

Energy–Frequency–Resonance: Phase-Locked Carriers and Holonomy Geometry

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Independent Research

Invariant carriers (photons, atoms, phonons, gravitational modes) phase-lock to the geometry of their trajectories. When the phase-closure (holonomy) condition holds, the accumulated phase is maximal and the observed effect appears as an effective interaction. In noisy media with time-dependent feedback, the carrier amplitude is governed by a closed-form kernel built from the exponential integral E_1 , enabling direct extraction of coherence time and feedback strength from data without assuming any predefined line shape. We provide the resonance rules, a holonomy unification formula, stable numerical evaluation, and interferometric predictions linking information flow to work at the $k_B T \ln 2$ limit.

FIGURE SUMMARY

I. RESONANCE AND PHASE LOCK

Define the phase error

$$\delta\phi(x, t, \omega; n) = \pi x + \omega t - 2\pi n, \quad n \in \mathbb{Z}. \quad (1)$$

Constructive interference occurs at $\delta\phi = 0$; path-closure yields stable standing modes. Around $\delta\phi = 0$, $|\Psi|$ attains a local maximum (resonance). Allowed loci:

$$\text{fixed } t: \quad x_n(t) = \frac{2n - \omega t / \pi}{2}, \quad \Delta x = 1, \quad (2)$$

$$\text{fixed } x: \quad t_n(x) = \frac{2\pi n - \pi x}{\omega}, \quad \Delta t = \frac{2\pi}{\omega}. \quad (3)$$

Ridge slope at fixed t gives a direct calibration,

$$\frac{d\omega}{dx} = -\frac{\pi}{t} \Rightarrow t = -\pi \left(\frac{d\omega}{dx} \right)^{-1}. \quad (4)$$

The resonance tolerance is

$$|\delta\phi| \lesssim \Delta\phi_{\text{tol}}, \quad \Delta\phi_{\text{tol}} \sim \max \left\{ \frac{1}{\text{SNR}}, \frac{t}{\tau} \right\}, \quad (5)$$

and a convenient quality estimate is

$$Q \approx \frac{\omega \tau}{1 + \alpha_{\text{eff}}}, \quad \alpha_{\text{eff}} \sim \overline{\epsilon(t) D_u[F(u(x))]} \quad (6)$$

Locking (PLL). To keep $\delta\phi \approx 0$,

$$\dot{x} = -\frac{\omega}{\pi} + k \delta\phi, \quad \dot{\omega} = -\frac{\pi}{t} + k \delta\phi. \quad (7)$$

II. OSCILLATOR IN NOISE WITH TIME-DEPENDENT FEEDBACK

We model the local carrier as

$$\Psi(x) = \int_0^\infty \frac{e^{-\epsilon(t)} \cos(\pi x + \omega t)}{1 + \epsilon(t) D_u[F(u(x))]} dt, \quad (8)$$

with causal damping $\epsilon(t) \geq 0$ and local functional sensitivity $D_u[F(u(x))]$. For $\epsilon(t) = t/\tau$ and locally constant $D := D_u[F(u(x))]$, define

$$\alpha(t) = \epsilon(t)D = \frac{t}{\tau}D, \quad a = \frac{1}{\tau} - i\omega, \quad \beta = \frac{D}{\tau}. \quad (9)$$

Then

$$\int_0^\infty \frac{e^{-at}}{1 + \beta t} dt = \frac{1}{\beta} e^{a/\beta} E_1\left(\frac{a}{\beta}\right), \quad (\text{Re } a > 0, \beta > 0), \quad (10)$$

and

$$\Psi(x) = \text{Re} \left\{ e^{i\pi x} \frac{\tau}{D} \exp\left(\frac{1 - i\omega\tau}{D}\right) E_1\left(\frac{1 - i\omega\tau}{D}\right) \right\}. \quad (11)$$

No line-shape is imposed; amplitude and phase are determined by $z = (1 - i\omega\tau)/D$.

Consistent limits. Let $S(z) = e^z E_1(z)$. For $|z| \rightarrow \infty$, use the asymptotic $S(z) \sim z^{-1} - z^{-2} + 2!z^{-3} - \dots$; for $|z| \ll 1$, $E_1(z) = -\gamma - \ln z + z - \frac{z^2}{4} + \dots$ (Euler γ).

III. HOLONOMY UNIFICATION

Effective interactions are holonomic phases collected by invariant carriers:

$$\Delta\varphi = \frac{1}{\hbar} \oint C_\mu dx^\mu. \quad (12)$$

Specializations:

$$\text{EM (U(1))} : C_\mu = qA_\mu, \quad \text{Weak (SU(2))} : C_\mu = gA_\mu^a T^a, \quad (13)$$

$$\text{Strong (SU(3))} : C_\mu = g_s A_\mu^a T^a, \quad \text{Grav. (eikonal/spin)} : C_\mu = p_\mu \text{ or } (14)$$

IV. INFORMATION TO WORK AT RESONANCE

At phase-lock and path-closure, visibility and accumulated phase are maximal; information transfer is optimal and the extractable work per bit obeys

$$W_{\text{max}} = k_B T \ln 2. \quad (15)$$

V. EXPERIMENTAL ROADMAP AND APPLICATIONS

- **Resonance maps:** scan (x, ω) ; read out $\Delta\omega = 2\pi/t$, slope $t = -\pi/(d\omega/dx)$, and bandwidth $\Delta\phi_{\text{tol}}$.
- **Feedback extraction:** fit amplitude and phase with Eq. (11) to obtain (τ, D) directly from data.
- **Holonomy tests:** vary closed loops in known A_μ (or rotation/eikonal geometry) and measure $\Delta\varphi$.
- **Platforms:** fiber Sagnac rings, integrated photonics, atom/neutron interferometers, superconducting microwave resonators, opto-/electro-mechanical phonon rings.
- **Control:** implement PLL on $\delta\phi$ to maintain resonance; operate near the $k_B T \ln 2$ information–work bound.

STABLE NUMERICS (NO OVERFLOW)

Evaluate the *scaled product* $S(z) = e^z E_1(z)$: use the asymptotic series for $|z| \gtrsim z_{\text{th}}$ and the local expansion for $|z| \lesssim z_{\text{th}}$ (consistent branch of $\ln z$). This avoids overflow and preserves amplitude/phase continuity in Eq. (11).

Appendix A: PLL Pseudocode for Phase Lock

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loop:
    delta_phi = wrap_to_pi(pi*x + omega*t - 2*pi*n)
    ctrl = Kp*delta_phi + Ki*Int(delta_phi) + Kd*d(delta_phi)/dt
    # lock in position (option A)
    x = x - (omega/pi)*dt + ctrl*dt
    # or lock in frequency (option B)
    # omega = omega - (pi/t)*dt + ctrl*dt
end

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AUTHOR CONTRIBUTIONS

[Jan Mikulik]: concept, derivations, figures, experiments; GPT-5 Thinking: formalization, unification, numerical stabilization, manuscript.

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- [1] Y. Aharonov and D. Bohm, Phys. Rev. **115**, 485 (1959).
[2] R. Colella, A. W. Overhauser, and S. A. Werner, Phys. Rev. Lett. **34**, 1472 (1975).
[3] M. Abramowitz and I. A. Stegun (eds.), *Handbook of Mathematical Functions* (Dover, 1972), § 5.1 (Exponential

Integral).

1) Oscillatory Field in Noise

$$\Psi(x) = \int_0^\infty \frac{e^{-\epsilon(t)} \cos(\pi x + \omega t)}{1 + \epsilon(t) D_u[F(u(x))]} dt$$

- $\epsilon(t)$... damping/decoherence (e.g. $\epsilon(t) = t/\tau$ or general noise function).
- $D_u[F(u(x))]$... derivative in the direction/functional sensitivity with respect to state $u(x)$ (your commutational substitution).
- πx ... harmonic periodicity in the “innate code”;
- ωt ... synchronization (carrier) phase.

Special cases

- **Pure wave (no noise):** $\epsilon(t) \rightarrow 0, D_u[\cdot] \rightarrow 0 \Rightarrow \Psi(x) = \int_0^\infty \cos(\pi x + \omega t) dt \rightarrow$ in distributional sense $\propto \delta(\omega) \cos(\pi x)$.
- **Exponential noise:** $\epsilon(t) = t/\tau \Rightarrow \Psi(x) = \Re \left\{ \frac{e^{i\pi x/(1+\alpha)} \int_0^\infty e^{-t/\tau} e^{i\omega t} dt}{1+\alpha} \right\} = \Re \left\{ \frac{e^{i\pi x/(1+\alpha)}}{1+\alpha} \cdot \frac{\tau}{1-i\omega\tau} \right\}$,
where $\alpha = \epsilon D_u[F(u(x))]$ (const.). \rightarrow Lorentzian profile and phase shift $\arg(1 - i\omega\tau)$.
- **Large noise:** denominator $\uparrow \rightarrow$ amplitude \downarrow (decoherence).

Meaning: $\Psi(x)$ is a *damped carrier oscillation* with information feedback via $D_u[F(u(x))]$. Noise controls both amplitude and phase delay.

2) One Formula for All “Forces”

$$\Delta\varphi = \frac{1}{\hbar} \oint C_\mu dx^\mu$$

where C_μ is the relevant connection:

- Electromagnetism (U(1)): $C_\mu = qA_\mu$.
- Weak (SU(2)): $C_\mu = gA_\mu^a T^a$.
- Strong (SU(3)): $C_\mu = g_s A_\mu^a T^a$.
- Gravity (ekonal/geodetic): $C_\mu = p_\mu$ (equiv. metric/spin connection).

Relation between the two planes:

- Holonomy ($\Delta\varphi$) gives *phase shift* between paths;
- Integral $\Psi(x)$ is the *local oscillator in noise*.
- Interferometrically, you measure $\Delta\varphi$ via phase shifts $\cos(\pi x + \omega t)$ and their damping $e^{-\epsilon(t)}$.

Yes — and here it is **cleanly typeset** (no changes of meaning, just formatting):

Exponential noise

$$\epsilon(t) = \frac{t}{\tau} \implies \Psi(x) = \Re \left\{ \frac{e^{i\pi x/(1+\alpha)}}{1+\alpha} \int_0^\infty e^{-t/\tau} e^{i\omega t} dt \right\} = \Re \left\{ \frac{e^{i\pi x/(1+\alpha)}}{1+\alpha} \cdot \frac{\tau}{1-i\omega\tau} \right\},$$

where $\alpha = \epsilon D_u[F(u(x))]$ (treated as a constant).

- **Lorentzian line shape (amplitude):**

$$|\Psi| \propto \frac{1}{|1+\alpha|} \frac{\tau}{\sqrt{1+(\omega\tau)^2}}.$$

- **Phase shift:**

$$\arg\left(\frac{\tau}{1-i\omega\tau}\right) = \arg(1-i\omega\tau)^{-1} = +\arctan(\omega\tau),$$

equivalently you can quote the factor’s argument as

$$\arg(1-i\omega\tau) = -\arctan(\omega\tau).$$

Total phase = $\pi x/(1+\alpha) + \arctan(\omega\tau) \pmod{2\pi}$.

Přesně tak — je to **řetězcové pravidlo** (a jednoduchá substituce parametrů).

- Integrál

$$\int_0^\infty e^{-t/\tau} e^{i\omega t} dt = \int_0^\infty e^{-(\frac{1}{\tau} - i\omega)t} dt = \frac{1}{\frac{1}{\tau} - i\omega} = \frac{\tau}{1 - i\omega\tau}.$$

Tady bereš exponenciálu $e^{g(t)} \leq g(t) = -(\frac{1}{\tau} - i\omega)t$; primitivní funkce je $-\frac{1}{g'(t)} e^{g(t)}$ a $g'(t) = -(\frac{1}{\tau} - i\omega)$ — přesně řetězcové pravidlo.

- Fázový faktor

$$\cos(\pi x + \omega t) \Big/ (1+\alpha) \rightsquigarrow \cos\left(\frac{\pi x}{1+\alpha} + \omega t\right),$$

tj. **reparametrizace** $x \mapsto x/(1+\alpha)$. Opět řetězcové pravidlo:

$$\frac{d}{dx} \cos\left(\frac{\pi x}{1+\alpha}\right) = -\sin(\cdot) \frac{\pi}{1+\alpha}.$$

Pro malé α : $\frac{1}{1+\alpha} \approx 1 - \alpha \Rightarrow$

phase $\approx \pi x(1 - \alpha)$ a amplitudový prefaktor $\approx (1 - \alpha)$.

- Výsledek:

$$\Psi(x) = \Re \left\{ \frac{e^{i\pi x/(1+\alpha)}}{1+\alpha} \cdot \frac{\tau}{1-i\omega\tau} \right\},$$

kde $\tau/(1-i\omega\tau)$ i rescaling $\pi x/(1+\alpha)$ jsou **přímým** důsledkem **chain rule** (a lineární substituce parametrů).

$$\alpha(t) = \epsilon(t) D_u[F(u(x))].$$

Pro exponenciální šum $\epsilon(t) = t/\tau$ tedy

$$\alpha(t) = \frac{t}{\tau} D, \quad D := D_u[F(u(x))] \text{ (lokálně konstantní v } x \text{)}.$$

Oprava integrálu (bez „konstantní α “)

$$\Psi(x) = \Re \left\{ e^{i\pi x} \int_0^\infty \frac{e^{-(\frac{1}{\tau} - i\omega)t}}{1 + \frac{D}{\tau} t} dt \right\}.$$

Označ $a = \frac{1}{\tau} - i\omega$, $\beta = \frac{D}{\tau}$. Pak platí uzavřený tvar

$$\int_0^\infty \frac{e^{-at}}{1 + \beta t} dt = \frac{1}{\beta} e^{a/\beta} E_1\left(\frac{a}{\beta}\right) \quad (\Re a > 0, \beta > 0),$$