# A FAMILY OF NATURAL EQUILIBRIUM MEASURES FOR SINAI BILLIARD FLOWS

#### JÉRÔME CARRAND

ABSTRACT. The Sinai billiard flow on the two-torus, i.e., the periodic Lorentz gaz, is a continuous flow, but it is not everywhere differentiable. Assuming finite horizon, we relate the equilibrium states of the flow with those of the Sinai billiard map T – which is a discontinuous map. We propose a definition for the topological pressure  $P_*(T,g)$  associated to a potential g. We prove that for any piecewise Hölder potential g satisfying a mild assumption,  $P_*(T,g)$  is equal to the definitions of Bowen using spanning or separating sets. We give sufficient conditions under which a potential gives rise to equilibrium states for the Sinai billiard map. We prove that in this case the equilibrium state  $\mu_g$  is unique, Bernoulli, adapted and gives positive measure to all nonempty open sets. For this, we make use of a well chosen transfer operator acting on anisotropic Banach spaces, and construct the measure by pairing its maximal eigenvectors. Last, we prove that the flow invariant probability measure  $\bar{\mu}_g$ , obtained by taking the product of  $\mu_g$  with the Lebesgue measure along orbits, is Bernoulli and flow adapted. We give examples of billiard tables for which there exists an open set of potentials satisfying those sufficient conditions.

#### 1. Introduction

1.1. Billiards and equilibrium states. In this work, we are concerned with a class of dynamics with singularities: the dispersing billiards introduced by Sinai [30] on the two-torus. A Sinai billiard on the torus is the periodic case of the planar Lorentz gaz (1905) model for the motion of a single dilute electron in a metal. The scatterers (corresponding to atoms of the metals) are assumed to be strictly convex (but not necessarily discs). Such billiards have become fundamental models in mathematical physics.

To be more precise, a Sinai billiard table Q on the two-torus  $\mathbb{T}^2$  is a set  $Q = \mathbb{T}^2 \setminus B$  with  $B = \bigsqcup_{i=1}^D B_i$  for some finite number  $D \geqslant 1$  of pairwise disjoint closed domains  $B_i$ , called scatterers, with  $C^3$  boundaries having strictly positive curvature – in particular, the scatterers are strictly convex. The billiard flow  $\phi_t$  is the motion of a point particle travelling at unit speed in Q with specular reflections off the boundary of the scatterers. Identifying outgoing collisions with incoming ones in the phase space, the billiard flow is continuous. However, the grazing collisions – those tangential to scatterers – give rise to singularities in the derivative [13]. The Sinai billiard map T – also called collision map – is the return map of the single

Date: November 22, 2022.

JC is grateful to ITS-ETH Zurich for their invitation in May 2022. Thanks to Viviane Baladi, Mark Demers and Alexey Korepanov for useful discussions and comments. Research supported by the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (grant agreement No 787304).

point particle to the scatterers. Because of the grazing collisions, the Sinai billiard map is a discontinuous map.

Sinai billiard maps and flows both preserve smooth invariant probability measures, respectively  $\mu_{\rm SRB}$  and  $\tilde{\mu}_{\rm SRB}$ , which have been extensively studied:  $(T, \mu_{\rm SRB})$  and  $(\phi_t, \tilde{\mu}_{\rm SRB})$  are uniformly hyperbolic, ergodic, K-mixing [30, 9, 31], and Bernoulli [20, 15]. The measure  $\mu_{\rm SRB}$  is T-adapted [24] in the sense of the integrability condition:

$$\int |\log d(x, \mathcal{S}_{\pm 1})| \, \mathrm{d}\mu_{\mathrm{SRB}} < \infty,$$

where  $S_{\pm 1}$  is the singularity set for  $T^{\pm 1}$ . Both systems enjoy exponential decay of correlations [33, 18]. Since the billiard has many periodic orbits, it thus has many other ergodic invariant measures, but until very recently most of the results apply to perturbations of  $\mu_{\text{SRB}}$  [12, 17].

In the case of an Anosov flow, it is known since the work of Bowen [7] that the Kolmogorov-Sinai entropy is upper-semicontinuous, which guarantees the existence of measures of maximal entropy, or more generally, of equilibrium states. Because of the singularities, billiard flows are not Anosov and therefore methods used in the context of Anosov flows cannot be applied easily. The upper-semicontinuity of the entropy is not known at the moment, and, more generally, the existence of equilibrium states has to be treated one potential at the time.

In a recent paper, Baladi and Demers [3] proved, under a mild technical assumption and assuming finite horizon, that there exists a unique measure of maximal entropy  $\mu_*$  for the billiard map, and that  $\mu_*$  is Bernoulli, T-adapted, charges all nonempty open sets and does not have atoms. Their construction of this measure relies on the use of a transfer operator acting on anisotropic Banach spaces, similar to those used by [18] in order to study  $\mu_{\rm SRB}$ . Combining their work with those of Lima–Matheus [25] and Buzzi [11], Baladi and Demers proved that their exists a positive constant C such that

(1.1) 
$$Ce^{h_*m} \leqslant \#\operatorname{Fix} T^m, \quad \forall m \geqslant 1,$$

where #Fix  $T^m$  denotes the number of fixed points of  $T^m$ , and  $h_*$  is the topological entropy of the map T from [3]. Baladi and Demers also give a condition under which  $\mu_*$  and  $\mu_{\text{SRB}}$  coincide.

In a subsequent paper, Baladi and Demers [4] constructed a family of equilibrium states  $\mu_t$  for T associated to the family of geometric potentials  $-t \log J^u T$ , where  $J^u T$  is the unstable Jacobian of T and  $t \in (0, t_*)$  for some  $t_* > 1$ . In the case t = 1,  $\mu_t = \mu_{\text{SRB}}$ . The construction again relies on the use of a family of transfer operators  $\mathcal{L}_t$  acting on anisotropic Banach spaces. For each  $t \in (0, t_*)$ , they proved that  $\mu_t$  is the unique equilibrium state associated with the potential  $-t \log J^u T$ , that  $\mu_t$  is mixing, T-adapted, has full support and does not have atoms. Baladi and Demers also showed that each transfer operator  $\mathcal{L}_t$  has a spectral gap, from which they deduced the exponential rate of mixing for each measure  $\mu_t$ , for  $C^1$  observables.

Even more recently, Demers and Korepanov [16] proved a polynomial decay of correlations for the measure  $\mu_*$  for Hölder observables.

In this paper, we give a sufficient condition under which a piecewise Hölder potential g admits equilibrium states for T. Under this assumption, we prove that the equilibrium state is in fact unique, Bernoulli, T-adapted and charges all nonempty open sets. We prove that its lift into a flow invariant measure is Bernoulli and flow-adapted. We also identify the potential

 $g = -h_{\text{top}}(\phi_1)\tau$  to be such that its corresponding equilibrium states for T – whenever they exist – are in bijection with measures of maximal entropy of the billiard flow.

Notice that the geometric potentials  $-t \log J^u T$  are not piecewise Hölder, and thus the work of Baladi and Demers [4] on those potentials is distinct from ours.

1.2. Statement of main results – Organization of the paper. Since transfer operator techniques are simpler to implement for maps than for flows, we will be concerned with the associated billiard map  $T: M \to M$  defined to be the first collision map on the boundary of Q, where  $M = \partial Q \times [-\pi/2, \pi/2]$ . We assume as in [33, 3], that the billiard table Q has finite horizon, in the sense that the billiard flow does not have any trajectories making only tangential collisions – in particular, this implies that the return time function  $\tau$  to a scatterers is bounded.

The first step is to find a suitable notion of topological pressure  $P_*(T,g)$  for the discontinuous map T and a potential  $g: M \to \mathbb{R}$ . In order to define it, we introduce as in [3], the following increasing family of partition of M. Let  $\mathcal{P}$  be the partition into maximal connected sets on which both T and  $T^{-1}$  are continuous, and let  $\mathcal{P}_{-k}^n = \bigvee_{i=-k}^n T^{-i}\mathcal{P}$ . Then the sequence  $\sum_{P \in \mathcal{P}_0^n} \sup_P e^{S_n g}$  is submultiplicative, where  $S_n g = \frac{1}{n} \sum_{i=0}^{n-1} g \circ T^i$  is the Birkhoff sum of g. We can thus define the topological pressure by

**Definition 1.1.** 
$$P_*(T,g) \coloneqq \lim_{n \to +\infty} \frac{1}{n} \log \sum_{P \in \mathcal{P}_0^n} \sup_{P} e^{S_n g}$$

Section 2 is dedicated to the study of this quantity. In particular, we prove (Proposition 2.2) that whenever the potential g is smooth enough – piecewise Hölder – and  $P_*(T,g)$  –  $\sup g > 0$  then  $P_*(T,g)$  coincides with both Bowen's definitions using spanning sets and separating sets. We also prove (Lemma 2.4) that for each T-invariant measure  $\mu$ , we have  $P_*(T,g) \ge h_{\mu}(T) + \int g \, d\mu$ . Finally, we show that if  $g = -h_{\rm top}(\phi_1)\tau$  admits an equilibrium state  $\mu_g$ , then the measure  $\bar{\mu}_g = (\int \tau \, \mathrm{d}\mu_g)^{-1} \mu_g \otimes \lambda$  is a measure of maximal entropy for the billiard flow, seen as a suspension flow over T, where  $\lambda$  is the Lebesgue measure in the flow direction.

To state our existence results (in Section 6), we need to quantify the recurrence to the singular set. Fix an angle  $\varphi_0$  close to  $\pi/2$  and  $n_0 \in \mathbb{N}$ . We say that a collision is  $\varphi_0$ -grazing if its angle with the normal is larger than  $\varphi_0$  in absolute value. Let  $s_0 = s_0(\varphi_0, n_0) \in (0, 1]$  denote the smallest number such that

(1.2) any orbit of length  $n_0$  has at most  $s_0 n_0$  collisions which are  $\varphi_0$ -grazing.

Due to the finite horizon condition, we can choose  $\varphi_0$  and  $n_0$  such that  $s_0 < 1$ . We refer to [3, §2.4] for further discussion on this quantity. From [13],  $\Lambda = 1 + \kappa_{\min} \tau_{\min} > 1$  is the expanding factor in the hyperbolicity of T, where  $\kappa_{\min}$  is the minimal curvature of the scatterers and  $\tau_{\min}$  is the minimum of the return time function  $\tau$ . Define  $\mathcal{S}_0 = \{(r, \varphi) \in M \mid |\varphi| = \pi/2\}$  the set of grazing collisions, and  $\mathcal{S}_{\pm n} = \bigcup_{i=0}^n T^{\mp i} \mathcal{S}_0$  the singular set of  $T^{\pm n}$ . Call  $\mathcal{N}_{\varepsilon}(\cdot)$  the  $\varepsilon$ -neighbourhood of a set. Then

**Theorem 1.2.** If g is a bounded, piecewise Hölder potential such that  $P_*(T, g)$ —sup  $g > s_0 \log 2$  and  $\log \Lambda > \sup g$  — inf g, then there exists a probability measure  $\mu_g$  such that

(i)  $\mu_g$  is T-invariant, T-adapted and for all  $k \in \mathbb{Z}$ , there exists  $C_k > 0$  such that  $\mu_g(\mathcal{N}_{\varepsilon}(\mathcal{S}_k)) \leqslant C_k |\log \varepsilon|^{-\gamma}$ .

- (ii)  $\mu_g$  the unique equilibrium state of T under g: that is  $P_*(T,g) = h_{\mu_g}(T) + \int g \, d\mu_g$  and  $P_*(T,g) > h_{\mu}(T) + \int g \, d\mu$  for all  $\mu \neq \mu_g$ .
- (iii)  $\mu_g$  is Bernoulli<sup>1</sup> and charges all nonempty open sets.

If the assumption  $\log \Lambda > \sup g - \inf g$  is weakened into the condition SSP.1 (as defined above Lemma 3.3), then item (i) still holds. If the assumption  $\log \Lambda > \sup g - \inf g$  is weakened into the condition SSP.2 (as defined above Corollary 3.6), then items (i), (ii) and (iii) hold.

The above theorem will follow from Proposition 6.1, Lemma 6.2, Corollary 6.14, and Propositions 6.18, 6.15, 6.10. Furthermore, assuming the sparse recurrence to singularities condition from [3], we provide in Remark 3.11 an open set of potentials, each having SSP.1 and SSP.2.

The tool used to construct the measure  $\mu_g$  is a transfer operator  $\mathcal{L}_g$  with  $\mathcal{L}_g f = (f e^g/J^s T) \circ T^{-1}$ , similar to the one used in [3] corresponding to the case  $g \equiv 0$ . This operator and the anisotropic Banach spaces on which it acts are defined in details in Section 4. Section 3 contains key combinatorial growth lemmas, controlling the growth in weighted complexity of the iterates of a stable curve. These lemmas will be crucial since the quantity they control appears in the norms of the iterates of  $\mathcal{L}_g$ . In Section 5, we prove a (degenerated) "Lasota–Yorke" type inequality (Proposition 5.1) – giving an upper bound on the spectral radius of  $\mathcal{L}_g$  – as well as a lower bound on the spectral radius (Theorem 5.3).

Section 6 is devoted to the construction and the properties of the measure  $\mu_g$ . From the estimates on the norms from the previous section, we are able to construct left and right maximal eigenvectors ( $\tilde{\nu}$  and  $\nu$ ) for  $\mathcal{L}_g$ . We construct the measure  $\mu_g$  by pairing these eigenvectors. We then prove the estimates on the measure of a neighbourhood of the singular sets (Lemma 6.2). Section 6.3 contains the key result of the absolute continuity of the stable and unstable foliations with respect to  $\mu_g$ , as well as the proof that  $\mu_g$  has total support – this is done by exploiting the  $\nu$ -almost everywhere positive length of unstable manifolds from Section 6.2. In Section 6.4, we show that  $\mu_g$  is ergodic, from which we bootstrap to K-mixing using Hopfargument. Adapting [15] with modifications from [3], we deduce from the hyperbolicity and the K-mixing that  $\mu_g$  is Bernoulli. Still in Section 6.4, we give an upper-bound on the measure of weighted Bowen balls, from which we deduce, using the Shannon–MacMillan–Breiman theorem, that  $\mu_g$  is an equilibrium state for T under the potential g (Corollary 6.14). Finally, the Section 6.5 is dedicated to the uniqueness of the equilibrium state  $\mu_g$ .

In Section 7, we prove using arguments from [13] that  $(\phi_t, \bar{\mu}_g)$  is K-mixing (Proposition 7.1), and again, using the hyperbolicity of the billiard flow, we adapt [15] in order to prove that  $(\phi_t, \bar{\mu}_g)$  is Bernoulli (Proposition 7.2). Finally, we prove that  $\bar{\mu}_g$  is flow adapted in the sense of the integrability condition formulated in Proposition 7.4. We summarize this results about the billiard flow in the following theorem.

**Theorem 1.3.** Let g be a potential satisfying the assumptions from Theorem 1.2, and let  $\bar{\mu}_g := (\int \tau \, \mathrm{d}\mu_g)^{-1} \mu_g \otimes \lambda$ . Then  $\bar{\mu}_g$  is a  $\phi_t$ -invariant Borel probability measure that is an equilibrium states for any potential  $\tilde{g}$  such that  $g = \lambda(\tilde{g}) - P(\phi_1, \tilde{g})\tau$ , where  $\lambda(\tilde{g})(x) = \int_0^{\tau(x)} \tilde{g}(\phi_t(x)) \, \mathrm{d}t$ . Furthermore,  $\bar{\mu}_g$  is flow adapted and  $(\phi_t, \bar{\mu}_g)$  is Bernoulli.

In a work in preparation with Baladi and Demers [2], we bootstrap from the results of the present paper to show that if  $h_{\text{top}}(\phi_1)\tau_{\text{min}} > s_0 \log 2$  then the potential  $-h_{\text{top}}(\phi_1)\tau$  satisfies

<sup>&</sup>lt;sup>1</sup>Recall that Bernoulli implies K-mixing, which implies strong mixing, which implies ergodic.

the sufficient assumptions SSP.1 and SSP.2 in our Theorem 1.2, thus constructing a measure of maximal entropy for the billiard flow. This is done by studying the family of potentials  $-t\tau$  and proving that the maximal value  $t_{\infty}$  of t such that  $-t'\tau$  has SSP.1 and SSP.2 for all  $0 \le t' \le t$ , satisfies  $t_{\infty} > h_{\text{top}}(\phi_1)$ . Indeed, recalling Remark 3.11 and Corollary 2.6, for every small enough |t|,  $-t\tau$  has SSP.1 and SSP.2, and the case  $t = h_{\text{top}}(\phi_1)$  corresponds to measures of maximal entropy for the billiard flow.

### 2. Topological Pressure, Variational Principle and Abramov Formula

In this section, we formulate definitions of topological pressure for the billiard map, and prove that – under some conditions – they are equivalent. Using a classical estimate, we then prove one direction of the variational principle. Finally, making use of the Abramov formula, we relate equilibrium states of T with the ones of the billiard flow. More specifically, we identified the potential for T which is related to – up to existence – the measures of maximal entropy of  $\phi_t$ .

We first introduce notation: Adopting the standard coordinates  $x=(r,\varphi)$  on each connected component  $M_i$  of

$$M := \partial Q \times \left[ -\frac{\pi}{2}, \frac{\pi}{2} \right] = \bigsqcup_{i=1}^{D} \partial B_i \times \left[ -\frac{\pi}{2}, \frac{\pi}{2} \right],$$

where r denotes arclength along  $\partial B_i$ ,  $\varphi$  is the angle the post-collisional trajectory makes with the normal to  $\partial B_i$  and  $M_i = \partial B_i \times \left[-\frac{\pi}{2}, \frac{\pi}{2}\right]$ . In these coordinates, the collision map  $T: M \to M$  preserve a smooth invariant probability measure  $\mu_{\text{SRB}}$  given by  $d\mu_{\text{SRB}} = (2|\partial Q|)^{-1}\cos\varphi \,drd\varphi$ .

We now define the sets where T and its iterates are discontinuous. Let  $S_0 := \{(r, \varphi) \in M \mid |\varphi| = \pi/2\}$  denote the set of grazing collisions. For each nonzero  $n \in \mathbb{N}$ , let

$$\mathcal{S}_{\pm n} \coloneqq \bigcup_{i=0}^n T^{\mp i} \mathcal{S}_0,$$

denote the singularity set for  $T^{\pm n}$ . It would be natural to study the map T restricted to the invariant set  $M \setminus \bigcup_{n \in \mathbb{Z}} \mathcal{S}_n$  where T is continuous, however the set of curves  $\bigcup_{n \in \mathbb{Z}} \mathcal{S}_n$  is dense in M [13, Lemma 4.55]. We thus introduce the classical family of partitions of M as follows.

For  $k, n \ge 0$ , let  $\mathcal{M}^n_{-k}$  denote the partition of  $M \setminus (\mathcal{S}_{-k} \cup \mathcal{S}_n)$  into its maximal connected components. Note that all elements of  $\mathcal{M}^n_{-k}$  are open sets on which  $T^i$  is continuous, for all  $-k \le i \le n$ . Since the thermodynamic sums over elements of  $\mathcal{M}^n_0$  of a potential g will play a key role in the estimates on the norms of the iterates of the transfer operator  $\mathcal{L}_g$  in Section 5, it should be natural – by analogy to the case of continuous maps – to define the topological pressure from these sums.

Another natural family of partitions is given as follows. Let  $\mathcal{P}$  denote the partition of M into maximal connected components on which both T and  $T^{-1}$  are continuous. Define  $\mathcal{P}_{-k}^n = \bigvee_{i=-k}^n T^{-i}\mathcal{P}$  and remark that  $T^i$  is continuous on every element of  $\mathcal{P}_{-k}^n$ , for all  $-k \leq i \leq n$ .

The interior of each element of  $\mathcal{P}$  corresponds to precisely one element of  $\mathcal{M}_{-1}^1$ , but its refinements  $\mathcal{P}_{-k}^n$  may also contain some isolated points – this happens if three or more scatterers have a common grazing collision. These partitions already appeared in the work of Baladi and

Demers, where they proved [3, Lemma 3.2] that the number of isolated points in  $\mathcal{P}_{-k}^n$  grows linearly in n + k.

Finally, denote  $\mathring{\mathcal{P}}_{-k}^n$  the collection of interior of elements of  $\mathcal{P}_{-k}^n$ . In [3, Lemma 3.3], Baladi and Demers proved that  $\mathring{\mathcal{P}}_{-k}^n = \mathcal{M}_{-k-1}^{n+1}$ , for all  $n, k \geq 0$ . It should be natural that the topological pressures obtained from these three families of partitions coincide. This is the content of Theorem 2.1.

In order to formulate the result on the equivalence between definitions of topological pressure for T, we need to be more specific about the definition of *piecewise* Hölder.

We say that a function g is  $(\mathcal{M}_0^1, \alpha)$ -Hölder,  $0 < \alpha < 1$ , if g is  $\alpha$ -Hölder continuous on each element of the partition  $\mathcal{M}_0^1$ . We define the  $C^{\alpha}$  norm  $|g|_{C^{\alpha}}$  of g to be the maximum, over all connected components U of the domain of continuity of g, of the usual  $C^{\alpha}$  norm  $|g|_{C^{\alpha}(U)}$ , that

 $|g|_{C^{\alpha}} = \max\{|g|_{C^{0}(U)} + H_{U}^{\alpha}(g) \mid U \text{ connected set on which } g \text{ is continuous}\},$ 

where

$$H_U^{\alpha}(g) = \sup_{x,y \in U} \frac{|g(x) - g(y)|}{d(x,y)^{\alpha}}.$$

Similarly, we say that a function g is  $\mathcal{M}_0^1$ -continuous if g is bounded and continuous on each element of the partition  $\mathcal{M}_0^1$ . We define the  $C^0$  norm  $|g|_{C^0}$  to be the maximum over all connected components U of the domain of continuity of g, of the usual  $C^0$  norm, that is

 $|g|_{C^0} = \max\{|g|_{C^0(U)} \mid U \text{ connected set on which } g \text{ is continuous}\}.$ 

**Theorem 2.1.** Let  $g: M \to \mathbb{R}$  be a potential bounded from above. Then

$$\lim_{n \to +\infty} \frac{1}{n} \log \sum_{A \in \mathcal{P}_n^n} \sup_{x \in A} e^{(S_n g)(x)} =: P_*(T, g)$$

exists and is called the pressure of g. Moreover, the map  $g \mapsto P_*(T,g)$  is convex.

When g is  $\mathcal{M}_0^1$ -continuous and  $P_*(T,g) - \sup g > 0$ , the following limits exist and are equal to  $P_*(T,g)$ .

- (i)  $\lim_{n \to +\infty} \frac{1}{n} \log \sum_{A \in \mathring{\mathcal{P}}_0^n} \sup_{x \in A} e^{(S_n g)(x)},$ (ii)  $\lim_{n \to +\infty} \frac{1}{n} \log \sum_{A \in \mathcal{M}_0^n} \sup_{x \in A} e^{(S_n g)(x)}.$

Furthermore, when g is also  $(\mathcal{M}_0^1, \alpha)$ -Hölder continuous, then the following limits are equal to  $P_*(T,g)$ .

- (iii)  $\lim_{n \to +\infty} \frac{1}{n} \log \sum_{A \in \mathcal{P}_0^n} \inf_{x \in A} e^{(S_n g)(x)},$
- (iv)  $\lim_{n \to +\infty} \frac{1}{n} \log \sum_{A \in \mathring{\mathcal{P}}_0^n} \inf_{x \in A} e^{(S_n g)(x)},$ (v)  $\lim_{n \to +\infty} \frac{1}{n} \log \sum_{A \in \mathcal{M}_0^n} \inf_{x \in A} e^{(S_n g)(x)},$

Finally, the sequence  $n \mapsto \log \sum_{A \in \mathcal{M}_0^n} \sup_{x \in A} e^{S_{n-1}g(x)}$  is subadditive.

**Proposition 2.2.** Let g be a  $\mathcal{M}_0^1$ -continuous potential. Let  $P_{\mathrm{span}}(T,g)$  and  $P_{\mathrm{sep}}(T,g)$  be the pressure obtained using Bowen's definition with, respectively, spanning sets and separating sets. Then  $P_{\mathrm{span}}(T,g) \leqslant P_*(T,g)$  and  $P_{\mathrm{sep}}(T,g) \leqslant P_*(T,g)$ . When  $P_*(T,g) - \sup g > 0$ , then  $P_*(T,g) = P_{\mathrm{sep}}(T,g)$ . Furthermore, when  $P_*(T,g) - \sup g > 0$  and g is  $(\mathcal{M}_0^1,\alpha)$ -Hölder,  $P_*(T,g) = P_{\mathrm{span}}(T,g)$ .

The proof of the last three forms of  $P_*(T,g)$  in Theorem 2.1 relies crucially on the following lemma.

**Lemma 2.3.** For every  $(\mathcal{M}_0^1, \alpha)$ -Hölder continuous potential g there exists a constant  $C_g$  such that for all  $n \ge 1$  and all  $P \in \mathcal{P}_0^n$ ,

$$\sup_{P} e^{S_n g} \leqslant C_g \inf_{P} e^{S_n g}.$$

The estimate still holds, for the same constant  $C_g$ , when  $\mathcal{P}_0^n$  is replaced by  $\mathcal{P}_{-l}^n$ ,  $\mathring{\mathcal{P}}_{-l}^n$  or  $\mathcal{M}_{-l}^n$ , for any fixed  $l \ge 0$ .

Before the proofs of these results, we first recall that T is uniformly hyperbolic in the sense [13] that the cones

(2.1) 
$$\mathcal{C}^{u} := \{ (dr, d\varphi) \in \mathbb{R}^{2} \mid \kappa_{\min} \leqslant d\varphi/dr \leqslant \kappa_{\max} + 1/\tau_{\min} \},$$

$$\mathcal{C}^{s} := \{ (dr, d\varphi) \in \mathbb{R}^{2} \mid -\kappa_{\min} \geqslant d\varphi/dr \geqslant -\kappa_{\max} - 1/\tau_{\min} \},$$

are strictly invariant under DT and  $DT^{-1}$ , respectively, whenever these derivatives exist. Here  $\kappa_{\max}$  is the maximum curvature of the scatterer boundaries,  $\kappa_{\min}$  the minimum, and  $\tau_{\min}$  is the minimum of the return time function  $\tau$ . Furthermore, there exists  $C_1 > 0$  such that for all  $n \ge 1$ ,

$$||D_x T^n(v)|| \geqslant C_1 \Lambda^n ||v||, \forall v \in \mathcal{C}^u, \quad ||D_x T^{-n}(v)|| \geqslant C_1 \Lambda^n ||v||, \forall v \in \mathcal{C}^s,$$

where  $\Lambda = 1 + 2\kappa_{\min}\tau_{\min}$  is the minimum hyperbolicity constant.

*Proof of Lemma 2.3.* Let  $d_n$  denote the Bowen distance, that is the dynamical distance given by

$$d_n(x,y) = \max_{0 \le i \le n} d(T^i x, T^i y),$$

where d(x, y) is the Euclidean metric on each  $M_i$ , with  $d(x, y) = 10D \max_i \operatorname{diam}(M_i)$  if x and y belong to different  $M_i$  (this definition ensure we have a compact set). Let  $\varepsilon_0 > 0$  be as in [3, eq (3.3)], that is: if  $d_n(x, y) < \varepsilon_0$  then x and y lie in the same element of  $\mathcal{M}_0^n$ . Therefore, by the uniform hyperbolicity of T, if  $d(T^i(x), T^i(y)) \leq \varepsilon_0/2$  for all  $|i| \leq n$  then  $d(x, y) \leq C_1 \Lambda^{-n} \varepsilon_0/2$ .

For all integer m and all potential g, define its variation by

$$\operatorname{Var}_m(g, T, \varepsilon) := \sup\{|g(x) - g(y)| \mid d(T^j x, T^j y) \leqslant \varepsilon, |j| \leqslant m\}.$$

When g is  $(\mathcal{M}_0^1, \alpha)$ -Hölder, we get that  $\operatorname{Var}_m(g, T, \frac{\varepsilon_0}{2C_1}) \leqslant C(\frac{\varepsilon_0}{2}\Lambda^{-m})^{\alpha}$ . Therefore

$$\sum_{m \geqslant 0} \operatorname{Var}_m(g, T, \frac{\varepsilon_0}{2C_1}) =: K < \infty.$$

By uniform hyperbolicity of T, there exists  $k_{\varepsilon}$  such that  $\operatorname{diam}(\mathcal{M}_{-k_{\varepsilon}}^{n+1}) < \varepsilon_0/2C_1$  for all  $n \geq k_{\varepsilon}$ . It then follows from the proof of [3, Lemma 3.5] that if x and y lie in the same element of  $\mathcal{P}_{-k_{\varepsilon}}^{k_{\varepsilon}+n}$ , then  $d_n(x,y) \leq \varepsilon_0/2C_1$ , for all  $n \geq 0$ .

Let  $P \in \mathcal{P}_{-k_{\varepsilon}}^{k_{\varepsilon}+n}$  and let  $x, y \in P$ . Let  $0 \leqslant k \leqslant n$ . Then for all  $|j| < m_k := \min(k, n-k)$ ,  $d(T^j(T^kx), T^j(T^ky)) < \varepsilon_0/2C_1$  and so  $|g(T^kx) - g(T^ky)| \leqslant \operatorname{Var}_{m_k}(g, T, \frac{\varepsilon_0}{2C_1})$ . Therefore

$$|S_n g(x) - S_n g(y)| \le 2 \sum_{m=0}^{\lfloor \frac{n}{2} \rfloor + 1} \operatorname{Var}_m(g, T, \frac{\varepsilon_0}{2C_1}) \le 2K < \infty.$$

Now, let  $P \in \mathcal{P}_0^n$  for some  $n \geqslant 2k_{\varepsilon}$ . Notice that  $\mathcal{P}_0^n = \bigvee_{i=k_{\varepsilon}}^{n-k_{\varepsilon}} T^{-i} \mathcal{P}_{-k_{\varepsilon}}^{k_{\varepsilon}}$ , in other words for all l such that  $k_{\varepsilon} \leqslant l \leqslant n - k_{\varepsilon}$ ,  $T^l P$  is included in an element of  $\mathcal{P}_{-k_{\varepsilon}}^{k_{\varepsilon}}$ . Finally, by decomposing each orbit into three parts, we get that for all  $x, y \in P$ ,

$$e^{S_n g(x) - S_n g(y)} = e^{S_{k\varepsilon} g(x) - S_{k\varepsilon} g(y)} e^{S_{n-2k\varepsilon} g(T^{k\varepsilon} x) - S_{n-2k\varepsilon} g(T^{k\varepsilon} y)} e^{S_{k\varepsilon} g(T^{n-k\varepsilon} x) - S_{k\varepsilon} g(T^{n-k\varepsilon} y)}$$

$$\leq e^{2k_\varepsilon (\sup g - \inf g)} e^{2K}.$$

The claim holds for  $n \ge 2k_{\varepsilon}$  by taking the sup over x and the inf over y in P. Since there are only finitely many values of n to correct for, by taking a larger constant, the claimed estimate holds for all  $n \ge 1$ .

Fix some  $l \geqslant 0$ . Since an element  $P \in \mathcal{P}_{-l}^n$  is contained in a unique element  $\tilde{P} \in \mathcal{P}_0^n$ , we get that

$$\sup_{P} e^{S_n g} \leqslant \sup_{\tilde{P}} e^{S_n g} \leqslant C \inf_{\tilde{P}} e^{S_n g} \leqslant C \inf_{P} e^{S_n g}.$$

Now, assume that  $\mathring{P} \neq \emptyset$ . Then, by the continuity of  $S_ng$  on P, the estimate also holds when the sup and the inf are taken over  $\mathring{P}$ . In other words, the claim is true for all  $P \in \mathring{\mathcal{P}}_{l}^{n}$ .

Since by [3, Lemma 3.3],  $\mathring{\mathcal{P}}_{-l}^n = \mathcal{M}_{-l-1}^{n+1}$ , the claim is true for all  $P \in \mathcal{M}_{-l}^n$ , for fixed  $l \ge 1$ . We finish the proof with the case  $P \in \mathcal{M}_0^n$ . Remark that letting  $A \in \mathcal{M}_{-1}^n$ , then  $T^{-1}A \in \mathcal{M}_0^{n+1}$ . Therefore

$$e^{-\sup g} \sup_{T^{-1}A} e^{S_{n+1}g} \leqslant \sup_{T^{-1}A} e^{S_{n+1}g-g} = \sup_{A} e^{S_{n}g} \leqslant C \inf_{A} e^{S_{n}g} = C \inf_{T^{-1}A} e^{S_{n+1}g-g}$$
$$\leqslant 2Ce^{-\inf g} \inf_{T^{-1}A} e^{S_{n+1}g}.$$

Only in this last case, we need to replace C by  $2Ce^{\sup g - \inf g} \geqslant C$ .

Proof of Theorem 2.1. Let  $p_n = \sum_{A \in \mathcal{P}_0^n} \sup_{x \in A} e^{(S_n g)(x)}$ . Then, for  $k \geqslant n$ ,

$$p_{n+k} = \sum_{B \cap C \in \mathcal{P}_0^n \bigvee T^{-n} \mathcal{P}_0^k} \sup_{x \in B \cap C} e^{(S_n g)(x) + (S_k g)(T^n x)} \leqslant p_n p_k.$$

Therefore  $(\log p_n)_n$  is a sub-additive sequence. It is then classical that  $\frac{1}{n}\log p_n$  converges to  $\inf_{n\geqslant 1}\frac{1}{n}\log p_n$ , hence  $P_*(T,g)$  exists. We now prove the statement about convexity. Let  $g_1$  and  $g_2$  be two potentials bounded from above and  $p\in[0,1]$ . Using the Hölder inequality, we get that for all  $n\geqslant 1$ 

$$\sum_{A \in \mathcal{P}_0^n} \sup_A e^{p S_n g_1 + (1 - p) S_n g_2} \leqslant \sum_{A \in \mathcal{P}_0^n} \sup_A \left( e^{S_n g_1} \right)^p \left( \sup_A e^{S_n g_2} \right)^{1 - p}$$

$$\leqslant \left( \sum_{A \in \mathcal{P}_0^n} \sup_A e^{S_n g_1} \right)^p \left( \sum_{A \in \mathcal{P}_0^n} \sup_A e^{S_n g_2} \right)^{1 - p}.$$

Taking the appropriate limits, we get that  $P_*(T, pg_1 + (1-p)g_2) \leq pP_*(T, g_1) + (1-p)P_*(T, g_2)$ , hence the claimed convexity.

For (i), consider  $\tilde{p}_n = \sum_{A \in \mathring{\mathcal{P}}_0^n} \sup_{x \in A} e^{(S_n g)(x)}$ . Notice that

$$\mathcal{P}_0^n = \{ A \in \mathcal{P}_0^n \mid \mathring{A} \neq \emptyset \} \sqcup \{ A \in \mathcal{P}_0^n \mid \mathring{A} = \emptyset \}.$$

Now, Baladi and Demers proved in [3, Lemma 3.2] that the cardinality of the second term in the right hand side grows at most linearly. Hence

$$\sum_{\substack{A \in \mathcal{P}_0^n \\ \mathring{a} - \emptyset}} \sup_{x \in A} e^{(S_n g)(x)} \leqslant C n e^{n \sup g}.$$

By the smoothness of g,

$$\sum_{\substack{A \in \mathcal{P}_0^n \\ \mathring{A} \neq \emptyset}} \sup_{x \in A} e^{(S_n g)(x)} = \sum_{\substack{A \in \mathring{\mathcal{P}}_0^n \\ \mathring{x} \neq \emptyset}} \sup_{x \in A} e^{(S_n g)(x)}.$$

Thanks to the assumption  $P_*(T,g) - \sup g > 0$ , the sum over elements  $A \in \mathcal{P}_0^n$  with  $\mathring{A} \neq \emptyset$  dominates the sum over A with  $\mathring{A} = \emptyset$ . Thus  $(\frac{1}{n} \log \tilde{p}_n)_n$  converges to the same limit as  $(\frac{1}{n} \log p_n)_n$  does.

For (ii), we use [3, Lemma 3.3] that  $\mathring{\mathcal{P}}_0^n = \mathcal{M}_{-1}^{n+1}$ . Hence

$$\sum_{A \in \mathcal{M}_0^{n+1}} \sup_{x \in A} e^{(S_{n+1}g)(x)} \leqslant \sum_{A \in \mathcal{M}_{-1}^{n+1}} \sup_{x \in A} e^{(S_{n+1}g)(x)} \leqslant \sum_{A \in \mathring{\mathcal{P}}_0^n} \sup_{x \in A} e^{(S_ng)(x)} \sup_{x \in M} e^{g(x)}.$$

Furthermore, since  $\mathcal{M}_{-1}^{n+1} = \mathcal{M}_0^{n+1} \bigvee \mathcal{M}_{-1}^0$ , each element of  $\mathcal{M}_0^{n+1}$  contains at most  $\#\mathcal{M}_{-1}^0$  elements of  $\mathcal{M}_{-1}^{n+1}$ . Hence

$$\sum_{A \in \mathcal{M}_0^{n+1}} \sup_{x \in A} e^{(S_{n+1}g)(x)} \geqslant \frac{1}{\# \mathcal{M}_{-1}^0} \sum_{A \in \mathring{\mathcal{P}}_0^n} \sup_{x \in A} e^{(S_ng)(x)} \inf_{x \in M} e^{g(x)}.$$

Point (iii) (resp. (iv), (v)) follows directly from the definition of  $P_*(T, g)$  (resp. from point (i), (ii)) and from Lemma 2.3 since

$$\inf_{A} e^{S_n g} \leqslant \sup_{A} e^{S_n g} \leqslant C \inf_{A} e^{S_n g},$$

for all A in  $\mathcal{P}_0^n$  (resp.  $\mathring{\mathcal{P}}_0^n$ ,  $\mathcal{M}_0^n$ ). For the final claim, we prove that  $\log \sum_{P \in \mathring{\mathcal{P}}_1^n} \sup_{P} e^{S_n g}$  is

subadditive. Take P a nonempty element of  $\mathring{\mathcal{P}}_1^{n+m}$ . It is the interior of an intersection of elements of the form  $T^{-j}A_j$  for some  $A_j \in \mathcal{P}$ , for  $j = 1, \ldots, n+m$ . This is equal to the

intersection of the interiors of  $T^{-j}A_j$ . But since P is nonempty, none of the  $T^{-j}A_j$  has empty interior, and so none of the  $A_j$  has empty interior. Thus the interiors of  $A_j$  are in  $\mathring{\mathcal{P}}$ . Now, splitting the intersection of the first n sets from the last m, we see that the intersection of the first n sets forms an element of  $\mathring{\mathcal{P}}_1^n$ . For the last m sets, we can factor out  $T^{-n}$  at the price of making the set slightly bigger:

$$int(T^{-n-j}A_{-n-j}) \subset T^{-n}(int(T^{-j}(A_{-n-j}))), \quad 1 \le j \le m$$

where int denotes the interior of a set. Thus

$$\begin{split} \sum_{P \in \mathring{\mathcal{P}}_{1}^{n+m}} \sup_{P} e^{S_{n+m}g} &\leqslant \sum_{\substack{A_{-j} \in \mathring{\mathcal{P}} \\ 1 \leqslant j \leqslant n+m}} \sup\{e^{S_{n}g + S_{m}g \circ T^{n}}(x) \mid x \in \bigcap_{j=1}^{n} T^{-j}A_{-j} \cap T^{-n} \bigcap_{j=1}^{m} T^{-j}A_{-n-j}\} \\ &\leqslant \sum_{\substack{A_{-j} \in \mathring{\mathcal{P}} \\ 1 \leqslant j \leqslant n}} \sup\{e^{S_{n}g}(x) \mid x \in \bigcap_{j=1}^{n} T^{-j}A_{-j}\} \sum_{\substack{A_{-j} \in \mathring{\mathcal{P}} \\ 1 \leqslant j \leqslant m}} \sup\{e^{S_{m}g}(x) \mid x \in \bigcap_{j=1}^{m} T^{-j}A_{-j}\} \\ &\leqslant \sum_{P \in \mathring{\mathcal{P}}_{1}^{n}} \sup_{P} e^{S_{n}g} \sum_{P \in \mathring{\mathcal{P}}_{1}^{m}} \sup_{P} e^{S_{m}g} \\ \end{split}$$

Taking logs, the sequence is subadditive. And then so is the sequence with  $\mathcal{M}_0^n$  in place of  $\mathring{\mathcal{P}}_1^{n-1}$ .

Proof of Proposition 2.2. We first prove the claim about separating sets. Let  $\varepsilon > 0$  and let  $k_{\varepsilon}$  be large enough so that  $\operatorname{diam}^{s}(\mathcal{M}_{-k_{\varepsilon}-1}^{0}) \leq C\Lambda^{-k_{\varepsilon}} < c_{1}\varepsilon$  for some constant  $c_{1}$  to be defined later. Therefore  $\operatorname{diam}^{u}(\mathcal{M}_{-k_{\varepsilon}-1}^{n+1}) \leq C\Lambda^{-k_{\varepsilon}} < c_{1}\varepsilon$  for all  $n \geq k_{\varepsilon}$ . By the uniform transversality between the stable and the unstable cones, we can choose  $c_{1}$  such that  $\operatorname{diam}(\mathcal{M}_{-k_{\varepsilon}-1}^{n+1}) < \varepsilon$  for all  $n \geq k_{\varepsilon}$ .

Let E be  $(n, \varepsilon)$ -separated, for some  $n \ge k_{\varepsilon}$ . It is shown in the proof of [3, Lemma 3.4] that if  $x, y \in E$  are distinct, then they cannot be contained in the same element of  $\mathcal{P}_{-k_{\varepsilon}}^{k_{\varepsilon}+n}$ . Thus

$$\sum_{x \in E} e^{S_n g(x)} \leqslant \sum_{A \in \mathcal{P}_{-k_{\varepsilon}}^{k_{\varepsilon}+n}} |e^{S_n g}|_{C^0(A)} = \sum_{A \in \mathcal{P}_0^{2k_{\varepsilon}+n}} |e^{S_n g \circ T^{-k_{\varepsilon}}}|_{C^0(A)}$$
$$\leqslant e^{k_{\varepsilon} (\sup g - 2\inf g)} \sum_{A \in \mathcal{P}_0^{k_{\varepsilon}+n}} |e^{S_{2k_{\varepsilon}+n}g}|_{C^0(A)}.$$

Therefore

$$\lim_{n\to\infty}\frac{1}{n}\log\sup\{\sum_{x\in E}e^{S_ng(x)}\mid\ E\ \text{is}\ (n,\varepsilon)\text{-separated}\}\leqslant P_*(T,g)\quad , \ \text{for all}\ \varepsilon>0.$$

Taking the limit  $\varepsilon \to 0$ , we get  $P_{\text{sep}}(T,g) \leqslant P_*(T,g)$ .

For the reverse inequality, assume that g is such that  $P_*(T,g) - \sup g > 0$ . From the proof of [3, Lemma 3.4], there exists  $\varepsilon_0 > 0$  such that for all  $\varepsilon < \varepsilon_0$ , any set E which contains only one point per element of  $\mathcal{M}_0^n$  is  $(n,\varepsilon)$ -separated. For all  $A \in \mathcal{M}_0^n$ , there exists  $x \in A$  such that

 $e^{S_n g(x)} \geqslant \frac{9}{10} \sup_A e^{S_n g}$ . Let E be the collection of such x. Thus

$$\sum_{x \in E} e^{S_n g(x)} \geqslant \frac{9}{10} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)}.$$

Therefore,

$$\lim_{n\to\infty}\frac{1}{n}\log\sup\{\sum_{x\in E}e^{S_ng(x)}\mid\ E\ \text{is}\ (n,\varepsilon)\text{-separated}\}\geqslant P_*(T,g)\,,\quad\text{for all}\ 0<\varepsilon<\varepsilon_0.$$

Taking the limit  $\varepsilon \to 0$ , we get  $P_{\text{sep}}(T,g) \geqslant P_*(T,g)$ , thus the claimed equality.

We now prove the claim concerning spanning sets. Let  $\varepsilon > 0$  and let  $k_{\varepsilon}$  be such that  $\operatorname{diam}(\mathcal{M}_{-k_{\varepsilon}-1}^{n+1}) < \varepsilon$  for all  $n \geq k_{\varepsilon}$ . Let F be a set containing one point in each element of  $\mathcal{P}_{-k_{\varepsilon}}^{n+1}$ . From the proof of [3, Lemma 3.5], F is  $(n, \varepsilon)$ -spanning. Since

$$\sum_{x \in F} e^{S_n g(x)} \leqslant e^{k_{\varepsilon}(\sup g - 2\inf g)} \sum_{A \in \mathcal{P}_0^{2k_{\varepsilon} + n}} |e^{S_{2k_{\varepsilon} + n}g}|_{C^0(A)}$$

we get that

$$\lim_{n\to\infty}\frac{1}{n}\log\inf\{\sum_{x\in F}e^{S_ng(x)}\mid F\text{ is }(n,\varepsilon)\text{-spanning}\}\leqslant P_*(T,g)\,,\quad\text{for all }\varepsilon>0\,.$$

Taking the limit  $\varepsilon \to 0$ , we get  $P_{\text{span}}(T,g) \leqslant P_*(T,g)$ .

For the reverse inequality, assume that g is a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential such that  $P_*(T, g) - \sup g > 0$ . Let  $\varepsilon < \varepsilon_0$  and let F be a  $(n, \varepsilon)$ -spanning set. By the proof of [3, Lemma 3.5], each element of  $\mathcal{M}_0^n$  contains at least one element of F. Thus

$$\sum_{x \in F} e^{S_n g(x)} \geqslant \sum_{A \in \mathcal{M}_n^n} \inf_A e^{S_n g}.$$

Taking the appropriate limits, we get that  $P_{\text{span}}(T,g) \ge P_*(T,g)$ , thus the claimed equality.  $\square$ 

2.1. Easy Direction of the Variational Principle for the Pressure. Recall that given a T-invariant probability measure  $\mu$  and a finite measurable partition  $\mathcal{A}$  of M, the entropy of  $\mathcal{A}$  with respect to  $\mu$  is defined by  $H_{\mu}(\mathcal{A}) = -\sum_{A \in \mathcal{A}} \mu(A) \log \mu(A)$ , and the entropy of T with respect to  $\mathcal{A}$  is  $h_{\mu}(T, \mathcal{A}) = \lim_{n \to \infty} \frac{1}{n} H_{\mu} \left( \bigvee_{i=0}^{n-1} T^{-i} \mathcal{A} \right)$ .

**Lemma 2.4.** Let  $\varphi: M \to \mathbb{R}$  be a measurable function. Then

$$P_*(T,\varphi) \geqslant P(T,\varphi) := \sup\{h_\mu(T) + \int \varphi \,\mathrm{d}\mu \mid \mu \text{ is a $T$-invariant Borel probability measure}\}$$

*Proof.* Let  $\mu$  be a T-invariant probability measure on M. Notice that  $\mathcal{P}$  is a generator for T since  $\bigvee_{i=-\infty}^{\infty} T^{-i}\mathcal{P}$  separates points in M. Thus  $h_{\mu}(T) = h_{\mu}(T,\mathcal{P})$  (see for example [32,

Theorem 4.17]). Then,

$$h_{\mu}(T) + \int \varphi \, d\mu = \lim_{n \to \infty} \frac{1}{n} \sum_{A \in \mathcal{P}_0^n} \left( -\mu(A) \log \mu(A) + \int_A S_n \varphi \, d\mu \right)$$

$$\leqslant \lim_{n \to \infty} \frac{1}{n} \sum_{A \in \mathcal{P}_0^n} \mu(A) (\sup_A (S_n \varphi) - \log \mu(A))$$

$$\leqslant \lim_{n \to \infty} \frac{1}{n} \log \sum_{A \in \mathcal{P}_0^n} \sup_A e^{S_n \varphi} \leqslant P_*(T, \varphi)$$

where we used [32, Lemma 9.9] for the second inequality.

2.2. Abramov Formula and Choice of the Potential g. In order to obtain the existence of MME for the billiard flow, we make use of the Abramov formula to relate equilibrium measure for T and some potential g, to the MME of the flow. First, we need the following lemma.

**Lemma 2.5.** Let  $\varphi$  be a bounded non-negative measurable function such that  $\varphi_0 := \inf\{\int \varphi \, d\mu \mid T_*\mu = \mu\} > 0$ . Then, there exists a unique real number  $c_{\varphi}$  such that  $P(T, -c_{\varphi}\varphi) = 0$ .

*Proof.* We first prove that the function  $t \mapsto P(T, t\varphi)$  is increasing. Let  $\varepsilon > 0$  and  $t_1 < t_2$ . There exists a T-invariant probability measure  $\mu_1$  such that

$$P(T, t_1 \varphi) \leqslant h_{\mu_1}(T) + t_1 \int \varphi \, \mathrm{d}\mu_1 + \varepsilon \leqslant P(T, t_2 \varphi) - (t_2 - t_1)\varphi_0 + \varepsilon.$$

By this computation, we also get that  $\lim_{t\to\pm\infty}P(T,t\varphi)=\pm\infty$  .

Now we prove that  $t \mapsto P(T, t\varphi)$  is continuous. Let  $\varepsilon > 0$  and  $t \in \mathbb{R}$ . By the previous computation, we get that  $\varepsilon \varphi_0 \leqslant P(T, (t+\varepsilon)\varphi) - P(T, t\varphi)$ . Let  $\mu_2$  be such that  $P(T, (t+\varepsilon)\varphi) \leqslant h_{\mu_2}(T) + (t+\varepsilon) \int \varphi \, d\mu_2 + \varepsilon$ . Thus

$$P(T, (t+\varepsilon)\varphi) - P(T, t\varphi) \le \varepsilon(1 + \sup \varphi).$$

Therefore  $t \mapsto P(T, t\varphi)$  is strictly increasing and continuous, so it must vanish at exactly one point, noted  $-c_{\varphi}$ .

We can now use this lemma with the Abramov formula to get the following

**Corollary 2.6.** Equilibrium measures of T under the potential  $-h_{top}(\phi_1)\tau$  and MME of the billiard flow (seen as a suspension flow) are in one-to-one correspondence through the bijection  $\mu \mapsto \mu_{\tau} := \frac{1}{\mu(\tau)}\mu \otimes \lambda$ , where  $\lambda$  is the Lebesgue measure.

*Proof.* Since  $\tau \ge \tau_{\min} > 0$ , the assumption of Lemma 2.5 is satisfied for  $\varphi = \tau$ . Let c be the constant given by Lemma 2.5 such that  $0 = P(T, -c\tau)$ . Then, for every equilibrium state m of T under the potential  $-c\tau$ , we get

$$0 = h_m(T) - c \int \tau \, \mathrm{d}m \geqslant h_\mu(T) - c \int \tau \, \mathrm{d}\mu,$$

for all T-invariant measure  $\mu$ . Thus

$$c = \frac{h_m(T)}{\int \tau \, \mathrm{d}m} \geqslant \frac{h_\mu(T)}{\int \tau \, \mathrm{d}\mu}.$$

Now, by the Abramov formula,  $c = h_{m_{\tau}}(\phi_1) \ge h_{\mu_{\tau}}(\phi_1)$ . In other words,  $m_{\tau}$  is a MME for the billiard flow. Furthermore, since  $\phi_1$  is a continuous map of a compact metric space, by [32, Theorem 8.6], we get that  $h_{\text{top}}(\phi_1) = \sup\{h_{\mu}(\phi_1) \mid (\phi_1)_* \mu = \mu\}$ . Thus  $c = h_{\text{top}}(\phi_1)$ .

To prove that the function is onto, we use that any  $\phi_t$ -invariant probability measure  $\mu_{\tau}$  must be of the form  $\frac{1}{\mu(\tau)}\mu \otimes \lambda$ , for some T-invariant probability measure  $\mu$ . Thus, reversing the above computations, we get that if  $\mu_{\tau}$  is a MME, then  $\mu$  is an equilibrium state for T under the potential  $-h_{\text{top}}(\phi_1)\tau$ .

Therefore, proving the existence and uniqueness of a MME for the billiard flow is equivalent to proving the existence and uniqueness of equilibrium states of T under the potential  $g = -h_{top}(\phi_1)\tau$ . Notice that in the second case, g is  $(\mathcal{M}_0^1, \frac{1}{2})$ -Hölder continuous and the condition  $P_*(T,g) - \sup g > 0$  from Theorem 2.1 is realised since  $P_*(T,g) - \sup g > P(T, -h_{top}(\phi_1)\tau) + h_{top}(\phi_1)\tau_{\min} > 0$ .

Remark 2.7. Using similar arguments as in Corollary 2.6, we can relate the equilibrium states of  $\phi_t$  under  $\tilde{g}: \Omega \to \mathbb{R}$  to the ones of T under  $g = \lambda(\tilde{g}) - P(\phi_1, \tilde{g})\tau$ , where  $\lambda(\tilde{g}): M \to \mathbb{R}$  is given by

$$\lambda(\tilde{g})(x) = \int_0^{\tau(x)} \tilde{g}(\phi_t(x)) dt.$$

## 3. Growth Lemma and Fragmentation Lemmas

This section contains growth lemmas, controlling the growth in complexity of the iterates of a stable curve, with a weight g. We also formulate the precise definitions of the conditions SSP.1 and SSP.2. The first condition will be used to prove the "Lasota–Yorke" bounds on the transfer operator  $\mathcal{L}_g$  in Proposition 5.1, as well as the lower bound on the spectral radius in Theorem 5.3, while SSP.2 will be crucial for the absolute continuity (Corollary 6.8) used to prove statistical properties (Propositions 6.12 and 6.15) and to compute the pressure (Corollary 6.14).

In view of deriving Lemma 3.4 from Lemma 3.3, we must work with a class of curves more general than stable manifolds. Recall the stable and unstable cones (2.1).

We define a set of cone-stable curves<sup>2</sup>  $\widehat{W}^s$  whose tangent vectors all lie in the stable cone  $C^s$  for the map, with length at most  $\delta_0$  (to be determined latter) and curvature bounded above so that  $T^{-1}\widehat{W}^s \subset \widehat{W}^s$ , up to subdivision of curves. We define a set of cone-unstable curves  $\widehat{W}^u$  similarly. These sets of curves will be relevant since  $S_n$  and  $S_{-k}$  are composed of curves in  $\widehat{W}^s$  and  $\widehat{W}^u$ , respectively.

For  $\delta \in (0, \delta_0]$  and  $W \in \widehat{\mathcal{W}}^s$ , let  $\mathcal{G}_0^{\delta}(W) := \{W\}$ . For  $n \geqslant 1$ , define the  $\delta$ -scaled subdivision  $\mathcal{G}_n^{\delta}(W)$  inductively as the collection of smooth components of  $T^{-1}(W')$  for  $W' \in \mathcal{G}_{n-1}^{\delta}(W)$ , where elements longer than  $\delta$  are subdivided to have length between  $\delta/2$  and  $\delta$ . Thus  $\mathcal{G}_n^{\delta}(W) \subset \widehat{\mathcal{W}}^s$  for each n and  $\bigcup_{U \in \mathcal{G}_n^{\delta}(W)} U = T^{-n}W$ . Moreover, if  $W \in \mathcal{W}^s$ , then  $\mathcal{G}_n^{\delta}(W) \subset \mathcal{W}^s$ .

Denote by  $L_n^{\delta}(W)$  those elements of  $\mathcal{G}_n^{\delta}(W)$  having length at least  $\delta/3$  (the long curves),  $S_n^{\delta}(W) \coloneqq \mathcal{G}_n^{\delta}(W) \setminus L_n^{\delta}(W)$  (the short curves), and define  $\mathcal{I}_n^{\delta}(W)$  to comprise those elements  $U \in \mathcal{G}_n^{\delta}(W)$  for which  $T^iU$  is contained in an element of  $S_{n-i}^{\delta}(W)$  for all  $0 \le i \le n-1$ .

<sup>&</sup>lt;sup>2</sup>The notation  $\mathcal{W}^s$  will be used for the stable manifolds. See Subsection 4.2.

A fundamental fact [14, Lemma 5.2] we will use is that the growth in complexity for the billiard is at most linear:

 $\exists K > 0$  such that  $\forall n \geq 0$ , the number of curves in  $S_{\pm n}$  that intersect (3.1)at a single point is at most Kn.

### 3.1. Growth Lemma.

**Lemma 3.1.** For any  $m \in \mathbb{N}$ , there exists  $\delta_0 = \delta_0(m) \in (0,1)$  such that for all  $W \in \widehat{\mathcal{W}}^s$ , if  $|W| < \delta_0$ , then for all  $0 \le l \le 2m$ ,  $T^{-l}W$  comprises at most Km + 1 connected components. Furthermore, for any  $\delta \in (0, \delta_0]$ , the  $\delta$ -scaled subdivisions satisfy the following estimates: for all  $n \ge 1$ , all  $\gamma \in [0, \infty)$ , all  $W \in \widehat{\mathcal{W}}^s$ , and all  $\mathcal{M}_0^1$ -continuous potential g, we have

a) 
$$\sum_{W_i \in \mathcal{I}_n^{\delta}(W)} \left( \frac{\log |W|}{\log |W_i|} \right)^{\gamma} |e^{S_n g}|_{C^0(W_i)} \leq 2^{((n \vee n_0)s_0 + 1)\gamma + 1} (Km + 1)^{n/m} e^{n \sup g}$$

b) 
$$\sum_{W_i \in \mathcal{G}_n^{\delta}(W)} \left( \frac{\log |W|}{\log |W_i|} \right)^{\gamma} |e^{S_n g}|_{C^0(W_i)} \leqslant \min \left\{ 2C\delta^{-1} 2^{((n \vee n_0)s_0 + 1)\gamma} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} \right\},$$

$$2^{2\gamma+1}C\delta^{-1}\sum_{j=1}^{n}2^{(j\vee n_0)s_0\gamma}(Km+1)^{j/m}e^{j\sup g}\sum_{A\in\mathcal{M}_0^{n-j}}|e^{S_{n-j}g}|_{C^0(A)}$$

where  $(n \vee n_0) = \max(n, n_0)$ .

Moreover, if  $|W| \ge \delta/2$ , then both factors  $2^{(ns_0+1)\gamma}$  can be replaced by  $2^{\gamma}$ .

*Proof.* By [13, Exercise 4.50], there exist constants  $\delta_{\rm CM} > 0$  and  $C \ge 1$  such that for all  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta_{\text{CM}}$ , then  $|T^{-1}W| \leqslant C|W|^{1/2}$ . Notice also that there exists  $\Lambda_1 := \Lambda_1(\varphi_0)$ such that for  $W \in \widehat{\mathcal{W}}^s$  with  $T^{-1}W \cap \{|\varphi| > \varphi_0\} = \emptyset$ , then  $|T^{-1}W| \leqslant \Lambda_1|W|$ . We want to iterate these bounds.

Let  $\delta \in (0, \delta_{\text{CM}}]$ ,  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta$ , and  $W_i \in \mathcal{I}_n^{\delta}(W)$ . Let  $V \subset W$  corresponding to  $W_i$ , that is  $V = T^n W_i$ . Thus, for all  $1 \leq j \leq n$ , we have  $|T^{-j}V| = |T^{n-j}W_i| \leq \delta/3$ . We can decompose  $V = \bigsqcup_{i_0 \in I_0} V_{i_0, \text{graz}}^0 \cup \bigsqcup_{j_0 \in J_0} V_{j_0, \text{exp}}^0$  such that: for all  $i_0 \in I_0$ ,  $T^{-1}V_{i_0, \text{graz}}^0 \subset \{|\varphi| \geqslant \varphi_0\}$ , and thus  $|T^{-1}V_{i_0, \text{graz}}^0| \leq C|V_{i_0, \text{graz}}^0|^{1/2}$ , and for all  $j_0 \in J_0$ ,  $T^{-1}V_{j_0, \text{exp}}^0 \subset \{|\varphi| < \varphi_0\}$ , and thus  $|T^{-1}V_{j_0, \text{exp}}^0| \leq \Lambda_1|V_{j_0, \text{exp}}^0|$ . We can perform the same decomposition for  $V_{i_0, \text{graz}}^0$  or  $V_{i_0, \text{graz}}^0$  instead of  $V_i$ .  $V_{j_0,\text{exp}}^0$  instead of V:

$$V_{i_0,\text{graz}}^0 = \bigsqcup_{i_1} V_{i_1,\text{graz}}^{1,i_0} \cup \bigsqcup_{j_1} V_{j_1,\text{exp}}^{1,i_0} , \qquad V_{j_0,\text{exp}}^0 = \bigsqcup_{i_1} V_{i_1,\text{graz}}^{1,j_0} \cup \bigsqcup_{j_1} V_{j_1,\text{exp}}^{1,j_0}.$$

We can iterate this decomposition until having a decomposition of  $T^{-n}V = W_i$ . Notice that since the stable curves  $T^{-j}V$  have length at most  $\delta/3 \leq \delta_{\rm CM}/3$  and are uniformly transverse to  $S_0$ , they can cross  $\{|\varphi| \geqslant \varphi_0\}$  at most B times, where B > 0 is a constant uniform in W.

Thus the number of pieces at each decomposition is bounded by 2B. Thus  $W_i = T^{-n}V = \bigcup_{\alpha_n, *} V_{\alpha_n, *}^{n,\alpha_0,\dots,\alpha_{n-1}}$ , where the union is made of at most  $(2B)^n$ Thus  $W_i = T^{-n}V = \bigsqcup_{* \in \{\text{graz}, \exp\}}$ 

elements we can estimate the length.

Consider first the case  $n \leqslant n_0$ . By definition,  $s_0$  is such that  $s_0 = \sup_{M} \frac{1}{n_0} \sum_{k=0}^{n_0-1} \mathbb{1}_{\{|\varphi| > \varphi_0\}} \circ T^k < 1$ . Thus, for each  $V_{\alpha_n,*}^{n,\alpha_0,\dots,\alpha_{n-1}}$  there are at most  $s_0 n_0$  indices  $\alpha_k \in I_k$ . Thus  $|V_{\alpha_n,*}^{n,\alpha_0,\dots,\alpha_{n-1}}| \leqslant C^2 \Lambda_1^{n_0} |V|^{2^{-s_0 n_0}}$ . Therefore

$$|W_i| = |T^{-n_0}V| \leqslant (2B)^{n_0} C^2 \Lambda_1^{n_0} |V|^{2^{-s_0n_0}} \leqslant \tilde{C} |W|^{2^{-s_0n_0}}, \quad \forall W_i \in \mathcal{I}_n^{\delta}(W), \ n \leqslant n_0, \ \delta \leqslant \delta_{\text{CM}}$$

Now, consider the case  $n=kn_0+l$ , for  $k\geqslant 1$  and  $0\leqslant l< n_0$ . By construction, if  $W_i\in\mathcal{I}_n^\delta(W)$ , then  $T^lW_i\subset W_i^0\in\mathcal{I}_{kn_0}^\delta(W)$  and  $T^{n_0}W_i^j\subset W_i^{j+1}\in\mathcal{I}_{(k-j-1)n_0}^\delta(W)$  for all  $0\leqslant j\leqslant k-1$ . Thus, we can iterate (3.2):

$$|W_i| \leqslant \tilde{C} |W_i^0|^{2^{-s_0 n_0}} \leqslant \tilde{C}^{\sum_{m=0}^j 2^{-m s_0 n_0}} |W_i^j|^{2^{-j s_0 n_0}} \leqslant \tilde{C}^2 |W|^{2^{-(k+1) s_0 n_0}}$$

and so  $|W_i| \leq \tilde{C}^2 |W|^{2^{-ns_0}}$  for all  $W_i \in \mathcal{I}_n^{\delta}(W)$ ,  $n \geq n_0$  and all  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta_{\text{CM}}$ . Therefore, if  $\delta \leq \min(\tilde{C}^{-2}, \delta_{\text{CM}})$ , we have

$$\left(\frac{\log|W|}{\log|W_i|}\right)^{\bar{\gamma}} \leqslant \left(2^{s_0 n} \left(1 - \frac{\log \tilde{C}^2}{\log|W_i|}\right)\right)^{\bar{\gamma}} \leqslant 2^{(ns_0 + 1)\bar{\gamma}}, \, \forall \, W_i \in \mathcal{I}_n^{\delta}(W),$$

since  $|W_i| \leq \delta$ .

(a) Let  $m \ge 1$  and  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta \le \min(\widetilde{C}^{-2}, \delta_{\mathrm{CM}})$ . First, we want to estimate the number of smooth components of  $T^{-l}W$ , for  $0 \le l \le 2m$ . The problem is the same as knowing the number of connected components of  $W \setminus \mathcal{S}_{-l}$ . Now, by (3.1), at most Kl curves in  $\mathcal{S}_{-l}$  can intersect at a given point. Since W and  $\mathcal{S}_{-l}$  are uniformly transverse, for each  $0 \le l \le 2m$  there exists  $\delta_{(l)}$  such that if  $|W| < \delta_{(l)}$  then  $W \setminus \mathcal{S}_{-l}$  has at most Kl+1 connected components. Let  $\delta_0 := \min\{\delta_{(l)} \mid 0 \le l \le 2m\}$ .

Let  $n \ge 1$ ,  $\delta \in (0, \delta_0]$  and  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta$ . We want to estimate  $\#\mathcal{I}_n^{\delta}(W)$ . We prove by induction that  $\#\mathcal{I}_{jm}^{\delta}(W) \le (Km+1)^j$ . For j=1, this follows from the choice of  $\delta_0$ . Since elements of  $\mathcal{I}_{(j+1)m}^{\delta}(W)$  are of the form  $V \in \mathcal{I}_m^{\delta}(W_i)$  for  $W_i \in \mathcal{I}_{jm}^{\delta}(W)$ , we have

$$\#\mathcal{I}_{(j+1)m}^{\delta}(W)\leqslant (Km+1)\#\mathcal{I}_{jm}^{\delta}(W)\leqslant (Km+1)^{j+1}.$$

Now for estimating  $\#\mathcal{I}_{im+l}^{\delta}(W)$ ,  $0 \leq l < m$ , we only need to modify the last step:

$$\#\mathcal{I}_{jm+l}^{\delta}(W) \leqslant (K(m+l)+1)\#\mathcal{I}_{(j-1)m}^{\delta}(W) \leqslant 2(Km+1)^{j}.$$

Therefore,  $\#\mathcal{I}_n^{\delta}(W) \leqslant 2(Km+1)^{n/m}$ , for all  $n \geqslant 1$ .

Finally, we have that for  $n \ge n_0$ 

$$\sum_{W_i \in \mathcal{I}_n^{\delta}(W)} \left( \frac{\log |W|}{\log |W_i|} \right)^{\gamma} |e^{S_n g}|_{C^0(W_i)} \leq e^{n \sup g} 2^{(ns_0 + 1)\bar{\gamma}} \# \mathcal{I}_n^{\delta}(W)$$

$$\leq 2^{(ns_0 + 1)\gamma + 1} (Km + 1)^{n/m} e^{n \sup g}$$

(b) Let 
$$\delta \leqslant \delta_0$$
, and  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta$ . We start by giving an estimate on  $\sum_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}$ .

Since the border of elements of  $\mathcal{M}_{-n}^0$  is contained in  $\mathcal{S}_{-l}$ , by uniform transversality, each element of  $\mathcal{M}_{-n}^0$  is crossed at most one time by W. Thus, each element of  $\mathcal{M}_0^n$  is crossed at most one time by  $T^{-n}W$ . Now, since the diameter of elements of  $\mathcal{M}_0^n$  is bounded uniformly in n by some constant C, there can be no more than  $2C\delta^{-1}$  elements of  $\mathcal{G}_n^{\delta}(W)$  in a single element of  $\mathcal{M}_0^n$ . Thus

(3.3) 
$$\sum_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)} \leq 2C\delta^{-1} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)}$$

First, in the case  $|W| \ge \delta/2$ , the estimate

$$\sum_{W_i \in \mathcal{G}_n^{\delta}(W)} \left( \frac{\log |W|}{\log |W_i|} \right)^{\bar{\gamma}} |e^{S_n g}|_{C^0(W_i)} \leqslant 2C \delta^{-1} 2^{((n \vee n_0) s_0 + 1)\bar{\gamma}} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)},$$

is enough for what we need.

Now, assume that  $|W| < \delta/2$ . Let  $F_1(W)$  denote those  $V \in \mathcal{G}_1^{\delta}(W)$  whose length is at least  $\delta/2$ . Inductively, define  $F_j(W)$ , for  $2 \leq j \leq n-1$ , to contain those  $V \in \mathcal{G}_j^{\delta}(W)$  whose length is at least  $\delta/2$ , and such that  $T^kV$  is contained in an element of  $\mathcal{G}_{j-k}^{\delta}(W) \setminus F_{j-k}(W)$  for any  $1 \leq k \leq j-1$ . Thus  $F_j(W)$  contains elements of  $\mathcal{G}_j^{\delta}(W)$  that are "long for the first time" at time j.

We group  $W_i \in \mathcal{G}_n^{\delta}(W)$  by its "first long ancestor" as follows. We say  $W_i$  has first long  $ancestor^3 \ V \in F_j(W)$  for  $1 \leq j \leq n-1$  if  $T^{n-j}W_i \subseteq V$ . Note that such a j and V are unique for each  $W_i$  if they exist. If no such j and V exist, then  $W_i$  has been forever short and so must belong to  $\mathcal{I}_n^{\delta}(W)$ . Denote by  $A_{n-j}(V)$  the set of  $W_i \in \mathcal{G}_n^{\delta}(W)$  corresponding to one  $V \in F_j(W)$ , that is

$$A_{n-j}(V) := \{ W_i \in \mathcal{G}_n^{\delta}(W) \mid T^{n-j}W_i \subset V \}.$$

By construction, we have the relation

$$\mathcal{G}_n^{\delta}(W) \setminus \left( \bigsqcup_{j=1}^{n-1} \bigsqcup_{V \in F_j(W)} A_{n-j}(V) \right) = \mathcal{I}_n^{\delta}(W).$$

Therefore

$$\begin{split} & \sum_{W_{i} \in \mathcal{G}_{n}^{\delta}(W)} \left( \frac{\log |W|}{\log |W_{i}|} \right)^{\gamma} |e^{S_{n}g}|_{C^{0}(W_{i})} \\ & = \sum_{j=1}^{n-1} \sum_{V_{l} \in F_{j}(W)} \sum_{W_{i} \in A_{n-j}(V_{l})} \left( \frac{\log |W|}{\log |W_{i}|} \right)^{\gamma} |e^{S_{n}g}|_{C^{0}(W_{i})} + \sum_{W_{i} \in \mathcal{I}_{n}(W)} \left( \frac{\log |W|}{\log |W_{i}|} \right)^{\gamma} |e^{S_{n}g}|_{C^{0}(W_{i})} \\ & \leq \sum_{j=1}^{n-1} \sum_{V_{l} \in F_{j}(W)} \left( \frac{\log |W|}{\log |V_{l}|} \right)^{\gamma} |e^{S_{j}g}|_{C^{0}(V_{l})} \sum_{W_{i} \in A_{n-j}(V_{l})} \left( \frac{\log |V_{l}|}{\log |W_{i}|} \right)^{\gamma} |e^{S_{n-j}g}|_{C^{0}(W_{i})} \\ & + 2^{((n \vee n_{0})s_{0}+1)\gamma} (Km+1)^{n/m} e^{n \sup g} \end{split}$$

<sup>&</sup>lt;sup>3</sup>Note that "ancestor" refers to the backwards dynamics mapping W to  $W_i$ .

$$\leq 2^{\gamma+1}C\delta^{-1} \sum_{j=1}^{n-1} \sum_{V_l \in F_j(W)} \left( \frac{\log |W|}{\log |V_l|} \right)^{\gamma} |e^{S_j g}|_{C^0(V_l)} \sum_{A \in \mathcal{M}_0^{n-j}} |e^{S_n g}|_{C^0(A)}$$

$$+ 2^{((n \vee n_0)s_0 + 1)\gamma} (Km + 1)^{n/m} e^{n \sup g}$$

$$\leq 2^{2\gamma+1}C\delta^{-1} \sum_{j=1}^{n} 2^{(j \vee n_0)s_0 \gamma} (Km + 1)^{j/m} e^{j \sup g} \sum_{A \in \mathcal{M}_0^{n-j}} |e^{S_n g}|_{C^0(A)}.$$

where we have applied part (a) from time 1 to time j and the first estimate in part (b) from time j to time n, since each  $|V_{\ell}| \ge \delta/2$ .

3.2. Fragmentation Lemmas. Here, given a potential g, we introduce the conditions of Small Singular Pressure (SSP.1 and SSP.2) which are crucial for the existence and the statistical properties of the equilibrium states  $\mu_g$  that will be constructed in Section 6. We prove in Lemma 3.3, Corollary 3.6 and Lemma 3.4 that there exist potentials satisfying simultaneously the conditions SSP.1 and SSP.2. These conditions and their consequences will be used in Section 3.3.

In what follows, we will always assume that the potential g is such that  $P_*(T,g) - \sup g > s_0 \log 2$ . Thus, there exists m large enough such that

$$\frac{1}{m}\log(Km+1) < P_*(T,g) - \sup g - s_0 \log 2,$$

and we choose  $\delta_0 = \delta_0(m)$  to be the corresponding length scale from Lemma 3.1. Notice that m, and therefore also  $\delta_0$ , depend on g.

In order to state the results of this subsection, we give a precise definition of SSP.1. First, we introduce some notations.

Let  $L_u^{\delta}(\mathcal{M}_{-n}^0)$  denote the elements of  $\mathcal{M}_{-n}^0$  whose unstable diameter<sup>4</sup> is at least  $\delta/3$ , for some  $\delta \in (0, \delta_0]$ . Similarly, let  $L_s^{\delta}(\mathcal{M}_0^n)$  denote the elements of  $\mathcal{M}_0^n$  whose stable diameter is at least  $\delta/3$ . Recall that the boundary of the partition formed by  $\mathcal{M}_0^n$  is comprised of stable curves belonging to  $\mathcal{S}_n = \bigcup_{j=0}^n T^{-j}(\mathcal{S}_0) \subset \widehat{\mathcal{W}}^s$ .

Define

$$\begin{split} \ell_n^s(g,\delta) &\coloneqq \inf\{\sum_{V \in L_n^\delta(W)} |e^{S_n g}|_{C^0(V)} \mid W \in \widehat{\mathcal{W}}^s, \ \tfrac{\delta}{3} \leqslant |W| \leqslant \delta\}, \\ \ell_n^u(g,\delta) &\coloneqq \inf\{\sum_{V \in L^\delta(W)} |e^{S_n^{-1} g}|_{C^0(V)} \mid W \in \widehat{\mathcal{W}}^u, \ \tfrac{\delta}{3} \leqslant |W| \leqslant \delta\}, \end{split}$$

where in the second line  $L_n^{\delta}(W)$ ,  $W \in \widehat{\mathcal{W}}^u$ , is defined similarly as in the case  $W \in \widehat{\mathcal{W}}^s$ , but replacing  $T^{-1}$  by T in the definitions, and  $S_n^{-1}g := \sum_{i=1}^n g \circ T^{-i} = S_n g \circ T^{-n}$ .

**Definition 3.2** (SSP.1). A potential g such that  $P(T,g) - \sup g > s_0 \log 2$  is said to have  $\varepsilon$ -SSP.1 (small singular pressure), for some  $\varepsilon > 0$ , if

(3.4) there exist 
$$\delta = \delta(\varepsilon) \in (0, \delta_0]$$
 and  $n_1 \in \mathbb{N}$  such that

<sup>&</sup>lt;sup>4</sup>Recall that the unstable diameter of a set is the length of the longest unstable curve contained in that set.

$$\frac{\sum\limits_{W_i \in L_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}}{\sum\limits_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}} \geqslant \frac{1 - 2\varepsilon}{1 - \varepsilon}, \quad \forall n \geqslant n_1 \ \forall W \in \widehat{\mathcal{W}}^s \ with \ |W| \geqslant \delta/3 \quad ,$$

(3.5) the sequences 
$$(e^{n \sup g} \ell_n^s(g, \delta)^{-1})_{n \geqslant n_1}$$
 and  $(e^{n \sup g} \ell_n^u(g, \delta)^{-1})_{n \geqslant n_1}$  are summable,

and the time reversal of (3.4) holds. More precisely, we call time reversal of (3.4) the same estimate but replacing  $S_n g$  and  $W \in \widehat{\mathcal{W}}^s$  with  $S_n^{-1} g = \sum_{i=1}^n g \circ T^{-i}$  and  $W \in \widehat{\mathcal{W}}^u$ . Notice that (3.4) (resp. its time reversal) implies that  $\ell_n^s(g,\delta)$  (resp.  $\ell_n^u(g,\delta)$ ) is nonzero for all  $n \ge n_1$ . A potential is said to have SSP.1 if it has  $\varepsilon$ -SSP.1 for some  $\varepsilon \le 1/4$ .

The following lemma bootstraps from Lemma 3.1 and will be crucial to get the lower bound on the spectral radius:

**Lemma 3.3.** If g is a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential such that  $P(T, g) - \sup g > s_0 \log 2$  and  $\log \Lambda > \sup g - \inf g$ , then g satisfies (3.4), as well as its time reversal, for all  $\varepsilon > 0$ .

Proof. Fix  $\varepsilon > 0$ . Choose  $n_1$  so large that  $6CC_1^{-1}n_1(Kn_1+1)e^{n_1(\sup g - \inf g - \log \Lambda)} < \varepsilon$ , where C is the constant from Lemma 2.3 and  $C_1$  is such that  $|T^{-n}W| > C_1\Lambda^n|W|$  whenever  $W \in \widehat{\mathcal{W}}^s$ . Next, choose  $\delta > 0$  sufficiently small that if  $W \in \widehat{\mathcal{W}}^s$  with  $|W| < \delta$ , then  $T^{-n}W$  comprises at most Kn + 1 smooth pieces of length at most  $\delta_0$  for all  $0 \le n \le 2n_1$ .

Let  $W \in \widehat{\mathcal{W}}^s$  with  $|W| \ge \delta/3$ . We shall prove the following equivalent inequality for  $n \ge n_1$ :

(3.6) 
$$\frac{\sum\limits_{W_i \in S_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}}{\sum\limits_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}} \leqslant \frac{\varepsilon}{1 - \varepsilon}.$$

For  $n \ge n_1$ , write  $n = kn_1 + l$  for some  $0 \le l < n_1$ . If k = 1, the above inequality is clear since  $S_{n_1+l}^{\delta}(W)$  contains at most  $K(n_1+l)+1$  components by assumption on  $\delta$  and  $n_1$ , while  $|T^{-n_1-l}W| \ge C_1\Lambda^{n_1+l}|W| \ge C_1\Lambda^{n_1+l}\delta/3$ . Thus  $\mathcal{G}_n^{\delta}(W)$  must contain at least  $C_1\Lambda^{n_1+l}/3$  curves since each has length at most  $\delta$ . Thus,

$$\frac{\sum\limits_{W_i \in S_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}}{\sum\limits_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}} \leqslant 3 \frac{K(n_1 + l) + 1}{C_1 \Lambda^{n_1 + l}} \frac{e^{(n_1 + l) \sup g}}{e^{(n_1 + l) \inf g}} \leqslant 6C_1^{-1} (Kn_1 + 1) e^{n_1 (\sup g - \inf g - \log \Lambda)} < \varepsilon,$$

where the last inequality holds by choice of  $n_1$ .

For k > 1, we split  $n = kn_1 + l$ ,  $0 \le l < n_1$ , into k - 1 blocks of length  $n_1$  and the last block of length  $n_1 + l$ . For each  $V \in \mathcal{G}_n^{\delta}(W) \setminus \mathcal{I}_n^{\delta}(W)$ , let j < n be the greatest integer such that  $T^{n-j}V$  is contained in an element  $V_a$  of  $L_j^{\delta}(W)$  and for all j < i < n,  $T^{n-i}V$  is contained in an element of  $S_i^{\delta}(W)$ . We call  $V_a$  the most recent long ancestor of V and j its age. If such a j does not exists, it means that for all i < n,  $T^{n-i}V$  is short, that is  $V \in \mathcal{I}_n^{\delta}(W)$  and we set j = 0 in this case.

We group elements of  $S_n^{\delta}(W)$  by their age in  $[0, n_1 - 1], [n_1, 2n_1 - 1], \dots, [(k-2)n_1, (k-1)n_1 - 1]$  and  $[(k-1)n_1, n-1]$ . In other words, we consider the following decomposition

$$(3.7) S_n^{\delta}(W) = \bigsqcup_{q=0}^{k-2} \left( \bigsqcup_{j=qn_1}^{(q+1)n_1-1} \bigsqcup_{V \in L_j^{\delta}(W)} \mathcal{I}_{n-j}^{\delta}(V) \right) \sqcup \left( \bigsqcup_{j=(k-1)n_1}^{n-1} \bigsqcup_{V \in L_j^{\delta}(W)} \mathcal{I}_{n-j}^{\delta}(V) \right).$$

We can therefore split the left hand side of (3.6) into two manageable parts. For this, we rely on Lemma 3.1 for  $\gamma = 0$  and the fact that

$$\mathcal{G}_n^{\delta}(W) \supset \bigsqcup_{V \in L_j^{\delta}(W)} \mathcal{G}_{n-j}^{\delta}(V), \quad \forall \, 0 < j < n.$$

Thus, using Lemma 2.3, we have

$$\frac{\sum_{q=0}^{k-2} \sum_{j=qn_{1}}^{(q+1)n_{1}-1} \sum_{V \in L_{j}^{\delta}(W)} \sum_{W_{i} \in \mathcal{I}_{n-j}^{\delta}(V)} |e^{S_{n}g}|_{C^{0}(W_{i})}}{\sum_{W_{i} \in \mathcal{G}_{kn_{1}+l}^{\delta}(W)} |e^{S_{n}g}|_{C^{0}(W_{i})}} \leq \sum_{q=0}^{k-2} \sum_{j=qn_{1}}^{(q+1)n_{1}-1} \frac{\sum_{V \in L_{j}^{\delta}(W)} |e^{S_{j}g}|_{C^{0}(V)} \sum_{W_{i} \in \mathcal{I}_{n-j}^{\delta}(V)} |e^{S_{n-j}g}|_{C^{0}(W_{i})}}{C^{-1} \sum_{V \in L_{j}^{\delta}(W)} |e^{S_{j}g}|_{C^{0}(V)} \sum_{W_{i} \in \mathcal{G}_{n-j}^{\delta}(V)} e^{(n-j)\inf g}}$$

$$\leq \sum_{q=0}^{k-2} 6CC_{1}^{-1} n_{1}(Kn_{1}+1)^{k-q} e^{(k-q)n_{1}(\sup g-\inf g-\log \Lambda)}$$

$$\leq \sum_{q=0}^{k-2} \varepsilon^{k-q} = \sum_{q=2}^{k} \varepsilon^{q}$$

Similarly, for the second part we have

$$\frac{\sum\limits_{j=(k-1)n_1}^{n-1}\sum\limits_{V\in L_j^{\delta}(W)}\sum\limits_{W_i\in\mathcal{I}_{n-j}^{\delta}(V)}|e^{S_ng}|_{C^0(W_i)}}{\sum\limits_{W_i\in\mathcal{G}_{kn_1+l}^{\delta}(W)}|e^{S_ng}|_{C^0(W_i)}} \leqslant \sum\limits_{j=(k-1)n_1}^{n-1}\frac{\sum\limits_{V\in L_j^{\delta}(W)}|e^{S_jg}|_{C^0(V)}\sum\limits_{W_i\in\mathcal{I}_{n-j}^{\delta}(V)}|e^{S_{n-j}g}|_{C^0(W_i)}}{C^{-1}\sum\limits_{V\in L_j^{\delta}(W)}|e^{S_jg}|_{C^0(V)}\sum\limits_{W_i\in\mathcal{G}_{n-j}^{\delta}(V)}|e^{S_{n-j}g}|_{C^0(W_i)}} \leqslant 6CC_1^{-1}n_1(Kn_1+1)e^{n_1(\sup g-\inf g-\log \Lambda)} \leqslant \varepsilon$$

Summing these two estimates, we obtain (3.6).

The time reversal is obtained from the same proof by changing the construction of the set  $\mathcal{G}_n^{\delta}(W)$  (and thus  $L_n^{\delta}(W)$ ,  $S_n^{\delta}(W)$  and  $\mathcal{I}_n^{\delta}(W)$ ) so that elements of  $\mathcal{G}_n^{\delta}(W)$  are contained in  $T^nW$  (instead of  $T^{-n}(W)$ ) for  $W \in \widehat{\mathcal{W}}^u$ .

Notice that if  $\varepsilon \leqslant 1/4$  and  $\delta_1 \leqslant \delta_0$  and  $n_1$  are the corresponding  $\delta$  and  $n_1$  from the  $\varepsilon$ -SSP.1 condition, then we have for all  $W \in \widehat{\mathcal{W}}^s$  with  $|W| \geqslant \delta_1/3$  and  $n \geqslant n_1$ 

(3.8) 
$$\sum_{W_i \in L_n^{\delta_1}(W)} |e^{S_n g}|_{C^0(W_i)} \geqslant \frac{2}{3} \sum_{W_i \in \mathcal{G}_n^{\delta_1}(W)} |e^{S_n g}|_{C^0(W_i)},$$

In particular, since  $\mathcal{G}_n^{\delta_1}(W) = L_n^{\delta_1}(W) \sqcup S_n^{\delta_1}(W)$ , we also get that

$$\sum_{W_i \in L_n^{\delta_1}(W)} |e^{S_n g}|_{C^0(W_i)} \geqslant 2 \sum_{W_i \in S_n^{\delta_1}(W)} |e^{S_n g}|_{C^0(W_i)}.$$

The following lemma will be used to get both lower and upper bounds on the spectral radius via Proposition 3.7:

**Lemma 3.4.** Let g be a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential such that  $P_*(T, g) - \sup g > s_0 \log 2$  and which has SSP.1. Let  $\delta_1$  and  $n_1$  be the corresponding parameters associated with SSP.1. Then there exist  $C_{n_1} > 0$  and  $n_2 \ge n_1$  such that for all  $n \ge n_2$ ,

$$\sum_{A \in L_{u}^{\delta_{1}}(\mathcal{M}_{-n}^{0})} |e^{S_{n}^{-1}g}|_{C^{0}(A)} \geqslant C_{n_{1}}\delta_{1} \sum_{A \in \mathcal{M}_{-n}^{0}} |e^{S_{n}^{-1}g}|_{C^{0}(A)},$$

$$\sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n})} |e^{S_{n}g}|_{C^{0}(A)} \geqslant C_{n_{1}}\delta_{1} \sum_{A \in \mathcal{M}_{0}^{n}} |e^{S_{n}g}|_{C^{0}(A)}.$$

Furthermore, if g is a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential with  $P_*(T, g) - \sup g > s_0 \log 2$  and  $\log \Lambda > \sup g - \inf g$ , then g has SSP.1.

*Proof.* We prove the lower bound for  $L_s^{\delta_1}(\mathcal{M}_0^n)$ . The lower bound for  $L_u^{\delta_1}(\mathcal{M}_{-n}^0)$  then follows by time reversal.

First, we need to define sets that will be relevant only here. Let

$$I_s(\mathcal{M}_0^n) := \{ A \in \mathcal{M}_0^n \mid \operatorname{diam}^s(A) < \delta_1/3 \}$$

be the complement of  $L_s^{\delta_1}(\mathcal{M}_0^n)$  in  $\mathcal{M}_0^n$ , and

 $I_s(T^{-j}S_0) := \{\text{unstable curves in } T^{-j}(S_0) \text{ with length less than } \delta_1/3\}.$ 

Define also  $L_s(T^{-j}S_0)$  as the complement of  $I_s(T^{-j}S_0)$  in  $\mathcal{G}_j^{\delta_1}(S_0)$ .

We will deduce the claim by estimating the sum of norms of  $e^{S_n g}$  over  $I_s(\mathcal{M}_0^n)$  by the one over  $L_s^{\delta_1}(\mathcal{M}_0^n)$ . To do so, we estimate the sum over  $I_s(\mathcal{M}_0^n)$  with the sums over  $I_s(T^{-j}S_0)$ . Then, using (3.4) we estimate the sum over  $I_s(T^{-j}S_0)$  with sums over  $L_s(T^{-j}S_0)$ . Finally, we estimate sums over  $L_s(T^{-j}S_0)$  with a sum over  $L_s^{\delta_1}(\mathcal{M}_0^n)$ .

In order to estimate the sum over  $I_s(\mathcal{M}_0^n)$ , first remark that if  $A \in \mathcal{M}_0^n$  then  $\partial A \subset \mathcal{S}_n = \bigcup_{j=0}^n T^{-j} \mathcal{S}_0$ . Let  $A \in I_s(\mathcal{M}_0^n)$ . We distinguish two cases:

- (a) For some  $1 \leq j \leq n$ ,  $\partial A$  contains a point of intersection between two curves of  $T^{-j}S_0$ . Since such intersection point is the image by  $T^{-j+1}$  of an intersection point between curves of  $T^{-1}S_0$ , which are finite, and thank to the linear complexity (3.1), we get that there are at most  $K_2n$  elements of  $I_s(\mathcal{M}_0^n)$  in this case.
- (b)  $\partial A$  only contains intersection points between curves belonging to  $T^{-j}\mathcal{S}_0$  for different j. Let  $j_A$  be the maximal  $1 \leq j \leq n$  such that  $A \cap T^{-j}\mathcal{S}_0 \neq \emptyset$ , and  $\gamma \in T^{-j_A}\mathcal{S}_0$  such that  $\gamma \cap A \neq \emptyset$ . Notice that  $\gamma$  must intersect other curves from  $\partial A$ . These curves belong to  $T^{-j}\mathcal{S}_0$  for some  $j < j_A$ . Applying  $T^j$ , it appears that  $\gamma$  must terminates at these intersection points, and thus  $\gamma \subset \partial A$ . Since  $\gamma$  is a stable curve,  $\gamma$  belongs to  $I_s(T^{-j_A}\mathcal{S}_0)$  by assumption on A. Finally, such a curve  $\gamma$  belong to at most 2 elements of  $I_s(\mathcal{M}_0^n)$ .

Therefore

$$(3.9) \qquad \sum_{A \in I_s(\mathcal{M}_0^n)} |e^{S_n g}|_{C^0(A)} \leqslant K_2 n e^{n \sup g} + C \sum_{j=1}^n \sum_{W \in I_s(T^{-j}\mathcal{S}_0)} |e^{S_n g}|_{C_+^0(W)} + |e^{S_n g}|_{C_-^0(W)},$$

where we have extended  $e^{S_n g}$  by Hölder continuity to W from both sides – and noted  $|\cdot|_{C^0_+(W)}$  and  $|\cdot|_{C^0_-(W)}$  the corresponding norms – and C is the constant from Lemma 2.3.

In order to use (3.4), we decompose  $S_0 = \bigsqcup_{i=1}^{l_0} U_i$  where each  $U_i$  is a connected curve such that  $\frac{\delta_1}{3} \leqslant |T^{-1}U_i| \leqslant \delta_1$ . But first we need to compare the sum indexed by  $I_s(T^{-j}S_0)$  with the one indexed by  $I_s(\mathcal{G}_{j-1}^{\delta_1}(U_i))$ . Let  $W \in I_s(T^{-j}S_0)$ . Thus, each  $W \cap T^{-j}U_i$  is a single maximal smooth component of length less than  $\delta_1/3$ . In other words,  $W \cap T^{-j}U_i \in I_s(\mathcal{G}_{j-1}^{\delta_1}(U_i))$ . Therefore

(3.10) 
$$\sum_{W \in I_s(T^{-j}S_0)} |e^{S_n g}|_{C^0_{\pm}(W)} \leqslant \sum_{i=1}^{l_0} \sum_{W \in I_s(G^{\delta_1}_{i,1}(T^{-1}U_i))} |e^{S_n g}|_{C^0_{\pm}(W)}.$$

Now, using SSP.1 (3.4), in the case  $j > n_1$ , we get that

$$(3.11) \qquad \sum_{W \in I_s(\mathcal{G}_{j-1}^{\delta_1}(T^{-1}U_i))} |e^{S_n g}|_{C_{\pm}^0(W)} \leqslant \frac{1}{2} e^{(n-j+1)\sup g} \sum_{W \in L_s(\mathcal{G}_{j-1}^{\delta_1}(T^{-1}U_i))} |e^{S_{j-1} g}|_{C_{\pm}^0(W)}.$$

In order to estimate this last sum with the sum indexed by  $L_s(\mathcal{G}_{n-1}^{\delta_1}(T^{-1}U_i))$ , notice that

$$L_s(\mathcal{G}_{n-1}^{\delta_1}(T^{-1}U_i)) \supset \bigsqcup_{V \in L_s(\mathcal{G}_{j-1}^{\delta_1}(T^{-1}U_i))} L_s(\mathcal{G}_{n-j}^{\delta_1}(V)).$$

Thus

$$\sum_{W \in L_{s}(\mathcal{G}_{n-1}^{\delta_{1}}(T^{-1}U_{i}))} |e^{S_{n}g}|_{C_{\pm}^{0}(W)} \geqslant \sum_{W \in L_{s}(\mathcal{G}_{j-1}^{\delta_{1}}(T^{-1}U_{i}))} \sum_{V \in L_{s}(\mathcal{G}_{n-j}^{\delta_{1}}(W))} |e^{S_{n-j}g + S_{j}g \circ T^{n-j}}|_{C_{\pm}^{0}(V)}$$

$$\geqslant C^{-2}e^{\inf g} \sum_{W \in L_{s}(\mathcal{G}_{j-1}^{\delta_{1}}(T^{-1}U_{i}))} |e^{S_{j-1}g}|_{C_{\pm}^{0}(W)} \sum_{V \in L_{s}(\mathcal{G}_{n-j}^{\delta_{1}}(W))} |e^{S_{n-j}g}|_{C_{\pm}^{0}(V)}$$

$$\geqslant C^{-2}e^{\inf g} \ell_{n-j}^{s}(g, \delta_{1}) \sum_{W \in L_{s}(\mathcal{G}_{j-1}^{\delta_{1}}(T^{-1}U_{i}))} |e^{S_{j-1}g}|_{C_{\pm}^{0}(W)},$$

where we used Lemma 2.3 for the second inequality, and the definition of  $\ell_{n-j}^s(g,\delta_1)$  for the third inequality. Notice however that (3.4) ensures that  $\ell_{n-j}^s(g,\delta_1) \neq 0$  only for  $n-j \geqslant n_1$ . We will treat these troublesome j in a second time. Assume for now that  $n_1 \leqslant j \leqslant n-n_1$ . Combining the above lower bound with (3.10) and (3.11), we get

$$(3.12) \qquad \sum_{W \in I_s(T^{-j}S_0)} |e^{S_n g}|_{C^0_{\pm}(W)} \leqslant \bar{C} e^{(n-j)\sup g} \ell^s_{n-j}(g, \delta_1)^{-1} \sum_{W \in L_s(T^{-n}S_0)} |e^{S_n g}|_{C^0_{\pm}(W)},$$

where we used that  $\bigsqcup_{i=1}^{l_0} L_s(\mathcal{G}_{j-1}^{\delta_1}(T^{-1}U_i)) \subset L_s(T^{-j}\mathcal{S}_0)$  – which is true if we choose the  $\delta_1$ -scaling  $\mathcal{G}_1(T^{-j}\mathcal{S}_0)$  to be adapted with the decomposition  $\mathcal{S}_0 = \bigsqcup_i U_i$ .

Now, if  $n - n_1 \leq j \leq n$ , then we obtain from similar computations

(3.13)

$$\sum_{W \in I_s(T^{-j}\mathcal{S}_0)} |e^{S_n g}|_{C^0_{\pm}(W)} \leqslant \frac{1}{2} C^2 e^{\inf g} e^{(n-j+1)\sup g} \ell^s_{n_1}(g,\delta_1)^{-1} \sum_{W \in L_s(T^{-j-n_1}\mathcal{S}_0)} |e^{S_{n_1+j}g}|_{C^0_{\pm}(W)}$$

Finally, we estimate the sum over  $L_s(T^{-n}S_0)$  with the sum over  $L_s^{\delta_1}(\mathcal{M}_0^n)$ . To do so, we use similar arguments than for the estimate (3.9). Let  $W \in L_s(T^{-n}S_0)$ . We distinguish the two following cases:

- (a) W intersects another curve from  $T^{-n}S_0$ . There are at most  $2K_2$  elements of  $L_s(T^{-n}S_0)$  in this case,
- (b) W does not intersect other curves from  $T^{-n}S_0$ . In that case, W must be contained in the boundary of an element of  $\mathcal{M}_0^n$ , and thus an element of  $L_s^{\delta_1}(\mathcal{M}_0^n)$ . Now, there are at most  $2C\delta_1^{-1}$  elements of  $L_s(T^{-n}S_0)$  in the boundary of a single element of  $L_s^{\delta_1}(\mathcal{M}_0^n)$ , where C is a large enough constant depending only on the billiard table.

Thus

(3.14) 
$$\sum_{W \in L_s(T^{-n}S_0)} |e^{S_n g}|_{C^0_{\pm}(W)} \leq 2K_2 e^{n \sup g} + C\delta_1^{-1} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_n^n)} |e^{S_n g}|_{C^0(A)}.$$

Similarly, for all  $n - n_1 \leq j \leq n$ ,

$$(3.15) \sum_{W \in L_s(T^{-n_1-j}\mathcal{S}_0)} |e^{S_{n_1+j}g}|_{C^0_{\pm}(W)} \leq 2K_2 e^{(n_1+j)\sup g} + C\delta_1^{-1} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^{n_1+j})} |e^{S_{n_1+j}g}|_{C^0(A)}.$$

Putting together (3.9), (3.12) and (3.14), as well as (3.13) and (3.15), we get

$$\begin{split} \sum_{A \in I_s(\mathcal{M}_0^n)} |e^{S_n g}|_{C^0(A)} &\leqslant K_2 n e^{n \sup g} + C \sum_{j=1}^{n_1-1} \sum_{W \in I_s(T^{-j}S_0)} |e^{S_n g}|_{C^0_+(W)} + |e^{S_n g}|_{C^0_-(W)} \\ &+ C \sum_{j=n_1}^{n-n_1} \sum_{W \in I_s(T^{-j}S_0)} |e^{S_n g}|_{C^0_+(W)} + |e^{S_n g}|_{C^0_-(W)} \\ &+ C \sum_{j=n-n_1+1}^{n} \sum_{W \in I_s(T^{-j}S_0)} |e^{S_n g}|_{C^0_+(W)} + |e^{S_n g}|_{C^0_-(W)} \\ &\leqslant (K_2 n + \bar{C}_{n_1}) e^{n \sup g} + \bar{C} \sum_{j=n_1}^{n-n_1} e^{j \sup g} \ell_j^s(g, \delta_1)^{-1} \sum_{W \in L_s(T^{-n}S_0)} |e^{S_n g}|_{C^0_+(W)} + |e^{S_n g}|_{C^0_-(W)} \\ &+ C \sum_{j=n-n_1+1}^{n} e^{(n-j) \sup g} \ell_{n_1}^s(g, \delta_1)^{-1} \sum_{W \in L_s(T^{-j-n_1}S_0)} |e^{S_j + n_1 g}|_{C^0_+(W)} + |e^{S_j + n_1 g}|_{C^0_-(W)} \\ &\leqslant (K_2 n + \bar{C}_{n_1}) e^{n \sup g} + \tilde{C} n \left( 2K_2 e^{n \sup g} + C \delta_1^{-1} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^n)} |e^{S_n g}|_{C^0(A)} \right) \end{split}$$

$$+\sum_{j=n-n_1+1}^{n} \bar{C}_{n_1} \left( 2K_2 e^{(n_1+j)\sup g} + C\delta_1^{-1} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^{j+n_1})} |e^{S_{j+n_1}g}|_{C^0(A)} \right),$$

where in the last inequality we used (3.5) and the fact that  $n - n_1 + 1 \le j \le n$  is equivalent to  $0 \le n - j \le n_1 - 1$ , that is, in the second sum over j after the second inequality symbol, the  $e^{(n-j)\sup g}$  are uniformly bounded.

We now relate the sum over  $L_s^{\delta_1}(\mathcal{M}_0^{j+n_1})$  to the sum over  $L_s^{\delta_1}(\mathcal{M}_0^n)$ . To do so, notice that if  $A \in L_s^{\delta_1}(\mathcal{M}_0^{j+n_1})$ , then it contains at most  $B^{j+n_1-n}$  elements of  $L_s^{\delta_1}(\mathcal{M}_0^n)$ , where  $B = |\mathcal{P}|$ . On the other hand, an element  $A' \in L_s^{\delta_1}(\mathcal{M}_0^n)$  is contained in exactly one element of  $L_s^{\delta_1}(\mathcal{M}_0^{j+n_1})$ . Thus

$$\sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{j+n_{1}})} |e^{S_{j+n_{1}}g}|_{C^{0}(A)} = \sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{j+n_{1}})} \sum_{A' \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n})} |e^{S_{j+n_{1}}g}|_{C^{0}(A)}$$

$$\leq \sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{j+n_{1}})} \sum_{A' \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n})} e^{n_{1} \sup g} |e^{S_{j}g}|_{C^{0}(A')}$$

$$\leq \sum_{A' \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n})} \sum_{A \in A'} e^{n_{1} \sup g} |e^{S_{j}g}|_{C^{0}(A')}$$

$$\leq B^{j+n_{1}-n} e^{n_{0} \sup g} \sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n})} |e^{S_{j}g}|_{C^{0}(A)},$$

and therefore,

$$\sum_{j=n-n_1+1}^{n} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^{j+n_0})} |e^{S_{j+n_1}g}|_{C^0(A)} \leqslant \sum_{j=n-n_1+1}^{n} B^{j+n_1-n} e^{n_1 \sup g} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^n)} |e^{S_{j}g}|_{C^0(A)}$$

$$\leqslant \sum_{j=n-n_1+1}^{n} B^{j+n_1-n} e^{n_1 \sup g} e^{(n-j) \inf g} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^n)} |e^{S_{n}g}|_{C^0(A)}$$

$$\leqslant \bar{C}_{n_1} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^n)} |e^{S_{n}g}|_{C^0(A)}.$$

Using this last estimates, we obtain

$$\sum_{A \in I_s(\mathcal{M}_0^n)} |e^{S_n g}|_{C^0(A)} \leqslant (K_2 n + \bar{C}_{n_1} + C_{g,\ell} n) e^{n \sup g} + (C_{g,\ell} + \bar{C}_{n_1}) \delta_1^{-1} \sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^n)} |e^{S_n g}|_{C^0(A)},$$

where  $C_{g,\ell}$  is a constant coming from the summability assumption (3.5), and  $\bar{C}_{n_1}$  depends only on  $n_1$  and g.

Finally, since  $I_s(\mathcal{M}_0^n) \sqcup L_s^{\delta_1}(\mathcal{M}_0^n) = \mathcal{M}_0^n$ , we get that

$$\sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^n)} |e^{S_n g}|_{C^0(A)} \geqslant \frac{\sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} - (K_2 n + \bar{C}_{n_1} + C_{g,\ell} n) e^{n \sup g}}{1 + (C_{g,\ell} + \bar{C}_{n_1}) \delta_1^{-1}}.$$

Since  $\lim_{n\to+\infty}\frac{1}{n}\log\sum_{A\in\mathcal{M}_0^n}|e^{S_ng}|_{C^0(A)}=P_*(T,g)$  and by the assumption  $P_*(T,g)>\sup g$ , there is an integer  $n_2$  such that for all  $n\geqslant n_2$ ,

$$\sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} - (K_2 n + \bar{C}_{n_1} + C_{g,\ell} n) e^{n \sup g} \geqslant \frac{1}{2} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)}.$$

Thus, there exists  $C_{n_1} > 0$  such that for all  $n \ge n_2$  the claim holds.

We now prove the second part of Lemma 3.4. Assume that g is a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential with  $P_*(T,g) - \sup g > s_0 \log 2$  and  $\log \Lambda > \sup g - \inf g$ . From the convexity of the topological pressure (Theorem 2.1), we get that  $t \mapsto P_*(T,tg)$  is a convex function. Thus, the map  $t \mapsto P_*(T,t(g-\sup g)) = P_*(T,tg) - t \sup g$  is continuous on [0,1]. Since for all s < t we have

$$\sum_{A \in \mathcal{M}_0^n} |e^{S_n t(g - \sup g)}|_{C^0(A)} \leqslant e^{n(t-s) \sup(g - \sup g)} \sum_{A \in \mathcal{M}_0^n} |e^{S_n s(g - \sup g)}|_{C^0(A)} = \sum_{A \in \mathcal{M}_0^n} |e^{S_n s(g - \sup g)}|_{C^0(A)},$$

the map is nonincreasing. Thus

$$P_*(T,g) - \sup g = P_*(T,g - \sup g) \leqslant P_*(T,0) = h_*,$$

where  $h_*$  is the topological entropy from [3]. Therefore we have  $h_* > s_0 \log 2$  and estimates from [3] can be used. For all  $W \in \widehat{\mathcal{W}}^s$  with  $\delta_1 \geqslant |W| \geqslant \delta_1/3$  and all  $n \geqslant n_1$ ,

$$\sum_{V \in L_n^{\delta_1}(W)} |e^{S_n g}|_{C^0(V)} \geqslant e^{n \inf g} \# L_n^{\delta_1}(W) \geqslant \frac{2}{3} e^{n \inf g} \# \mathcal{G}_n^{\delta_1}(W) \geqslant \frac{2}{3} c_0 e^{n \inf g} \# \mathcal{M}_0^n$$
$$\geqslant \frac{2}{3} c_0 e^{n (\inf g + P_*(T, 0))}$$

where we used [3, Lemma 5.2] for the second inequality, and Propositions 4.6 and 5.5 from [3] in the third inequality<sup>5</sup>.

Thus we get that  $\ell_n^s(g,\delta_1) \geqslant \frac{2}{3}c_0e^{n(\inf g+P_*(T,0))}$ . Since  $\ell_n^s(T,0) = h_* \geqslant \log \Lambda$ , we then get the summability of the sequence  $(e^{n\sup g}\ell_n^s(g,\delta_1)^{-1})_{n\geqslant n_1}$ . The summability of  $e^{n\sup g}\ell_n^u(g,\delta_1)^{-1}$  is obtained similarly by considering lower bounds on  $\#L_u^{\delta_1}(W)$ , also given in [3].

We now introduce the precise definition of SSP.2:

<sup>&</sup>lt;sup>5</sup>We can choose the scale  $\delta_1$  from [3] to agree with the one here. The constant  $c_0$  comes from [3, Proposition 5.5] and depends on  $\delta_1$ .

 $<sup>^6</sup>$ log  $\Lambda$  is a lower bound on the unstable Lyapunov exponent of T. Integrating against  $\mu_{SRB}$  gives the desired inequality.

**Definition 3.5** (SSP.2). A potential g is said to have  $\varepsilon$ -SSP.2 if it has  $\varepsilon$ -SSP.1, if there exists  $\bar{n}_1:(0,+\infty)\to\mathbb{N}$  such that

$$(3.16) \qquad \frac{\sum\limits_{W_i \in L_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}}{\sum\limits_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)}} \geqslant \frac{1 - 3\varepsilon}{1 - \varepsilon}, \quad \forall W \in \widehat{\mathcal{W}}^s, \, \forall n \geqslant \bar{n}_1(|W|),$$

and if the time reversal<sup>7</sup> of (3.16) holds, where  $\delta$  is the corresponding constant from  $\varepsilon$ -SSP.1. A potential is said to have SSP.2 if it has  $\varepsilon$ -SSP.2 for some  $\varepsilon \leq 1/4$ .

Corollary 3.6. If g is a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential such that  $P(T, g) - \sup g > s_0 \log 2$  and  $\log \Lambda > \sup g - \inf g$ , then there exists  $C_2 > 0$  such that g has  $\varepsilon$ -SSP.2 for all  $\varepsilon > 0$  and  $\bar{n}_1(|W|) = C_2 n_1 \frac{|\log(|W|/\delta)|}{|\log \varepsilon|}$ , where  $\delta$  and  $n_1$  are the corresponding constants from Lemma 3.3.

*Proof.* From the Lemmas 3.3 and 3.4, such a potential has SSP.1. We thus only prove (3.16). The proof is essentially the same as the one for Lemma 3.3, except that for curves shorter than  $\delta/3$  one must wait  $n \lesssim |\log(|W|/\delta)|$  for at least one component of  $\mathcal{G}_n^{\delta}(W)$  to belong to  $L_n^{\delta}(W)$ .

More precisely, fix  $\varepsilon > 0$  and the corresponding  $\delta$  and  $n_1$  from Lemma 3.3. Let  $W \in \widehat{W}^s$  with  $|W| < \delta/3$  and take  $n > n_1$ . Decomposing  $\mathcal{G}_n^{\delta}(W)$  and  $\mathcal{S}_n^{\delta}(W)$  as in the proof of Lemma 3.3, we estimate the second part as before. For the first part, we have to split the sum between  $\mathcal{I}_n^{\delta}(W)$  and the rest, which is estimated as before.

For the first part, concerning  $\mathcal{I}_n^{\delta}(W)$ , for  $\delta$  sufficiently small, notice that since the flow is continuous, either  $\#\mathcal{G}_l^{\delta}(W) \leqslant Kl+1$  by (3.1) or at least one element of  $\mathcal{G}_l^{\delta}(W)$  has length at least  $\delta/3$ . Let  $n_2$  denote the first iterate l at which  $\mathcal{G}_l^{\delta}(W)$  contains at least one element of length more than  $\delta/3$ . By the complexity estimate (3.1) and the fact that  $|T^{-n_2}W| \geqslant C_1\Lambda^{n_2}|W|$  by hyperbolicity of T, there exists  $\overline{C}_2 > 0$ , independent of  $W \in \widehat{\mathcal{W}}^s$ , such that  $n_2 \leqslant \overline{C}_2 |\log(|W|/\delta)|$ .

Now, for  $n \ge n_2$ ,

$$\begin{split} \sum_{W_i \in \mathcal{I}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)} &\leqslant \sum_{W' \in \mathcal{G}_{n_2}^{\delta}(W)} |e^{S_{n_2} g}|_{C^0(W')} \sum_{W_i \in \mathcal{I}_{n-n_2}^{\delta}(W')} |e^{S_{n-n_2} g}|_{C^0(W_i)} \\ &\leqslant K(K n_2 + 1) e^{n_2 \sup g} \times 2(K n_1 + 1)^{\frac{n-n_2}{n_1}} e^{(n-n_2) \sup g} \end{split}$$

and by hyperbolicity and Lemma 2.3,

$$\begin{split} \sum_{W_i \in \mathcal{G}_n^{\delta}(W)} |e^{S_n g}|_{C^0(W_i)} \geqslant C^{-1} |e^{S_{n_2} g}|_{C^0(W')} \sum_{W_i \in \mathcal{G}_{n-n_2}^{\delta}(W')} e^{(n-n_2)\inf g} \\ \geqslant \frac{1}{3} C_1 C^{-1} e^{n_2\inf g} e^{(n-n_2)(\inf g + \log \Lambda)} \end{split}$$

<sup>&</sup>lt;sup>7</sup>As for (3.4), we call time reversal of (3.16) the same estimate but with  $S_n g$  and  $W \in \widehat{\mathcal{W}}^s$  replaced by  $S_n^{-1} g$  and  $W \in \widehat{\mathcal{W}}^u$ .

where  $W' \in \mathcal{G}_{n_2}^{\delta}(W)$  is such that  $|W'| > \delta/3$ . Therefore,

$$\frac{\sum\limits_{W_{i}\in\mathcal{I}_{n}^{\delta}(W)}|e^{S_{n}g}|_{C^{0}(W_{i})}}{\sum\limits_{W_{i}\in\mathcal{G}_{n}^{\delta}(W)}|e^{S_{n}g}|_{C^{0}(W_{i})}} \leq 6C_{1}^{-1}Ce^{n_{2}(\sup g-\inf g)}K(Kn_{2}+1)(Kn_{1}+1)^{\frac{n-n_{2}}{n_{1}}}e^{(n-n_{2})(\sup g-\inf g-\log \Lambda)}$$

$$\leq 2c_0^{-1}C^2e^{n_2(\sup g-\inf g)}K(Kn_2+1)\varepsilon^{n/n_1}.$$

Since  $n_2 \leqslant \overline{C}_2 |\log(|W|/\delta)|$ , we can bound this expression by  $\varepsilon$  by choosing some  $C_2 > 0$  and n large enough so that  $n/n_1 \geqslant C_2 \frac{\log(|W|/\delta)}{\log \varepsilon}$ . For such n, the left hand side of (3.6) is bounded by  $\varepsilon + \frac{\varepsilon}{1-\varepsilon} \leqslant \frac{2\varepsilon}{1-\varepsilon}$ , which completes the proof of the corollary.

As usual, the time reversal of (3.16) is obtained by performing the same proof, but with the time reversal counterpart of  $\mathcal{G}_n^{\delta}(W)$ , for unstable curves W.

3.3. Exact Exponential Growth of Thermodynamic Sums – Cantor Rectangles. It follows from the submultiplicativity in the characterisation of  $P_*(T,g)$  that

$$e^{nP_*(T,g)} \leqslant e^{-\inf g} \sum_{A \in \mathcal{M}_0^n} \sup_{x \in A} e^{(S_ng)(x)}$$

for all n. In this subsection, we shall prove a supermultiplicativity statement (Lemma 3.9) from which we deduce the upper bound for  $\sum_{A \in \mathcal{M}_0^n} \sup_{x \in A} e^{(S_n g)(x)}$  in Proposition 3.10 giving the upper bound in Proposition 5.1, and ultimately the upper bound on the spectral radius of  $\mathcal{L}_g$  on  $\mathcal{B}$ .

The following key estimate is a lower bound on the weighted rate of growth of stable curves having a certain length. The proof will crucially use the fact that the SRB measure is mixing in order to bootstrap from SSP.1.

**Proposition 3.7.** Let g be a  $(\mathcal{M}_0^1, \alpha)$ -Hölder potential with  $P_*(T, g) - \sup g > s_0 \log 2$  and which has SSP.1. Let  $\delta_1$  be the value of  $\delta$  from the condition SSP.1. Then there exists  $c_0 > 0$  such that for all  $W \in \widehat{\mathcal{W}}^s$  with  $|W| \ge \delta_1/3$  and  $n \ge 1$ , we have

$$\sum_{W_i \in \mathcal{G}_n^{\delta_0}(W)} |e^{S_n g}|_{C^0(A)} \geqslant c_0 \sum_{A \in \mathcal{M}_{-n}^0} |e^{S_n^{-1} g}|_{C^0(A)}.$$

The constant  $c_0$  depends on  $\delta_1$ .

The proof relies crucially on the notion of *Cantor rectangles*. We introduce this notion as in [3, Definition 5.7]. Let  $W^s(x)$  and  $W^u(x)$  denote the maximal smooth components of the local stable and unstable manifolds of  $x \in M$ .

**Definition 3.8.** A solid rectangle D in M is a closed connected set whose boundary comprises precisely four nontrivial curves: two stable manifolds and two unstable manifolds. Given a solid rectangle D, the (locally maximal) Cantor rectangle R in D is formed by taking the points in D whose local stable and unstable manifolds completely cross D. Cantor rectangles have a natural product structure: for any  $x, y \in R$ , then  $W^s(x) \cap W^u(y) \in R$ . In [13, Section 7.11], Cantor rectangles are proved to be closed, and thus contain their outer boundaries, which are contained in the boundary of D. With a slight abuse, we will call this pairs of stable and

unstable manifolds the stable and unstable boundaries of R. In this case, we denote D by D(R) to emphasize that it is the smallest solid rectangle containing R.

*Proof.* Using [13, Lemma 7.87], we may cover M by Cantor rectangles  $R_1, ..., R_k$  satisfying

(3.17) 
$$\inf_{x \in R_i} \frac{m_{W^u}(W^u(x) \cap R_i)}{m_{W^u}(W^u(x) \cap D(R_i))} \geqslant 0.9, \quad \forall 1 \leqslant i \leqslant k,$$

whose stable and unstable boundaries have lengths at most  $\frac{1}{10}\delta_1$ , with the property that any stable curve of length at least  $\delta_1/3$  properly crosses at least one of them. A stable curve  $W \in \widehat{W}^s$  is said to properly cross R if W crosses both unstable sides of R, W does not cross any stable manifolds  $W^s(x) \cap D(R)$  for  $x \in R$ , and the point  $W \cap W^u(x)$  subdivides the curve  $W^u(x) \cap D(R)$  in a ratio between 0.1 and 0.9 (i.e. W does not come to close to either stable boundary of R). The cardinality K is fixed, depending only on K1.

Recall that  $L_u^{\delta_1}(\mathcal{M}_{-n}^0)$  denotes the elements of  $\mathcal{M}_{-n}^0$  whose unstable diameter is longer than  $\delta_1/3$ . We claim that for all  $n \in \mathbb{N}$ , at least one  $R_i$  is fully crossed in the unstable direction by a subset  $\tilde{L}$  of  $\mathcal{M}_{-n}^0$  such that

(3.18) 
$$\sum_{A \in \tilde{L}} |e^{S_n^{-1}g}|_{C^0(A)} \geqslant \frac{1}{k} \sum_{A \in L_u^{\delta_1}(\mathcal{M}_{-n}^0)} |e^{S_n^{-1}g}|_{C^0(A)}.$$

Notice that if  $A \in \mathcal{M}_{-n}^0$ , then  $\partial A$  is comprised of unstable curves belonging to  $\bigcup_{i=1}^n T^i \mathcal{S}_0$ , and possibly  $\mathcal{S}_0$ . By definition of unstable manifolds,  $T^i \mathcal{S}_0$  cannot intersect the unstable boundaries of the  $R_i$ ; thus if  $A \cap R_i \neq \emptyset$ , then either  $\partial A$  terminates inside  $R_i$  or A fully crosses  $R_i$ . Thus elements of  $L_u^{\delta_1}(\mathcal{M}_{-n}^0)$  fully cross at least one  $R_i$  and so at least one  $R_i$  must be fully crossed by a large fraction  $\tilde{L}$  of  $L_u^{\delta_1}(\mathcal{M}_{-n}^0)$  in the sense of (3.18), proving the claim.

For each  $n \in \mathbb{N}$ , denote by  $i_n$  the index of a rectangle  $R_{i_n}$  which is fully crossed by a large enough subset  $\tilde{L}_n$  of  $L_u(\mathcal{M}_{-n}^0)$ , in the sense of (3.18).

Fix  $\delta_* \in (0, \delta_1/10)$  and for i=1,...k, choose a "high density" subset  $R_i^* \subset R_i$  satisfying the following conditions:  $R_i^*$  has a non-zero Lebesgue measure, and for any unstable manifold  $W^u$  such that  $W^u \cap R_i^* \neq \emptyset$  and  $|W^u| < \delta_*$ , we have  $\frac{m_{W^u}(W^u \cap R_i^*)}{|W^u|} \geqslant 0.9$ . (Such a  $\delta_*$  and  $R_i^*$  exist due to the fact that  $m_{W^u}$ -almost every  $y \in R_i$  is a Lebesgue density point of the set  $W^u(y) \cap R_i$  and the unstable foliation is absolutely continuous with respect to  $\mu_{\text{SRB}}$  or, equivalently, Lebesgue.)

Due to the mixing property of  $\mu_{\text{SRB}}$  and the finiteness of the number of rectangles  $R_i$ , there exist  $\varepsilon > 0$  and  $n_3 \in \mathbb{N}$  such that for all  $1 \leq i, j \leq k$  and all  $n \geq n_3$ ,  $\mu_{\text{SRB}}(R_i^* \cap T^{-n}R_j) \geq \varepsilon$ . If necessary, we increase  $n_3$  so that the unstable diameter of the set  $T^{-n}R_i$  is less than  $\delta_*$  for each i, and  $n \geq n_3$ .

Now let  $W \in \widehat{\mathcal{W}}^s$  with  $|W| \ge \delta_1/3$  be arbitrary. Let  $R_j$  be a Cantor rectangle that is properly crossed by W. Let  $n \in \mathbb{N}$  and let  $i_n$  be as above. By mixing,  $\mu_{\text{SRB}}(R_{i_n}^* \cap T^{-n_3}R_j) \ge \varepsilon$ . By [13, Lemma 7.90], there is a component of  $T^{-n_3}W$  that fully crosses  $R_{i_n}^*$  in the stable direction. Call this component  $V \in \mathcal{G}_{n_3}^{\delta_0}(W)$ . Thus

$$\sum_{W_i \in \mathcal{G}_n^{\delta_0}(V)} |e^{S_n g}|_{C^0(W_i)} = \sum_{W_i \in \mathcal{G}_n^{\delta_0}(V)} |e^{S_n^{-1} g}|_{C^0(T^n W_i)} \geqslant \sum_{A \in \tilde{L}_n} \inf_A |e^{S_n^{-1} g}| \geqslant \frac{1}{C_g} \sum_{A \in \tilde{L}_n} \sup_A |e^{S_n^{-1} g}|$$

$$\geqslant \frac{1}{k} \sum_{A \in L_u^{\delta_1}(\mathcal{M}_{-n}^0)} |e^{S_n^{-1}g}|_{C^0(A)}.$$

We now have to relate the lhs to the analogous quantity where V is replace by W.

$$\begin{split} & \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} |e^{S_{n}g}|_{C^{0}(W_{i})} = \sum_{V_{j} \in \mathcal{G}_{n+n_{3}}^{\delta_{0}}(W)} \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} \frac{|e^{S_{n}g}|_{C^{0}(W_{i})}}{\#\{V_{j} \in \mathcal{G}_{n+n_{3}}^{\delta_{0}}(W) \mid T^{n_{3}}V_{j} \subset W_{i}\}} \\ & \geqslant \sum_{V_{j} \in \mathcal{G}_{n+n_{3}}^{\delta_{0}}(W)} |e^{S_{n}g \circ T^{n_{3}}}|_{C^{0}(V_{j})} \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W) \atop T^{n_{3}}V_{j} \subset W_{i}} \frac{1}{\#\{V_{j} \in \mathcal{G}_{n+n_{3}}^{\delta_{0}}(W) \mid T^{n_{3}}V_{j} \subset W_{i}\}} \\ & \geqslant \frac{C\delta_{0}}{\#\mathcal{M}_{0}^{n_{3}}} e^{-n_{3}\sup g} \sum_{V_{j} \in \mathcal{G}_{n+n_{3}}^{\delta_{0}}(W)} |e^{S_{n+n_{3}}g}|_{C^{0}(V_{j})} \geqslant \frac{C\delta_{0}}{\#\mathcal{M}_{0}^{n_{3}}} e^{-n_{3}\sup g} \sum_{V_{j} \in \mathcal{G}_{n}^{\delta_{0}}(V)} |e^{S_{n+n_{3}}g}|_{C^{0}(V_{j})} \\ & \geqslant \frac{C\delta_{0}}{\#\mathcal{M}_{0}^{n_{3}}} e^{-n_{3}(\sup g - \inf g)} \sum_{V_{j} \in \mathcal{G}_{n}^{\delta_{0}}(V)} |e^{S_{n}g}|_{C^{0}(V_{j})} \geqslant \frac{1}{kC_{g}} \frac{C\delta_{0}}{\#\mathcal{M}_{0}^{n_{3}}} e^{-n_{3}(\sup g - \inf g)} \sum_{A \in L_{u}^{\delta_{1}}(\mathcal{M}_{-n}^{0})} |e^{S_{n}^{-1}g}|_{C^{0}(A)} \\ & \geqslant C_{n_{1}}\delta_{1} \frac{1}{kC_{g}} \frac{C\delta_{0}}{\#\mathcal{M}_{0}^{n_{3}}} e^{-n_{3}(\sup g - \inf g)} \sum_{A \in \mathcal{M}_{-n}^{0}} |e^{S_{n}^{-1}g}|_{C^{0}(A)}, \end{split}$$

for all  $n \ge \max\{n_2, n_3\}$ , where we used Lemma 3.4 for the last inequality. Thus the proposition holds for all  $n \ge \max\{n_2, n_3\}$ . It extends to all  $n \in \mathbb{N}$  since there are finitely many values of n to correct for.

**Lemma 3.9** (Supermultiplicativity). There exists a constant  $c_1$  such that for all  $n \in \mathbb{N}$ , and all 0 < j < n, we have

$$\sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} \geqslant c_1 \sum_{A \in \mathcal{M}_0^{n-j}} |e^{S_{n-j} g}|_{C^0(A)} \sum_{A \in \mathcal{M}_{-j}^0} |e^{S_j^{-1} g}|_{C^0(A)}.$$

*Proof.* Fix  $n, j \in \mathbb{N}$  with j < n. First, notice that

$$\sum_{A \in \mathcal{M}_{0}^{n}} |e^{S_{n}g}|_{C^{0}(A)} \geqslant \sum_{A \in \mathcal{M}_{0}^{n}} \inf_{A} e^{(S_{n-j}g + S_{j}^{-1}g) \circ T^{j}} \geqslant \sum_{A \in \mathcal{M}_{-j}^{n-j}} \inf_{A} e^{S_{n-j}g} \inf_{A} e^{S_{j}^{-1}g}$$

$$\geqslant \sum_{A \in \mathcal{M}_{0}^{n-j}} \inf_{A} e^{S_{n-j}g} \sum_{\substack{B \in \mathcal{M}_{-j}^{0} \\ B \cap A \neq \emptyset}} \inf_{B} e^{S_{j}^{-1}g}$$

$$\geqslant C_{g}^{2} \sum_{A \in \mathcal{M}_{0}^{n-j}} |e^{S_{n-j}g}|_{C^{0}(A)} \sum_{\substack{B \in \mathcal{M}_{-j}^{0} \\ B \cap A \neq \emptyset}} |e^{S_{j}^{-1}g}|_{C^{0}(B)}$$

$$\geqslant C_{g}^{2} \sum_{A \in \mathcal{M}_{0}^{n-j}} |e^{S_{n-j}g}|_{C^{0}(A)} \sum_{\substack{B \in \mathcal{M}_{-j}^{0} \\ B \cap A \neq \emptyset}} |e^{S_{j}^{-1}g}|_{C^{0}(B)},$$

where we used Lemma 2.3 twice for the forth inequality.

Recall that  $L_u^{\delta_1}(\mathcal{M}_{-j}^0)$  denotes the elements of  $\mathcal{M}_{-j}^0$  whose unstable diameter is longer than  $\delta_1/3$ . Similarly,  $L_s^{\delta_1}(\mathcal{M}_0^{n-j})$  denotes those elements of  $\mathcal{M}_0^{n-j}$  whose stable diameter is larger than  $\delta_1/3$ . By Lemma 3.4

$$\sum_{A \in L_s^{\delta_1}(\mathcal{M}_0^{n-j})} |e^{S_{n-j}g}|_{C^0(A)} \geqslant C_{n_1} \delta_1 \sum_{A \in \mathcal{M}_0^{n-j}} |e^{S_{n-j}g}|_{C^0(A)}, \quad \text{for } n-j \geqslant n_2.$$

Let  $A \in L_s(\mathcal{M}_0^{n-j})$  and let  $V_A \in \widehat{\mathcal{W}}^s$  be a stable curve in A with length at least  $\delta_1/3$ . By Proposition 3.7,

$$\sum_{W_i \in \mathcal{G}_i^{\delta_0}(V_A)} |e^{S_j g}|_{C^0(W_i)} \geqslant c_0 \sum_{A \in \mathcal{M}_{-j}^0} |e^{S_j^{-1} g}|_{C^0(W_i)}.$$

Each component of  $\mathcal{G}_{j}^{\delta_{0}}(V_{A})$  corresponds to one component of  $V \setminus \mathcal{S}_{-j}$  (up to subdivision of long pieces in  $\mathcal{G}_{j}^{\delta_{0}}(V_{A})$ ). Thus

$$\sum_{A \in \mathcal{M}_{0}^{n-j}} |e^{S_{n-j}g}|_{C^{0}(A)} \sum_{\substack{B \in \mathcal{M}_{-j}^{0} \\ B \cap A \neq \emptyset}} |e^{S_{j}^{-1}g}|_{C^{0}(B)} \geqslant \sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n-j})} |e^{S_{n-j}g}|_{C^{0}(A)} \sum_{W_{i} \in \mathcal{G}_{j}^{\delta_{0}}(V_{A})} |e^{S_{j}^{-1}g}|_{C^{0}(T^{j}W_{i})}$$

$$\geqslant \sum_{A \in L_{s}^{\delta_{1}}(\mathcal{M}_{0}^{n-j})} |e^{S_{n-j}g}|_{C^{0}(A)} \sum_{W_{i} \in \mathcal{G}_{j}^{\delta_{0}}(V_{A})} |e^{S_{j}g}|_{C^{0}(W_{i})}$$

$$\geqslant C \sum_{A \in \mathcal{M}_{0}^{n-j}} |e^{S_{n-j}g}|_{C^{0}(A)} \sum_{B \in \mathcal{M}_{-j}^{0}} |e^{S_{j}g}|_{C^{0}(B)},$$

proving the lemma with  $c_1 = c_0 C_{n_1} C^2 \delta_1$  when  $n - j \ge n_2$ . For  $n - j \le n_2$ , since

$$\sum_{A \in \mathcal{M}_0^{n-j}} |e^{S_{n-j-1}g}|_{C^0(A)} \leqslant \left(\sum_{A \in \mathcal{M}_0^1} |e^g|_{C^0(A)}\right)^{n-j}$$

we obtain the lemma by decreasing  $c_1$  since there are only finitely many values to correct for.

**Proposition 3.10** (Exact Exponential Growth). Let g be a  $(\mathcal{M}_0^1, \alpha)$ -Hölder continuous potential such that  $P_*(T, g) - \sup g > 0$  and which has SSP.1. Let  $c_1$  be the constant given by Lemma 3.9. Then for all  $n \in \mathbb{N}$ , we have

$$\sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} \leqslant \frac{2}{c_1} e^{nP_*(T,g)}.$$

Proof. Let  $\psi(n) := e^{-nP_*(T,g)} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)}$ . Suppose there exists  $n_1 \in \mathbb{N}$  such that  $\psi(n_1) \geqslant 1$ 

 $2/c_1$ , where  $c_1$  is the constant from Lemma 3.9. Then

$$\psi(2n_1) \geqslant c_1 \psi(n_1)^2 = \frac{1}{c_1} (c_1 \psi(n_1))^2.$$

Integrating this bound, we have inductively for any  $k \ge 1$ ,

$$\psi(2^k n_1) \geqslant \frac{1}{c_1} (c_1 \psi(n_1))^{2^k}.$$

This implies that  $\lim_{k\to+\infty} \frac{1}{2^k n_1} \log \psi(2^k n_1) \geqslant \frac{1}{n_1} \log 2 > 0$ , which contradicts the definition of  $\psi(n)$  (since  $\lim_{n\to+\infty} \frac{1}{n} \log \psi(n) = 0$ ). We conclude that  $\psi(n) \leqslant 2/c_1$  for all  $n \geqslant 1$ .

Remark 3.11. Notice that for g=0, the condition  $P_*(T,g)-\sup g>s_0\log 2$  becomes  $h_*>s_0\log 2$ , where  $h_*$  is the topological entropy of T defined in [3]. This is precisely the condition of sparse recurrence to singularities from [3], and as discussed there, we don't know any example of billiard table not satisfying this condition. Notice that by continuity, if  $h_*>s_0\log 2$  holds, then  $P_*(T,g)-\sup g>s_0\log 2$  holds for all g in a neighbourhood of the zero potential. For potential g close enough to 0, we have  $\log \Lambda>\sup g-\inf g$ . Therefore, by Lemmas 3.3, 3.4 and Corollary 3.6, there exists a neighbourhood of g=0 (in the  $(\mathcal{M}_0^1,\alpha)$ -Hölder topology) in which every potential has SSP.1 and SSP.2, and thus all the consequent results from the present section also hold.

In particular, for any  $t \in \mathbb{R}$  with |t| close enough to zero, the potential  $-t\tau$  has SSP.1 and SSP.2.

3.4. Estimates on norms of the potential. In Section 6, we will need similar estimates as in the present section but with the  $C^0$  norm replaced by the  $C^{\beta}$  norm,  $0 < \beta < 1/3$ . The following lemma shows that previous estimates are still valid up to a multiplicative constant.

**Lemma 3.12.** For every bounded  $(\mathcal{M}_0^1, \alpha)$ -Hölder continuous potential g, there exists C > 0 such that for all  $W \in \mathcal{W}^s$ , all  $n \ge 0$  and all  $W_i \in \mathcal{G}_n^{\delta}(W)$ ,  $|e^{S_n g}|_{C^{\alpha}(W_i)} \le C|e^{S_n g}|_{C^0(W_i)}$ , where  $\delta \in (0, \delta_0]$ .

*Proof.* Let g be such a potential. Let c be such that  $g \geqslant c$ . Let  $W_i \in \mathcal{G}_n^{\delta}(W)$ . Then

$$\begin{split} H_{W_i}^{\alpha}(e^{S_n g}) &\leqslant \sum_{k=0}^{n-1} |e^{-g \circ T^k + S_n g}|_{C^0(W_i)} H_{W_i}^{\alpha}(g \circ T^k) \\ &\leqslant |e^{S_n g}|_{C^0(W_i)} \sum_{k=0}^{n-1} e^{-c} C \Lambda^{-\alpha k} |g|_{C^{\alpha}(M)} \\ &\leqslant |e^{S_n g}|_{C^0(W_i)} C \frac{1}{1 - \Lambda^{\alpha}} e^{-c} |g|_{C^{\alpha}(M)}, \end{split}$$

where for the second inequality we adapted the argument from [3, eq (6.2)], so that

$$\frac{g(T^kx)-g(T^ky)}{d_W(T^kx,T^ky)^\alpha}\frac{d_W(T^kx,T^ky)^\alpha}{d_W(x,y)^\alpha}\leqslant CH^\alpha_{T^kW_i}(g)|J^sT^k|^\alpha_{C^0(W_i)}\leqslant C\Lambda^{-\alpha k}|g|_{C^\alpha(M)}.$$

4. The Banach Spaces  $\mathcal B$  and  $\mathcal B_w$  and the Transfer Operators  $\mathcal L_g$ 

In Section 6, we construct the equilibrium state  $\mu_g$  for T under the potential g out of left and right eigenvectors,  $\tilde{\nu}$  and  $\nu$ , of a transfer operator  $\mathcal{L}_g$  associated with the billiard map and the potential g, acting on suitable Banach spaces  $\mathcal{B}$  and  $\mathcal{B}_w$  of anisotropic distributions. In this section, we define the Banach spaces  $\mathcal{B}$  and  $\mathcal{B}_w$  as well as the transfer operator  $\mathcal{L}_g$ .

4.1. Motivation and heuristics. The spaces  $\mathcal{B}$  and  $\mathcal{B}_w$  are the same as in [3], but we recall their construction not only for completeness, but also to introduce notations. The norms we introduce below are defined by integrating along stable manifolds in  $\mathcal{W}^s$ . We define precisely the notion of distance  $d_{\mathcal{W}^s}(\cdot,\cdot)$  between such curves as well as a distance  $d(\cdot,\cdot)$  defined among functions supported on these curves.

In the setup of uniform hyperbolic dynamic, the relevant transfer operator to study equilibrium states associated to a potential g – see for example [1] – can be defined on measurable function f by

$$\mathcal{L}_g f = \left( e^g \frac{f}{J^s T} \right) \circ T^{-1}$$

where  $J^sT$  is the stable Jacobian of T. Ignoring first the low regularity of  $J^sT$ , we see from the hyperbolicity of T that the composition with  $T^{-1}$  should increase the regularity of f in the unstable direction, while decreasing the regularity in the stable direction. By integrating along stable manifold against the arclength measure, we hope to recover some regularity along the stable manifold – notice that by a change of variable,  $J^sT$  does disappear. Morally, the weak norm  $|\cdot|_w$  and the strong stable norm  $||\cdot||_s$  measure the regularity of the averaged action of  $\mathcal{L}_g$ . On the other hand, the strong unstable norm  $||\cdot||_u$  capture the regularity when passing from a stable manifold to another one. Here, this regularity should be though as a log-scaled Hölder regularity.

4.2. Definition of the Banach spaces and embeddings into distribution. Let  $W^s$  denote the set of all nontrivial connected subsets W of stable manifolds for T so that W has length at most  $\delta_0$ . Such curves have bounded curvature above by fixed constant [13, Prop. 4.29]. Thus  $T^{-1}W^s = W^s$ , up to subdivision of curves. Obviously,  $W^s \subset \widehat{W}^s$ . We define  $W^u$  similarly from unstable manifolds of T.

Given a curve  $W \in \mathcal{W}^s$ , we denote by  $m_W$  the unnormalized Lebesgue (arclength) measure on W, so that  $|W| = m_W(W)$ . Since the stable cone  $C^s$  (2.1) is bounded away from the vertical, we may view each stable manifolds  $W \in \mathcal{W}^s$  as the graph of a function  $\varphi_W(r)$  of the arclength coordinate r ranging over some interval  $I_W$ , that is

$$W = \{G_W(r) := (r, \varphi_W(r)) \mid r \in I_W\}.$$

Given two curves  $W_1, W_2 \in \mathcal{W}^s$ , we may use this representation to define a "distance" between them. Define

$$d_{\mathcal{W}^s}(W_1, W_2) = |I_{W_1} \triangle I_{W_2}| + |\varphi_{W_1} - \varphi_{W_2}|_{C^1(I_{W_1} \cap I_{W_2})}$$

if  $I_{W_1} \cap I_{W_2} \neq \emptyset$ , and  $d_{\mathcal{W}^s}(W_1, W_2) = +\infty$  otherwise.

Similarly, given two test functions  $\psi_1$  on  $W_1$ , and  $\psi_2$  on  $W_2$ , we define a distance between them by

$$d(\psi_1, \psi_2) = |\psi_1 \circ G_{W_1} - \psi_2 \circ G_{W_2}|_{C^0(I_{W_1} \cap I_{W_2})},$$

whenever  $d_{\mathcal{W}^s}(W_1, W_2)$  is finite, and  $d(\psi_1, \psi_2) = +\infty$  otherwise.

<sup>&</sup>lt;sup>8</sup>Actually,  $d_{W^s}$  is not a metric since it does not satisfies the triangle inequality. It is nonetheless sufficient for our purpose to produce a usable notion of a distance between stable manifolds.

We can now introduce the norms used to define the spaces  $\mathcal{B}$  and  $\mathcal{B}_w$ . These norms will depend on the constants  $\epsilon_0 > 0$  and  $\delta_0 \in (0,1)$ , as well as on four positive real numbers  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\zeta$  so that

$$0 < \beta < \alpha \le \min\{1/3, \alpha_q\}, \quad 1 < 2^{s_0 \gamma} < e^{P_*(T, g) - \sup g}, \quad 0 < \zeta < \gamma$$

where g is a given, bounded  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential such that  $P_*(T, g) - \sup g > s_0 \log 2$ .

Remark 4.1. The condition  $\alpha \leq 1/3$  is needed for [3, Lemma 4.4], which is used to prove the embedding into distributions. The number 1/3 comes from the regularity of the density function of the conditional measures in the disintegration of  $\mu_{\text{SRB}}$  against the stable foliation. The bound  $\alpha \leq \alpha_g$  will make possible to see g as an element of  $\mathcal{B}$ . The upper bound on  $\gamma$  arises from the use of the growth lemma 3.1. The dependence on  $\delta_0$  comes from the definition of  $\mathcal{W}^s$ .

For  $f \in C^1(M)$ , define the weak norm of f by

$$|f|_{w} = \sup_{W \in \mathcal{W}^{s}} \sup_{\substack{\psi \in C^{\alpha}(W) \\ |\psi|_{C^{\alpha}(W)} \le 1}} \int_{W} f \,\psi \,\mathrm{d}m_{W}.$$

Similarly, define the strong stable norm of f by  $^9$ 

$$||f||_s = \sup_{W \in \mathcal{W}^s} \sup_{\substack{\psi \in C^{\beta}(W) \\ |\psi|_{C^{\beta}(W)} \le |\log|W||^{\gamma}}} \int_W f \,\psi \,\mathrm{d}m_W,$$

(note that  $|f|_w \leq \max\{1, |\log \delta_0|^{-\gamma}\} ||f||_s$ ). Finally, for  $\varsigma \in (0, \gamma)$ , define the strong unstable norm<sup>10</sup> of f by

$$||f||_{u} = \sup_{\varepsilon \le \varepsilon_{0}} \sup_{\substack{W_{1}, W_{2} \in \mathcal{W}^{s} \\ d_{\mathcal{W}^{s}}(W_{1}, W_{2}) \le \varepsilon \\ |\psi_{i}|_{C^{\alpha}(W_{i})} \le 1 \\ d(\psi_{1}, \psi_{2}) = 0}} |\log \varepsilon|^{\varsigma} \left| \int_{W_{1}} f \psi_{1} dm_{W_{1}} - \int_{W_{2}} f \psi_{2} dm_{W_{2}} \right|.$$

In order to use functional analysis results, we need to work with complete spaces. Since  $C^1(M)$  is not complete for the norms<sup>11</sup>  $|\cdot|_w$  and  $|\cdot|_s + |\cdot|_w$ , we will use the corresponding completed spaces.

**Definition 4.2** (The Banach spaces). The space  $\mathcal{B}_w$  is the completion of  $C^1(M)$  with respect to the weak norm  $|\cdot|_w$ , while  $\mathcal{B}$  is the completion of  $C^1(M)$  with respect to the strong norm,  $||\cdot||_{\mathcal{B}} = ||\cdot||_s + ||\cdot||_u$ . Notice that since  $|\cdot|_w \leq ||\cdot||_{\mathcal{B}}$ , there is a canonical map  $\mathcal{B} \to \mathcal{B}_w$ .

Since the main purpose of the spaces  $\mathcal{B}$  and  $\mathcal{B}_w$  is to contain left and right eigenvectors of a transfer operator acting on those spaces, a crucial feature of  $\mathcal{B}$  and  $\mathcal{B}_w$  is that we can see them as subspaces of the distributional space  $(C^1(M))^*$ . Thanks to this property, we will be able

<sup>&</sup>lt;sup>9</sup>The logarithmic modulus of continuity in  $||f||_s$  is used to obtain a finite spectral radius.

<sup>&</sup>lt;sup>10</sup>The logarithmic modulus of continuity appears in  $||f||_u$  because of the logarithmic modulus of continuity in  $||f||_s$ . Its presence in  $||f||_u$  causes the loss of the spectral gap.

<sup>&</sup>lt;sup>11</sup>For example, the sequence  $\left(\frac{1}{n}\sin 2\pi n^2\frac{r}{|\Gamma_i|}\right)_n$  is a Cauchy sequence of  $C^1(M)$  functions with respect to  $|\cdot|_w$ , but diverges in the  $C^1$ -norm.

to construct a positive distribution by pairing the left and right eigenvectors, and to extend it into the desired equilibrium measure. In order to state this result, we need to introduce some other spaces, on which the transfer operator will be naturally defined (and then extended to  $\mathcal{B}$  and  $\mathcal{B}_w$ ).

Define the usual homogeneity strips

$$\mathbb{H}_k := \left\{ (r, \varphi) \in M_i \mid \frac{\pi}{2} - \frac{1}{k^2} \leqslant \varphi \leqslant \frac{\pi}{2} - \frac{1}{(k+1)^2} \right\}, \quad k \geqslant k_0,$$

and analogously for  $k \leq -k_0$ . Define  $\mathcal{W}_{\mathbb{H}}^s \subset \mathcal{W}^s$  as the set of stable manifolds  $W \in \mathcal{W}^s$  such that  $T^nW$  lies in a single homogeneity strip for all  $n \geq 0$ . We write  $\psi \in C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s)$  if  $\psi \in C^{\alpha}(W)$  for all  $W \in \mathcal{W}_{\mathbb{H}}^s$  with uniformly bounded Hölder norm. The norm of  $\psi$  in  $C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s)$  is defined to be the sup over all the  $C^{\alpha}(W)$  norms, with W ranging in  $\mathcal{W}_{\mathbb{H}}^s$ . Similarly, define the space  $C^{\alpha}_{\cos}(\mathcal{W}_{\mathbb{H}}^s)$  containing the functions  $\psi$  such that  $\psi \cos \varphi \in C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s)$ . The norm of  $\psi$  in  $C^{\alpha}_{\cos}(\mathcal{W}_{\mathbb{H}}^s)$  is defined to be the norm of  $\psi \cos \varphi$  in  $C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s)$ . Clearly,  $C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s) \subset C^{\alpha}_{\cos}(\mathcal{W}_{\mathbb{H}}^s)$ .

The canonical map  $\mathcal{B}_w \to (\mathcal{F})^*$  (for  $\mathcal{F} = C^1(M)$ , or  $\mathcal{F} = C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s)$ ) is understood in the following sense: for  $f \in \mathcal{B}_w$ , there exists  $C_f < \infty$  such that letting  $f_n \in C^1(M)$  be a sequence converging to f in the  $\mathcal{B}_w$  norm, for every  $f \in \mathcal{F}$  the following limit exists

$$f(\psi) := \lim_{n \to +\infty} \int f_n \psi \, \mathrm{d}\mu_{\mathrm{SRB}}$$

and satisfies  $|f(\psi)| \leq C_f ||\psi||_{\mathcal{F}}$ .

We summarize the properties of these Banach spaces obtained in [3] in the following proposition.

**Proposition 4.3.** The spaces  $\mathcal{B}_w$  and  $\mathcal{B}$  are such that:

(i) The following canonical maps are all continuous

$$C^1(M) \to \mathcal{B} \to \mathcal{B}_w \to (C^{\alpha}(\mathcal{W}_{\mathbb{H}}^s))^* \to (C^1(M))^*,$$

and the first two maps are injective. In particular, we also have the two injective and continuous maps

$$(\mathcal{B}_w)^* \to \mathcal{B}^* \to (C^1(M))^*.$$

(ii) The inclusion map  $\mathcal{B} \hookrightarrow \mathcal{B}_w$  is compact.

*Proof.* The point (i) is the content of [3, Proposition 4.2]. We detail the proof of the injectivity of the map  $\mathcal{B} \to \mathcal{B}_w$ . To do so, we prove that the formula defining  $|\cdot|_w$  (respectively  $||\cdot||_s$  and  $||\cdot||_w$ ) can be extended when  $f \in \mathcal{B}_w$  (respectively  $f \in \mathcal{B}$ ), and that it coincides with the norm of f.

First, notice that when  $f \in C^1(M)$ , then for given  $W \in \mathcal{W}^s$  and  $\psi \in C^{\alpha}(W)$  we have  $\int_W f\psi \, \mathrm{d}m_W \leqslant |f|_w |\psi|_{C^{\alpha}(W)}$ . Thus the map  $f \mapsto \int_W f\psi \, \mathrm{d}m_W$  can be extended uniquely to  $\mathcal{B}_w$ .

Now, let  $f \in \mathcal{B}_w$  and  $\varepsilon > 0$ . Let  $f_n$  be a Cauchy sequence of  $C^1(M)$  functions converging to f in  $\mathcal{B}_w$ . Thus, there exists some  $n_{\varepsilon}$  such that for all  $n \ge n_{\varepsilon}$ ,  $|f - f_n|_w \le \varepsilon$ . Let  $W \in \mathcal{W}^s$  and  $\psi \in C^{\alpha}(W)$  with  $|\psi|_{C^{\alpha}(W)} \le 1$ . By definition of  $|f_n|_w$ , for all n, there exist  $W_n$  and  $\psi_n \in C^{\alpha}(W_n)$  with  $|\psi_n|_{C^{\alpha}(W_n)} \le 1$  such that

$$\left| \int_{W_n} f_n \psi_n \, \mathrm{d} m_{W_n} - |f_n|_w \right| \leqslant \varepsilon.$$

Thus, we have

$$\left| \int_{W_n} f \psi_n \, \mathrm{d} m_{W_n} - \int_{W_n} f_n \psi_n \, \mathrm{d} m_{W_n} \right| \leqslant |f - f_n|_w |\psi_n|_{C^{\alpha}(W_n)} \leqslant \varepsilon, \quad \forall n \geqslant n_{\varepsilon},$$

and so  $\left||f_n|_w - \int_{W_n} f\psi_n \, \mathrm{d} m_{W_n}\right| \leqslant 2\varepsilon$ . In particular, we get

$$\sup_{W \in \mathcal{W}^s} \sup_{\substack{\psi \in C^{\alpha}(W) \\ |\psi|_{C^{\alpha}(W)} \leq 1}} \int_W f \, \psi \, \mathrm{d} m_W \geqslant |f|_w.$$

We now prove the reverse inequality. Using the same notations as above, there exist  $V \in \mathcal{W}^s$  and  $\varphi \in C^{\alpha}(V)$  with  $|\varphi|_{C^{\alpha}(V)} \leq 1$  such that

$$\left| \int_{V} f \varphi \, dm_{V} - \sup_{W \in \mathcal{W}^{s}} \sup_{\substack{\psi \in C^{\alpha}(W) \\ |\psi|_{C^{\alpha}(W)} \leq 1}} \int_{W} f \, \psi \, dm_{W} \right| \leq \varepsilon.$$

Now, since

$$\left| \int_{V} f_{n} \varphi \, dm_{V} - \int_{V} f \varphi \, dm_{V} \right| \leq |f - f_{n}|_{w} \leq \varepsilon, \quad \forall n \geqslant n_{\varepsilon},$$

we have that  $|\sup_{W \in \mathcal{W}^s} \sup_{\substack{\psi \in C^{\alpha}(W) \\ |\psi|_{C^{\alpha}(W)} \leq 1}} \int_W f \psi \, \mathrm{d} m_W - \int_V f_n \varphi \, \mathrm{d} m_V| \leq 2\varepsilon$  for all large enough n.

In particular

$$\sup_{W \in \mathcal{W}^s} \sup_{\substack{\psi \in C^{\alpha}(W) \\ |\psi|_{C^{\alpha}(W)} \le 1}} \int_W f \, \psi \, \mathrm{d} m_W \le |f_n|_w + 2\varepsilon.$$

Taking the limit in n, we get the claimed inequality.

The corresponding results for  $f \in \mathcal{B}$  and norms  $||\cdot||_s$  and  $||\cdot||_u$  are obtained similarly, noticing that for all  $f \in C^1(M)$ ,

$$\int_{W} f\psi \, \mathrm{d}m_{W} \leqslant ||f||_{s} |\psi|_{C^{\beta}(W)} |\log |W||^{-\gamma} \leqslant ||f||_{\mathcal{B}} |\psi|_{C^{\beta}(W)} |\log |W||^{-\gamma}, \quad \forall W \in \mathcal{W}^{s}, \forall \psi \in C^{\beta}(W)$$

Thus the integrals against  $C^{\beta}(W)$  functions in the definition of  $||\cdot||_s$  makes sense even when  $f \in \mathcal{B}$ . On the other hand, since  $|\cdot|_w \leq ||\cdot||_{\mathcal{B}}$ , the integrals in the definition of  $||\cdot||_u$  can be extended to  $f \in \mathcal{B}$  as in the above case where  $f \in \mathcal{B}_w$ .

We can now show the injectivity of the canonical map  $\mathcal{B} \to \mathcal{B}_w$ . Let  $f \in \mathcal{B}$  with  $||f||_{\mathcal{B}} \neq 0$ . If  $||f||_s \neq 0$ , then the fact that  $|f|_w \neq 0$  follows from the definition of  $C^{\beta}(W)$  as the closure of  $C^1(W)$  in the  $C^{\beta}$  norm, so that  $C^{\alpha}(W)$  is dense in  $C^{\beta}(W)$ . Now, if  $||f||_u \neq 0$ , then by definition of  $||\cdot||_u$ , we can find some  $W \in \mathcal{W}^s$  and  $\psi \in C^{\alpha}(W)$  so that  $\int_w f\psi \, \mathrm{d}m_W > 0$ . Thus  $|f|_w \neq 0$ .

The point (ii) is precisely the content of [3, Proposition 6.1].

4.3. The transfer operators. We may define the transfer operator  $\mathcal{L}_g: (C_{\cos}^{\alpha}(\mathcal{W}_{\mathbb{H}}^s))^* \to$  $(C^{\alpha}(\mathcal{W}^s))^*$ , for a given weight function g by

$$\mathcal{L}_g f(\psi) = f\left(e^g \frac{\psi \circ T}{J^s T}\right), \quad \psi \in C^{\alpha}(\mathcal{W}^s).$$

This operator is well defined because, if  $\psi \in C^{\alpha}(\mathcal{W}^s)$  then  $e^g \psi \circ T \in C^{\alpha}(\mathcal{W}^s)$ . Furthermore, since  $J^sT$  and  $\cos\varphi$  are 1/3-log-Hölder on homogeneous stable manifolds, and  $\cos\varphi/J^sT$  is bounded away from 0 and  $+\infty$  also on homogeneous stable manifolds, we get that  $1/J^sT \in$  $C_{\cos}^{\alpha}(\mathcal{W}_{\mathbb{H}}^{s})$ . Thus  $e^{g}\frac{\psi \circ T}{J^{s}T} \in C_{\cos}^{\alpha}(\mathcal{W}_{\mathbb{H}}^{s})$ . When  $f \in C^{1}(M)$ , we identify f with the measure<sup>12</sup>

$$(4.1) f\mu_{\text{SRB}} \in (C_{\text{cos}}^{\alpha}(\mathcal{W}_{\mathbb{H}}^{s}))^{*}.$$

The measure above is (abusively) still denoted by f. For  $f \in C^1(M)$ , we have

$$\mathcal{L}_g(f\mu_{\text{SRB}})(\psi) = \int f \, e^g \frac{\psi \circ T}{J^s T} \, d\mu_{\text{SRB}} = \int \left( e^g \frac{f}{J^s T} \right) \circ T^{-1} \psi \, d\mu_{\text{SRB}}$$
$$= \left( \left( e^g \frac{f}{J^s T} \right) \circ T^{-1} \, \mu_{\text{SRB}} \right) (\psi).$$

Thus, due to our identification (4.1) we have  $\mathcal{L}_q f = (e^g f/J^s T) \circ T^{-1}$ , as claimed above.

**Proposition 4.4.** For any fixed  $(\mathcal{M}_0^1, \alpha_q)$ -Hölder potential g and corresponding spaces  $\mathcal{B}$  and  $\mathcal{B}_w$ :

- (i) If  $f \in C^1(M)$ , then  $\mathcal{L}_g(f) \in \mathcal{B}$ . (ii) The operators  $\mathcal{L}_g : (C^1(M), |\cdot|_w) \to \mathcal{B}_w$  and  $\mathcal{L}_g : (C^1(M), ||\cdot||_{\mathcal{B}}) \to \mathcal{B}$  are continuous. In particular,  $\mathcal{L}_q$  extends uniquely into operators on both  $\mathcal{B}_w$  and  $\mathcal{B}$ .

*Proof.* The proof of (i) proceeds similarly as in its analogous result [4, Lemma 4.3] by introducing a mollification  $f_{\eta} \in C^1(M)$  of  $\mathcal{L}_g(f)$ , where  $f \in C^1(M)$  and  $\eta > 0$ . As noticed in [4, Remark 4.11, the proof of [4, Lemma 4.3] can be adapted to the case g = 0. The modification mainly relies on giving a nonhomogenous version of [4, Lemma 4.9], which can easily be done for a weaker – but sufficiently tight – upperbound. The corresponding bounds on  $||\mathcal{L}_q(f) - f_{\eta}||_s$ and  $||\mathcal{L}_g(f) - f_{\eta}||_u$  decrease to 0 as  $\eta$  goes to zero. Now, for any  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential g, one can use same techniques and decompositions to get the bound  $||\mathcal{L}_g(f) - f_{\eta}||_s \leq C_{f,g} \eta^{\alpha_g/6}$ . Concerning the unstable norm, we follow the modifications described in [4, Remark 4.11], but distinguishing between the cases  $\eta^{\alpha_g/6} < |\log \varepsilon|^{-2\zeta}$  and  $\eta^{\alpha_g/6} \ge |\log \varepsilon|^{-2\zeta}$  (instead of  $\eta^{1/6}$ ). In the first case, we obtain a bound  $||\mathcal{L}_g(f) - f_{\eta}||_u \leqslant C_{f,g}\eta^{\alpha_g/12}$ , whereas in the second case (in which  $\varepsilon$  is superexponentially small in  $\eta$ ) we get an upperbound of the form  $C_{f,g}\eta^{-\frac{18+\alpha_g}{12}}\exp\left(-\min(\alpha-\beta,\frac{1}{2})\eta^{-\frac{\alpha_g}{12\zeta}}\right)$ . Therefore, both  $||\mathcal{L}_g(f) - f_{\eta}||_s$  and  $||\mathcal{L}_g(f) - f_{\eta}||_u$ converge to 0 as  $\eta$  goes to zero.

Point (ii) follows from Proposition 5.1, in the case 
$$n = 1$$
.

<sup>&</sup>lt;sup>12</sup>To show the claimed inclusion just use that  $d\mu_{SRB} = (2|\partial Q|)^{-1}\cos\varphi\,drd\varphi$ .

#### 5. Norm Estimates and Spectral Radius

The purpose of this section is to state and prove sharp upper and lower bounds on the norm of the iterated operator  $\mathcal{L}_q^n$ , both in  $\mathcal{B}_w$  and  $\mathcal{B}$ .

**Proposition 5.1.** Let g be a  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder continuous potential. Assume that  $P_*(T, g) - \sup g > s_0 \log 2$  and that SSP.1 holds. Then there exist  $\delta_0$  and C > 0 such that for all  $f \in \mathcal{B}$ ,

(5.1) 
$$|\mathcal{L}_g^n f|_w \leqslant \frac{C}{\delta_0} e^{nP_*(T,g)} |f|_w, \quad \forall n \ge 0 ;$$

(5.2) 
$$\|\mathcal{L}_{g}^{n} f\|_{s} \leqslant \frac{C}{\delta_{0}} e^{nP_{*}(T,g)} \|f\|_{s}, \quad \forall n \geq 0 ;$$

(5.3) 
$$\|\mathcal{L}_g^n f\|_u \leqslant \frac{C}{\delta_0} (\|f\|_u + \|f\|_s) e^{nP_*(T,g)}, \quad \forall n \ge 0.$$

It follows that the spectral radius of  $\mathcal{L}_q$  on  $\mathcal{B}$  and  $\mathcal{B}_w$  is at most  $e^{P_*(T,g)}$ .

Remark 5.2. It is possible to obtain similar estimates without the assumption SSP.1, however an additional factor  $e^{n\varepsilon}$  appears on the right hand sides, for any arbitrary  $\varepsilon > 0$ . We indicate places in the proof where it happens and how to correct for. The conclusion about the upper bound of the spectral radius still holds. Nonetheless, in order to construct nontrivial maximal eigenvectors, we will need the estimates from Proposition 5.1.

**Theorem 5.3.** Let g be a  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder continuous potential. Assume that  $P_*(T, g) - \sup g > s_0 \log 2$  and that SSP.1 holds. Then there exists C such that

$$||\mathcal{L}_g^n 1||_s \geqslant |\mathcal{L}_g^n 1|_w \geqslant C \frac{\delta_1}{2} e^{nP_*(T,g)}.$$

Proof of Proposition 5.1. Let  $\delta_0$  be the scale associated to g as in the beginning of Section 3.2. The set  $\mathcal{W}^s$  is defined with respect to the scale  $\delta_0$ .

We start with the weak norm estimate (5.1). Let  $f \in C^1(M)$ ,  $W \in \mathcal{W}^s$  and  $\psi \in C^{\alpha}(W)$  be such that  $|\psi|_{C^{\alpha}(W)} \leq 1$ . For  $n \geq 0$  we use the definition of the weak norm on each  $W_i \in \mathcal{G}_n^{\delta_0}(W)$  to estimate

$$\int_{W} \mathcal{L}_{g}^{n} f \psi \, \mathrm{d}m_{W} = \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} \int_{W_{i}} f e^{S_{n}g} \psi \circ T^{n} \mathrm{d}m_{W_{i}} \leqslant |f|_{w} \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} |e^{S_{n}g}|_{C^{\alpha}(W_{i})} |\psi \circ T^{n}|_{C^{\alpha}(W_{i})}.$$

Clearly,  $\sup |\psi \circ T^n|_{W_i} \leq \sup_W |\psi|$ . For  $x, y \in W_i$ , we have,

$$(5.4) \qquad \frac{|\psi(T^n x) - \psi(T^n y)|}{d_W(T^n x, T^n y)^{\alpha}} \cdot \frac{d_W(T^n x, T^n y)^{\alpha}}{d_W(x, y)^{\alpha}} \leqslant C|\psi|_{C^{\alpha}(W)} |J^s T^n|_{C^0(W_i)}^{\alpha}$$
$$\leqslant C\Lambda^{-\alpha n} |\psi|_{C^{\alpha}(W)},$$

so that  $H_{W_i}^{\alpha}(\psi \circ T^n) \leqslant C\Lambda^{-\alpha n}H_W^{\alpha}(\psi)$  and thus  $|\psi \circ T^n|_{C^{\alpha}(W_i)} \leqslant C|\psi|_{C^{\alpha}(W)}$ . By Lemma 3.12, we get

$$\int_{W} \mathcal{L}_{g}^{n} f \psi \, dm_{W} \leqslant C |f|_{w} |\psi|_{C^{\alpha}(W)} \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} |e^{S_{n}g}|_{C^{0}(W_{i})} \leqslant \frac{2C}{\delta_{0}} |f|_{w} |\psi|_{C^{\alpha}(W)} \sum_{A \in \mathcal{M}_{0}^{n}} |e^{S_{n}g}|_{C^{0}(A)},$$

$$\leqslant \frac{2C}{c_{1}\delta_{0}} |f|_{w} |\psi|_{C^{\alpha}(W)} e^{nP_{*}(T,g)},$$

where the second inequality uses that there are no more than  $2\delta_0^{-1}$  curves  $W_i$  of  $\mathcal{G}_n^{\delta_0}(W)$  per element of  $\mathcal{M}_0^n$ , and the third inequality uses the Exact Exponential Growth from Proposition  $3.10^{13}$ .

Now we prove the strong stable norm estimate (5.2). We can choose m so large that  $2^{s_0\gamma}(Km+1)^{1/m} < e^{P_*(T,g)-\sup g}$ . Let  $W \in \mathcal{W}^s$ ,  $\psi \in C^{\beta}(W)$  such that  $|\psi|_{C^{\beta}(W)} \leq |\log |W||^{\gamma}$ . Then, by definition of the strong norm

$$\int_{W} \mathcal{L}_{g}^{n} f \psi dm_{W} = \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} \int_{W_{i}} f \psi \circ T^{n} e^{S_{n}g} dm_{W_{i}}$$

$$\leq \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} ||f||_{s} |\psi \circ T^{n}|_{C^{\beta}(W_{i})} |e^{S_{n}g}|_{C^{\beta}(W_{i})} |\log |W_{i}||^{-\gamma}$$

$$\leq C||f||_{s} \sum_{W_{i} \in \mathcal{G}_{n}^{\delta_{0}}(W)} \left(\frac{\log |W|}{\log |W_{i}|}\right)^{\gamma} |e^{S_{n}g}|_{C^{\beta}(W_{i})}$$

$$\leq C||f||_{s} 2^{2\gamma+1} \delta_{0}^{-1} \sum_{j=1}^{n} 2^{js_{0}\gamma} (Km+1)^{j/m} e^{j\sup g} \sum_{A \in \mathcal{M}_{0}^{n-j}} |e^{S_{n-j}g}|_{C^{0}(A)}$$

where for the last line we used Lemma 3.1(b) and Lemma 3.12. Let

$$D_n := C2^{2\gamma+1}\delta_0^{-1} \sum_{j=1}^n 2^{js_0\gamma} (Km+1)^{j/m} e^{j\sup g} \sum_{A \in \mathcal{M}_0^{n-j}} |e^{S_{n-j}g}|_{C^0(A)}.$$

From Proposition 3.10, for all  $n \ge 1$ ,  $\sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} \le \frac{2}{c_1} e^{nP_*(T,g)}$ . Let  $\varepsilon_1 = P_*(T,g) - \sup g - \log(2^{s_0 \gamma} (Km + 1)^{1/m})$ . Thus<sup>14</sup>,

$$D_n \leqslant 2^{2\gamma + 1} \frac{C}{c_1 \delta_0} \sum_{j=1}^n e^{(P_*(T,g) - \varepsilon_1)j} e^{(n-j)P_*(T,g)} \leqslant 2^{2\gamma + 1} \frac{1}{1 - e^{-\varepsilon_1}} \frac{C}{c_1 \delta_0} e^{nP_*(T,g)}.$$

This concludes the proof of (5.2).

Finally, we now prove the strong unstable norm estimate (5.3). Fix  $\tilde{\varepsilon} < \varepsilon_0$ , and consider two curves  $W^1$ ,  $W^2 \in \mathcal{W}^s$  with  $d_{\mathcal{W}^s}(W^1, W^2) < \tilde{\varepsilon}$ .

For  $n \ge 1$ , we describe how to partition  $T^{-n}W^{\ell}$  into "matched" pieces  $U_j^{\ell}$  and "unmatched" pieces  $V_i^{\ell}$ ,  $\ell = 1, 2$ .

Let  $\omega$  be a connected component of  $W^1 \setminus \mathcal{S}_{-n}$ . To each point  $x \in T^{-n}\omega$ , we associate a vertical line segment  $\gamma_x$  of length at most  $C\Lambda^{-n}\tilde{\varepsilon}$  such that its image  $T^n\gamma_x$ , if not cut by a singularity, will have length  $C\tilde{\varepsilon}$ . By [13, §4.4], all the tangent vectors to  $T^i\gamma_x$  lie in the

<sup>&</sup>lt;sup>13</sup>without the assumptions SSP.1 and SSP.3, Proposition 3.10 might not hold. Still, for  $\varepsilon > 0$  and all  $n \ge 1$ ,  $\sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)} \leqslant C_\varepsilon \, e^{n(P_*(T,g)+\varepsilon)} \text{ because of the subadditivity from Theorem 2.1.}$ 

<sup>&</sup>lt;sup>14</sup>Here, again, conclusion from Proposition 3.10 can be replaced.

unstable cone  $C^u(T^ix)$  for each  $i \ge 1$  so that they remain uniformly transverse to the stable cone and enjoy the minimum expansion given by  $\Lambda$ .

Doing this for each connected component of  $W^1 \setminus \mathcal{S}_{-n}$ , we subdivide  $W^1 \setminus \mathcal{S}_{-n}$  into a countable collection of subintervals of points for which  $T^n \gamma_x$  intersects  $W^2 \setminus \mathcal{S}_{-n}$  and subintervals for which this is not the case. This in turn induces a corresponding partition on  $W^2 \setminus \mathcal{S}_{-n}$ .

We denote by  $V_i^{\ell}$  the pieces in  $T^{-n}W^{\ell}$  which are not matched up by this process and note that the images  $T^nV_i^{\ell}$  occur either at the endpoints of  $W^{\ell}$  or because the vertical segment  $\gamma_x$  has been cut by a singularity. In both cases, the length of the curves  $T^nV_i^{\ell}$  can be at most  $C\tilde{\varepsilon}$  due to the uniform transversality of  $\mathcal{S}_{-n}$  with the stable cone and of  $C^s(x)$  with  $C^u(x)$ .

In the remaining pieces the foliation  $\{T^n\gamma_x\}_{x\in T^{-n}W^1}$  provides a one-to-one correspondence between points in  $W^1$  and  $W^2$ . We further subdivide these pieces in such a way that the lengths of their images under  $T^{-i}$  are less than  $\delta_0$  for each  $0\leqslant i\leqslant n$  and the pieces are pairwise matched by the foliation  $\{\gamma_x\}$ . We call these matched pieces  $U_j^\ell$ . Since the stable cone is bounded away from the vertical direction, we can adjust the elements of  $\mathcal{G}_n^{\delta_0}(W^\ell)$  created by artificial subdivisions due to length so that  $U_j^\ell\subset W_i^\ell$  and  $V_k^\ell\subset W_{i'}^\ell$  for some  $W_i^\ell, W_{i'}^\ell\in\mathcal{G}_n^{\delta_0}(W^\ell)$  for all  $j,k\geqslant 1$  and  $\ell=1,2$ , without changing the bounds on sums over  $\mathcal{G}_n^{\delta_0}(W^\ell)$ . There is at most one  $U_j^\ell$  and two  $V_j^\ell$  per  $W_i^\ell\in\mathcal{G}_n^{\delta_0}(W^\ell)$ .

In this way we write  $W^{\ell} = (\cup_j T^n U_j^{\ell}) \cup (\cup_i T^n V_i^{\ell})$ . Note that the images  $T^n V_i^{\ell}$  of the unmatched pieces must be short while the images of the matched pieces  $U_j^{\ell}$  may be long or short.

We have arranged a pairing of the pieces  $U_i^{\ell} = G_{U_i^{\ell}}(I_j), \ \ell = 1, 2$ , with the property:

(5.5) If 
$$U_j^1 = \{(r, \varphi_{U_i^1}(r)) \mid r \in I_j\}$$
 then  $U_j^2 = \{(r, \varphi_{U_i^2}(r)) \mid r \in I_j\}$ ,

so that the point  $x = (r, \varphi_{U_j^1}(r))$  is associated with the point  $\bar{x} = (r, \varphi_{U_j^2}(r))$  by the vertical segment  $\gamma_x \subset \{(r, s)\}_{s \in [-\pi/2, \pi/2]}$ , for each  $r \in I_j$ .

Given  $\psi_{\ell}$  on  $W^{\ell}$  with  $|\psi_{\ell}|_{C^{\alpha}(W^{\ell})} \leq 1$  and  $d(\psi_{1}, \psi_{2}) \leq \tilde{\varepsilon}$ , we must estimate

$$\left| \int_{W^{1}} \mathcal{L}_{g}^{n} f \psi_{1} dm_{W_{1}} - \int_{W^{2}} \mathcal{L}_{g}^{n} f \psi_{2} dm_{W_{2}} \right| \leq \sum_{l,i} \left| \int_{V_{i}^{l}} f \psi_{l} \circ T^{n} e^{S_{n}g} dm \right| 
+ \sum_{j} \left| \int_{U_{j}^{1}} f \psi_{1} \circ T^{n} e^{S_{n}g} dm - \int_{U_{j}^{2}} f \psi_{2} \circ T^{n} e^{S_{n}g} dm \right|.$$

We first estimate the differences of matched pieces  $U_j^l$ . The function  $\phi_j = (\psi_1 \circ T^n e^{S_n g}) \circ G_{U_j^1} \circ G_{U_j^2}^{-1}$  is well defined on  $U_j^2$ , and we can estimate each difference by

(5.7) 
$$\left| \int_{U_{j}^{1}} f \psi_{1} \circ T^{n} e^{S_{n}g} dm - \int_{U_{j}^{2}} f \psi_{2} \circ T^{n} e^{S_{n}g} dm \right| \leq \left| \int_{U_{j}^{1}} f \psi_{1} \circ T^{n} e^{S_{n}g} dm - \int_{U_{j}^{2}} f \phi_{j} dm \right| + \left| \int_{U_{j}^{2}} f (\phi_{j} - \psi_{2} \circ T^{n} e^{S_{n}g}) dm \right|.$$

We bound the first term in equation (5.7) using the strong unstable norm. We have that  $|G_{U_j^1} \circ G_{U_j^2}^{-1}|_{C^1} \leqslant C_g$ , for some  $C_g > 0$  due to the fact that each curve  $U_j^l$  has uniformly bounded curvature and slopes bounded away from infinity. Thus  $|\phi_j|_{C^{\alpha}(U_j^2)} \leqslant CC_g |\psi_1|_{C^{\alpha}(W^1)} |e^{S_n g}|_{C^{\alpha}(W^1)}$ . Moreover,  $d(\psi_1 \circ T^n e^{S_n g}, \phi_j) = |\psi_1 \circ T^n e^{S_n g} \circ G_{U_j^1} - \phi_j \circ G_{U_j^2}| = 0$  by definition of  $\phi_j$ . To complete the bound on the first term, we need the following estimate from [18, Lemma 4.2]: There exists C > 0, independent of  $W^1$  and  $W^2$ , such that

(5.8) 
$$d_{\mathcal{W}^s}(U_i^1, U_i^2) \leqslant C\Lambda^{-n} n\tilde{\varepsilon} =: \varepsilon_1, \quad \forall j.$$

Then we apply the definition of the strong unstable norm with  $\varepsilon_1$  instead of  $\tilde{\varepsilon}$ . Thus,

$$(5.9) \quad \sum_{j} \left| \int_{U_{j}^{1}} f \psi_{1} \circ T^{n} e^{S_{n}g} dm - \int_{U_{j}^{2}} f \phi_{j} dm \right| \leq 2\delta_{0}^{-1} C C_{g}^{2} |\log \varepsilon_{1}|^{-\zeta} ||f||_{u} \sum_{A \in \mathcal{M}_{0}^{n}} |e^{S_{n}g}|_{C^{0}(A)},$$

where we used Lemmas 3.12 and 3.1(b) with  $\gamma=0$  since there is at most one matched piece  $U_j^1$  corresponding to each component  $W_i^1\in\mathcal{G}_n^{\delta_0}(W^1)$  of  $T^{-n}W^1$ .

It remains to estimate the second term using the strong stable norm.

$$\left| \int_{U_j^2} f(\phi_j - \psi_2 \circ T^n e^{S_n g}) dm \right| \leq ||f||_s |\log |U_j^2||^{-\gamma} |\phi_j - \psi_2 \circ T^n e^{S_n g}|_{C^{\beta}(U_j^2)}.$$

In order to estimate this last  $C^{\beta}$ -norm, we use that  $|G_{U_i^2}|_{C^1} \leqslant C_g$  and  $|G_{U_i^2}^{-1}|_{C^1} \leqslant C_g$ .

$$\begin{aligned} |\phi_{j} - \psi_{2} \circ T^{n} e^{S_{n}g}|_{C^{\beta}(U_{j}^{2})} &\leqslant C|(\psi_{1} \circ T^{n} e^{S_{n}g}) \circ G_{U_{j}^{1}} - (\psi_{2} \circ T^{n} e^{S_{n}g}) \circ G_{U_{j}^{2}}|_{C^{\beta}(I_{j})} \\ &\leqslant C|(\psi_{1} \circ T^{n} \circ G_{U_{j}^{1}} - \psi_{2} \circ T^{n} \circ G_{U_{j}^{2}})(e^{S_{n}g} \circ G_{U_{j}^{1}}) \\ &+ (\psi_{2} \circ T^{n} \circ G_{U_{j}^{2}})(e^{S_{n}g} \circ G_{U_{j}^{1}} - e^{S_{n}g} \circ G_{U_{j}^{2}})|_{C^{\beta}(I_{j})} \\ &\leqslant C|(\psi_{1} \circ T^{n} \circ G_{U_{j}^{1}} - \psi_{2} \circ T^{n} \circ G_{U_{j}^{2}})|_{C^{\beta}(I_{j})}|e^{S_{n}g}|_{C^{0}(U_{j}^{1})} \\ &+ C|\psi_{2}|_{C^{0}(U_{j}^{2})}|e^{S_{n}g} \circ G_{U_{j}^{1}} - e^{S_{n}g} \circ G_{U_{j}^{2}}|_{C^{\beta}(I_{j})}. \end{aligned}$$

It follows from [18, Lemma 4.4] that

$$|\psi_1 \circ T^n \circ G_{U_i^1} - \psi_2 \circ T^n \circ G_{U_i^2}|_{C^{\beta}(I_j)} \leqslant C\tilde{\varepsilon}^{\alpha - \beta}.$$

Now, we need to estimate  $|e^{S_ng} \circ G_{U_j^1} - e^{S_ng} \circ G_{U_j^2}|_{C^{\beta}(I_j)}$ . Since  $d(T^i(G_{U_j^1}(r)), T^i(G_{U_j^2}(r))) \leq C\Lambda^{-(n-i)}\tilde{\varepsilon}$  for all  $r \in I_j$  and  $0 \leq i \leq n$ , we get

$$|e^{S_{n}g} \circ G_{U_{j}^{1}}(r) - e^{S_{n}g} \circ G_{U_{j}^{2}}(r)| = e^{S_{n}g(G_{U_{j}^{1}}(r))} |1 - e^{S_{n}g(G_{U_{j}^{2}}(r)) - S_{n}g(G_{U_{j}^{1}}(r))}|$$

$$\leq 2|e^{S_{n}g}|_{C^{0}(U_{j}^{1})}|S_{n}g(G_{U_{j}^{2}}(r)) - S_{n}g(G_{U_{j}^{1}}(r))|$$

$$\leq 2C \frac{\Lambda^{\alpha_{g}}}{\Lambda^{\alpha_{g}} - 1}|g|_{C^{\alpha_{g}}}(C\tilde{\varepsilon})^{\alpha_{g}}|e^{S_{n}g}|_{C^{0}(U_{j}^{1})}$$

We estimate the  $\beta$ -Hölder constant in two ways. First, using (5.11) twice, we have for all  $r, s \in I_j$  that

$$|e^{S_ng}\circ G_{U_j^1}(r)-e^{S_ng}\circ G_{U_j^2}(r)-e^{S_ng}\circ G_{U_j^1}(s)+e^{S_ng}\circ G_{U_j^2}(s)|\leqslant C\tilde{\varepsilon}^{\alpha_g}|e^{S_ng}|_{C^0(U_j^1)}.$$

On the other hand, using that  $G_{U_{\cdot}^{\ell}}(r)$  and  $G_{U_{\cdot}^{\ell}}(s)$  lie on the same stable curve,

$$\begin{split} |e^{S_ng} \circ G_{U^1_j}(r) - e^{S_ng} \circ G_{U^2_j}(r) - e^{S_ng} \circ G_{U^1_j}(s) + e^{S_ng} \circ G_{U^2_j}(s)| \\ & \leqslant |e^{S_ng} \circ G_{U^1_j}(r) - e^{S_ng} \circ G_{U^1_j}(s) + e^{S_ng} \circ G_{U^2_j}(s) - e^{S_ng} \circ G_{U^2_j}(r)| \\ & \leqslant |e^{S_ng}|_{C^{\alpha_g}(U^1_j)} d(G_{U^1_j}(r), G_{U^1_j}(s))^{\alpha_g} + |e^{S_ng}|_{C^{\alpha_g}(U^2_j)} d(G_{U^2_j}(r), G_{U^2_j}(s))^{\alpha_g} \\ & \leqslant C|e^{S_ng}|_{C^0(U^1_i)} |r-s|^{\alpha_g}. \end{split}$$

Thus, this quantity is bounded by the min of the two estimates. This min is maximal when the two upperbounds are equal, that is when  $\tilde{\varepsilon} = C|r-s|$ . Therefore, the  $\beta$ -Hölder constant satisfies

$$H_{I_j}^{\beta}(e^{S_n g} \circ G_{U_j^1} - e^{S_n g} \circ G_{U_j^2}) \leqslant C\tilde{\varepsilon}^{\alpha_g - \beta}|e^{S_n g}|_{C^0(U_j^1)}.$$

We therefore have proved that

$$|e^{S_n g} \circ G_{U_j^1} - e^{S_n g} \circ G_{U_j^2}|_{C^{\beta}(I_j)} \leqslant C \tilde{\varepsilon}^{\alpha_g - \beta} |e^{S_n g}|_{C^0(U_j^1)} \,.$$

Combining the above estimates inside (5.10), we finally have

$$|\phi_j - \psi_2 \circ T^n e^{S_n g}|_{C^{\beta}(U_j^2)} \le C \tilde{\varepsilon}^{\alpha - \beta} |e^{S_n g}|_{C^0(U_j^1)}.$$

Summing over j yields

$$\left| \sum_{j} \left| \int_{U_j^2} f(\phi_j - \psi_2 \circ T^n e^{S_n g}) dm \right| \leqslant C |\log \delta_0|^{-\gamma} ||f||_s \tilde{\varepsilon}^{\alpha - \beta} 2\delta_0^{-1} \sum_{A \in \mathcal{M}_0^n} |e^{S_n g}|_{C^0(A)},$$

where we used Lemma 3.1(b) with  $\gamma = 0$  since there is at most one matched piece  $U_j^l$  corresponding to each component  $W_i^l \in \mathcal{G}_n^{\delta_0}(W^l)$  of  $T^{-n}W^l$ . Since  $\delta_0 < 1$  is fixed, this completes the estimate on the second term of the matched pieces (originating from (5.6)).

We now turn to the estimate of the first sum in (5.6) concerning the unmatched pieces.

We say an unmatched curve  $V_i^1$  is created at time j,  $1 \le j \le n$ , if j is the first time that  $T^{n-j}V_i^1$  is not part of a matched element of  $\mathcal{G}_j^{\delta_0}(W^1)$ . Indeed, there may be several curves  $V_i^1$  (in principle exponentially many in n-j) such that  $T^{n-j}V_i^1$  belongs to the same unmatched element of  $\mathcal{G}_j^{\delta_0}(W^1)$ . Define

$$A_{j,k} = \{i \mid V_i^1 \text{ is created at time } j$$
 and  $T^{n-j}V_i^1$  belongs to the unmatched curve  $W_k^1 \subset T^{-j}W^1\}$ .

Due to the uniform hyperbolicity of T, and, again, uniform transversality of  $\mathcal{S}_{-n}$  with the stable cone and of  $C^s(x)$  with  $C^u(x)$ , we have  $|W_k^1| \leq C\Lambda^{-j}\tilde{\varepsilon}$ .

Recall that from Lemma 3.1(a) for  $\bar{\gamma} = 0$ , if for a certain time q, every element of  $\mathcal{G}_q^{\delta_0}(W_k^1)$  have length less than  $\delta_0/3$  – that is, if  $\mathcal{G}_q^{\delta_0}(W_k^1) = \mathcal{I}_q^{\delta_0}(W_k^1)$  – then we have the subexponential growth

(5.12) 
$$\sum_{V \in \mathcal{G}_q^{\delta_0}(W_k^1)} |e^{S_q g}|_{C^0(V)} \leq 2(Km+1)^{q/m} e^{q \sup g}.$$

We would like to establish a lower bound on the value of q as a function of j.

More precisely, we want to find q(j), as large as possible, so that

(a)  $\mathcal{G}_{q(j)}^{\delta_0}(W_k^1) = \mathcal{I}_{q(j)}^{\delta_0}(W_k^1);$ 

(b) 
$$\frac{|\log |V||^{-\gamma}}{|\log \tilde{\varepsilon}|^{-\varsigma}} \le 1, \text{ for all } V \in \mathcal{I}_{q(j)}^{\delta_0}(W_k^1).$$

This is the content of the next two lemmas.

**Lemma 5.4.** If  $W \in \widehat{\mathcal{W}}^s$  is such that  $\widetilde{C}^2|W|^{2^{-kn_0s_0}} < \delta_0/3$  for some  $k \geqslant 1$ , where  $\widetilde{C}$  is the constant from (3.2). Then  $\mathcal{G}_{kn_0}^{\delta_0}(W) = \mathcal{I}_{kn_0}^{\delta_0}(W)$ , and for all  $1 \leqslant l \leqslant k$ , and all  $W_i \in \mathcal{G}_{ln_0}^{\delta_0}(W)$ ,  $|W_i| \leqslant \widetilde{C}^2|W|^{2^{-ln_0s_0}}$ .

Proof. We prove the lemma by induction on k. We start with the case k=1. Let  $1\leqslant l\leqslant n_0$  and  $W_i\in\mathcal{G}_l^{\delta_0}(W)$ . Denote  $V=T^lW_i\subset W$ . Then, for all  $0\leqslant j\leqslant l, |T^jW_i|\leqslant \delta_0$ . Decomposing  $T^{-l}V=W_i$  as in the beginning of the proof of Lemma 3.1, we get that  $|W_i|\leqslant \tilde{C}|W|^{2^{-n_0s_0}}$ , which is less than  $\delta_0/3$  by assumption. Thus,  $\mathcal{G}_l^{\delta_0}(W)=S_l^{\delta_0}(W)$  for each  $0\leqslant l\leqslant n_0$ . Therefore  $\mathcal{G}_{n_0}^{\delta_0}(W)=\mathcal{I}_{n_0}^{\delta_0}(W)$ , with the claimed estimate.

Consider now the case k > 1. Notice that, by construction, we have

$$\mathcal{G}_{(k+1)n_0}^{\delta_0}(W) = \bigcup_{W_i \in \mathcal{G}_{kn_0}^{\delta_0}(W)} \mathcal{G}_{n_0}^{\delta_0}(W).$$

Thus, we can apply the same method to estimate the length of an element  $W_i \in \mathcal{G}_{kn_0+l}^{\delta_0}(W)$  from the length of its *parent* in  $\mathcal{G}_{kn_0}^{\delta_0}(W)$ , iterating the estimates in the same fashion as for (3.2).

**Lemma 5.5.** The above conditions (a) and (b) are satisfied for  $q(j) := \frac{(\gamma - \zeta) \log(j - j_0)}{\gamma s_0 \log 2} - 1$ , for all  $j \ge j_1$ , where  $j_1 > j_0 \ge 0$  are constants (uniform in  $\tilde{\varepsilon}$  and  $W_k^1$ ). For  $j < j_1$ , set q(j) = 0.

*Proof.* Since  $|W_k^1| \le C\tilde{\varepsilon}\Lambda^{-j}$  and using Lemma 5.4, the condition (a) will be satisfied whenever  $\tilde{C}^2(C\tilde{\varepsilon}\Lambda^{-j})^{2^{-qs_0}} \le \delta_0/3$ .

Let  $j_0$  be such that  $C\Lambda^{-j_0} < 1$ . Then (a) is satisfied whenever  $\tilde{C}^2\Lambda^{-(j-j_0)2^{-q}} \leqslant \delta_0/3$ , that is

(5.13) 
$$q \leqslant \frac{\log(j-j_0)}{s_0 \log 2} - C_2, \text{ with } C_2 := \frac{1}{s_0 \log 2} \log \frac{\log \frac{3\tilde{C}^2}{\delta_0}}{\log \Lambda}.$$

Note that  $C_2$  is uniform, and that the right-hand-side of (5.13) is larger than q(j) for all j large enough, say  $j \ge j_1$ .

Using the estimate from Lemma 5.4, condition (b) is satisfied whenever  $|\log \tilde{C}^2 (C\tilde{\epsilon}\Lambda^{-j})^{2^{-q}}|^{\gamma} > |\log \tilde{\epsilon}|^{\zeta}$ . Now, we have that

$$|\log \tilde{C}^2 (C\tilde{\varepsilon}\Lambda^{-j})^{2^{-q}}| = |\log \tilde{C}^2 + 2^{-q}\log(C\tilde{\varepsilon}\Lambda^{-j})| > \frac{1}{2}|2^{-q}\log(C\tilde{\varepsilon}\Lambda^{-j})|,$$

whenever

(5.14) 
$$q + 1 \leqslant \frac{\log(j - j_0)}{s_0 \log 2} + C_3, \text{ with } C_3 = \frac{1}{s_0 \log 2} \log \frac{\log \Lambda}{\log \tilde{C}^2}$$

Note that  $C_3$  is uniform, and that the right-hand-side of (5.14) is larger than q(j) for all j large enough, say  $j \ge j_1$  (up to increasing the value of  $j_1$ ).

We thus have to prove that  $|\log C\tilde{\varepsilon}\Lambda^{-j}|^{\gamma} > 2^{(q+1)\gamma}|\log \tilde{\varepsilon}|^{\zeta}$  (which implies (b)). Notice that, from the definition of q(j), we have  $2^{(q(j)+1)\gamma} \leq (j-j_0)^{\gamma-\zeta}$ . We distinguish two cases.

Assume first that  $(j - j_0) \log \Lambda \ge |\log \tilde{\varepsilon}|$ . Therefore

$$2^{(q(j)+1)\gamma} |\log \tilde{\varepsilon}|^{\zeta} \leqslant (j-j_0)^{\gamma-\zeta} |\log \tilde{\varepsilon}|^{\zeta} \leqslant (j-j_0)^{\gamma} (\log \Lambda)^{\zeta} \leqslant ((j-j_0)\log \Lambda)^{\gamma}$$

$$\leqslant ((j-j_0)\log \Lambda + |\log \tilde{\varepsilon}| + |\log C\Lambda^{-j_0}|)^{\gamma}$$

$$\leqslant |-(j-j_0)\log \Lambda + \log \tilde{\varepsilon} + \log C\Lambda^{-j_0}|^{\gamma}$$

$$\leqslant |\log C\tilde{\varepsilon}\Lambda^{-j}|^{\gamma}.$$

On the other hand, if  $(j - j_0) \log \Lambda \leq |\log \tilde{\varepsilon}|$ , then

$$2^{(q(j)+1)\gamma} |\log \tilde{\varepsilon}|^{\zeta} \leqslant (j-j_0)^{\gamma-\zeta} |\log \tilde{\varepsilon}|^{\zeta} \leqslant \frac{|\log \tilde{\varepsilon}|^{\gamma-\zeta}}{(\log \Lambda)^{\gamma-\zeta}} |\log \tilde{\varepsilon}|^{\zeta} \leqslant |\log \tilde{\varepsilon}|^{\gamma}$$

$$\leqslant ((j-j_0) \log \Lambda + |\log \tilde{\varepsilon}| + |\log C\Lambda^{-j_0}|)^{\gamma}$$

$$\leqslant |-(j-j_0) \log \Lambda + \log \tilde{\varepsilon} + \log C\Lambda^{-j_0}|^{\gamma}$$

$$\leqslant |\log C\tilde{\varepsilon}\Lambda^{-j}|^{\gamma}.$$

Thus, the choice q(j) satisfies (a) and (b) for all  $j \ge j_1$ .

We next estimate<sup>15</sup> over the unmatched pieces  $V_i^l$  in (5.6), using the strong stable norm. Since cases l=1 and l=2 are similar here, we only deal with the case l=1.

$$\begin{split} &\sum_{V_i^1} \left| \int_{V_i^1} f \psi_1 \circ T^n e^{S_n g} \mathrm{d} m_{V_i^1} \right| = \sum_{j=1}^n \sum_k \sum_{i \in A_{j,k}} \left| \int_{T^{n-j}V_i^1} (\mathcal{L}_g^{n-j} f) \psi_1 \circ T^j e^{S_j g} \right| \\ &\leqslant \sum_{j=1}^n \sum_k \sum_{V_l \in \mathcal{G}_{q(j)}^{\delta_0}(W_k^1)} \left| \int_{V_l} (\mathcal{L}_g^{n-j-q(j)} f) \psi_1 \circ T^{j+q(j)} e^{S_{j+q(j)}g} \right| \\ &\leqslant \sum_{j=1}^n \sum_k \sum_{V_l \in \mathcal{G}_{q(j)}^{\delta_0}(W_k^1)} \left| |\mathcal{L}_g^{n-j-q(j)} f| |_s C |\log |V_l||^{-\gamma} |\psi_1 \circ T^{j+q(j)}|_{C^{\beta}(V_l)} |e^{S_{j+q(j)}g}|_{C^{\beta}(V_l)} \right| \\ &\leqslant C ||f||_s \sum_{j=1}^n \frac{C}{c_1 \delta_0} e^{(n-j-q(j))P_*(T,g)} |\log \tilde{\varepsilon}|^{-\zeta} \sum_k \sum_{V_l \in \mathcal{G}_{q(j)}^{\delta_0}(W_k^1)} |e^{S_{j+q(j)}g}|_{C^{\beta}(V_l)} \\ &\leqslant \frac{C}{c_1 \delta_0} ||f||_s \sum_{j=1}^n e^{(n-j-q(j))P_*(T,g)} |\log \tilde{\varepsilon}|^{-\zeta} \sum_{W_k^1 \subset T^{-j}W^1} \sum_{V_l \in \mathcal{G}_{q(j)}^{\delta_0}(W_k^1)} |e^{S_j g \circ T^{q(j)} + S_{q(j)}g}|_{C^{\beta}(V_l)} \\ &\leqslant \frac{C}{c_1 \delta_0} ||f||_s \sum_{j=1}^n e^{(n-j-q(j))P_*(T,g)} |\log \tilde{\varepsilon}|^{-\zeta} \sum_{W_k^1 \subset T^{-j}W^1} |e^{S_j g}|_{C^0(W_k^1)} \sum_{V_l \in \mathcal{G}_{q(j)}^{\delta_0}(W_k^1)} |e^{S_{q(j)}g}|_{C^0(V_l)} \end{aligned}$$

<sup>&</sup>lt;sup>15</sup>For the 4<sup>th</sup> and 6<sup>th</sup> inequalities, we use Proposition 3.10. Here again,  $P_*(T,g)$  can be replaced by  $P_*(T,g)+\varepsilon$  up to a larger multiplicative constant.

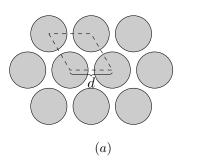
$$\leqslant \frac{C}{c_1 \delta_0} ||f||_s \sum_{j=1}^n e^{(n-j-q(j))P_*(T,g)} |\log \tilde{\varepsilon}|^{-\zeta} \frac{4C}{c_1 \delta_0} e^{jP_*(T,g)} e^{q(j)\sup g} (Km+1)^{q(j)/m} 
\leqslant \frac{C}{(c_1 \delta_0)^2} ||f||_s |\log \tilde{\varepsilon}|^{-\zeta} e^{nP_*(T,g)} \sum_{j=1}^n e^{-q(j)(P_*(T,g)-\sup g - \frac{1}{m}\log(Km+1))}.$$

Now, for  $\tilde{\varepsilon} > 0$ , fixed, since we assume that  $P_*(T,g) - \sup g > s_0 \log 2$ , we can chose m large enough and  $\zeta$  small enough such that  $\varepsilon_1 := P_*(T,g) - \sup g - \frac{1}{m} \log(Km+1) - \frac{\gamma}{\gamma-\zeta} s_0 \log 2 > 0$ . By definition of g(j), we obtain that

$$\sum_{j=j_1}^n e^{-q(j)(P_*(T,g)-\sup g-\frac{1}{m}\log(Km+1))} = \sum_{j=j_1}^n e^{-\frac{(\gamma-\zeta)\log(j-j_0)}{\gamma s_0\log 2}(\varepsilon_1 + \frac{\gamma}{\gamma-\zeta}s_0\log 2)} = \sum_{j=j_1}^n (j-j_0)^{-1-\frac{\gamma-\zeta}{\gamma s_0\log 2}\varepsilon_1},$$

is bounded. The bound (5.3) then follows by combining all the above estimates into (5.6) and taking the appropriate suprema.

Remark 5.6. In the case  $g = -h_{\rm top}(\phi_1)\tau$ , the assumption  $P_*(T,g) - \sup g > s_0 \log 2$  in Proposition 5.1 is implied by the condition  $h_{\rm top}(\phi_1)\tau_{\rm min} > s_0 \log 2$ , which is itself implied by  $\tau_{\rm min}h_{\mu_{\rm SRB}}(T)/\mu_{\rm SRB}(\tau) > s_0 \log 2$  thanks to the Abramov formula. This latter condition appears to be satisfied for billiards studied by Baras and Gaspard [22] and by Garrido [21], as long as  $\tau_{\rm min}$  is not too small.



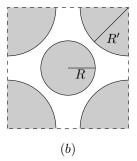


FIGURE 1. (a) The Sinai billiard on a triangular lattice studied in [22] with angle  $\pi/3$ , scatterer of radius 1, and distance d between the centers of adjacent scatterers. (b) The Sinai billiard on a square lattice with scatterers of radius R < R' studied in [21]. The boundary of a single cell is indicated by dashed lines in both tables.

Indeed, Garrido [21] studied the Sinai billiard corresponding to the periodic Lorentz gas with two scatterers of radius R < R' on the unit square lattice (Figure 1(b)). Setting R' = 0.4, Garrido computed  $h_{\mu_{\text{SRB}}}(T)$  and  $\mu_{\text{SRB}}(\tau)$  for about 20 values of R ranging from R = 0.1 (when the horizon becomes infinite) to  $R = \frac{\sqrt{2}}{2} - 0.4$  (when the scatterers touch:  $\tau_{\text{min}} = 0$ ). According to  $[3, \S 2.4]$ , in those examples we can always find  $\varphi_0$  and  $n_0$  such that  $s_0 \leqslant \frac{1}{2}$ . Furthermore,  $\tau_{\text{min}} = \frac{\sqrt{2}}{2} - 0.4 - R$ . Now, for  $R = 0.1^+$ , we find that

$$\tau_{\min} h_{\mu_{\text{SRB}}}(T)/\mu_{\text{SRB}}(\tau) \geqslant (\frac{\sqrt{2}}{2} - 0.5) \frac{1.7}{0.5} \geqslant 0.7 > \frac{1}{2} \log 2 \geqslant s_0 \log 2$$

and for R = 0.2, we find that

$$\tau_{\min} h_{\mu_{\text{SRB}}}(T) / \mu_{\text{SRB}}(\tau) \geqslant (\frac{\sqrt{2}}{2} - 0.6) \frac{1.4}{0.3} \geqslant 0.48 > \frac{1}{2} \log 2 \geqslant s_0 \log 2$$
.

Since for  $R \in (0.1, 0.2]$ ,  $R \mapsto \tau_{\min}(R)$  is a linear function, and according to Garrido Figures 6 and 8,  $R \mapsto \mu_{\text{SRB}}(\tau)(R)$  is well approximated by an affine function and  $R \mapsto h_{\mu_{\text{SRB}}}(T)(R)$  is lower bounded by an affine function joining the values at R = 0.1 and 0.2, it appears that the condition  $\tau_{\min}h_{\mu_{\text{SRB}}}(T)/\mu_{\text{SRB}}(\tau) > s_0 \log 2$  is satisfied for all  $R \in (0.1, 0.2]$ .

Baras and Gaspard studied the Sinai billiard corresponding to the Lorentz gas with disks of radius 1 centered in a triangular lattice (Figure 1(a)). The distance d between points on the lattice is varied from d=2 (when the scatterers touch:  $\tau_{\min}=0$ ) to  $d=4/\sqrt{3}$  (when the horizon becomes infinite). We have that  $\tau_{\min}=d-2$  and, still according to [3, § 2.4], in those examples we can always find  $\varphi_0$  and  $n_0$  such that  $s_0 \leq \frac{1}{2}$ . The computed values are the average Lyapunov exponent of the billiard flows given in [22], provide a lower bound directly on  $h_{\mu_{\text{SRB}}}(T)/\mu_{\text{SRB}}(\tau)$ . For d=0.2, we find

$$\tau_{\min} h_{\mu_{SRB}}(T)/\mu_{SRB}(\tau) \geqslant (\frac{4}{\sqrt{3}} - 2) \, 1.8 \geqslant 0.55 > \frac{1}{2} \log 2 \geqslant s_0 \log 2$$
.

The condition  $h_{\text{top}}(\phi_1)\tau_{\text{min}} > s_0 \log 2$  is a little bit more restrictive than the one used by Baladi and Demers in [3] since, by the Abramov formula,  $h_* = h_{\text{top}}(\phi_1)\mu_*(\tau) \ge h_{\text{top}}(\phi_1)\tau_{\text{min}}$ . (Also, we do not know any example of billiard for which the condition  $h_* > s_0 \log 2$  is not satisfied.)

We now turn to the condition SSP.1. Unfortunately, we don't know any billiard table such that the potential  $g = -h_{\rm top}(\phi_1)\tau$  satisfies a simple condition implying SSP.1. By simple, we mean a sufficient condition that does not involve topological entropies, since they are notoriously hard to estimate numerically. First, recall from Lemmas 3.3 and 3.4 that  $\log \Lambda > h_{\rm top}(\phi_1)(\tau_{\rm max} - \tau_{\rm min})$  implies SSP.1. Remark that since g and  $\frac{1}{n}S_ng$  are cohomologous, they would give rise to the same equilibrium states. It is then advantageous to work with the Birkhoff average instead of g because  $\max \frac{1}{n}S_ng \leqslant \tau_{\rm max}$  and  $\min \frac{1}{n}S_ng = \tau_{\rm min}$  (notice that  $\tau_{\rm min}$  is achieved on an orbit of period 2). Now, taking advantage of the Abramov formula and of the variational principle, we get that  $\max \frac{1}{n}S_ng < 2\tau_{\rm min}$  implies  $h_* > h_{\rm top}(\phi_1)(\max \frac{1}{n}S_ng - \tau_{\rm min})$  (recall that  $h_* > \log \Lambda$  is the topological entropy of T, as defined in [3]). The condition  $\max \frac{1}{n}S_ng < 2\tau_{\rm min}$  involves quantities that are easy to estimate numerically, however, we don't know any billiard table satisfying this condition.

We now deduce the bounds of Theorem 5.3 from the rate of growth of stable curves proved in Proposition 3.7.

Proof of Theorem 5.3. To prove this lower bound on  $|\mathcal{L}_g^n 1|_w$ , recall the choice of  $\delta_1 > 0$  from Lemma 3.3 for  $\varepsilon = 1/4$ . Let  $W \in \mathcal{W}^s$  with  $|W| \ge \delta_1/3$  and set the test function  $\psi \equiv 1$ . For  $n \ge n_1$ ,

$$\int_W \mathcal{L}_g^n 1 \mathrm{d} m_W = \sum_{W_i \in \mathcal{G}_n^{\delta_1}(W)} \int_{W_i} e^{S_n g} \mathrm{d} m_{W_i} \geqslant \sum_{W_i \in \mathcal{G}_n^{\delta_1}(W)} \frac{\delta_1}{2} \inf_{W_i} e^{S_n g} \geqslant \frac{\delta_1}{2} C^{-1} \sum_{W_i \in \mathcal{G}_n^{\delta_1}(W)} \sup_{W_i} e^{S_n g},$$

where we used Lemma 2.3 for the second inequality, since for each  $W_i \in \mathcal{G}_n^{\delta_1}(W)$  there exists  $A \in \mathcal{M}_0^n$  such that  $W_i \subset A$  and

$$\sup_{W_i} e^{S_n g} \leqslant \sup_{A} e^{S_n g} \leqslant C \inf_{A} e^{S_n g} \leqslant C \inf_{W_i} e^{S_n g}.$$

We can now use Proposition 3.7 to get

(5.15) 
$$\int_{W} \mathcal{L}_{g}^{n} 1 dm_{W} \geqslant \frac{\delta_{1}}{2C} c_{0} \sum_{A \in \mathcal{M}_{-n}^{0}} |e^{S_{n}^{-1}g}|_{C^{0}(A)} \geqslant \frac{\delta_{1}}{2C} c_{0} e^{nP_{*}(T,g)}.$$

Thus

$$||\mathcal{L}_{g}^{n}1||_{s} \geqslant |\mathcal{L}_{g}^{n}1|_{w} \geqslant \frac{\delta_{1}}{2}c_{0}e^{nP_{*}(T,g)}.$$

Letting n tend to infinity, one obtains  $\lim_{n\to\infty} ||\mathcal{L}_g^n \mathbf{1}||_{\mathcal{B}}^{1/n} \geqslant e^{P_*(T,g)}$ .

# 6. The measure $\mu_a$

This section is devoted to the construction, the properties and the uniqueness of an equilibrium state  $\mu_q$  for T, associated to a potential g.

We will assume throughout that g is a  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential such that  $P_*(T, g)$ -sup  $g > s_0 \log 2$  and that the conditions SSP.1 and SSP.2 are satisfied.

6.1. Construction of the measure  $\mu_g$  – Measure of Singular Sets. In this section, we construct a T-invariant probability measure  $\mu_g$  on M by combining in (6.1) a maximal eigenvector of  $\mathcal{L}_g$  on  $\mathcal{B}$  and a maximal eigenvector of its dual, obtained in Proposition 6.1. In addition, the information on these left and right eigenvectors will give Lemma 6.2 and Corollary 6.3, which imply that  $\mu_g$  is T-adapted.

We first show that such maximal eigenvectors exist and are in fact nonnegative Radon measures – that is, elements of the dual of  $C^0(M)$ .

**Proposition 6.1.** If g is a  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder continuous potential such that  $P_*(T, g) - \sup g > s_0 \log 2$  and  $\log \Lambda > \sup g - \inf g$ , then there exist  $\nu \in \mathcal{B}_w$  and  $\tilde{\nu} \in \mathcal{B}_w^*$  such that  $\mathcal{L}_g \nu = e^{P_*(T,g)} \nu$  and  $\mathcal{L}_g^* \tilde{\nu} = e^{P_*(T,g)} \tilde{\nu}$ . In addition,  $\nu$  and  $\tilde{\nu}$  take nonnegative values on nonnegative  $C^1$  functions on M and are thus nonnegative Radon measures. Finally,  $\tilde{\nu}(\nu) \neq 0$  and  $||\nu||_u \leq \overline{C}$ .

It is easy to see that  $|f\varphi|_w \leq |\varphi|_{C^1}|f|_w$  (use  $|\varphi\psi|_{C^{\alpha}(W)} \leq |\varphi|_{C^1}|\psi|_{C^{\alpha}(W)}$ ). Clearly, if  $f \in C^1$  and  $\varphi \in C^1$  then  $f\varphi \in C^1$ . Therefore, if  $P_*(T,g) - \sup g > s_0 \log 2$  and both SSP conditions are satisfied, a bounded linear map  $\mu_g$  from  $C^1(M)$  to  $\mathbb C$  can be defined by taking  $\nu$  and  $\tilde{\nu}$  from Proposition 6.1 and setting

(6.1) 
$$\mu_g(\varphi) = \frac{\tilde{\nu}(\varphi\nu)}{\tilde{\nu}(\nu)}.$$

This map is nonnegative for all nonnegative  $\varphi$  and thus defines a nonnegative measure  $\mu_g \in (C^0(M))^*$ , with  $\mu_g(1) = 1$ . Clearly,  $\mu_g$  is a T invariant probability measure since for every  $\varphi \in C^1$  we have

$$\tilde{\nu}(\varphi\nu) = e^{-P_*(T,g)}\tilde{\nu}(\varphi\mathcal{L}_g(\nu)) = e^{-P_*(T,g)}\tilde{\nu}(\mathcal{L}_g((\varphi\circ T)\nu)) = \tilde{\nu}((\varphi\circ T)\nu) = \tilde{\nu}(\nu)\mu_g(\varphi\circ T).$$

*Proof.* Let 1 denote the constant function equal to 1 on M. We will take this as a seed in our construction of a maximal eigenvector. By Theorem 5.3, we see that  $\|\mathcal{L}_g^n 1\|_{\mathcal{B}} \ge \|\mathcal{L}_g^n 1\|_{s} \ge \|\mathcal{L}_g^n 1\|_{w} \ge Ce^{nP_*(T,g)}$ . Now consider

(6.2) 
$$\nu_n := \frac{1}{n} \sum_{k=0}^{n-1} e^{-kP_*(T,g)} \mathcal{L}_g^k 1, \quad n \geqslant 1.$$

By construction, the  $\nu_n$  are nonnegative, and thus Radon measures. By Proposition 5.1, they satisfy  $||\nu_n||_{\mathcal{B}} \leq \overline{C}$ , so using the relative compactness of  $\mathcal{B}$  in  $\mathcal{B}_w$  ([3, Proposition 6.1]), we extract a subsequence  $(n_j)$  such that  $\lim_j \nu_{n_j} = \nu$  is a nonnegative Radon measure, and the convergence is in  $\mathcal{B}_w$ . Since  $\mathcal{L}_g$  is continuous on  $\mathcal{B}_w$ , we may write,

$$\mathcal{L}_{g}\nu = \lim_{j \to \infty} \frac{1}{n_{j}} \sum_{k=0}^{n_{j}-1} e^{-kP_{*}(T,g)} \mathcal{L}_{g}^{k+1} 1$$

$$= \lim_{j \to \infty} \frac{e^{P_{*}(T,g)}}{n_{j}} \sum_{k=0}^{n_{j}-1} e^{-kP_{*}(T,g)} \mathcal{L}_{g}^{k} 1 - \frac{1}{n_{j}} e^{P_{*}(T,g)} 1 + \frac{1}{n_{j}} e^{(n_{j}-1)P_{*}(T,g)} \mathcal{L}_{g}^{n_{j}} 1$$

$$= e^{P_{*}(T,g)} \nu.$$

where we used that the second and third terms go to 0 (in the  $\mathcal{B}$ -norm). We thus obtain a nonnegative measure  $\nu \in \mathcal{B}_w$  such that  $\mathcal{L}_g \nu = e^{P_*(T,g)} \nu$ .

Although  $\nu$  is not a priori an element of  $\mathcal{B}$ , it does inherit bounds on the unstable norm from the sequence  $\nu_n$ . The convergence of  $(\nu_{n_i})$  to  $\nu$  in  $\mathcal{B}_w$  implies that

$$\lim_{j \to \infty} \sup_{W \in \mathcal{W}^s} \sup_{\substack{\psi \in C^{\alpha}(W) \\ |\psi|_{C^{\alpha}(W)} \leq 1}} \left( \int_W \nu \psi \, \mathrm{d} m_W - \int_W \nu_{n_j} \psi \, \mathrm{d} m_W \right) = 0.$$

Since  $||\nu_{n_i}||_u \leq \overline{C}$ , it follows that  $||\nu||_u \leq \overline{C}$ , as claimed.

Next, recalling the bound  $|\int f d\mu_{SRB}| \leq \hat{C}|f|_w$  from [3, Proposition 4.2], setting  $d\mu_{SRB} \in (\mathcal{B}_w)^*$  to be the functional defined on  $C^1(M) \subset \mathcal{B}_w$  by  $d\mu_{SRB}(f) = \int f d\mu_{SRB}$  and extended by density, we define

(6.3) 
$$\tilde{\nu}_n := \frac{1}{n} \sum_{k=0}^{n-1} e^{-kP_*(T,g)} (\mathcal{L}_g^*)^k (\mathrm{d}\mu_{\mathrm{SRB}}).$$

Then, we have  $|\tilde{\nu}_n(f)| \leq C|f|_w$  for all n and all  $f \in \mathcal{B}_w$ . So  $\tilde{\nu}_n$  is bounded in  $(\mathcal{B}_w)^* \subset \mathcal{B}^*$ . By compactness of the embedding ([3, Proposition 6.1]), we can find a subsequence  $\tilde{\nu}_{\tilde{n}_j}$  converging to  $\tilde{\nu} \in \mathcal{B}^*$ . By the argument above, we have  $\mathcal{L}_q^* \tilde{\nu} = e^{P_*(T,g)} \tilde{\nu}$ .

We next check that  $\tilde{\nu}$ , which in principle lies in the dual of  $\mathcal{B}$ , is in fact an element of  $(\mathcal{B}_w)^*$ . For this, it suffices to find  $\tilde{C} < \infty$  so that for any  $f \in \mathcal{B}$  we have

(6.4) 
$$\tilde{\nu}(f) \leqslant \tilde{C}|f|_w.$$

Now, for  $f \in \mathcal{B}$  and any  $n_i \ge 1$ , we have

$$|\tilde{\nu}(f)| \leq |(\tilde{\nu}_{n_j} - \tilde{\nu})(f)| + |\tilde{\nu}_{n_j}(f)| \leq |(\tilde{\nu}_{n_j} - \tilde{\nu})(f)| + |f|_w.$$

Since  $\tilde{\nu}_{\tilde{n}_j} \to \tilde{\nu}$  in  $\mathcal{B}^*$ , we conclude  $|\tilde{\nu}(f)| \leq |f|_w$  for all  $f \in \mathcal{B}$ . Since  $\mathcal{B}$  is dense in  $\mathcal{B}_w$ , by [29, Thm I.7]  $\tilde{\nu}$  extends uniquely to a bounded linear functional on  $\mathcal{B}_w$  satisfying (6.4). It only remains to prove that  $\tilde{\nu}(\nu) > 0$ .

Let  $(n_j)$  (resp.  $(\tilde{n}_j)$ ) denote the subsequence such that  $\nu = \lim_j \nu_{n_j}$  (resp.  $\tilde{\nu} = \lim_j \tilde{\nu}_{\tilde{n}_j}$ ). Since  $\tilde{\nu}$  is continuous on  $\mathcal{B}_w$ , we have on the one hand

$$\tilde{\nu}(\nu) = \lim_{j \to \infty} \tilde{\nu}(\nu_{n_j}) = \lim_{j \to \infty} \frac{1}{n_j} \sum_{k=0}^{n_j - 1} e^{-kP_*(T,g)} \tilde{\nu}(\mathcal{L}_g^k 1) = \lim_{j \to \infty} \frac{1}{n_j} \sum_{k=0}^{n_j - 1} \tilde{\nu}(1) = \tilde{\nu}(1),$$

where we have used that  $\tilde{\nu}$  is an eigenvector of  $\mathcal{L}_q^*$ . On the other hand,

$$\tilde{\nu}(1) = \lim_{j \to \infty} \frac{1}{\tilde{n}_j} \sum_{k=0}^{\tilde{n}_j - 1} e^{-kP_*(T,g)} (\mathcal{L}_g^*)^k d\mu_{\text{SRB}}(1) = \lim_{j \to \infty} \frac{1}{\tilde{n}_j} \sum_{k=0}^{\tilde{n}_j - 1} e^{-kP_*(T,g)} \int \mathcal{L}_g^k 1 d\mu_{\text{SRB}}.$$

Next, we disintegrate  $d\mu_{SRB}$  as in the proof of [3, Lemma 4.4] into conditional measure  $\mu_{SRB}^{W_{\xi}}$  on maximal homogeneous stable manifolds  $W_{\xi} \in \mathcal{W}_{\mathbb{H}}^{s}$  and a factor measure  $d\hat{\mu}_{SRB}(\xi)$  on the index set  $\Xi$  of stable manifolds. Recall that  $\mu_{SRB}^{W_{\xi}} = |W_{\xi}|^{-1} \rho_{\xi} dm_{W}$ , where  $\rho_{\xi}$  is uniformly log-Hölder continuous so that

$$0 < c_{\rho} \leqslant \inf_{\xi \in \Xi} \inf_{W_{\xi}} \rho_{\xi} \leqslant \sup_{\xi \in \Xi} |\rho_{\xi}|_{C^{\alpha}(W_{\xi})} \leqslant C_{\rho} < \infty.$$

Let  $\Xi^{\delta_1}$  denote those  $\xi \in \Xi$  such that  $|W_{\xi}| \geq \delta_1/3$  and note that  $\hat{\mu}_{SRB}(\Xi^{\delta_1}) > 0$ . Then, disintegrating as usual, we get by (5.15) for  $k \geq n_1$ ,

$$\int \mathcal{L}_g^k 1 d\mu_{\text{SRB}} = \int_{\Xi} \int_{W_{\xi}} \mathcal{L}_g^k 1 \rho_{\xi} |W_{\xi}|^{-1} dm_{W_{\xi}} d\hat{\mu}_{\text{SRB}}(\xi)$$

$$\geqslant \int_{\Xi^{\delta_1}} \int_{W_{\xi}} \mathcal{L}_g^k 1 dm_{W_{\xi}} c_{\rho} 3\delta_1^{-1} d\hat{\mu}_{\text{SRB}}(\xi) \geqslant c_{\rho} \frac{2c_0}{3} e^{kP_*(T,g)} \hat{\mu}_{\text{SRB}}(\Xi^{\delta_1}) > 0.$$

Thus  $\tilde{\nu}(\nu) = \tilde{\nu}(1) \geqslant c_{\rho} \frac{2c_0}{3} \hat{\mu}_{SRB}(\Xi^{\delta_1}) > 0$  as required.

**Lemma 6.2.** For any  $\gamma > 0$  such that  $2^{s_0\gamma} < e^{P_*(T,g)-\sup g}$  and any  $k \in \mathbb{Z}$  there exists  $C_k > 0$  such that

(6.5) 
$$\mu_g(\mathcal{N}_{\varepsilon}(\mathcal{S}_k)) \leqslant C_k |\log \varepsilon|^{-\gamma}, \quad \forall \varepsilon > 0.$$

In particular, for any  $p > 1/\gamma$  (one can choose p < 1 for  $\gamma > 1$ ),  $\eta > 0$ , and  $k \in \mathbb{Z}$ , for  $\mu_g$ -almost every  $x \in M$ , there exists C > 0 such that

(6.6) 
$$d(T^n x, \mathcal{S}_k) \geqslant C e^{-\eta n^p}, \quad \forall n \geqslant 0.$$

*Proof.* First, for each  $k \ge 0$ , we claim that there exists  $C_k > 0$  such that for all  $\varepsilon > 0$ ,

(6.7) 
$$|\nu(\mathcal{N}_{\varepsilon}(\mathcal{S}_k))| \leqslant C|1_{k,\varepsilon}\nu|_w \leqslant C_k|\log\varepsilon|^{-\gamma},$$

where  $1_{k,\varepsilon} = \mathbb{1}_{\mathcal{N}_{\varepsilon}(\mathcal{S}_k)}$ . The proof of the first inequality in (6.7) is formally the same as in the proof of [3, Lemma 7.3].

We now prove the second inequality in (6.7). Let  $W \in \mathcal{W}^s$  and  $\psi \in C^{\alpha}(W)$  with  $|\psi|_{C^{\alpha}(W)} \leq$  1. Due to the uniform transversality of curves in  $\mathcal{S}_{-k}$  with the stable cone, the intersection  $W \cap \mathcal{N}_{\varepsilon}(\mathcal{S}_{-k})$  can be expressed as a finite union with cardinality bounded by a constant  $A_k$ 

(depending only on  $S_{-k}$ ) of stable manifolds  $W_i \in W^s$ , of lengths at most  $C\varepsilon$ . Therefore, for any  $f \in C^1(M)$ ,

$$\int_{W_{\xi}} f 1_{k,\varepsilon} \psi \, \mathrm{d} m_W = \sum_i \int_{W_i} f \psi \, \mathrm{d} m_{W_i} \leqslant \sum_i |f|_w |\psi|_{C^{\alpha}(W_i)} \leqslant C A_k |f|_w.$$

It follows that  $|1_{k,\varepsilon}f|_w \leqslant A_k|f|_w$  for all  $f \in \mathcal{B}_w$ . Similarly, we have  $|1_{k,\varepsilon}f|_w \leqslant A_k||f||_s|\log \varepsilon|^{-\gamma}$  for all  $f \in \mathcal{B}$ . Now, recalling  $\nu_n$ , we estimate,

$$|1_{k,\varepsilon}\nu|_w \leqslant |1_{k,\varepsilon}(\nu-\nu_n)|_w + |1_{k,\varepsilon}\nu_n|_w \leqslant A_k|\nu-\nu_n|_w + C_k'|\log\varepsilon|^{-\gamma}||\nu_n||_{\mathcal{B}}.$$

Since  $||\nu_n||_{\mathcal{B}} \leqslant \overline{C}$  for all  $n \geqslant 1$ , we take the limit as  $n \to \infty$  to conclude that  $|1_{k,\varepsilon}\nu|_w \leqslant C_k|\log \varepsilon|^{-\gamma}$ , concluding the proof of (6.7).

Next, applying (6.4), we have

$$\tilde{\nu}(\nu) \, \mu_g(\mathcal{N}_{\varepsilon}(\mathcal{S}_{-k})) = \tilde{\nu}(1_{k,\varepsilon}\nu) \leqslant \tilde{C}|1_{k,\varepsilon}\nu|_w \leqslant \tilde{C}C_k|\log \varepsilon|^{-\gamma} \quad \forall k \geqslant 0.$$

To obtain the analogous bound for  $\mathcal{N}_{\varepsilon}(\mathcal{S}_k)$ , for k > 0, we use the invariance of  $\mu_g$ . It follows from [13, Exercise 4.50] that  $T(\mathcal{N}_{\varepsilon}(\mathcal{S}_1)) \subset \mathcal{N}_{C_{\varepsilon}^{1/2}}(\mathcal{S}_{-1})$ . Thus,

$$\mu_g(\mathcal{N}_{\varepsilon}(\mathcal{S}_1)) \leqslant \mu_g(\mathcal{N}_{C\varepsilon^{1/2}}(\mathcal{S}_{-1})) \leqslant C_1 |\log C\varepsilon^{1/2}|^{-\gamma} \leqslant C_1' |\log \varepsilon|^{-\gamma}$$
.

The estimate for  $\mathcal{N}_{\varepsilon}(\mathcal{S}_k)$ , for  $k \geq 2$ , follows similarly since  $T^k \mathcal{S}_k = \mathcal{S}_{-k}$ .

Finally, fix  $\eta > 0$ ,  $k \in \mathbb{Z}$  and  $p > 1/\gamma$ . Since

(6.8) 
$$\sum_{n\geq 0} \mu_g(\mathcal{N}_{e^{-\eta n^p}}(\mathcal{S}_k)) \leqslant \tilde{C}C_k \eta^{-\gamma} \sum_{n\geq 1} n^{-p\gamma} < \infty,$$

by the Borel-Cantelli Lemma, the positive orbit of  $\mu_g$ -almost every  $x \in M$  visits  $\mathcal{N}_{e^{-\eta n^p}}(\mathcal{S}_k)$  only finitely many times, and the last part of the lemma follows.

- Corollary 6.3. a) For any  $\gamma > 0$  so that  $P_*(T,g) \sup g > \gamma s_0 \log 2$  and any  $C^1$  curve S uniformly transverse to the stable cone, there exists C > 0 such that  $\nu(\mathcal{N}_{\varepsilon}(S)) \leq C |\log \varepsilon|^{-\gamma}$  and  $\mu_g(\mathcal{N}_{\varepsilon}(S)) \leq C |\log \varepsilon|^{-\gamma}$  for all  $\varepsilon > 0$ .
  - b) The measures  $\nu$  and  $\mu_g$  have no atoms, and  $\mu_g(W) = 0$  for all  $W \in \mathcal{W}^s$  and  $W \in \mathcal{W}^u$ .
  - c) The measure  $\mu_g$  is adapted:  $\int |\log d(x, \mathcal{S}_{\pm 1})| d\mu_g < \infty$ .
  - d)  $\mu_a$ -almost every point in M has a stable and unstable manifold of positive length.

*Proof.* The proof is identical to the one of [3, Corollary 7.4], where  $\mu_*$  should be replace by  $\mu_g$ .

6.2.  $\nu$ -Almost Everywhere Positive Length of Unstable Manifolds. In this section, we establish almost everywhere positive length of unstable manifolds in the sense of the measure  $\nu$  – the maximal eigenvector of  $\mathcal{L}_g$  in  $\mathcal{B}_w$ , extended into a measure since it is nonnegative distribution. To do so, we will view elements of  $\mathcal{B}_w$  as leafwise measure (Definition 6.4). Indeed, in Lemma 6.6, we make a connection between the disintegration of  $\nu$  as a measure, and the family of leafwise measures on the set of stable manifolds  $\mathcal{W}^s$ .

**Definition 6.4** (Leafwise distribution and leafwise measure). For  $f \in C^1(M)$  and  $W \in W^s$ , the map defined on  $C^{\alpha}(W)$  by

$$\psi \mapsto \int_W f \psi \, \mathrm{d} m_W,$$

can be viewed as a distribution of order  $\alpha$  on W. Since  $|\int_W f\psi \, dm_W| \leq |f|_w |\psi|_{C^{\alpha}(W)}$ , we can extend the map sending  $f \in C^1(M)$  to this distribution of order  $\alpha$ , to  $f \in \mathcal{B}_w$ . We denote this extension by  $\int_W f\psi \, dm_W$  or  $\int_W \psi f$ , and we call the corresponding family of distributions  $(f,W)_{W\in \mathcal{W}^s}$  the leafwise distribution associated to  $f \in \mathcal{B}_w$ .

Note that if  $\int_W f \psi dm_W \geqslant 0$  for all  $\psi \geqslant 0$ , then the leafwise distribution on W can be extended into a bounded linear functional on  $C^0(W)$ , or in other words, a Radon measure. If this holds for all  $W \in \mathcal{W}^s$ , the leafwise distribution is called a leafwise measure.

**Lemma 6.5** (Almost Everywhere Positive Length of Unstable Manifolds, for  $\nu$ ). For  $\nu$ -almost every  $x \in M$  the stable and unstable manifolds have positive length. Moreover, viewing  $\nu$  as a leafwise measure, for every  $W \in \mathcal{W}^s$ ,  $\nu$ -almost every  $x \in W$  has an unstable manifold of positive length.

**Lemma 6.6.** Let  $\nu^{W_{\xi}}$  and  $\hat{\nu}$  denote the conditional measures and factor measure obtained by disintegrating  $\nu$  on the set of homogeneous stable manifolds  $W_{\xi} \in \mathcal{W}_{\mathbb{H}}^s$ ,  $\xi \in \Xi$ . Then for any  $\psi \in C^{\alpha}(M)$ ,

$$\int_{W_{\xi}} \psi \, \mathrm{d}\nu^{W_{\xi}} = \frac{\int_{W_{\xi}} \psi \rho_{\xi} \nu}{\int_{W_{\xi}} \rho_{\xi} \nu} \quad \forall \xi \in \Xi, \ and \ \mathrm{d}\hat{\nu}(\xi) = |W_{\xi}|^{-1} \left( \int_{W_{\xi}} \rho_{\xi} \nu \right) \mathrm{d}\hat{\mu}_{SRB}(\xi).$$

Moreover, viewed as a leafwise measure,  $\nu(W) > 0$  for all  $W \in \mathcal{W}^s$ .

*Proof.* First, we establish the following claim: For  $W \in \mathcal{W}^s$ , we let  $n_2 \leqslant \overline{C}_2 |\log(|W|/\delta)|$  be the constant from the proof of Corollary 3.6 (This is the first time l such that  $\mathcal{G}_l^{\delta_1}(W)$  has at least one element of length at least  $\delta_1/3$ .) Then there exists  $\overline{C} > 0$  such that for all  $W \in \mathcal{W}^s$ ,

(6.9) 
$$\int_{W} \nu \geqslant \overline{C} |W|^{(P_{*}(T,g)-\sup g)\overline{C}_{2}}.$$

Indeed, recalling (6.2) and using Theorem 5.3, we have for  $\overline{C} = \frac{c_0}{2C} \delta_1^{1-(P_*(T,g)-\inf g)\overline{C}_2}$ 

$$\int_{W} \nu = \lim_{n_{j}} \frac{1}{n_{j}} \sum_{k=0}^{n_{j}-1} e^{-kP_{*}(T,g)} \int_{W} \mathcal{L}_{g}^{k} 1 \, dm_{W}$$

$$\geqslant \lim_{n_{j}} \frac{1}{n_{j}} \sum_{k=n_{2}}^{n_{j}-1} e^{-kP_{*}(T,g)} \sum_{W_{i} \in \mathcal{G}_{n_{2}}^{\delta_{1}}(W)} \int_{W_{i}} e^{S_{n_{2}}g} \mathcal{L}_{g}^{k-n_{2}} 1 \, dm_{W_{i}}$$

$$\geqslant \lim_{n_{j}} \frac{1}{n_{j}} \sum_{k=n_{2}}^{n_{j}-1} e^{-kP_{*}(T,g)} e^{n_{2} \inf g} \frac{\delta_{1}}{2C} c_{0} e^{P_{*}(T,g)(k-n_{2})}$$

$$\geqslant \frac{\delta_{1}}{2C} c_{0} e^{-n_{2}(P_{*}(T,g)-\inf g)} \geqslant \overline{C} |W|^{(P_{*}(T,g)-\sup g)\overline{C}_{2}}.$$

This proves the last statement of the lemma.

Next, for any  $f \in C^1(M)$ , according to our convention, we view f as an element of  $\mathcal{B}_w$  by considering it as a measure integrated against  $\mu_{SRB}$ . Now let  $(\nu_{n_i})_i$  be the sequence of

functions defined by (6.2) such that  $|\nu_{n_j} - \nu|_w \to 0$ . For any  $\psi \in C^{\alpha}(M)$ , we have

$$\begin{split} \nu_{n_{j}}(\psi) &= \int_{M} \nu_{n_{j}} \psi \, \mathrm{d}\mu_{\mathrm{SRB}} = \int_{\Xi} \int_{W_{\xi}} \nu_{n_{j}} \psi \rho_{\xi} \, \mathrm{d}m_{W_{\xi}} |W_{\xi}|^{-1} \, \mathrm{d}\hat{\mu}_{\mathrm{SRB}}(\xi) \\ &= \int_{\Xi} \frac{\int_{W_{\xi}} \nu_{n_{j}} \psi \rho_{\xi} \, \mathrm{d}m_{W_{\xi}}}{\int_{W_{\xi}} \nu_{n_{j}} \rho_{\xi} \, \mathrm{d}m_{W_{\xi}}} \, \mathrm{d}(\hat{\mu}_{\mathrm{SRB}})_{n_{j}}(\xi) \end{split}$$

where  $d(\hat{\mu}_{SRB})_{n_j}(\xi) = |W_{\xi}|^{-1} \int_{W_{\xi}} \nu_{n_j} \rho_{\xi} dm_{W_{\xi}} d\hat{\mu}_{SRB}(\xi)$ . By definition of convergences in  $\mathcal{B}_w$  since  $\psi$ ,  $\rho_{\xi} \in C^{\alpha}(W_{\xi})$ , the ratio of integrals converges (uniformly in  $\xi$ ) to  $\int_{W_{\xi}} \psi \rho_{\xi} \nu / \int_{W_{\xi}} \rho_{\xi} \nu$ , and the factor measure converges to  $|W_{\xi}|^{-1} (\int_{W_{\xi}} \rho_{\xi} \nu) d\hat{\mu}_{SRB}(\xi)$ . Note that since  $\rho_{\xi}$  is uniformly log-Hölder, and due to (6.9), we have  $\int_{W_{\xi}} \rho_{\xi} \nu > 0$  with lower bound depending only on the length of  $W_{\xi}$ . Finally, by [3, Proposition 4.2] and [3, Lemma 4.4], we have  $\nu_{n_j}(\psi)$  converging to  $\nu(\psi)$ . Disintegrating  $\nu$  according the statement of the lemma yields to the claimed identifications.

Proof of Lemma 6.5. The statement about stable manifolds of positive length follows from the characterization of  $\hat{\nu}$  in Lemma 6.6, since the set of points with stable manifolds of zero length has zero  $\hat{\mu}_{SRB}$ -measure [13].

We fix  $W \in \mathcal{W}^s$  and prove the statement about  $\nu$  as a leafwise measure. This will imply the statement regarding unstable manifolds for the measure  $\nu$  by Lemma 6.6.

Fix  $\varepsilon > 0$  and  $\hat{\Lambda} \in (1, \Lambda)$ , and define  $O = \bigcup_{n \geq 1} O_n$ , where

$$O_n := \{ x \in W \mid n = \min\{ j \geqslant 1 \mid d_u(T^{-j}x, \mathcal{S}_1) < \varepsilon C_e \hat{\Lambda}^{-j} \} \},$$

and  $d_u$  denotes distance restricted to the unstable cone. By [13, Lemma 4.67], any  $x \in W \setminus \mathcal{O}$  has an unstable manifold of length at least  $2\varepsilon$ . We now estimate  $\nu(O) = \sum_{n\geqslant 1} \nu(O_n)$ , where equality holds since the  $\mathcal{O}_n$  are disjoint. Since each  $O_n$  is a finite union of open subcurves of W, we have

(6.10) 
$$\int_{W} \mathbb{1}_{O_{n}} \nu = \lim_{j \to \infty} \int_{W} \mathbb{1}_{O_{n}} \nu_{n_{j}} = \lim_{j \to \infty} \frac{1}{n_{j}} \sum_{k=0}^{n_{j}-1} e^{-kP_{*}(T,g)} \int_{W} \mathbb{1}_{O_{n}} \mathcal{L}_{g}^{k} 1 \, dm_{W}.$$

We give estimates in two cases.

Case I: k < n. Write  $\int_{W \cap O_n} \mathcal{L}_g^k 1 \, \mathrm{d} m_W = \sum_{W_i \in \mathcal{G}_k^{\delta_0}(W)} \int_{W_i \cap T^{-k}O_n} e^{S_k g} \, \mathrm{d} m_{W_i}$ .

If  $x \in T^{-k}O_n$ , then  $y = T^{-n+k}x$  satisfies  $d_u(y, \mathcal{S}_1) < \varepsilon C_e \hat{\Lambda}^{-n}$  and thus we have  $d_u(Ty, \mathcal{S}_{-1}) \le C\varepsilon^{1/2}\hat{\Lambda}^{-n/2}$ . Due to the uniform transversality of stable and unstable cones, as well as the fact that elements of  $\mathcal{S}_{-1}$  are uniformly transverse to the stable cone, we have  $d_s(Ty, \mathcal{S}_{-1}) \le C\varepsilon^{1/2}\hat{\Lambda}^{-n/2}$  as well, with possibly a larger constant C.

Let  $r_{-j}^s(x)$  denote the distance from  $T^{-j}x$  to the nearest endpoint of  $W^s(T^{-j}x)$ , where  $W^s(T^{-j}x)$  is the maximal local stable manifold containing  $T^{-j}x$ . From the above analysis, we see that  $W_i \cap T^{-k}O_n \subseteq \{x \in W_i : r_{-n+k+1}^s(x) \le C\varepsilon^{1/2}\hat{\Lambda}^{-n/2}\}$ . The time reversal of the growth lemma [13, Thm 5.52] gives  $m_{W_i}(r_{-n+k+1}^s(x) \le C\varepsilon^{1/2}\hat{\Lambda}^{-n/2}) \le C'\varepsilon^{1/2}\hat{\Lambda}^{-n/2}$  for a constant C'

that is uniform in n and k. Thus, using Proposition 3.10, we find

$$\int_{W\cap O_n} \mathcal{L}_g^k 1 \, \mathrm{d} m_W \leqslant C' \varepsilon^{1/2} \hat{\Lambda}^{-n/2} \sum_{W_i \in \mathcal{G}_k^{\delta_0}(W)} |e^{S_k g}|_{C^0(W_i)} \leqslant C e^{kP_*(T,g)} \varepsilon^{1/2} \hat{\Lambda}^{-n/2}.$$

Case II:  $k \ge n$ . Using the same observation as in Case I, if  $x \in T^{-n+1}O_n$ , then x satisfies  $d_s(x, \mathcal{S}_{-1}) \le C\varepsilon^{1/2}\hat{\Lambda}^{-n/2}$ . We change variables to estimate the integral precisely at time -n+1, and then use Theorem 5.3 and Proposition 3.10,

$$\int_{W \cap O_n} \mathcal{L}^k 1 \, dm_W = \sum_{W_i \in \mathcal{G}_{n-1}^{\delta_0}(W)} \int_{W_i \cap T^{-n+1}O_n} e^{S_{n-1}g} \mathcal{L}_g^{k-n+1} 1 \, dm_{W_i}$$

$$\leq \sum_{W_i \in \mathcal{G}_{n-1}^{\delta_0}(W)} \int_{W_i \cap (r_1^s \leq C\varepsilon^{1/2}\hat{\Lambda}^{-n/2})} e^{S_{n-1}g} \mathcal{L}_g^{k-n+1} 1 \, dm_{W_i}$$

$$\leq \sum_{W_i \in \mathcal{G}_{n-1}^{\delta_0}(W)} |\log |W_i \cap (r_1^s \leq C\varepsilon^{1/2}\hat{\Lambda}^{-n/2})||^{-\gamma} |e^{S_{n-1}g}|_{C^{\beta}(W_i)} ||\mathcal{L}_g^{k-n+1} 1||_s$$

$$\leq \sum_{W_i \in \mathcal{G}_{n-1}^{\delta_0}(W)} |\log (C\varepsilon^{1/2}\hat{\Lambda}^{-n/2})|^{-\gamma} C|e^{S_{n-1}g}|_{C^0(W_i)} e^{(k-n+1)P_*(T,g)}$$

$$\leq |\log (C\varepsilon^{1/2}\hat{\Lambda}^{-n/2})|^{-\gamma} Ce^{kP_*(T,g)}.$$

Using the estimates of Cases I and II in (6.10) and using the weaker bound, we see that,

$$\int_{W} \mathbb{1}_{O_n} \nu_{n_j} \leqslant C |\log(C\varepsilon^{1/2}\hat{\Lambda}^{-n/2})|^{-\gamma}.$$

Summing over n, we have,  $\int_W 1_O \nu_{n_j} \leq C' |\log \varepsilon|^{1-\gamma}$ , uniformly in j. Since  $\nu_{n_j}$  converges to  $\nu$  in the weak norm, this bound carries over to  $\nu$ . Since  $\varepsilon > 0$  was arbitrary and  $\gamma > 1$ , this implies  $\nu(O) = 0$ , completing the proof of the lemma.

6.3. Absolute Continuity of  $\mu_g$  – Full Support. In this subsection, we will assume that  $\gamma > 1$ , which is possible since  $P_*(T,g) - \sup g > s_0 \log 2$ . In the next subsection, we prove that  $\mu_g$  is Bernoulli. This proof relies on showing first that  $\mu_g$  is K-mixing. As a first step, we will prove that  $\mu_g$  is ergodic, using a Hopf-type argument. This will require the absolute continuity of the stable and the unstable foliations for  $\mu_g$ , which will be deduce from SSP.2 and the following absolute continuity for  $\nu$ :

**Proposition 6.7.** Let R be a Cantor rectangle. Fix  $W^0 \in W^s(R)$  and for  $W \in W^s(R)$ , let  $\Theta_W$  denote the holonomy map from  $W^0 \cap R$  to  $W \cap R$  along unstable manifolds in  $W^u(R)$ . Then for any  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential with  $P_*(T,g) - \sup g > s_0 \log 2$  and having SSP.1,  $\Theta_W$  is absolutely continuous with respect to the leafwise measure  $\nu$ .

*Proof.* Since by Lemma 6.5 unstable manifolds comprise a set of full  $\nu$ -measure, it suffices to fix a set  $E \subset W^0 \cap R$  with  $\nu$ -measure zero, and prove that the  $\nu$ -measure of  $\Theta_W(E) \subset W$  is also zero.

Since  $\nu$  is a regular measure on  $W^0$ , for  $\varepsilon > 0$ , there exists an open set  $O_{\varepsilon} \subset W^0$ ,  $O_{\varepsilon} \supset E$ , such that  $\nu(O_{\varepsilon}) \leqslant \varepsilon$ . Indeed, since  $W^0$  is compact, we may choose  $O_{\varepsilon}$  to be a finite union of

intervals. Let  $\psi_{\varepsilon}$  be a smooth function which is 1 on  $O_{\varepsilon}$  and 0 outside of an  $\varepsilon$ -neighbourhood of  $O_{\varepsilon}$ . We may choose  $\psi_{\varepsilon}$  so that  $\int_{W^0} \psi_{\varepsilon} \nu < 2\varepsilon$ .

Using (5.4), we choose  $n = n(\varepsilon)$  such that  $|\psi_{\varepsilon} \circ T^n|_{C^1(T^{-n}W^0)} \leq 1$  and  $\Lambda^{-n} \leq \varepsilon$ . Following the procedure described in the proof of the estimate on the unstable norm in Proposition 5.1, we subdivide  $T^{-n}W^0$  and  $T^{-n}W$  into matched pieces  $U_j^0$ ,  $U_j$  and unmatched pieces  $V_i^0$ ,  $V_i$ . With this construction, none of the unmatched pieces  $T^nV_i^0$  intersect an unstable manifold in  $W^u(R)$  since unstable manifolds are not cut under  $T^{-n}$ .

Indeed, on matched pieces, we may choose a foliation  $\Gamma_j = \{\gamma_x\}_{x \in U_i^0}$  such that:

- i)  $T^n\Gamma_j$  contains all unstable manifolds in  $\mathcal{W}^u(R)$  that intersect  $T^nU_i^0$ ;
- ii) between unstable manifolds in  $\Gamma_j \cap T^{-n}(\mathcal{W}^u(R))$ , we interpolate via unstable curves;
- iii) the resulting holonomy  $\Theta_j$  from  $T^nU_j^0$  to  $T^nU_j$  has uniformly bounded Jacobian<sup>16</sup> with respect to arc-length, with bound depending on the unstable diameter of D(R), by [5, Lemmas 6.6, 6.8];
- iv) pushing forward  $\Gamma_j$  to  $T^n\Gamma_j$  in D(R), we interpolate in the gaps using unstable curves; call  $\overline{\Gamma}$  the resulting foliation of D(R);
- v) the associated holonomy map  $\overline{\Theta}_W$  extends  $\Theta_W$  and has uniformly bounded Jacobian, again by [5, Lemmas 6.6 and 6.8].

Using the map  $\overline{\Theta}_W$ , we define  $\widetilde{\psi}_{\varepsilon} = \psi_{\varepsilon} \circ \overline{\Theta}_W^{-1}$ , and note that  $|\widetilde{\psi}_{\varepsilon}|_{C^1(W)} \leq C |\psi_{\varepsilon}|_{C^1(W^0)}$ , where we write  $C^1(W)$  for the set of Lipschitz functions on W, i.e.,  $C^{\alpha}$  with  $\alpha = 1$ .

Next, we modify  $\psi_{\varepsilon}$  and  $\widetilde{\psi}_{\varepsilon}$  as follows: We set them equal to 0 on the images of unmatched pieces,  $T^nV_i^0$  and  $T^nV_i$ , respectively. Since these curves do not intersect unstable manifolds in  $\mathcal{W}^u(R)$ , we still have  $\psi_{\varepsilon} = 1$  on E and  $\widetilde{\psi}_{\varepsilon} = 1$  on  $\Theta_W(E)$ . Moreover, the set of points on which  $\psi_{\varepsilon} > 0$  (resp.  $\widetilde{\psi}_{\varepsilon} > 0$ ) is a finite union of open intervals that cover E (resp.  $\Theta_W(E)$ ).

Since  $\int_{W^0} \psi_{\varepsilon} \nu < 2\varepsilon$ , in order to estimate  $\int_W \widetilde{\psi}_{\varepsilon} \nu$ , we estimate the following difference, using matched pieces

$$\int_{W^0} \psi_{\varepsilon} \, \nu - \int_{W} \widetilde{\psi}_{\varepsilon} \, \nu = e^{-nP_{*}(T,g)} \left( \int_{W^0} \psi_{\varepsilon} \, \mathcal{L}^{n} \nu - \int_{W} \widetilde{\psi}_{\varepsilon} \, \mathcal{L}^{n} \nu \right) 
= e^{-nP_{*}(T,g)} \sum_{j} \int_{U_{j}^{0}} \psi_{\varepsilon} \circ T^{n} \, e^{S_{n}g} \, \nu - \int_{U_{j}} \phi_{j} \, \nu + \int_{U_{j}} (\phi_{j} - \widetilde{\psi}_{\varepsilon} \circ T^{n} \, e^{S_{n}g}) \, \nu ,$$

where  $\phi_j = (\psi_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_j^0} \circ G_{U_j^1}^{-1}$ , and  $G_{U_j^0}$  and  $G_{U_j}$  represent the functions defining  $U_j^0$  and  $U_j$ , respectively, defined as in (5.5). Next, since  $d(\psi_{\varepsilon} \circ T^n e^{S_n g}, \phi_j) = 0$  by construction, and using (5.8) and the assumption that  $\Lambda^{-n} \leq \varepsilon$ , we have by (5.9),

(6.12) 
$$e^{-nP_*(T,g)} \left| \sum_j \int_{U_j^0} \psi_{\varepsilon} \circ T^n \, \nu - \int_{U_j} \phi_j \, \nu \right| \leqslant C |\log \varepsilon|^{-\varsigma} ||\nu||_u \, .$$

 $<sup>^{16}</sup>$ Indeed, [5] shows the Jacobian is Hölder continuous, but we shall not need this here.

It remains to estimate the last term in (6.11). This we do using the weak norm,

(6.13) 
$$\int_{U_j} (\phi_j - \widetilde{\psi}_{\varepsilon} \circ T^n e^{S_n g}) \nu \leq |\phi_j - \widetilde{\psi}_{\varepsilon} \circ T^n e^{S_n g}|_{C^{\alpha}(U_j)} |\nu|_w.$$

By (5.10), we have

$$|\phi_j - \widetilde{\psi}_{\varepsilon} \circ T^n e^{S_n g}|_{C^{\alpha}(U_j)} \leqslant C|(\psi_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_i^0} - (\widetilde{\psi}_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_j}|_{C^{\alpha}(I_j)},$$

where  $I_j$  is the common r-interval on which  $G_{U_j^0}$  an  $G_{U_j}$  are defined.

Fix  $r \in I_j$ , and let  $x = G_{U_j^0}(r) \in U_j$  and  $\bar{x} = G_{U_j}(r)$ . Since  $U_j^0$  and  $U_j$  are matched, there exist  $y \in U_j^0$  and an unstable curve  $\gamma_y \in \Gamma_j$  such that  $\gamma_y \cap U_j = \bar{x}$ . By definition of  $\widetilde{\psi}_{\varepsilon}$ , we have  $\widetilde{\psi}_{\varepsilon} \circ T^n(\bar{x}) = \psi_{\varepsilon} \circ T^n(y)$ . Thus,

$$\begin{split} |(\psi_{\varepsilon} \circ T^n \, e^{S_n g}) \circ G_{U_j^0}(r) - (\widetilde{\psi}_{\varepsilon} \circ T^n \, e^{S_n g}) \circ G_{U_j}(r)| \\ & \leqslant |\psi_{\varepsilon} \circ T^n(x) - \widetilde{\psi}_{\varepsilon} \circ T^n(\bar{x})||e^{S_n g(x)}| + |\widetilde{\psi}_{\varepsilon} \circ T^n(\bar{x})||e^{S_n g(x)} - e^{S_n g(\bar{x})}| \\ & \leqslant (|\psi_{\varepsilon} \circ T^n(x) - \psi_{\varepsilon} \circ T^n(y)| + |\psi_{\varepsilon} \circ T^n(y) - \widetilde{\psi}_{\varepsilon} \circ T^n(\bar{x})|)e^{n\sup g} + |e^{S_n g(x)} - e^{S_n g(\bar{x})}| \\ & \leqslant \left(|\psi_{\varepsilon} \circ T^n|_{C^1(U_j^0)} d(x,y) + |g|_{C^{\alpha g}} \frac{\Lambda^{\alpha_g}}{\Lambda^{\alpha_g} - 1} (C\varepsilon)^{\alpha_g}\right) e^{n\sup g} \\ & \leqslant (C\Lambda^{-n} + C\varepsilon^{\alpha_g})e^{n\sup g} \leqslant C(\varepsilon + \varepsilon^{\alpha_g})e^{n\sup g} \,, \end{split}$$

where we have used the fact that  $d(x,y) \leq C\Lambda^{-n}$  due to the uniform transversality of stable and unstable curves. We also used the fact that, by definition, the vertical segment  $\gamma_x$  connecting x to  $\bar{x}$  is such that  $|T^n\gamma_x| < C\varepsilon$ . Since each  $T^i\gamma_x$  lies in the extended unstable cone, for all  $0 \leq i \leq n$ , we get that  $d(T^i(x), T^i(\bar{x})) \leq C\Lambda^{-(n-i)}\varepsilon$ , hence the bound

$$\begin{split} |e^{S_ng(x)} - e^{S_ng(\bar{x})}| &\leqslant |e^{S_ng(x)}| \cdot |1 - e^{S_ng(\bar{x}) - S_ng(x)}| \leqslant 2e^{n\sup g} |S_ng(\bar{x}) - S_ng(x)| \\ &\leqslant \frac{\Lambda^{\alpha_g}}{\Lambda^{\alpha_g} - 1} (C\varepsilon)^{\bar{\alpha}} |g|_{C^{\alpha_g}} e^{n\sup g} \end{split}$$

where we used that  $|1 - e^x| \le 2|x|$  when x is near 0.

Now given  $r, s \in I_i$ , we have on the one hand,

$$|(\psi_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_j^0}(r) - (\widetilde{\psi}_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_j}(r) - (\psi_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_j^0}(s) + (\widetilde{\psi}_{\varepsilon} \circ T^n e^{S_n g}) \circ G_{U_j}(s)|$$

$$\leq 2C \varepsilon^{\widetilde{\alpha}} e^{n \sup g},$$

while on the other hand,

$$\begin{split} &|(\psi_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j^0}(r)-(\widetilde{\psi}_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j}(r)-(\psi_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j^0}(s)+(\widetilde{\psi}_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j}(s)|\\ &=|(\psi_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j^0}(r)-(\psi_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j^0}(s)-((\widetilde{\psi}_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j}(r)-(\widetilde{\psi}_{\varepsilon}\circ T^n\,e^{S_ng})\circ G_{U_j}(r))|\\ &\leqslant |\psi_{\varepsilon}|_{C^0(W^0)}|e^{S_ng}|_{C^{\alpha_g}}\,d(G_{U_j^0}(r),G_{U_j^0}(s))^{\alpha_g}+|\psi_{\varepsilon}\circ T^n|_{C^1(W^0)}d(G_{U_j^0}(r),G_{U_j^0}(s))|e^{S_ng}|_{C^0}\\ &+|\widetilde{\psi}_{\varepsilon}|_{C^0(W)}|e^{S_ng}|_{C^{\alpha_g}}\,d(G_{U_j^0}(r),G_{U_j^0}(s))^{\alpha_g}+|\widetilde{\psi}_{\varepsilon}\circ T^n|_{C^1(W)}d(G_{U_j^0}(r),G_{U_j^0}(s))|e^{S_ng}|_{C^0}\\ &\leqslant (C|r-s|+C'|r-s|^{\alpha_g})e^{n\sup g}\leqslant C|r-s|^{\alpha_g}e^{n\sup g}\,, \end{split}$$

where we have used Lemma 3.12 and the fact that  $G_{U_j}^{-1}$  and  $G_{U_j}^{-1}$  have bounded derivatives since the stable cone is bounded away from the vertical.

The difference between evaluation at r and s is bounded by the minimum of these two expressions. This is greatest when the two are equal, i.e., when  $|r-s|=C\varepsilon$ . Thus  $H^{\alpha}((\psi_{\varepsilon}\circ T^ne^{S_ng})\circ G_{U_j^0}-(\widetilde{\psi}_{\varepsilon}\circ T^ne^{S_ng})\circ G_{U_j})\leqslant C\varepsilon^{\alpha_g-\alpha}e^{n\sup g}$ , and so  $|\phi_j-\widetilde{\psi}_{\varepsilon}\circ T^ne^{S_ng}|_{C^{\alpha}(U_j)}\leqslant C\varepsilon^{\alpha_g-\alpha}e^{n\sup g}$ . Putting this estimate together with (6.12) and (6.13) in (6.11), we conclude,

$$(6.14) \qquad \left| \int_{W^0} \psi_{\varepsilon} \, \nu - \int_{W} \widetilde{\psi}_{\varepsilon} \, \nu \right| \leqslant C |\log \varepsilon|^{-\varsigma} \|\nu\|_{u} + C \varepsilon^{\alpha_{g} - \alpha} |\nu|_{w} e^{-n(P_{*}(T,g) - \sup g)}.$$

Now since  $\int_{W^0} \psi_{\varepsilon} \nu \leq 2\varepsilon$ , we have

(6.15) 
$$\int_{W} \widetilde{\psi}_{\varepsilon} \, \nu \leqslant C' |\log \varepsilon|^{-\varsigma} \,,$$

where C' depends on  $\nu$ . Since  $\widetilde{\psi}_{\varepsilon} = 1$  on  $\Theta_W(E)$  and  $\widetilde{\psi}_{\varepsilon} > 0$  on an open set containing  $\Theta_W(E)$  for every  $\varepsilon > 0$ , we have  $\nu(\Theta_W(E)) = 0$ , as required.

Corollary 6.8 (Absolute Continuity of  $\mu_g$  with Respect to Unstable Foliations). Let R be a Cantor rectangle with  $\mu_g(R) > 0$ . Fix  $W^0 \in \mathcal{W}^s(R)$  and for  $W \in \mathcal{W}^s(R)$ , let  $\Theta_W$  denote the holonomy map from  $W^0 \cap R$  to  $W \cap R$  along unstable manifolds in  $\mathcal{W}^u(R)$ . Then  $\Theta_W$  is absolutely continuous with respect to the measure  $\mu_g$ .

In order to deduce the corollary from the Proposition 6.7, we introduce the set  $M^{\text{reg}}$ , as in [3], of regular points and a countable cover of this set by Cantor rectangles. The set  $M^{\text{reg}}$  is defined by

$$M^{\text{reg}} = \{ x \in M \mid d(x, \partial W^s(x)) > 0, \ d(x, \partial W^u(x)) > 0 \}.$$

At each  $x \in M^{\text{reg}}$ , we can apply [13, Prop 7.81] and construct a closed locally maximal Cantor rectangle  $R_x$  containing x, which is the direct product of local stable and unstable manifolds. Furthermore, by trimming the sides, we may arrange it so that  $\frac{1}{2}\text{diam}^s(R_x) \leq \text{diam}^u(R_x) \leq 2 \text{diam}^s(R_x)$ .

**Lemma 6.9** (Countable Cover of  $M^{\text{reg}}$  by Cantor Rectangle). There exists a countable set  $\{x_j\}_{j\in\mathbb{N}}\subset M^{\text{reg}}$  such that  $\cup_j R_{x_j}=M^{\text{reg}}$  and each  $R_j:=R_{x_j}$  satisfies (3.17).

*Proof.* This lemma is exactly the content of [3, Lemma 7.10].

Let  $\{R_j \mid j \in \mathbb{N}\}$  be the family of Cantor rectangles constructed in Lemma 6.9, discarding the ones with zero  $\mu_g$ -measure. Then  $\mu_g(\cup_j R_j) = \mu_g(M^{\text{reg}}) = 1$ , by Corollary 6.3(d). In the rest of the paper, we shall work with this countable collection of rectangles.

Given a Cantor rectangle R, define  $\mathcal{W}^s(R)$  to be the set of stable manifolds that completely cross D(R), and similarly for  $\mathcal{W}^u(R)$ .

Proof of Corollary 6.8. In order to prove absolute continuity of the unstable foliation with respect to  $\mu_g$ , we will show that the conditional measures  $\mu_g^W$  of  $\mu_g$  are equivalent to  $\nu$  on  $\mu_g$ -almost every  $W \in \mathcal{W}^s(R)$ .

Fix a Cantor rectangle R satisfying (3.17) with  $\mu_g(R) > 0$ , and  $W^0$  as in the statement of Corollary 6.8. Let  $E \subset W^0 \cap R$  satisfy  $\nu(E) = 0$ , for the leafwise measure  $\nu$ .

For any  $W \in \mathcal{W}^s(R)$ , we have the holonomy map  $\Theta_W : W^0 \cap R \to W \cap R$  as in the proof of Proposition 6.7. For  $\varepsilon > 0$ , we approximate E, choose n and construct a foliation  $\overline{\Gamma}$  of the solid rectangle D(R) as before. Define  $\psi_{\varepsilon}$  and use the foliation  $\overline{\Gamma}$  to define  $\widetilde{\psi}_{\varepsilon}$  on D(R). We

have  $\widetilde{\psi}_{\varepsilon} = 1$  on  $\overline{E} = \bigcup_{x \in E} \overline{\gamma}_x$ , where  $\overline{\gamma}_x$  is the element of  $\overline{\Gamma}$  containing x. We extend  $\widetilde{\psi}_{\varepsilon}$  to M by setting it equal to 0 on  $M \setminus D(R)$ .

It follows from the proof of Proposition 6.7, in particular (6.15), that  $\widetilde{\psi}_{\varepsilon}\nu \in \mathcal{B}_w$ , and  $|\widetilde{\psi}_{\varepsilon}\nu|_w \leqslant C'|\log \varepsilon|^{-\varsigma}$ . Now,

(6.16) 
$$\widetilde{\nu}(\nu) \,\mu_g(\widetilde{\psi}_{\varepsilon}) = \widetilde{\nu}(\widetilde{\psi}_{\varepsilon}\nu) = \lim_{j \to \infty} \frac{1}{n_j} \sum_{k=0}^{n_j - 1} e^{-kP_*(T,g)} (\mathcal{L}_g^*)^k \,\mathrm{d}\mu_{\mathrm{SRB}}(\widetilde{\psi}_{\varepsilon}\nu)$$

$$= \lim_{j \to \infty} \frac{1}{n_j} \sum_{k=0}^{n_j - 1} e^{-kP_*(T,g)} \mu_{\mathrm{SRB}}(\mathcal{L}_g^k(\widetilde{\psi}_{\varepsilon}\nu)).$$

For each k, using the disintegration of  $\mu_{SRB}$  as in the proof of Lemma 6.6 with the same notation as there, we estimate,

$$\mu_{\text{SRB}}(\mathcal{L}_g^k(\widetilde{\psi}_{\varepsilon}\nu)) = \int_{\Xi} \int_{W_{\xi}} \mathcal{L}_g^k(\widetilde{\psi}_{\varepsilon}\nu) \, \rho_{\xi} \, dm_{W_{\xi}} \, |W_{\xi}|^{-1} \, d\hat{\mu}_{\text{SRB}}(\xi)$$

$$\leqslant C \int_{\Xi} |\mathcal{L}_g^k(\widetilde{\psi}_{\varepsilon}\nu)|_w \, |W_{\xi}|^{-1} \, d\hat{\mu}_{\text{SRB}}(\xi)$$

$$\leqslant C e^{kP_*(T,g)} |\widetilde{\psi}_{\varepsilon}\nu|_w \leq C e^{kP_*(T,g)} |\log \varepsilon|^{-\varsigma},$$

where we have used (5.1) in the last line. Thus  $\mu_g(\widetilde{\psi}_{\varepsilon}) \leq C |\log \varepsilon|^{-\varsigma}$ , for each  $\varepsilon > 0$ , so that  $\mu_g(\bar{E}) = 0$ .

Disintegrating  $\mu_g$  into conditional measures  $\mu_g^{W_{\xi}}$  on  $W_{\xi} \in \mathcal{W}^s$  and a factor measure  $d\hat{\mu}_g(\xi)$  on the index set  $\Xi_R$  of stable manifolds in  $\mathcal{W}^s(R)$ , it follows that  $\mu_g^{W_{\xi}}(\bar{E}) = 0$  for  $\hat{\mu}_g$ -almost every  $\xi \in \Xi_R$ . Since E was arbitrary, the conditional measures of  $\mu_g$  on  $\mathcal{W}^s(R)$  are absolutely continuous with respect to the leafwise measure  $\nu$ .

To show that in fact  $\mu_g^W$  is equivalent to  $\nu$ , suppose now that  $E \subset W^0$  has  $\nu(E) > 0$ . For any  $\varepsilon > 0$  such that  $C' |\log \varepsilon|^{-\varsigma} < \nu(E)/2$ , where C' is from (6.15), choose  $\psi_\varepsilon \in C^1(W^0)$  such that  $\nu(|\psi_\varepsilon - 1_E|) < \varepsilon$ , where  $1_E$  is the indicator function of the set E. As above, we extend  $\psi_\varepsilon$  to a function  $\widetilde{\psi}_\varepsilon$  on D(R) via the foliation  $\overline{\Gamma}$ , and then to M by setting  $\widetilde{\psi}_\varepsilon = 0$  on  $M \setminus D(R)$ .

We have  $\psi_{\varepsilon}\nu \in \mathcal{B}_w$  and by (6.14)

(6.17) 
$$\nu(\widetilde{\psi}_{\varepsilon} 1_{W}) \geqslant \nu(\psi_{\varepsilon} 1_{W^{0}}) - C' |\log \varepsilon|^{-\varsigma}, \quad \text{for all } W \in \mathcal{W}^{s}(R).$$

Now following (6.16) and disintegrating  $\mu_{SRB}$  as usual, we obtain,

To estimate this last expression, we estimate the thermodynamic sum over the curves  $W_{\xi,i}$  which properly cross the rectangle R.

By SSP.2 and the choice of  $\delta_1$  in (3.8), there exists  $k_0$ , depending only on the minimum length of  $W \in \mathcal{W}^s(R)$ , such that

$$\sum_{W_i \in L_k^{\delta_1}(W_\xi)} |e^{S_k g}|_{C^0(W_i)} \geqslant \frac{1}{3} \sum_{W_i \in \mathcal{G}_k^{\delta_1}(W_\xi)} |e^{S_k g}|_{C^0(W_i)} \,, \quad \text{for all } k \geqslant k_0.$$

By choice of our covering  $\{R_i\}$  from Lemma 6.9, all  $W_{\xi,j} \in L_k^{\delta_1}(W_{\xi})$  properly cross one of finitely many  $R_i$ . By the topological mixing property of T, there exists  $n_0$ , depending only on the length scale  $\delta_1$ , such that some smooth component of  $T^{-n_0}W_{\xi,j}$  properly crosses R. Thus, letting  $\mathcal{C}_k(W_{\xi})$  denote those  $W_{\xi,i} \in \mathcal{G}_k^{\delta_1}(W_{\xi})$  which properly cross R, we have

$$\sum_{W_{i} \in \mathcal{C}_{k+n_{0}}(W_{\xi})} |e^{S_{k}g}|_{C^{0}(W_{i})} \geqslant \sum_{W_{\xi,i} \in L_{k}^{\delta_{1}}(W_{\xi})} \sum_{\tilde{W} \subset \mathcal{G}_{n_{0}}^{\delta_{1}}(W_{\xi,i}) \cap \mathcal{C}_{k+n_{0}}(W_{\xi})} e^{n_{0}\inf g} |e^{S_{k}g}|_{C^{0}(W_{\xi,i})}$$

$$\geqslant e^{n_{0}\inf g} \sum_{W_{\xi,i} \in L_{k}^{\delta_{1}}(W_{\xi})} |e^{S_{k}g}|_{C^{0}(W_{\xi,i})}$$

$$\geqslant \frac{1}{3} e^{n_{0}\inf g} \sum_{W_{\xi,i} \in \mathcal{G}_{k}^{\delta_{1}}(W_{\xi})} |e^{S_{k}g}|_{C^{0}(W_{\xi,i})} \geqslant \frac{1}{3} c e^{n_{0}\inf g} e^{kP_{*}(T,g)},$$

for all  $k \ge k_0$ , where c > 0 depends on  $c_0$  from Proposition 3.7 as well as the minimum length of  $W \in \mathcal{W}^s(R)$ .

Using this lower bound on the sum together with (6.17) yields,

$$\mu_g(\widetilde{\psi}_{\varepsilon}) \geqslant \frac{1}{3} c e^{-n_0 P_*(T,g)} \left( \nu(\psi_{\varepsilon}) - C' |\log \varepsilon|^{-\varsigma} \right) \geqslant C'' \left( \nu(E) - |\log \varepsilon|^{-\varsigma} \right).$$

Taking  $\varepsilon \to 0$ , we have  $\mu_g(\bar{E}) \geqslant C''\nu(E)$ , and so  $\mu_g^W(\bar{E}) > 0$  for almost every  $W \in \mathcal{W}^s(R)$ .  $\square$ 

**Proposition 6.10** (Full Support). We have  $\mu_g(O) > 0$  for any open set O.

*Proof.* The proof is the same as the one of [3, Proposition 7.11], replacing  $\mu_*$  by  $\mu_g$ .

6.4. Bernoulli property of  $\mu_g$  and Variational Principle. In this section, we use the absolute continuity results on the holonomy map from Section 6.3 to establish that  $\mu_g$  is K-mixing. We also prove an upper bound on the  $\mu_g$ -measure of weighted dynamical Bowen balls. Using these estimates, we are able to prove that  $\mu_g$  is an equilibrium state for T under the potential g – that is,  $\mu_g$  realizes the sup in the definition of P(T,g) – and  $\mu_g$  satisfies the variational principle:  $P_*(T,g) = P(T,g)$ . Finally, using again the absolute continuity along side with Cantor rectangles and the bound (6.5) on the neighbourhoods of the singular sets, we can bootstrap from the K-mixing to prove that  $\mu_g$  is Bernoulli.

**Lemma 6.11** (Single Ergodic Component). If R is a Cantor rectangle with  $\mu_g(R) > 0$ , then all the stable manifolds  $W^s(R)$  are contained in a single ergodic component of  $\mu_g$ .

*Proof.* Replacing  $\mu_*$  by  $\mu_g$ , the proof of the analogous result [3, Lemma 7.15] can be applied verbatim. The proof there follows the Hopf strategy.

**Proposition 6.12.** For all  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential g such that  $P_*(T, g) - \sup g > s_0 \log 2$  and having SSP.1 and SSP.2,  $(T, \mu_g)$  is K-mixing.

*Proof.* Replacing  $\mu_*$  by  $\mu_g$ , the proof of the analogous result [3, Proposition 7.16] can be applied verbatim. We outline the steps of the proof.

First, Baladi and Demers show that  $(T^n, \mu_*)$  is ergodic for all  $n \ge 1$ . To do so, they use the topological mixing of T to prove that any two Cantor rectangles belong to the same ergodic component of  $T^n$ .

Then, they prove that T is K-mixing. To do so, they construct a measurable partition out of the stable and unstable manifolds, that is finer than the Pinsker partition  $\pi(T)$ . Using the covering of  $M^{\text{reg}}$  by Cantor rectangles  $\{R_i\}$ , and the absolute continuity of the holonony map, they prove that each  $R_i$  belongs to a single component of  $\pi(T)$ . From this, they deduce that  $\pi(T)$  contains finitely many elements on which T acts by permutation. Since  $\pi(T)$  is T-invariant and  $(T^n, \mu_*)$  is ergodic for all  $n \ge 1$ ,  $\pi(T)$  must be trivial.

**Proposition 6.13** (Upper Bounds on Weighted Dynamical Balls). Assume that  $P_*(T,g) - \sup g > s_0 \log 2$  and that SSP.1 holds. There exists  $A < \infty$  such that for all  $\varepsilon > 0$  sufficiently small,  $x \in M$ , and  $n \ge 1$ , the measure  $\mu_g$  constructed in (6.1) satisfies

$$\mu_q(e^{-S_n^{-1}g}\mathbb{1}_{B_n(x,\varepsilon)}) \leqslant Ae^{-nP_*(T,g)}.$$

*Proof.* The inequality follows from the beginning of the proof of [3, Proposition 7.12], where  $\mu_*, \mathcal{L}$  and  $h_*$  should be replaced by respectively  $\mu_g, \mathcal{L}_g$  and  $P_*(T,g)$ .

Corollary 6.14. For all  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential g such that  $P_*(T, g) - \sup g > s_0 \log 2$  and having SSP.1 and SSP.2, the measure  $\mu_g$  is an equilibrium state of T under the potential g: we have  $P_*(T, g) = h_{\mu_g}(T) + \int g \, \mathrm{d}\mu_g$ .

*Proof.* For all  $x \in M$ , let  $\mathcal{P}_0^n(x)$  denotes the element of  $\mathcal{P}_0^n$  containing x. By the Shannon–MacMillan–Breiman theorem, we have

$$-\lim_{n\to\infty} \frac{1}{n} \log \mu_g(\mathcal{P}_0^n(x)) = h_{\mu_g}(\mathcal{P}, T) = h_{\mu_g}(T) \quad \text{for } \mu_g\text{-a.e. } x \in M,$$

where the last equality follows from the Kolmogorov–Sinai theorem (because T is expansive [3, Lemma 3.4]). Furthermore, since by Proposition 6.12  $\mu_g$  is ergodic, then  $\frac{1}{n} \log e^{-S_n^{-1}g}$  converges to  $-\mu_g(g)$  as n goes to infinity. Thus

(6.19) 
$$-\lim_{n \to \infty} \frac{1}{n} \log \left( e^{-S_n^{-1} g(x)} \mu_g(\mathcal{P}_0^n(x)) \right) = h_{\mu_g}(T) + \int g \, \mathrm{d}\mu_g \quad \text{for } \mu_g\text{-a.e. } x \in M.$$

Now, by Lemma 2.3, there exists a constant C such that for all  $x \in M$  and all  $y, z \in \mathcal{P}_0^n(x)$ , we have  $|S_n^{-1}g(y) - S_n^{-1}g(z)| \leq C$ . Thus

$$e^{-C} \leqslant \frac{\mu_g \left( e^{-S_n^{-1}g} \mathbb{1}_{\mathcal{P}_0^n(x)} \right)}{e^{-S_n^{-1}g(x)} \mu_g(\mathcal{P}_0^n(x))} \leqslant e^C,$$

and so we can replace  $e^{-S_n^{-1}g(x)}\mu_g(\mathcal{P}_0^n(x))$  in (6.19) by  $\mu_g\left(e^{-S_n^{-1}g}\mathbb{1}_{\mathcal{P}_0^n(x)}\right)$ .

Now, we want to replace  $\mathcal{P}_0^n(x)$  with a dynamical ball and use Proposition 6.13. To do so, recall that for all  $\varepsilon < \varepsilon_0$ , the dynamical ball  $B_n(x,\varepsilon)$  is included in a single element of  $\mathcal{M}_0^n$ , which is itself included in at most C elements of  $\mathcal{P}_0^n$ , for some C independent of x. Thus,

$$-\liminf_{n\to\infty} \frac{1}{n} \log \mu_g \left( e^{-S_n^{-1}g} \mathbb{1}_{B_n(x,\varepsilon)} \right) \leqslant h_{\mu_g}(T) + \int g \, \mathrm{d}\mu_g,$$

On the other hand, for  $\varepsilon$  small enough, we get by Proposition 6.13,

$$-\liminf_{n\to\infty}\frac{1}{n}\log\mu_g\left(e^{-S_n^{-1}g}\mathbbm{1}_{B_n(x,\varepsilon)}\right)\geqslant P_*(T,g).$$

Combining these last two inequalities, we get  $h_{\mu_g}(T) + \int g \, d\mu_g \geqslant P_*(T,g)$ , which ends the proof.

**Proposition 6.15.** Under the assumptions of Proposition 6.12,  $(T, \mu_q)$  is Bernoulli.

*Proof.* The proof follows the arguments in Section 5 and 6 in [15], relying on the notion vwB partition introduced by Ornstein in [27]. Actually, we can apply the same modifications as in the proof of the analogous result [3, Proposition 7.19], replacing  $\mu_*$  by  $\mu_q$ .

6.5. Uniqueness of the equilibrium state. This subsection is devoted to the uniqueness of the equilibrium state  $\mu_g$  (Proposition 6.18). The proof relies on exploiting the fact that while the lower bound on weighted Bowen balls (or thermodynamic sum over elements of  $\mathcal{M}_{-n}^0$ ) cannot be improved for  $\mu_g$ -almost every x, yet if one fixes n, most elements of  $\mathcal{M}_{-n}^0$  (in the sense of thermodynamic sums) should either have unstable diameter of a fixed length, or have previously been contained in an element of  $\mathcal{M}_{-j}^0$  with this property, for some j < n (Lemma 6.16). Such elements collectively satisfy stronger lower bounds on their measure, when weighted accordingly (Lemma 6.17). Since we have establish good control of the sums over  $\mathcal{M}_{-n}^0$  and  $\mathcal{M}_0^n$  in Section 3, we will work with these partitions instead of Bowen balls.

Recalling (3.1), choose  $m_1$  large enough so that  $(Km_1 + 1)^{1/m_1} < e^{\frac{1}{4}(P_*(T,g) - \sup g)}$ . Now, choose  $\delta_2 > 0$  sufficiently small that for all  $n, k \in \mathbb{N}$ , if  $A \in \mathcal{M}_{-n}^k$  is such that

$$\max\{\operatorname{diam}^u(A), \operatorname{diam}^s(A)\} \leq \delta_2,$$

then  $A \setminus S_{\pm m_1}$  consists of at most  $Km_1 + 1$  connected components.

For  $n \ge 1$ , define

$$B_{-2n}^0 := \{ A \in \mathcal{M}_{-2n}^0 \mid \forall \, 0 \leqslant j \leqslant n/2 \,,$$
$$T^{-j}A \subset E \in \mathcal{M}_{-n+j}^0 \text{ such that } \operatorname{diam}^u(E) < \delta_2 \},$$

and its time reversal

$$\begin{split} B_0^{2n} \coloneqq \{ A \in \mathcal{M}_0^{2n} \mid \forall \, 0 \leqslant j \leqslant n/2 \,, \\ T^j A \subset E \in \mathcal{M}_0^{n-j} \text{ such that } \mathrm{diam}^s(E) < \delta_2 \}. \end{split}$$

Next, set  $B_{2n} = \{A \in \mathcal{M}_{-2n}^0 \mid \text{ either } A \in B_{-2n}^0 \text{ or } T^{-2n}A \in B_0^{2n} \}$ . Define  $G_{2n} = \mathcal{M}_{-2n}^0 \setminus B_{2n}$ . Our first lemma shows that the thermodynamic sum over elements of  $B_{2n}$  is small relative to the one over elements of  $\mathcal{M}_{-2n}^0$ , for large n. Let  $n_1 \geq 2m_1$  be chosen so that for all  $A \in \mathcal{M}_{-n}^0$ , diam<sup>s</sup> $(A) \leq C\Lambda^{-n} \leq \delta_2$  for all  $n \geq n_1$ .

**Lemma 6.16.** There exists C > 0 such that for all  $n \ge n_1$ ,

$$\sum_{A \in B_{2n}} |e^{S_{2n}^{-1}g}|_{C^0(A)} \leqslant Ce^{\frac{3}{2}nP_*(T,g)}e^{\frac{1}{2}n\sup g}(Km_1+1)^{\frac{n}{m_1}+1} \leqslant Ce^{\frac{7}{4}nP_*(T,g)+\frac{1}{4}n\sup g}.$$

Notice that since  $P_*(T,g) - \sup g > 0$ , we have that  $\frac{7}{4}P_*(T,g) + \frac{1}{4}\sup g < 2P_*(T,g)$ .

*Proof.* Let  $n \ge n_1$  and  $A \in B^0_{-2n} \subset \mathcal{M}^0_{-2n}$ . For all  $0 \le j \le \lfloor n/2 \rfloor$ , call  $A_j \in \mathcal{M}^0_{-\lceil 3n/2 \rceil - j}$  the unique element containing  $T^{-\lfloor n/2 \rfloor + j}A$ . By definition of  $B^0_{-2n}$ , we have that  $\operatorname{diam}^u(A_j) \le \delta_2$ , meanwhile  $\operatorname{diam}^s(A_j) \le \delta_2$  by choice of  $n_1$ .

By choice of  $\delta_2$ , we have that  $A_0$  is the union of at most  $Km_1 + 1$  elements of  $\mathcal{M}_{-\lceil 3n/2 \rceil}^{m_1}$ . Thus the number of connected components of  $T^{m_1}A_0$  is at most  $Km_1 + 1$ . Notice that this fact not only applies to  $A_0$ , but also to  $A_{m_1}, \ldots, A_{lm_1}, A_{\lfloor n/2 \rfloor}$ , where  $\lfloor n/2 \rfloor = lm_1 + i$ ,  $0 \le i < m_1$ . Thus, we get

$$\#\{A' \in B_{-2n}^0 \mid T^{-\lfloor n/2 \rfloor}A' \subset A_0\} \leqslant (Km_1 + 1)^{l+1} \leqslant (Km_1 + 1)^{\frac{n}{m_1} + 1}.$$

We are now able to estimate the thermodynamic sum over  $B_{-2n}^0$ :

$$\begin{split} \sum_{A \in B_{-2n}^{0}} |e^{S_{2n}^{-1}g}|_{C^{0}(A)} &= \sum_{A_{0} \in \mathcal{M}_{-\lceil 3n/2 \rceil}^{0}} \sum_{\substack{A' \in B_{-2n}^{0} \\ T^{-\lfloor n/2 \rfloor}A' \subset A_{0}}} |e^{S_{2n}^{-1}g}|_{C^{0}(A')} = \sum_{A_{0}} \sum_{A'} \left| e^{S_{\lceil 3n/2 \rceil}^{-1}g \circ T^{\lfloor n/2 \rfloor} + S_{\lfloor n/2 \rfloor}^{-1}g} \right|_{C^{0}(A')} \\ &\leqslant \sum_{A_{0}} \left| e^{S_{\lceil 3n/2 \rceil}^{-1}g} \right|_{C^{0}(A_{0})} \sum_{A'} \left| e^{S_{\lfloor n/2 \rfloor}^{-1}g} \right|_{C^{0}(A')} \\ &\leqslant e^{\frac{1}{2}n \sup g} (Km_{1} + 1)^{\frac{n}{m_{1}} + 1} \sum_{A_{0}} \left| e^{S_{\lceil 3n/2 \rceil}^{-1}g} \right|_{C^{0}(A_{0})} \\ &\leqslant Ce^{\frac{7}{4}nP_{*}(T,g) + \frac{1}{4}n \sup g}. \end{split}$$

Now, notice that  $B_0^{2n}$  is the time reversal of  $B_{-2n}^0$ , thus

$$\sum_{A \in B_0^{2n}} |e^{S_{2n}g}|_{C^0(A)} \leqslant Ce^{\frac{7}{4}nP_*(T^{-1},g) + \frac{1}{2}n\sup g} = Ce^{\frac{7}{4}nP_*(T,g) + \frac{1}{4}n\sup g}.$$

Hence

$$\sum_{A \in B_0^{2n}} |e^{S_{2n}^{-1}g}|_{C^0(T^{-2n}A)} = \sum_{A \in B_0^{2n}} |e^{S_{2n}g}|_{C^0(A)} \leqslant Ce^{\frac{7}{4}nP_*(T,g) + \frac{1}{4}n\sup g}.$$

Finally, we get

$$\sum_{A \in B_{2n}} |e^{S_{2n}^{-1}g}|_{C^0(A)} \leqslant 2Ce^{\frac{7}{4}nP_*(T,g) + \frac{1}{4}n\sup g}.$$

Next, the following lemma establishes the importance of long pieces in providing good lower bounds on the measure of weighted elements of the partition.

**Lemma 6.17.** There exists  $C_{\delta_2} > 0$  such that for all  $j \ge 1$  and all  $A \in \mathcal{M}_{-j}^0$  such that  $diam^u(A) \ge \delta_2$  and  $diam^s(T^{-j}A) \ge \delta_2$ , we have

$$\mu_q(e^{-S_j^{-1}g}\mathbb{1}_A) \geqslant C_{\delta_2}e^{-jP_*(T,g)}$$
.

*Proof.* Let  $R_1, \ldots, R_k$  be Cantor rectangles such that  $\mu_g(R_i) > 0$  for all  $1 \le i \le k$ , and such that any unstable or stable curve of length more than  $\delta_2$  crosses at least one of them. Note  $\mathcal{R}_{\delta_2} = \{R_1, \ldots, R_k\}$  this family.

Let j > 0 and  $A \in \mathcal{M}_{-j}^0$  such that  $\operatorname{diam}^u(A) \geq \delta_2$  and  $\operatorname{diam}^s(T^{-j}A) \geq \delta_2$ . By choice of  $\mathcal{R}_{\delta_2}$ , A crosses some rectangle  $R_i$  and  $T^{-j}A$  also crosses some rectangle  $R_{i'}$ . Note  $\Xi_i$  the

index set for the family of stable manifolds  $W_{\xi}$  of  $R_i$ . For  $\xi \in \Xi_i$ , let  $W_{\xi,A} := W_{\xi} \cap A$ . Since  $T^{-j}A$  properly crosses  $R_{i'}$  in the stable direction, and that  $T^{-j}$  is smooth on A, it follows that  $T^{-j}W_{\xi,A}$  is a single curve containing a stable manifold of  $R_{i'}$ .

Let  $l_{\delta_2}$  denote the length of the smallest stable manifold among the ones in the family of Cantor rectangles  $\mathcal{R}_{\delta_2}$ . Thus, for all  $\xi \in \Xi_i$ 

$$\int_{W_{\xi,A}} e^{-S_j^{-1}g} \nu = e^{-jP_*(T,g)} \int_{T^{-j}W_{\xi,A}} \nu \geqslant e^{-jP_*(T,g)} \bar{C} l_{\delta_2}^{\bar{C}_2(P_*(T,g)-\sup g)}.$$

Finally, let  $D(R_i)$  be the smallest solid rectangle containing  $R_i$ . Since  $\mu_g^W$  and  $\nu$  are equivalent on  $\mu_g$ -a.e.  $W \in \widehat{\mathcal{W}}^s$ , we get

$$\mu_g(e^{-S_j^{-1}g}\mathbb{1}_A) \geqslant \mu_g(e^{-S_j^{-1}g}\mathbb{1}_{A\cap D(R_i)}) \geqslant \int_{\Xi_i} \mu_g^{W_\xi}(e^{-S_j^{-1}g}\mathbb{1}_A) \,\mathrm{d}\hat{\mu}_g(\xi)$$
$$\geqslant C \int_{\Xi_i} \nu(e^{-S_j^{-1}g}\mathbb{1}_{A\cap W_\xi}) \,\mathrm{d}\hat{\mu}_g(\xi) \geqslant C'_{\delta_2}\hat{\mu}_g(\Xi_i)e^{-jP_*(T,g)}.$$

Since the family  $\mathcal{R}_{\delta_2}$  is finite, this proves the lemma.

**Proposition 6.18.** If g is a  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential with  $P_*(T, g) - \sup g > s_0 \log 2$ , having SSP.1 and SSP.2, then the measure  $\mu_g$  is the unique equilibrium state for T under the potential g.

*Proof.* Usually, given a known equilibrium state (thus ergodic)  $\mu_g$ , in order to prove uniqueness it suffices to check that for all T-invariant measure  $\mu$  singular with respect to  $\mu_g$ , we have  $h_{\mu}(T) + \mu(g) < h_{\mu_g}(T) + \mu_g(g)$  – see for example [23, Section 20.3]. This is the strategy we adopt.

Let  $\mu$  be a T-invariant Borel probability measure, singular with respect to  $\mu_g$ , that is there exists a Borel set  $F \subset M$  with  $T^{-1}F = F$  and  $\mu_g(F) = 0$  but  $\mu(F) = 1$ .

For each  $n \in \mathbb{N}$ , we consider the partition  $\mathcal{Q}_n$  of maximal connected components of M on which  $T^{-n}$  is continuous. By [3, Lemma 3.2 and 3.3],  $\mathcal{Q}_n$  is  $\mathcal{M}_{-n}^0$  plus isolated points whose cardinality grows at most linearly with n. Thus  $G_{2n} \subset \mathcal{Q}_{2n}$  for each n. Define  $\tilde{B}_{2n} = \mathcal{Q}_{2n} \setminus G_{2n}$ . The set  $\tilde{B}_{2n}$  contains  $B_{2n}$  plus isolated points, and so its associated thermodynamic sum is bounded by the expression in Lemma 6.16 plus  $\#\{\text{isolated points}\}e^{2n\sup g}$ . Since  $P_*(T,g) - \sup g > 0$ , we have that  $\frac{7}{4}P_*(T,g) + \frac{1}{4}\sup g > 2\sup g$ , and thus the contribution of isolated points is small compared with the one of  $B_{2n}$ .

By uniform hyperbolicity of T, the diameters of the elements of  $T^{\lfloor n/2 \rfloor} \mathcal{Q}_n$  tend to zero as n goes to infinity. This implies the following fact.

**Sublemma 6.19.** For each  $n \ge n_1$ , there exists a finite union  $C_n$  of elements of  $Q_n$  such that

$$\lim_{n \to +\infty} (\mu + \mu_g)(F \triangle T^{\lfloor n/2 \rfloor} \mathcal{C}_n) = 0.$$

*Proof.* The proof is essentially the same as [3, Sublemma 7.24] where the role of  $\mu_*$  is played by  $\mu_g$ . Since notations are introduced in this proof, we write it down for completeness and latter use.

Let  $\bar{\mu} = \mu + \mu_g$  and  $\tilde{\mathcal{Q}}_n = T^{-\lfloor n/2 \rfloor} \mathcal{Q}_n$ . For  $\delta > 0$ , by regularity of Radon measures, pick compact sets  $K_1 \subset F$  and  $K_2 \subset M \setminus F$  such that  $\max\{\bar{\mu}(F \setminus K_1), \bar{\mu}((M \setminus F) \setminus K_2)\} < \delta$ . Since

 $K_1$  and  $K_2$  are disjoint and compact, we have  $\eta = \eta_{\delta} \coloneqq d(K_1, K_2) > 0$ . If  $\operatorname{diam}(\tilde{Q}) < \eta/2$ , then either  $\tilde{Q} \cap K_1 = \emptyset$  or  $\tilde{Q} \cap K_2 = \emptyset$ . Let  $n_{\delta}$  be large enough so that the diameter of  $\tilde{Q}_k$  is smaller than  $\eta_{\delta}/2$  for all  $k \geqslant n_{\delta}$ . Fix  $n = 2n_{\delta}$  and set  $\tilde{C}_n$  to be the union of  $\tilde{Q} \in \tilde{Q}_n$  such that  $\tilde{Q} \cap K_1 \neq \emptyset$ . By construction,  $K_1 \subset \tilde{C}_n$  and  $\tilde{C}_n \cap K_2 = \emptyset$ . Hence  $\bar{\mu}(F \triangle \tilde{C}_n) \leqslant \delta + \bar{\mu}(K_1 \triangle \tilde{C}_n) \leqslant \delta + \bar{\mu}(M \setminus (K_1 \cup K_2)) \leqslant 3\delta$ . Defining  $C_n = T^{\lfloor n/2 \rfloor} \tilde{C}_n$  completes the proof.

Remark that since  $T^{-1}F = F$ , it follows that  $(\mu + \mu_g)(\mathcal{C}_n \triangle F)$  also tends to zero as  $n \to +\infty$ . Since  $\mathcal{Q}_{2n}$  is generating for  $T^{2n}$ , we have

$$h_{\mu}(T^{2n}) = h_{\mu}(T^{2n}, \mathcal{Q}_{2n}) \leqslant H_{\mu}(\mathcal{Q}_{2n}) = -\sum_{Q \in \mathcal{Q}_{2n}} \mu(Q) \log \mu(Q).$$

Thus,

$$2nP_{\mu}(T,g) = 2nh_{\mu}(T) + 2n\mu(g) = h_{\mu}(T^{2n}) + \mu(S_{2n}^{-1}g) \leqslant H_{\mu}(Q_{2n}) + \mu(S_{2n}^{-1}g)$$

$$\leqslant \sum_{Q \in Q_{2n}} \mu(Q) \left( -\log \mu(Q) + S_{2n}^{-1}g(x_Q) + C_g \right),$$

where  $x_Q \in Q$  and  $C_q$  is the constant from Lemma 2.3.

Now, we want to distinguish elements of  $\mathcal{Q}_{2n}$ . From the proof of Sublemma 6.19, for each n, there exists a compact set  $K_1(n)$  that defines  $\tilde{\mathcal{C}}_n = T^{-\lfloor n/2 \rfloor} \mathcal{C}_n$ , and satisfying  $(\mu + \mu_g)(\cup_n K_1(n)) = (\mu + \mu_g)(F)$ . We group elements  $Q \in \mathcal{Q}_{2n} \subset \mathcal{Q}_n$  according to whether  $T^{-n}Q \subset \tilde{\mathcal{C}}_n$  or  $T^{-n}Q \cap \tilde{\mathcal{C}}_n = \emptyset$ . This dichotomy is well defined because if Q is not an isolated point, and if  $T^{-n}Q \cap \tilde{\mathcal{C}}_n \neq \emptyset$ , then  $T^{-n}Q \in \mathcal{M}_{-n}^n$  is contained in an element of  $\mathcal{M}_{-\lceil n/2 \rceil}^{\lfloor n/2 \rfloor}$  that intersect  $K_1(n)$ . Thus  $Q \subset T^n \tilde{\mathcal{C}}_n = T^{\lceil n/2 \rceil} \mathcal{C}_n$  – the case where Q is an isolated point is obvious. Therefore,

$$2nP_{\mu}(T,g) \le C_{g} + \sum_{Q \subset T^{n}\tilde{\mathcal{C}}_{n}} \mu(Q) \left( -\log \mu(Q) + S_{2n}^{-1}g(x_{Q}) \right) + \sum_{Q \in \mathcal{Q}_{2n} \times T^{n}\tilde{\mathcal{C}}_{n}} \mu(Q) \left( -\log \mu(Q) + S_{2n}^{-1}g(x_{Q}) \right)$$

$$\le C_{g} + \frac{2}{e} + \mu(T^{n}\tilde{\mathcal{C}}_{n}) \log \left( \sum_{Q \subset T^{n}\tilde{\mathcal{C}}_{n}} e^{S_{2n}^{-1}g(x_{Q})} \right) + \mu(M \setminus T^{n}\tilde{\mathcal{C}}_{n}) \log \left( \sum_{Q \in \mathcal{Q}_{2n} \times T^{n}\tilde{\mathcal{C}}_{n}} e^{S_{2n}^{-1}g(x_{Q})} \right)$$

where we used in the last line that the convexity of  $x \log x$  implies that for all  $p_j > 0$  with  $\sum_{j=1}^{N} p_j \leq 1$ , and all  $a_j \in \mathbb{R}$ , we have (see [23, (20.3.5)])

$$\sum_{j=1}^{N} p_j(-\log p_j + a_j) \leqslant \frac{1}{e} + \sum_{j=1}^{N} p_j \log \sum_{i=1}^{N} e^{a_i}.$$

Then, since  $-2nP_{\mu_q} = (\mu(T^n\tilde{\mathcal{C}}_n) + \mu(M \setminus T^n\tilde{\mathcal{C}}_n))e^{-2nP_*(T,g)}$ , we write for  $n \ge n_1$ 

$$(6.20)$$

$$2n(P_{\mu}(T,g) - P_{\mu_{g}}(T,g)) - \frac{2}{e} - C_{g}$$

$$\leq \mu(T^{-n}\tilde{\mathcal{C}}_{n}) \log \left( \sum_{Q \subset T^{n}\tilde{\mathcal{C}}_{n}} e^{S_{2n}^{-1}g(x_{Q}) - 2nP_{*}(T,g)} \right) + \mu(M \setminus T^{-n}\tilde{\mathcal{C}}_{n}) \log \left( \sum_{Q \in \mathcal{Q}_{2n} \setminus T^{n}\tilde{\mathcal{C}}_{n}} e^{S_{2n}^{-1}g(x_{Q}) - 2nP_{*}(T,g)} \right)$$

$$\leq \mu(\mathcal{C}_{n}) \log \left( \sum_{\substack{Q \subset T^{n}\tilde{\mathcal{C}}_{n} \\ Q \in G_{2n}}} e^{S_{2n}^{-1}g(x_{Q}) - 2nP_{*}(T,g)} + \sum_{\substack{Q \subset T^{n}\tilde{\mathcal{C}}_{n} \\ Q \in \tilde{B}_{2n}}} e^{S_{2n}^{-1}g(x_{Q}) - 2nP_{*}(T,g)} \right)$$

$$+ \mu(M \setminus \mathcal{C}_{2n}) \log \left( \sum_{\substack{Q \in G_{2n} \setminus T^{n}\tilde{\mathcal{C}}_{n}}} e^{S_{2n}^{-1}g(x_{Q}) - 2nP_{*}(T,g)} + \sum_{\substack{Q \in \tilde{B}_{2n} \setminus T^{n}\tilde{\mathcal{C}}_{n}}} e^{S_{2n}^{-1}g(x_{Q}) - 2nP_{*}(T,g)} \right)$$

where we used that  $Q_{2n} = G_{2n} \sqcup \tilde{B}_{2n}$ . By Lemma 6.16 (and the remark concerning the contribution of isolated points), both sums over elements of  $\tilde{B}_{2n}$  are bounded by  $Ce^{-\frac{1}{4}n(P_*(T,g)-\sup g)}$ .

It remains to estimate both sums over elements of  $G_{2n}$ . To do so, we want to use Lemma 6.17, that is for each  $Q \in G_{2n}$ , we want to assign a set  $\bar{E}$  satisfying the assumptions of the lemma. Let  $Q \in G_{2n}$ . Thus  $Q \notin B_{-2n}^0$ , and so there exists  $0 \le j \le \lfloor n/2 \rfloor$  such that  $T^{-j}Q \subset E_j \in \mathcal{M}_{2n+j}^0$  with  $\dim^u(E_j) \ge \delta_2$ . Also, since  $T^{-2n}Q \notin B_0^{2n}$ , there exists  $0 \le k \le \lfloor n/2 \rfloor$  such that  $T^{-2n+k}Q \subset \tilde{E}_k \in \mathcal{M}_0^{2n-k}$  with  $\dim^s(\tilde{E}_k) \ge \delta_2$ . Thus, both  $\tilde{E}_k \in \mathcal{M}_0^{2n-k}$  and  $T^{-2n+j+k}E_j \in \mathcal{M}_{-k}^{2n-j-k}$  contain  $T^{-2n+k}Q$ . In particular, there exists  $\tilde{E} \in \mathcal{M}_0^{2n-j-k}$  containing both  $\tilde{E}_k$  and  $T^{-2n+j+k}E_j$ . Let  $\bar{E} = T^{2n-j-k}\tilde{E} \in \mathcal{M}_{-2n+j+k}^0$ . Notice that by construction  $E_j \subset \bar{E}$  and  $\tilde{E}_k \subset T^{-2n+j+k}\bar{E}$ , therefore  $\bar{E}$  satisfies  $\dim^u(\bar{E}) \ge \delta_2$  and  $\dim^s(T^{-2n+j+k}\bar{E}) \ge \delta_2$ , the assumption from Lemma 6.17. Thus,

$$\mu_q(e^{-S_{2n-j-k}^{-1}g}\mathbb{1}_{\bar{E}}) \geqslant C_{\delta_2}e^{-(2n-j-k)P_*(T,g)}$$
.

We call  $(\bar{E}, j, k)$  an admissible triple for  $Q \in G_{2n}$  if  $0 \leq j, k \leq \lfloor n/2 \rfloor$  and  $\bar{E} \in \mathcal{M}^0_{-2n+j+k}$ , with  $T^{-j}Q \subset \bar{E}$  and  $\min\{\operatorname{diam}^u(\bar{E}), \operatorname{diam}^s(T^{-2n+j+k}\bar{E})\} \geqslant \delta_2$ . By the above construction, such admissible triples always exist, but there may be many associated to a given  $Q \in G_{2n}$ . However, we can define the unique maximal triple for Q by taking first the maximum j, and then the maximum k over all admissible triples for Q.

Let  $\mathcal{E}_{2n}$  be the set of maximal triples obtained in this way from elements of  $G_{2n}$ . For  $(\bar{E}, j, k) \in \mathcal{E}_{2n}$ , let  $\mathcal{A}_M(\bar{E}, j, k)$  denote the set of  $Q \in G_{2n}$  for which the maximal triple is  $(\bar{E}, j, k)$ . The importance of the set  $\mathcal{E}_{2n}$  lies in [3, Sublemma 7.25], which we state, and prove, as follows for completeness.

**Sublemma 6.20.** Suppose that  $(\bar{E}_1, j_1, k_1)$ ,  $(\bar{E}_2, j_2, k_2) \in \mathcal{E}_{2n}$  with  $j_2 \geqslant j_1$  and  $(\bar{E}_1, j_1, k_1) \neq (\bar{E}_2, j_2, k_2)$ . Then  $T^{-(j_2-j_1)}\bar{E}_1 \cap \bar{E}_2 = \emptyset$ .

Proof. By contradiction, let  $(\bar{E}_1, j_1, k_1)$ ,  $(\bar{E}_2, j_2, k_2) \in \mathcal{E}_{2n}$  with  $j_2 \geqslant j_1$ ,  $(\bar{E}_1, j_1, k_1) \neq (\bar{E}_2, j_2, k_2)$  and  $T^{-(j_2-j_1)}\bar{E}_1 \cap \bar{E}_2 \neq \emptyset$ . Notice that  $T^{-(j_2-j_1)}\bar{E}_1 \in \mathcal{M}^{j_2-j_1}_{-2n+j_2+k_1}$  while  $\bar{E}_2 \in \mathcal{M}^0_{-2n+j_2+k_2}$ .

Consider first the case  $k_1 \leq k_2$ . Therefore  $T^{-(j_2-j_1)}\bar{E}_1 \subset E_2$ . In particular, any element  $Q \in \mathcal{A}_M(\bar{E}_1, j_1, k_1)$  satisfies  $T^{-j_2}Q \subset \bar{E}_2$ , and so  $Q \in \mathcal{A}_M(\bar{E}_2, j_2, k_2)$ , a contradiction.

Consider now the case  $k_1 > k_2$ . Therefore  $T^{-(j_2-j_1)}\bar{E}_1$  and  $\bar{E}_2$  are both contained in an element  $\bar{E}' \in \mathcal{M}^0_{-2n+j_2+k_1}$ . Since  $\bar{E}_2 \subset \bar{E}'$ , we have that  $\operatorname{diam}^u(\bar{E}') \geqslant \delta_2$ . Also, since  $T^{-2n+j_1+k_1}\bar{E}_1 \subset T^{-2n+j_2+k_1}\bar{E}'$ , we have that  $\operatorname{diam}^s(T^{-2n+j_2+k_1}\bar{E}') \geqslant \delta_2$ . Note that if  $Q \in \mathcal{A}_M(\bar{E}_1,j_1,k_1) \cup \mathcal{A}_M(\bar{E}_2,j_2,k_2)$ , then  $(\bar{E}',j_2,k_1)$  is an admissible triple for Q. Thus, if  $j_1=j_2$ , then  $\bar{E}'=\bar{E}_1$ . For  $Q \in \mathcal{A}_M(\bar{E}_2,j_2,k_2)$ , then  $Q \subset \bar{E}_1$  and so  $(\bar{E}_1,j_1,k_1)$  is an admissible triple for Q, which contradicts the maximality of  $(\bar{E}_2,j_2,k_2)$  since  $k_1>k_2$ . Similarly, if  $j_2>j_1$ , then for  $Q \in \mathcal{A}_M(\bar{E}_1,j_1,k_1)$ , the triple  $(\bar{E}',j_2,k_1)$  is admissible for Q, which contradicts the maximality of  $(\bar{E}_1,j_1,k_1)$ .

We now prove that if  $T^n\tilde{\mathcal{C}}_n\cap\mathcal{A}_M(\bar{E},j,k)\neq\emptyset$ , then  $\mathcal{A}_M(\bar{E},j,k)\subset T^n\tilde{\mathcal{C}}_n$  and  $\bar{E}\subset T^{n-j}\tilde{\mathcal{C}}_n$ . Let  $Q\in\mathcal{A}_M(\bar{E},j,k)$  be such that  $Q\cap T^n\tilde{\mathcal{C}}_n\neq\emptyset$ . Then, by definition of  $(\bar{E},j,k),\,T^{-n}Q\subset T^{-n+j}\bar{E}\in\mathcal{M}_{-n+k}^{n-j}$ . Since  $0\leqslant j,k\leqslant \lfloor n/2\rfloor$ , there exists  $E'\in\mathcal{M}_{-\lfloor n/2\rfloor}^{\lfloor n/2\rfloor}$  such that  $T^{-n+j}\bar{E}\subset E'$ . In particular, we have  $E'\in\tilde{\mathcal{Q}}_n$  and  $E'\cap\tilde{\mathcal{C}}_n\neq\emptyset$ . Thus, by construction of  $\tilde{\mathcal{C}}_n$ , we have  $\tilde{\mathcal{C}}_n\supset E'\supset T^{-n+j}\bar{E}\supset T^{-n}Q$ . In particular, we get  $Q\subset T^n\tilde{\mathcal{C}}_n$ , and thus  $\mathcal{A}_m(\bar{E},j,k)\subset T^n\tilde{\mathcal{C}}_n$ . We also get  $\bar{E}\subset T^{n-j}\tilde{\mathcal{C}}_n$ .

On the other hand, we prove that if  $T^n\tilde{\mathcal{C}}_n\cap\mathcal{A}_M(\bar{E},j,k)=\emptyset$ , then  $\mathcal{A}_M(\bar{E},j,k)\subset M\smallsetminus T^n\tilde{\mathcal{C}}_n$  and  $T^{-n+j}\bar{E}\subset M\smallsetminus\tilde{\mathcal{C}}_n$ . Let  $Q\in\mathcal{A}_M(\bar{E},j,k)$ . Then, by assumption,  $T^{-n}Q\cap\tilde{\mathcal{C}}_n=\emptyset$ . As above, there exists  $E'\in\mathcal{M}_{-\lfloor n/2\rfloor}^{\lfloor n/2\rfloor}$  containing both  $T^{-n}Q$  and  $T^{-n+j}\bar{E}$ . In particular,  $E'\in\tilde{\mathcal{Q}}_n$  and  $E'\cap\tilde{\mathcal{C}}_n=\emptyset$ . By construction of  $\tilde{\mathcal{C}}_n$ , we get that  $E'\in M\smallsetminus\tilde{\mathcal{C}}_n$ . Thus  $Q\in M\smallsetminus T^n\tilde{\mathcal{C}}_n$ , and so  $\mathcal{A}_M(\bar{E},j,k)\subset M\smallsetminus T^n\tilde{\mathcal{C}}_n$ . Also,  $T^{-n+j}\bar{E}\subset M\smallsetminus\tilde{\mathcal{C}}_n$ .

The only last step we have to do before estimating the sums over  $G_{2n}$  is to prove that for each  $(\bar{E}, j, k) \in \mathcal{E}_{2n}$ , we have

(6.21) 
$$\sum_{Q \in \mathcal{A}_M(\bar{E},j,k)} |e^{S_{2n}^{-1}g}|_{C^0(Q)} \leq C e^{(j+k)P_*(T,g)} |e^{S_{2n-j-k}^{-1}g}|_{C^0(\bar{E})}$$

where C>0 is a constant depending only on the potential g. To do so, notice that if  $Q\in \mathcal{A}_M(\bar{E},j,k)$ , then by construction,  $T^{-j}Q\subset \bar{E}$ . Thus  $T^{-n}Q\in \mathcal{M}^n_{-n}$  is a subset of  $T^{-(n-j)}\bar{E}\in \mathcal{M}^{n-j}_{-n+k}$ . Decomposing  $T^{-n}Q=Q_-\cap Q_+$  with  $Q_-\in \mathcal{M}^0_{-n}$  and  $Q_+\in \mathcal{M}^n_0$ , and  $T^{-(n-j)}\bar{E}=E_-\cap E_+$  with  $E_-\in \mathcal{M}^0_{-n+k}$  and  $E_+\in \mathcal{M}^{n-j}_0$ , we see that  $Q_-\subset E_-$  and  $Q_+\subset E_+$ . Thus

$$\sum_{Q \in \mathcal{A}_{M}(\bar{E},j,k)} |e^{S_{2n}^{-1}g}|_{C^{0}(Q)} = \sum_{Q \in \mathcal{A}_{M}(\bar{E},j,k)} |e^{S_{2n}^{-1}g \circ T^{n}}|_{C^{0}(T^{-n}Q)}$$

$$\leqslant \sum_{\substack{Q_{-} \in \mathcal{M}_{-n}^{0} \\ Q_{-} \subset E_{-}}} \sum_{\substack{Q_{+} \in \mathcal{M}_{0}^{n} \\ Q_{+} \subset E_{+}}} |e^{S_{n}^{-1}g + S_{n}g \circ T}|_{C^{0}(Q_{-} \cap Q_{+})}$$

$$\leqslant \sum_{\substack{Q_{-} \in \mathcal{M}_{-n}^{0} \\ Q_{-} \subset E_{-}}} |e^{S_{n}^{-1}g}|_{C^{0}(Q_{-})} \sum_{\substack{Q_{+} \in \mathcal{M}_{0}^{n} \\ Q_{+} \subset E_{+}}} |e^{S_{n}g \circ T}|_{C^{0}(Q_{+})}$$

$$\leqslant \sum_{\substack{Q_{-} \in \mathcal{M}_{-n}^{0} \\ Q_{-} \subset E_{-}}} |e^{S_{n}^{-1}g \circ T^{n-k}}|_{C^{0}(T^{-n+k}Q_{-})} \sum_{\substack{Q_{+} \in \mathcal{M}_{0}^{0} \\ Q_{+} \subset E_{+}}} |e^{S_{n}g \circ T \circ T^{-(n-j)}}|_{C^{0}(T^{n-j}Q_{+})}.$$

Now, notice that  $T^{-n+k}Q_- \in \mathcal{M}_{-k}^{n-k}$  is a subset of  $T^{-n+k}E_- \in \mathcal{M}_0^{n-k}$ . Thus  $T^{-n+k}Q_-$  must be of the form  $\tilde{Q}_- \cap T^{-n+k}E_-$  for some  $\tilde{Q}_- \in \mathcal{M}_{-k}^0$ . Similarly,  $T^{n-j}Q_+$  must be of the form  $\tilde{Q}_+ \cap T^{n-j}E_+$  for some  $\tilde{Q}_+ \in \mathcal{M}_0^j$ . Thus

$$\begin{split} & \sum_{Q \in \mathcal{A}_{M}(\bar{E}, j, k)} |e^{S_{2n}^{-1}g}|_{C^{0}(Q)} \\ & \leqslant \sum_{\tilde{Q}_{-} \in \mathcal{M}_{-k}^{0}} |e^{S_{2n}^{-1}g \circ T^{n-k}}|_{C^{0}(\tilde{Q}_{-} \cap T^{-n+k}E_{-})} \sum_{\tilde{Q}_{+} \in \mathcal{M}_{0}^{j}} |e^{S_{n}g \circ T \circ T^{-(n-j)}}|_{C^{0}(\tilde{Q}_{+} \cap T^{n-j}E_{+})} \\ & \leqslant \sum_{\tilde{Q}_{-} \in \mathcal{M}_{-k}^{0}} |e^{S_{k}^{-1}g}|_{C^{0}(\tilde{Q}_{-})} |e^{S_{2n-j-k}^{-1}g - S_{n-j}^{-1}g}|_{C^{0}(T^{n-j}E_{-})} \sum_{\tilde{Q}_{+} \in \mathcal{M}_{0}^{j}} |e^{S_{j}g \circ T}|_{C^{0}(\tilde{Q}_{+})} |e^{S_{n-j}^{-1}g}|_{C^{0}(T^{n-j}E_{+})} \,. \end{split}$$

Now, using Lemma 2.3, the supermultiplicativity from Lemma 3.9 and the exact exponential growth from Proposition 3.10, we get the upper bound (6.21) with  $C = 2 C_g e^{\sup g - \inf g}$ . We can now estimate the sums over elements of  $G_{2n}$ .

$$\begin{split} \sum_{\substack{Q \in G_{2n} \\ Q \subset T^{n} \tilde{\mathcal{C}}_{n}}} e^{S_{2n}^{-1} g(x_{Q}) - 2nP_{*}(T,g)} \leqslant \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset T^{n-j} \tilde{\mathcal{C}}_{n}}} \sum_{\substack{Q \in \mathcal{A}_{M}(\bar{E},j,k) \\ \bar{E} \subset T^{n-j} \tilde{\mathcal{C}}_{n}}} e^{S_{2n}^{-1} g(x_{Q}) - 2nP_{*}(T,g)} \\ \leqslant \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset T^{n-j} \tilde{\mathcal{C}}_{n}}} Ce^{-(2n-j-k)P_{*}(T,g)} |e^{S_{2n-j-k}^{-1} g}|_{C^{0}(\bar{E})} \\ \leqslant \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset T^{n-j} \tilde{\mathcal{C}}_{n}}} CC_{\delta_{2}}^{-1} \mu_{g}(e^{-S_{2n-j-k}^{-1} g} \mathbb{1}_{\bar{E}}) |e^{S_{2n-j-k}^{-1} g}|_{C^{0}(\bar{E})} \\ \leqslant CC_{\delta_{2}}^{-1} C_{g} \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset T^{n-j} \tilde{\mathcal{C}}_{n}}} \mu_{g}(\bar{E}) \leqslant CC_{\delta_{2}}^{-1} C_{g} \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset T^{n-j} \tilde{\mathcal{C}}_{n}}} \mu_{g}(T^{-n+j}\bar{E}) \\ \leqslant C' \mu_{g}(\tilde{\mathcal{C}}_{n}) \end{split}$$

where  $C' = CC_{\delta_2}^{-1}C_g$ . Similarly,

$$\begin{split} \sum_{Q \in G_{2n} \smallsetminus T^n \tilde{\mathcal{C}}_n} e^{S_{2n}^{-1} g(x_Q) - 2nP_*(T,g)} \leqslant \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset M \smallsetminus T^{n-j} \tilde{\mathcal{C}}_n}} \sum_{Q \in \mathcal{A}_M(\bar{E},j,k)} e^{S_{2n}^{-1} g(x_Q) - 2nP_*(T,g)} \\ \leqslant \sum_{\substack{(\bar{E},j,k) \in \mathcal{E}_{2n} \\ \bar{E} \subset M \smallsetminus T^{n-j} \tilde{\mathcal{C}}_n}} C e^{-(2n-j-k)P_*(T,g)} |e^{S_{2n-j-k}^{-1} g}|_{C^0(\bar{E})} \end{split}$$

$$\leq \sum_{\substack{(\bar{E},j,k)\in\mathcal{E}_{2n}\\ \bar{E}\subset M\setminus T^{n-j}\tilde{C}_{n}}} CC_{\delta_{2}}^{-1}\mu_{g}(e^{-S_{2n-j-k}^{-1}g}\mathbb{1}_{\bar{E}})|e^{S_{2n-j-k}^{-1}g}|_{C^{0}(\bar{E})}$$

$$\leq CC_{\delta_{2}}^{-1}C_{g}\sum_{\substack{(\bar{E},j,k)\in\mathcal{E}_{2n}\\ \bar{E}\subset M\setminus T^{n-j}\tilde{C}_{n}}} \mu_{g}(\bar{E}) \leq CC_{\delta_{2}}^{-1}C_{g}\sum_{\substack{(\bar{E},j,k)\in\mathcal{E}_{2n}\\ \bar{E}\subset M\setminus T^{n-j}\tilde{C}_{n}}} \mu_{g}(T^{-n+j}\bar{E})$$

$$\leq C'\mu_{g}(M\setminus\tilde{C}_{n})$$

Putting these bounds together allows us to complete our estimate of (6.20),

$$2n(P_{\mu}(T,g) - P_{\mu_g}(T,g)) - \frac{2}{e} - C_g \leqslant \mu(C_n) \log \left( C' \mu_g(C_n) + C e^{-\frac{1}{4}n(P_*(T,g) - \sup g)} \right) + \mu(M \setminus C_n) \log \left( C' \mu_g(M \setminus C_n) + C e^{-\frac{1}{4}n(P_*(T,g) - \sup g)} \right).$$

Since by Sublemma 6.19  $\mu(C_n)$  tends to 1 as  $n \to +\infty$ , while  $\mu_g(C_n)$  tends to 0 as  $n \to +\infty$ , the limit of the right-hand side tends to  $-\infty$ . This yields a contradiction unless  $P_{\mu}(T,g) < P_{\mu_g}(T,g)$ .

### 7. The Billiard Flow

Throughout this section, we see the billiard flow  $\phi_t$  as the vertical flow in the space

$$\tilde{\Omega} = \{(x, t) \in M \times \mathbb{R} \mid 0 \leqslant t \leqslant \tau(x)\} / \sim,$$

where the equivalence relation is defined by  $(x, \tau(x)) \sim (T(x), 0)$ . In other words, we see  $\phi_t$  as the suspension flow over T under the return time  $\tau$ . Furthermore, transporting the Euclidean metric on  $\mathcal{Q} \times \mathbb{S}^1$  onto  $\tilde{\Omega}$ , the flow  $\phi_t$  is uniformly hyperbolic.

**Proposition 7.1.** Let g be a  $(\mathcal{M}_0^1, \alpha_g)$ -Hölder potential such that  $P_*(T, g) - \sup g > s_0 \log 2$ , with SSP.1 and SSP.2. Let  $\bar{\mu}_g = (\mu_g(\tau))^{-1} \mu_g \otimes \lambda$ , where  $\lambda$  is the Lebesgue measure. Then  $(\phi_t, \bar{\mu}_g)$  is a K-system.

*Proof.* The ergodicity of  $(\phi_t, \bar{\mu}_g)$  follows directly from the ergodicity of  $(T, \mu_g)$  proved in Proposition 6.12.

To prove the K-mixing, we follow closely the method used in Sections 6.9, 6.10 and 6.11 from [13]. In fact, replacing  $\mu$  and  $\mu_{\Omega}$  with  $\mu_g$  and  $\bar{\mu}_g$  throughout these sections, we only have to check that [13, Exercise 6.35] is still true in order to apply verbatim the arguments. This is what we prove here.

To do so, we first need to recall some of the construction done in [13, Section 6.9]. If  $x_1$  and  $x_3$  are two nearby points in M such that

$$(7.1) \{x_2\} := W^u(x_1) \cap W^s(x_3) \neq \emptyset, \quad \{x_4\} := W^s(x_1) \cap W^u(x_3) \neq \emptyset,$$

we then construct the 4-loop  $Y_1$ ,  $Y_2$ ,  $Y_3$ ,  $Y_4$ ,  $Y_5 \in \Omega$  as follow. Let  $Y_1 = X_1 = (x_1, t)$  and  $X_3 = (x_3, t)$ . Define

$$Y_2 = W^u(Y_1) \cap W_{loc}^{ws}(X_3), \quad Y_3 = W^s(Y_2) \cap W_{loc}^{wu}(X_3),$$
  
 $Y_4 = W^u(Y_3) \cap W_{loc}^{ws}(X_1), \quad Y_5 = W^s(Y_4) \cap W_{loc}^{wu}(X_1),$ 

where  $W^u$  and  $W^s$  are unstable and stable manifolds for the flow, and  $W^{wu}_{loc}$  and  $W^{ws}_{loc}$  are local weak unstable and local weak stable manifolds for the flow. We always assume that this

construction stays under the ceiling function  $\tau$ . Actually, as proven in [13, Lemma 6.40] there exists  $\sigma$  such that  $Y_5 = \phi_{\sigma}(Y_1)$ , with  $|\sigma| = \mu_{\text{SRB}}(K)$  where K is the rectangle in M with corners  $x_1, x_2, x_3, x_4$ . Thus the 4-loops are always open.

For  $x \in M$ , let  $\mathcal{L}_x = \{\phi_t(x) \mid 0 < t < \tau(x)\}$ . Then the partition  $\{\mathcal{L}_x \mid x \in M\}$  of  $\tilde{\Omega}$  is measurable and the conditional measures of  $\bar{\mu}_g$  on  $\mathcal{L}_x$  are uniform. Call  $\lambda_x$  the Lebesgue probability measure on  $\mathcal{L}_x$ . Let  $D \subset \Omega$  be such that  $\bar{\mu}_g(D) = 1$  and let  $E_1 = \{x \in M \mid \lambda_x(\mathcal{L}_x \setminus D) = 0\}$ . Clearly,  $\mu_g(E_1) = 1$ .

We call a point  $x_1 \in E_1$  rich if for any  $\varepsilon > 0$  there exists another point  $x_3 \in E_1$  such that  $0 < d(x_1, x_3) < \varepsilon$  and (7.1) holds with  $x_2$  and  $x_4 \in E_1$ . Denote  $E_2 \subset E_1$  the set of rich points.

The analogous of [13, Exercise 6.35] is to prove that  $\mu_g(E_2) = 1$ . Let  $\{R_j\}_{j\geqslant 1}$  be the cover of  $M^{\text{reg}}$  into Cantor rectangles (discarding the ones with zero  $\mu_g$ -measure). Let R be one of those Cantor rectangle and denote  $\mu_R$  the conditional measure of  $\mu_g$  on R. It is enough to prove that  $\mu_R(E_2) = 1$ . Since  $\mu_g(E_1) = 1$  we have that  $\mu_R(E_1) = 1$ . Furthermore, since the partition of R into stable manifolds is measurable, we can disintegrate  $\mu_R$  with respect to this partition, with conditional measure  $\mu_s^W$  on  $W \in R \cap \mathcal{W}^s$ . It follows that for  $\mu_R$ -a.e. point  $x \in E_1 \cap R$ , if  $W = W(x) \in \mathcal{W}^s$  contains x then  $\mu_s^W(W \cap E_1) = 1$ . Similarly, for  $\mu_R$ -a.e. point  $x \in E_1 \cap R$ , if  $W = W(x) \in \mathcal{W}^u$  contains  $x \in \mathcal{U}^u$  then  $x \in \mathcal{U}^u$  is the conditional measure on  $x \in \mathcal{U}^u$  in the disintegration of  $x \in \mathcal{U}^u$  with respect to the measurable partition  $x \in \mathcal{U}^u$  of  $x \in \mathcal{U}^u$ . Then  $x \in \mathcal{U}^u$  is the conditional measure on  $x \in \mathcal{U}^u$  in the disintegration of  $x \in \mathcal{U}^u$  with respect to the measurable partition  $x \in \mathcal{U}^u$  of  $x \in \mathcal{U}^u$  denotes the set of points  $x \in \mathcal{U}^u$  in  $x \in \mathcal{U}^u$  that both stable and unstable conditional measure on leaves containing  $x \in \mathcal{U}^u$  give measure 1 to  $x \in \mathcal{U}^u$ 

Let  $E_2^R \subset E_2$  be the set of rich points  $x_1$  such that  $x_3$  belongs to  $R \cap E_1$  (and therefore  $x_2$  and  $x_4$  also belong to  $R \cap E_1$  by the properties of a Cantor rectangle). By contradiction, assume that  $\mu_R(E_2^R) \neq 1$ . Define the sets

$$C_2^R = \{x_1 \in E_1 \cap R \mid \exists \varepsilon > 0, \forall x_3 \in E_1 \cap R, \text{ if } 0 < d(x_1, x_3) < \varepsilon \text{ then } x_2 \notin E_1 \cap R\},\$$
  
 $C_4^R = \{x_1 \in E_1 \cap R \mid \exists \varepsilon > 0, \forall x_3 \in E_1 \cap R, \text{ if } 0 < d(x_1, x_3) < \varepsilon \text{ then } x_4 \notin E_1 \cap R\}.$ 

Note that we don't have to introduce in these definitions the condition (7.1) since it is automatically satisfied by the construction of Cantor rectangles. Thus, we have  $(E_1 \cap R) \setminus E_2^R = C_2^R \cup C_4^R$ , so that  $\mu_R(C_2^R \cup C_4^R) > 0$ . Assume first that  $\mu_R(C_2^R) > 0$ . Define the family of sets

$$C_{2,n}^R = \{x_1 \in C_2^R \mid \varepsilon \geqslant \frac{1}{n}\}.$$

Since  $\bigcup_{n\geqslant 1}C_{2,n}^R=C_2^R$  is an increasing union, there is some n such that  $\mu_R(C_{2,n}^R)>0$ . Let  $x_1\in C_{2,n}^R\cap E_R$  and  $W\in \mathcal{W}^u$  be such that  $x_1\in W$ . Let  $x_3\in E_1\cap R\cap E_R$  be such that  $0< d(x_1,x_3)<\frac{1}{n}$ . Let  $W_0\in \mathcal{W}^u$  be the unstable manifold containing  $x_3$ . By construction of  $E_R$ , we have  $\mu_u^{W_0}(W_0\cap E_1)=1$ , and since  $\mu_u^{W_0}$  have support  $W_0$  (otherwise,  $\mu_g$  would not have total support because of the absolute continuity of the holonomy), in fact we have that

$$\mu_u^{W_0}(W_0 \cap E_1 \cap B(x_1, \frac{1}{n})) > 0.$$

Thus  $\mu_u^W(\Theta_W(W_0\cap E_1\cap B(x_1,\frac{1}{n})))>0$ . Now, if  $\tilde{x}_3\in W_0\cap E_1\cap B(x_1,\frac{1}{n})$ , then  $\tilde{x}_2\notin E_1$ . In other words,  $E_1\cap\Theta_W(W_0\cap E_1\cap B(x_1,\frac{1}{n}))=\emptyset$ . Since  $x_1\in E_R$ , we have that  $\mu_u^W(W\cap E_1)=1$ , so that  $\mu_u^W(\Theta_W(W_0\cap E_1\cap B(x_1,\frac{1}{n})))=0$ , a contradiction. Thus  $E_R\cap C_{2,n}^R=\emptyset$ , so that  $\mu_R(C_2^R)=0$ . We proceed similarly, exchanging the role of  $\mathcal{W}^s$  and  $\mathcal{W}^u$ , in order to prove that  $\mu_R(C_2^R)=0$ . Finally, we get that  $\mu_R(E_2^R)=1$ , the contradiction closing the proof.

**Proposition 7.2.** Under the assumptions of Proposition 7.1,  $(\phi_t, \bar{\mu}_g)$  is Bernoulli.

*Proof.* The idea of the proof is to bootstrap from the K-mixing following the techniques of [15] with modifications similar to those in [3, Proposition 7.19]. The proof in [15] proceeds in two steps.

Step 1. Construction of  $\delta$ -regular coverings. Given  $\delta > 0$ , the idea is to cover  $\tilde{\Omega}$ , up to a set of  $\bar{\mu}_g$ -measure at most  $\delta$ , by small Cantor boxes – essentially a set of the form Cantor rectangle times interval – such that  $\bar{\mu}_g$  restricted to each Cantor box is arbitrarily close to a product measure. The basis of the boxes will be very similar to the covering  $\{R_i\}_{i\in\mathbb{N}}$  from Lemma 6.9, however, some adjustments must be made in order to guarantee uniform properties of the Jacobian of the relevant holonomy map.

Above a Cantor rectangle R with  $\mu_g(R) > 0$ , we construct a Cantor box B following the construction of P-sets from [28, Section 3]. Let  $W_1^s$  and  $W_2^s$  be the stable sides of the smallest solid rectangle D(R) containing R. Let W be a stable manifold for  $\phi_t$  projecting on  $W_1^s$  through the map  $P_-: (x,t) \in \Omega \mapsto x \in M$  if  $t < \tau(x)$ , and being such that  $W \subset \tilde{\Omega}_0 := \{(x,t) \mid 0 < t < \tau(x)\}$ . Consider the set  $W_R \subset W$  of points  $(x,t) \in W$  such that  $x \in R$ . Let  $t_0$  be small enough so that  $S = \bigcup_{t=0}^{t_0} \phi_t(W_R) \subset \Omega_0$ . Now,  $B_0$  is obtained by moving S along the unstable manifolds of  $\phi_t$  to another surface of that type, spanned by  $W_2^s$ . That is, for each  $(x,t) \in S$ , take the unstable manifold W(x,t) of  $\phi_t$  passing by (x,t), and projecting on the unstable manifold for T passing by  $x \in R$ . Let  $B_0 = \bigcup_{(x,t) \in S} W(x,t)$  and let  $B \subset B_0$  be the set of points  $(x,t) \in B_0$  such that  $x \in R$ . Notice that, up to subdividing R into smaller rectangles taking a smaller  $t_0$ , we can assume that  $B \subset \tilde{\Omega}_0$ . Thus, by construction, the set B has the property that for all  $x, y \in B$ , the local unstable manifold of x and the local weakly stable manifold of x intersect each other at a single point which lies in x. This is the crucial property of what Ornstein and Weiss, in [28], called a rectangle.

Since  $\mu_g(R) > 0$ , we have  $\bar{\mu}_g(B) = t_0 \mu_g(R) > 0$ , so that the conditional measure  $\bar{\mu}_B$  of  $\bar{\mu}_g$  restricted to B makes sense. Now, we want to equip B with a product measure, absolutely continuous with respect to  $\bar{\mu}_B$ . We proceed as follows. Since the partition of B into unstable manifolds is measurable, we can disintegrate  $\bar{\mu}_B$  into conditional measures  $\bar{\mu}^{W_\xi}$ , on  $W_\xi \cap B$  with  $\xi \in Z_\phi$ , and a factor measure  $\hat{\mu}$  on the set  $Z_\phi$  parametrizing the unstable manifolds of B. Fix a point  $z \in B$ , and consider B as the product of  $W^u(z) \cap B$  with  $W^{ws}(z) \cap B$ , where  $W^u(z)$  is the local unstable manifold of z and  $W^{ws}(z)$  is the local weak stable manifold of z. Define  $\bar{\mu}_B^p = \bar{\mu}^{W^u(z)} \otimes \hat{\mu}$ , and note that we can view  $\hat{\mu}$  as inducing a measure on  $W^{ws}(z)$ . We still have to prove that  $\bar{\mu}_B^p \ll \bar{\mu}_B$ .

Similarly, let  $\mu_R$  be the conditional measure of  $\mu_g$  restricted to R. Since the partition into unstable manifolds  $W_{\xi}$ ,  $\xi \in Z$ , is measurable, we can disintegrate  $\mu_R$  into the conditional measures  $\mu^W$  on  $W \cap R$  and a factor measure  $\hat{\mu}$  on Z. We want to relate the disintegration  $\bar{\mu}_B$  with the one of  $\mu_R$ . Notice that we can view  $Z_{\phi}$  as the set  $Z \times [0, t_0]$ , where Z parametrize the set of unstable manifolds of R through the map associating  $\xi_{\phi} \in Z_{\phi}$  with the pair  $(\xi, t)$  where  $\xi \in Z$  is such that  $P_{-}(W_{\xi_{\phi}}) = W_{\xi} \subset D(R)$  and t is the value in the definition of S where  $W_{\xi_{\phi}}$  and S intersect. Considering sets  $A \subset B$  of the form  $A = P_{-}(A) \times [t_{-}, t_{+}]$ , we get that

$$\int_{\xi_{\phi} \in Z_{\phi}} \bar{\mu}^{W_{\xi_{\phi}}}(A) \, \mathrm{d}\hat{\bar{\mu}}(\xi_{\phi}) = \bar{\mu}_{B}(A) = \int_{t_{-}}^{t_{+}} \mu_{R}(P_{-}(A)) \, \mathrm{d}t = \int_{t_{-}}^{t_{+}} \int_{\xi \in Z} \mu^{W_{\xi}}(P_{-}(A)) \, \mathrm{d}\hat{\mu}(\xi) \, \mathrm{d}t.$$

Thus, we can identify  $\bar{\mu}^{W_{\xi_{\phi}}}$  with  $\mu^{P_{-}(W_{\xi_{\phi}})}$ , and  $\mathrm{d}\hat{\mu}$  with  $\mathrm{d}\hat{\mu}\mathrm{d}t$ . From this identifications, we deduce that the projection map  $P_{W,-}$  from some W to  $P_{-}(W)$ , and its inverse are absolutely continuous. The absolute continuity of the holonomy map  $\bar{\Theta}_{W}$  between unstable manifolds  $W_{0}$  and W in B thus follows directly from the absolute continuity of the holonomy map between unstable manifolds in R since  $\bar{\Theta}_{W} = P_{W,-}^{-1} \circ \theta_{P_{-}(W)} \circ P_{W,-}$ . This implies that  $\bar{\mu}_{B}^{p}$  is absolutely continuous with respect to  $\bar{\mu}_{B}$ , and thus, also with respect to  $\bar{\mu}_{g}$ . The following definition is taken from [15].

**Definition 7.3.** For  $\delta > 0$ , a  $\delta$ -regular covering of  $\Omega$  is a finite collection  $\mathfrak{B}$  of disjoint Cantor boxes for which<sup>17</sup>,

- a)  $\bar{\mu}_g(\bigcup_{B\in\mathfrak{B}}B)\geqslant 1-\delta$ .
- b) Every  $B \in \mathfrak{B}$  satisfies  $\left|\frac{\bar{\mu}_B^p(B)}{\bar{\mu}_g(B)} 1\right| < \delta$ . Moreover, there exists  $G \subset B$ , with  $\bar{\mu}_g(G) > (1 \delta)\bar{\mu}_g(B)$ , such that  $\left|\frac{\mathrm{d}\bar{\mu}_B^p}{\mathrm{d}\bar{\mu}_g}(x) 1\right| < \delta$  for all  $x \in G$ .

By [15, Lemma 5.1], such coverings exist for any  $\delta > 0$ , and for Cantor boxes arbitrarily small. The proof essentially uses the covering of  $M^{\text{reg}}$  from Lemma 6.9 to build Cantor boxes, up to finite subdivision of the covering to ensure a). To get b), subdivide the boxes into smaller ones on which the Jacobian of the holonomy map between unstable manifolds is nearly 1. This argument relies on Lusin's theorem and goes through in our setting with no changes.

Step 2. Proof that  $\bar{\alpha}_i$  is vwb. First, define  $\bar{\alpha}_i$  to be the partition of  $\tilde{\Omega}$  into sets of the form  $\tilde{\Omega}_0 \cap (A \times [\frac{l}{2^i}, \frac{l+1}{2^i}))$ , where  $A \in \mathcal{M}_{-1}^1$  and  $l \in \mathbb{N}$ . Then  $\bar{\alpha}_0 \leqslant \bar{\alpha}_1 \leqslant \bar{\alpha}_2 \leqslant \ldots$  is such that  $\bigvee_{i=1}^{\infty} \bigvee_{n=-\infty}^{+\infty} \phi_n \bar{\alpha}_i$  generates the whole  $\sigma$ -algebra on  $\tilde{\Omega}$ . Using Theorems 4.1 and 4.2 from [15], we only need to prove that each partition  $\bar{\alpha}_i$  is vwB in order to prove that  $(\phi_t, \bar{\mu}_g)$  is Bernoulli.

Using  $\mathcal{M}_{-1}^1$  as the basis elements of  $\bar{\alpha}_i$  allows us to apply the bounds (6.5) directly since  $\partial \mathcal{M}_{-1}^1 = \mathcal{S}_1 \cup \mathcal{S}_{-1}$ . We can now apply the same arguments as in [15, Section 6.2] with the modifications described in the second part of the proof of [3, Proposition 7.19]. Actually, the only place where we need to be careful is [3, Eq. (7.33)] because of our additional horizontal cuttings. We finish the proof by dealing with this equation. We first have to recall some notations from [3] first.

Fix some  $i \in \mathbb{N}$ , and let  $\bar{\alpha} = \bar{\alpha}_i$ . Let  $\varepsilon > 0$  and define  $\delta = e^{-(\varepsilon/C')^{2/(1-\gamma)}}$  (recalling that  $\gamma > 1$ ), where C' > 0 is the constant from (7.2) below.

Let  $\mathfrak{B} = \{B_1, B_2, \dots, B_k\}$  be a  $\delta$ -regular cover of  $\tilde{\Omega}$  such that the diameters of the  $B_i$ 's are less than  $\delta$ . Define the partition  $\pi = \{B_0, B_1, B_2, \dots, B_k\}$ , where  $B_0 = \Omega \setminus \bigcup_{i=1}^k B_i$ . For each  $i \geq 1$ , let  $G_i \subset B_i$  denote the set identified in Definition 7.3(b). Since  $(\phi_{-1}, \bar{\mu}_g)$  is K-mixing, there exists an even number N = 2m, such that for any integers  $N_0$ ,  $N_1$  such that  $N < N_0 < N_1$ ,  $\delta$ -almost every atom A of  $\bigvee_{N_0 - m}^{N_1 - m} \phi_{-i}\bar{\alpha}$ , satisfies

$$\left|\frac{\bar{\mu}_g(B|A)}{\bar{\mu}_g(B)} - 1\right| < \delta, \quad \text{for all } B \in \pi,$$

where  $\bar{\mu}_g(\cdot|A)$  is the measure  $\bar{\mu}_g$  conditioned to A. Now let  $m, N_0, N_1$  be given as above and define  $\omega = \bigvee_{N_0-m}^{N_1-m} \phi_{-i}\bar{\alpha}$ . Since the estimates on the  $\bar{\mu}_g$ -measure of the bad sets  $\hat{F}_1$  and  $\hat{F}_2$ 

<sup>&</sup>lt;sup>17</sup>The corresponding definition in [15] has a third condition, but it is satisfied in our setting since the stable and unstable manifolds are one-dimensional and have bounded curvature.

do not change, we skip them and define directly the set  $\hat{F}_3$ . Define  $F_3$  to be the set of all points  $x \in \Omega \setminus B_0$  such that  $W^s(x)$  intersects the boundary of the element  $\omega(x)$  before it fully crosses the rectangle  $\pi(x)$ . Thus, if  $x \in F_3$ , there exists a subcurve of  $W^s(x)$  connecting x to the boundary of  $(\phi_{-i}\bar{\alpha})(x)$  for some  $i \in [N_0 - m, N_1 - m]$ . Then since  $\pi(x)$  has diameter less than  $\delta$ ,  $\phi_i(x)$  lies within a distance  $C\tilde{\Lambda}^{-i}\delta$  of the boundary of  $\bar{\alpha}$  – where  $\tilde{C}_1$  and  $\tilde{\Lambda} > 1$  come from the hyperbolicity of the billiard flow. Using the bound (6.5), the total measure of such points must add up to at most

$$(7.2) \qquad \sum_{i=N_0-m}^{N_1-m} \left( \frac{C}{|\log(\tilde{C}_1\tilde{\Lambda}^{-i}\delta)|^{\gamma}} + C_{\bar{\alpha}}\tilde{C}_1\tilde{\Lambda}^{-i}\delta) \right) \leqslant C_1' |\log \delta|^{1-\gamma} + C_2' \delta \leqslant C' |\log \delta|^{1-\gamma},$$

for some C'>0. Letting  $\hat{F}_3$  denote the union of atoms  $A\in\omega$  such that  $\bar{\mu}_g(F_3|A)>|\log\delta|^{\frac{1-\gamma}{2}}$ , it follows that  $\bar{\mu}_g(F_3)\leqslant C'|\log\delta|^{\frac{1-\gamma}{2}}$ . This is at most  $\varepsilon$  by choice of  $\delta$ .

The same precaution allows us to get the same bound on  $\bar{\mu}_q(\hat{F}_4)$  as in [3].

Finally, the bad set to be avoided in the construction of the joining is  $B_0 \cup (\cup_{i=1}^4 \hat{F}_i)$ . Its measure is less than  $c\varepsilon$  by choice of  $\delta$ . From this point, once the measure of the bad set is controlled, the rest of the proof in Section 6.3 of [15] can be repeated verbatim. This proves that  $\bar{\alpha}$  is vwB.

**Proposition 7.4.** Under the assumptions of Proposition 7.1, the measure  $\bar{\mu}_q$  is flow adapted<sup>18</sup>.

*Proof.* Let  $\Omega = \{(x, y, \theta) \in Q \times \mathbb{S}^1\} \subset \mathbb{T}^3$  denote the phase space for the billiard flow  $\Phi_t$  with the usual Euclidean metric denoted by  $d_{\Omega}$ . Let  $\nu_g$  be the flow invariant measure obtained as the image of  $\bar{\mu}_g$  by the conjugacy map between  $\Omega$  and  $\tilde{\Omega}$ . Let

$$\mathcal{S}_0^- = \{ \Phi_{-t}(z) \in \Omega \mid z \in \mathcal{S}_0 \text{ and } t \leqslant \tau(T^{-1}z) \}$$

denote the flow surface obtained by flowing  $S_0$  backward until its first collision under the inverse flow. Similarly, let

$$\mathcal{S}_0^+ = \{ \Phi_t(z) \in \Omega \mid z \in \mathcal{S}_0 \text{ and } t \leqslant \tau(z) \}$$

denote the forward flow of  $S_0$  until its first collision. To show that the measure  $\nu_t$  is flow-adapted, it suffices to show that  $\int_{\Omega} |\log d_{\Omega}(x, S_0^{\pm})| d\nu_g(x) < \infty$ . For then this implies that  $\log \|D\Phi_t\|$  is integrable for each  $t \in [-\tau_{\min}, \tau_{\min}]$  and then by subadditivity for each  $t \in \mathbb{R}$ .

Let  $P^{\pm}(\cdot)$  denote the projection under the forward (backward) flow of a subset of  $\Omega$  until first collision. Let  $N_{\varepsilon}^{M}(\cdot)$  denote the  $\varepsilon$ -neighborhood of a set in M in the Euclidean metric  $d_{M}$  and let  $N_{\varepsilon}^{\Omega}(\cdot)$  denote the  $\varepsilon$ -neighborhood of a set in  $\Omega$  in the metric  $d_{\Omega}$ . It follows from [13, Exercise 3.15], that there exists C > 0 such that for any  $\varepsilon > 0$ ,

$$(7.3) P^-(N_\varepsilon^\Omega(\mathcal{S}_0^-)) \subset N_{C\varepsilon^{1/2}}^M(\mathcal{S}_1) \text{and similarly} P^+(N_\varepsilon^\Omega(\mathcal{S}_0^+)) \subset N_{C\varepsilon^{1/2}}^M(\mathcal{S}_{-1})$$

From (6.5), there exist  $C_q > 0$  and  $\gamma > 1$  such that

(7.4) 
$$\mu_q(N_{\varepsilon}^M(\mathcal{S}_{\pm 1})) \leqslant C_q |\log \varepsilon|^{-\gamma} \quad \text{for all } \varepsilon > 0.$$

Putting together (7.3) and (7.4) yields

(7.5) 
$$\nu_g(N_{\varepsilon}^{\Omega}(\mathcal{S}_0^-)) \leqslant \tau_{\max} C_g |\log C \varepsilon^{1/2}|^{-\gamma} \le C' \tau_{\max} |\log \varepsilon|^{-\gamma}.$$

<sup>&</sup>lt;sup>18</sup>This result is due to Mark Demers. I thank him for allowing me to use his proof.

For p > 1 to be chosen below, define for  $n \ge 1$ ,  $A_n = N_{e^{-np}}^{\Omega}(\mathcal{S}_0^-) \setminus N_{e^{-(n+1)p}}^{\Omega}(\mathcal{S}_0^-)$ . If  $x \in A_n$ , then  $|\log d_{\Omega}(x, \mathcal{S}_0^-)| \le (n+1)^p$ . Thus we estimate using (7.5),

$$\begin{split} \int_{\Omega} |\log d_{\Omega}(x, \mathcal{S}_{0}^{-})| \, \mathrm{d}\nu_{g} &\leqslant 1 + \log \mathrm{diam}(\Omega) + \sum_{n \geq 1} \int_{A_{n}} |\log d_{\Omega}(x, \mathcal{S}_{0}^{-})| \, \mathrm{d}\nu_{g} \\ &\leqslant 1 + \log \mathrm{diam}(\Omega) + \sum_{n \geq 1} (n+1)^{p} C' \tau_{\max} n^{-\gamma p} \,, \end{split}$$

and the series converges as long as  $p > 1/(\gamma - 1)$ . A similar argument shows that  $\log d_{\Omega}(x, \mathcal{S}_0^+)$  is  $\nu_g$  integrable so that  $\nu_g$  is flow adapted.

## APPENDIX A. MOTIVATIONS FROM UNIFORM HYPERBOLIC DYNAMICS

We start this note by presenting the usual method the existence of measures of maximal entropy is proved in the case of uniform hyperbolicity. First, we consider a hyperbolic transformation of a compact set, and then the case of an Anosov flow.

A.1. **Hyperbolic maps.** Let X be a compact Riemannian manifold and let  $T: X \to X$  be a  $\mathcal{C}^r$  diffeomorphism. Assume that T is uniformly hyperbolic, that is

$$\exists \lambda > 1, \exists C > 0, \exists E^s, E^u \subset TX \text{ such that}$$

$$(i) TX = E^s \oplus E^u, DT(E^s) \subset E^s, DT^{-1}(E^u) \subset E^u,$$

$$(ii) ||D_x T^n v_s|| \leqslant C\lambda^{-n} ||v_s||, \quad \forall n \geqslant 0, \forall v_s \in E^s_x \subset T_x X,$$

$$(iii) ||D_x T^{-n} v_u|| \leqslant C\lambda^{-n} ||v_u||, \quad \forall n \geqslant 0, \forall v_u \in E^u_x \subset T_x X.$$

One fundamental theorem about hyperbolic dynamic is the Hadamard–Perron Theorem [23, Theorem 6.2.8] which states that there exist two unique families of  $C^r$  manifolds,  $\{W_m^+\}_{m\in\mathbb{Z}}$  and  $\{W_m^-\}_{m\in\mathbb{Z}}$ , everywhere tangent respectively to  $E^s$  and to  $E^u$ , obtained as the graph of some functions, and satisfying some stability and contraction properties. A key tool in the proof is the construction of families of stable and unstable cones.

As a consequence [23, Corollary 6.4.10], all such diffeomorphisms are expansive, that is

(A.1) 
$$\exists \delta > 0, \, \forall x, y \in X, \, [d(T^n(x), T^n(y)) < \delta, \, \forall n \in \mathbb{Z} \Rightarrow x = y].$$

From the expansive property, it follows from [32, Theorem 8.2] that the metric entropy  $\mu \mapsto h_{\mu}(T)$  is upper semi-continuous, hence the existence of equilibrium states for every continuous potential – and in particular existence of measures of maximal entropy for the zero potential. In the proof of [32, Theorem 8.2], expansiveness is only use to get the equality  $h_{\mu}(T) = h_{\mu}(T, \mathcal{A})$  for partition  $\mathcal{A}$  with diam( $\mathcal{A}$ ) <  $\delta$  (the expansivity constant of T) and any T-invariant measure  $\mu$ .

As proved by Bowen [6, Theorem 3.5], the expansiveness assumption of [32, Theorem 8.2] can be weakened to *entropy-expansiveness* (the proof remains unchanged). This weakening will be relevant in the case of Anosov flows.

A.2. **Anosov flows.** Let X be a compact manifold and  $\phi = \{\varphi^t\} : \mathbb{R} \times X \to X$  be a smooth flow. Assume that  $\phi$  is an Anosov flow, that is

$$\exists \lambda > 1, \exists C > 0, \exists E^s, E^u, E^c \subset TX \text{ such that}$$

$$(i) TX = E^c \oplus E^s \oplus E^u,$$

$$(ii) D\varphi^t(E^{s/u}) = E^{s/u}, \dim E_x^c = 1, \frac{d}{dt} \Big|_{t=0} \varphi^t(x) \in E_x^c \setminus \{0\},$$

$$(iii) ||D_x \varphi^t_{|E_x^s|}| \leq C\lambda^{-t}, ||D_x \varphi^{-t}_{|E_x^u|}| \leq C\lambda^{-t}, \forall t \geq 0.$$

In [7, Proposition 1.6], Bowen proves that an Anosov flow is *flow expansive* (in the sense of Bowen–Walters), that is – as defined in [8] in the case of a fixed-point free flow,

(A.2) 
$$\forall \varepsilon > 0, \ \exists \delta > 0, \ \forall x, y \in X, \ \forall h \in \mathcal{C}^0(\mathbb{R}) \text{ with } h(0) = 0, \\ [d(\varphi^t(x), \varphi^{h(t)}(y)) < \delta, \ \forall t \in \mathbb{R} \Rightarrow y \in \varphi^{]-\varepsilon,\varepsilon[}(x)].$$

The key ingredient of the proof is the local product structure for hyperbolic flows. From (A.2), it is easy to see, for h = id, that an Anosov flow satisfies the following weaker property

(A.3) 
$$\exists \varepsilon > 0, \ \exists s > 0, \ \forall x \in X, \\ \Gamma_{\varepsilon}(x) \coloneqq \{ y \in X \mid \forall t \in \mathbb{R}, \ d(\varphi^{t}(x), \varphi^{t}(y)) < \varepsilon \} \subset \varphi^{[-s,s]}(x).$$

Bowen proved [6, Example 1.6] that (A.3) is a sufficient condition so that every time  $\varphi^t$  of the flow is entropy-expansive. Therefore the map  $\mu \in \mathcal{M}_X(\varphi^1) \mapsto h_{\mu}(\varphi^1)$  is upper semi-continuous, and so is its restriction to  $\mathcal{M}_X(\phi) \subset \mathcal{M}_X(\varphi^1)$ . Hence, Anosov flows have equilibrium states for every continuous potential, and in particular for the zero potential, measures of maximal entropy.

## APPENDIX B. OBSTRUCTIONS FOR THE BILLIARD FLOW

In the previous section, in both situations, proofs of existence of MME use some sort of expansiveness. However, the existence of a local product structure is a key ingredient in order to establish the expansivity property: it gives a scale used as the  $\delta$  in (A.1) and the  $\varepsilon$  in (A.3). Furthermore, the uniform contraction of stable (resp. unstable) manifolds for large positive (resp. negative) times is used, and not some estimates of their lengths in negative (resp. positive) times (such as fragmentation or growth lemmas, see for example [13]).

B.1. Entropy expansiveness. In Bowen's proof, the local product structure is the main tool in order to prove flow expansiveness. In the case of the billiard flow, their is no such structure. Indeed, stable and unstable manifolds exist only for Lebesgue-almost every point and there is no deterministic control of their length (hence no uniform scale for a local structure). One might argue that a billiard flow admits invariant "cone" fields [5, Section 2] and construct stable and unstable curves, but then the control on the length of those curves when applying the flow is in term of expansion, not in term of contraction.

It then seems that h-expansiveness of each time  $\varphi^t$  of the flow is too much to ask for. Still, one might hope that each  $\varphi^t$  is asymptotically h-expansive, that is  $h^*(\varphi^t) := \lim_{\varepsilon \to 0} h^*(\varphi^t, \varepsilon) = 0$ , where  $h^*(\varphi^t, \varepsilon) = \sup_{x \in X} h(\varphi^t, \overline{B}(x, \varepsilon))$ . This definition was first introduced by Misiurewicz in

[26] where he proved that the metric entropy of an asymptotic h-expansive transformation is upper semi-continuous.

The quantity  $h^*(\varphi^t)$  is usually referred to as the topological tail entropy of  $\varphi^t$  [19]. In the context of smooth dynamics, Buzzi [10] has shown that if  $f \in \mathcal{C}^r(M)$ , then  $h^*(f) \leq \frac{\dim(M)R(f)}{r}$  for some constant R(f). In particular, the metric entropy of a  $\mathcal{C}^{\infty}$  transformation is upper semi-continuous. Clearly, this result does not apply to billiard flows.

Proving that the topological tail entropy of the billiard flow is zero is enough to prove the upper semi-continuity of the metric entropy, hence the existence of some measure of maximal entropy.

B.2. Relations with the Collision Map. In [8, Theorem 6], Bowen and Walters prove that the special flow constructed over a continuous transformation and under a continuous return time function, is flow expansive if and only if the base map is expansive. Since flow expansiveness is an invariant for flow under reparametrization, without loss of generality, the return time function can be chosen constant.

In [3], Baladi and Demers show that the collision map is expansive. However, since the return time is only piecewise continuous, it is not easy to relate the expansivity of the collision map to flow expansiveness of the billiard flow. As shown in Figure 2, two trajectories can be easily separated by the collision map, but they remain *close* in the phase space of the flow. We see that for a  $\delta$  too large in (A.2) (and a natural choice of h), the two trajectories cannot be distinguished. What could be a good choice for  $\delta$ ? The main problem being to find a  $\delta$  independent of trajectories (it is *easier* to find a  $\delta$  for specific trajectories, such as those ones in Figure 2, but the inf of those  $\delta$  over all trajectories might be 0). If such  $\delta$  existed, we expect it is controlled in some way by  $\tau_{min}$ .

For similar reasons, it appears that it is not a simple consequence of the collision map expansiveness for the flow to satisfy condition (A.3) (which is a weaker than *flow expansiveness*). For example, the two orbits shown in Figure 2 (b) are close in the phase space of the flow, but far apart in the phase space of the collision map (since the collisions they make are distinct).

#### References

- [1] V. Baladi. Dynamical zeta functions and dynamical determinants for hyperbolic maps, volume 68 of Ergebnisse der Mathematik und ihrer Grenzgebiete. 3. Folge. A Series of Modern Surveys in Mathematics [Results in Mathematics and Related Areas. 3rd Series. A Series of Modern Surveys in Mathematics]. Springer, Cham, 2018. A functional approach.
- [2] V. Baladi, J. Carrand, and M. F. Demers. Measure of maximal entropy for finite horizon Sinai billiard flow. *in preparation*.
- [3] V. Baladi and M. F. Demers. On the measure of maximal entropy for finite horizon Sinai billiard maps. J. Amer. Math. Soc., 33(2):381–449, 2020.
- [4] V. Baladi and M. F. Demers. Thermodynamic formalism for dispersing billiards. *Journal of Modern Dynamics*, 18(0):415–493, 2022.
- [5] V. Baladi, M. F. Demers, and C. Liverani. Exponential decay of correlations for finite horizon Sinai billiard flows. *Invent. Math.*, 211(1):39–177, 2018.
- [6] R. Bowen. Entropy-expansive maps. Transactions of the American Mathematical Society, 164:323–331, 1972.
- [7] R. Bowen. Periodic orbits for hyperbolic flows. American Journal of Mathematics, 94(1):1–30, 1972.
- [8] R. Bowen and P. Walters. Expansive one-parameter flows. *Journal of differential Equations*, 12(1):180–193, 1972

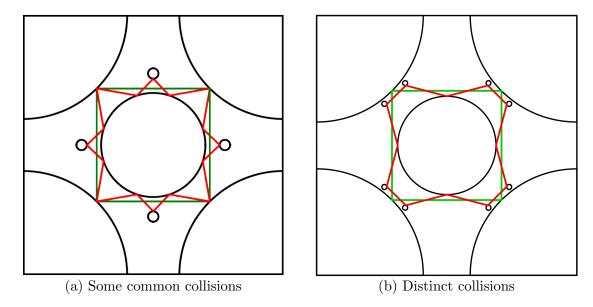


Figure 2. Two examples of two periodic trajectories.

- [9] L. A. Bunimovič and J. G. Sinaĭ. The fundamental theorem of the theory of scattering billiards. *Mat. Sb.* (N.S.), 90(132):415–431, 479, 1973.
- [10] J. Buzzi. Intrinsic ergodicity of smooth interval maps. Israel J. Math., 100:125-161, 1997.
- [11] J. Buzzi. The degree of Bowen factors and injective codings of diffeomorphisms. J. Mod. Dyn., 16:1–36, 2020
- [12] J. Chen, F. Wang, and H.-K. Zhang. Markov partition and thermodynamic formalism for hyperbolic systems with singularities, 2017.
- [13] N. Chernov and R. Markarian. Chaotic billiards. Number 127. American Mathematical Soc., 2006.
- [14] N. I. Chernov. Sinai billiards under small external forces. Ann. Henri Poincaré, 2(2):197–236, 2001.
- [15] N. I. Chernov and C. Haskell. Nonuniformly hyperbolic K-systems are Bernoulli. Ergodic Theory Dynam. Systems, 16(1):19–44, 1996.
- [16] M. F. Demers and A. Korepanov. Rates of mixing for the measure of maximal entropy of dispersing billiard maps, 2022.
- [17] M. F. Demers, L. Rey-Bellet, and H.-K. Zhang. Fluctuation of the entropy production for the Lorentz gas under small external Forces. Comm. Math. Phys., 363(2):699-740, 2018.
- [18] M. F. Demers and H.-K. Zhang. Spectral analysis of the transfer operator for the Lorentz gas. J. Mod. Dyn., 5(4):665–709, 2011.
- [19] T. Downarowicz. Entropy in dynamical systems, volume 18 of New Mathematical Monographs. Cambridge University Press, Cambridge, 2011.
- [20] G. Gallavotti and D. S. Ornstein. Billiards and Bernoulli schemes. Comm. Math. Phys., 38:83–101, 1974.
- [21] P. L. Garrido. Kolmogorov-Sinai entropy, Lyapunov exponents, and mean free time in billiard systems. J. Statist. Phys., 88(3-4):807–824, 1997.
- [22] P. Gaspard and F. Baras. Chaotic scattering and diffusion in the Lorentz gas. Phys. Rev. E (3), 51(6, part A):5332–5352, 1995.
- [23] A. Katok and B. Hasselblatt. Introduction to the modern theory of dynamical systems, volume 54. Cambridge university press, 1997.
- [24] A. Katok, J.-M. Strelcyn, F. Ledrappier, and F. Przytycki. Invariant manifolds, entropy and billiards; smooth maps with singularities, volume 1222 of Lecture Notes in Mathematics. Springer-Verlag, Berlin, 1986.

- [25] Y. Lima and C. Matheus. Symbolic dynamics for non-uniformly hyperbolic surface maps with discontinuities. Ann. Sci. Éc. Norm. Supér. (4), 51(1):1–38, 2018.
- [26] M. Misiurewicz. Diffeomorphism without any measure with maximal entropy. Bull. Acad. Polon. Sci. Sér. Sci. Math. Astronom. Phys., 21:903–910, 1973.
- [27] D. S. Ornstein. Imbedding Bernoulli shifts in flows. In Contributions to Ergodic Theory and Probability (Proc. Conf., Ohio State Univ., Columbus, Ohio, 1970), pages 178–218. Springer, Berlin, 1970.
- [28] D. S. Ornstein and B. Weiss. Geodesic flows are Bernoullian. Israel J. Math., 14:184–198, 1973.
- [29] M. Reed and B. Simon. Methods of modern mathematical physics. I. Academic Press, Inc. [Harcourt Brace Jovanovich, Publishers], New York, second edition, 1980. Functional analysis.
- [30] J. G. Sinaĭ. Dynamical systems with elastic reflections. Ergodic properties of dispersing billiards. Uspehi Mat. Nauk, 25(2 (152)):141–192, 1970.
- [31] Y. G. Sinaĭ and N. I. Chernov. Ergodic properties of some systems of two-dimensional disks and three-dimensional balls. Uspekhi Mat. Nauk, 42(3(255)):153–174, 256, 1987.
- [32] P. Walters. An introduction to ergodic theory, volume 79 of Graduate Texts in Mathematics. Springer-Verlag, New York-Berlin, 1982.
- [33] L.-S. Young. Statistical properties of dynamical systems with some hyperbolicity. *Ann. of Math.* (2), 147(3):585–650, 1998.

LABORATOIRE DE PROBABILITÉS, STATISTIQUES ET MODÉLISATION (LPSM), CNRS, SORBONNE UNIVERSITÉ, UNIVERSITÉ DE PARIS, 4, PLACE JUSSIEU, 75005 PARIS, FRANCE

Email address: jcarrand@lpsm.paris