



TUM School of Natural Sciences

Technische Universität München

Nuclear Structure Investigations of Light Nuclei
with the R3B Experiment

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Vollständiger Abdruck der von der TUM School of Natural Sciences der Technischen
Universität München zur Erlangung des akademischen Grades eines
Doktors der Naturwissenschaften (Dr. rer. nat.)
genehmigten Dissertation.

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Die Dissertation wurde am DATUM-HIER bei der Technischen Universität München
eingereicht und durch die TUM School of Natural Sciences am DATUM-HIER
angenommen.



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Techinsche Universität München
Fachbereich: Dense and Strange Hadronic Matter (E62)

2023

Abstract

The nucleosynthesis of heavy elements ($A > 56$) beyond iron, primarily via the rapid neutron-capture process (r-process), represents one of the most fascinating and complex phenomena in nuclear astrophysics. This process, which is responsible for the majority of heavy element production in the universe, provides a critical interface between astrophysics and nuclear physics, offering a unique "laboratory" to study nuclear reactions and properties far from stability. Neutron stars (NS), with their extreme densities and exotic compositions, are of exceptional interest in this context. These astrophysical objects serve as a natural laboratory to explore nuclear structure under extreme conditions, such as those described by the equation of state (EoS). Understanding the EoS, which governs the macroscopic properties of neutron stars, remains a formidable challenge. While astrophysical observations provide macroscopic constraints on NS models, nuclear experiments in terrestrial laboratories play a vital role in constraining the microscopic nuclear physics parameters embedded within these models. Moreover, neutron star mergers (NSMs), which are stellar collisions involving NSs, are recognized as the primary astrophysical sites for the r-process. These events generate conditions favorable for the production of neutron-rich nuclei, driving the synthesis of heavy elements. A comprehensive understanding of the r-process requires detailed knowledge of nuclear structure and reaction dynamics in regimes far from nuclear stability. However, experimental constraints on critical observables, such as fission yields and fission barriers, are limited due to the challenges associated with producing and studying these neutron-rich nuclei in laboratory settings. This doctoral thesis, titled "Nuclear Structure Investigations of Light Nuclei with the R3B Experiment", addresses these challenges by presenting the analysis of two key reaction channels using carbon isotopes. These investigations were conducted with the R3B (Reactions with Relativistic Radioactive Beams) setup during the commissioning experiment S444 as part of the FAIR Phase-0 campaign at GSI. Specifically, this work focuses on:

1. Charge-changing and total interaction cross sections: The study of $^{12}\text{C}+^{12}\text{C}$ collisions via the transmission method. These measurements provide critical insights into the nuclear structure and reaction mechanisms of light nuclei, which serve as essential benchmarks for understanding nuclear matter in astrophysical environments.
2. Quasi-free scattering (QFS) reactions: The investigation of the $^{12}\text{C}(\text{p},2\text{p})^{11}\text{B}$ reaction as a tool to probe nuclear structure. This approach demonstrates its potential for studying the dynamics of the r-process through fission in QFS experiments, offering an innovative method to address experimental limitations in current r-process studies.

The results of these analyses contribute to our understanding of the structure of light nuclei and provide critical experimental constraints relevant to nuclear astrophysics. By advancing our knowledge of fundamental nuclear properties and reaction dynamics, this work bridges the gap between laboratory-based nuclear physics and astrophysical processes, shedding light on the origin of heavy elements in the universe.

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1 Introduction

Talk something about first ideas from Greeks of the structure of matter /historical part Comparisons between atomic model and nuclear models And give a short overview about what you will talk in this introduction part and how related to your thesis

-

1.1 Nuclear Structure Models

Influenced by the good results/predictions of the Fermi gas model -> same model applied to nucleons in nucleus -> talk about the key features and limitations of this model -> liquid drop model -> can well postulate the binding energy (model from experimental observations) From the success of the atomic shell model -> nuclear shell model -> describes well shell closures->magic numbers ->picture and description of shell model ->limitations -> maybe mention optical model -> as it is used to predict cross section in elastic scattering reactions ->mention that there are also other models as mean field, cluster models, ab initio models, collective models which exceed the topics of this thesis, but give some references to it.

1.2 R3B Experiment - "The Universe in the Lab"

Central formula is the equation of states (EOS). We want to study it within a wide range.

- in astrophysics NS are therefore of special interest. Measuring their mass and radius -> this relation depends on the EOS.
- since NS are highly asymmetric matter (neutrons only) they are a good playground to test the different models. Since they are all consistent at rho0 (real world) they diverge for higher matter asymmetry.
- plot of some models
- formula of tailer expansion of Equation of state
- at FAIR we want to study nuclear matter under extreme conditions,highly asymmetric matter.
- constraining the symmetry energy
- this can be done via neutron removal cross section, cite T. Aumann, Bertulani etc.

- show picture of sensitiveness of this method, shortly mention that L was previously constrained in PREX experiments(give some literature for further readings)
- the predicted models rely on the input from reaction mechanisms -> $^{12}\text{C}+^{12}\text{C}$ good to probe the underlying reaction mechanisms-> one part of the thesis

Moreover to study the single particle states of nucleons inside nucleus, especially when going toward exotic nuclei-> quasi free scattering experiments.-> refer to the theory of those scattering experiments This was the first time we had our CALIFA CALORIMETER in its final frame. Refer to the application section.

To point out is the pilot experiment we had in 2021 to study fission via qfs. This could have a large impact in the understanding of the dynamics of fission far off the line of stability.

Explain the fission via qfs a little bit more in detail here.

This could explain maybe better the element abundances we have and if there is any island of stability at $Z > 126$.

2 Nuclear Scattering Reactions in the Glauber Framework

Explain path to go:

- short intro why we use scattering reactions: to study structure of nuclei (mostly via e. scattering as they only interact via electromagn. force) and refer to chapter. Introduce problems qe account body problem etc
- starting from elastic scattering at low energies going to medium to high energies
- explain glauber: eikonal wave approximations, optical limit which leads to optical potential model, description via nucleon nucleon interaction

2.1 Elastic scattering at low energies

- Rutherford scattering pioniering works (alpha particle on gold atoms) description in classical mechanics
- Rutherford scattering only valid if coulomb potential
- at low energies below pion production only elastic cross section
- introduce here the Scattering problem in quantum mechanics (see Schindler, p. 23)

- from classical physics we go to quantum mechanics and use partial wave decomposition with theta being the scattering angle
- include in this picture coulomb and nuclear interaction (as done in Kuk, page 120)

2.2 Nuclear Density distribution studies via elastic scattering

- Rutherford, Mott and Rosenbluth
- show something about form Factors
- are charge radii and nuclear radii measured?
- note approximation for ^{12}C that neutron radius same as proton radius

2.3 Glauber Model for nuclear scattering at high energies

- explain model assumption
- eikonal wave approximation: high incoming momentum, low scattering angle
- Show picture of scattering
- optical limit: - Nucleons at high energy are undeflected due to large momentum (linear trajectory) -Nucleus large compared to nucleon-nucleon force -Motion of nucleons independent of nucleus -overall cross-section described in terms of nucleon-nucleon cross section
- Description in the Probability Approach (also here I use all the eiconal and optical limit approximation)
- Description in the Eikonal optical-limit approximation
- Comparison of both descriptions methods for total interaction cross section should end up with same. Advantage of PA is that you can calculate the cross section for a defined number of removed projectile nucleons (Schindler, p. 49)

2.3.1 Extensions to the Glauber model

- things which cannot be neglected and influence the cross section
- in medium modifications (see Lukas thesis)
- Coulomb interaction

- Pauli blocking

2.4 Cross sections for $^{12}C + ^{12}C$

- Total reaction cross section
- charge changing cross section → maybe I have to include lise++ calculations
- isotope corrections- neutron removal → maybe I have to include lise++ calculations

2.5 Quasi-Free Scattering (QFS) Reactions

Quasi-free scattering (QFS) reactions, as a subset of direct reactions, are processes where a projectile nucleon interacts with a target nucleon in a single, fast, and highly localized event. The timescale of these interactions is extremely short, approximately 10^{-21} s, implying that the relative kinetic energies of the participants are high, typically $\gtrsim 100$ AMeV.

In QFS experiments conducted in direct kinematics, a proton serves as the projectile, colliding with a nucleon or a cluster within the target nucleus. Conversely, in inverse kinematics, the nucleus of interest becomes the projectile, while a proton or proton-like particle is used as the target. Despite these differing setups, both approaches are fundamentally equivalent, differing only due to a Lorentz transformation of the reference frame.

As the name implies, QFS reactions are conceptually similar to free nucleon-nucleon scattering, with the primary approximation being that the influence of the residual nucleus is neglected to first order. This approximation simplifies the theoretical description of the process, allowing to treat the interaction as a two-body problem within a (only weakly interfering) nuclear environment.

The first experiments confirming the existence of quasi-free scattering processes were conducted at the University of California, Berkeley, in 1952 by O. Chamberlain and Emilio Segrè[1]. In their study, lithium nuclei were bombarded with 350 MeV protons, and coincident proton pairs were observed with an opening angle of approximately 90°. That same year, J.B. Cladis, W.N. Hess, and B.J. Moyer published results on the scattering of 340 MeV protons on deuterium and carbon targets[2], further substantiating the phenomenon.

In 1957, Tyren, Maris, and Hillman designed an experiment aimed at fully characterizing proton-proton collisions within the quasi-free scattering framework[3]. Their results validated the assumption of a direct and clean interaction between the projec-

tile and the target nucleon, free from significant distortions caused by the surrounding nucleus. Furthermore, these experiments demonstrated that QFS reactions could serve as a powerful tool for probing nuclear structure and testing predictions of the shell model. Specifically, they enabled the study of key nuclear parameters such as spin-orbit splitting and energy differences between nuclear shells and hence to probe the shell evolution.

In QFS reactions, these parameters can be extracted by analyzing the reaction products, which include the two correlated outgoing protons, the residual nucleus, and any gamma rays emitted during de-excitation of the residual nucleus. The detailed measurements of these observables provide critical insights into the underlying nuclear structure and dynamics.

The experimental discoveries and theoretical insights from these early studies catalyzed significant advancements in the theoretical modeling of QFS reactions. These models have since become essential tools for understanding nucleon-nucleon interactions within the nuclear medium and for refining our knowledge of nuclear structure and reaction mechanisms.

2.5.1 Kinematics of QFS Reactions

A simplified picture, which however explains the essential physics of the QFS reaction, can be found in figure 1: we have an incoming proton knocking out a nuclear constituent (proton/neutron) in a free nucleon-nucleon collision ending with a final state having the scattered proton, the scattered off nuclear constituent and the rest nucleus ($A-1$)*. In this picture index 0 is assigned to the incoming proton, 1 to the knocked out nucleon, 2 to the outgoing projectile proton, A to the initial nucleus and $A-1$ to the final nucleus accordingly. From energy -momentum conservation the reaction can be expressed as:

$$P_A + P_0 = P_1 + P_2 + P_{A-1} \quad (1)$$

with P_i the four momentum (\mathbf{p}_i, E_i).

In direct kinematics as presented in figure 1 the separation of the ejected nucleon for a certain final state of the nucleus $A-1$ is given by¹:

$$S = T_0 - (T_1 + T_2 + T_{A-1}); \text{ with } T_i \text{ the kinetic energy of particle } i \quad (2)$$

In the idealized shell model the separation energy equals to the (negative) energy of the nucleus' single-particle state. For the case the proton knocks out the least bound

¹In inverse kinematics the four-momentum vectors need to be boosted to the center of mass frame of the initial nucleus via Lorentz transformation.

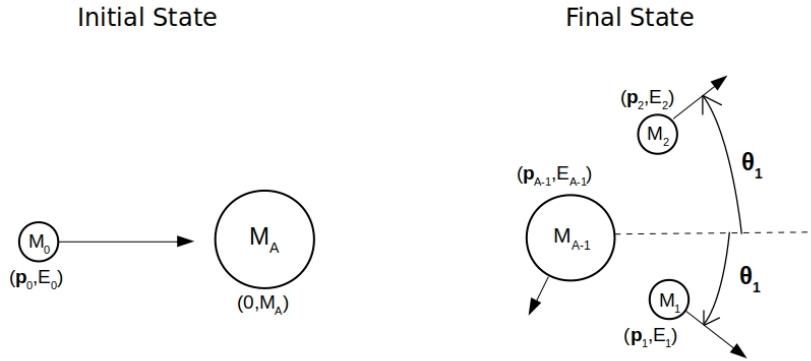


Figure 1: Simplified picture of the QFS reaction process in direkt kinematics.

nucleon (proton/neutron) the final nucleus is lying in the ground state with $E_{A-1} = M_{A-1}c^2 + T_{A-1}$.

If a nucleon has been ejected from an inner shell resulting in a hole state, the final nucleus will be in an excited state:

$$E_{A-1}^* = M_{A-1}c^2 + T_{A-1} + E_{exe} \quad (3)$$

The excitation energy E_{exe} is reflected in the difference in the separation energy of the least bound nucleon and the ejected one. From the experimental point of view E_{exe} is directly accessible via gamma detection from the transition of the final nucleus from the excited to the ground state. According to the picture of having a nucleon nucleon scattering process with no influence of the residual nucleus A-1 we can approximate $\mathbf{p}_A \approx \mathbf{p}_i + \mathbf{p}_{A-1}$ where \mathbf{p}_i is the initial nucleon momentum inside the initial nucleus A. Since the initial nucleus is at rest ($\mathbf{p}_A = 0$), the recoil momentum of the nucleus in final state \mathbf{p}_{A-1} equals to $-\mathbf{p}_i$ the momentum of the initial nucleon pointing in opposite direction.

In addition the four-momentum of the inner nucleon can be deduced from momentum measurement of the initial and final state free nucleons:

$$P_i \approx P_{miss} \equiv P_1 + P_2 - P_0 \quad (4)$$

where P_{miss} is the so called "measured missing four-momentum of the reaction"^[4]². Thereupon, the separation energy measurement and the recoil momentum distribution

² P_{miss} is only equal to P_i for the unperturbed QFS (no ISI/FSI) case.

fully describe the single particle state in the various shell levels. In addition and as complementary method γ rays can be measured in coincidence with the reaction and consequently exclusive cross section and momentum distribution measurements of the single particle states are accessible as illustrated in figure blabla.

Considering the removal of a single nucleon from an initial nucleus with A nucleons and initial spin I Ψ_i^A and final nucleus state Ψ_f^{A-1} with final spin I_f an overlap function between intial and final state many-body wave function can be written as:

$$\langle \vec{r}, \Psi_f^{A-1} | \Psi_i^A \rangle = \sum_j C_j^{if} \psi_j(\vec{r}), \quad \text{with } |I - I_f| < j < |I + I_f| \quad (5)$$

where $S_j^{if} = |C_j^{if}|^2$ is the commonly named spectroscopic factor, see ref [5] sec.2.1. S_j^{if} is summed over all final single particle states m (from $-j$ to j). It is unity for nucleon removal from a pure single particle state and equals $(2j+1)$ when the nucleon was removed from a filled j -subshell. The spectroscopic factor is linked to the exclusive experimental cross section measurement of a single particle state and the theoretical predicted one, as in ref [5]:

$$\sigma_{th}^{if} = \sum_j S_j^{if} \sigma_{sp}(nlj) \quad (6)$$

where σ_{sp} are the theoretical cross sections for the normalized wave functions Ψ_j of the final state nucleus A-1 with the appropriate quantum numbers.

From this considerations the spectroscopic factor can be used to probe the theoretical shell predictions. In past this was already done with direct qfs-reactions ($e, e'p$). Results for the spectroscopic factor with data obtained at the NIKHEF facility is shown in figure 2. The substantial reduction of the spectroscopic factor with respect to the independent particle model (IPM) or mean field of $\approx 35\%$ indicates a substantial depletion of the single-hole states and inferring from this a refined model prediction has to be applied³.

While qfs-reaction with electrons ($e, e'p$) in direct kinematics is a valuable method to make precise measurements for stable (target) nuclei it is not suitable for the experimental analysis with exotic (neutron or proton rich) nuclei. Due to their instability they have a short lifetime (e.g. ^{52}Ca with $\tau = 4.6s$) they can hardly serve as targets. Inducing the reaction in inverse kinematics - having electrons as targets and the exotic nuclei of interest as projectiles - is also not feasible since no free electrons can be captured as target. The alternative approach is to use proton induced quasi-free scattering in inverse kinematics where the exotic beam impinges on an extended proton target such

³The mean field potential does not consider spin-spin interaction V_{ss} , non central tensor-potential V_T or spin-orbit interaction V_{LS} .

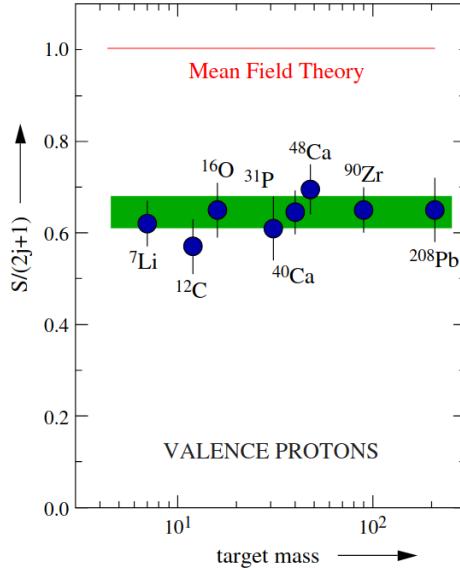


Figure 2: Normalized spectroscopic factors from the $(e, e'p)$ reaction as a function of target mass taken from [6]. As input for the theoretical cross sections the prediction from mean field models were used.

as liquid hydrogen H_2 or a fixed proton rich target such as CH_2 .⁴

Summarizing, the qfs-reaction technique in inverse kinematics with exotic nuclei as in-flight projectiles and a proton-like target opens new possibilities to probe theoretical model predictions of deformed nuclei far off stability which haven't been accessible before.

For precise cross section measurements of the single particle states σ_{sp} the kinematical characteristics of the qfs-reaction products have to be considered for correct identification of qfs-events and clear background subtraction. The following descriptions are already encoded in equation 1.

As starting point one compares the qfs-reaction of the two nucleons with the two dimensional non-central collision of free pointlike particles in nonrelativistic kinematics. Since both kinetic energy and momenta are conserved a clear signature is expected: the opening angle of the scattered particles is exactly 90° .

However, at kinetic energies of 400 AMeV and more, relativistic effects affects the opening angle of the scattered particles.⁵ put in here the formula for the opening angle in realtivisti kinematics, to be derived

For the final description of the process it has to be considered that the projectile

⁴Herefore a carbon target reference run is used to extract the qfs-reaction events from the CH_2 target run.

⁵Relativistic considerations become relevant at velocities $\beta \gtrsim 0.1c$. This corresponds to ≈ 4.5 MeV.

nucleon inside the nucleus has a non negligible momentum (see P_i in equation 4) and in addition the separation energy $S_{p/n}$ has to be expended to remove the scattered nucleon of the nuclear potential of the projectile nucleus.

To account for the nucleon momentum the picture of the fermi gas model where protons and neutrons freely move inside the nucleus' potential is applied. In the ground state of the nucleus the nucleons can reach a momentum up to the Fermi momentum p_F . Except for the light nuclei, the Fermi momentum is almost independent of A and amounts to $\approx 250 \text{ MeV}/c$ ⁶. The mean quadratic momenta of the nucleons is related to the Fermi momentum by:

$$\langle \mathbf{p}^2 \rangle = \frac{3}{5} p_F^2$$
 (7)

The width of the momentum distribution of the nucleons is given in the Goldhaber model by:

$$\sigma^2 = \sigma_0^2 K(A - K)/(A - 1), \text{ with } \sigma_0 \approx 90 \text{ MeV}/c$$
 (8)

where A is the mass number of the projectile nucleus and K the mass number of the fragment after scattering. For a detailed derivation and further readings see ref. [7] and [8]. In the impulse approximation, where we assume that the interaction between the projectile and target nucleons has approximately the same form as the interaction between two nucleons in free space, the inner momenta of the scattered nucleons smear out the angular correlations of the outgoing fragment - which has the same momentum as the knocked out nucleon in the cms system of the nucleus but points in opposite direction - as well as of the two nucleons involved in the scattering. As consequence the opening angle distribution of the two scattered nucleons gets broadened as well as the azimuthal angular correlation which for the assumption of zero nucleon momentum sharply peaked at $\Delta\phi = 180^\circ$.

In view of the added nucleon momentum the kinematics get expanded from a two dimensional scattering reaction to a three dimensional process where the reaction plane is determined by the plane spanned by the momentum vector of the projectile nucleus \mathbf{p}_A and the scattered nucleon after the reaction⁸. With the momentum measurement of the two correlated outgoing nucleons and the momentum of the incoming projectile nucleus it is possible to retrieve the internal momentum of the knocked out nucleon perpendicular to the reaction plane via the formula derived in reference[9]:

$$Q_{\perp i} = \sin(\theta_{1/2}) \cdot \sin(\phi_{1/2} - \phi_{2/1}) \cdot \mathbf{p}_{2/1}$$
 (9)

⁶For light nuclei like ^{12}C , $p_F \approx 230 \text{ MeV}/c$ can be assumed

⁷From the derivation of the mean kinetic energy of the nucleons $\langle E_{\text{kin}} \rangle = \frac{\int_0^{p_F} E_{\text{kin}} p^2 dp}{\int_0^{p_F} p^2 dp} = \frac{3}{5} \cdot \frac{p_F^2}{2M}$

⁸For (p,2p) reactions there is an ambiguity which of the scattered protons \mathbf{p}_1 or \mathbf{p}_2 origin from the projectile nucleus or the proton like target.

In case of (p,2p) reactions it is impossible to track back the origin of each nucleon - from the projectile nucleus or the proton-like target. Iferring from this there are two possible solutions for Q_{\perp} . The ambiguity can be resolved by incorporating the momentum of the fragment nucleus perpendicular to the reaction plane $Q_{\perp A-1}$. The fragment momentum vector should be of the same amout but point in opposite direction to $Q_{\perp i}$. TODO: insert here the correlation plot of VALERII, check if that is correct.

Up to this point the scattering process was described similar to the free elastic scattering in relativistic kinematics with the projectile nucleon having a non negligible momentum. However the projectile nucleon is an off-shell particle bound inside the nucleus.

$$\mu_i = \sqrt{m_A^2 + m_{A-1}^2 + 2m_A\sqrt{m_{A-1}^2 + |\mathbf{p}_i|^2}} \quad (10)$$

in the c.m.s. of the projectile nucleus, with

$$0 < \mu_i < m_n \quad (11)$$

where m_n is the mass of the appropriate free nucleon. To knock the bound nucleon out of the nucleus the separation energy has to be overcome which can be observed in the reduced opening angle of the two outgoing correlated nucleons.

- put in the L.V. Chulkov formula and plot to get the transversal mom etc
- Last correction: whe have to consider the separation energy (and NOT binding energy) -> makes angle again smaller etc show also on the formula that it gets smaller...this is then also an explanation why it gets smaller if nucleon is knocked out from a deeper shell.

2.5.2 Cross Sections for QFS Reactions - Qualitative Considerations

see more in the standard work, TO BE DONE

2.5.3 Application Fields of QFS Reactions

Over several decades of experimental and theoretical research, quasi-free scattering (QFS) reactions have been firmly established as a powerful and direct tool for probing the microscopic structure of atomic nuclei. These reactions provide critical insights into nuclear correlations, single-particle properties, and the momentum distributions of nucleons within the nucleus.

The versatility of QFS extends across various applications, including the study of short-range nucleon-nucleon interactions, exotic nuclear states, and the modification of nucleon properties in nuclear matter. Advances in experimental facilities and state-of-the-art detector systems, have significantly improved the precision and scope of QFS

measurements. These developments have enabled in-depth investigations of nuclear dynamics, the structure of unstable isotopes, and fundamental aspects of quantum many-body systems, contributing to a deeper understanding of nuclear and subnuclear phenomena.

This subsection will point out the most exciting and promising application fields of QFS reactions with focus on the applicability in the R3B experimental setup. A detailed review can be found in ref [10].

Single-Particle Spectroscopic Strength As already mentioned in section 2.5.1 in many experiments a reduction in the spectroscopic strength of about 35% with respect to the independent particle model (IPM) and shell model predictions was observed. In one-nucleon removal reactions - experiments with isotope beams impinging on a composite nuclear target - the extraction of the missing spectroscopic strength is challenging as the kinematical pattern is highly complex. In contrast, the qfs reactions mechanism in inverse kinematics retains a clear kinematical signature which makes it a valuable tool to study the quenching of the spectroscopic factor which originates from residual correlations between bound nucleons inside the nucleus.

Several experiments were carried out at the GSI Facility to study the reduction factor of the spectroscopic strength and its evolution over a broad range in isotopic chains, such as for oxygen shown in figure 3. With the commmissioning of the SIS100 and SuperFRS at GSI it will be achievable to study (p,2p) reactions with very short lived nuclei at reasonable intensities which will presumably draw even more attention to qfs reaction approach.

Gamma Spectroscopy Alongside to the study of the single-particle spectroscopic strengths qfs reactions provide valuable information of the shell structure and deformation of the fragment ($A-1$) via gamma spectroscopy. While quasi-free scattering with nucleons in the outermost shell yield a fragment in the ground state, the scattering with nucleons in inner shells result in bound excited states of the ($A-1$) fragment. This excited fragment instantaneously transits to the ground state by the emission of a Doppler shifted γ ray.

A big advantage of the (p,2p) reaction is ability to make precise vertex reconstruction by accurate measurement of the two outgoing correlated protons. This is of particular importance for extended targets and determines the emission point within the target, enabling an accurate calculation of the fragment's velocity and therefore a precise Doppler correction.

From gamma spectroscopy the excitation energy $E(2+)$, the first excited state of a nucleus, can be observed which is a fundamental observable of nuclear struc-

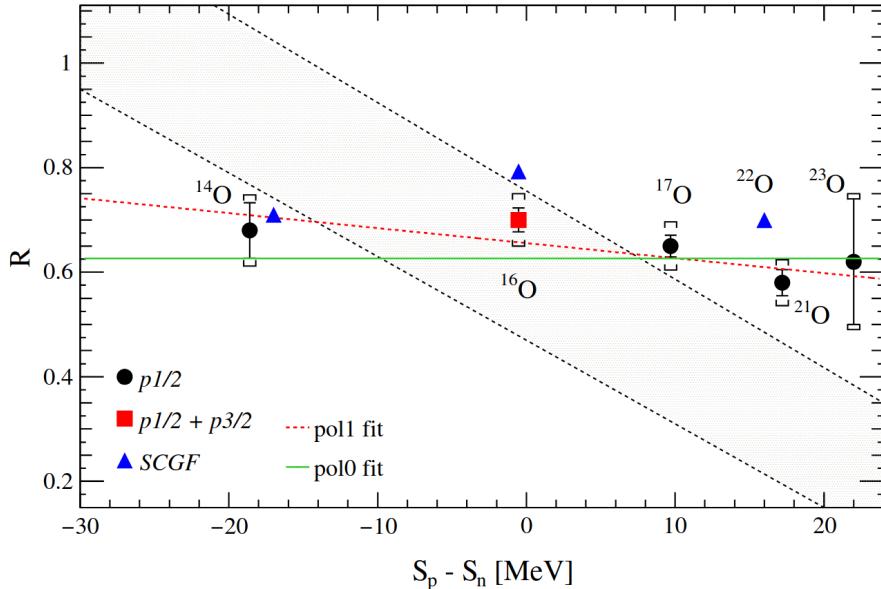


Figure 3: Extracted reduction factors of the spectroscopic strength from (p,2p) measurements (circles and square) as a function of the difference between proton and neutron separation energy S_p/S_n . The blue triangles correspond to the predictions from self-consistent green functions (SCGF). The shaded area indicates the trend from an analysis of one-nucleon removal cross sections. From Ref. [11].

ture, providing key insights into the shell configuration and deformation of the nucleus[10]. Especially when going towards exotic neutron rich isotopes the study of the gamma emission lines can strongly contribute to the understanding of nuclear shell evolution.

QFS to probe inner clustering and halo formation Clusterization inside nuclei is a phenomena widely observed in experimental astrophysics with large implications on processes such as the synthesis of carbon inside stars via triple α clusterization [12], α - radioactivity, or the evolution of core collapse supernovae[13], just to mention some. Its presence has also to be implemented in the state of the art equation of state calculations.

Single-particle cluster states can be directly accessed via qfs ($p,p\alpha$) reactions. In the low energy regime of 100 AMeV the $^{12}\text{C}(p,p\alpha)^8\text{Be}$ reaction, see ref [14], as benchmark study approved the reaction mechanism and up to a scaling factor the measured cross sections aligned exceptionally well to the free $p-\alpha$ scattering measurements.

Moreover the qfs mechanism allows to acces light nuclei going towards the neutron drip line via (p,pn) reaction. These light nuclei are mostly weakly bound with an extended low density neutron-matter distribution forming a so-called *nuclear-*

halo.

The most prominent candidate for such is ^{11}Li , often called *Borromean three-body system*[15], consisting of two neutrons interacting with a core (9Li) via weak, short-range interactions. First evidence for the correlation between the two neutron was strongly pushed by experimental studies at GSI, see [16], and theoretical work by Bertulani et. al.[17].

The latest experiment with the focus on the study of multi-neutron correlations in drip-line nuclei was carried out at R3B in 2022 (experiment S509) probing broad isotopic chains of Li, Be, B, C and N via the qfs mechanism.

Short Range Correlations(SRC) Short range correlations refer to elementary nucleon-nucleon interactions within a nucleus and occur over very short distances, typically on the order of 1-2 femtometers. These correlations are characterized by pairs of nucleons with high relative momentum ($> k_F$, Fermi momentum) and low total momentum. These NN- interactions are treated as good explanatory candidate for observed deviations from mean field approximation. Many phenomena, such as the above discussed reduction in the spectroscopic strength and the EMC effect[18], the observation that the cross section for deep inelastic scattering from an atomic nucleus is different from that of the same number of free protons and neutrons, are associated to the short range correlations inside the nucleus.

From isotopic chain studies[19] it has been shown that there is an indication of SRC dependency on isospin, see figure blabla, where SRC are predominantly preferred by p-n pairs than by nn or pp pairs. This again can have significant impact for asymmetric nuclei and imply a stronger quenching of the spectroscopic strength for proton single particle states below the Fermi momentum for neutron rich matter.

Since pioneering experiments were made at JLab and Brookhaven via $(e, e'p)$ and $(e, e'n)$ reactions several experimental campaigns were carried out at GSI with the R3B setup. In 2022 the S509 experiment was performed to exploit for the first time the use of short-lived nuclei scattering off a proton probe in inverse kinematics at R3B, followed by the S091 experiment in 2024 with the focus on probing NN-correlations in atomic nuclei via (p, pd) QFS reactions[20].

Fission via QFS Fission induced via quasifree scattering (QFS) provides a powerful method to extract fission barriers on an event-by-event basis. The fundamental concept involves the knockout of a proton or neutron from an incoming exotic ion beam. By measuring the energy of the correlated emitted nucleons, it becomes possible to probe the excitation state of the daughter nucleus. In cases where deeply bound nucleons are knocked out, the resulting daughter nucleus of-

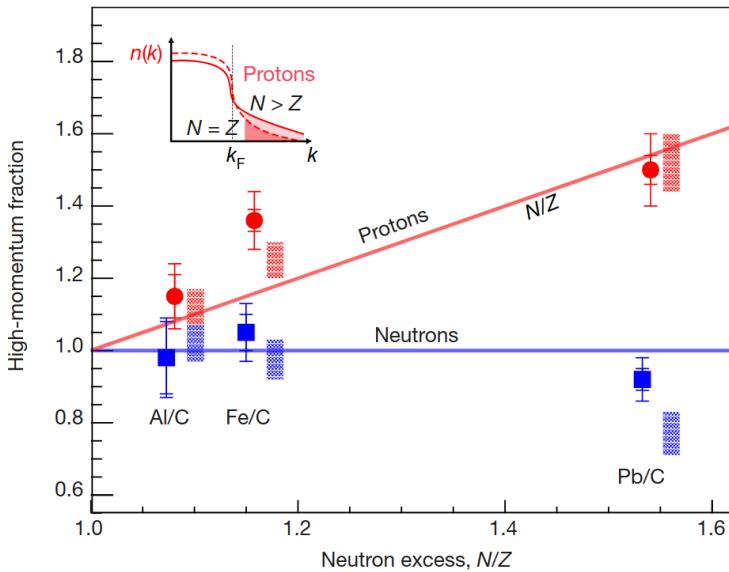


Figure 4: Fraction of high-momentum neutrons and protons with respect to the N/Z ratio for several carbon isotopes. A clear trend towards high momentum proton distribution for neutron rich carbon isotopes (red line) is shown, while the high momentum neutron distribution stays constant with respect to the N/Z ratio (blue line). From reference [19].

ten populates an excited state, which can be experimentally observed. The R3B setup is particularly well-suited for such investigations, as it enables full kinematic reconstruction of the reaction products. This approach allows for the detailed exploration of the potential-energy surface and the fission dynamics across a broad range of fissility and excitation energies, taking advantage of relativistic radioactive beams.

A pilot experiment was conducted in 2021 at R3B using a stable ^{238}U beam. The data collected from this experiment provide simultaneous information on several key fission observables, which can be used to constrain theoretical calculations. These calculations aim to describe both the static and dynamic properties of the fission process and enable comparisons with predictions from various theoretical frameworks, including phenomenological approaches, macroscopic-microscopic models, and fully microscopic statistical or time-dependent Hartree-Fock calculations.

A critical aspect of this experimental setup is the simultaneous measurement of the two outgoing fission fragments. This is achieved using a dedicated ionization chamber (TWIM-MUSIC) and advanced tracking detectors. The success of this methodology has already been demonstrated by promising results, see [21],[22]

and figure 5, highlighting the potential of this technique in advancing our understanding of nuclear fission dynamics.

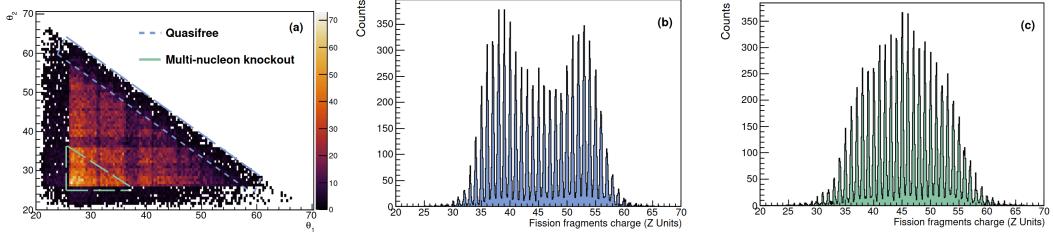


Figure 5: (a) Polar angle correlations of the two protons from fission via qfs event. In (b) the charge distribution for quasi-free ($p,2p$) fission reactions is selected while (c) displays the charge distribution for multi-nucleon knockout reactions of FSI. From Ref [21], S455 Experiment in 2021 at R3B, GSI.

3 Experiment

The present commissioning experiment was performed in 2020 at the FAIR Facility at GSI (Gesellschaft für Schwerionenforschung) in Darmstadt (Germany). The GSI Helmholtzzentrum für Schwerionenforschung operates a unique accelerator facility for heavy ions and focuses on several cutting-edge research fields. These include:

- 1. Nuclear Physics:** Studying the properties of atomic nuclei, exploring the forces that bind protons and neutrons, and investigating exotic nuclei far from stability.
- 2. Hadron and Quark Matter:** Investigating the behavior of hadrons (particles made of quarks) and the state of matter under extreme conditions, such as those found in neutron stars or during the early universe.
- 3. Atomic Physics:** Examining the structure and dynamics of atoms, including highly charged ions, to understand fundamental atomic interactions and refine quantum electrodynamics.
- 4. Plasma Physics:** Creating and analyzing high-energy-density plasmas to simulate conditions found in stellar interiors and other astrophysical phenomena.
- 5. Biophysics and Medical Research:** Exploring the effects of ion beams on biological systems for applications in cancer therapy, particularly using heavy ion therapy, and studying radiation protection for space missions.

6. **Materials Research:** Investigating the response of materials to high radiation doses to develop more resilient materials for use in various technologies, including nuclear reactors and space exploration.

3.1 GSI facility

The GSI Helmholtzzentrum für Schwerionenforschung located at Darmstadt has a long history of research.... tell something about the beginnnings, first really heavy elements found there.

The GSI Helmholtzzentrum für Schwerionenforschung GmbH was founded in 1969 (as "Gesellschaft für Schwerionenforschung mbH) looks back on a successful research history. In the time between 1981 and 2010 six new superheavy elements were discovered. In the medical research field GSI has developed advanced cancer therapy techniques using heavy ion beams which target tumors with high precision, minimizing damage to surrounding healthy tissues.

Along with those groundbreaking discoveries in research the facility at GSI has always been an inspiring source of drive for new technologies.

The key devices/apparatus which enable to carry out experiments with heavy ions at GSI are: important to mention: GSI is the only facility with heavy ions in the world. The starting point for the production of relativistic heavy ions at GSI is the ion source where ions are generated by stripping electrons off the shell of the atoms. Depending on the experimental needs the ion sources at GSI are able to produce ions of many different kinds of elements (up to Uranium).

On the first acceleration stage the stable primary ions are injected from the ion source into the UNIversal Linear Accelerator (UNILAC). On a length of 120 meters ions are accelerated up to maximum energy of 11.4 AMeV. The low energy beam is now injected into the ring accelerator SIS18 (Schwerionensynchotron 18). Here the ion beam is further accelerated up to 4.7 GeV/u (for protons) / 1 GeV/u (for Uranium). The magnets and the ultra-high vacuum ($\sim 10^{-9}$ Pa) keep the ions well on their circular path (SIS18 has a circumference of 216 meters). For the production of rare heavy isotopes the primary ion beam from SIS18 can be impinged on a light nuclear target, e.g beryllium, the so called production target. These secondary beams of radioactive isotopes can be either stored in the experimental storage ring (ESR) for later use or transferred to the FRagment Separator (FRS). The FRS as a high-resolution magnetic spectrometer is capable to precisely select specific isotopes and to forward the desired beam of exotic relativistic nuclei to the various experiments or direct it to the ESR for later use.

3.1.1 FAIR Project

The FAIR (Facility for Antiproton and Ion Research) situated next to the GSI will be one of the most complex and largest accelerator facilities in the world. The construction of the superconducting ring accelerator SIS100 with a circumference of 1.1 km, storage rings and experiment sites begun in the summer of 2017. Commissioning is planned in 2025 (?). Early Science. Before the commmissioning of the ring accelerator SIS100 several prioritized experiments with large impact in the scientific world will take place in the newly built experimental halls, such as experiments with the R3BSetup in the High Energy Cave (HEC).

3.2 R3B Setup

The R3B (Reactions with Relativistic Radioactive Beams) experiment in Cave C at the GSI Helmholtz Centre for Heavy Ion Research in Germany is a cutting-edge research experiment focused on the study of nuclear reactions and structure using high-energy radioactive ion beams. The experiment aims to investigate exotic nuclei far from stability, offering insights into the fundamental properties of nuclear matter, nucleosynthesis processes, and the forces governing nuclear interactions. A schematic overview of the R3B Setup can be seen in Figure blabla.

The short living (neutron rich) isotopes are injected to the Cave C from the FRS, which preselects as mass spectrometer the isotopes of interest, and impinge on a fixed target. The R3B setup is designed for kinematically complete reaction studies. To fulfill this requirement the incoming ions are tracked and identified on an event-by-event basis by dedicated detectors in the FRS via time-of-flight and deltaE measurement techniques. Depending on the settings and composition of the incoming ion beam different type of reactions take place in the target area with a large variety of reaction products: heavy ions (as producs from fission/spallation reactions), neutrons, light charged particles and gamma rays. For the detection of gammas and light charged ions from reactions with the target the dedicated CALIFA calorimeter (see more in section blabla) and various tracking detectors are installed in the target region. The GLAD (GSI Large Acceptance Dipole) magnet, located at the center of the Cave C, acts as mass spectrometer for the forward boosted charged reaction residues. The magnetic rigidity of the charged reaction residues is measured by a combination tracking detectors and a time of flight wall after the GLAD magnet. This allows to identify the charged reaction residues and their momenta. For the detection of the neutrons, not deflected by the magnetic field of the GLAD magnet, the new array neutron detector (NeuLAND) is positioned after

GLAD on the zero degree line with the incoming ion beam.

The high flexibility of the R3B Setup, it can be operated with The combination of the large spectrum of incoming ion beams in a broad energy range provided by the FRS facility and the high flexibility of the R3B Setup with state of the art detectors for the specific physics-studies of interest makes it to an attractive play-ground for experimental astro-physics.

3.3 Detector Setup in S444 Commissioning Experiment 2020

The S444 Experiment (successor experiment of the FAIR Phase-0 program in 2019, ref to Lukas Ponnath Thesis) for the commissioning of the CALIFA Calorimeter in its final mechanical design took place in February 2020. The choice to operate with stable ^{12}C primary beam with four beam energy settings - 400/550/650/800 AMeV gave the opportunity to use it as preparation for the following up S467 experimental run with neutron-rich Ca isotopes as medium-heavy incoming beam. The detectors for positional tracking, charge identification and time measurement were provided by the SOFIA(Study on Fission with Aladin, make footnote that ALADIN was the predecessor or GLAD) collaboration. These detectors are optimized for fission experiments with medium to heavy reaction fragments. As for the S444 experiment with primary ^{12}C incoming beam no fission reaction with multiple heavy charged fragments is expected the Sofia detectors were adapted accordingly (e.g. only one of the four sections of the Twin-Music Ionisation chamber was operated, see more in chapter Twin).

For this commissioning experiment most detectors and parts of the setup were operated in air. The target chamber was evacuated by gaseous helium at room temperature as well as the GLAD magnet. The fact that the ions interact with particles in air causes angular straggling in the flightpath reconstruction and can limit the resolution of reconstructed momenta from the reaction on the target.

3.3.1 Multi Wire Proportional Chambers (MWPC)

The positional tracking of the incoming ions as well as the charged reaction products were performed by using Multi Wire Proportional Chambers (MWPC). A MWPC operates on the principle of proportional counters that are arranged side by side in a plane, thereby providing spatial resolution for particle radiation. The multi wire proportional chambers were developed in late 1960s by George Charpak⁹ at CERN[23].

⁹George Charpak received the Nobel Prize in Physics in 1992 for his invention and development of particle detectors, in particular, the multiwire proportional chambers.

The MWPC operates in the same way as aligned proportional counters with the difference of not having dividing walls between the anode wires. This reduces the material budget, hence improving the spatial resolution and reducing reactions with the detected particle.

In the general design the MWCP is made up of a plane of anode wires enclosed between two cathode planes which are aligned parallel or vertical to the anode wires. Depending on the beam conditions the anode wires are set to high voltage (~ 1100 V) while the cathode planes are grounded.

The volume between the two cathode planes is filled by a gas mixture of 84% Argon and 16% CO₂. The decision of the gas mixture is driven by a balanced ratio between amplification and quenching properties of the gas.

When a charged particle passes through the detector it ionizes the gas. Primary electrons are created followed by a secondary ionization via electron avalanche. The electron avalanche drifts towards the wires (anodes) while the positive ions drift towards the grounded cathode planes. As the MWPCs are operated in the proportional region, the number of created electrons/ions is proportional to the initial ionization. Instead of reading out the signal from the wires it is read out from the strips of the cathode plane. This improves the position resolution in case the cathode planes are aligned perpendicular to the wires. In case multiple (neighboring) strips give signal the signal distribution over the strips is analyzed and fitted to provide the position information.

In the R3B setup for the S444 experiment four MWPCs were installed:

1. MWPC0: right at the beginning of the beam entrance in Cave C, 184 cm upstream to the target position to detect x- and y positions of the incoming ions.
2. MWPC1: 88 cm downstream to the target for positional tracking in x and y of the outgoing reaction fragment
3. MWPC2: 154 cm downstream also for positional tracking of the fragment
4. MWPC3: after the GLAD magnet. The x position of this detector gives the information about the magnetic rigidity of the reaction fragment.

Despite having the same mode of operation, they slightly differ in their construction design and positional resolution. For the technical specifications of the individual MWPCs, see table 1. Still to do: put in plot with potential field of mwpc and one with crosssign charged particle.

Common MWPC Settings	
Gas	84% Ar, 16% CO ₂
Windows	Mylar®
Anode wires voltage	1100 V
Cathode planes voltage	Ground
Wire pitch	2.5 mm
Wire diameter	5 μm
Width of X pads	3.125 mm

MWPC0	
X pads	64 pads, vertically segmented into two equal parts
Y pads	64 pads, horizontally segmented (3.125 mm width)
Active surface	200 × 200 mm ²

MWPC1 & MWPC2	
X pads	64 pads, vertically segmented into two equal parts
Y pads	40 pads (5 mm width), horizontally segmented
Active surface	200 × 200 mm ²

MWPC3	
X pads	288 pads
Y pads	120 pads (5 mm width)
Active surface	900 × 600 mm ²

Table 1: SOFIA MWPCs - Technical specifications

3.3.2 Ionisation Chambers - R3BMusic/TWIM Music

For the S444 experiment at R3B two types of multi sampling ionisation chambers (MUSICs) were installed: the R3B MUSIC, centered 153 cm upstream to the target, and the TWIN-MUSIC, 132 cm downstream to the target. Like the MWPCs (see 3.3.1) the ionisation chambers are gas-filled detectors for tracking down charged particles. While MWPCs consist only of a few mm of active gaseous volume, the ionsiation chambers have an expanded gaseous volume which allows to make precise energy loss measurements from the ionisation process of the gas. The multi sampling ionsiation chambers consist of a cathode plane and an anode plane, consisting of multiple anode strips. When a charged particle crosses the chamber the gas gets ionized and the created electrons and ions are separated by the strong electric field. While the ions drift

towards the cathode plane the electrons move to the anodes where each anode is read out separately. Since the energy loss of the passing through particle is proportional to the square of its charge ($\Delta E \sim Z^2$) the signal from the anodes allow to precisely measure the charge of the particle. Moreover multi-sampling ionisation chambers measure the drift time of the electrons created during the ionisation process on each anode (compared to one or more reference anodes). Assuming a constant electron drift velocity ($\sim 40\text{mm}/\mu\text{s}$) over the gaseous volume the time information of each anode signal can be used to reconstruct the x-position of the passing through particle).

R3B MUSIC

The R3B MUSIC, installed 153 cm upstream to the target, is used to measure both the charge of the incoming ion before impinging on the target and the angle of the particle's trajectory. The detector has an active gaseous dimension of $20 \times 20 \times 40 \text{ cm}^3$, confined on one side by a cathode plane and on the other side by an anode plane consisting of 10 anodes (8 readout anodes and 2 screen anodes). For the technical specifications, see table 2.

Dimensions

Detector dimension:	$51 \times 54 \times 53 \text{ cm}^3$
Active dimension:	$20 \times 20 \times 40 \text{ cm}^3$
Dimension of one anode:	$20 \times 20 \times 5 \text{ cm}^3$
Dimension of one screen anode:	$20 \times 20 \times 2 \text{ cm}^3$

Gas

P75 (Ar 25%, CH4 75%)

Voltage

Cathode (left to beam direction): $-(2 - 6)\text{kV}$

Anode (right to beam direction): $+300\text{V}$

Resoultions

still to do!

Table 2: R3B MUSIC - Technical specifications

TWIN MUSIC

The TWIN MUSIC is a double ionisation chamber with one central cathode plane and two independent drift volumes and anode planes on each side. Each of the anode planes consists of 16 anodes for readout plus two screen anodes. Furthermore each anode is again segmented into up/down which splits the detector into four dedicated sections. As the TWIN MUSIC is placed 132 cm downstream to the target it is employed to measure charge and angular direction of the outgoing medium-to-heavy fragments. The detector

was in particular designed for fission experiments where two or more fission fragments are created. If each fragment is flying through one of the four sections (which is mostly the case due to momentum conservation rules) charge and angle of each fragment can be measured independently.

To fulfill the required permanence of the field in both extended gaseous volumes (of dimension 11x22x40 cm³) a Frisch grid is located 3 mm from the anode planes. The Frisch grid is metal mesh grid that shields the anode from the movement of ions produced during ionization process in the chamber ensuring that only the electrons that reach the anode contribute to the signal. Additionally, the shielding of the anodes by the Frisch grid account for the fast rise time of the signal at the anodes which diminishes pile-up effects and makes the detector high beam-rate capable (up to 100kHz). Further technical specifications you can find in table (bla bla)

Dimensions

Detector dimension:	43 x48 x55 cm ³
Active dimension:	two halves each 11x22x40 cm ³
Distance central cathode - Frisch grid:	11 cm
Distance Frisch grid from anode planes:	3mm

Gas

CH4 [79%], Ar [20%] and CO2 [1%]

Voltage

Central cathode:	-(2 – 6)kV
Anode planes:	+600V
Frisch Grid:	+250V

Resoultions

$\Delta E/E$	< 5% FWHM, total < 2% FWHM
ΔX	< 40 μ m

Table 3: TWIN MUSIC - Technical specifications, see also [24]

3.3.3 Sofia Start Detector

The SOFIA Start detector is positioned right after the R3B Music ionisation chamber and gives a time reference for the incoming ion. It is a 1 mm thin scintillating plastic blade attached with a photo multiplier tube on each side. The scintillator light from excitation of the incoming ions produce a clear signal on both photomultiplier tubes

used for the time measurement:

$$t_{start} = 0.5 \cdot (t_{left} + t_{right})$$

To shield the plastic detector from daylight it is wrapped in mylar foil ($300\mu\text{m}$ thickness).

3.3.4 GLAD Magnet

The **GSI Large Acceptance zero degree superconducting Dipole magnet GLAD** sits in the center of the R3B Setup in the cave C hall⁶. With an adjustable field integral up to 5 Tm it has a high acceptance range in magnetic rigidity which is crucial for the identification of highly asymmetric reaction fragments. The homogeneous magnetic field in GLAD allows to achieve momentum resolutions $\Delta p / p$ of 10^{-3} in combination with the dedicated tracking system.

The large opening angle of ± 80 mrad makes the GLAD magnet highly transmissive for evaporated or scattered neutrons in the reaction process which will be subsequently detected in the NeuLAND detector.

The default bending angle of the beam with respect to the beam line was set to 18° . Herefore the currents where adjusted according to the beam energy:

1. 400 AMeV beam: 1444 Ampere
2. 550 AMeV beam: 1778 Ampere
3. 650 AMeV beam: 1957 Ampere
4. 800 AMeV beam: 2223 Ampere

3.3.5 CALIFA Calorimeter

The **CALOrimeter for In Flight detection of γ -rays and high energy charged pArticles**, CALIFA, is one of the main detector components of the R3B setup. It surrounds the target area and covers the full azimuthal range and a polar angular acceptance from 7° up to 140° in the target region. The calorimeter serves for the detection of gamma rays in the energy region $100 \text{ keV} \lesssim E_\gamma \lesssim 30 \text{ MeV}$ and light charged particles, mostly protons, up to $E_p \lesssim 700 \text{ MeV}$. To fulfill the demands requested by the different experimental campaigns an energy resoultion of $\frac{\Delta E}{E}(@1\text{MeV}) \sim 6\%$ in the gamma-ray energy regime $\frac{\Delta E}{E}(@100\text{MeV}) \sim 1\%$ in the proton range regime is required.

Geometry



Figure 6: Upstream view of GLAD magnet in the center of Cave C after installation in February 2016. Picture from [25]

The CALIFA detector is a highly segmented detector with more than 2500 CsI crystals installed in the final design. Since experiments in the R3B setup operate in relativistic kinematics both the incoming ions as well as the measured particles originating from reactions inside the source experience relativistic effects, more precisely the so called relativistic Doppler effect. The emitted gammas and protons are not isotropically distributed around the source region but are boosted in forward direction. Moreover, the energy measured in the lab frame is different from the kinetic energy in the rest frame of the incoming ion.

The relativistic Doppler effect has a huge impact on the geometric design and requirements of CALIFA. Therefore the detector was subdivided into three polar angle ranges⁷:

- $7^\circ \leq \theta \leq 19^\circ$ - CEPA (CALIFA Endcap Phoswich Array): The most forward segment consists of 96 CsI(Tl) crystals. Due to the aforementioned rel. Doppler effect this area will have the highest intensities and energies. For high beam energies most of the particles will not be stopped inside the crystal and will escape as "punch-throughs". Despite the "punch-through" ions deposit only a fraction of their kinetic energy (ΔE) in CALIFA it is possible to reconstruct the initial energy of the particle¹⁰. In CEPA crystals with a length of 15 cm are used and cover each a polar angle of $\approx 2^\circ$. This finer segmentation in polar angular range has the benefit to compensate for the high rate.

- $19^\circ \leq \theta \leq 43^\circ$ - Intristic Phoswich (iPhos): In conjunction with the CEPA, the

¹⁰This is done by exploiting the distinct scintillation components of CsI, see more in chapter 5 of [27]

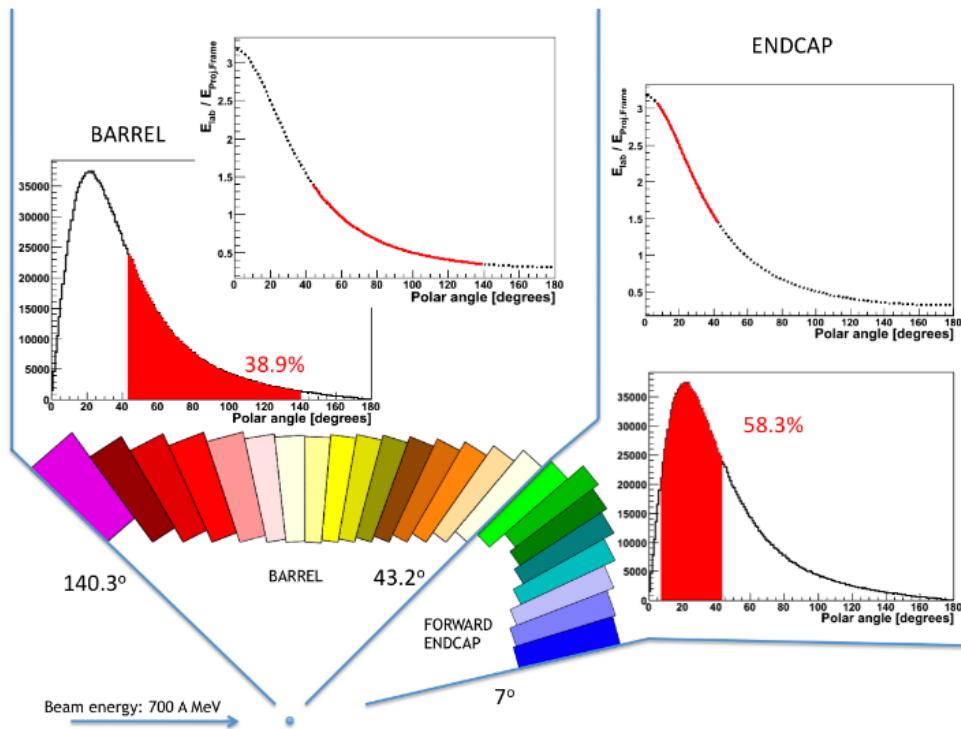


Figure 7: Schematic view of the Barrel and Endcap(iPhos and CEPA) segments of CALIFA and the according angular and energy distribution of emitted γ rays (monoenergetic in the projectile frame at beam energy of 700 AMeV). From [26]

iPhos region forms a part of the CALIFA Endcap. The iPhos region is, same as for the CEPA, affected by hight rates. Protons reaching this region have high kinetic energies ($E_{kin,p} \leq 600\text{MeV}$) and herefore a large fraction of "punch-throughs" are expected. In the iPhos region 480 CsI(Tl) crystals with a length of 22 cm¹¹ are installed and cover each a polar angle of $\approx 3^\circ$. (For more information see Endcap TDR:[28]).

- $43^\circ \leq \theta \leq 140^\circ$ - Barrel: This segment covers the region where the lowest rates and energies are expected. The Barrel region contains 1952 CsI(Tl) crystals, due to the large polar angular coverage and the requirement of Doppler correction via angle measurement. The most forward crystals have a length of 22 cm (which allow to stop protons with $E_{kin,p} \leq 315\text{MeV}$). This length is reduced down to 12 cm for the most backward crystals (For more information see Barrel TDR:[26]).

The crystals are arranged in carbon fibre alveoli with a nominal wall thickness of 230 μm [26] that provide a support structure for the crystals and keep the material bugged as low as possible. The alveoli in turn are hold and covered by individual aluminium tiles. The volume enclosed by the alveoli and the aluminium tiles is flooded with nitrogen to keep humidity low on the surface of the crystals. For a sufficient suspension of the aluminium cover a robust external holding structure was designed.

In 2019 CALIFA was for the first time integrated into the R3B setup in form of the CALIFA demonstrator, a prototype consisting of seven mechanically separate petals, each of it containing a set of 64 crystals.

At the end of 2019 the CALIFA frame in its final design was installed and the forward barrel part ($43^\circ \leq \theta \leq 90^\circ$ and full azimuthal coverage) was equipped with 1024 crystals.

For the S444 and the S467 experiment in 2020 CALIFA was equipped with 180 more crystals in the iPhos region ($19^\circ \leq \theta \leq 43^\circ$) which corresponds a coreage of 37.5 % in azimuthal angle for that region. Right before the S455 fissioning experiment [22] the full installation of the iPhos region was completed.

In February 2024 the full CEPA region ($7^\circ \leq \theta \leq 19^\circ$) with 96 crystals was commissioned for the first time together with a new equipped part of the backward barrel ($90^\circ \leq \theta \leq 102.5^\circ$, 128 crystals).

Energy and particle reconstruction with CsI scintillator crystals

Scintillator material, as caesium iodide doped with thallium (CsI(Tl)), is widely used

¹¹To fully stop protons with $E_{kin,p} \approx 600\text{AMeV}$ crystals with a lenght of 60 cm would be needed. Such long crystals would have multiple drawbacks: reduced energy resolution due to worse scintillator light transport, enhanced nuclear reactions inside the crystals and challenging demands on stability of the detector holding structure

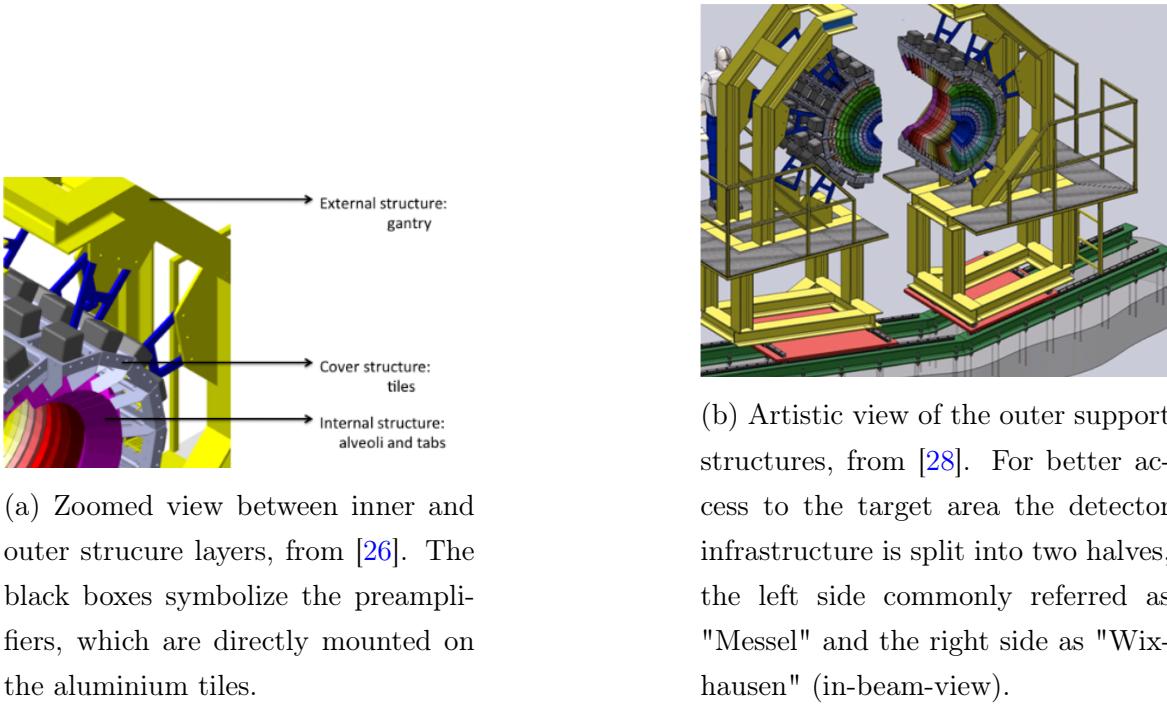


Figure 8: CALIFA internal and external holding structure

in experimental physics to detect ionizing radiation from γ - rays or charged particles. The energy from the incoming radiation excites the electrons in the CsI crystal from their ground state to higher energy levels (excited states). The deexcitation, followed by photon emission (scintillation light), occurs via various complex mechanisms, not all of them completely understood yet (for more information see TODO:references!). The efficiency and properties of the scintillation process can be fine-tuned by doping CsI with small amounts of other materials, such as thallium. Thallium-doped Cesium Iodide produces light with a peak emission around 550 nm (green light) and has high light output¹². The high density of CsI with 4.51 g/cm^3 makes it to an optimal scintillator material to efficiently absorb γ - rays and other high-energy particles. Moreover the CsI(Tl) crystal is well transparent to its own scintillation light, which is essential for the detection of the scintillator light. CsI(Tl) crystals are in addition relatively robust compared to other crystals and only slightly hygroscopic making them suitable for long-term use in experimental setups.

In a first approximation, the total amount of emitted light is proportional to the energy deposited in the scintillator. For γ -rays this is valid for $E_\gamma \gtrsim 400 \text{ keV}$ [30]. However, for charged particles significant deviations from linearity are observed, a so called "quenching" (TODO: add here more information, also the thesis of Max Winkel).

¹²The light output per MeV deposited energy in CsI(Tl), measured in [29], resulted in $5.2 \cdot 10^4$ (scintillation) photons/MeV.

Although the energy calibration of CsI(Tl) crystals for charged particles is challenging, CsI(Tl) as such has the beneficial property of having a complex time dependend light emission consisting of two distinct exponential components. The time dependend light emission response of CsI(Tl) $L(t)$ can be approximated as:

$$L(t) = \frac{N_f}{\tau_s} \exp(-\frac{t}{\tau_f}) + \frac{N_s}{\tau_s} \exp(-\frac{t}{\tau_s}) \quad (12)$$

Where N_f is the amplitude of the fast component and N_s the amplitude of the slow component. Accordingly τ_f the life time of the fast component ($\tau_f \approx 650 - 770\text{ns}$) and τ_s the lifetime of the slow component ($\tau_s \approx 3.2 - 3.5\mu\text{s}$). It has been found that the proportion between the two components is energy and particle dependend. This property can be used to identify isotopes by extracting the N_f and N_s values from pulse shape analysis (PSA) on the according light emission response¹³. TODO: put in a picture of nf vs ns an one with the scintillation light impulse

From scintillator light to electrical signal

The scintillator light produced by the ionisation processes has first to be transported to the back-end of the crystal. The optimum design has been determined to be frustum-shaped crystals, wrapped into enhanced specular reflector (ESR) foil which provides excellent reflectivity. Finally at the backside of the crystals a large area avalanche photodiode (LAAPD) is attached¹⁴. Avalanche photodiodes have the same working principle as photodiodes to convert (scintillator) light into electricity. As a result of an additional highly doted p-layer a region with very high field is formed which accounts for amplification factors up to ≈ 100 .

For the next amplification step the electric signal is forwarded via thin ribbon cables to the back end of the preamplifiers from Mesytec[34] which can serve up to 32 input channels. For CALIFA two general types of peramplifiers are in use:

1. Dual Range (DR) Preamplifiers: They are used in the iPhos and CEPA region where both high energetic protons as well as gammas are expected. They cover two amplification ranges in parallel: the *gamma range* with low input signal and high amplification and the *proton range* with high input signal with low amplification. Following from this they have 64 channel differential signal output.
2. Single Range (SR) Preamplifiers: In the Barrel region, where mostly γ rays are expected, only one amplification range is needed. Depending on the experimental demands these preamplifiers can be switched to *gamma* or *proton range*. These peamplifiers have a 32 channel differential signal output.

¹³The method has been implemented in the CALIFA Firmware as *QuickParticleIdentification – QPID*. For more information see [31] and [32]

¹⁴Detailed information about the crystal wrapping and LAAPD gluing can be found in this work:[33]

The fall time for the preamplifiers τ_{RC} has been chosen to $\approx 35 \mu\text{s}$. This is a trade-off between the ballistic deficit on one side (reduction of the signal amplitude due to low τ_{RC} , see also [31], chapter 3.4.5) and rate capability (restricted by large τ_{RC} value) on the other side.

The differential signal output of the preamplifiers is then transmitted over shielded and twisted line pairs to the input of the FEBEX Addon Boards (FAB) for further processing.

Signal Processing and readout system

The central hardware module for the signal processing in CALIFA is the FEBEX 3B Module (Front End Board with optical link EXtension[35]). Attached on it is a so-called AddOn board developed by TUM. The signal from the preamplifier gets here first filtered by a low pass two pole bessel filter ($f_c \approx 25\text{MHz}$). Furthermore, since the input of the FEBEX ADCs cover a range of $\pm 0.9 \text{ V}$ while the signal output from the preamplifier only has one polarity, an offset to the signal is applied to use the full range of the ADCs. The signal from the ADCs is read out continuously on the FEBEX card and split up into two branches:

1. Fast/trigger branch: After being fed to a trapezoidal filter the signal is examined by three leading edge discriminators with configurable thresholds. Depending on the experimental requirements a coincident matrix between one or more discriminators and optionally external triggers validates the signal as event ready for data recording.
2. Slow branch: The validated signal gets delayed over a ring buffer and a pulse shape analysis is performed via various steps - signal decimation, moving average unit -technique, baseline subtraction and moving window deconvolution (MWD) - to recall the major steps¹⁵. From the resulting pulse shape pulse height measurement the energy deposited in the scintillator is determined. In addition the algorithm for the quick particle identification (QPID) is applied on the incoming signal which provides the fast(N_f) and slow(N_s) component of the signal for isotope identification and differentiation of stopped and punch-through particles. The CALIFA Firmware also allows to make time over threshold (TOT) measurements which is convenient for energy reconstruction when the incoming signals exceed the ADC range (which might happen when the preamplifier is set to *gamma range*)¹⁶.

¹⁵A really detailed description of the pulse shape analysis in CALIFA can be found in Philipp Klenze's[36] and Max Winkel's thesis[32].

¹⁶The TOT energy-reconstruction method has the drawback of being really sensitive to pile-up events

Central hub for internal and external trigger forwarding are the Exploder modules[37]. The FEBEX crates are connected over an eight fold flat cable to the Exploder to the triggerbus. The trigger bus between the data acquisition PCs and the FEBEX cards is controlled by the TRIXOR card. This card sits inside the data acquisition PCs and is connected via ECL-lines to the Exploders and the PEXOR(PCIe Optical Link Interface) card via 26 fold flat cable. The PEXOR card is responsible for the data transfer between FEBEX cards and the data acquisition PCs. It is connected via glass fibre cables to the FEBEX crates and stores locally in a 576 MB large RLDRAM (Reduced Latency Dynamic Random Access Memory) the data which consecutively gets transmitted via DMA (Direct Memory Access - data transmission speed up to 560 MB/s) to the RAM of the data acquisition PCs.

Since the Exploder provides various input and output lines and an internal switchable bypass matrix the CALIFA calorimeter can be operated as free running system with internal event validation only or by (additional) external validation. In case of internal validation only the signal from the preamplifiers has to exceed predefined threshold(s) to be accepted. This kind of configuration can be used for the purpose of calibration (with γ -ray sources like ^{22}Na or ^{60}Co) and expected low event rate experiments. For high event rates and event coincidence with other subdetectors additional external validation (e.g. clean CFD signal from START detector) is implemented.

To overcome dead time initiated by the readout/data transmission procedure each FEBEX channel allows to store up to 254 recorded events on the local memory. Each FEBEX channel has two available memory banks. Whenever one FEBEX channel reaches its preconfigured maximum number of events the full FEBEX crate is read out. To avoid dead time all channels force a memory bank switch thus allowing continuous dead time free event recording.

The Multi Branch System (MBS), developed at the GSI Helmholtz Center for Heavy Ion Research, is used on the data acquisition PCs. This software consists of several components that control and read out detectors, store the data, or forward it via various network protocols. The system also allows for the joint readout of multiple systems. For this, the trigger modules (TRIXOR) are connected via a special trigger bus to exchange trigger signals and dead time information. The triggered data is then collected, time sorted and cross-detector events are built by the dedicated MBS event builder¹⁷.

overestimating the energy deposition. Hence more suitable for regions with low event rates, such as Barrel region.

¹⁷A more detailed explanation about the readout system and the critical FEBEX timing topic can be found in Philipp Klenze's thesis [36].



Figure 9: Sofia ToFW in Cave C, from [24].

Plastic	EJ-232, no quencher
Plastic dimensions	5x32x600 mm ³
Detector dimension	5x900x600 mm ³ (28 plastics)
Photo-multiplier tubes	Hamamatsu 6533 and 10580
Total number of PMTs	56 (two per plastic - top and bottom)

Table 4: Sofia ToFW - Technical specifications

3.3.6 Sofia Time of Flight Wall

The Sofia Time of Flight Wall (or "Stop detector") is positioned at the very end of the experiment setup at approximately 6.6 m distance from the target position. It consists of a plane of 28 vertically aligned scintillator bars, each of dimension 32x600x5 mm. The scintillator plastics are numbered from 0 to 28 from left to right (when looking in beam direction). The time of flight of the ions between Start and ToFW can be measured by subtracting the time measurement of the Start detector from the ToFW. The combined resolution of the Start and TOFW detector is at 40 ps, for an average time of flight of 30 ns[24]. For the technical specifications of the Sofia ToFW, see table 4 and reference [38].

3.3.7 NeuLAND Detector

For the detection of knocked-out or evaporated neutrons the **New Large-Area Neutron Detector (NeuLAND)** is installed at zero degrees after GLAD. In its final design it will consist of 30 double planes with 100 plastic scintillators of size 5x5x250 cm³ providing an active detector surface of 2.5x2.5 m² and thickness of 3m. Its high detection efficiency, a time resolution of $\sigma_t \leq 150\text{ps}$ and a mulit-neutron efficiency of 50% to 70% for four-neutron events are crucial detector features for complete kinematics experiments at R3B. For more detailed information, see [39] and [40].

For the S444 commissioning experiment in 2020 eight double-planes of the NeuLAND detector have been used.

3.3.8 Calibration of the Detector Systems

4 Analysis - Total Interaction Cross Section of $^{12}\text{C} + ^{12}\text{C}$

This chapter will go through the analysis step by step from the unpacking stage to the final measurement of the total interaction cross section. It will start by a short overview of the transmission method used for the cross section measurements. The next step is the selection of clean incoming ^{12}C isotopes. Following the identification of the carbon isotopes after the target - for the measurement of the charge changing cross section - and as final step the interaction cross section measurement.

All relevant detector related geometrical and efficiency corrections will be addressed and their influence to the final result and its uncertainty will be discussed.

4.1 Cross Section Measurement via Transmission Method

In its most generic form cross sections give a measure of the probability that a specific reaction will take place when two or more particles collide. The cross sections measured in scattering experiments, as well as the energy and angular distribution of the reaction products, provide information about the dynamics of the interaction between the projectile and the target particle, i.e., about the shape of the interaction potential and the coupling strength.

The cross section σ can be derived by looking at the relation between the number of incoming particles (N_1) and unreacted particles after the target (N_2). For an experiment with fixed target with thickness z and volumetric number density n the number of reacted particles in the infinitesimal thin target layer dz can be expressed as:

$$\frac{dN_2}{dz} = -n\sigma N_2 \quad (13)$$

Solving this differential equation for N_2 (with the condition $N_2 = N_1$ for $z = 0$) discloses an exponential relation:

$$N_2 = N_1 e^{-n\sigma z} = N_1 e^{-N_t \sigma} \quad (14)$$

Where $n \cdot z$ can be summarized as N_t , the total number of scattering centers per unit area. The relation (N_2/N_1) , number of unreacted particles after the target versus number of incoming particles, is often called survival probability P_{surv} . For an idealistic experimental setup with full detector efficiency and no interactions in the setup material the cross section could simply be deduced from equation 14. To account for reactions of the projectile that occur within the setup material and first order detector specific distortions of output signals the survival probability (N_2/N_1) has to be divided by the survival probability for an empty target run $P_{surv}^E = (N_2^E/N_1^E)$, where N_1^E is

the number of incoming particles and N_2^E the number of unreacted particles after the target for an empty run respectively. Thereby the setup specific efficiency(ϵ_{setup}) and transmission factor(t_{setup}) are cancelled out to obtain the underlying number of unreacted particles after the target \tilde{N}_2 :

$$N_2 = \tilde{N}_2 \cdot t_{setup} \cdot \epsilon_{setup}, \text{ with } \tilde{N}_2$$

$$N_2^E = \tilde{N}_2^E \cdot \epsilon_{setup} \text{ with } \frac{\tilde{N}_2^E}{\tilde{N}_1^E} \text{ the setup specific transmission factor } t_{setup}$$

The final formula for the cross section for a so called transmission measurement is:

$$\begin{aligned} \sigma &= -\frac{1}{N_t} \ln\left(\frac{N_1^E}{N_2^E} \cdot \frac{N_2}{N_1}\right) = -\frac{1}{N_t} \ln\left(\frac{N_1^E}{\tilde{N}_2^E \cdot \epsilon_{setup}} \cdot \frac{\tilde{N}_2 \cdot t_{setup} \cdot \epsilon_{setup}}{N_1}\right) \\ &\quad \text{With } \frac{\tilde{N}_2^E}{\tilde{N}_1^E} = t_{setup} \quad (15) \\ \sigma &= -\frac{1}{N_t} \ln\left(\frac{1}{t_{setup}} \cdot \frac{\tilde{N}_2 \cdot t_{setup}}{N_1}\right) = -\frac{1}{N_t} \ln\left(\frac{\tilde{N}_2}{N_1}\right) \end{aligned}$$

From the above formula 15 it is evident that for cross section measurements with the transmission method three types of observables have to be measured:

■ Number of scattering centers N_t

The number of scattering centers per unit area of the target is a target specific number. It depends from the target thickness and and its density. The values herefore are taken from [41]¹⁸:

(a) Thin target:

target thickness $d = 0.5451 \text{ cm}$; $N_t = 5.0588795 \cdot 10^{22}$; $\Delta N_t = 0.0648\%$

(b) Medium target:

target thickness $d = 1.0793 \text{ cm}$; $N_t = 1.0016600 \cdot 10^{23}$; $\Delta N_t = 0.2620\%$

(c) Thick target:

target thickness $d = 2.1928 \text{ cm}$; $N_t = 2.0350598 \cdot 10^{23}$; $\Delta N_t = 0.0322\%$

where N_t was calculated by:

$$N_t = \frac{\rho \cdot d \cdot N_A}{M} \quad (16)$$

with ρ the target density¹⁹, N_A the Avogadro constant ($6.02214076 \cdot 10^{23} \text{ mol}^{-1}$) and M the molar mass of the target (for carbon $M = 12.011 \text{ g} \cdot \text{mol}^{-1}$).

¹⁸For the purpose of this work the target thicknesses were remeasured at GSI with a chromatic sensor giving 2D depth profiles of each target.

¹⁹ $\rho = 1.851 \text{ g/cm}^3$, from [41]

■ Number of incoming projectiles (^{12}C) N_1

For the measurement only events with well identified incoming ^{12}C projectiles are chosen. Therefore strict cuts on the detectors upstream the target area are set. This strict event selection makes sure that we only consider events with single ^{12}C . This will be discussed in more detail in section 4.2.

■ Number of unreacted projectiles (^{12}C) N_2 after the target

Detectors downstream the target area are used to count the number of unreacted projectiles ^{12}C . To reduce detector specific influences which could distort the result it is advisable to use only as few as required detectors for the clear identification of unreacted projectiles. Moreover detector specific efficiencies are cancelled out by including both empty and target runs in the cross section calculation(see equation15). For all downstream detectors used in this analysis it is critical to minimize any selection cuts and systematically check their effects on N_2 .

4.2 Event Selection

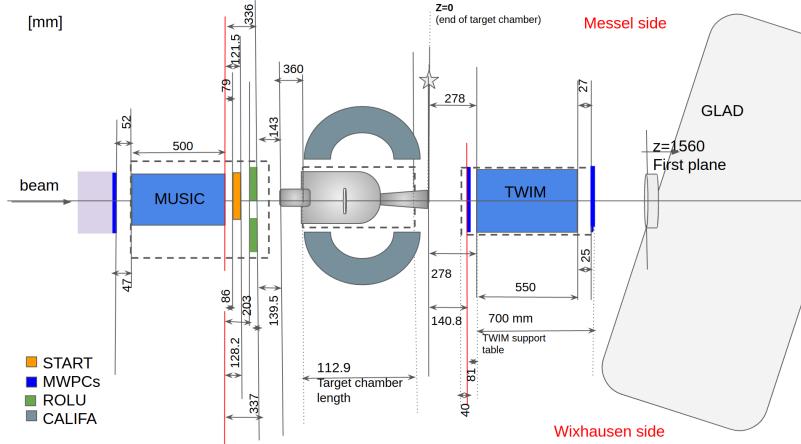


Figure 10: R3B Setup for the S444 experiment in the target region. TODO: select overview with less numbers/measures...

For event selection, all three upstream detectors are utilized: the MWPC0, the R3BMUSIC Ionization Chamber, and the start detector. To ensure a clean incoming event selection, the following prerequisites must be met:

1. ^{12}C identification of incoming projectile by upstream detectors:

In the S444 experiment the incoming beam was directly delivered by the SIS18 ring accelerator, which is operated in ultra-high vacuum. The level of contamination is low.(TODO: up to which stage do we have vacuum? Until MW0?)

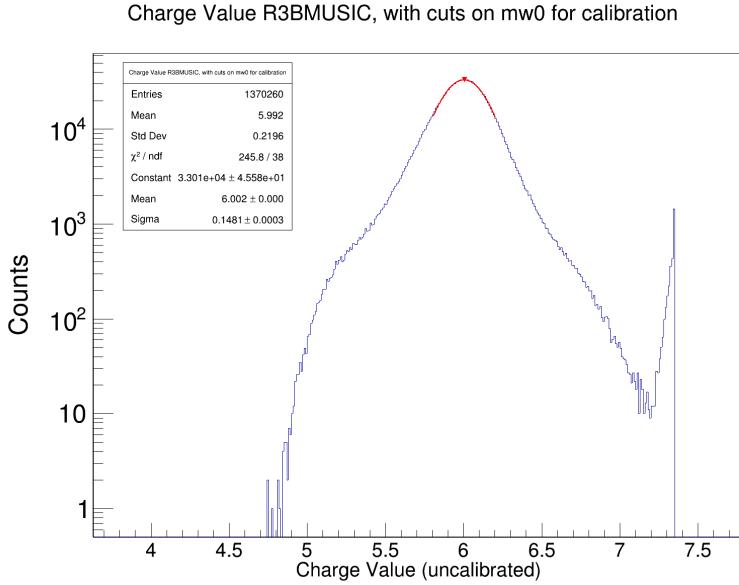


Figure 11: Charge distribution on R3BMUSIC with predefined calibration parameters with already applied positional cuts on MWPC0 - positioned upstream to the ionisation chamber. The rise beyond $Z \geq 7.2$ comes from pile-up events. TODO: more explanation needed? 2D plot needed?

For the charge identification of the incoming ion the R3BMUSIC ionisation chamber is used which is positioned directly after the MWPC0 at the beam entrance in Cave C, see figure 10. The R3BMUSIC detector measures anode-wise the energy loss of the passing-through ion which in the first order is proportional to the square of its charge ($\Delta E \sim Z^2$). Therefore the calibration parameters from the online analysis are used²⁰. Figure 11 shows the measured charge distribution in R3BMUSIC. To select $Z = 6$ incoming ions the distribution is fitted with a gaussian fit function. All ions with charge within the $\pm 1\sigma$ range are accepted. Figure 12 summarizes the $\pm 1\sigma$ cuts on the R3BMUSIC charge for empty/target runs for all beam energies.

2. Pileup rejection and TPat selection:

The overall recoding and merging of the data from various subdetectors is one of the tasks of the Data AcQuisition (DAQ) system. Whether an event is recorded or not depends on the pre-established trigger logic. Various detectors can send out triggers to the main DAQ when certain conditions are given (e.g. CALIFA can be configured to send out a trigger when a hit with more than 20 MeV is recorded in the calorimeter). The different triggers are processed by the trigger logic and summarized as a defined trigger pattern, so called TPat, which is stored

²⁰These are generic parameter values used to the detector performance during the experiment phase.

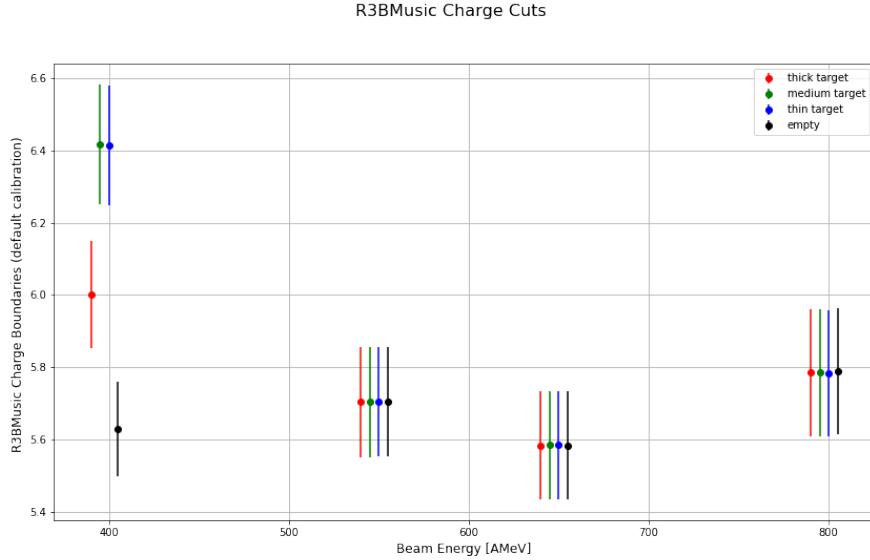


Figure 12: Strict $\pm 1\sigma$ charge cuts with R3BMUSIC for incoming particle selection. Fixed predefined calibration parameters were used which do not compensate different gain settings between runs. This is in particular the case for the 400 AMeV beam energy runs.

in a 16-bit mask for each event. Table 5 gives an overview of the trigger logic and the trigger patterns set in the S444 experiment. For this analysis the "*Min. Bias*" trigger is required²¹.

Since the TPat selection itself does not necessary set any pileup constraints it

Bit Position	TPat Name	Description
0	Min Bias	Hit in Start detector
1	Reaction	"CalifaOR" -high energy hit in CALIFA
2	Neutron	Hit in NeuLAND
3	p+n	Hit in CALIFA and Neuland
8	Califa	high energy hit in califa - off-spill
9	NeuLAND	Hit in NeuLAND - off-spill

Table 5: List of TPats set for S444 experiment. As for the selected runs low beam rates ($< 10kHz$) were expected no dead time issues should arise for the in-beam detectors, therefore no downscaling of the *Min. Bias* TPat was deployed.

is important to analyse the signals of the detectors upstream carefully to insure

²¹This includes also "*Reaction*" and "*Neutron*" TPat since these patterns contain also "*Min. Bias*" TPat as necessary condition.

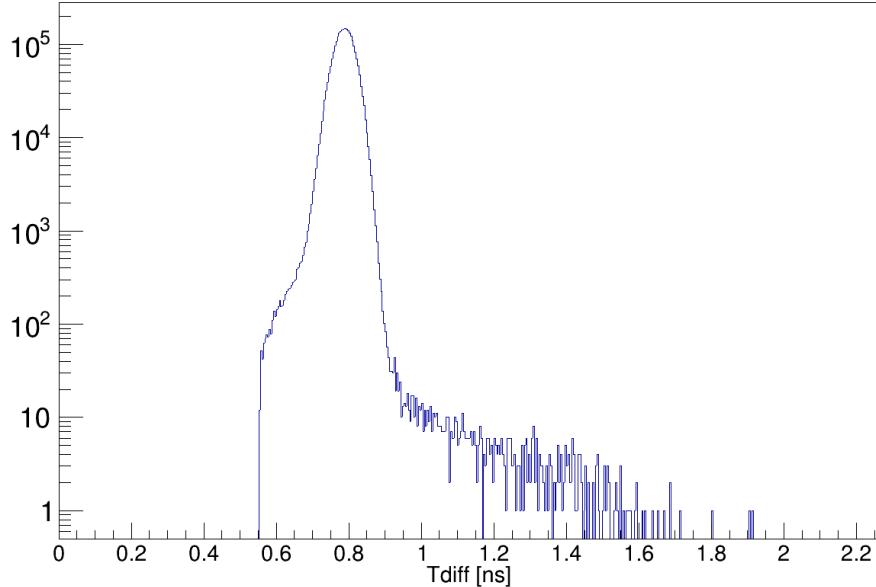


Figure 13: $\Delta t_{right-left}$ between hits in the Start detector for events with exactly one hit on the left and right preamplifier and limiting the time difference in the range 0.555 ns to 1.946ns.

yourself that only events with one incoming ^{12}C ion at a time get selected. Therefore events with incoming ions with charge $Z = 6 \pm 1\sigma$ are chosen, as discussed in the previous point. Moreover it is required that both left and right preamplifiers of the start detector have seen a coincident signal within a time-window of 1.391 ns.—TODO: why this time window?— The overall searching window of the start detector was set to 2 μs , see figure 13. For the MWPC0 which is mounted right at the beam entrance of Cave C no hit multiplicity cuts were applied considering its operating mode, which is designed for charge sharing between the readout pads.

3. Projectile's focus on the active target region:

To assure that the incoming ^{12}C ion hits the target it is necessary to select only events where the projectile is focussed to the active target region. Therefore strict cuts on the MWPC0 x and y position are applied. This was achieved by fitting the x and y distribution of the MWPC0 (without any restrictions on it) by a gaussian function. The selection of focused incoming projectiles was then restricted to events with hits in MWPC0 within the $\pm 1\sigma$ region in the x and y position, see figure 14 and 15.

The MWPC0 x-position and the available projectile angle in the x-y plane from the R3BMUSIC is used to propagate the corresponding x-position on the target location to further check that the selected projectiles hit the target parallel to the z-position (= beam direction) and do only have a minimal incident angle, see

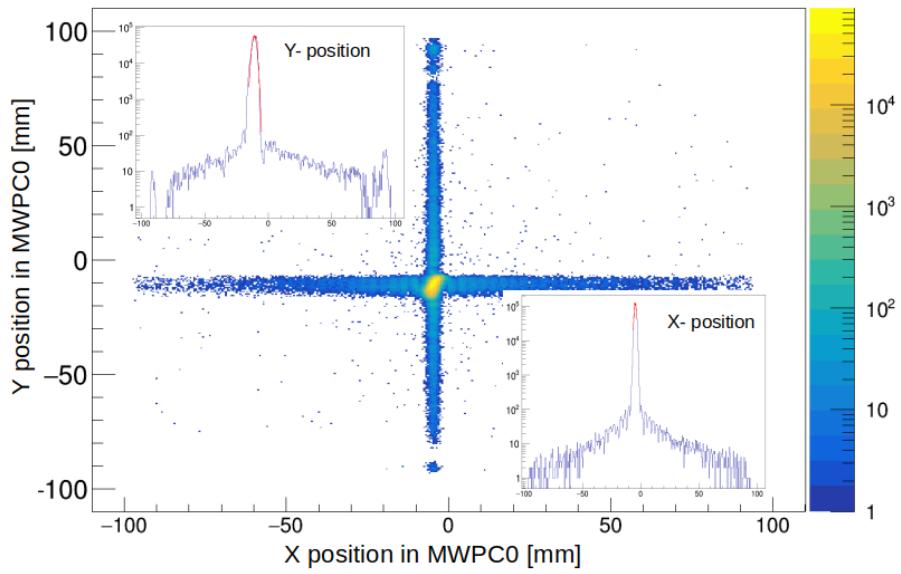


Figure 14: x-y position of incoming ion on MWPC0.

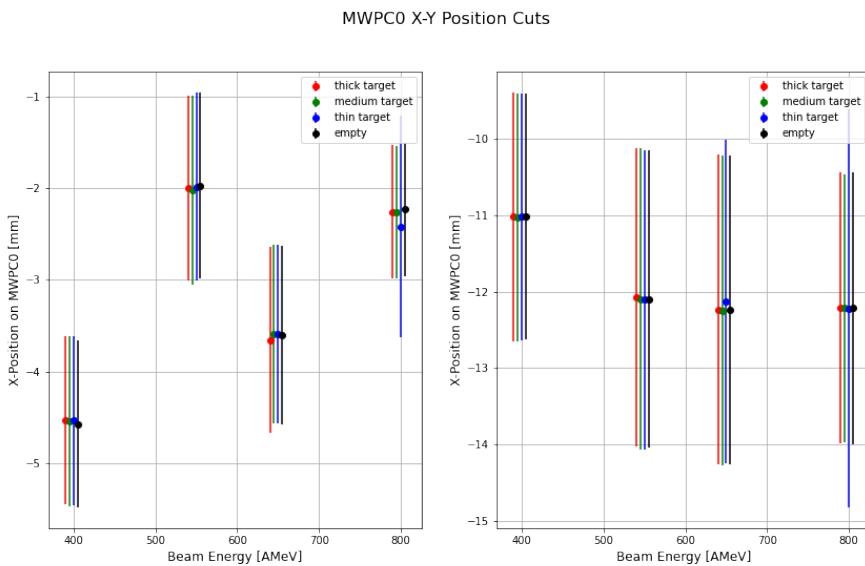


Figure 15: Overview of $\pm 1\sigma$ cuts in x and y in MWPC0 for empty/target runs. TODO: labelling on the right side is wrong! (should be y!)

figure 16.

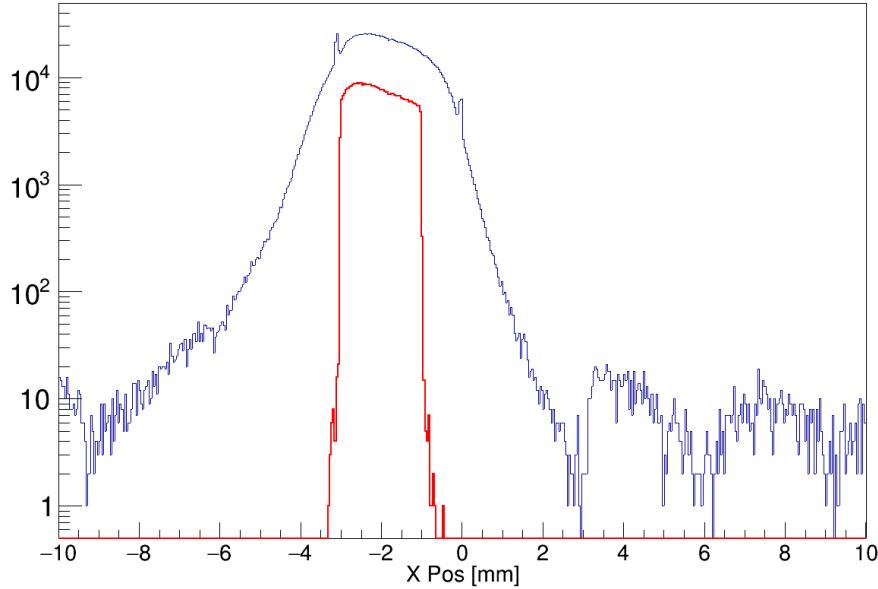


Figure 16: Propagated x-position on target location from measured x-value on MWPC0 and x-y plane angle measurement from R3BMUSIC. The target area is 3 x 3 cm. In red the selected events with $\pm 1\sigma$ cut in x and y position in MWPC0, in blue all events. TODO: which run is this?

4.3 Charge Changing Cross Section Measurement

The charge changing cross section refers to a measure of the probability that the incoming projectile will undergo a reaction inside the target that changes its charge. To measure the charge changing cross section it can be referred to formula 15 where in this case N_2 is the number of survived carbon isotopes, i.e. projectiles which did not change their charge state. For this measurement only the data from the double ionisation chamber TWIN Music (see section 3.3.2) needs to be read out and analyzed.

While for the event selection before the target the cut conditions can be arbitrarily strict (it will only have an impact to the statistics and the derived statistical error), cuts on the downstream detectors need to be avoided if at all possible. Too selective cuts on the identification of N_2 can distort the measurement.

4.3.1 TWIN MUSIC Calibration

For the analysis of data in TWIN MUSIC - different to the upstream detectors, where calibrated data with default calibration parameters is used - the so called *mapped* raw

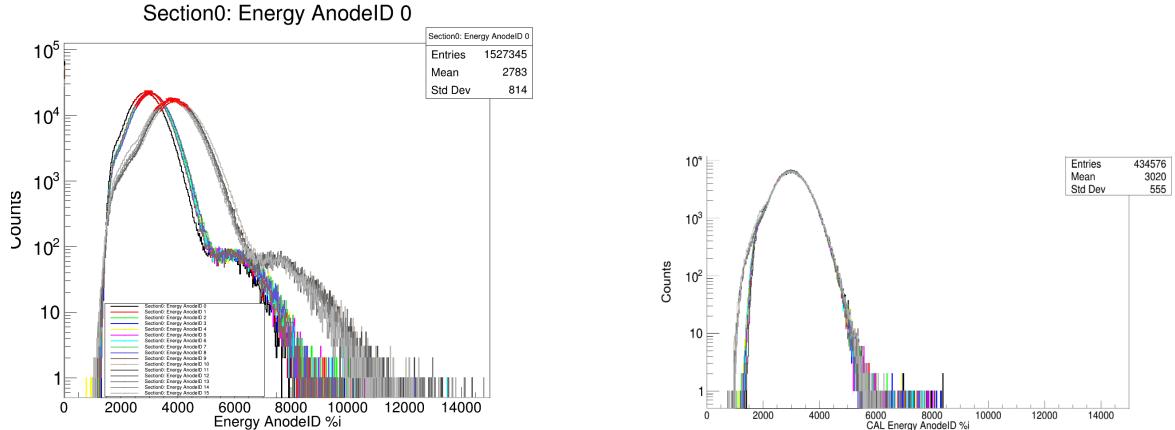
level data is processed. In the mapped level TWIN MUSIC provides following information:

- **SectionID:** The detector is a double ionisation chamber and as such divided into four parts (in beam perspective): section 1 - right down; section 2 - right up; section 3 - left down; section 4 - left up. For the S444 experiment only section 1 was operated and accordingly centered on the beam spot.
- **AnodeID:** TWIN MUSIC has 16 anodes for energy-loss readout and one reference anode (anodeID = 17).
- **Time:** Each hit in each anode gets assigned to a time. Each time has individually no meaning. The drift time (in ns) of the electrons from the ionisation process of the gas by the inflying projectile (or the fragments of it) to the anode is calculated by subtracting the individual anode time by the time of the reference anode. The reference anode receives its clean signal from a constant fraction discriminator of the start detector.
- **Energy:** Each hit in each anode gets assigned to an energy - except the hit in the reference anode. To reconstruct the charge of the crossing through charged particle anodewise or detectorwise the parametrization formula $Z = [0] + [1]*\sqrt{E} + [2]*E$ is used.

The calibration of the TWIN energy for each anode is done run-wise. For the TWIN MUSIC only events where all anodes having exactly one hit (including the reference anode) are chosen. The most prominent peak ($Z = 6$) was fitted with gaussian function. The calibration was then done by determining the scaling factor for each anode that shifts the mean of the gaussian fits to the same position, see figure 17. For this analysis the peaks were shifted to $\Delta E = 6$. Since $\Delta E \propto Z^2$ holds, the scaled ΔE value, even though peaking at 6, is not equate the charge $Z = 6$. TODO: describe better the calibration steps!!! this is not good!

4.3.2 TWIN MUSIC Event Selection

As previously stated cuts on the downstream detectors are avoided. However, events which have hits in one or several anodes in TWIN MUSIC but no signal in the reference anode are discarded as a whole neither contributing to N_1 (incoming selected ions) nor to N_2 (unreacted ions). If no reference time from start CFD signal is available it is not possible to measure the drift time in the individual anodes which makes it not possible to distinguish between signal and noise hits for multi-hit anode events in TWIN MUSIC. The number of events affected by this cut is in the region of few tens. This is negligible



(a) Uncalibrated raw ΔE distributions for all 16 anodes for the thick target run, 550 AMeV beam energy. The last six anodes have a slightly different electronics amplification chain.

(b) Gaussian fit applied to prominent peak and shifted to same position.

Figure 17: Fitting procedure in TWIN.TODO: nicer labeling!

to the number of incoming ions N_1 and should not have any dependence whether the projectile reacted or not.

4.3.3 Carbon Identification

The identification of carbon isotopes in TWIN is done by reconstructing fragments with charge $Z = 6$ from 2D plots where coincident mean energy losses ΔE for different anode combinations are plotted. Since the TWIN MUSIC is multi-hit capable various strategies were developed to deal with multi-hit events, i.e. when having anodes with multiple hits, decide which hit originates from the final state products from the reaction and which from background and noise.

The default strategy is to use the time information of each hit for selection. It has to be remarked that for the S444 experiment the TWIN MUSIC was read out by two independent MDPP modules[42]. The signals from the first reference anode and the first eight upstream anodes were sent to module 1, the ones from the last eight downstream anodes and the second reference anode were forwarded to module 2. For the first eight upstream anodes the drift time is calculated by subtracting the hit time in each anode by the reference time from the first reference anode and for the last eight downstream anodes accordingly the second reference anode was used.

The time based selection algorithm for multi-hit anodes works as follows:

# incoming projectiles N_1	400 MeV/nucleon	550 MeV/nucleon	650 MeV/nucleon	800 MeV/nucleon
Empty	574279(*451*)	453729(*34*)	522451(*44*)	395451(*52*)
thin	569503(*422*)	476323(*33*)	538037(*43*)	481459(*36*)
medium	606578(*431*)	451137(*27*)	500688(*40*)	345654(*46*)
thick	655762(*497*)	436457(*30*)	530869(*29*)	479679(*61*)

(a) Number of clean selected incoming ^{12}C ions. In brackets number of rejected events because of missing tref in TWIN MUSIC. TODO: change the bracket notation, looks like error number!!

# survived carbon isotopes N_2	400 MeV/nucleon	550 MeV/nucleon	650 MeV/nucleon	800 MeV/nucleon
Empty	563382(1.898%)	444618(2.008%)	511923(2.015%)	387513(2.007%)
thin	538245(5.489%)	449422(5.648%)	507557(5.665%)	454099(5.683%)
medium	552763(8.872%)	410376(9.035%)	455159(9.093%)	314119(9.123%)
thick	553935(15.528%)	368004(15.684%)	446115(15.965%)	402696(16.049%)

(b) Number of survived carbon isotopes after the target identified via 2D gaussian fit with borders within 3.5σ cut. In brackets the percentage of projectiles with a charge state of $Z < 6$ after the target.

Table 6: Numbers of incoming projectiles N_1 and survived carbon isotopes N_2 for all energy and target runs.

1. Get the mean drift time for the eight upstream anodes(t_{mean_up}) and the eight downstream anodes(t_{mean_down})²². Anodes with multiple hits do not contribute to this calculation.
2. If there are anodes with muliple hits compare the hit time with the accorging mean drift time (t_{mean_up} for any of the eight upstream anodes, t_{mean_down} for any of the eight downstream anodes). Calculate herefore the absolute difference between mean drift time and each hit time:

$$\Delta t = |\bar{t} - t_{drift}^i|; i = \text{anodeID } (1-16) \text{ with } \bar{t} = \begin{cases} t_{mean_up} & \text{for } i \leq 8 \\ t_{mean_down} & \text{for } i \geq 9 \end{cases} \quad (17)$$

3. For each anodes with muliple hits select the hit with lowest drift time differnce to the mean drift time.

²²For the case all eight downstream anodes have multiple hits, set $t_{mean_down} = t_{mean_up}$ and vice versa

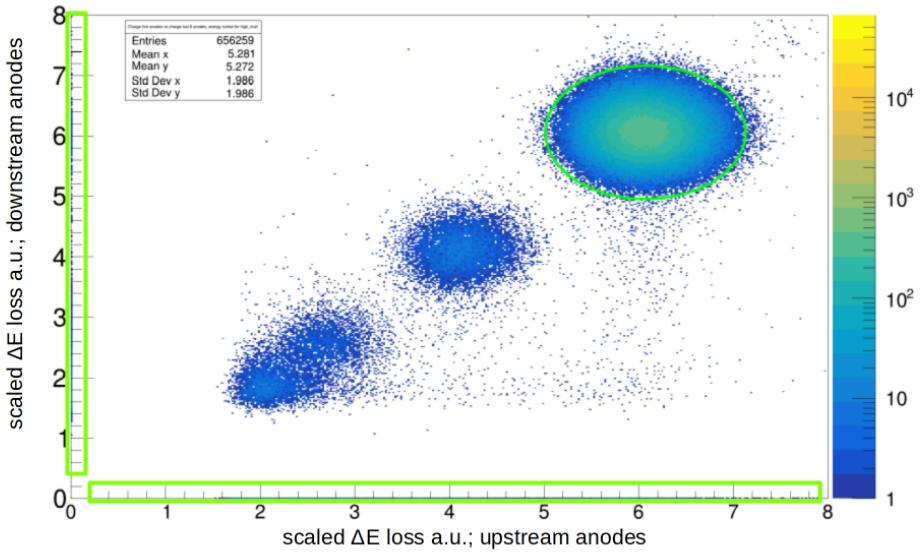


Figure 18: Two dimensional gaussian fit with 3.5σ cut on identified carbon isotopes in TWIN MUSIC. The horizontal and vertical side bars contain events where either the eight upstream anodes or downstream anodes have no hit entry. The cluster at scaled ΔE loss ≈ 4.5 corresponds to boron isotopes ($Z=5$). The clusters for $Z=4$ (Be) and $Z=3$ (Li) overlap.

After having selected the appropriate hit for single and multi-hit anodes the mean value for the pre-calibrated ΔE loss, see figure 17b, for the eight upstream and accordingly for the eight downstream anodes is determined. Finally, to select the number of survived carbon isotopes the mean ΔE of the eight upstream anodes versus the mean ΔE of the eight downstream anodes is plotted. To retrieve the number of survived carbon isotopes following two-dimensional gaussian fit is applied on the 2D plot on the charge $Z = 6$ blob, see figure 18:

$$f(x) = Ae^{-\frac{1}{2}((\frac{x-\bar{x}}{\sigma_x})^2 + (\frac{y-\bar{y}}{\sigma_y})^2)} \quad (18)$$

where x is the mean rescaled energy loss of the first upstream anodes and y the according eight downstream anodes. The number of survived carbon isotopes is given by the integral of events within the 2D gaussian fit. Since the anodes were read out by two independent MDPP modules with slightly different thresholds also events along the histogram axes with no hit entry in either the upstream anodes or downstream anodes are analyzed. For those events a one dimensional gaussian cut is applied using the parameters from equation 18 (see horizontal and vertical bars in figure 18).

To get the charge changing cross section values equation 15 has to be applied where both the number of survived carbon isotopes for target run and empty run are determined via the 2D gaussian fit as in figure 18. The number of target particles, are defined by

the target thickness and its density and are listed in section 4.1. The resulting charge

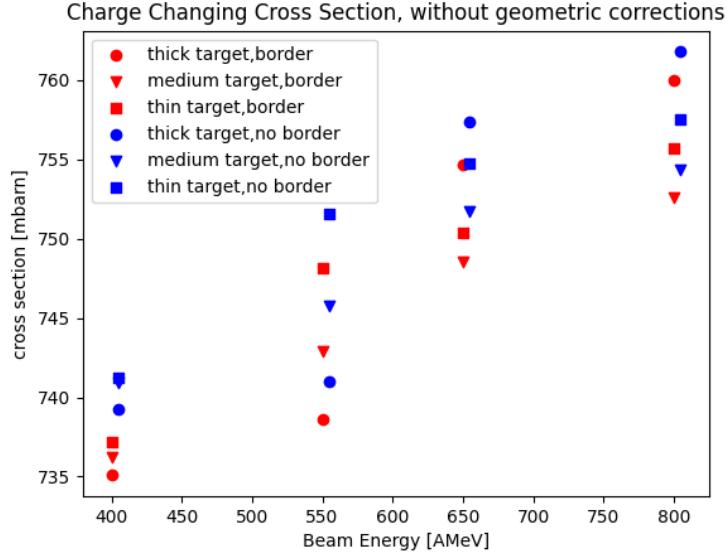


Figure 19: Charge changing cross section without geometry corrections. The red data points result from considering also events with only hits in the upstream or downstream anodes, the blue data points don't take these events in consideration.

changing cross sections are summarized in figure 19 once with consideration of the vertical/horizontal bars in figure 18 and once without. Hits within the 3.5σ gaussian fit are identified as carbon isotopes. To get the optimal σ cut on the two dimensional gaussian fit on the energy losses of the upstream anodes versus downstream anodes the charge changing cross section for all targets and all energies was systematically measured for σ -cuts in the range of 1 to 5 σ , see figure 20. In the region $\sim 3.5\sigma$ the variation of the cross section is minimal. Another method to assert the number survived

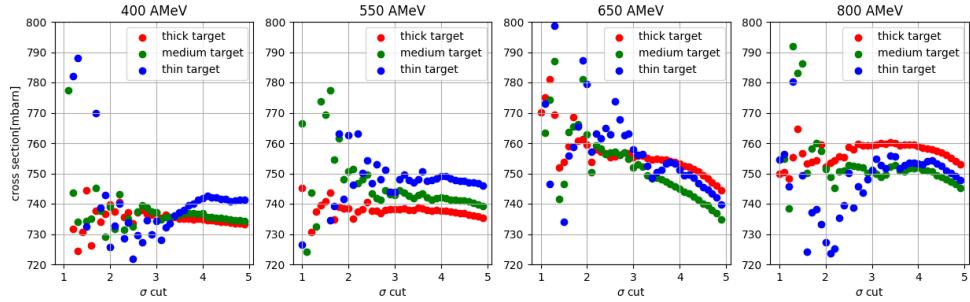


Figure 20: Measured charge changing cross sections according to the σ cut applied on the figure 18 (with borders) for the different target thicknesses and beam energies.

carbon ions is to apply a diagonal cut on the 2D ΔE histogram. To set the slope and offset of the diagonal cut line firstly the two dimensional gaussian fit is applied, same

as for the previous method. Then the intersection point between the 3.5σ ellipse and the identity line (ΔE upstream anodes = ΔE downstream anodes) is found. Through this point, perpendicular to the identity line, the diagonal line is drawn. Everything above the diagonal line is considered as survived carbon ions. Moreover the borders are considered within the 3.5σ cut, see figure 21. The effects of the two different methods

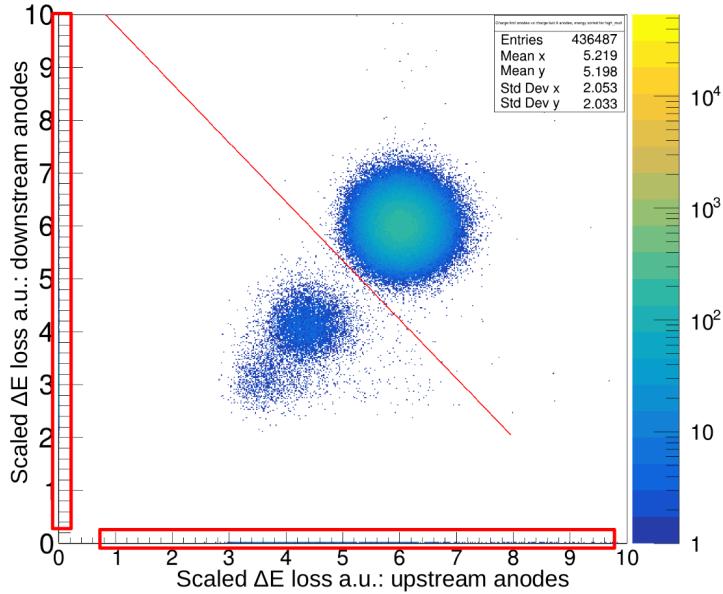


Figure 21: Diagonal cut on identified carbon isotopes along the gaussian 3.5σ cut with borders. All hits above the diagonal line are counted as carbon isotopes. Histogram from thick target run, 550 AMeV beam energy.

used for the identification of the carbon isotopes for the charge changing cross section is summarized in figure 22. The differences in the measured cross sections are within the margin of error herefore both methods are comparable, as expected. To check wether single anodes or groups of anodes are malfunctioning the charge changing cross section measurement was repeated using only certain anodes for the charge identification:

- a) anodes 2-8 versus anodes 9-15 (omitting first and last anode)
- b) anodes 1-4 versus anodes 5-8 (upstream anodes)
- c) anodes 5-8 versus anodes 9-12 (central anodes)
- d) anodes 9-12 versus anodes 13-16 (downstream anodes)

The results from the measurement are summarized in figure 23. The difference between the default gaussian fit method (with 3.5σ cut and considering the borders) considering all 16 anodes and applying the same method but omitting the first and last anode is minimal over all four beam energies. When selecting only 8 out of 16 anodes instead the

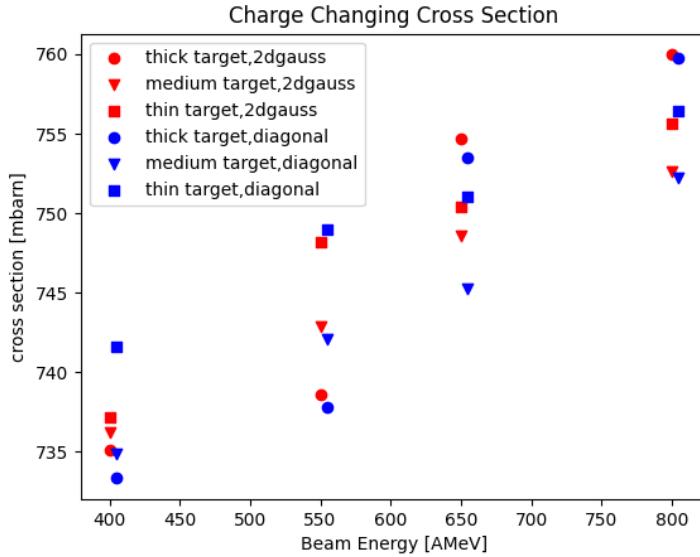


Figure 22: Comparison of charge changing cross section measured via 2D gaussian fit cut and diagonal cut. The differnces are within the margin of error.

cross sections are systematically lower when going to high beam energies. The energy loss inside the TWIN MUSIC decreases with higher beam intensities, according to the Bethe-Bloch formula:

$$-\frac{dE}{dx} = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left(\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right) \quad (19)$$

where:

$$K = 4\pi N_A r_e^2 m_e c^2 \approx 0.307 \text{ MeV cm}^2 \text{ g}^{-1},$$

z = charge of the incident particle (in elementary charge units),

Z = atomic number of the target material,

A = atomic mass of the target material,

$\beta = \frac{v}{c}$ = velocity of the particle relative to the speed of light,

$\gamma = \frac{1}{\sqrt{1 - \beta^2}}$ = Lorentz factor,

T_{\max} = maximum kinetic energy transferable to an electron in one collision,

I = mean excitation potential of the target material,

$\delta(\beta\gamma)$ = density effect correction.

The behaviour of dE/dx for small β - values are dominated by the $1/\beta^2$ term. The decrease of deposited energy for larger beam energies has as consequence a lower relative resolution in the two dimensional ΔE loss histogram (see figure 18) reflecting the

poissonian distribution properties. In addition reducing the number of readout anodes by a factor two degrades the resolution by a factor $\sqrt{2}$ ²³. This has as consequence that the ellipsis with 3.5σ cut incorporates a non negligible amount of boron isotopes which are counted as survived carbon isotopes which in turn reduces the measured charge chaning cross section.

While in the above measurements a time based secelction algorithm for multi-hit an-

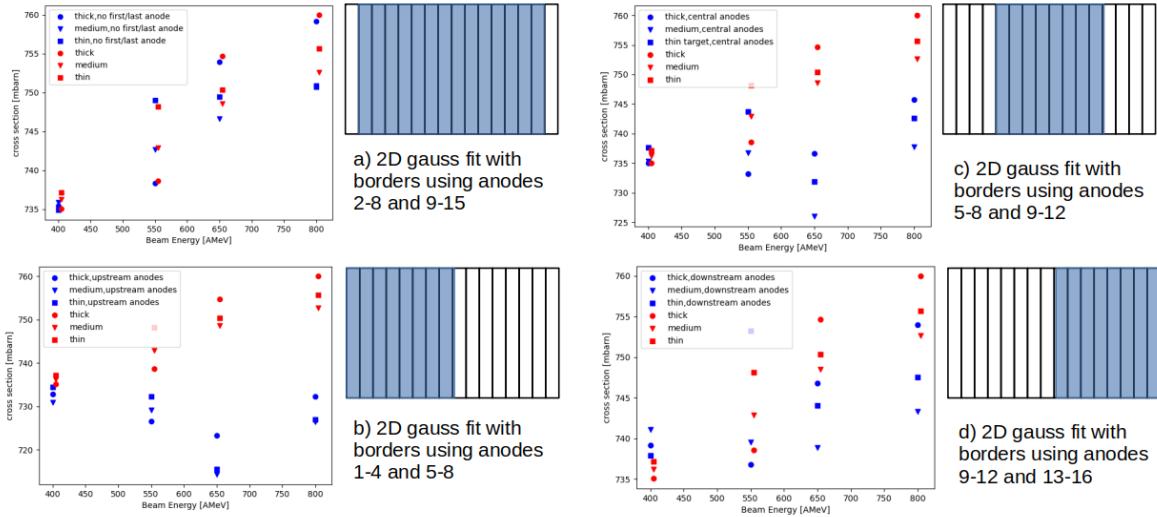


Figure 23: Measurement of charge changing cross sections using different anode sections to make the two dimensional gaussian fit on the identified carbon isotopes. Red: using all 16 anodes. Blue: the various combinations.

odes was used also an energy based selection was tested. This algorithm selects for multi-hit anodes the hit with the highest energy as physical hit and discards all others, as they are considered as background/noise. Figure 24 compares the time based method versus the energy based method. In both cases a two dimensional gaussian 3.5σ cut is applied as in figure 18 and the borders are counted as well. The difference in the outcome is negligible. This can be explained since noise or background signal should be both uncorrelated to the event time and at a low energy level and are therefore filtered by both algorithms. The final charge changing cross section measurements with 2D gaussian fit applying a 3.5σ cut and including the borders of the histogram are summarized in figure 25. At this stage also the statistical errors are incorporated.

²³It can be assumed a similar ΔE distribution for all anodes. Hence the central limit theorem can be applied where $\sigma = \frac{\sigma_{anode}}{\sqrt{n}}$ with $n = \text{number of anodes}$.

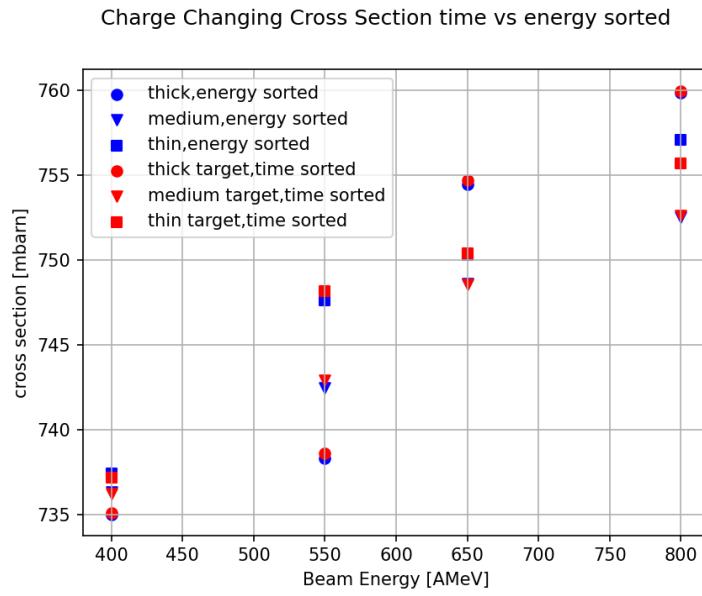


Figure 24: Comparison of charge changing cross section measurements when using time sorting algorithm(red) and energy sorting algorithm(blue) for multi-hit anodes.



Figure 25: Measurement of charge changing cross sections using all 16 anodes of the TWIN MUSIC applying the 2D gaussian fit and considering the borders as in figure 18.

4.4 Geometric Corrections

For the S444 experiment only section 1 (right down) of the TWIN MUSIC, which was centered on the beam spot, was operated. As consequence full geometric efficiency could not be assumed. To visualize the restricted geometric efficiency of the TWIN MUSIC the position in x and y (perpendicular plane to the beam direction) on the MWPC1 in front of the ionisation chamber was plotted, once without any conditions on the TWIN MUSIC and once with the condition of having identified a carbon isotope (with the 2D gauss-fit method as described in chapter 4.3.3), see figure 26. The large active surface area of $200 \times 200 \text{ mm}^2$ of the MWPC1 affirms that all the carbon fragments are detected²⁴ whereas the TWIN MUSIC behind it, with an active surface of $55 \times 110 \text{ mm}^2$ (section 1), is not sensitive to the scattered fragments with large deflection angle.

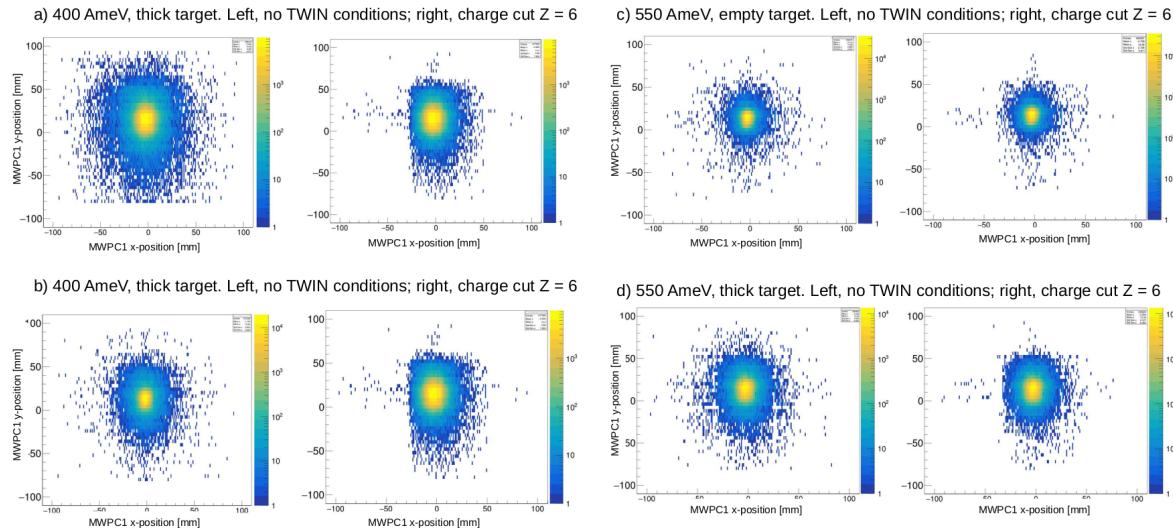


Figure 26: Distribution in x and y on MWPC1 for different energies with and without target. TODO: subplot b) wrong naming,it's the empty run 400 AMeV.

The efficiency loss depends on the target thickness - for thicker targets the geometric distribution of the fragments is broader and therefore the efficiency loss larger - and the beam energies - for larger beam energies the scattering angles decrease due to the boost effects, the efficiency loss gets smaller. This means, that the efficiency loss needs to be compensated runwise. These efficiency effects can be observed in figure 26.

To compensate correctly for the geometric efficiency it has to be considered that for the charge changing cross section measurement only the carbon isotopes after the target are counted in the TWIN MUSIC. Therefore the correction should only be applied to

²⁴This statement does not hold for light fragments as protons or deuterons. Their deflection angle exceeds the geometric acceptance of the MWPC1.

the carbon isotopes ($Z = 6$) on the x-y distribution on the MWPC1, see figure 26. The geometric efficiency correction is done graphically on the x-y distribution of the MWPC1 for carbon isotopes by following procedure:

1. Correction for the x-position distribution:

- (a) First fit x-distribution with double-gaussian function with five free parameters and common mean value μ_x

$$f(x) = A \cdot \exp\left(-\frac{(x - \mu_x)^2}{a^2}\right) + B \cdot \exp\left(-\frac{(x - \mu_x)^2}{b^2}\right) \quad (20)$$

- (b) Fit again within range $\mu_x \pm \epsilon_x$. The parameter ϵ_x is fixed by educated guess, TODO. As μ_x take the value from the fit in the previous step. A fit for the central region of the x-distribution is obtained,

$$f(x)_{central}(A_{central}, a_{central}, B_{central}, b_{central}, \mu_{central}).$$

- (c) The obtained fit function $f(x)_{central}$ is then used to compare with the data distribution($f(x)_{data}$) in the border regions $[-100, \mu_{central} - \epsilon_x]$ and $[\mu_{central} + \epsilon_x, 100]$. Since only the left border region (low x-positions) is affected by the limited geometric acceptance, the right border region can be used for correction:

$$\Delta_{xcorr} = \int_{\mu_{central} + \epsilon_x}^{100} f(x)_{data} - f(x)_{central} - \int_{-100}^{\mu_{central} - \epsilon_x} f(x)_{data} - f(x)_{central} \quad (21)$$

2. Correction for the y-position distribution:

- (a) First fit y-distribution with double-gaussian function with five free parameters and common mean value μ_y

$$f(y)_{fit} = C \cdot \exp\left(-\frac{(y - \mu_y)^2}{c^2}\right) + D \cdot \exp\left(-\frac{(y - \mu_y)^2}{d^2}\right) \quad (22)$$

- (b) The obtained fit function $f(y)$ is then used to compare the data distribution($f(y)_{data}$) in the border regions $[-100, \mu_y - \epsilon_y]$ and $[\mu_y + \epsilon_y, 100]$. The parameter ϵ_y is fixed by educated guess, TODO. As μ_y take the value from the fit in the previous step. Same as for the x-correction both border regions are compared. The high border region (high y-positions) affected by the limited geometric acceptance while the low border region (low y-positions) has full geometric acceptance.

$$\Delta_{ycorr} = \int_{-100}^{\mu_{central} - \epsilon_y} f(y)_{data} - f(y)_{fit} - \int_{\mu_y + \epsilon_y}^{100} f(y)_{data} - f(y)_{fit} \quad (23)$$

3. To correct the number of survived carbon isotopes $N_2 = N_{carbon}$ both corrections in x and y are applied²⁵

$$N_2^{corr} = N_2 + \frac{\Delta_{xcorr} + \Delta_{ycorr}}{N_2} \quad (24)$$

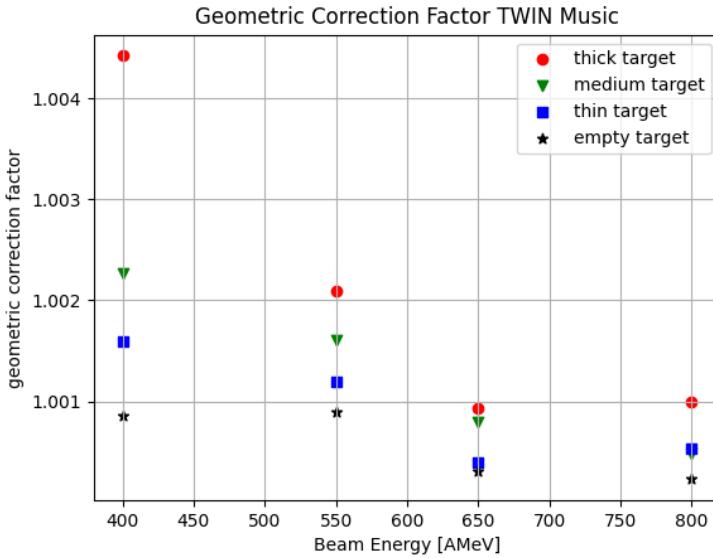


Figure 27: Geometric correction factors from limited geometric efficiency of TWIN Music.

Figure 27 summarizes the geometric correction factors ϵ_{geo_corr} obtained from the graphical reconstruction of missed TWIN Music events as described above. The correction factor is subsequently applied to the charge-changing cross-section, resulting in the final corrected charge-changing cross-section:

$$\sigma_{geo_corr} = -\frac{1}{N_t} \ln\left(\frac{N_1^E N_2}{N_2^E N_1} \cdot \epsilon_{geo_corr}\right) = -\frac{1}{N_t} \left(\ln\left(\frac{N_1^E N_2}{N_2^E N_1}\right) + \ln(\epsilon_{geo_corr}) \right) \quad (25)$$

Figure 28 shows the measured charge changing cross section once without considering the the limited geometric acceptance of the TWIM Music and once applying the geometric correction factors as presented in equation 25. As expected significantly affected by the geometric correction are runs with 400 and 550 AMeV beam energy whereas at beam energies of 650 and 800 AMeV the effect is exceptionally small since at high beam energies the fragments after the target preceive a strong boost effect in beam direction which constrains the distribution in the x-y-plane.

²⁵Under the assumption that x and y are uncorrelated where the x-y distribution on the MWPC1 is given by $f(x, y) = f(x) \cdot f(y)$.

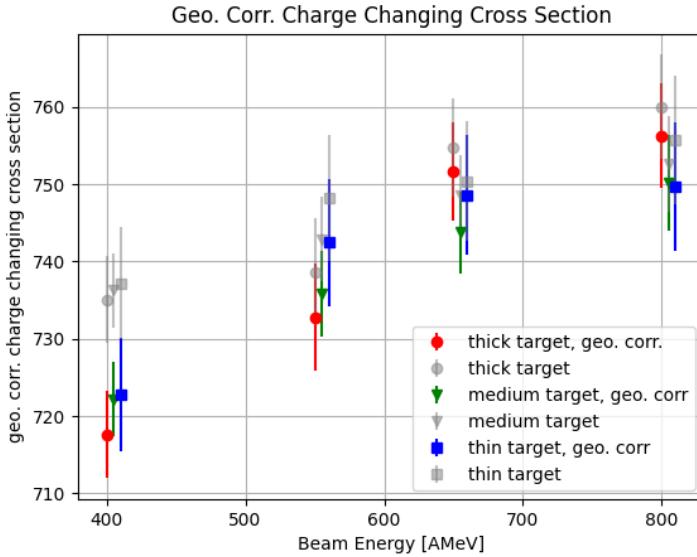


Figure 28: Charge changing cross section with applied geometry corrections. In gray: charge changing cross section measurements before applying corrections, as in figure 23.

4.5 Isotope Correction - Total Interaction Cross Section

The general formulation for the calculation of cross sections in equation 15 can be used to determine the cross section for specific channels depending on the definition of N_2 . In the previous subsection 4.3 where the charge changing cross section was measured, N_2 had to be sensitive to the charge of the outgoing fragment. Therefore N_2 was defined as the number of survived carbon isotopes, i.e. $N_2 = N_2^{Z=6}$. For the measurement of the total interaction cross section N_2 has to be sensitive to both proton and neutron number of the fragments. N_2 is therefore restricted to the number of survived ^{12}C isotopes, i.e. $N_2 = N_2^{^{12}\text{C}}$. Since $N_2^{^{12}\text{C}}$ is a subset of $N_2^{Z=6}$, $N_2^{^{12}\text{C}}$ can be determined by identifying and disentangling the number of survived ^{12}C isotopes from the set of events with carbon isotopes $N_2^{Z=6}$. For that reason the positional correlations on the x-coordinate of MWPC2 (upstream to GLAD) and the MWPC3 (downstream to GLAD) are exploited. The GLAD magnet, which acts as a mass spectrometer, deflects the passing through fragment. Depending on its proton to neutron ratio the fragment is deflected more or less, described by the formula for magnetic rigidity:

$$B \cdot \rho = \frac{\gamma \cdot m \cdot v}{q} \quad (26)$$

where:

- B is the strength of the magnetic field,
- ρ is the radius of curvature of the particle's trajectory,
- γ is the Lorentz factor,
- v is the velocity of the particle,
- m is its mass,
- q is its charge.

Figure 30 shows the x distribution on MWPC2 versus MWPC3 for the thick target run at a beam energy of 400 AMeV. The main diagonal line corresponds to the ^{12}C isotopes. From all isotopes they have the largest mass to proton ratio ($n+p/p$) and are therefore less deflected by the magnetic field of GLAD. The less prominent line below corresponds to ^{11}C isotopes and on the lower edge few events with ^{10}C are visible. To identify the number of survived ^{12}C fragments a graphical selection of the ^{12}C isotopes - the main diagonal line - is applied.

Graphical Selection Algorithm:

- Fit the main diagonal line, which corresponds to the ^{12}C isotopes, with a linear fit function $f(x_{mw2}) = a \cdot x_{mw2} + b$.
- To get the most accurate intersection line between ^{12}C isotopes and all lighter carbon isotopes the linear fit function from the previous step is taken as starting point. Iteratively the offset value b is reduced by $b_i = b_{i-1} - 1$. For all iteration steps the ratio $r_{^{12}\text{C}(i)}$ of hits below the linear fit function and the total number of hits in the histogram is calculated.
- The derivative $\frac{dr_{^{12}\text{C}}(i)}{db_i}$ is calculated.
- Finally the offset b_i with the largest value for $\frac{dr_{^{12}\text{C}}(i)}{db_i}$ is selected as cutting line between ^{12}C isotopes and $^{11}\text{C}/^{10}\text{C}$ isotopes²⁶.

The ratio $r_{^{12}\text{C}}$ is unaffected by detector efficiencies of MWPC2 and MWPC3²⁷ and therefore the ratio $r_{^{12}\text{C}}$ can be applied directly as isotope correction to calculate the total interaction cross section:

$$\sigma_I = \sigma_{geo_corr} + \sigma_{iso}$$

With the isotope correction cross section σ_{iso} :

$$\sigma_{iso} = -\frac{1}{N_t} \ln(r_{^{12}\text{C}})$$

²⁶For empty target runs the offset b is manually selected

²⁷Under the assumption of constant efficiency over x_{mw2} and x_{mw3}

To calculate the isotope correction cross section six different methods were employed and compared with each other:

- **MWPC2 and MWPC3 hit-level data:**

For all MWPCs the standard *cal-to-hit* step sorts the calibrated hits in the detector according to the calibrated charge deposited in the pads. The final position (in mm) is determined by selecting the hit with the highest charge deposition Q_{max} and its left (Q_L) and right neighbour (Q_R) pads²⁸. These charge and position values are inserted in the "hyperbolic squared secant" function [43] with the following charge distribution function:

$$Q(x) = \frac{a_1}{\cosh^2\left(\frac{\pi(x-a_2)}{a_3}\right)}$$

where a_1 is the amplitude of the distribution Q_{max} , a_2 its centroid, and a_3 derives as follows:

$$a_3 = \frac{\pi\omega}{\cosh^{-1}\left(0.5 \times \left(\sqrt{\frac{Q_{max}}{Q_L}} + \sqrt{\frac{Q_{max}}{Q_R}}\right)\right)}$$

ω being the width of the pads. The centroid of the distribution, which is used as final hit position in the *hit-data* level, can be deduced from:

$$a_2 = \frac{a_3}{\pi} \times \tanh^{-1}\left(\frac{\sqrt{\frac{Q_{max}}{Q_L}} - \sqrt{\frac{Q_{max}}{Q_R}}}{2 \sinh\left(\frac{\pi\omega}{a_3}\right)}\right)$$

Figure 29 shows the "hyperbolic squared secant" function with the inserted values for Q_{max} , Q_R and Q_L . The "hyperbolic squared secant" function is used to determine the x hit position as well as the y hit position for all MWPCs. Figure 30 shows the x_{mw2} versus x_{mw3} distribution of carbon isotopes for the 400 AMeV run with the thick target. The two correlated lines corresponding to the ^{12}C and ^{11}C isotopes can clearly be distinguished. The vertical line can be interpreted as amount of events where the incoming centered carbon fragment gets scattered by air or the detector material in place between MWPC2 and MWPC3. The horizontal wide spread line has no physical interpretation and can rather be explained by the *cal-to-hit* step in MWPC2: For events where there is not a spatially constrained hit cluster but sparse hits the hyperbolic squared secant function may pick the wrong Q_{max} and therefore wrongly reconstructs the x position in MWPC2.

- **MWPC2 and MWPC3 data with own "hit-clustering" level:**

To overcome the issue with potentially wrong x-position reconstruction in the

²⁸In case no charge deposition in one or both neighbours the charge value is set to 1 respectively.

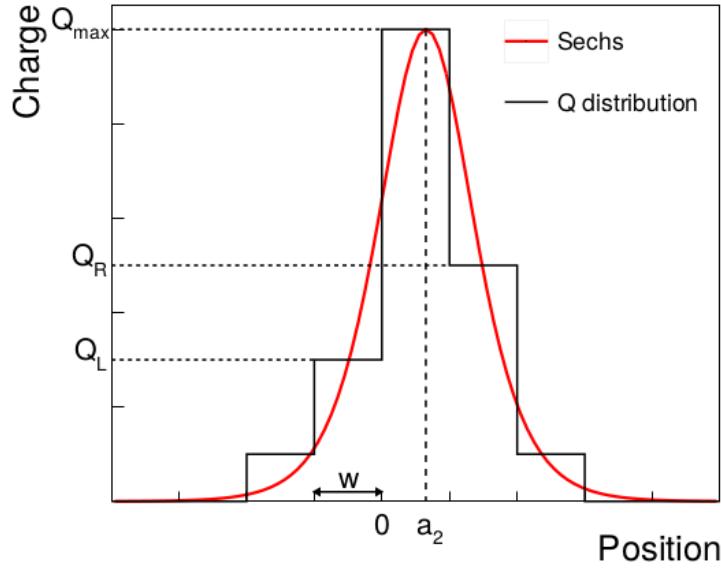


Figure 29: Figure taken from [24], with w being the width of the cathode pads of the MWPC and a_2 the final position value of the hit determined by the hyperbolic squared secant function (in red). In black the measured charge deposition distribution in the MWPC.

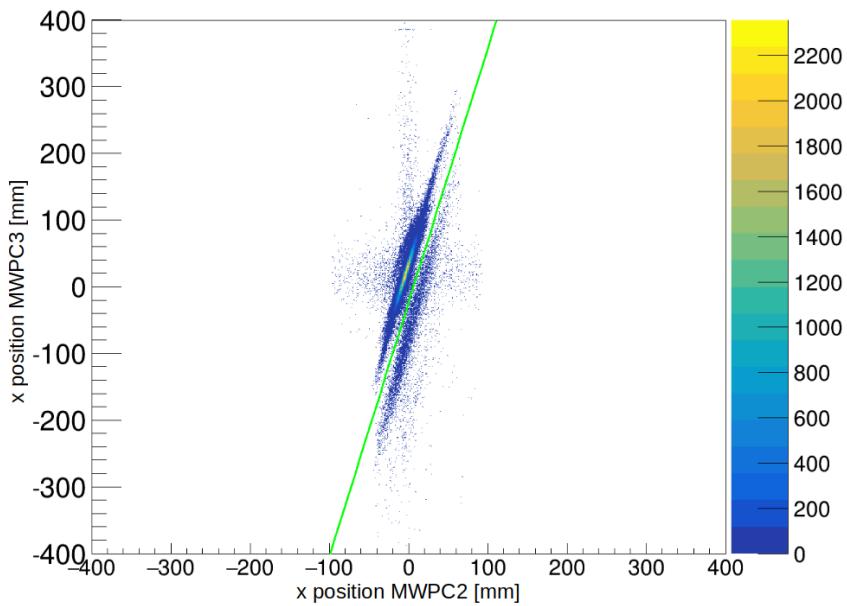
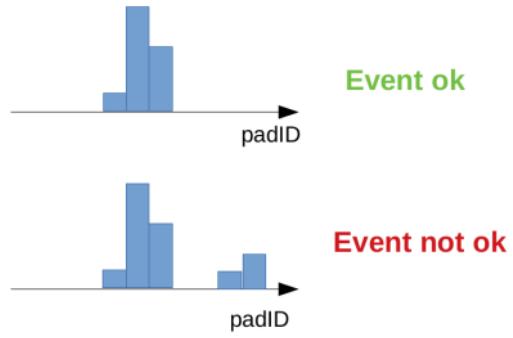


Figure 30: Distribution of x in MWPC2 and MWPC3 for the 400 AMeV run with thick target. The green line corresponds to the intersection line between ^{12}C and $^{11}\text{C}/^{10}\text{C}$ isotopes fixed on the graphical selection algorithm.

Figure 31: Restricted event selection for MWPC2 and MWPC3: only events with one single coherent (i.e. without any holes) cluster are accepted.



MWPCs the event selection on MWPC2 and MWPC3 was restricted to events where both MWPC2 and MWPC3 have only one spatially constrained cluster (see figure 31) to avoid ambiguities in the position determination. This method strongly retains uncorrelated hits in MWPC2 and MWPC3.

Figure 32 shows the distribution of x in MWPC2 and MWPC3 using the own hit-clustering reconstruction. This reconstruction method removes the uncorrelated horizontal line which was observed in figure 30. However the statistics are reduced by $\approx 35\%$ ²⁹.

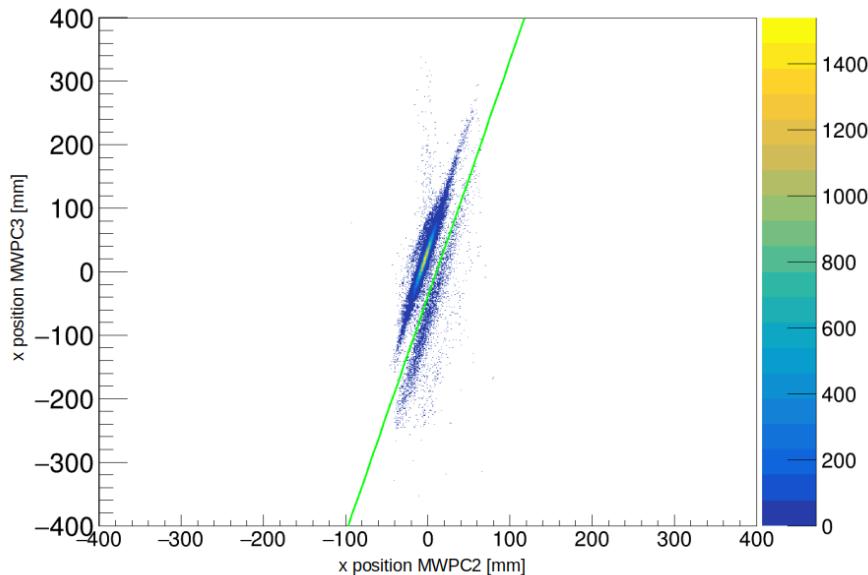


Figure 32: Distribution of x in MWPC2 and MWPC3 using the own hit clustering reconstruction. Thick target run, beam energy 400 AMeV. The green line corresponds to the intersection line between ^{12}C and $^{11}\text{C}/^{10}\text{C}$ isotopes fixed on the graphical selection algorithm.

²⁹Number of entries in the 2D plot for default reconstruction method: 533816, for the own hit clustering reconstruction: 346315 for the 400 AMeV run with thick target.

- **MWPC1 and MWPC3 hit-level data:**

To get the ratio r_{12C} it is necessary to correlate the x positon before and after the GLAD magnet. This task can be completed by MWPC3 with respect to MWPC2 or MWPC1. Since the MWPC1 is upstream to MWPC2 the positional distribution of the carbon fragments narrower which as consequence makes it more difficult to disentangle ^{12}C and $^{11}C/^{10}C$ isotopes, see figure 33. Moreover MWPC1 had two noisy pads which distorts the distribution when using the standard *cal-to-hit* step to get the position value, see figure 33.

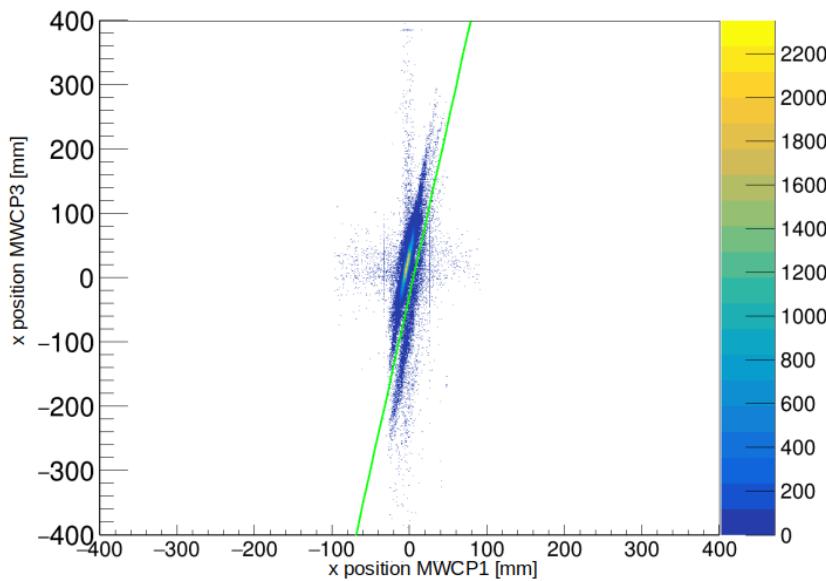


Figure 33: Distribution of x in MWPC1 and MWPC3. Thick target run, beam energy 400 AMeV. The two vertical lines stem from two noisy pads in MWPC1.

- **MWPC1 and MWPC3 data with own "hit-clustering" level:**

Again, to overcome the issue with potentially wrong x-position reconstruction in the MWPCs the own hit clustering reconstruction method, as described above, was applied to MWPC1 and MWPC3 resulting in the 2D plot 34.

- **MWPC2 and MWPC3 with own clustering, quadrant selection in MWPC2:³⁰**

The limited geometric acceptance of TWIN MUSIC, described in section 4.4, affects the isotope correction too. The x distribution on the MWPC1/2/3 is cut off on the lower end and the y distribution on the higher end, see figure 26. The $^{11}C/^{10}C$ isotopes are expected to have a broader x and y distribution. Missing the lower and higher edges in the x and y distribution respectively distorts the r_{12C} ratio towards higher values. This results in a lower cross section contribution

³⁰I did the quadrant selection also in MWPC2, outcome really similar, TODO: add to appendix...

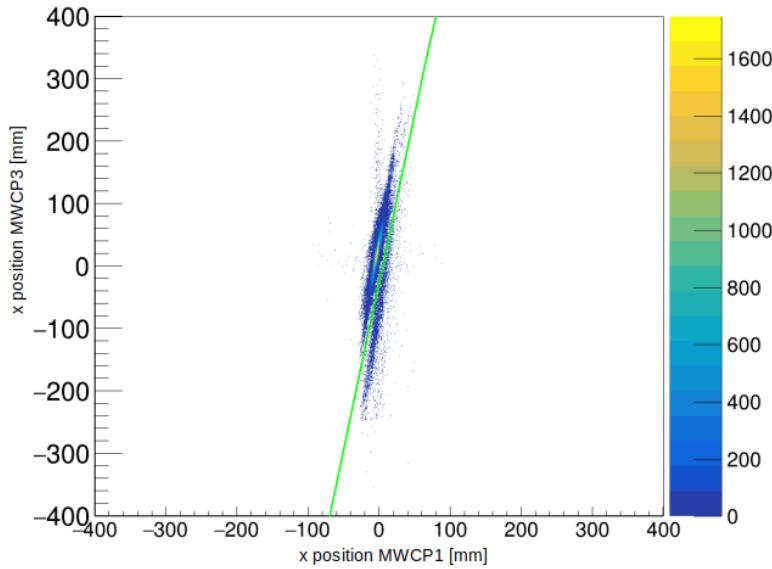


Figure 34: Distribution of x in MWPC1 and MWPC3 using the own hit clustering reconstruction. Thick target run, beam energy 400 AMeV.

of the isotope correction, especially for the low energy runs. To correct for this the x-y distribution in MWPC1 was split into four quadrants, see figure 35. The intersection lines were derived by the mean of the gaussian fit of the x and y distribution. The cross section contribution of the isotope correction was measured for all four quadrants.

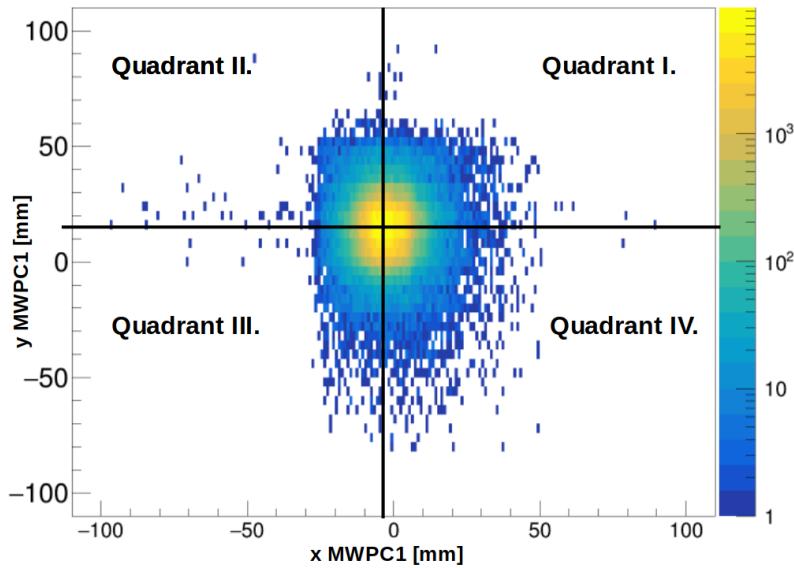


Figure 35: Distribution of x and y in MWPC1 split up in four quadrants. Thick target run, beam energy 400 AMeV.

4.5.1 Results for Isotope Correction Methods

- MWPC2 and MWPC3 hit-level data:

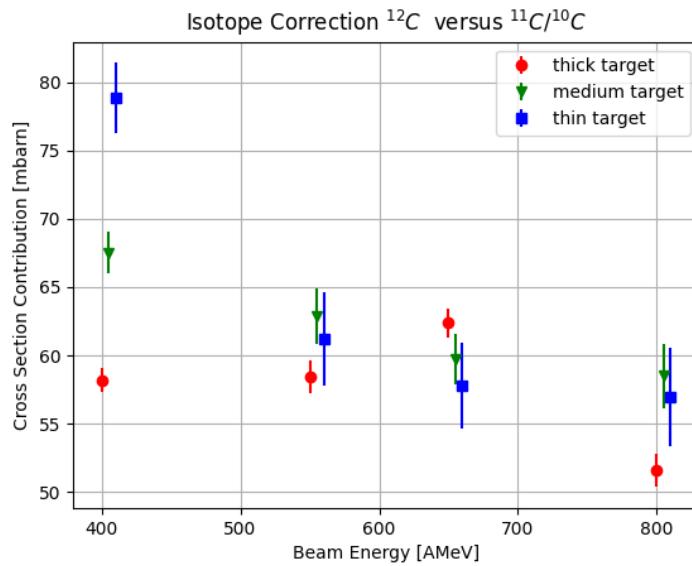


Figure 36: Isotope correction contribution to the total interaction cross section using standard hit level data in MWPC2 and MWPC3.

- MWPC2 and MWPC3 data with own "hit-clustering" level:

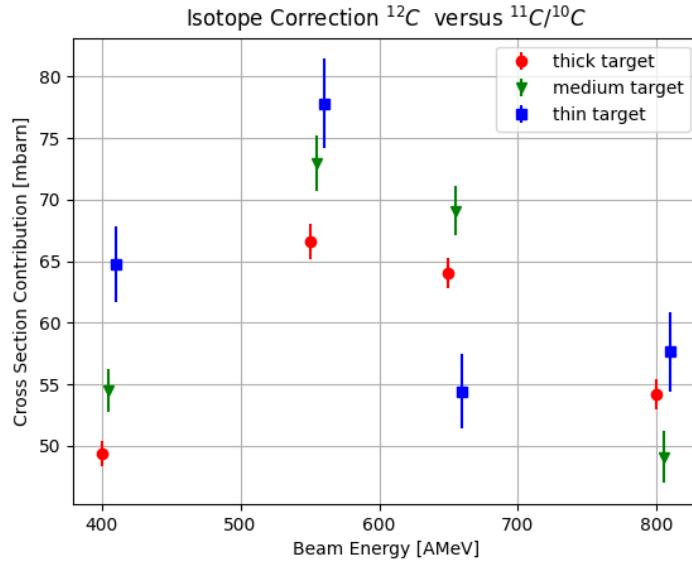


Figure 37: Isotope correction contribution to the total interaction cross section using own hit clustering in MWPC2 and MWPC3.

- MWPC1 and MWPC3 hit-level data:

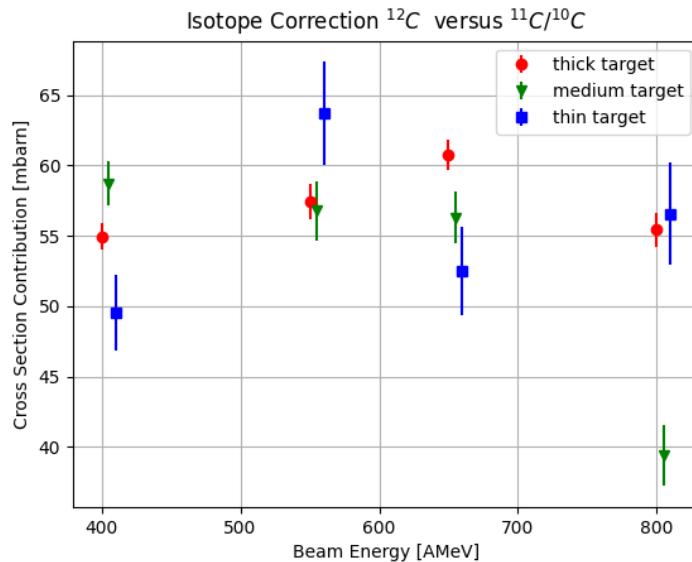


Figure 38: Isotope correction contribution to the total interaction cross section using standard hit level data in MWPC1 and MWPC3.

- MWPC1 and MWPC3 data with own "hit-clustering" level:

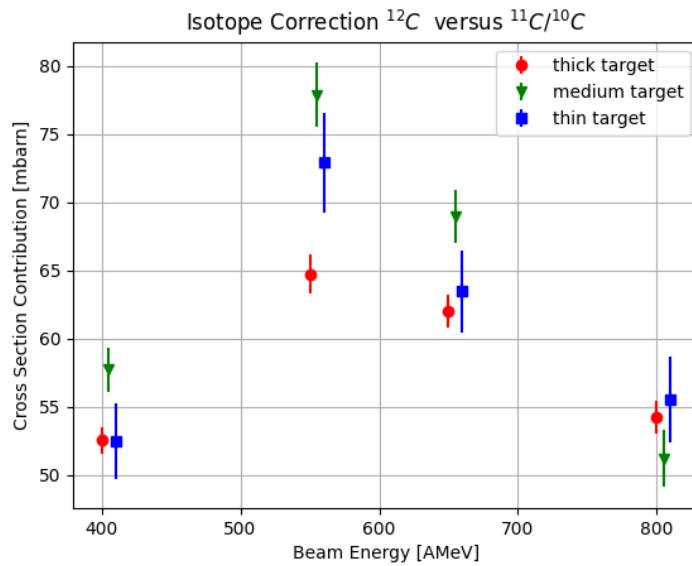


Figure 39: Isotope correction contribution to the total interaction cross section using own hit clustering in MWPC1 and MWPC3.

- MWPC2 and MWPC3 with own clustering, quadrant selection in MWPC1:

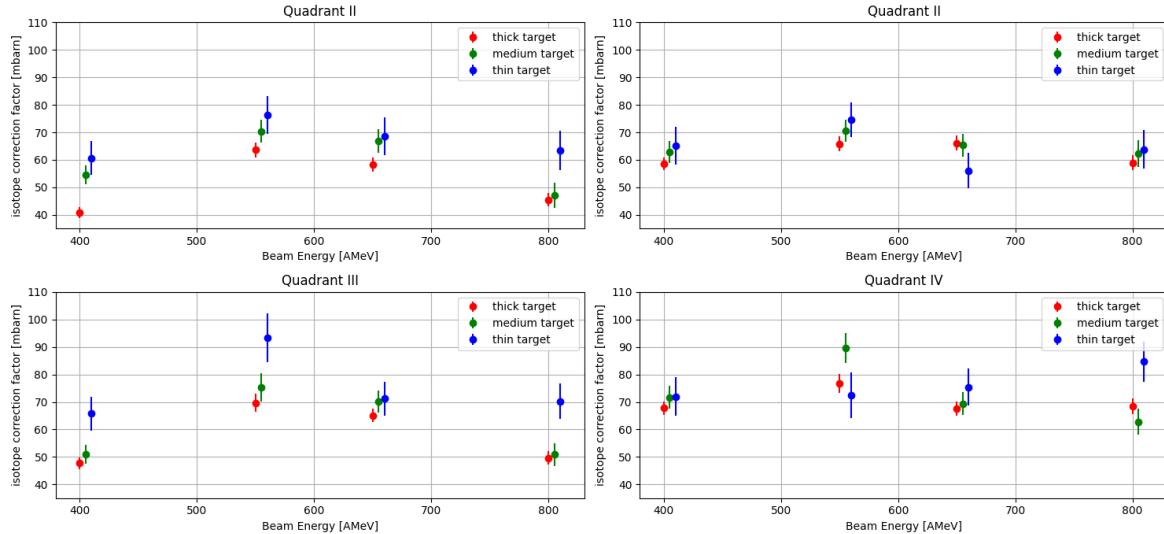


Figure 40: Isotope correction contribution to the total interaction cross section using own hit clustering in MWPC2 and MWPC3. Comparison for different quadrant selection in MWPC1. Quadrant four is the preferred one as it is not affected by limited geometric acceptance of TWIN MUSIC.

4.6 Statistical and Systematic Error Analysis

All three measurements presented in this section, the charge changing cross section, the isotopic cross section correction cross section and the total interaction cross section, rely on the transmission method and the error analysis of those measurements are treated the same accordingly. The generic formula for all three measurements was presented in section 4.1, equation 14.

Since all quantities in equation 14 are mutually independent, the gaussian error propagation for statistical and systematical uncertainties is given by:

$$\Delta_{stat./syst.} = \sqrt{\sum_{i=1}^n \left(\frac{\partial f}{\partial x_i} \cdot \Delta_{x_i} \right)^2} \quad (27)$$

with:

$f(x_1, x_2, \dots, x_n)$ the error prone cross section function,

x_i the independent variables (N_1, N_2, N_t ; see equation 14),

Δ_{x_i} stat./syst. uncertainties associated to independent variable x_i

4.6.1 Statistical Uncertainties

For the charge changing cross section the combined statistical error is:

$$\Delta\sigma_{\text{stat.}|cc} = \sqrt{\left(\frac{\partial\sigma_{cc}}{\partial N_t}\Delta N_t\right)^2 + \left(\frac{\partial\sigma_{cc}}{\partial P_{\text{surv}|cc}}\Delta P_{\text{surv}|cc}\right)^2 + \left(\frac{\partial\sigma_{cc}}{\partial P_{\text{surv}|cc}^E}\Delta P_{\text{surv}|cc}^E\right)^2} \quad (28)$$

where ΔN_t accounts for the uncertainty in the measurement of the target thickness. $P_{\text{surv}|cc}$, the surviving probability of having a carbon isotope in the final state, follows a binomial distribution with N_1 independent experiments and a boolean valued reaction output (survived vs. non-survived). Therefore $P_{\text{surv}|cc}$ follows the standard error of binomial distributed variable:

$$\Delta P_{\text{surv}|cc}^{(E)} = \sqrt{\frac{P_{\text{surv}|cc}^{(E)} \cdot (1 - P_{\text{surv}|cc}^{(E)})}{N_1^{(E)}}} \quad (29)$$

Inserting in equation 28 gives:

$$\Delta\sigma_{\text{stat.}|cc} = \sqrt{\left(\frac{\sigma_{cc}\Delta N_t}{N_t}\right)^2 + \frac{1}{N_t^2} \left(\frac{1 - P_{\text{surv}|cc}}{P_{\text{surv}|cc}N_1} + \frac{1 - P_{\text{surv}|cc}^{(E)}}{P_{\text{surv}|cc}^{(E)}N_1^E}\right)} \quad (30)$$

Similar for the isotopic correction cross section with $P_{\text{surv}|iso}$ the ratio of reconstructed ^{12}C isotopes to all identified carbon isotopes in the x-position correlation plots for MWPC2 and MWPC3:

$$\Delta\sigma_{\text{stat.}|iso} = \sqrt{\left(\frac{\sigma_{iso}\Delta N_t}{N_t}\right)^2 + \frac{1}{N_t^2} \left(\frac{1 - P_{\text{surv}|iso}}{P_{\text{surv}|iso}N_1} + \frac{1 - P_{\text{surv}|iso}^{(E)}}{P_{\text{surv}|iso}^{(E)}N_1^E}\right)} \quad (31)$$

For viewing the statistical uncertainties of the total interaction cross section measurements the uncertainties in the target thickness as well as the uncertainties of the carbon isotope surviving probability $P_{\text{surv}|cc}$ and the ^{12}C surviving probability $P_{\text{surv}|iso}$ have to be accounted for:

$$\Delta\sigma_{\text{stat.}|tot} = \sqrt{\left(\frac{\sigma_{tot}\Delta N_t}{N_t}\right)^2 + \frac{1}{N_t^2} \left(\frac{1 - P_{\text{surv}|cc}}{P_{\text{surv}|cc}N_1} + \frac{1 - P_{\text{surv}|cc}^{(E)}}{P_{\text{surv}|cc}^{(E)}N_1^E}\right) + \frac{1}{N_t^2} \left(\frac{1 - P_{\text{surv}|iso}}{P_{\text{surv}|iso}N_1} + \frac{1 - P_{\text{surv}|iso}^{(E)}}{P_{\text{surv}|iso}^{(E)}N_1^E}\right)} \quad (32)$$

4.6.2 Systematic Uncertainties

4.7 (Preliminary) Results Total Interaction Cross Section

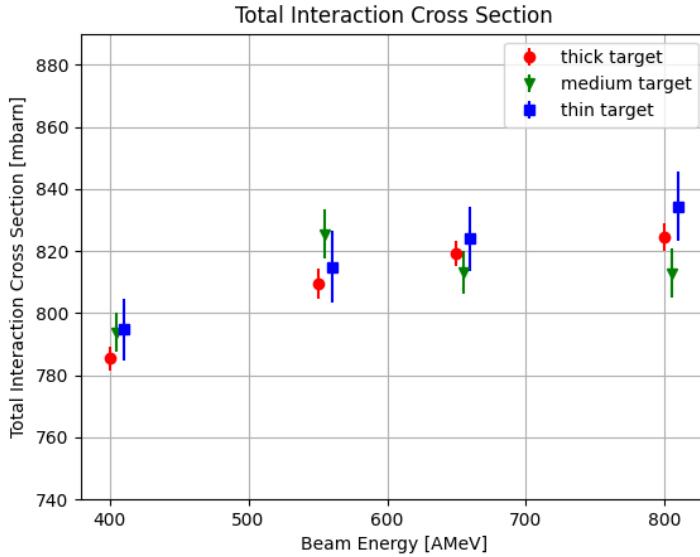


Figure 41: Total interaction cross section using the quadrant IV. selection in MWPC1 x-y plot (see figure 35) for the isotope correction and applying the geometric correction on the charge changing cross section.

5 Qualitative Analysis - Quasi-Free Scattering $^{12}\text{C}(\text{p},2\text{p})^{11}\text{B}$

Until the S444 experiment in 2020 CALIFA consisted out of a prototype frame filled with up to 64 CsI(Tl) crystals. The geometric coverage was therefore poor. For the S444 experiment CALIFA got its final frame and was fully filled in the forward barrel and 35%filled in the iPhos region, CEPA was not installed yet. With these improvements it was possible to commission qfs-experiments with CALIFA at R3B. In the follow up experiment S467, also in 2020, the first experimental run to study single-particle structures of neutron-rich Ca isotopes via qfs-reactions was carried out.

Even though great improvements in the detector development were achieved the correction factors to correct for geometric acceptance would be much too high (≈ 10) for precise cross section measurements for qfs-reactions since the correction factors would in turn rely on a simplified reaction model. A precise analysis of the acceptance correction factor and the development of a more sophisticated and data driven reaction model is out of the scope of this work. Therefore this analysis focusses on the methods of qfs-reaction identification and the extraction of the key informations dicussed in section 2.5. TODO: add more stuff here...

5.1 Setup-Calibration

For setup description refer to section 4. For all detectors except the SOFIA (Study On FIission with Aladin) Time of Flight Wall the calibration parameters investigated by the respective detector-expert group were adopted. Herefore we will subsequently only describe the Time of Flight Wall calibration in this subsection.

5.1.1 Flight-Path Reconstruction - SOFIA Time of Flight Wall

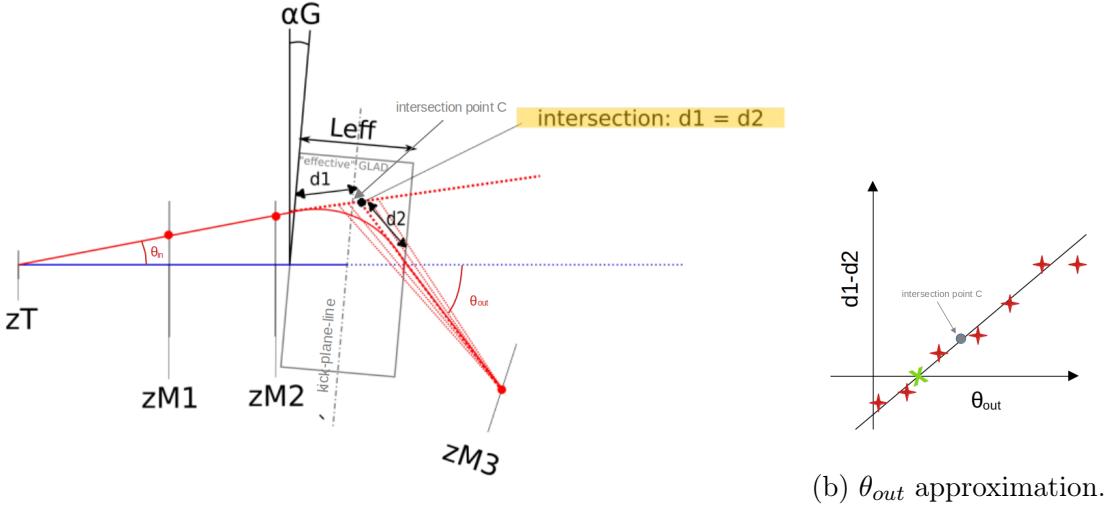
The procedure to calibrate the Time of Flight (ToF) Wall involves beforehand a precise flight-path reconstruction of the projectile from the entrance of Cave C downstream to the ToF Wall. Since only one tracking detector downstream to GLAD was in operation for the S444 experiment no angle of the deflected fragment/beam θ_{out} could be directly extracted. Herefore an advanced tracking algorithm was developed, motivated from ref. [44] (section 3.4).

- The first step is to measure the scattering angle after target, θ_{in} , and draw an extended line from the target position through the effective magnetic field of GLAD.
- Draw a trajectory line from the hit position in MWPC3 to the intersection point C (see figure 42a) of the "kick-plane-line" and the reconstructed and extended track line upstream to GLAD.
- Now sweep along the reconstructed track line upstream to GLAD. For each step a value for θ_{out} and $d1 - d2$ is gathered, see figure 42b.
- As in figure 42b shown, fit the data-points from previous step with a linear fit function and find the corresponding θ_{out} value where the fit line intersects with the abscissa. This corresponds to the case where $d1 = d2$. This is the approximated "kickpoint" in GLAD.
- Previous steps need to be executed for all events accordingly.

Now that the scattering angle after the target θ_{in} and the angle after GLAD, θ_{out} is known and the position outside the magnetic field of GLAD are fixed, with (z_1, x_1) the entrance point on the GLAD field and (z_2, x_2) the exit point, the bending radius r from the magnetic deflection can be determined:

$$r = \frac{L_{eff}}{2 \cdot \sin\left(\frac{\theta_{in} + \theta_{out}}{2}\right) \cdot \cos\delta} \quad (33)$$

where L_{eff} is the effective active width of the magnetic field of GLAD, which corresponds $\approx 2.06\text{ m}$. This value could also be verified by extracting L_{eff} from the formula for



(a) Flighpath reconstruction with "Leff" being the effective active For detailed information see text.
width of the magnetic field of GLAD

Figure 42: Flighpath tracking and reconstruction of the fragment/beam after the target.

the magnetic rigidity $B \cdot r = \gamma\beta m/q$ for empty target runs with known values of the B-field³¹.

The angle δ in equation 33 is given by the trajectory line going through $(z_1, x_1)/(z_2, x_2)$ and the line parallel to the GLAD magnet width L_{eff} .

A detailed derivation of equation 33 can be found in appendix B. The arc trajectory l_{GLAD} of the fragment within GLAD can be reconstructed using the bending radius r and the entry and exit points $(z_1, x_1)/(z_2, x_2)$:

$$l_{GLAD} = r \cdot \omega \quad \text{with} \quad \omega = 2 \cdot \arcsin(t_{1/2}(2 \cdot r)) \quad (34)$$

where ω is the central angle and $t_{1/2}$ is the cord length between (z_1, x_1) and (z_2, x_2) .

At this stage the full pathlength L form the Start detector to the ToFW is fixed:

From Start to the target l_{ST} : For this path section a straight flighpath parallel to the z-coordinate is assumed. The pathlength is taken from the position measurements of the Start detector and the target accordingly ($= 118.3 \text{ cm}$).

Target to GLAD entrance point (z_1, x_1) , l_{in} : For the exact position assignment of (z_1, x_1) both MWPC1 and MWPC2 position measurements have been calibrated with empty target runs by including an offset value in order to center the x-position in both detectors around zero. From the position measurement of the

³¹For detector calibration, primarily for the ToFWall, we had several "empty sweep runs", with empty target and different B-field strength settings. Those runs were optimal to validate L_{eff} .

central position of GLAD, its tilting angle α and L_{eff} the intersection point of the fragment trajectory before GLAD and the "effective GLAD magnetic field rectangle", (z_1, x_1) , can be determined and with it the according flightpath section l_{in} .

Arc trajectory within GLAD, l_{GLAD} : This flightpath passage is determined by reconstructing (z_1, x_1) and (z_2, x_2) , as it has been described in detail in the previous section. Hence the magnetic bending radius can be determined, see equation 33, and finally the arc trajectory l_{GLAD} as in equation 34.

GLAD exit point (z_2, x_2) to ToFW, l_{out} : The hit position (z_{MWPC3}, x_{MWPC3}) in MWPC3 is given by the reconstructed hit position in this detector and the position measurement of the detector itself. The exit point (z_2, x_2) was reconstructed in the previous steps. Hence the trajectory line from (z_2, x_2) to (z_{MWPC3}, x_{MWPC3}) is fixed and is expanded to the intersection with the ToFW plane for the concluding l_{out} measurement.

The resulting flightpath L recombines from:

$$L = l_{ST} + l_{in} + l_{GLAD} + l_{out}$$

In the flightpath reconstruction the deflections in the y-dimension were omitted as the angular straggling in the target is small (TODO: give sigma value) and since the deflection of the fragment within GLAD is independent from its y-position this contribution can be disregarded.

5.1.2 Time of Flight Calibration - SOFIA Start & Time of Flight Wall

For the time of flight measurement in the S444 setup the time is recorded and digitized by the VFTX, VME-FPGA Time-to-Digital Converter (TDC) Modules based on tapped delay line (TDL) TDCs[45]. These modules provide for each detected signal a coarse time, is determined by counting cycles of a 200 MHz readout clock, resulting in a 5 ns binning resolution, and a fine time, which is obtained using an FPGA-based Time-to-Digital Converter (TDC), which employs a tapped delay line (TDL). In this approach, the signal propagates through a series of delayed logic modules within the FPGA until the subsequent clock cycle terminates the sampling process. The number of delay elements traversed by the signal is used to compute the time difference between the signal onset and the end of the clock cycle. The translation of the resulting fine time, with reasonable assumption of an uniform distribution, is achieved via a calibrated linear function. This procedure assigns to each preamplifier signal in the start detector

(left/right) and the ToF Wall (up/down) a calibrated raw time raw_t :

$$raw_t = coarse_time_clocks \cdot 5\text{ns} + offset[fine_time]$$

Finally, the raw time of flight is reconstructed by combining all four time measurements:

$$RawTof = 0.5 * (raw_t_{i,down} + raw_t_{i,up}) - 0.5 * (raw_t_{start,right} + raw_t_{start,left}) \quad (35)$$

where i ($\in 0 \dots 27$) refers to the scintillator number of the ToF Wall. Since the mentioned time measurements are standalone and not synchronized the $RawTof$ has to be corrected by an offset which has to be determined again for each scintillator bar i of the ToF Wall:

$$\Delta_{ToF}[i] = \overline{L[i]} / v_{beam} - \overline{RawTof[i]}$$

where $\overline{L[i]}$ is the mean reconstructed path length for all events in empty target runs which hit scintillator bar i of the ToF Wall. The beam velocity v_{beam} is directly taken from the given beam settings, e.g. 400 AMeV beam, empty target corresponds to $v_{beam} = 214,2\text{mm/ns}$. The mean raw ToF $\overline{RawTof[i]}$ again results from all events which hit scintillator bar i and is extracted from the mean value of a gaussian fit to the raw ToF. The resulting calibrated ToF can then be expressed as:

$$ToF[i] = 0.5 * (raw_t_{i,down} + raw_t_{i,up}) - 0.5 * (raw_t_{start,right} + raw_t_{start,left}) + \Delta_{ToF}[i]$$

To estimate the time of flight resolution for events with hit in ToF Wall bar i it has to be noted that the ToF is affected by the flight path. Hence the the time of flight should be written as:

$$\widetilde{ToF} = \frac{\overline{L[i]}}{\frac{\overline{L[i]}}{ToF[i]}}$$

$\overline{L[i]}$ is the mean pathlength, whereas $L[i]$ and $ToF[i]$ are eventwise selected. By reconstructing the flight-path as described in subsection 5.1.1 and employment of the mentioned time calibration steps the average time resolution σ_t results $\approx 90\text{ps}$ ³².

5.2 Event Selection

For the precise measurement of the total interaction cross section $^{12}C + ^{12}C$, as detailed in the section 4, event selection was of critical importance, requiring stringent cuts on

³²To remove events with large angular straggling, a cut of $\pm 20\text{mm}$ on the beam spot for the y-position on the MWPC3 was applied.

the TPats (see table 5), as well as on the charge and position of the incoming particles. In contrast, for this qualitative QFS analysis, these factors played a minor role, allowing for only minimal cuts on the incoming ions:

- Both left and right preamplifiers of the start detector have seen a coincident signal.
- Exactly one hit in MWPC0 in the hit-level data.

However, for the identification of fragments downstream of the target, various detector signals and event parameters are required:

- One hit (in the hit-level data) in the MWPC tracking detectors, MWPC1 and MWPC2 upstream to GLAD, MWPC3 downstream to GLAD.
- Charge measurement in the TWIN-MUSIC.
- One hit scintillation bar in ToF Wall with signal from both up/down PMTs.

5.3 Fragment Identification

The first step for the identification of the QFS-reaction channel $^{12}C(p, 2p)^{11}B$ is the identification of the fragment ^{11}B via the formula for the magnetic rigidity: $B \cdot r = \gamma\beta m/q$. From the flightpath reconstruction and the ToF measurement as described in previous sections r, β and γ are obtained whereas the charge of the fragment q is directly measured in the TWIN MUSIC detector.

-> mention magnetic rigidity and lorentz transformation stuff...! -> plot A/q plot, explain it well also the peak at Z 8 -> make also a sketch of the qfs- event with the two protons and the fragment and the gamma

5.4 QFS-Protons

->polar angle correlation/opening angle peak ->Azimuthal angle distribution, maybe I can clean it up making a cut on the energy... -> opening angle... do I have a nicer plot? ->Califa front view plot of detector, explain that it was only partly filled etc... ->some simulations to compare.... -> reconstruction of inner momentum. Give also some values -> Correlations between fragment and proton pair

5.4.1 Proton Separation Energy S_p

The proton separation energy S_p of ^{12}C from the reaction channel $^{12}C(p, 2p)^{11}B$ is defined by the masses and energies of the initial and final state particles in the ^{12}C rest

frame. One can write the full energy conservation in the rest frame of ^{12}C :

$$M_{^{12}C} + \gamma \cdot m_p = E_1^* + E_2^* + (M_{^{11}B}} + E_{ex}) + T_{^{11}B^*} \quad (36)$$

where E_1^* and E_2^* are total energies of the two protons in the ^{12}C rest frame, E_{ex} is residual excitation in ^{11}B and $T_{^{11}B^*}$ the kinetic energy of the ^{11}B in the ^{12}C rest frame. Using this equation, one can express the total binding energy B of the knocked out nucleon:

$$B = S_p - E_{ex} = M_{^{12}C} - M_{^{11}B} - E_{ex} - m_p = E_1^* + E_2^* + T_{^{11}B^*} - \gamma \cdot m_p - m_p \quad (37)$$

Using Lorentz transformation from lab to ^{12}C rest frame one can obtain for E_1^*/E_2^* :

$$E_{1/2}^* = \gamma \cdot m_p + \gamma \cdot T_{1/2} - \beta \cdot \gamma \cdot p_{1/2} \cdot \cos(\theta_{1/2})$$

The kinetic energy of the ^{11}B in the ^{12}C rest frame can be written as:

$$T_{^{11}B^*} = \frac{p_i^2}{2 \cdot M_{^{11}B}}$$

with p_i the inner momentum of the proton knocked out of the ^{12}C projectile, see equation 4. Putting this in the previous equation one finally gets for the binding energy B :

$$B = (\gamma - 1) \cdot m_p + \gamma \cdot (T_1 + T_2) - \beta \cdot \gamma \cdot (p_1 * \cos(\theta_1) + p_2 * \cos(\theta_2)) + T_{^{11}B} \quad (38)$$

For $E_{ex} = 0$, i.e. the fragment ^{11}B in the ground state, the binding energy B is equal to the one proton separation energy S_p ³³.

As for the S444 experiment no target detector tracking system was available. Consequently the energies as well as the azimuthal (ϕ) and polar(θ) angles of the two correlated protons had to be fully reconstructed using the CALIFA calorimeter. The calorimeter provided an energy resolution $\frac{\Delta E}{E}(@100MeV) \lesssim 1\%$, along with angular resolutions of approximately $\Delta\phi \approx \frac{6}{\sqrt{12}}^\circ$, and $\Delta\theta \approx \frac{2}{\sqrt{12}}^\circ$. The best way to visualize the separation energy S_p is to plot it against the summed energy of the two protons in the ^{12}C rest frame as shown in figure 43. Two vertical lines are visible. They correspond to the two QFS-reaction types within the CH_2 target: the proton within the ^{12}C projectile can either scatter on the hydrogen (proton-like) part or on the carbon part of the plastic target. For the first case only the separation energy S_p as derived in equation 5.4.1 is necessary to remove the proton of the projectile within the ^{12}C nucleus. In the second

³³Since the ^{11}B fragment predominantly remains in its ground state—implying that, in most cases, the outermost protons are removed—the binding energy B and the proton separation energy S_p are used interchangeably, also due to the limited energy resolution.

case the QFS-reaction is between two protons both bound within a carbon nucleus – the projectile carbon and the target carbon part. Therefore more energy is needed to free both protons from their nuclear bond. This corresponds to the left vertical line in figure 43. It should be noted that the reconstructed one-proton separation energy S_p is shifted with respect to the mean value of $\approx 16\text{MeV}$. For a precise measurement the accurate position of the target and the reaction vertex would be needed, as well as a precise measurement of the kinetic energy of the incoming ¹²C would be required (for the actual measurements the beam energy value was set to 400AMeV .

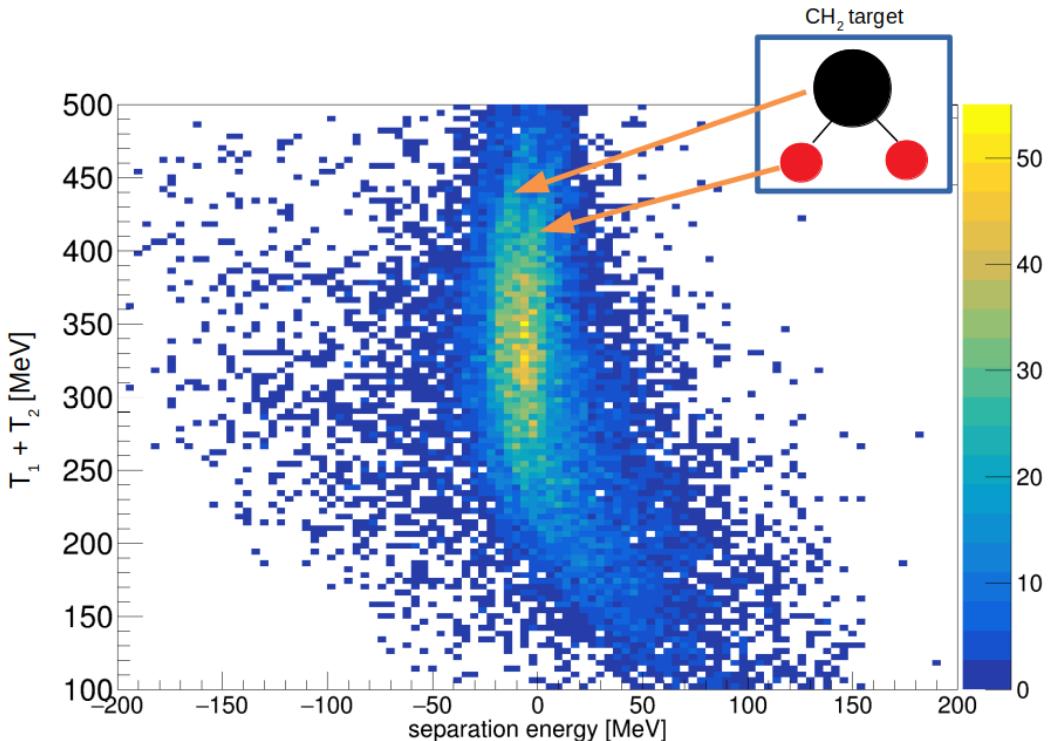


Figure 43: Sum of the kinetic energies of the two correlated protons (T_1/T_2) versus the one proton separation energy S_p . Since this energy is needed to solve the proton from the nucleus' core, its value has conventionally a negative sign.

5.5 Reconstruction of excited¹¹B states

->Show level scheme of ¹¹B, explain excited states and why it is not expected to see any peaks at 4.4 MeV. ->Explain Doppler correction (this is really the key point of the CALIFA, why it is highly segmented) -> show all conditions that you apply -> This comm. exp. has shown that we can really pin down reaction mechanism and we are ready for exciting p2p experiments -> Limitations: CALIFA was only partly filled for this exp.

6 Results and Discussion

Appendices

A TWIN MUSIC Geometric Acceptance Correction via Efficiency Measurement

Instead of correcting the limited geometric acceptance of TWIN MUSIC via graphical fitting (see section 4.4) it is also feasible correcting via TWIN MUSIC efficiency measurement. The correction factor is given by:

$$\epsilon_{geo_corr} = \frac{N_{MWPC1,MWPC2}}{N_{MWPC1,MWPC2,TWIN}} \quad (39)$$

where $N_{MWPC1,MWPC2}$ corresponds to the number of events with a hit in MWPC1 and MWPC2 whereas $N_{MWPC1,MWPC2,TWIN}$ imposes the further condition having a hit in TWIN MUSIC too.

The corresponding correction factor ϵ_{geo_corr} is consequently applied on all target and empty runs. The resulting corrected charge changing cross section is shown in figure 44. The same correction factor can be applied to the total interaction cross section as

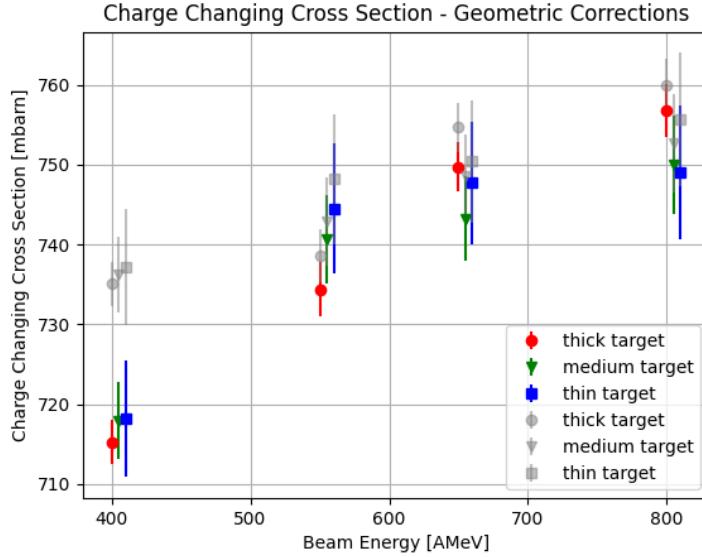


Figure 44: Charge changing cross section correction due to limited geometric acceptance of TWIN MUSIC via efficiency correction with MWPC1 and MWPC2. In gray: charge changing cross section measurements before applying corrections, as in figure 23

in figure 45.

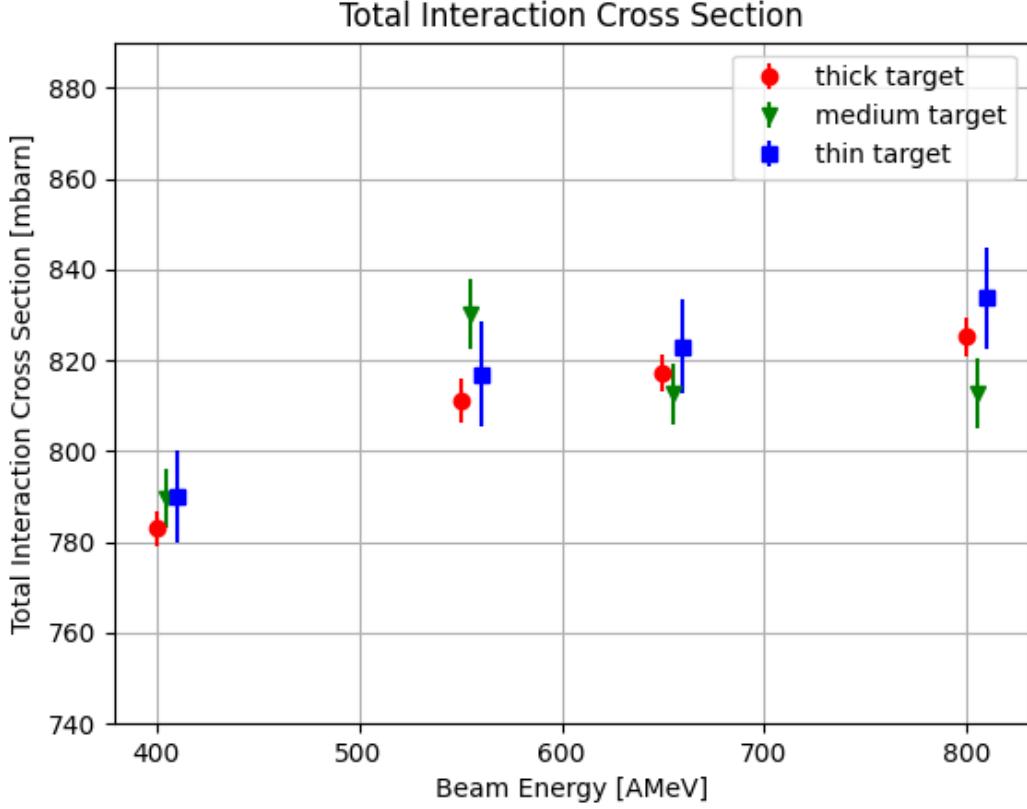


Figure 45: Total interaction cross section of $^{12}\text{C} + ^{12}\text{C}$ using the TWIN MUSIC efficiency correction factor, see equation 39, to compensate for the limited geometric acceptance in TWIN MUSIC.

B Flight-Path Reconstruction

The first step in the radius and flightpath reconstruction is expressing the entrance point on the GLAD field (z_1, x_1) and the exit point (z_2, x_2) with the center of the circle path (z_0, x_0) as reference, see figure 46:

$$\begin{aligned}
 z_1 &= z_0 - r \cos(90^\circ - \theta_i) = z_0 - r \sin(\theta_i) \\
 x_1 &= x_0 + r \sin(90^\circ - \theta_i) = x_0 + r \cos(\theta_i) \\
 z_2 &= z_0 + r \cos(90^\circ - \theta_o) = z_0 + r \sin(\theta_o) \\
 x_2 &= x_0 + r \sin(90^\circ - \theta_o) = x_0 + r \cos(\theta_o)
 \end{aligned}$$

The slope m_1 of the intersection line between (z_1, x_1) and (z_2, x_2) is given by:

$$m_1 = \frac{x_2 - x_1}{z_2 - z_1} = \frac{\cos(\theta_o) - \cos(\theta_i)}{\sin(\theta_o) + \sin(\theta_i)}$$

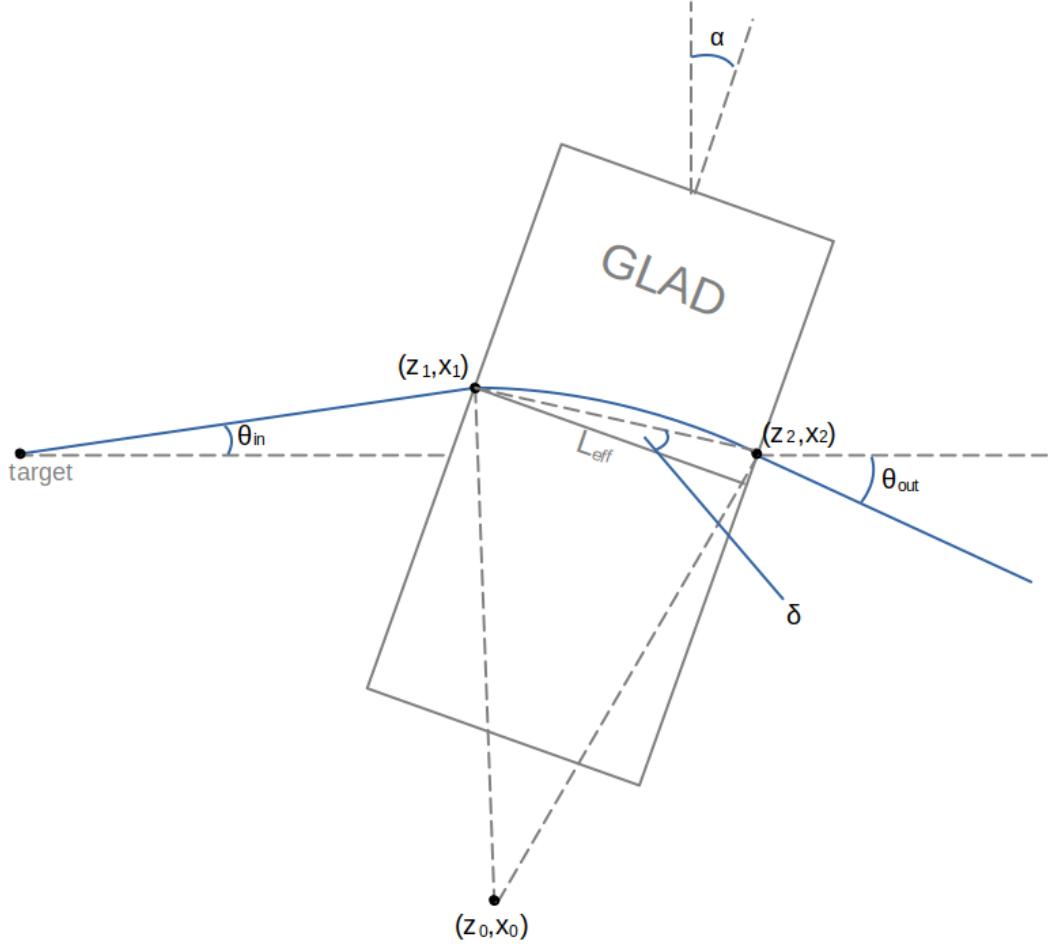


Figure 46: Flightpath reconstruction with reference positions (z_0, x_0) , (z_1, x_1) and (z_2, x_2) . The GLAD magnet is tilted by $\alpha = 14^\circ$.

and with the distance between the two points given by:

$$\begin{aligned}\Delta_{i/o}^2 &= r^2 \left[(\cos\theta_o - \cos\theta_i)^2 + (\sin\theta_o + \sin\theta_i)^2 \right] \\ &= 4r^2 \sin^2\left(\frac{\theta_i}{2} + \frac{\theta_o}{2}\right) \\ \Rightarrow \Delta_{i/o} &= 2r \sin\left(\frac{\theta_i}{2} + \frac{\theta_o}{2}\right)\end{aligned}$$

To describe the distance between (z_1, x_1) and (z_2, x_2) with the given effective GLAD length L_{eff} ($= 2.06\text{ m}$) the tilting angle α (see figure 46) of GLAD in relation to the incoming beam line direction has to be considered. Consequential the angle δ between the trajectory connecting (z_1, x_1) and (z_2, x_2) and the line parallel to the GLAD magnet

width L_{eff} can be determined as:

$$\tan(\delta) = \left| \frac{m_1 - m_2}{1 + m_1 \cdot m_2} \right|$$

with $m_2 = -\tan(\alpha)$:

$$\delta = \text{atan} \left(\frac{\frac{\cos\theta_o - \cos\theta_i}{\sin\theta_o + \sin\theta_i} + \tan\alpha}{1 - \frac{\cos\theta_o - \cos\theta_i}{\sin\theta_o + \sin\theta_i} \cdot \tan\alpha} \right)$$

The final relation between L_{eff} and the bending radius r of the fragment within GLAD can be written as:

$$\begin{aligned} L_{eff} &= 2r \sin\left(\frac{\theta_i}{2} + \frac{\theta_o}{2}\right) \cdot \cos\delta \\ &= 2r \sin\left(\frac{\theta_i}{2} + \frac{\theta_o}{2}\right) \cdot \frac{1}{\sqrt{\delta^2 + 1}} \end{aligned} \quad (40)$$

C second stuff I want to add

hello here to add, end.

References

- [1] O. Chamberlain and E. Segrè, “Proton-proton collisions within lithium nuclei,” *Physical Review*, vol. 87, no. 1, p. 81, 1952.
- [2] J. Cladis, W. Hess, and B. Moyer, “Nucleon momentum distributions in deuterium and carbon inferred from proton scattering,” *Physical Review*, vol. 87, no. 3, p. 425, 1952.
- [3] T. Maris, P. Hillman, and H. Tyrāln, “Quasi-elastic scattering and nuclear structure,” *Nuclear Physics*, vol. 7, pp. 1–9, 1958.
- [4] M. Patsyuk, J. Kahlbow, G. Laskaris, M. Duer, V. Lenivenko, E. Segarra, T. Atovullaev, G. Johansson, T. Aumann, A. Corsi, *et al.*, “Unperturbed inverse kinematics nucleon knockout measurements with a carbon beam,” *Nature physics*, vol. 17, no. 6, pp. 693–699, 2021.
- [5] P. Hansen and J. Tostevin, “Direct reactions with exotic nuclei,” *Annual Review of Nuclear and Particle Science*, vol. 53, no. 1, pp. 219–261, 2003.
- [6] W. Dickhoff and C. Barbieri, “Self-consistent green’s function method for nuclei and nuclear matter,” *Progress in Particle and Nuclear Physics*, vol. 52, no. 2, pp. 377–496, 2004.

- [7] A. S. Goldhaber, “Statistical models of fragmentation processes,” *Physics Letters B*, vol. 53, no. 4, pp. 306–308, 1974.
- [8] H. Feshbach and K. Huang, “Fragmentation of relativistic heavy ions,” *Physics Letters B*, vol. 47, no. 4, pp. 300–302, 1973.
- [9] L. Chulkov, F. Aksouh, A. Bleile, O. Bochkarev, D. Cortina-Gil, A. Dobrovolsky, P. Egelhof, H. Geissel, M. Hellström, N. Isaev, *et al.*, “Quasi-free scattering with 6, 8he beams,” *Nuclear Physics A*, vol. 759, no. 1-2, pp. 43–63, 2005.
- [10] V. Panin, T. Aumann, and C. Bertulani, “Quasi-free scattering in inverse kinematics as a tool to unveil the structure of nuclei: A tribute to mahir s. hussein,” *The European Physical Journal A*, vol. 57, pp. 1–13, 2021.
- [11] L. Atar, S. Paschalis, C. Barbieri, C. A. Bertulani, P. Díaz Fernández, M. Holl, M. A. Najafi, V. Panin, H. Alvarez-Pol, T. Aumann, *et al.*, “Quasifree ($p, 2 p$) reactions on oxygen isotopes: observation of isospin independence of the reduced single-particle strength,” *Physical review letters*, vol. 120, no. 5, p. 052501, 2018.
- [12] M. Hjorth-Jensen, “The carbon challenge,” *Physics*, vol. 4, p. 38, 2011.
- [13] K. Sumiyoshi and G. Roepke, “Appearance of light clusters in post-bounce evolution of core-collapse supernovae,” *Physical Review CâĂŤNuclear Physics*, vol. 77, no. 5, p. 055804, 2008.
- [14] J. Mabiala, A. Cowley, S. Förtsch, E. Buthelezi, R. Neveling, F. Smit, G. Steyn, and J. Van Zyl, “Analyzing power and cross section distributions of the $c\ 12 (p, p \alpha) be\ 8$ cluster knockout reaction at an incident energy of 100 mev,” *Physical Review CâĂŤNuclear Physics*, vol. 79, no. 5, p. 054612, 2009.
- [15] L. Johannsen, A. Jensen, and P. Hansen, “The ^{11}li nucleus as a three-body system,” *Physics Letters B*, vol. 244, no. 3-4, pp. 357–362, 1990.
- [16] H. Simon, D. Aleksandrov, T. Aumann, L. Axelsson, T. Baumann, M. Borge, L. Chulkov, R. Collatz, J. Cub, W. Dostal, *et al.*, “Direct experimental evidence for strong admixture of different parity states in ^{11}li ,” *Physical review letters*, vol. 83, no. 3, p. 496, 1999.
- [17] C. Bertulani and M. S. Hussein, “Geometry of borromean halo nuclei,” *Physical Review CâĂŤNuclear Physics*, vol. 76, no. 5, p. 051602, 2007.
- [18] W. Xing, X.-G. Wang, and A. W. Thomas, “Electromagnetic form factors for nucleons in short-range correlations,” *Physics Letters B*, vol. 846, p. 138195, 2023.

- [19] “Probing high-momentum protons and neutrons in neutron-rich nuclei,” *Nature*, vol. 560, no. 7720, pp. 617–621, 2018.
- [20] S. Paschalis, M. Petri, A. Macchiavelli, O. Hen, and E. Piasetzky, “Nucleon-nucleon correlations and the single-particle strength in atomic nuclei,” *Physics Letters B*, vol. 800, p. 135110, 2020.
- [21] G. García-Jiménez, H. Alvarez-Pol, A. Graña-González, J. Rodríguez-Sánchez, J. Benlliure, D. Cortina-Gil, L. Atar, L. Audouin, G. Authelet, A. Besteiro, *et al.*, “Study of (p, 2p) fission reactions in inverse kinematics using the r3b set-up,” in *EPJ Web of Conferences*, vol. 290, p. 02009, EDP Sciences, 2023.
- [22] J. Graña-González, J. Rodríguez-Sánchez, J. Benlliure, G. García-Jiménez, H. Alvarez-Pol, D. Cortina-Gil, L. Atar, L. Audouin, G. Authelet, A. Besteiro, *et al.*, “Fission studies in inverse kinematics with the r3b setup,” in *EPJ Web of Conferences*, vol. 290, p. 02015, EDP Sciences, 2023.
- [23] G. Charpak, R. Bouclier, T. Bressani, J. Favier, and Č. Zupančič, “The use of multiwire proportional counters to select and localize charged particles,” *Nuclear Instruments and Methods*, vol. 62, no. 3, pp. 262–268, 1968.
- [24] J.-F. Martin, J. Taïeb, G. Boutoux, A. Chatillon, T. Gorbinet, E. Pellereau, L. Audouin, A. Heinz, H. Alvarez-Pol, Y. Ayyad, *et al.*, “Fission-fragment yields and prompt-neutron multiplicity for coulomb-induced fission of u 234, 235 and np 237, 238,” *Physical Review C*, vol. 104, no. 4, p. 044602, 2021.
- [25] GSI Helmholtzzentrum für Schwerionenforschung GmbH, “Large-acceptance dipole magnet glad.” https://www.gsi.de/fileadmin/Kernreaktionen/GLAD_in_CaveC.jpg, 2013. [Online; accessed September 09, 2024].
- [26] “Technical report for the design, construction and commissioning of the califa barrel.” https://igfae.usc.es/~r3b/documentos/TDR/CALIFA_BARREL_TDR.pdf. Accessed: 2024-09-24.
- [27] M. Bendel, *Entwicklung einer neuartigen Nachweismethode hochenergetischer Teilchen im CALIFA-Kalorimeter*. Dr., Technische Universität München, München, 2014. Technische Universität München, Diss., 2014.
- [28] “Technical report for the design, construction and commissioning of the califa endcap.” https://edms.cern.ch/ui/file/1833748/2/TDR_R3B_CALIFA_ENDCAP_public.pdf. Accessed: 2024-09-24.

- [29] I. Holl, E. Lorenz, and G. Mageras, “A measurement of the light yield of common inorganic scintillators,” *IEEE Transactions on Nuclear Science*, vol. 35, no. 1, pp. 105–109, 1988.
- [30] A. Syntfeld-Kazuch, L. Swiderski, W. Czarnacki, M. Gierlik, W. Klamra, M. Moszynski, and P. Schotanus, “Non-proportionality and energy resolution of csi (tl),” *IEEE Transactions on Nuclear Science*, vol. 54, no. 5, pp. 1836–1841, 2007.
- [31] M. Winkel, *Implementierung und Erprobung einer digitalen Pulsformanalyse zur Auslese von Kalorimetern*. PhD thesis, Diplomarbeit, 2011.
- [32] M. O. Winkel, *Komplexe Pulsformalgorithmen und Teilchenidentifikation zur Echtzeit-Implementierung in CALIFA*. PhD thesis, Technische Universität München, 2016.
- [33] A.-L. Hartig, “Evolution of califa: From single detector modules to benchmark reactions,”
- [34] “Mesytec.” <https://www.mesytec.com/>. [Online; accessed September 29, 2024].
- [35] “Febex 3b-front end board with optical link extension.” <https://www.gsi.de/fileadmin/EE/Module/FEBEX/febex3b.pdf>. [Online; accessed September 30, 2024].
- [36] K. Philipp, *Gamma detection with CALIFA*. PhD thesis, Technische Universität München, 2024.
- [37] GSI Helmholtzzentrum für Schwerionenforschung GmbH, “Exploder.” <https://www.gsi.de/en/work/forschung/experimentelektronik/digitalelektronik/digitalelektronik/module/lwl/exploder>. [Online; accessed September 30, 2024].
- [38] A. Bail, J. Taieb, A. Chatillon, G. Belier, B. Laurent, and E. Pellereau, “Time of flight measurement on the sofia experiment,” in *2011 2nd International Conference on Advancements in Nuclear Instrumentation, Measurement Methods and their Applications*, pp. 1–4, IEEE, 2011.
- [39] K. Boretzky, I. Gašparić, M. Heil, J. Mayer, A. Heinz, C. Caesar, D. Kresan, H. Simon, H. T. Törnqvist, D. Körper, *et al.*, “Neuland: The high-resolution neutron time-of-flight spectrometer for r3b at fair,” *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 1014, p. 165701, 2021.

- [40] K. Boretzky, G. Alkazov, L. Atar, T. Aumann, and C. Beinrucker, “Neuland-from double-planes to the demonstrator,” *GSI Sci. Rep.*, vol. 200, p. 2015, 2014.
- [41] L. Ponnath, *Precise Measurement of Nuclear Interaction Cross Sections with R3B*. PhD thesis, Technische Universität München, 2023.
- [42] “Mdpp-16 detector readout systems, mesytec.” <https://www.mesytec.com/products/nuclear-physics/MDPP-16.html>. [Online; accessed October 18, 2024].
- [43] K. Lau and J. Pyrlik, “Optimization of centroid-finding algorithms for cathode strip chambers,” *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment*, vol. 366, no. 2-3, pp. 298–309, 1995.
- [44] O. Bertini, *Study of 3 Lambda H and 4 Lambda H in the reaction of 6 Li+ 12 C at 2 A GeV*. PhD thesis, Dissertation, Mainz, Johannes Gutenberg-Universität Mainz, 2013.
- [45] E. Bayer and M. Traxler, “Development of a high-resolution (< 10 ps rms) 32-channel tdc in a field-programmable-gate-arrays,” *Instruments-Methods-45, GSI Scientific Report 2009*, p. 325, 2009.