## DPI and sufficiency, $\alpha > 1$

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Throughout these notes, we will assume that  $\alpha > 1$  and  $\max\{\alpha/2, \alpha - 1\} \leq z \leq \alpha$ . Let  $\psi, \varphi \in \mathcal{M}_*^+$  be such that  $D_{\alpha,z}(\psi||\varphi) < \infty$ , we will also assume that  $\varphi$  is faithul.

We put  $p := \frac{z}{\alpha}$  and  $q := \frac{z}{\alpha-1}$ , so that  $1/2 \le p \le 1 \le q$ . By the assumptions, there is some unique  $y \in L_{2z}(\mathcal{M})$  such that

$$h_{\psi}^{\frac{1}{2p}} = y h_{\varphi}^{\frac{1}{2q}}.$$
 (1)

By [5, 6], we have the variational formula

$$Q_{\alpha,z}(\psi \| \varphi) = \sup_{a \in \mathcal{M}_{+}} \alpha \operatorname{Tr} \left( h_{\psi}^{\frac{1}{2p}} a h_{\psi}^{\frac{1}{2p}} \right)^{p} - (\alpha - 1) \operatorname{Tr} \left( h_{\varphi}^{\frac{1}{2q}} a h_{\varphi}^{\frac{1}{2q}} \right)^{q}$$

$$= \sup_{a \in \mathcal{M}_{+}} \alpha \operatorname{Tr} \left( y h_{\varphi}^{\frac{1}{2q}} a h_{\varphi}^{\frac{1}{2q}} y^{*} \right)^{p} - (\alpha - 1) \operatorname{Tr} \left( h_{\varphi}^{\frac{1}{2q}} a h_{\varphi}^{\frac{1}{2q}} \right)^{q}$$

$$= \sup_{w \in L_{q}(\mathcal{M})^{+}} \alpha \operatorname{Tr} \left( y w y^{*} \right)^{p} - (\alpha - 1) \operatorname{Tr} w^{q},$$

this follows from the fact that  $h_{\varphi}^{\frac{1}{2q}}\mathcal{M}_{+}h_{\varphi}^{\frac{1}{2q}}$  is dense in  $L_{q}(\mathcal{M})^{+}$ . The supremum is attained at a unique point  $\bar{w} = (y^{*}y)^{\alpha-1} \in L_{1}(\mathcal{M})^{+}$ , uniqueness follows from strict concavity of the function  $w \mapsto \alpha \operatorname{Tr}(ywy^{*})^{p} - (\alpha - 1)\operatorname{Tr} w^{q}$ .

Let  $\Phi: \mathcal{M}_* \to \mathcal{N}_*$  be a 2-positive trace preserving map and let  $\varphi_0 := \Phi(\varphi)$ ,  $\psi_0 := \Phi(\psi)$ . Assume that also  $\varphi_0$  is faithful. Let  $\Phi_{\varphi}: \mathcal{N}_* \to \mathcal{M}_*$  be the Petz dual of  $\Phi$  with respect to  $\varphi$ , then we have

$$\Phi(h_{\varphi}^{1/2}ah_{\varphi}^{1/2}) = h_{\varphi_0}^{1/2}\Phi_{\varphi}^*(a)h_{\varphi_0}^{1/2}, \qquad \Phi_{\varphi}(h_{\varphi_0}^{1/2}bh_{\varphi_0}^{1/2}) = h_{\varphi}^{1/2}\Phi^*(b)h_{\varphi}^{1/2}, \qquad a \in \mathcal{M}, \ b \in \mathcal{N},$$

here  $\Phi^*: \mathcal{N} \to \mathcal{M}$  and  $\Phi_{\varphi}^*: \mathcal{M} \to \mathcal{N}$  are the 2-positive unital normal maps that are adjoints of  $\Phi$  resp.  $\Phi_{\varphi}$ . More generally, since for any  $r \geq 1$ ,  $\Phi$  is a contraction  $L_r(\mathcal{M}, \varphi)$  to  $L_r(\mathcal{N}, \varphi_0)$ , and similarly for  $\Phi_{\varphi}$ , there are positive contractions  $\Phi_{r,\varphi}: L_r(\mathcal{M}) \to L_r(\mathcal{N})$  and  $\Phi_{r,\varphi_0}: L_r(\mathcal{N}) \to L_r(\mathcal{M})$  such that

$$\Phi(h_{\varphi}^{\frac{1}{2r'}}ah_{\varphi}^{\frac{1}{2r'}}) = h_{\varphi_0}^{\frac{1}{2r'}}\Phi_{r,\varphi}(a)h_{\varphi_0}^{\frac{1}{2r'}}, \qquad \Phi_{\varphi}(h_{\varphi_0}^{\frac{1}{2r'}}bh_{\varphi_0}^{\frac{1}{2r'}}) = h_{\varphi}^{\frac{1}{2r'}}\Phi_{r,\varphi_0}(b)h_{\varphi}^{\frac{1}{2r'}}, \qquad a \in L_r(\mathcal{M}), \ b \in L_r(\mathcal{N})$$

here r' is such that  $\frac{1}{r} + \frac{1}{r'} = 1$ .

By DPI, we have  $D_{\alpha,z}(\psi_0\|\varphi_0) \leq D_{\alpha,z}(\psi\|\varphi) < \infty$ , so that there is some unique  $y_0 \in L_{2z}(\mathcal{N})$  such that

$$h_{\psi_0}^{\frac{1}{2p}} = y_0 h_{\varphi_0}^{\frac{1}{2q}}.$$

**Lemma 1.** Keeping the above assumptions and notations, we have for any  $w_0 \in L_q(\mathcal{N})^+$ 

$$\operatorname{Tr} \Phi_{q,\varphi_0}(w_0)^q \le \operatorname{Tr} w_0^q, \qquad \operatorname{Tr} (y\Phi_{q,\varphi_0}(w_0)y^*)^p \ge \operatorname{Tr} (y_0w_0y_0^*)^p.$$

*Proof.* The first inequality is immediate from the fact that  $\Phi_{q,\varphi_0}$  is a contraction. For the second inequality, let us first assume that  $w_0 = h_{\varphi_0}^{\frac{1}{2q}} b h_{\varphi_0}^{\frac{1}{2q}}$  for some  $b \in \mathcal{N}_+$ . Then

$$h_{\varphi}^{\frac{1}{2q'}}\Phi_{q,\varphi_0}(w_0)h_{\varphi}^{\frac{1}{2q'}} = \Phi_{\varphi}(h_{\varphi_0}^{\frac{1}{2q'}}w_0h_{\varphi_0}^{\frac{1}{2q'}}) = \Phi_{\varphi}(h_{\varphi_0}^{\frac{1}{2}}bh_{\varphi_0}^{\frac{1}{2}}) = h_{\varphi}^{\frac{1}{2q'}}h_{\varphi}^{\frac{1}{2q}}\Phi^*(b)h_{\varphi}^{\frac{1}{2q}}h_{\varphi}^{\frac{1}{2q'}},$$

so that  $\Phi_{q,\varphi_0}(w_0) = h_{\varphi}^{\frac{1}{2q}} \Phi^*(b) h_{\varphi}^{\frac{1}{2q}}$ . Therefore

$$\operatorname{Tr} (y \Phi_{q,\varphi_0}(w_0) y^*)^p = \operatorname{Tr} (y h_{\varphi}^{\frac{1}{2q}} \Phi^*(b) h_{\varphi}^{\frac{1}{2q}} y^*)^p = \operatorname{Tr} (h_{\psi}^{\frac{1}{2p}} \Phi^*(b) h_{\psi}^{\frac{1}{2p}})^p \ge \operatorname{Tr} (h_{\psi_0}^{\frac{1}{2p}} b h_{\psi_0}^{\frac{1}{2p}})^p$$

$$= \operatorname{Tr} (y_0 h_{\varphi_0}^{\frac{1}{2q}} b h_{\varphi_0}^{\frac{1}{2q}} y_0^*)^p = \operatorname{Tr} (y_0 w_0 y_0^*)^p,$$

here the inequality was proved in [2]. Since  $h_{\varphi_0}^{\frac{1}{2q}} \mathcal{N}_+ h_{\varphi_0}^{\frac{1}{2q}}$  is dense in  $L_q(\mathcal{N})^+$ , the statement follows.

**Theorem 1.** Let  $\Phi: \mathcal{M}_* \to \mathcal{N}_*$  be a 2-positive trace preserving map and let  $\psi, \varphi \in \mathcal{M}_*^+$  be such that  $D_{\alpha,z}(\psi \| \varphi) < \infty$ . Then  $D_{\alpha,z}(\Phi(\psi) \| \Phi(\varphi)) = D_{\alpha,z}(\psi \| \varphi)$  if and only if  $\Phi_{\varphi} \circ \Phi(\psi) = \psi$ .

*Proof.* By usual arguments, we may assume that both  $\varphi$  and  $\varphi_0$  are faithful. Then there is a conditional expectation  $\mathcal{E}$  onto the set of fixed points of  $\Phi^* \circ \Phi_{\varphi}^*$  such that  $\varphi \circ \mathcal{E} = \varphi$  and  $\Phi_{\varphi} \circ \Phi(\psi) = \psi$  if and only if also  $\psi \circ \mathcal{E}$ . This is what we are going to prove, using the extensions of conditional expectations to the Haagerup  $L_p$ -spaces in [4], see also [3, Sec. 1].

So assume that  $D_{\alpha,z}(\psi_0\|\varphi_0) = D_{\alpha,z}(\psi\|\varphi)$ . Let  $\bar{w} \in L_q(\mathcal{M})^+$  and  $\bar{w}_0 \in L_q(\mathcal{N})^+$  be the unique elements such that the suprema in the variational formulas for  $D_{\alpha,z}(\psi\|\varphi)$  resp.  $D_{\alpha,z}(\psi_0\|\varphi_0)$  are attained. We have by Lemma 1

$$D_{\alpha,z}(\psi||\varphi) \ge \alpha \operatorname{Tr} (y\Phi_{q,\varphi_0}(\bar{w}_0)y^*)^p - (\alpha - 1)\operatorname{Tr} \Phi_{q,\varphi_0}(\bar{w}_0)^q$$
  
 
$$\ge \alpha \operatorname{Tr} (y_0\bar{w}_0y_0^*)^p - (\alpha - 1)\operatorname{Tr} \bar{w}_0^q = D_{\alpha,z}(\psi_0||\varphi_0) = D_{\alpha,z}(\psi||\varphi),$$

so that both inequalities must be equalities. This implies that in particular

$$\operatorname{Tr} \bar{w}_0^q = \operatorname{Tr} \Phi_{q,\varphi_0}(\bar{w}_0)^q.$$

By uniqueness, we must also have  $\bar{w} = \Phi_{q,\varphi_0}(\bar{w}_0)$ . Let now  $\omega \in \mathcal{M}_*^+$  be given by  $h_\omega = h_\varphi^{\frac{1}{2q'}} \bar{w} h_\varphi^{\frac{1}{2q'}}$  and similarly  $h_{\omega_0} = h_{\varphi_0}^{\frac{1}{2q'}} \bar{w}_0 h_{\varphi_0}^{\frac{1}{2q'}}$ , then we get  $\Phi_{\varphi}(\omega_0) = \omega$  and also by definition of the sandwiched Rényi divergence,

$$\tilde{D}_{\alpha}(\omega_0 \| \varphi_0) = \operatorname{Tr} \bar{w}_0^q = \operatorname{Tr} \Phi_{a,\varphi_0}(\bar{w}_0)^q = \tilde{D}_{\alpha}(\Phi_{\varphi}(\omega_0) \| \Phi_{\varphi}(\varphi_0)).$$

By [1], this implies that  $\Phi_{\varphi}$  is sufficient with respect to  $\{\omega_0, \varphi_0\}$  and hence  $\Phi \circ \Phi_{\varphi}(\omega_0) = \omega_0$ . It follows that

$$\Phi_{\varphi} \circ \Phi(\omega) = \Phi_{\varphi} \circ \Phi \circ \Phi_{\varphi}(\omega_0) = \Phi_{\varphi}(\omega_0) = \omega,$$

which implies that  $\omega \circ \mathcal{E} = \omega$ . Using the extensions of  $\mathcal{E}$  and their properties, we get

$$h_{\varphi}^{\frac{1}{2q'}} \bar{w} h_{\varphi}^{\frac{1}{2q'}} = h_{\omega} = \mathcal{E}(h_{\omega}) = h_{\varphi}^{\frac{1}{2q'}} \mathcal{E}(\bar{w}) h_{\varphi}^{\frac{1}{2q'}},$$

which implies that  $\bar{w} = \mathcal{E}(\bar{w}) \in L_q(\mathcal{E}(\mathcal{M}))^+$ . But then also

$$|y| = \bar{w}^{\frac{1}{2(\alpha-1)}} = \bar{w}^{\frac{q}{2z}} \in L_{2z}(\mathcal{E}(\mathcal{M}))^+$$

Let y = u|y| be the polar decomposition of y, then we obtain from (1) that  $uu^* = s(\psi)$ . Further,

$$u^*h_{\psi}^{\frac{1}{2p}} = |y|h_{\varphi}^{\frac{1}{2q}} \in L_{2p}(\mathcal{E}(\mathcal{M}))$$

and by uniqueness of the polar decomposition in  $L_{2p}(\mathcal{M})$  and  $L_{2p}(\mathcal{E}(\mathcal{M}))$ , we obtain that  $h_{\psi}^{\frac{1}{2p}} \in L_{2p}(\mathcal{E}(\mathcal{M}))^+$ ,  $u \in \mathcal{E}(\mathcal{M})$ . Hence we must have  $h_{\psi} \in L_1(\mathcal{E}(\mathcal{M}))$  so that  $\psi \circ \mathcal{E} = \psi$ .

## References

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