

$$1) \quad \epsilon_{\mu\alpha\beta} \epsilon_{\mu\sigma\tau} = \delta_{\alpha\sigma} \delta_{\beta\tau} - \delta_{\alpha\tau} \delta_{\beta\sigma}$$

Let $a, b, c \in \mathbb{R}^3$, then

$$(b \times c)_i = \epsilon_{ijk} b_j c_k$$

$$\begin{aligned} (a \times (b \times c))_i &= \epsilon_{ijk} \epsilon_{k\ell n} a_j b_\ell c_n \\ &= \epsilon_{kij} \epsilon_{k\ell n} a_j b_\ell c_n \end{aligned}$$

$$(b(a \cdot c) - c(a \cdot b))_i$$

$$= b_i a_k c_k - c_i a_k b_k$$

$$= a_j b_\ell c_n \delta_{jn} \delta_{i\ell} - a_j b_\ell c_n \delta_{i\ell} \delta_{jn}$$

$$= a_j b_\ell c_n (\delta_{jn} \delta_{i\ell} - \delta_{i\ell} \delta_{jn})$$

$$= a_j b_\ell c_n (\delta_{i\ell} \delta_{jn} - \delta_{in} \delta_{j\ell})$$

and since a, b, c were arbitrary,

we must have

$$\epsilon_{kij} \epsilon_{k\ell n} = \delta_{i\ell} \delta_{jn} - \delta_{in} \delta_{j\ell}$$

2)

We first recall the action of J^2, J_z, J_{\pm} on $|j m\rangle$

$$J^2 |j m\rangle = j(j+1) \hbar^2 |j m\rangle$$

$$J_z |j m\rangle = m \hbar |j m\rangle$$

$$J_{\pm} |j m\rangle = \hbar \underbrace{\sqrt{j(j+1) - m(m \pm 1)}}_{\lambda_{jm \pm}} |j m \pm 1\rangle$$

$$J_{\pm} = J_x \pm i J_y$$

$$J_x = \frac{1}{2} (J_+ + J_-)$$

$$J_y = \frac{1}{2i} (J_+ - J_-)$$

$$\langle J_x \rangle = \langle j m | J_x | j m \rangle$$

$$\langle J_x \rangle = \frac{1}{2} \langle j m | J_+ | j m \rangle + \frac{1}{2} \langle j m | J_- | j m \rangle$$

$$\langle J_x \rangle = \frac{1}{2} (\lambda_{jm+} \langle j m | j m+1 \rangle + \lambda_{jm-} \langle j m | j m-1 \rangle)$$

$$\boxed{\langle J_x \rangle = 0}$$

$$\begin{aligned}
\langle J_x^2 \rangle &= \langle j m | J_x^2 | j m \rangle \\
&= \frac{\hbar^2}{4} \left(\langle j m | J_+^2 + J_-^2 + J_+ J_- + J_- J_+ | j m \rangle \right) \\
&= \frac{\hbar^2}{4} \left(\langle j m | J_+ J_- | j m \rangle + \langle j m | J_- J_+ | j m \rangle \right) \\
&= \frac{\hbar^2}{4} \left(\sqrt{j(j+1) - m(m-1)} \sqrt{j(j+1) - (m-1)m} \right. \\
&\quad \left. + \sqrt{j(j+1) - m(m+1)} \sqrt{j(j+1) - (m+1)m} \right)
\end{aligned}$$

$$\langle J_x^2 \rangle = \frac{\hbar^2}{2} (j^2 + j - m^2)$$

$$\sigma_{J_x} = \sqrt{\langle J_x^2 \rangle - \langle J_x \rangle^2} = \sqrt{\langle J_x^2 \rangle}$$

$$\sigma_{J_x} = \frac{\hbar}{\sqrt{2}} \sqrt{j^2 + j - m^2}$$

$$\begin{aligned}
 \langle J_y^2 \rangle &= -\frac{1}{4} \langle jm | -J_+ J_- | jm \rangle \\
 &\quad - \frac{1}{4} \langle jm | -J_- J_+ | jm \rangle \\
 &= \langle J_x^2 \rangle
 \end{aligned}$$

So

$$\sigma_{J_y} = \sigma_{J_x}$$

Therefore,

$$\sigma_{J_x} \sigma_{J_y} = \frac{\hbar^2}{2} (j^2 + j - m^2)$$

The uncertainty principle states that

$$\sigma_{J_x} \sigma_{J_y} \geq \frac{1}{2} |\langle [J_x, J_y] \rangle|$$

$$\text{and } [J_x, J_y] = i\hbar J_z$$

So

$$\sigma_{J_x} \sigma_{J_y} \geq \frac{\hbar^2}{2} |\langle J_z \rangle| = \frac{\hbar^2}{2} |m|$$

Since $-j \leq m \leq j$,

$$m^2 \leq j^2.$$

And since $j \geq 0$,

$$|m| \leq j$$

So

$$|m| + m^2 \leq j + j^2$$

and

$$j^2 + j - m^2 \geq |m|$$

Therefore,

$$\sigma_{J_x} \sigma_{J_y} = \frac{\hbar^2}{2} (j^2 + j - m^2) \geq \frac{\hbar^2}{2} |m|$$

so the uncertainty principle is satisfied.

3)

$$df(r, \theta, \phi) = d\vec{r} \cdot \nabla f(\vec{r})$$

in Spherical coordinates

$$d\vec{r} = dr + r d\phi + r^2 \sin\theta d\theta$$

and

$$df(r, \theta, \phi) = \partial_r f dr + \partial_\phi f d\phi + \partial_\theta f d\theta$$

So

$$(\nabla f \cdot \hat{r}) dr = \partial_r f dr$$

$$(\nabla f \cdot \hat{\phi}) r d\phi = \partial_\phi f d\phi$$

$$(\nabla f \cdot \hat{\theta}) r^2 \sin\theta d\theta = \partial_\theta f d\theta$$

so, since \hat{r} , $\hat{\phi}$, and $\hat{\theta}$ form an orthogonal coordinate system, we must have

$$\nabla = \partial_r \hat{r} + \frac{1}{r} \partial_\phi \hat{\phi} + \frac{1}{r^2 \sin\theta} \partial_\theta \hat{\theta}$$