

# Matrix product state methods for excitations

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A Panorama of Tensor Networks, Krakow

9 October 2025

# Infinite matrix product states

$$|\Psi\rangle = \begin{array}{ccccccc} & A_1^{s_1} & A_2^{s_2} & A_3^{s_3} & & \dots & A_N^{s_N} \\ & \bigcirc & \bigcirc & \bigcirc & & & \bigcirc \\ | & | & | & | & & & | \end{array}$$

Enforce translation invariance and take  $N \rightarrow \infty$ :

$$|\Psi\rangle = \begin{array}{ccccccc} & A^{s_1} & A^{s_2} & A^{s_3} & A^{s_4} & A^{s_5} & \\ & \bigcirc & \bigcirc & \bigcirc & \bigcirc & \bigcirc & \\ | & | & | & | & | & | & \end{array} \dots$$

- Observables can be calculated by fixed-point relations
- No finite-size effects: only finite entanglement (bond dimension)
- How to get information about excitations? (Dynamics vs statics)

# Single-mode approximation

Given a creation operator  $\hat{a}^\dagger$ , get an approximate excited state

$$|\Phi_k\rangle = \sum_n e^{ikn} \hat{a}_n^\dagger |\Psi\rangle = \sum_n e^{ikn} \dots \overset{A}{\circ} \overset{A}{\circ} \overset{\hat{a}^\dagger A}{\bullet} \overset{A}{\circ} \overset{A}{\circ} \dots$$

$n$

MPO:

$$\hat{W}_k = \begin{pmatrix} e^{ik} \hat{I} & \hat{a}^\dagger \\ 0 & \hat{I} \end{pmatrix}$$

Can write as a block-triangular MPS

$$|\Phi_k\rangle = \dots \overset{\mathcal{A}_k}{\circ} \overset{\mathcal{A}_k}{\circ} \overset{\mathcal{A}_k}{\circ} \overset{\mathcal{A}_k}{\circ} \overset{\mathcal{A}_k}{\circ} \dots, \quad \mathcal{A}_k = \begin{pmatrix} e^{ik} A & \hat{a}^\dagger A \\ 0 & A \end{pmatrix}$$

## MPS excitation ansatz

$$|\Phi_k[B]\rangle = \sum_n e^{ikn} \dots \text{---} \overset{A}{\bigcirc} \text{---} \overset{A}{\bigcirc} \text{---} \overset{B}{\bullet} \text{---} \overset{\tilde{A}}{\bigcirc} \text{---} \overset{\tilde{A}}{\bigcirc} \text{---} \dots$$

$$= \dots \text{---} \overset{\mathcal{A}}{\bigcirc} \text{---} \overset{\mathcal{A}}{\bigcirc} \text{---} \overset{\mathcal{A}}{\bigcirc} \text{---} \overset{\mathcal{A}}{\bigcirc} \text{---} \overset{\mathcal{A}}{\bigcirc} \text{---} \dots, \quad \mathcal{A} = \begin{pmatrix} e^{ik} A & B \\ 0 & \tilde{A} \end{pmatrix}$$

- Generalization of single-mode approximation  $B = \hat{a}^\dagger A$
- $B$  can be optimized for each  $k$
- Block-triangular structure: reminiscent of MPOs
- Reuse ground state data instead of starting from scratch
- Non-injective: needed to differentiate excitation from background

# Excitation ansatz fixed-point equations

$$E^{\omega\alpha\beta}(n+1) \left[ \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right] = E^{\omega'\alpha'\beta'}(n) \left[ \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right]$$

- Only need  $\omega' \leq \omega$ ,  $\alpha' \leq \alpha$ ,  $\beta' \leq \beta \Rightarrow$  solve recursively
- Same algorithm as iMPS, but with extra indices  $\alpha, \beta$

$$\mathcal{H}_{\text{eff}}^k B = E^{\omega\alpha 1} \left[ \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right]$$

JJO & McCulloch, Phys. Rev. Research **7**, 023018 (2025)

Compare Michel & McCulloch, arXiv:1008.4667

# Multi-site windows

For broader excitations, increase size of window:

$$|\Phi_k[\mathbf{B}]\rangle = \sum_n e^{ikn} \dots \text{---} \underset{n}{\begin{array}{c} A \quad A \quad B_1 \quad B_2 \quad B_3 \quad \tilde{A} \quad \tilde{A} \\ \circ \quad \circ \quad \bullet \quad \bullet \quad \bullet \quad \circ \quad \circ \\ | \quad | \quad | \quad | \quad | \quad | \quad | \end{array}} \dots,$$

$$\mathcal{A} = \begin{pmatrix} e^{ik} A & B_1 & & \\ & 0 & B_2 & \\ & & 0 & B_3 \\ & & & \tilde{A} \end{pmatrix}$$

Optimize one tensor at a time (EA DMRG):

$$\mathcal{H}_{\text{eff}}^{n,k} B_n = E^{\omega\alpha n} \left[ \text{Diagram} \right] F^{\omega'\alpha'(n+1)}$$

Compare Vanderstraeten et al., Phys. Rev. B **101**, 115138 (2020)

# Expectation values

Expectation values are polynomials in system size  $L$ :

$$\langle \hat{H} \rangle = \frac{\langle \Phi | \hat{H} | \Phi \rangle}{\langle \Phi | \Phi \rangle} = \frac{EL^2 + \Delta L}{L} = EL + \Delta$$

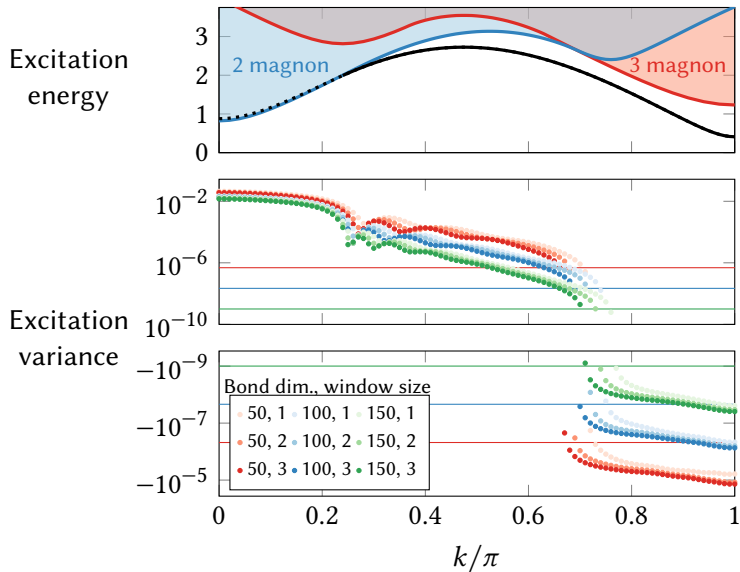
Can calculate the variance:

$$\langle \hat{H}^2 \rangle - \langle \hat{H} \rangle^2 = \sigma_E^2 L + \sigma_\Delta^2$$

Since we only have an approximation to the GS:

- $\Delta$  can be less than the actual value
- $\sigma_\Delta^2$  can be negative
- Bound by error in GS

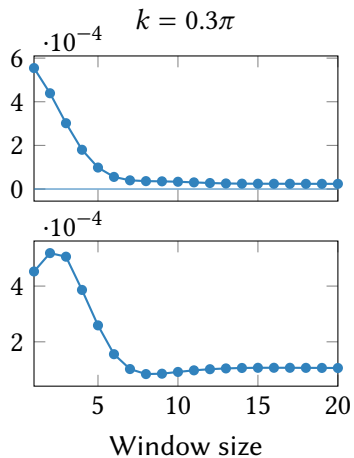
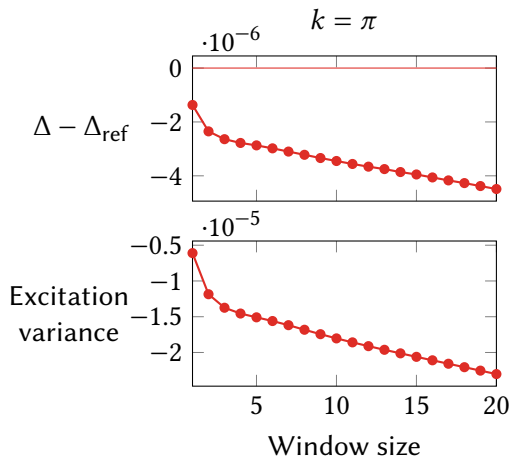
# Spin-1 Heisenberg model





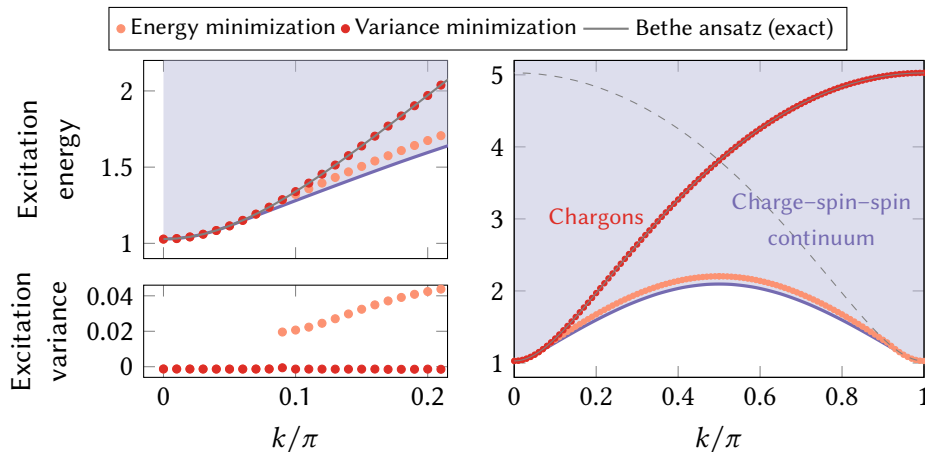
# Spin-1 Heisenberg model

Tuning window size (bond dim. 50)



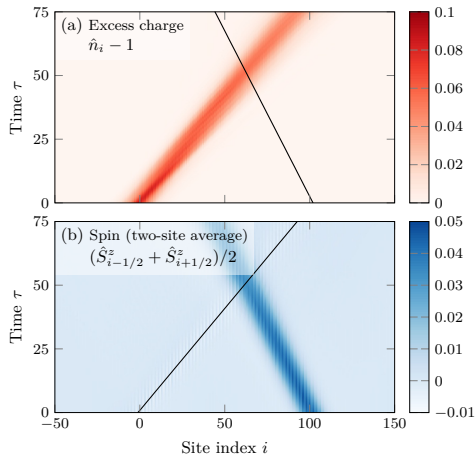
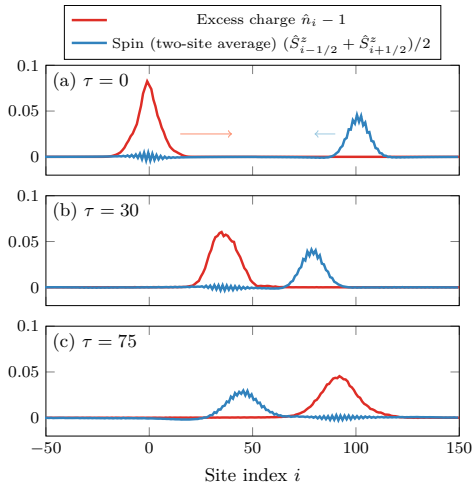
# Minimizing energy variance

Chargons in the Hubbard model,  $U/t = 5$



Compare Zauner–Stauber et al., Phys. Rev. B **97**, 235155 (2018)

# Wave packet preparation: spin-charge separation



# Multi-particle states

$$\begin{aligned}
 |Y_{k_1, k_2}\rangle = \sum_n & \left[ q_1 \sum_{m>n} e^{i(k_1 n + k_2 m)} \dots \text{---} \underset{n}{\overset{B}{\bullet}} \text{---} \dots \text{---} \underset{m}{\overset{C}{\bullet}} \text{---} \dots \right. \\
 & + q_2 \sum_{m<n} e^{i(k_2 m + k_1 n)} \dots \text{---} \underset{m}{\overset{C}{\bullet}} \text{---} \dots \text{---} \underset{n}{\overset{B}{\bullet}} \text{---} \dots \\
 & \left. + e^{i(k_1 + k_2)n} \dots \text{---} \underset{n}{\overset{\text{'Correction window' } W}{\bullet \bullet \bullet}} \text{---} \dots \right],
 \end{aligned}$$

$$\mathcal{A} = \begin{pmatrix} e^{i(k_1 + k_2)A} & q_1 B & q_2 C & W \\ & e^{ik_2 A} & & C \\ & & e^{ik_1 A} & B \\ & & & A \end{pmatrix}$$

Generalizes to  $\geq 3$  excitations

Vanderstraeten et al., Phys. Rev. B **92**, 125136 (2015)

# Conclusion

- Excitation ansatz wave functions are triangular MPSs
- Can calculate fixed-point relations recursively
- Works for arbitrary operators; variance minimization
- Directly handles extensions to the EA (multi-site windows, multi-particle states)