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Problem 1. Consider two observables Q and R , such that $[Q, R] = 0$. In this case the uncertainty principle says that $\sigma_Q \sigma_R \geq 0$. Is it *always* the case that $\sigma_Q \sigma_R = 0$? If yes, prove it; if not, give a counterexample.

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1)

Suppose $Q = R = S_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$

And let $\begin{pmatrix} a \\ b \end{pmatrix} \in \mathbb{C}^2$ be a

vector in the spin- z basis

Where $|a|^2 + |b|^2 = 1$.

We compute $\overline{\sigma_{S_y}} \overline{\sigma_{S_y}} = \overline{\sigma_{S_y}^2}$

$$\overline{\sigma_{S_y}^2} = \langle S_y^2 \rangle - \langle S_y \rangle^2$$

$$= \frac{\hbar^2}{4} \left((a^* \ b^*) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} - \left((a^* \ b^*) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} \right)^2 \right)$$

$$= \frac{\hbar^2}{4} \left(1 + (ab^* - ba^*)^2 \right)$$

Which is nonzero if, for example,
 $a = b$.

This shows that while

$[Q, R] = 0$ implies $\sqrt{a} \sqrt{R} \geq 0$,

$\sqrt{Q} \sqrt{R}$ is not always zero.

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Problem 2. Prove the generalized uncertainty relation

$$\Delta P \Delta Q \geq \frac{1}{4} |\langle [P, Q] \rangle|^2 + \frac{1}{4} |\langle [P, Q]_+ \rangle - 2\langle P \rangle \langle Q \rangle|^2, \quad (1)$$

where $[A, B]_+ := AB + BA$ is the anticommutator.

Claim: $\Delta P \Delta Q \geq \frac{1}{4} |\langle [P, Q] \rangle|^2$
 $+ \frac{1}{4} |\langle [P, Q]_+ \rangle - 2\langle P \rangle \langle Q \rangle|^2$

Proof:

We know from class that

for a pair of observables P, Q ,

$$\langle P^2 \rangle \langle Q^2 \rangle \geq |\langle PQ \rangle|^2 \checkmark$$

and that this is preserved

✓ under $P \rightarrow P - \langle P \rangle, Q \rightarrow Q - \langle Q \rangle$.

So we have

$$\begin{aligned} \Delta P \Delta Q &= \langle (P - \langle P \rangle)^2 \rangle \langle (Q - \langle Q \rangle)^2 \rangle \\ &\geq |\langle (P - \langle P \rangle)(Q - \langle Q \rangle) \rangle|^2 \end{aligned}$$

We proceed by rewriting the right hand side of the eq. above.

$$|\langle PQ - \langle P \rangle Q - P \langle Q \rangle + \langle P \rangle \langle Q \rangle \rangle|^2$$

$$= |\langle PQ \rangle - \langle P \rangle \langle Q \rangle|^2$$

(since $\langle \rangle$ is linear)

We now use the fact that

$$PQ = \frac{1}{2} [P, Q] + \frac{1}{2} [P, Q]_+ \quad \checkmark$$

$$= \left| \frac{1}{2} \langle [P, Q] \rangle \checkmark + \frac{1}{2} (\langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle) \right|^2$$

$$= \frac{1}{4} \left| \langle [P, Q] \rangle + \langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle \right|^2 \quad \checkmark$$

$$= \frac{1}{4} \left(\langle [P, Q] \rangle + \langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle \right) \cdot \left(\langle [P, Q] \rangle^* + \langle [P, Q]_+ \rangle^* - 2 \langle P \rangle \langle Q \rangle \right)$$

Since $\langle P \rangle, \langle Q \rangle \in \mathbb{R}$.

$$\begin{aligned}
= & \frac{1}{4} \left(|\langle [P, Q] \rangle|^2 + \langle [P, Q]_+ \rangle \langle [P, Q]_+ \rangle^* \right. \\
& + \langle [P, Q]_+ \rangle \langle [P, Q] \rangle^* \\
& + \langle [P, Q]_+ \rangle^* \langle [P, Q] \rangle \\
& - 2 \langle P \rangle \langle Q \rangle \langle [P, Q] \rangle^* \\
& - 2 \langle P \rangle \langle Q \rangle \langle [P, Q]_+ \rangle^* \quad \leftarrow \\
& + 4 \langle P \rangle^2 \langle Q \rangle^2 \quad \leftarrow \\
& - 2 \langle P \rangle \langle Q \rangle \langle [P, Q] \rangle \\
& \left. - 2 \langle P \rangle \langle Q \rangle \langle [P, Q]_+ \rangle \right) \quad \leftarrow
\end{aligned}$$

Now since $[P, Q] = PQ - QP$

and $\langle PQ \rangle = \langle QP \rangle^*$,

We have $\langle [P, Q] \rangle = - \langle [P, Q] \rangle^*$

So the mess above simplifies

to :

$$\begin{aligned}
&= \frac{1}{4} \left(|\langle [P, Q] \rangle|^2 \right. \\
&\quad + \langle [P, Q]_+ \rangle \langle [P, Q]_+ \rangle^* \\
&\quad - 2 \langle P \rangle \langle Q \rangle \langle [P, Q]_+ \rangle^* \\
&\quad + 4 \langle P \rangle^2 \langle Q \rangle^2 \\
&\quad \left. - 2 \langle P \rangle \langle Q \rangle \langle [P, Q]_+ \rangle \right)
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{4} |\langle [P, Q] \rangle|^2 \\
&\quad + \frac{1}{4} \left(\langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle \right) \\
&\quad \cdot \left(\langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle \right)^*
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{4} |\langle [P, Q] \rangle|^2 \\
&\quad + \frac{1}{4} |\langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle|^2
\end{aligned}$$

So

$$\begin{aligned}
V_P V_Q \geq & \frac{1}{4} |\langle [P, Q] \rangle|^2 \\
& + \frac{1}{4} |\langle [P, Q]_+ \rangle - 2 \langle P \rangle \langle Q \rangle|^2
\end{aligned}$$

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Problem 3. Any Hermitian operator P that satisfies $P^2 = P$ is called a **projection operator** or **projector**.

Suppose that P_1 and P_2 are projectors. Then show that the product $P_1 P_2$ is a projector if and only if $[P_1, P_2] = 0$.

Claim: If P_1 and P_2 are projectors,
then $P_1 P_2$ is a projector
if and only if $[P_1, P_2] = 0$

Proof:

(i) Suppose $P_1 P_2$ is a projector. ✓

I'm not sure how to prove this
implication and Wikipedia says
it is not necessarily true.

$$P_1 P_2 = P_2 P_1 \text{ via Hermiticity} \leadsto [P_1, P_2] = 0$$

(ii) Now Suppose $[P_1, P_2] = 0$

Then

$$P_1 P_2 = P_2 P_1$$

and

$$(P_1 P_2)^2 = (P_2 P_1)^2$$

$$(P_1 P_2)^2 = P_2 P_1 P_2 P_1 = P_1 P_2^2 P_1$$

$$\text{but } P_1 = P_1^2 \text{ and } P_2 = P_2^2$$

$$\text{and } P_1 P_2 = P_2 P_1, \text{ so}$$

$$(P_1 P_2)^2 = P_1 P_2$$

So $P_1 P_2$ is a projector. ✓



Problem 4. For the operator

$$A = \epsilon(|1\rangle\langle 1| - |0\rangle\langle 0|) + \gamma|0\rangle\langle 1| + \gamma^*|1\rangle\langle 0| \quad (2)$$

on the Hilbert space $\{|0\rangle, |1\rangle\}$, where $\epsilon \geq 0$ and $\gamma \in \mathbb{C}$, show that the eigenvectors may be written as

$$\begin{aligned} |+\rangle &= \sin \theta |0\rangle + e^{i\phi} \cos \theta |1\rangle \\ |-\rangle &= \cos \theta |0\rangle - e^{i\phi} \sin \theta |1\rangle, \end{aligned} \quad (3)$$

where

$$\tan 2\theta = \frac{|\gamma|}{\epsilon} \quad \left(0 \leq \theta < \frac{\pi}{2}\right). \quad (4)$$

Also find the corresponding eigenvalues.

4)

First, we identify $\{|0\rangle, |1\rangle\}$
with $\left\{\begin{pmatrix} 0 \\ 1 \end{pmatrix}, \begin{pmatrix} 1 \\ 0 \end{pmatrix}\right\}$ since
it makes A easier to deal
with. (In my opinion)

So

$$A = \begin{pmatrix} \epsilon & \gamma^* \\ \gamma & -\epsilon \end{pmatrix}$$

Which has eigenvalues given

by the solutions, λ , to the equation:

$$\lambda^2 - \epsilon^2 - |\gamma|^2 = 0$$

$$\text{So } \lambda = \pm \sqrt{\epsilon^2 + |\gamma|^2} \checkmark$$

Then for some $\vec{a}, \vec{b} \in \text{span}\left\{\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \end{pmatrix}\right\}$

We have

$$A \vec{a} = \sqrt{\epsilon^2 + |\gamma|^2} \vec{a}$$

and

$$A \vec{b} = -\sqrt{\epsilon^2 + |\gamma|^2} \vec{b}$$

letting $\vec{a} = \begin{pmatrix} a \\ 1 \end{pmatrix}$, $\vec{b} = \begin{pmatrix} 1 \\ -a^* \end{pmatrix}$, $a \in \mathbb{C}$

Note: It is clear that A is Hermitian
so we know that $\langle \vec{a}, \vec{b} \rangle = 0$,
and that we are free to
scale eigenvectors

To compute a we solve
the following equation:

$$\begin{pmatrix} \epsilon & \gamma^* \\ \gamma & -\epsilon \end{pmatrix} \begin{pmatrix} a \\ 1 \end{pmatrix} = \sqrt{\epsilon^2 + |\gamma|^2} \begin{pmatrix} a \\ 1 \end{pmatrix}$$

This yields

$$a = \frac{-\gamma}{\epsilon - \sqrt{\epsilon^2 + |\gamma|^2}}$$

so the eigenvectors of A
are

$$\left\{ \begin{pmatrix} \frac{-\gamma}{\epsilon - \sqrt{\epsilon^2 + |\gamma|^2}} \\ 1 \end{pmatrix}, \begin{pmatrix} \frac{1}{\gamma^*} \\ \epsilon - \sqrt{\epsilon^2 + |\gamma|^2} \end{pmatrix} \right\}$$

with eigen values $\sqrt{\epsilon^2 + |\gamma|^2}$, $-\sqrt{\epsilon^2 + |\gamma|^2}$
respectively.

Since we are free to scale \vec{a}
and \vec{b} , we will use the new
eigenbasis:

$$\left\{ \begin{pmatrix} e^{i\phi} \\ -\frac{\epsilon + \sqrt{\epsilon^2 + |\gamma|^2}}{|\gamma|} \end{pmatrix}, \begin{pmatrix} \frac{-\epsilon + \sqrt{\epsilon^2 + |\gamma|^2}}{|\gamma|} \\ -e^{-i\phi} \end{pmatrix} \right\}$$

Where $\phi = \text{Arg}(\gamma)$ ✓

Now let $X = \frac{-\epsilon + \sqrt{\epsilon^2 + |\gamma|^2}}{|\gamma|}$

the eigenvectors become

$$\left\{ \begin{pmatrix} e^{i\phi} \\ X \end{pmatrix}, \begin{pmatrix} X \\ -\bar{e}^{i\phi} \end{pmatrix} \right\}$$

we make one more rescaling

$$\left\{ \begin{pmatrix} \frac{e^{i\phi}}{\sqrt{1+X^2}} \\ \frac{X}{\sqrt{1+X^2}} \end{pmatrix}, \begin{pmatrix} \frac{-X}{\sqrt{1+X^2}} e^{i\phi} \\ \frac{1}{\sqrt{1+X^2}} \end{pmatrix} \right\}$$

Now, using the substitution,

$\tan(2\theta) = \frac{|\gamma|}{\epsilon}$, we will write \vec{a} and \vec{b} in terms of θ .

From the substitution and the fact that

$$\frac{2 \tan(\theta)}{1 - \tan^2(\theta)} = \tan(2\theta),$$

We have

$$\theta = \arctan \left(\frac{-\epsilon + \sqrt{\epsilon^2 + |\chi|^2}}{|\chi|} \right)$$

$$\theta = \arctan(x)$$

since $\frac{1}{\sqrt{1+x^2}} = \cos(\arctan(x))$

and $\frac{x}{\sqrt{1+x^2}} = \sin(\arctan(x))$ ✓

the eigenvectors of A become

$$\left\{ \begin{pmatrix} e^{i\phi} \cos(\theta) \\ \sin(\theta) \end{pmatrix}, \begin{pmatrix} -e^{i\phi} \sin(\theta) \\ \cos(\theta) \end{pmatrix} \right\}$$

with eigenvalues $\sqrt{\epsilon^2 + |\chi|^2}$, $-\sqrt{\epsilon^2 + |\chi|^2}$

which is what was to be shown.

In Dirac notation, $\vec{a} = |+\rangle$, $\vec{b} = |-\rangle$

$$\{ \sin(\theta) |0\rangle + e^{i\phi} \cos(\theta) |1\rangle,$$

$$\cos(\theta) |0\rangle - e^{i\phi} \sin(\theta) |1\rangle \}$$

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Problem 5. Both the Poisson bracket and the quantum commutator satisfy the following properties (these aren't hard to verify; so if they aren't obvious to you, then go ahead and verify them).

1. antisymmetry: $[A, B] = -[B, A]$
2. null bracket with scalars: $[c, A] = 0$ ($c \in \mathbb{C}$)
3. associativity: $[A + B, C] = [A, B] + [B, C]$; $[A, B + C] = [A, B] + [A, C]$
4. product rules: $[AB, C] = A[B, C] + [A, C]B$; $[A, BC] = [A, B]C + B[A, C]$
5. Jacobi identity: $[A, [B, C]] + [B, [C, A]] + [C, [A, B]] = 0$

Suppose that we haven't yet settled on a quantum bracket $[A, B]$, but we wish to choose one that satisfies the above properties. Show that the *only* possible choice is the quantum commutator $[A, B] = AB - BA$ (up to an overall real factor).

Hint: starting with $[AA', BB']$ for arbitrary operators A, A', B , and B' , apply the two product rules to write this expression as the sum of four brackets. Applying the product rules in either order gives two expressions; you should require that they be equivalent.

5)

We will show that we must have $[A, B] = AB - BA$

for two arbitrary operators, A, B .
If $[,]$ is to satisfy the relations above.

Let A, A', B, B' be arbitrary operators and consider $[AA', BB']$

By applying the first and then the second product rule, we have

$$\begin{aligned} [AA', BB'] &= A[A', BB'] + [A, BB']A' \\ &= A([A', B]B' + B[A', B']) \quad (1) \\ &\quad + ([A, B]B' + B[A, B'])A' \checkmark \end{aligned}$$

Applying the second and then the first yields

$$\begin{aligned} [AA', BB'] &= [AA', B]B' + B[AA', B'] \\ &= (A[A', B] + [A, B]A')B' \quad (2) \\ &\quad + B(A[A', B'] + [A, B']A') \checkmark \end{aligned}$$

But (1) must equal (2) so

$$\begin{aligned} & (A[A', B] + [A, B]A')B' \\ & + B(A[A', B'] + [A, B']A') \\ & = A([A', B]B' + B[A', B']) \\ & + ([A, B]B' + B[A, B'])A' \end{aligned}$$

simplifying a little...

$$\begin{aligned} & A[A', B]B' + [A, B]A'B' \\ & + BA[A', B'] + B[A, B']A' \\ & - A[A', B]B' - AB[A', B'] \\ & - [A, B]B'A' - B[A, B']A' = 0 \end{aligned}$$

cancelling some terms, we have

$$\begin{aligned} & [A, B](A'B' - B'A') \quad \checkmark \\ & = (AB - BA)[A', B'] \quad (*) \end{aligned}$$

Since this must be true for

all operators A, A', B, B' ,

we can fix A' , B' and then
each side of (*) only depends
on A and B . Since the
action of each side is the same
for all test functions, f ,
we must have

$$[A, B] = AB - BA \quad \checkmark$$

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