

TAM 470 - HW 5 Solutions

Problem 1

$$\theta''(t) + \frac{g}{l} \theta(t) = 0 \quad - \quad (3.1)$$

$$\theta(0) = \theta_0 \quad \text{and} \quad \theta'(0) = \omega_0 \quad - \quad (3.2)$$

(a) Let $\theta'(t) = \beta(t) \Rightarrow \theta''(t) = \beta'(t)$

Using these in (3.1) and (3.2),

$$\beta'(t) + \frac{g}{l} \theta(t) = 0$$

$$\theta(0) = \theta_0 \quad \text{and} \quad \beta(0) = \omega_0.$$

In the matrix form,

$$\begin{bmatrix} \theta'(t) \\ \beta'(t) \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{g}{l} & 0 \end{bmatrix} \begin{bmatrix} \theta(t) \\ \beta(t) \end{bmatrix}$$

$$\text{with } \theta(0) = \theta_0 \quad \text{and} \quad \beta(0) = \omega_0.$$

(b) Let λ be the eigen values.

$$\det |A - I\lambda| = 0 \quad \text{and} \quad A = \begin{bmatrix} 0 & 1 \\ -\frac{g}{l} & 0 \end{bmatrix}$$

$$\Rightarrow \begin{vmatrix} 0 - \lambda & 1 \\ -\frac{g}{l} & 0 - \lambda \end{vmatrix} = 0$$

$$\Rightarrow \lambda^2 + \frac{g}{l} = 0$$

or

$$\lambda = i\sqrt{\frac{g}{l}}, \quad -i\sqrt{\frac{g}{l}}$$

(c) The eigen values are purely imaginary
 \Rightarrow The forward Euler method is unstable. Refer to Figure 4.2 in Moyn

(d) The backward Euler scheme is unconditionally stable for all eigenvalues with a non-positive real part.

(e) Referring to Moyn, Page 59 and 60, when $\lambda_R = 0$,
 $A = B \Rightarrow |T| = 1$. This is called neutral stability (i.e. no error in amplitude).
Can also interpret as unconditionally stable w.r.t amplitude

(f) (i) Forward Euler formulation

$$\begin{bmatrix} \frac{\theta_{n+1} - \theta_n}{h} \\ \frac{\beta_{n+1} - \beta_n}{h} \end{bmatrix} = [A] \begin{bmatrix} \theta_n \\ \beta_n \end{bmatrix}$$

$$\Rightarrow \begin{bmatrix} \theta_{n+1} \\ \beta_{n+1} \end{bmatrix} = ([I] + h[A]) \begin{bmatrix} \theta_n \\ \beta_n \end{bmatrix}$$

(ii) Backward Euler formulation:

$$\begin{bmatrix} \frac{\theta_{n+1} - \theta_n}{h} \\ \frac{\beta_{n+1} - \beta_n}{h} \end{bmatrix} = [A] \begin{bmatrix} \theta_{n+1} \\ \beta_{n+1} \end{bmatrix}$$

$$\begin{bmatrix} \theta_{n+1} \\ \beta_{n+1} \end{bmatrix} - h[A] \begin{bmatrix} \theta_{n+1} \\ \beta_{n+1} \end{bmatrix} = \begin{bmatrix} \theta_n \\ \beta_n \end{bmatrix}$$

$$\Rightarrow ([I] - h[A]) \begin{bmatrix} \theta_{n+1} \\ \beta_{n+1} \end{bmatrix} = \begin{bmatrix} \theta_n \\ \beta_n \end{bmatrix}$$

(iii) Trapezoidal rule formulation:-

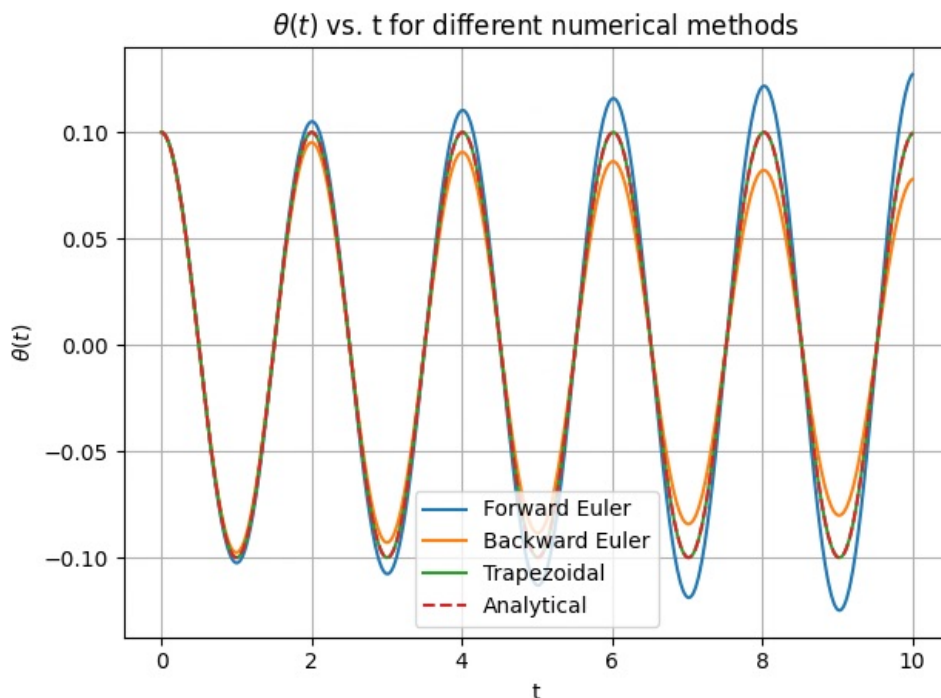
$$\begin{bmatrix} \frac{\theta_{n+1} - \theta_n}{h} \\ \frac{\beta_{n+1} - \beta_n}{h} \end{bmatrix} = [A] \begin{bmatrix} \frac{\theta_{n+1} + \theta_n}{2} \\ \frac{\beta_{n+1} + \beta_n}{2} \end{bmatrix}$$

Rearranging,

$$\left[[I] - \frac{h}{2} [A] \right] \begin{bmatrix} \theta_{n+1} \\ \beta_{n+1} \end{bmatrix} = \left[[I] + \frac{h}{2} [A] \right] \begin{bmatrix} \theta_n \\ \beta_n \end{bmatrix}$$

where $[I]$ is the (2×2) identity matrix.

(g) The $\theta(t)$ vs t plot is shown below.



- The amplitude of forward euler is increasing which demonstrates unstable behavior.
 - We see the dampening behavior of Backward Euler. (Stable but not accurate).
 - Trapezoid rule captures the oscillatory behavior and is stable. The plot overlaps with analytical solution the most.
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Problem 2

$$y' = e^{\sin y} - ty, \quad y(0) = 1$$

$$f(y, t) = e^{\sin y} - ty$$

$$\frac{\partial f}{\partial y} = f_y = e^{\sin y} \cdot \cos y - t$$

(a) Trapezoidal rule,

$$\frac{y_{n+1} - y_n}{h} = \frac{f(y_{n+1}, t_{n+1}) + f(y_n, t_n)}{2}$$

⇒ Bring all terms to LHS,

$$F(y_{n+1}) = y_{n+1} - y_n - \frac{h}{2} f(y_{n+1}, t_{n+1}) - \frac{h}{2} f(y_n, t_n) = 0$$

$$F(y_{n+1}) = y_{n+1} - y_n - \frac{h}{2} (e^{\sin y_{n+1}} - t_{n+1} y_{n+1} + e^{\sin y_n} - t_n y_n) = 0$$

$$(b) DF(y_{n+1}) = \frac{\partial F}{\partial y_{n+1}} = 1 - \frac{h}{2} f_y(y_{n+1}, t_{n+1})$$

$$DF(y_{n+1}) = 1 - \frac{h}{2} (e^{\sin y_{n+1}} \cdot \cos y_{n+1} - t_{n+1})$$

(c) Newton-Raphson:

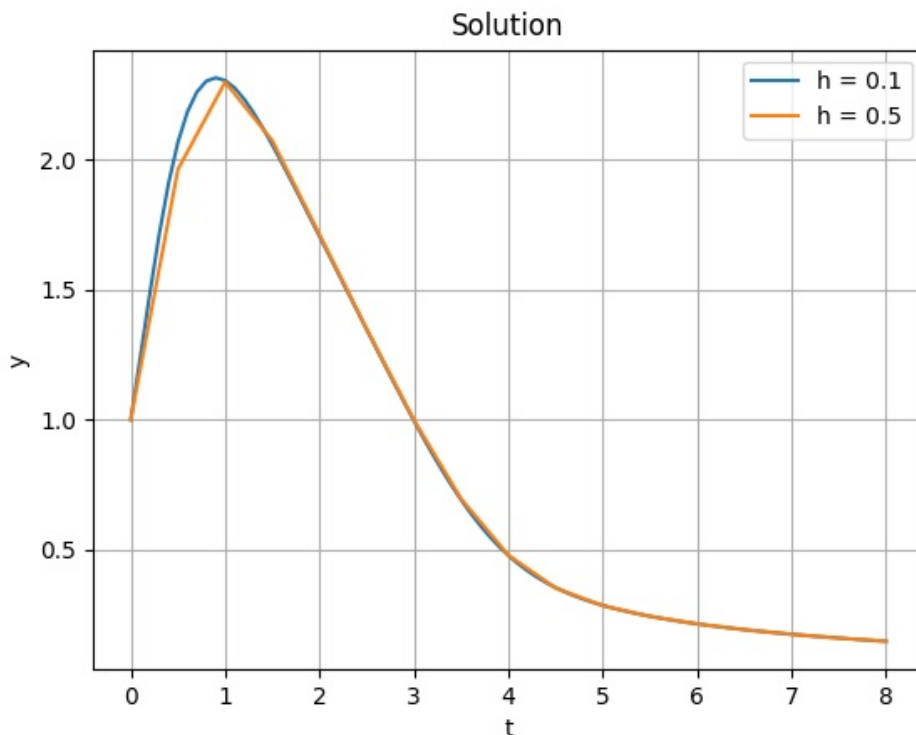
$$y_{n+1}^{k+1} = y_{n+1}^k - \frac{F(y_{n+1}^k)}{DF(y_{n+1}^k)}$$

$$\therefore y_{n+1}^{k+1} = y_{n+1}^k - \frac{\left(y_{n+1}^k - y_n - \frac{h}{2} \left(e^{\sin y_{n+1}^k} - t_{n+1} y_{n+1}^k + e^{\sin y_n} - t_n y_n \right) \right)}{1 - \frac{h}{2} \left(e^{\sin y_{n+1}^k} \cos y_{n+1}^k - t_{n+1} \right)}$$

Note that y_n is not updated and we can just write y_n instead of y_n^k as it does not change with 'k' iteration counter.

(d) Implement on PL

(e)



Problem 3

$$y' = \lambda(y - \cos t) - \sin t, \quad y(0) = y_0$$

$$y_{\text{exact}} = e^{\lambda t} (y_0 - 1) + \cos t$$

(a) Let the RHS be $f = \lambda(y - \cos t) - \sin t$

$$\text{Trapezoid rule} \rightarrow \frac{y_{n+1} - y_n}{h} = \frac{1}{2} (f_n + f_{n+1})$$

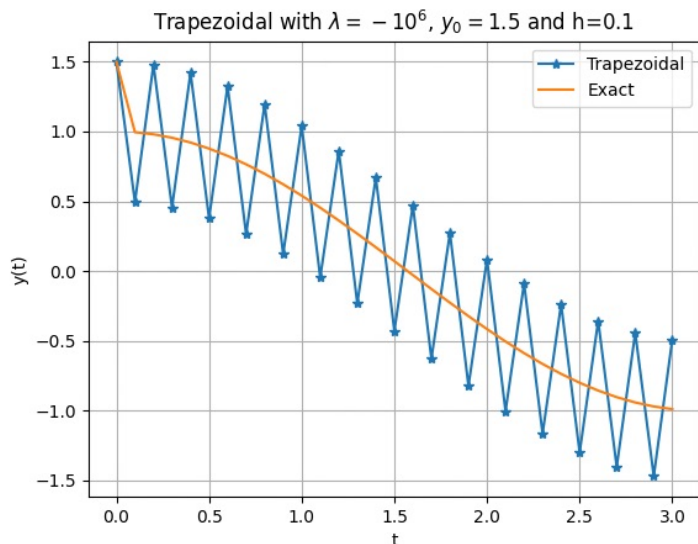
$$\therefore \frac{y_{n+1} - y_n}{h} = \frac{1}{2} \left[\lambda (y_{n+1} + y_n - \cos t_n - \cos t_{n+1}) - \sin t_n - \sin t_{n+1} \right]$$

Rearranging.

$$y_{n+1} \left[1 - \frac{\lambda h}{2} \right] = \left[1 + \frac{\lambda h}{2} \right] y_n - \frac{\lambda h}{2} (\cos t_n + \cos t_{n+1}) - \frac{h}{2} (\sin t_n + \sin t_{n+1})$$

(b) We use the expression from (a) with $\lambda = -10^6$, $h = 0.1$, $t \in [0, 3]$ and $y_0 = 1.5$

Plot shown below.



$$(c) |h\lambda| \gg 1 \quad \text{and} \quad |\lambda| \gg 1$$

Recall trapezoidal formula. from (a),

$$\cancel{2h} y_{n+1} \left[\frac{1}{h\lambda} - \frac{1}{2} \right] = h \left[\left(\frac{1}{h\lambda} + \frac{1}{2} \right) y_n - \frac{1}{2} (\cos t_n + \cos t_{n+1}) - \frac{1}{2\lambda} (\sin t_n + \sin t_{n+1}) \right]$$

$$|h\lambda| \gg 1 \Rightarrow \frac{1}{|h\lambda|} \ll 1 \rightarrow 0$$

$$|\lambda| \gg 1 \Rightarrow \frac{1}{|\lambda|} \ll 1 \rightarrow 0$$

Using above approximations,

$$-\frac{y_{n+1}}{2} = \frac{y_n}{2} - \frac{1}{2} (\cos t_n + \cos t_{n+1})$$

$$\Rightarrow \underline{\underline{y_{n+1} \approx -y_n + \cos t_n + \cos t_{n+1}}}$$

(d) From (c),

$$y_{n+1} - \cos t_{n+1} \approx -y_n + \cos t_n$$

$$n=0, y_1 - \cos t_1 \approx -y_0 + 1 = (-1)(y_0 - 1)$$

$$n=1, y_2 - \cos t_2 \approx -y_1 + \cos t_1 = (-1)(y_1 - \cos t_1) \\ (-1)^2 (y_0 - 1)$$

$$n=2, y_3 - \cos t_3 \approx -y_2 + \cos t_2 = (-1)[y_2 - \cos t_2] \\ (-1)^3 (y_0 - 1)$$

When $n = (k-1)$, where k is some integer,

$$y_k - \cos t_k \approx (-1)^k (y_0 - 1)$$

or

$$y_n \approx (-1)^n (y_0 - 1) + \cos t_n$$

(e) The oscillating behavior is due to the term $(-1)^n (y_0 - 1)$ which changes sign for each increment of n .

The oscillation will not occur with $y_0 = 1$ as this makes that oscillatory term 0 and $y_n \approx \cos t_n$.
