

Exercise 1:

$$H = \frac{p_e^2}{2m_e} + \frac{p_n^2}{2m_n} - \frac{e^2}{4\pi\epsilon_0 r}$$

$$P = p_e + p_n \quad p = \frac{m_n p_e - m_e p_n}{M}$$

$$M = m_e + m_n \quad \mu = \frac{m_e m_n}{M}$$

$$\frac{p_e^2}{2m_e} + \frac{p_n^2}{2m_n}$$

$$\begin{aligned} \frac{p^2}{2M} + \frac{p^2}{2\mu} &= \frac{p_e^2 + p_n^2 + p_e p_n + p_n p_e}{2M} \\ &\quad + \frac{m_n^2 p_e^2 + m_e^2 p_n^2 - m_e m_n p_e p_n}{2\mu M^2} \\ &\quad - \frac{m_e m_n p_n p_e}{2\mu M^2} \\ &= \frac{p_e^2}{2M} + \frac{p_n^2}{2M} + \frac{\cancel{p_e p_n}}{2M} + \frac{\cancel{p_n p_e}}{2M} \\ &\quad + \frac{m_n^2 p_e^2}{2\mu M^2} + \frac{m_e^2 p_n^2}{2\mu M^2} \\ &\quad - \frac{\cancel{p_e p_n}}{2M} - \frac{\cancel{p_n p_e}}{2M} \\ &= \frac{p_e^2}{2M} + \frac{p_n^2}{2M} + \frac{m_n p_e^2}{2m_e M} + \frac{m_e p_n^2}{2m_n M} \end{aligned}$$

$$\begin{aligned}
&= \frac{p_e^2}{2} \left( \frac{m_n}{m_e} \cdot \frac{1}{m_e + m_n} + \frac{1}{m_e + m_n} \right) \\
&\quad \frac{p_n^2}{2} \left( \frac{m_e}{m_n} \cdot \frac{1}{m_e + m_n} + \frac{1}{m_e + m_n} \right) \\
&= \frac{p_e^2}{2m_e} + \frac{p_n^2}{2m_n}
\end{aligned}$$

Exercise 2:

In order for there to be a potential well, we need that, for

all  $\lambda \gg 0$ , there is a minimum

for  $\frac{\lambda}{r^2} - \frac{\gamma}{r}$

but  $-\frac{2\lambda}{r^3} + \frac{\gamma}{r^2} = 0 \rightarrow r^2 = \frac{2\lambda}{\gamma}$

Which always exists since

$\lambda > 0$ ,  $\gamma > 0$ .

Exercise 3:

$$\begin{aligned}\delta E_1 &= \langle \psi_0 | V | \psi_0 \rangle, \\ &= \langle 100 | V | 100 \rangle\end{aligned}$$

$$\begin{aligned}\langle r \theta \phi | 100 \rangle &= R_{10}(r) Y_0^0(\Omega) \\ &= \frac{1}{\sqrt{\pi} a_0^3} e^{-r/a_0}\end{aligned}$$

$$\begin{aligned}V(r) &= \frac{\hbar c \alpha}{2a_n} \left( \frac{r^2}{a_n^2} - 3 \right) + \frac{\hbar c \alpha}{r} \\ &= \frac{\hbar c \alpha}{2a_n} \left( \frac{r^2}{a_n^2} - 3 + \frac{2a_n}{r} \right)\end{aligned}$$

$$\begin{aligned}\delta E_1 &= \frac{1}{\pi a_0^3} \int d\Omega \int_0^{a_n} \frac{\hbar c \alpha}{2a_n} \left( \frac{r^2}{a_n^2} - 3 + \frac{2a_n}{r} \right) e^{-\frac{2r}{a_0}} r^2 dr \\ &= \frac{2 \hbar c \alpha}{a_0^3 a_n} \int_0^{a_n} r^2 \left( \frac{r^2}{a_n^2} - 3 + \frac{2a_n}{r} \right) e^{-\frac{2r}{a_0}} dr\end{aligned}$$

Exercise 4:

$$\delta E_1 = \frac{1}{\pi a_0^3} \int d\Omega \int_0^{a_n} r^2 \frac{\hbar c \alpha}{2a_n} \left( \frac{r^2}{a_n^2} - 3 + \frac{2a_n}{r} \right) e^{-2r/a_0} dr$$

$$\approx \frac{2 \hbar c \alpha}{a_0^3 a_n} \int_0^{a_n} \left( \frac{r^4}{a_n^2} - 3r^2 + 2a_n r \right) dr$$

$$= \frac{2 \hbar c \alpha}{a_0^3 a_n} \left( \frac{a_n^5}{5a_n^2} - a_n^3 + a_n^3 \right)$$

$$\delta E_1 = \frac{2 \hbar c \alpha}{5 a_0^3 a_n} a_n^3 = \frac{2 \hbar c \alpha a_n^2}{5 a_0^3}$$

# DC Polarizability of Hydro. Atom

Exercise 1: Since  $V = e\mathcal{E}z$  and

$$\langle \psi_0 | V | \psi_0 \rangle = \iiint d\Omega dr \psi_0^* \psi_0 e\mathcal{E}z$$

is an integral of an odd function over the symmetric interval  $[0, \pi]$ ,  $\langle \psi_0 | V | \psi_0 \rangle = 0$

Exercise 2:

$$[z, f(p_z)] = i\hbar f'(p_z)$$

$$\begin{aligned} \rightarrow [z, p^2] &= [z, p_z^2] = 2i\hbar p_z \\ &= -2\hbar^2 \partial_z \end{aligned}$$

Exercise 3:

$$\begin{aligned} [p^2, r] \psi &= -\hbar^2 [\nabla^2, r] \psi \\ &= -\hbar^2 (\nabla^2(r\psi) - r\nabla^2\psi) \\ &= -\hbar^2 (2\nabla\psi \cdot \nabla r + \psi\nabla^2 r + r\nabla^2\psi - r\nabla^2\psi) \\ &= -\hbar^2 (2\nabla\psi \cdot \nabla r + \psi\nabla^2 r) \\ &= -\hbar^2 (2\partial_r + \frac{2}{r}) \psi \\ &= -2\hbar^2 (\partial_r + \frac{1}{r}) \end{aligned}$$

Exercise 4:

$$\langle \vec{r} | (\partial_r + \frac{1}{r}) z | 2_c \rangle$$

$$\begin{aligned}\partial_r (z | 2_c) &= | 2_c \rangle \partial_r z + z \partial_r | 2_c \rangle \\ &= -\frac{z}{a_0} | 2_c \rangle + | 2_c \rangle \partial_r z \\ &= -\frac{z}{a_0} | 2_c \rangle + \frac{z}{r} | 2_c \rangle\end{aligned}$$

$$\begin{aligned}\rightarrow \langle \vec{r} | (\partial_r + \frac{1}{r}) z | 2_c \rangle \\ &= \langle \vec{r} | \left( \frac{z}{r} - \frac{z}{a_0} + \frac{z}{r} \right) | 2_c \rangle \\ &= \langle \vec{r} | \left( \frac{2z}{r} - \frac{1}{a_0} \right) z | 2_c \rangle\end{aligned}$$

Exercise 5:

$$\langle \vec{r} | (r + 2a_0) \partial_z | 2_c \rangle$$

$$\partial_z | 2_c \rangle = \partial_r 2_c \partial_z r = -\frac{1}{a_0} \frac{z}{r} | 2_c \rangle$$

$$\begin{aligned}\langle \vec{r} | (r + 2a_0) \partial_z | 2_c \rangle \\ &= \langle \vec{r} | (r + 2a_0) \left( -\frac{1}{a_0} \frac{z}{r} \right) | 2_c \rangle \\ &= \langle \vec{r} | \left( -\frac{2z}{r} - \frac{2z}{a_0} \right) z | 2_c \rangle\end{aligned}$$

Exercise 6:

$$\begin{aligned}
 \langle \vec{r} | H_c(r+2a_c) z | \psi_c \rangle &= \langle \vec{r} | (r+2a_c) z H_c | \psi_c \rangle + \frac{2\hbar^2}{\mu a_c} \langle \vec{r} | z | \psi_c \rangle \\
 \frac{\mu a_c}{2\hbar^2} \left( \langle \psi | H_c(r+2a_c) z | \psi_c \rangle - \langle \psi | (r+2a_c) z H_c | \psi_c \rangle \right) \\
 &= \langle \psi | z | \psi_c \rangle = z_{\psi_0} \\
 &= \frac{\mu a_c}{2\hbar^2} \left( \langle \psi | (H_c(r+2a_c) z - (r+2a_c) z H_c) | \psi_c \rangle \right) \\
 &= \frac{\mu a_c}{2\hbar^2} \left( E_\psi \langle \psi | (r+2a_c) z | \psi_c \rangle - E_c \langle \psi | (r+2a_c) z | \psi_c \rangle \right) \\
 &= \frac{\mu a_c}{2\hbar^2} (E_\psi - E_c) \langle \psi | (r+2a_c) z | \psi_c \rangle
 \end{aligned}$$

$$\begin{aligned}
 \rightarrow \\
 \delta E_2 &= e^2 \mathcal{E}^2 \sum_{\psi \neq 0} \frac{|z_{\psi_0}|^2}{E_{\psi_0}} \\
 &= -e^2 \mathcal{E}^2 \sum_{\psi \neq 0} \frac{\mu a_c}{2\hbar^2} \langle \psi_c | z | \psi \rangle \langle \psi | (r+2a_c) z | \psi_c \rangle \\
 &= -\frac{e^2 \mathcal{E}^2 \mu a_c}{2\hbar^2} \left( \langle \psi_c | z (1 - |\psi_c\rangle \langle \psi_c|) (r+2a_c) z | \psi_c \rangle \right)
 \end{aligned}$$

$$\delta E_2 = -\frac{e^2 \mathcal{E}^2 \mu a_c}{2\hbar^2} \left[ \langle \psi_c | (r+2a_c) z^2 | \psi_c \rangle - \langle \psi_c | z | \psi_c \rangle \langle \psi_c | (r+2a_c) z | \psi_c \rangle \right]$$

$$\delta E_2 = -\frac{e^2 \mathcal{E}^2 \mu a_c}{2\hbar^2} \langle \psi_c | z^2 (r+2a_c) | \psi_c \rangle$$

$$\delta E_2 = -\frac{e^2 \mathcal{E}^2 \mu a_c}{6\hbar^2} \langle \psi_c | r^2 (r+2a_c) | \psi_c \rangle$$

Exercise 7: Since  $\psi_0$  is spherically symmetric,  $\hat{z}$  is the same as  $\hat{x}, \hat{y}$  so  $\langle \psi_0 | \hat{z} | \psi_0 \rangle$

$$= \frac{1}{3} \langle \psi_0 | x^2 + y^2 + z^2 | \psi_0 \rangle$$

$$= \frac{1}{3} \langle \psi_0 | r^2 | \psi_0 \rangle$$

$$\langle \psi_0 | r^2 (r + 2a_0) | \psi_0 \rangle$$

$$= 4\pi \frac{1}{4a_0^3} \int_0^\infty dr r^2 e^{-2r/a_0} r^2 (r + 2a_0)$$

$$= \frac{27a_0^3}{2}$$

$$SE_2 = - \frac{qe^2 \mu a_0^4 E^2}{4 \hbar^2} = -q\pi \epsilon_0 a_0^3 E^2$$