

Notes on two kinds of incompatibility witnesses

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Here the two definitions of incompatibility witnesses are briefly described. I call them the d-witness (d as "discrimination") and the g-witness (g as "geometric"), but perhaps better names can be found.

1 Some notations

Let us first fix some notations. Let \mathcal{H} , $\dim \mathcal{H} < \infty$ be a Hilbert space, $\mathcal{L}_s(\mathcal{H})$ the set of self-adjoint operators, $\mathcal{L}(\mathcal{H})^+$ the cone of positive operators. For a finite set X , let $\mathcal{O}(X)$ denote the set of observables with values in X .

Let X_1, \dots, X_k be finite sets and let $\mathcal{O}_i = (\mathcal{O})(X_i)$. Further, let Δ_i be the simplex with vertices $\{\delta_j^i, j \in X_i\}$. Let $\mathbf{S} = \Pi_i \Delta_i$ denote the Cartesian product with pointwise convex structure and let

$$\mathbf{s}_{n_1, \dots, n_k} = (\delta_{n_1}^1, \dots, \delta_{n_k}^k)$$

be the vertices of \mathbf{S} . We put

$$\begin{aligned} A(\mathbf{S}) &= \{\text{affine functions } f : \mathbf{S} \rightarrow \mathbb{R}\} \\ A(\mathbf{S})^+ &= \text{positive elements in } A(\mathbf{S}) \\ V(\mathbf{S}) &= A(\mathbf{S})^* \\ V(\mathbf{S})^+ &= \{\psi \in V(\mathbf{S}), \langle \psi, f \rangle \geq 0, \forall f \geq 0\} \end{aligned}$$

Note that $V(\mathbf{S})^+$ is a closed convex cone with base (isomorphic to) \mathbf{S} . Let

$$\mathbf{m}^i : \mathbf{S} \rightarrow \Delta_i$$

be the projection map and let $\mathbf{m}_j^i \in A(\mathbf{S})^+$ be such that

$$\mathbf{m}^i(\mathbf{s}) = \sum_{j \in X_i} \mathbf{m}_j^i(\mathbf{s}) \delta_j^i.$$

Note that \mathbf{m}_j^i generate the extremal rays of $A(\mathbf{S})^+$ (cf. [2, Sec. III]).

2 Collections of observables and compatibility

Let $\mathcal{O} = \mathcal{O}_1 \times \dots \times \mathcal{O}_k$ be the set of collections of observables. We can identify collections $(O^1, \dots, O^k) \in \mathcal{O}$, with affine maps $F : \mathcal{S}(\mathcal{H}) \rightarrow \mathbf{S}$ in an obvious way: we identify O^i with the affine map

$$O^i : \mathcal{S}(\mathcal{H}) \ni \rho \mapsto \sum_{j \in X_i} \text{Tr}[O_j^i \rho] \delta_j^i$$

and put

$$F(\rho) = (O^1(\rho), \dots, O^k(\rho)).$$

Conversely, any affine map $F : \mathcal{S}(\mathcal{H}) \rightarrow \mathbf{S}$ determines a collection of observables as

$$\mathrm{Tr}[O_j^i \rho] = \langle \mathbf{m}_j^i, F(\rho) \rangle, \quad \rho \in \mathfrak{S}(\mathcal{H}). \quad (1)$$

We can also determine the map F as

$$F(\rho) = \sum_{n_1, \dots, n_k} \mathrm{Tr}[F_{n_1, \dots, n_k} \rho] \mathbf{s}_{n_1, \dots, n_k}$$

for some elements $F_{n_1, \dots, n_k} \in \mathcal{L}_s(\mathcal{H})$. Note that these elements are not unique, but by (1) we must have

$$O_j^i = \sum_{n_1, \dots, n_{i-1}, n_{i+1}, \dots, n_k} F_{n_1, \dots, n_{i-1}, j, n_{i+1}, \dots, n_k}, \quad \forall i, j$$

and therefore the elements F_{n_1, \dots, n_k} must sum to identity. The collection is compatible if and only if F_{n_1, \dots, n_k} may be chosen positive, equivalently, the map F is entanglement breaking, since it can be written as a composition of a measurement and preparation.

3 The g -witnesses

Let \mathcal{A} be the vector space of all linear maps $\mathcal{L}_s(\mathcal{H}) \rightarrow V(\mathbf{S})$. Since any affine map on $\mathfrak{S}(\mathcal{H})$ can be extended to a linear map on $\mathcal{L}_s(\mathcal{H})$, the set \mathcal{O} can be identified with a subset of positive maps in \mathcal{A} . Let $\mathcal{P} \subset \mathcal{O}$ be the set of compatible collections, which corresponds to etb elements in \mathcal{O} . We now want to witness the set \mathcal{P} by affine functionals on \mathcal{O} . Any such affine functional is given in the form

$$p = \psi + t,$$

where $\psi \in \mathcal{A}^*$ and $t \in \mathbb{R}$. The dual space \mathcal{A}^* can be identified with the space of linear maps $\{W : V(\mathbf{S}) \rightarrow \mathcal{L}_s(\mathcal{H})\}$, with duality

$$\langle F, W \rangle = \mathrm{tr}(F \circ W),$$

here tr denotes the usual trace of a linear map on a vector space.

Any linear map $W : V(\mathbf{S}) \rightarrow \mathcal{L}_s(\mathcal{H})$ is determined by the values

$$\psi_{n_1, \dots, n_k} := W(\mathbf{s}_{n_1, \dots, n_k}) \quad (2)$$

of W on the extremal points $\mathbf{s}_{n_1, \dots, n_k}$ (since \mathbf{S} generates $V(\mathbf{S})$). Moreover, the map W is positive if and only if $\psi_{n_1, \dots, n_k} \in \mathcal{L}(\mathcal{H})^+$. Alternatively, since the elements \mathbf{m}_j^i generate $A(\mathbf{S})$, W is determined as

$$W(\mathbf{s}) = \sum_{i,j} \psi_j^i \mathbf{m}_j^i(\mathbf{s}), \quad (3)$$

where $\psi_j^i \in \mathcal{L}_s(\mathcal{H})$ are some (nonunique) elements satisfying

$$\psi_{n_1, \dots, n_k} = \sum_i \psi_{n_i}^i.$$

We then have

$$\begin{aligned}\langle F, W \rangle &= \text{tr}(F \circ W) = \sum_{i,j} \text{Tr} [\psi_j^i O_j^i] \\ &= \sum_{n_1, \dots, n_k} \text{Tr} [F_{n_1, \dots, n_k} \psi_{n_1, \dots, n_k}]\end{aligned}$$

Further, let $u_i := \sum_j |X_i|^{-1} \delta_j^i \in \Delta_i$ be the uniform probability distributions in Δ_i . Then $\bar{s} = (u_1, \dots, u_k)$ is the barycenter of \mathbf{S} . We have

$$\bar{\psi} := W(\bar{s}) = \frac{1}{\prod_i |X_i|} \sum_{n_1, \dots, n_k} \psi_{n_1, \dots, n_k} = \sum_{i,j} \frac{1}{|X_i|} \psi_j^i,$$

this element will be called the barycenter of W .

It now follows from known properties of cones of positive maps and their duals that $\{O^1, \dots, O^k\}$ is incompatible (that is, the corresponding map F is non-etb) if and only if there is a positive map W such that $\langle F, W \rangle < 0$. Such maps are called the g -witnesses.

Note that if all ψ_j^i are positive, W is etb and $\langle F, W \rangle$ is positive for all F . However, not every non-etb map is a witness (this is because the collections of observables generate only a section of the cone of positive maps, see [2, Thm. 2]).

4 g -witnesses and d -witnesses

A d -witness can be defined as an ensemble $\rho_j^i \in \mathcal{L}(\mathcal{H})^+$ such that $\sum_{i,j} \text{Tr} [\rho_j^i] = 1$ and

$$\pi(\{\rho_j^i\}) := \sup_{(O^1, \dots, O^k) \in \mathcal{O}} \sum_{i,j} \text{Tr} [O_j^i \rho_j^i] > \sup_{(O^1, \dots, O^k) \in \mathcal{P}} \sum_{i,j} \text{Tr} [O_j^i \rho_j^i] =: \pi_c(\{\rho_j^i\}).$$

This has some operational interpretation, see [3].

We now show the relation of these two types of witnesses. First, let $\{\rho_j^i\}$ be a d -witness. Define the map $\tilde{W} : \mathbf{S} \rightarrow \mathcal{L}_s(\mathcal{H})$ by

$$\tilde{W}(s) = \sum_{i,j} \rho_j^i m_j^i(s),$$

then for $(O^1, \dots, O^k) \in \mathcal{O}$ and the corresponding map F , we have

$$\sum_{i,j} \text{Tr} [O_j^i \rho_j^i] = \text{tr}(F \circ \tilde{W}).$$

As remarked before, the map \tilde{W} is positive and etb, so cannot be a g -witness itself. Let us denote $\pi_c := \pi_c(\{\rho_j^i\})$. Note that

$$\pi_c = \sup_{O \in \mathcal{O}(X_1 \times \dots \times X_k)} \sum_{n_1, \dots, n_k} \text{Tr} [O_{n_1, \dots, n_k} \tilde{\psi}_{n_1, \dots, n_k}],$$

where $\tilde{\psi}_{n_1, \dots, n_k} := \tilde{W}(s_{n_1, \dots, n_k}) = \sum_{i=1}^k \rho_{n_i}^i$. This was observed in [1] and follows also from the above considerations. Using the dual SDP program, we obtain that there is some state ρ_c such that $\tilde{\psi}_{n_1, \dots, n_k} \leq \pi_c \rho_c$ for all $n_i \in X_i$.

Let $\rho \in \mathfrak{S}(\mathcal{H})$ be arbitrary and let W_ρ be the constant map $\mathbf{S} \rightarrow \rho$. We have $\text{tr}(W_\rho \circ F) = 1$ for any $F \in \mathcal{O}$. Put

$$W := \pi_c W_{\rho_c} - \tilde{W}.$$

Then W is a g -witness: W is a positive map, since all images of the vertices are positive. If $F = (O^1, \dots, O^k)$ represents a collection of measurements such that $\text{Tr}[O_j^i \rho_j^i] > \pi_c$, then

$$\text{tr}(F \circ W) = \pi_c - \text{Tr}[O_j^i \rho_j^i] < 0.$$

In fact, $\{\rho_j^i\}$ is a d -witness if and only if \tilde{W} is a g -witness.

Conversely, let W be a g -witness, with the corresponding expressions given by Eqs. (2) and (3). For each i , let $\psi^i \in \mathcal{L}_s(\mathcal{H})$ be such that $\psi_j^i \leq \psi^i$ for all $j \in X_i$. Put

$$\tilde{\rho}_j^i := \psi^i - \psi_j^i, \quad c := \sum_{i,j} \text{Tr} \tilde{\rho}_j^i, \quad \rho_j^i := c^{-1} \tilde{\rho}_j^i.$$

Then $\{\rho_j^i\}$ is an ensemble, moreover, we have for $F = (O^1, \dots, O^k) \in \mathcal{P}$,

$$\begin{aligned} \sum_{i,j} \text{Tr}[O_j^i \rho_j^i] &= c^{-1} \left(\sum_i \text{Tr} \psi^i - \text{Tr}(F \circ W) \right) \leq c^{-1} \sum_i \text{Tr} \psi^i \\ &< c^{-1} \left(\sum_i \text{Tr} \psi^i - \inf_{F \in \mathcal{O}} \text{Tr}(F \circ W) \right) = \sup_{F \in \mathcal{O}} \text{Tr}[O_j^i \rho_j^i]. \end{aligned}$$

5 Two-outcome measurements

An advantage of g -witnesses may be that, at least in the case of two-outcome measurements, it is not difficult to find witnesses of this form, or to establish whether a given positive map is a witness or not:

If $|X_i| = 2$ for all i , we have the following characterization of g -witnesses (cf. [2, Coro 3]): Let $X_i = \{0, 1\}$ for all i and let $W : \mathbf{S} \rightarrow \mathcal{L}(\mathcal{H})^+$. Let ψ_{n_1, \dots, n_k} , $n_i \in \{0, 1\}$ and ψ_j^i be determined by (3) and (2). Put

$$\begin{aligned} e^i &:= \psi_{1, \dots, 1, 0, 1, \dots, 1} - \psi_{1, \dots, 1} \quad 0 \text{ in } i\text{-th place} \\ &= \psi_0^i - \psi_1^i \end{aligned}$$

Note that e^i are the images under W of vectors given by edges of the hypercube \mathbf{S} , adjacent to a fixed vertex. Further, the barycenter of W is

$$\bar{\psi} = \frac{1}{2} \sum_{i,j} \psi_j^i.$$

Then W is a witness if and only if

$$\sum_{i=1}^k \|e^i\|_1 > 2 \text{Tr} \bar{\psi}$$

which means that

$$\sum_i \|\psi_0^i - \psi_1^i\|_1 > \sum_{i,j} \text{Tr}[\psi_j^i].$$

(Note that the ψ_j^i cannot be all positive). The two-outcome measurements O^1, \dots, O^k attaining $\inf \text{tr}(F \circ W)$ are given by support projections of $(\psi_0^i - \psi_1^i)_+$.

Example 1. (Pairs and triples of two-outcome qubit measurements) Any incompatibility witness for pairs of two-outcome measurements is described by a parallelogram whose vertices $\psi_{i,j}$ are positive operators. We now look at a special case of squares such that the vertices are pure qubit states. Let

$$|y_k(\theta)\rangle = \cos\left(\frac{\theta}{2}\right)|0\rangle + i^k \sin\left(\frac{\theta}{2}\right)|1\rangle, \quad k = 0, \dots, 3, \quad \theta \in (0, \pi)$$

The witness $W(\theta)$ has vertices $\psi_{i,j}(\theta)$ given by

$$\psi_{0,0}(\theta) = y_0(\theta), \quad \psi_{0,1}(\theta) = y_1(\theta), \quad \psi_{1,1}(\theta) = y_2(\theta), \quad \psi_{1,0}(\theta) = y_3(\theta)$$

By the above condition, W is a witness iff

$$\|\psi_{1,1} - \psi_{0,1}\|_1 + \|\psi_{1,1} - \psi_{1,0}\|_1 > 2$$

which gives the condition

$$\sin(\theta) > \frac{1}{\sqrt{2}}.$$

Similarly, consider cubes with vertices on the Bloch sphere, given by

$$\begin{aligned} |y_{k,1}(\theta)\rangle &= \cos\left(\frac{\theta}{2}\right)|0\rangle + i^k \sin\left(\frac{\theta}{2}\right)|1\rangle \\ |y_{k,0}(\theta)\rangle &= \sin\left(\frac{\theta}{2}\right)|0\rangle + i^k \cos\left(\frac{\theta}{2}\right)|1\rangle, \quad k = 0, \dots, 3, \quad \theta \in (0, \pi/2) \end{aligned}$$

(the elements $y_{k,0}$ are vertices of the base of the cube). Here we get the condition

$$\sqrt{2} \sin(\theta) + \cos(\theta) > 1$$

References

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- [2] A. Jenčová, Incompatible measurements in a class of general probabilistic theories, arxiv:1705.08008, to appear in PRA
- [3] A. Toigo, Preliminary notes on incompatibility witnesses