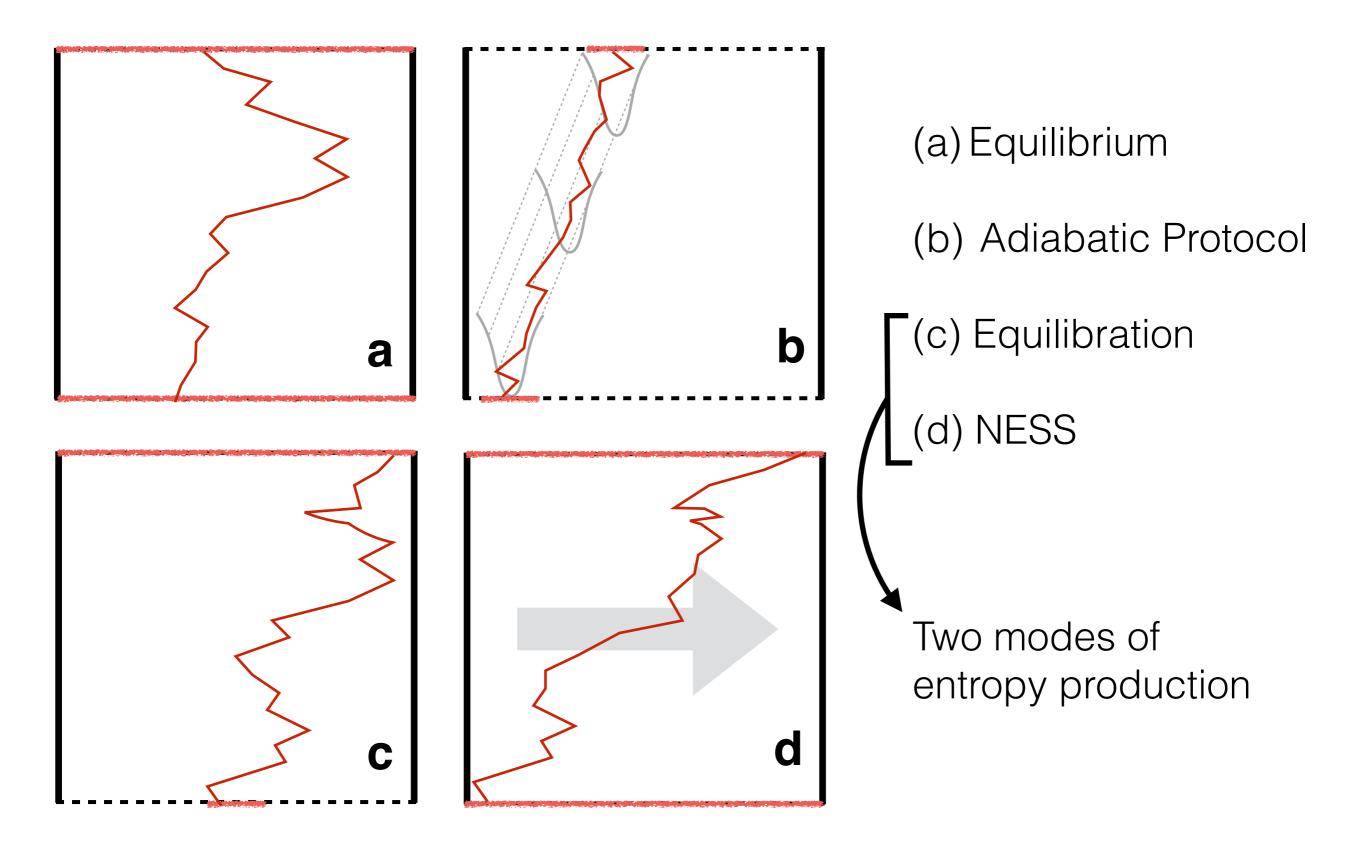












$$\Pr\left[\xi_{x}|1\right] \sim p(x(0)) \cdot \exp\left[-\int_{0}^{T} \frac{\xi_{x}(t)^{2}}{4D} dt\right] = p(x(0)) \cdot \exp\left[-\int_{0}^{T} \frac{\left(\dot{x}(t) - \mu F(x(t), t)\right)^{2}}{4D} dt\right]$$

$$\Pr[\xi_x|2] \sim q(x(t)) \cdot \exp\left[-\int_0^T \frac{(-\dot{x}(T-t) - \mu F(x(T-t), T-t)^2}{4D}dt\right]$$

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Work given to disordered degrees of freedom = heat

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$$\Pr\left[\xi_{x}|2\right] \sim q(x(t)) \cdot \exp\left[-\int_{0}^{T} \frac{\left(-\dot{x}(T-t) - \mu F(x(T-t), T-t)^{2}}{4D} dt\right]$$

$$\ln\left[\frac{\Pr\left[\xi_{x}|1\right]}{\Pr\left[\xi_{x}|1\right]}\right] = \ln\left[\frac{p(x(0))}{q(x(t))}\right] + \frac{\mu}{D} \int_{0}^{T} F(x(t),t)\dot{x}(t)dt$$

Work given to disordered degrees of freedom = heat In steady-state with bath at T:

$$\Delta S = \frac{Q}{T}$$

To Do:

Actually compute entropy production for active matter systems and use a local steady-state hypothesis to derive hydrodynamics.

Understand how microscopic irreversibility gets reflected at the macroscale.

