

Some specialized deformations we discussed last class are translations & rotations which are special cases of homogeneous deformations ( $F$  is a constant field).

We know that translations and rotations preserve distances, i.e.

$$\xrightarrow{(1)}$$

$$\| \underline{p} - \underline{q} \| = \| \underline{f}(\underline{p}) - \underline{f}(\underline{q}) \| \quad \forall \underline{p}, \underline{q} \in B,$$

whenever  $f$  is a translation or a rotation. You may now ask if a converse statement is true:

Is every deformation that preserves distance, a rotation or a translation?

- True.

Theorem: The following are equivalent

- (a)  $f$  satisfies ①
- (b) There exists a constant  $R \in \text{Orth}^+$ , and  $\underline{y} \in V$  such that

$$\underline{f}(\underline{p}) = \underline{y} + R \underline{p}.$$

We will now see how a deformation changes volumes of infinitesimal volume elements. An infinitesimal volume element at a point  $\underline{p} \in \mathcal{B}$  is given by the three mutually orthogonal basis vectors  $\underline{\varepsilon}_u, \underline{\varepsilon}_v, \underline{\varepsilon}_w$  :

$$\underline{\varepsilon}_u \cdot (\underline{\varepsilon}_v \times \underline{\varepsilon}_w) = \varepsilon^3 \underline{u} \cdot (\underline{v} \times \underline{w}) \quad \text{--- (2)}$$

(2) is transformed to

$$\begin{aligned} F(\underline{\varepsilon}_u) \cdot (F(\underline{\varepsilon}_v) \times F(\underline{\varepsilon}_w)) &= \varepsilon^3 F\underline{u} \cdot (F\underline{v} \times F\underline{w}) \\ &= \varepsilon^3 (\det F) \underline{u} \cdot (\underline{v} \times \underline{w}) \end{aligned}$$

Therefore, the volume is scaled by  $\det F$ .

### Example

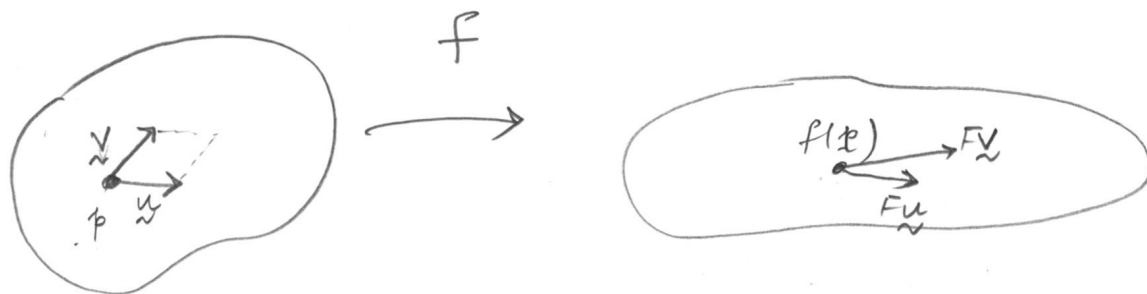
Consider the shear deformation

$$F(\underline{p}) = \begin{bmatrix} 1 & \gamma & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix}$$

$$(x_1, x_2, x_3) \mapsto F(\underline{p}) = \begin{bmatrix} 1 & \gamma & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix}$$

$\det F \equiv 1 \Rightarrow$  Shear deformation is volume preserving.

Let us now see how areas change due to deformation. (3)



Reference area :  $\underline{u} \times \underline{v} \rightarrow$  A vector whose magnitude describes area and direction " the normal.

Deformed area :  $\underline{F_u} \times \underline{F_v}$

$$\begin{aligned} (\underline{F_u} \times \underline{F_v})_i &= \epsilon_{ijk} (F_u)_j (F_v)_k \\ &= \epsilon_{ijk} F_{jL} u_L F_{kM} v_M \quad - (3) \end{aligned}$$

Multiplying (3) by  $F_{iN}$ , we have

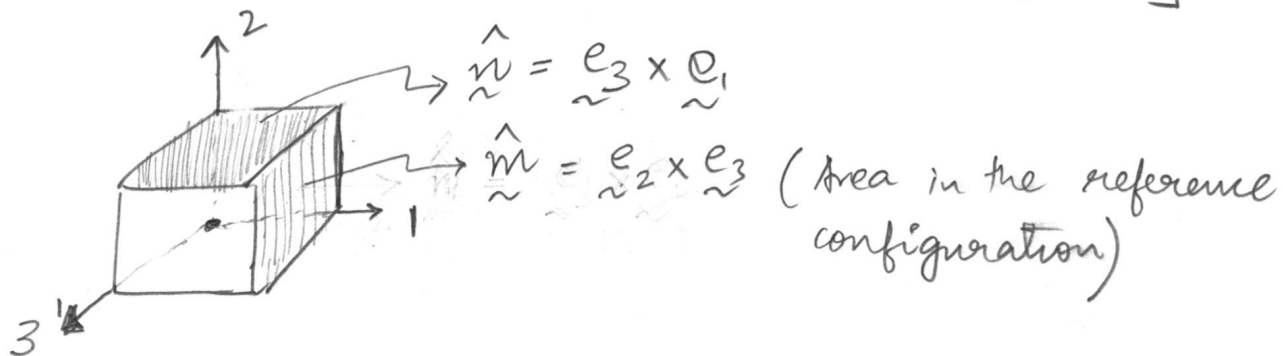
$$\begin{aligned} \left[ F^T (\underline{F_u} \times \underline{F_v}) \right]_N &= \underbrace{\epsilon_{ijk} F_{iN} F_{jL} F_{kM}}_{(\det F) \epsilon_{NLM}} u_L v_M \\ &= (\det F) (\underline{u} \times \underline{v})_N \end{aligned}$$

$$\therefore \boxed{\underline{F_u} \times \underline{F_v} = (\det F) F^{-T} (\underline{u} \times \underline{v})} \quad - (4)$$

(4)

Eq (4) is referred to as Nanson's formula, named after British mathematician Edward J. Nanson (1850-1936).

Example: Shear deformation  $f(\underline{p}) = \begin{bmatrix} 1 & \gamma & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix}$ .



$$F \underline{e}_3 \times F \underline{e}_1 = (\det F) F^{-T} \hat{\underline{n}}$$

$$= F^{-T} \hat{\underline{n}} \quad (\because \text{shear is volume preserving})$$

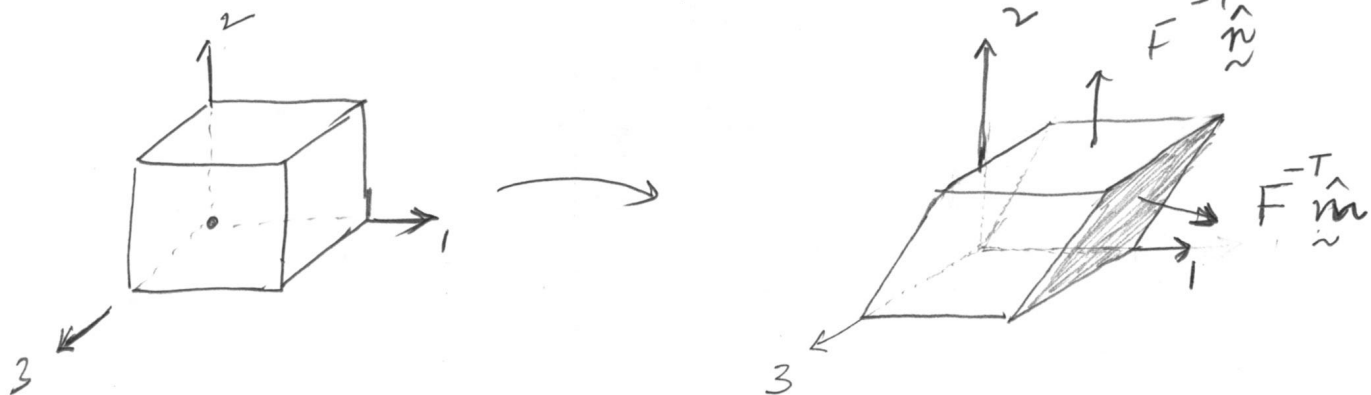
$$= \begin{bmatrix} 1 & \gamma & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}^{-T} \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}$$

$$= \begin{bmatrix} 1 & 0 & 0 \\ -\gamma & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}$$

$$= \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}$$

$$F \underline{e}_2 \times F \underline{e}_3 = \begin{bmatrix} 1 & 0 & 0 \\ -\gamma & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ -\gamma \\ 0 \end{bmatrix}$$

(5)



The unit area described by  $\hat{\mathbf{n}}$  transform to an area of magnitude  $\sqrt{1 + \mathbf{r}^2}$ , and direction along  $\begin{bmatrix} 1 \\ -\mathbf{r} \\ 0 \end{bmatrix}$ .

### Material and spatial tensor fields

A deformation maps the reference configuration (also referred to as a body) to a deformed configuration. Any property of the body may now be described using either the reference configuration or the deformed configuration. Consider a scalar field  $g$  that describes temperature.

$g: \mathcal{f}(\mathcal{B}) \rightarrow \mathbb{R}$ $g(\underline{x})$	vs	$\tilde{g}: \mathcal{B} \rightarrow \mathbb{R}$ $\tilde{g}(\underline{\mathbf{p}})$
$\downarrow$	vs	$\downarrow$
$\underline{\text{Relation:}} \quad g(\mathcal{f}(\underline{\mathbf{p}}))$	=	$\tilde{g}(\underline{\mathbf{p}})$
$\downarrow$		$\downarrow$
Spatial / Eulerian description.		Material / reference / Lagrangian description