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1) Bremsstrahlung

a) Let e be the particle's charge and (et I(t) be the particle's velocity at time t. Since the particle has no acceleration for too and for t > T, and an acceleration of -a for oxtxT, ve have $a(t) = \dot{v}(t) = -a\theta(t(T-t))$ Which we can Ferrier transform to get v(u) $\dot{v}(\omega) = -\alpha \int \Theta(t(T-t)) e^{i\omega t} dt$

$$\dot{v}(\omega) = -\alpha \int_{0}^{\infty} e^{i\omega t} dt$$

$$\dot{v}(\omega) = -\alpha \left[\frac{e^{i\omega t}}{i\omega} \right]_{0}^{\infty}$$

$$\dot{v}(\omega) = -\frac{\alpha}{i\omega} \left[e^{i\omega T_{2}} - e^{-i\omega T_{2}} \right]$$

$$\dot{v}(\omega) = -2\alpha e^{i\omega T_{2}} \left[e^{i\omega T_{2}} - e^{-i\omega T_{2}} \right]$$

$$\dot{v}(\omega) = -2\alpha e^{i\omega T_{2}} \frac{\sin(\omega T_{2})}{\omega}$$

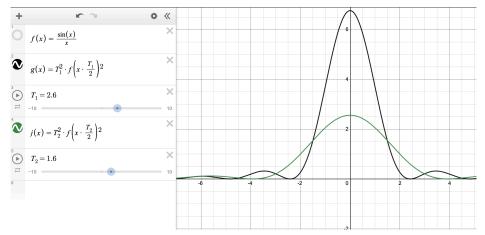
$$\rightarrow |\dot{v}(u)|^2 = \alpha^2 T^2 \sin c^2 (U T_2)$$

where $\operatorname{Sinc}(x) = \frac{\sin(x)}{x}$

So We have

$$\frac{dU}{d\omega} = \frac{2e^2a^2T^2}{3\pi c^3} \sin^2(\omega T_2)$$

A Sketch of the function is below



The radiated frequencies are concentrated around $\omega = 0$ and decay to roughly zero for $\omega > 2\pi$. Thus, the spectrum is dominated by frequencies ω in -2π (ω < 2π .

If we wanted to produce radiation of higher frequency we would decrease T.

The total radiated energy is given by

$$U = \int_{0}^{8} d\omega \, dU = \int_{0}^{8} dt \, P(t)$$

$$U = \frac{2e^{2}a^{2}T^{2}}{3\pi c^{3}} \int_{0}^{8} d\omega \, \sin^{2}(\omega T_{2})$$

$$U = \frac{2e^{2}a^{2}T^{2}}{3\pi c^{3}} \int_{0}^{8} d\omega \, \frac{\sin^{2}(\omega T_{2})}{\omega^{2}T_{2}^{2}}$$

$$X = \omega T_{2}, \, dx = T_{2} d\omega$$

$$U = \frac{2e^{2}a^{2}T^{2}}{3\pi c^{3}} \cdot \frac{Z}{T} \int_{0}^{8} dx \, \frac{\sin^{2}(x)}{x^{2}}$$

$$U = \frac{2e^{2}a^{2}T^{2}}{3\pi c^{3}} \cdot \frac{\pi}{Z}$$

$$U = \frac{2e^{2}a^{2}}{3c^{3}} \cdot \frac{\pi}{Z}$$

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But since v(t) is zero for to T, we have

$$U = \int_{0}^{\infty} dt P(t) = \int_{0}^{\infty} dt P(t)$$
 $J = \frac{1}{3} \frac{d^{2}}{3} \frac{d^{2}$

$$P = \frac{dU}{dT} = \frac{2ea^2}{3c^3}$$

Which is precisely the Larmor formula.

2) Pulsar

$$T = 7.5 c$$
, $\dot{T} = 8 \times 10^{-11}$
 $M = 3 \times 10^{33} g$, $R = 10^6 cm$

a) $P = \frac{2}{3c^3} \left(\dot{\vec{m}}(t) \right)^2$

If we assume the star loses energy due purely to magnetic dipole radiation, then

 $-\frac{dU}{dt} = P$ where $P = \frac{2}{3c^3} \dot{\vec{m}}(t)^2$

and $U = \frac{1}{2} \cdot \frac{\Omega^2}{2}$, $T = 2 \cdot \frac{10}{5} \cdot \frac{R^2}{5}$

Since $\dot{\vec{m}}(t)$ lies in the xy-plane we can parameterize it by $\dot{\vec{m}}(t) = m \left(\cos \left(\frac{2\pi}{100} t \right), \sin \left(\frac{2\pi}{100} t \right), a \right)$

Since
$$T \ll 1$$
, we can assume $\vec{m} \approx -\Omega^2 \vec{m}$ so that $D = 2 \vec{m}^2 = 2 \Omega^4 m^2$

$$P = \frac{2}{3C^3} \dot{\vec{m}}^2 = \frac{2}{3C^3} \Omega^4 m^2$$

and
$$\Omega(t) = \frac{2\pi}{T + \dot{T}t}$$

Next, We have

$$U(t) = \frac{MR^2}{5} \Omega^2$$

$$S_0 - \frac{du}{dt} = 2 \underline{MR}^2 \Omega(t) \dot{\Omega}(t)$$

$$\hat{P} = -\frac{du}{dt}$$
, we have

$$\frac{2}{3 c^3} \Omega^4 M^2 = -\frac{2}{5} M R^2 \Omega(t) \dot{\Omega}(t)$$

$$\Rightarrow M^2 = -\frac{3}{5} M R^2 c^3 \frac{\Omega(t)}{\Omega^3}$$

$$M^2 = \frac{3}{5} M R^2 c^3 \frac{\dot{\Gamma} \Gamma(t)}{4\pi^2}$$

$$\Rightarrow m^2 \approx \frac{3}{20\pi^2} M R^2 c^3 + T$$

$$\vec{B}(\vec{x}) = 3(\hat{x} \cdot \vec{m}) \hat{x} - \vec{m} + O(\vec{r}_{4})$$
Since this quantity is maximized
$$\vec{f}_{er} \quad \hat{x} || \vec{m}, \quad \text{we have}$$

$$\vec{B}_{max}(R) = \frac{2M}{R^{3}}$$

$$B_{\text{max}} = \frac{4 \text{ m}^2}{R^6}$$

$$B_{\text{max}} = \int \frac{3}{5\pi^2} \frac{\text{M}^2}{R^4} c^3 \overrightarrow{T} \overrightarrow{T}$$

Uhich seems quite wrong since I can't get the units to work out.