

1 Divergence and Vorticity?

The divergence and vorticity of the perturbation velocity in cylindrical coordinates is:

$$\begin{aligned}
\vec{\nabla} \cdot \vec{v}' &= \frac{1}{r} \frac{\partial}{\partial r} (r v'_r) + \frac{1}{r} \frac{\partial v'_\theta}{\partial \theta} + \frac{\partial v'_x}{\partial x} \\
&= \frac{1}{r_T} \left(\frac{1}{\tilde{r}} \frac{\partial}{\partial \tilde{r}} (\tilde{r} v'_r) + \frac{1}{\tilde{r}} \frac{\partial v'_\theta}{\partial \theta} + \frac{\partial v'_x}{\partial \tilde{x}} \right) \\
\vec{\nabla} \times \vec{v}' &= \begin{pmatrix} \left(\frac{1}{r} \frac{\partial v'_x}{\partial \theta} - \frac{\partial v'_\theta}{\partial x} \right) \vec{e}_r \\ + \left(\frac{\partial v'_r}{\partial x} - \frac{\partial v'_x}{\partial r} \right) \vec{e}_\theta \\ + \frac{1}{r} \left(\frac{\partial}{\partial r} (r v'_\theta) - \frac{\partial v'_r}{\partial \theta} \right) \vec{e}_x \end{pmatrix} \\
&= \frac{1}{r_T} \begin{pmatrix} \left(\frac{1}{\tilde{r}} \frac{\partial v'_x}{\partial \theta} - \frac{\partial v'_\theta}{\partial \tilde{x}} \right) \vec{e}_r \\ + \left(\frac{\partial v'_r}{\partial \tilde{x}} - \frac{\partial v'_x}{\partial \tilde{r}} \right) \vec{e}_\theta \\ + \frac{1}{\tilde{r}} \left(\frac{\partial}{\partial \tilde{r}} (\tilde{r} v'_\theta) - \frac{\partial v'_r}{\partial \theta} \right) \vec{e}_x \end{pmatrix}
\end{aligned}$$

Remembering the definitions:

$$\begin{aligned}
v'_r &= v_r(r) e^{i(k_x x + m\theta - \omega t)} \\
v'_\theta &= v_\theta(r) e^{i(k_x x + m\theta - \omega t)} \\
v'_x &= v_x(r) e^{i(k_x x + m\theta - \omega t)} \\
p' &= p(r) e^{i(k_x x + m\theta - \omega t)}
\end{aligned}$$

with the nondimensional counterparts:

$$\begin{aligned}
v'_r &= A \tilde{v}_r(\tilde{r}) e^{i(\bar{\gamma} \tilde{x} + m\theta - \omega t)} \\
v'_\theta &= A \tilde{v}_\theta(\tilde{r}) e^{i(\bar{\gamma} \tilde{x} + m\theta - \omega t)} \\
v'_x &= A \tilde{v}_x(\tilde{r}) e^{i(\bar{\gamma} \tilde{x} + m\theta - \omega t)} \\
p' &= \bar{\rho} A^2 \tilde{p}(\tilde{r}) e^{i(\bar{\gamma} \tilde{x} + m\theta - \omega t)}
\end{aligned}$$

where:

$$\bar{\gamma} = k_x r_T$$

The axial derivatives are:

$$\begin{aligned}
\frac{\partial v'_r}{\partial x} &= \frac{1}{r_T} \frac{\partial v'_r}{\partial \tilde{x}} \\
&= i\tilde{\gamma}\tilde{v}_r \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right) \\
\frac{\partial v'_\theta}{\partial x} &= i\tilde{\gamma}\tilde{v}_\theta \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right) \\
\frac{\partial v'_x}{\partial x} &= i\tilde{\gamma}\tilde{v}_x \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right)
\end{aligned}$$

The azimuthal derivatives are:

$$\begin{aligned}
\frac{1}{r} \frac{\partial v'_r}{\partial \theta} &= \frac{im\tilde{v}_r}{\tilde{r}} \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right) \\
\frac{1}{r} \frac{\partial v'_\theta}{\partial \theta} &= \frac{im\tilde{v}_\theta}{\tilde{r}} \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right) \\
\frac{1}{r} \frac{\partial v'_x}{\partial \theta} &= \frac{im\tilde{v}_x}{\tilde{r}} \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right)
\end{aligned}$$

To do the radial derivatives, note that:

$$\frac{\partial A}{\partial r} = \frac{A}{r_T} \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}}$$

The radial derivatives (remembering that $A = A(r)$!) are:

$$\begin{aligned}
\frac{\partial v'_r}{\partial r} &= \left(\tilde{v}'_r \frac{\partial A}{\partial r} + \frac{A}{r_T} \frac{\partial v'_r}{\partial \tilde{r}} \right) e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \\
&= \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_r + \frac{\partial \tilde{v}_r}{\partial \tilde{r}} \right) \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right) \\
\frac{\partial v'_\theta}{\partial r} &= \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_\theta + \frac{\partial \tilde{v}_\theta}{\partial \tilde{r}} \right) \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right) \\
\frac{\partial v'_x}{\partial r} &= \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_x + \frac{\partial \tilde{v}_x}{\partial \tilde{r}} \right) \left(\frac{A}{r_T} e^{i(\tilde{\gamma}\tilde{x}+m\theta-\omega t)} \right)
\end{aligned}$$

gives:

$$\begin{aligned}\vec{\nabla} \cdot \vec{v}' &= \left(\frac{\tilde{v}_r}{\tilde{r}} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2 \tilde{v}_r}{\tilde{r}} + \frac{\partial \tilde{v}_r}{\partial \tilde{r}} + \frac{im}{\tilde{r}} \tilde{v}_\theta + i\gamma \tilde{v}_x \right) \left(\frac{A}{r_T} e^{i(\gamma \tilde{x} + m\theta - \omega t)} \right) \\ \vec{\nabla} \times \vec{v}' &= \begin{pmatrix} \left(\frac{im\tilde{v}_x}{\tilde{r}} - i\gamma \tilde{v}_\theta \right) \vec{e}_r \\ + \left(i\gamma \tilde{v}_r - \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_x + \frac{\partial \tilde{v}_x}{\partial r} \right) \right) \vec{e}_\theta \\ + \left(\frac{\tilde{v}_\theta}{\tilde{r}} + \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_\theta + \frac{\partial \tilde{v}_\theta}{\partial r} \right) - \frac{im\tilde{v}_r}{\tilde{r}} \right) \vec{e}_x \end{pmatrix} \left(\frac{A}{r_T} e^{i(\gamma \tilde{x} + m\theta - \omega t)} \right)\end{aligned}$$

If all three vorticity components are zero, the perturbations will be irrotational. From the first two velocity components, this gives:

$$\begin{aligned}\tilde{v}_\theta &= \frac{m}{\gamma \tilde{r}} \tilde{v}_x \\ \tilde{v}_r &= \frac{-i}{\gamma} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_x + \frac{\partial \tilde{v}_x}{\partial \tilde{r}} \right)\end{aligned}$$

As a test, the last vorticity component is:

$$\begin{aligned}\frac{\tilde{v}_\theta}{\tilde{r}} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_\theta + \frac{\partial \tilde{v}_\theta}{\partial \tilde{r}} - \frac{im}{\tilde{r}} \tilde{v}_r &= \begin{pmatrix} \frac{m}{\gamma \tilde{r}} \left(\frac{1}{\tilde{r}} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \right) \tilde{v}_x \\ + \frac{\partial}{\partial \tilde{r}} \left(\frac{m}{\gamma \tilde{r}} \tilde{v}_x \right) \\ - \frac{m}{\gamma \tilde{r}} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_x + \frac{\partial \tilde{v}_x}{\partial \tilde{r}} \right) \end{pmatrix} \\ &= \begin{pmatrix} \frac{m}{\gamma \tilde{r}} \left(\frac{1}{\tilde{r}} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \right) \tilde{v}_x \\ - \frac{m}{\gamma \tilde{r}^2} \tilde{v}_x \\ + \frac{m}{\gamma \tilde{r}} \frac{\partial \tilde{v}_x}{\partial \tilde{r}} \\ - \frac{m}{\gamma \tilde{r}} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{\tilde{r}} \tilde{v}_x + \frac{\partial \tilde{v}_x}{\partial \tilde{r}} \right) \end{pmatrix} \\ &= 0\end{aligned}$$

showing that these are the irrotational relations between the perturbation velocities.

1.1 Velocity decomposition

Let's decompose the perturbation velocity field into irrotational and divergence-free components:

$$\begin{aligned}\tilde{v}_x &= \tilde{v}_{x,\omega} + \tilde{v}_{x,\phi} \\ \tilde{v}_r &= \tilde{v}_{r,\omega} + \tilde{v}_{r,\phi} \\ \tilde{v}_\theta &= \tilde{v}_{\theta,\omega} + \tilde{v}_{\theta,\phi}\end{aligned}$$

with the relations:

$$\begin{aligned}\tilde{v}_{\theta,\phi} &= \frac{m}{\gamma\tilde{r}}\tilde{v}_{x,\phi} \\ \tilde{v}_{r,\phi} &= \frac{-i}{\gamma}\left(\left(\frac{\gamma-1}{2}\right)\frac{M_\theta^2}{\tilde{r}}\tilde{v}_{x,\phi} + \frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}}\right)\end{aligned}$$

The velocity divergence gives one relation for the rotational components of the perturbation velocities:

$$\begin{aligned}\tilde{v}_{x,\omega} &= \frac{i}{\gamma}\left(\frac{\tilde{v}_{r,\omega}}{\tilde{r}} + \left(\frac{\gamma-1}{2}\right)\frac{M_\theta^2\tilde{v}_{r,\omega}}{\tilde{r}} + \frac{\partial\tilde{v}_{r,\omega}}{\partial\tilde{r}}\right) - \frac{m}{\gamma\tilde{r}}\tilde{v}_{\theta,\omega} \\ &= i\left(\frac{1+\Gamma}{\gamma\tilde{r}}\right)\tilde{v}_{r,\omega} + \frac{i}{\gamma}\frac{\partial\tilde{v}_{r,\omega}}{\partial\tilde{r}} - \frac{m}{\gamma\tilde{r}}\tilde{v}_{\theta,\omega}\end{aligned}$$

As a check, the velocity perturbations are put into the velocity divergence equation:

$$\begin{aligned}\vec{\nabla} \cdot \vec{v}' &= \left(\frac{\tilde{v}_r}{\tilde{r}} + \left(\frac{\gamma-1}{2}\right)\frac{M_\theta^2\tilde{v}_r}{\tilde{r}} + \frac{\partial\tilde{v}_r}{\partial\tilde{r}} + \frac{im}{\tilde{r}}\tilde{v}_\theta + i\gamma\tilde{v}_x\right) \\ &= \left(\frac{\tilde{v}_{r,\phi}}{\tilde{r}} + \left(\frac{\gamma-1}{2}\right)\frac{M_\theta^2\tilde{v}_{r,\phi}}{\tilde{r}} + \frac{\partial\tilde{v}_{r,\phi}}{\partial\tilde{r}} + \frac{im}{\tilde{r}}\tilde{v}_{\theta,\phi} + i\gamma\tilde{v}_{x,\phi}\right) \\ &\quad + \left(\frac{\tilde{v}_{r,\omega}}{\tilde{r}} + \left(\frac{\gamma-1}{2}\right)\frac{M_\theta^2\tilde{v}_{r,\omega}}{\tilde{r}} + \frac{\partial\tilde{v}_{r,\omega}}{\partial\tilde{r}} + \frac{im}{\tilde{r}}\tilde{v}_{\theta,\omega} + i\gamma\tilde{v}_{x,\omega}\right)\end{aligned}$$

$$\begin{aligned}
&= \left(\begin{aligned} &\frac{1}{r} \left(1 + \left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \frac{-i}{\bar{\gamma}} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{r} \tilde{v}_{x,\phi} + \frac{\partial \tilde{v}_{x,\phi}}{\partial r} \right) \\ &+ \frac{\partial}{\partial r} \left(\frac{-i}{\bar{\gamma}} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{r} \tilde{v}_{x,\phi} + \frac{\partial \tilde{v}_{x,\phi}}{\partial r} \right) \right) \\ &+ \frac{im^2}{\bar{\gamma} r^2} \tilde{v}_{x,\phi} \\ &+ i\bar{\gamma} \tilde{v}_{x,\phi} \\ &+ \frac{\tilde{v}_{r,\omega}}{r} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2 \tilde{v}_{r,\omega}}{r} + \frac{\partial \tilde{v}_{r,\omega}}{\partial r} + \frac{im}{r} \tilde{v}_{\theta,\omega} \\ &+ i\bar{\gamma} \left(\frac{i}{\bar{\gamma}} \left(\frac{\tilde{v}_{r,\omega}}{r} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2 \tilde{v}_{r,\omega}}{r} + \frac{\partial \tilde{v}_{r,\omega}}{\partial r} \right) - \frac{m}{\bar{\gamma} r} \tilde{v}_{\theta,\omega} \right) \end{aligned} \right) \\
&= \left(\begin{aligned} &\frac{-i}{\bar{\gamma} r^2} \left(1 + \left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \left(\left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \tilde{v}_{x,\phi} \\ &- \frac{i}{\bar{\gamma} r} \left(1 + \left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \left(\frac{\partial \tilde{v}_{x,\phi}}{\partial r} \right) \\ &- \frac{i}{\bar{\gamma}} \frac{\partial}{\partial r} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{r} \tilde{v}_{x,\phi} \right) \\ &- \frac{i}{\bar{\gamma}} \frac{\partial^2 \tilde{v}_{x,\phi}}{\partial r^2} \\ &+ i\bar{\gamma} \left(1 + \frac{m^2}{\bar{\gamma}^2 r^2} \right) \tilde{v}_{x,\phi} \\ &+ \frac{\tilde{v}_{r,\omega}}{r} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2 \tilde{v}_{r,\omega}}{r} + \frac{\partial \tilde{v}_{r,\omega}}{\partial r} + \frac{im}{r} \tilde{v}_{\theta,\omega} \\ &- \left(\frac{\tilde{v}_{r,\omega}}{r} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2 \tilde{v}_{r,\omega}}{r} + \frac{\partial \tilde{v}_{r,\omega}}{\partial r} \right) - \frac{im}{r} \tilde{v}_{\theta,\omega} \end{aligned} \right) \\
&= \left(\begin{aligned} &i\bar{\gamma} \left(\left(1 + \frac{m^2}{\bar{\gamma}^2 r^2} \right) - \frac{1}{\bar{\gamma}^2 r^2} \left(1 + \left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \left(\left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \right) \tilde{v}_{x,\phi} \\ &- \frac{i}{\bar{\gamma} r} \left(1 + \left(\frac{\gamma-1}{2} \right) M_\theta^2 \right) \left(\frac{\partial \tilde{v}_{x,\phi}}{\partial r} \right) \\ &+ \frac{i}{\bar{\gamma}} \left(\left(\frac{\gamma-1}{2} \right) \frac{M_\theta^2}{r^2} \tilde{v}_{x,\phi} \right) \\ &- \frac{i}{\bar{\gamma}} \frac{\partial M_\theta}{\partial r} \left(\left(\frac{\gamma-1}{2} \right) \frac{2M_\theta}{r} \tilde{v}_{x,\phi} \right) \\ &- \frac{i}{\bar{\gamma} r} \left(\left(\frac{\gamma-1}{2} \right) M_\theta^2 \frac{\partial \tilde{v}_{x,\phi}}{\partial r} \right) \\ &- \frac{i}{\bar{\gamma}} \frac{\partial^2 \tilde{v}_{x,\phi}}{\partial r^2} \end{aligned} \right) \\
&= \left(\begin{aligned} &i \left(\bar{\gamma} + \frac{1}{\bar{\gamma} r^2} \left(m^2 - \Gamma^2 - \tilde{r} \frac{\partial \Gamma}{\partial r} \right) \right) \tilde{v}_{x,\phi} \\ &- \frac{i}{\bar{\gamma} r} (1 + 2\Gamma) \frac{\partial \tilde{v}_{x,\phi}}{\partial r} \\ &- \frac{i}{\bar{\gamma}} \frac{\partial^2 \tilde{v}_{x,\phi}}{\partial r^2} \end{aligned} \right)
\end{aligned}$$

where

$$\Gamma = \left(\frac{\gamma-1}{2} \right) M_\theta^2$$

Note that the divergence of velocity can be written solely in terms of $\tilde{v}_{x,\phi}$:

$$\left(\frac{\tilde{v}_r}{\tilde{r}} + \left(\frac{\gamma - 1}{2} \right) \frac{M_\theta^2 \tilde{v}_r}{\tilde{r}} + \frac{\partial \tilde{v}_r}{\partial \tilde{r}} + \frac{im}{\tilde{r}} \tilde{v}_\theta + i\gamma \tilde{v}_x \right) = \Phi_0 \tilde{v}_{x,\phi} + \Phi_1 \frac{\partial \tilde{v}_{x,\phi}}{\partial \tilde{r}} + \Phi_2 \frac{\partial^2 \tilde{v}_{x,\phi}}{\partial \tilde{r}^2}$$

where

$$\begin{aligned} \Phi_0 &= i \left(\bar{\gamma} + \frac{1}{\bar{\gamma} \tilde{r}^2} \left(m^2 - \Gamma^2 - \tilde{r} \frac{\partial \Gamma}{\partial \tilde{r}} \right) \right) \\ \Phi_1 &= -\frac{i}{\bar{\gamma} \tilde{r}} (1 + 2\Gamma) \\ \Phi_2 &= -\frac{i}{\bar{\gamma}} \end{aligned}$$

1.2 Back to business...

The SWIRL code equations are:

$$\begin{aligned} & -i \left(\frac{k}{\tilde{A}} - \frac{mM_\theta}{\tilde{r}} - \bar{\gamma}M_x \right) \tilde{v}_r - \frac{2M_\theta}{\tilde{r}} \tilde{v}_\theta = -\frac{\partial \tilde{p}}{\partial \tilde{r}} - \frac{(\gamma - 1) M_\theta^2}{\tilde{r}} \tilde{p} \\ & -i \left(\frac{k}{\tilde{A}} - \frac{mM_\theta}{\tilde{r}} - \bar{\gamma}M_x \right) \tilde{v}_\theta + \left(\frac{M_\theta}{\tilde{r}} + \frac{\partial M_\theta}{\partial \tilde{r}} + \left(\frac{\gamma - 1}{2} \right) \frac{M_\theta^3}{\tilde{r}} \right) \tilde{v}_r = -\frac{im}{\tilde{r}} \tilde{p} \\ & -i \left(\frac{k}{\tilde{A}} - \frac{mM_\theta}{\tilde{r}} - \bar{\gamma}M_x \right) \tilde{v}_x + \left(\frac{\partial M_x}{\partial \tilde{r}} + \left(\frac{\gamma - 1}{2} \right) \frac{M_x M_\theta^2}{\tilde{r}} \right) \tilde{v}_r = -i\gamma \tilde{p} \\ & -i \left(\frac{k}{\tilde{A}} - \frac{mM_\theta}{\tilde{r}} - \bar{\gamma}M_x \right) \tilde{p} + \frac{\partial \tilde{v}_r}{\partial \tilde{r}} + \left(\left(\frac{\gamma + 1}{2} \right) \frac{M_\theta^2}{\tilde{r}} + \frac{1}{\tilde{r}} \right) \tilde{v}_r + \frac{im}{\tilde{r}} \tilde{v}_\theta + i\gamma \tilde{v}_x = 0 \end{aligned}$$

Defining:

$$\alpha = \frac{k}{\tilde{A}} - \frac{mM_\theta}{\tilde{r}} - \bar{\gamma}M_x$$

and expanding the pressure equation to isolate the divergence of velocity term gives:

$$-i\alpha \tilde{v}_r - \frac{2M_\theta}{\tilde{r}} \tilde{v}_\theta = -\frac{\partial \tilde{p}}{\partial \tilde{r}} - \frac{(\gamma - 1) M_\theta^2}{\tilde{r}} \tilde{p}$$

$$\begin{aligned}
-i\alpha\tilde{v}_\theta + \left(\frac{M_\theta}{\tilde{r}} + \frac{\partial M_\theta}{\partial \tilde{r}} + \left(\frac{\gamma-1}{2} \right) \frac{M_\theta^3}{\tilde{r}} \right) \tilde{v}_r &= -\frac{im}{\tilde{r}}\tilde{p} \\
-i\alpha\tilde{v}_x + \left(\frac{\partial M_x}{\partial \tilde{r}} + \left(\frac{\gamma-1}{2} \right) \frac{M_x M_\theta^2}{\tilde{r}} \right) \tilde{v}_r &= -i\tilde{\gamma}\tilde{p} \\
-i\alpha\tilde{p} + \frac{M_\theta^2}{\tilde{r}}\tilde{v}_r + \left(\Phi_0\tilde{v}_{x,\phi} + \Phi_1\frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} + \Phi_2\frac{\partial^2\tilde{v}_{x,\phi}}{\partial\tilde{r}^2} \right) &= 0
\end{aligned}$$

Rewriting this in terms of the rotational and irrotational velocities,

$$\begin{aligned}
-i\alpha(\tilde{v}_{r,\phi} + \tilde{v}_{r,\omega}) - \frac{2M_\theta}{\tilde{r}}(\tilde{v}_{\theta,\phi} + \tilde{v}_{\theta,\omega}) &= -\frac{\partial\tilde{p}}{\partial\tilde{r}} - \frac{2\Gamma}{\tilde{r}}\tilde{p} \\
-i\alpha(\tilde{v}_{\theta,\phi} + \tilde{v}_{\theta,\omega}) + \left((1+\Gamma)\frac{M_\theta}{\tilde{r}} + \frac{\partial M_\theta}{\partial\tilde{r}} \right)(\tilde{v}_{r,\phi} + \tilde{v}_{r,\omega}) &= -\frac{im}{\tilde{r}}\tilde{p} \\
-i\alpha(\tilde{v}_{x,\phi} + \tilde{v}_{x,\omega}) + \left(\frac{\partial M_x}{\partial\tilde{r}} + \Gamma\frac{M_x}{\tilde{r}} \right)(\tilde{v}_{r,\phi} + \tilde{v}_{r,\omega}) &= -i\tilde{\gamma}\tilde{p} \\
-i\alpha\tilde{p} + \frac{M_\theta^2}{\tilde{r}}(\tilde{v}_{r,\phi} + \tilde{v}_{r,\omega}) + \left(\Phi_0\tilde{v}_{x,\phi} + \Phi_1\frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} + \Phi_2\frac{\partial^2\tilde{v}_{x,\phi}}{\partial\tilde{r}^2} \right) &= 0
\end{aligned}$$

Substituting in:

$$\begin{aligned}
\left(-i\alpha \left(\frac{-i}{\tilde{\gamma}} \left(\frac{\Gamma}{\tilde{r}}\tilde{v}_{x,\phi} + \frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} \right) + \tilde{v}_{r,\omega} \right) - \frac{2M_\theta}{\tilde{r}} \left(\frac{m}{\tilde{\gamma}\tilde{r}}\tilde{v}_{x,\phi} + \tilde{v}_{\theta,\omega} \right) \right) &= -\frac{\partial\tilde{p}}{\partial\tilde{r}} - \frac{2\Gamma}{\tilde{r}}\tilde{p} \\
\left(-i\alpha \left(\frac{m}{\tilde{\gamma}\tilde{r}}\tilde{v}_{x,\phi} + \tilde{v}_{\theta,\omega} \right) + \left((1+\Gamma)\frac{M_\theta}{\tilde{r}} + \frac{\partial M_\theta}{\partial\tilde{r}} \right) \left(\frac{-i}{\tilde{\gamma}} \left(\frac{\Gamma}{\tilde{r}}\tilde{v}_{x,\phi} + \frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} \right) + \tilde{v}_{r,\omega} \right) \right) &= -\frac{im}{\tilde{r}}\tilde{p} \\
\left(-i\alpha \left(\tilde{v}_{x,\phi} + i \left(\frac{1+\Gamma}{\tilde{\gamma}\tilde{r}} \right) \tilde{v}_{r,\omega} + \frac{i}{\tilde{\gamma}} \frac{\partial\tilde{v}_{r,\omega}}{\partial\tilde{r}} - \frac{m}{\tilde{\gamma}\tilde{r}}\tilde{v}_{\theta,\omega} \right) + \left(\frac{\partial M_x}{\partial\tilde{r}} + \Gamma\frac{M_x}{\tilde{r}} \right) \left(\frac{-i}{\tilde{\gamma}} \left(\frac{\Gamma}{\tilde{r}}\tilde{v}_{x,\phi} + \frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} \right) + \tilde{v}_{r,\omega} \right) \right) &= -i\tilde{\gamma}\tilde{p} \\
\left(-i\alpha\tilde{p} + \frac{M_\theta^2}{\tilde{r}} \left(\frac{-i}{\tilde{\gamma}} \left(\frac{\Gamma}{\tilde{r}}\tilde{v}_{x,\phi} + \frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} \right) + \tilde{v}_{r,\omega} \right) + \left(\Phi_0\tilde{v}_{x,\phi} + \Phi_1\frac{\partial\tilde{v}_{x,\phi}}{\partial\tilde{r}} + \Phi_2\frac{\partial^2\tilde{v}_{x,\phi}}{\partial\tilde{r}^2} \right) \right) &= 0
\end{aligned}$$

Defining:

$$\begin{aligned}
\Theta &= (1 + \Gamma) \frac{M_\theta}{\tilde{r}} + \frac{\partial M_\theta}{\partial \tilde{r}} \\
\tau_\phi &= \frac{\Gamma}{\tilde{r}} \tilde{v}_{x,\phi} + \frac{\partial \tilde{v}_{x,\phi}}{\partial \tilde{r}} \\
\tau_\omega &= \frac{\Gamma}{\tilde{r}} \tilde{v}_{r,\omega} + \frac{\partial \tilde{v}_{r,\omega}}{\partial \tilde{r}} \\
S &= \Gamma \frac{M_x}{\tilde{r}} + \frac{\partial M_x}{\partial \tilde{r}}
\end{aligned}$$

and gathering the rotational and irrotational components together:

$$\begin{aligned}
\begin{pmatrix} \frac{\alpha}{\tilde{\gamma}} \tau_\phi - \frac{2mM_\theta}{\tilde{\gamma} \tilde{r}^2} \tilde{v}_{x,\phi} \\ -i\alpha \tilde{v}_{r,\omega} - \frac{2M_\theta}{\tilde{r}} \tilde{v}_{\theta,\omega} \end{pmatrix} &= -\frac{\partial \tilde{p}}{\partial \tilde{r}} - \frac{2\Gamma}{\tilde{r}} \tilde{p} \\
\begin{pmatrix} -\frac{i\Theta}{\tilde{\gamma}} \tau_\phi - \frac{im\alpha}{\tilde{\gamma} \tilde{r}} \tilde{v}_{x,\phi} \\ +\Theta \tilde{v}_{r,\omega} - i\alpha \tilde{v}_{\theta,\omega} \end{pmatrix} &= -\frac{im}{\tilde{r}} \tilde{p} \\
\begin{pmatrix} -i\alpha \tilde{v}_{x,\phi} + \frac{\alpha}{\tilde{\gamma}} \tau_\omega + \frac{\alpha}{\tilde{\gamma} \tilde{r}} \tilde{v}_{r,\omega} + i\frac{m\alpha}{\tilde{\gamma} \tilde{r}} \tilde{v}_{\theta,\omega} \\ -i\frac{S}{\tilde{\gamma}} \tau_\phi + S \tilde{v}_{r,\omega} \end{pmatrix} &= -i\tilde{\gamma} \tilde{p} \\
\begin{pmatrix} -i\alpha \tilde{p} \\ -i\frac{M_\theta^2}{\tilde{\gamma} \tilde{r}} \tau_\phi \\ +\Phi_0 \tilde{v}_{x,\phi} + \Phi_1 \frac{\partial \tilde{v}_{x,\phi}}{\partial \tilde{r}} + \Phi_2 \frac{\partial^2 \tilde{v}_{x,\phi}}{\partial \tilde{r}^2} \\ -i\frac{M_\theta^2}{\tilde{r}} \tilde{v}_{r,\omega} \end{pmatrix} &= 0
\end{aligned}$$

I can't help but notice these groupings:

$$\begin{aligned}
A &= \frac{1}{\tilde{\gamma}} \tau_\phi - i\tilde{v}_{r,\omega} \\
B &= \frac{m}{\tilde{\gamma} \tilde{r}} \tilde{v}_{x,\phi} + \tilde{v}_{\theta,\omega}
\end{aligned}$$

which gives:

$$\begin{aligned}
\alpha A - \frac{2M_\theta}{\tilde{r}} B &= -\frac{\partial \tilde{p}}{\partial \tilde{r}} - \frac{2\Gamma}{\tilde{r}} \tilde{p} \\
i\Theta A - i\alpha B &= -\frac{im}{\tilde{r}} \tilde{p}
\end{aligned}$$

$$\begin{pmatrix} -i\alpha\tilde{v}_{x,\phi} + \frac{\alpha}{\tilde{\gamma}}\tau_\omega + \frac{\alpha}{\tilde{\gamma}r}\tilde{v}_{r,\omega} + i\frac{m\alpha}{\tilde{\gamma}r}\tilde{v}_{\theta,\omega} \\ -iSA \end{pmatrix} = -i\tilde{\gamma}\tilde{p}$$

$$\begin{pmatrix} -i\alpha\tilde{p} \\ -i\frac{M_\theta^2}{r}A \\ +\Phi_0\tilde{v}_{x,\phi} + \Phi_1\frac{\partial\tilde{v}_{x,\phi}}{\partial r} + \Phi_2\frac{\partial^2\tilde{v}_{x,\phi}}{\partial r^2} \end{pmatrix} = 0$$

I can't help but notice these groupings: