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# How to get down to microkelvin

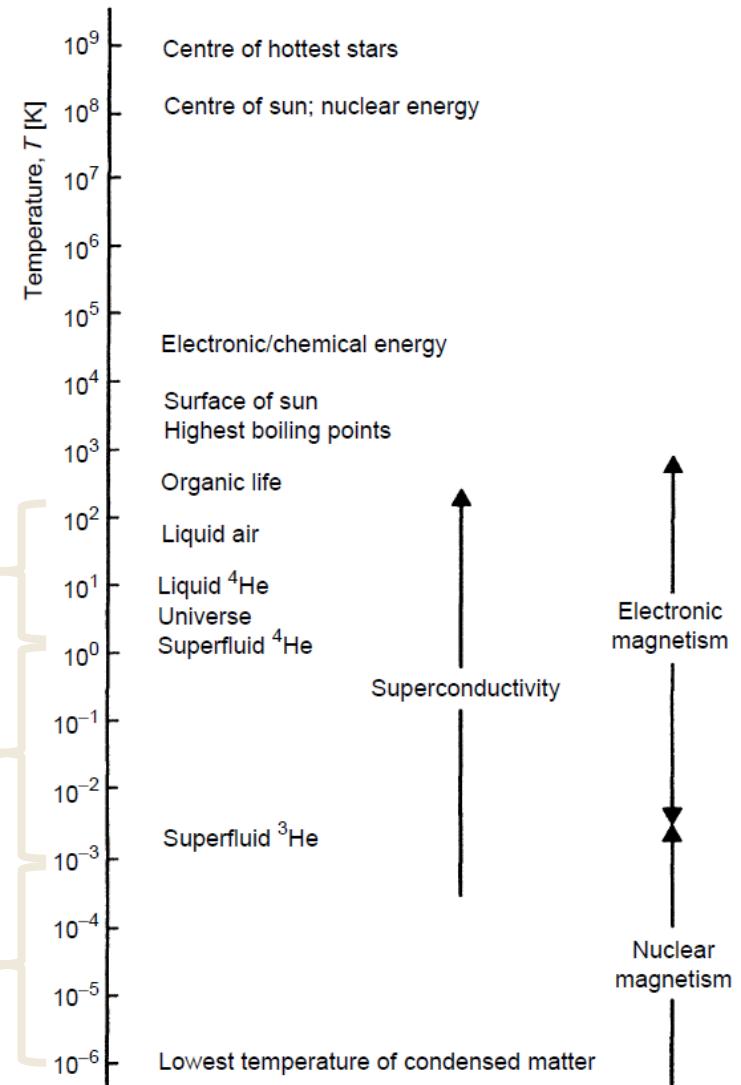
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Matthäus Krantz

HighRR Seminar, 17.05.2017

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- History
- Liquification of Cryoliquids
- $^4\text{He}$  Cryostat (wet)
- Pulse Tube Cooler (dry)
- $^3\text{He}/^4\text{He}$  Dilution Refrigerator
- Adiabatic Demagnetization
- Nuclear Demagnetization



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# History

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# Historic Introduction

1845: Faraday successfully liquefied most known gases, except for six, which became known as the permanent gases: oxygen, hydrogen, nitrogen, carbon monoxide, methane and nitric oxide (helium not found as gas yet).

1877: Oxygen and nitrogen were liquefied by Louis Cailletet (France) and Raoul Pictet (Switzerland), 80 K reached.

1898: Liquefaction and solidification of hydrogen by James Dewar (13 K reached), Cambridge.

1908: Last permanent gas, helium, was liquefied by Kamerlingh Onnes, Leiden. He reached 0.83 K in 1922.



Michael Faraday (1791 – 1867)



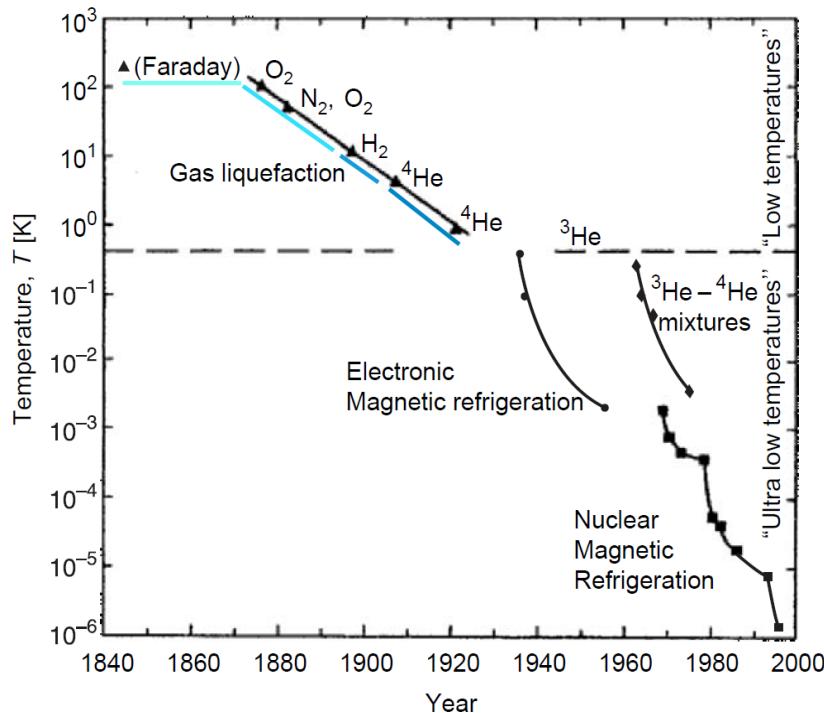
Raoul Pictet  
(1846 – 1929)



James Dewar  
(1842 – 1923)



Heike Kamerlingh Onnes  
(1853 – 1926)



# Historic Introduction

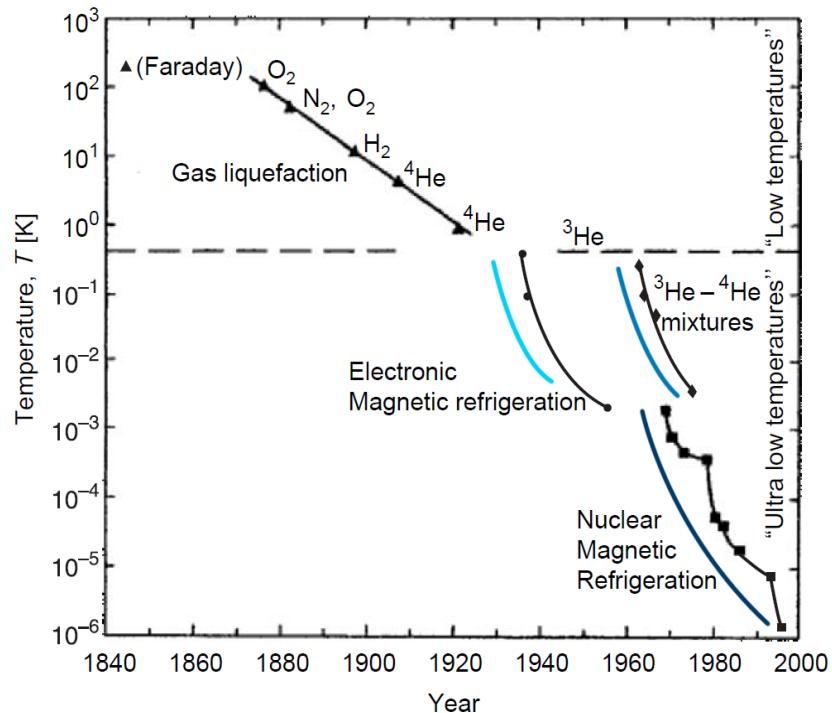
1930's: Adiabatic Demagnetization opens up lower millikelvin range for experiments  
(single shot cooling)

Today:  $T = 5 - 30 \text{ mK}$  (record: 1 mK)

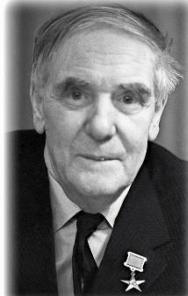
1960's: Development of Dilution Refrigerators which allow continuous cooling to millikelvin.  
Today:  $T = 5 - 20 \text{ mK}$  (record: 2 mK)

1970's: Nuclear Demagnetization makes microkelvin temperatures available for experiments  
(single shot cooling)

Today:  $T < 100 \mu\text{K}$  (record 1.5  $\mu\text{K}$ )



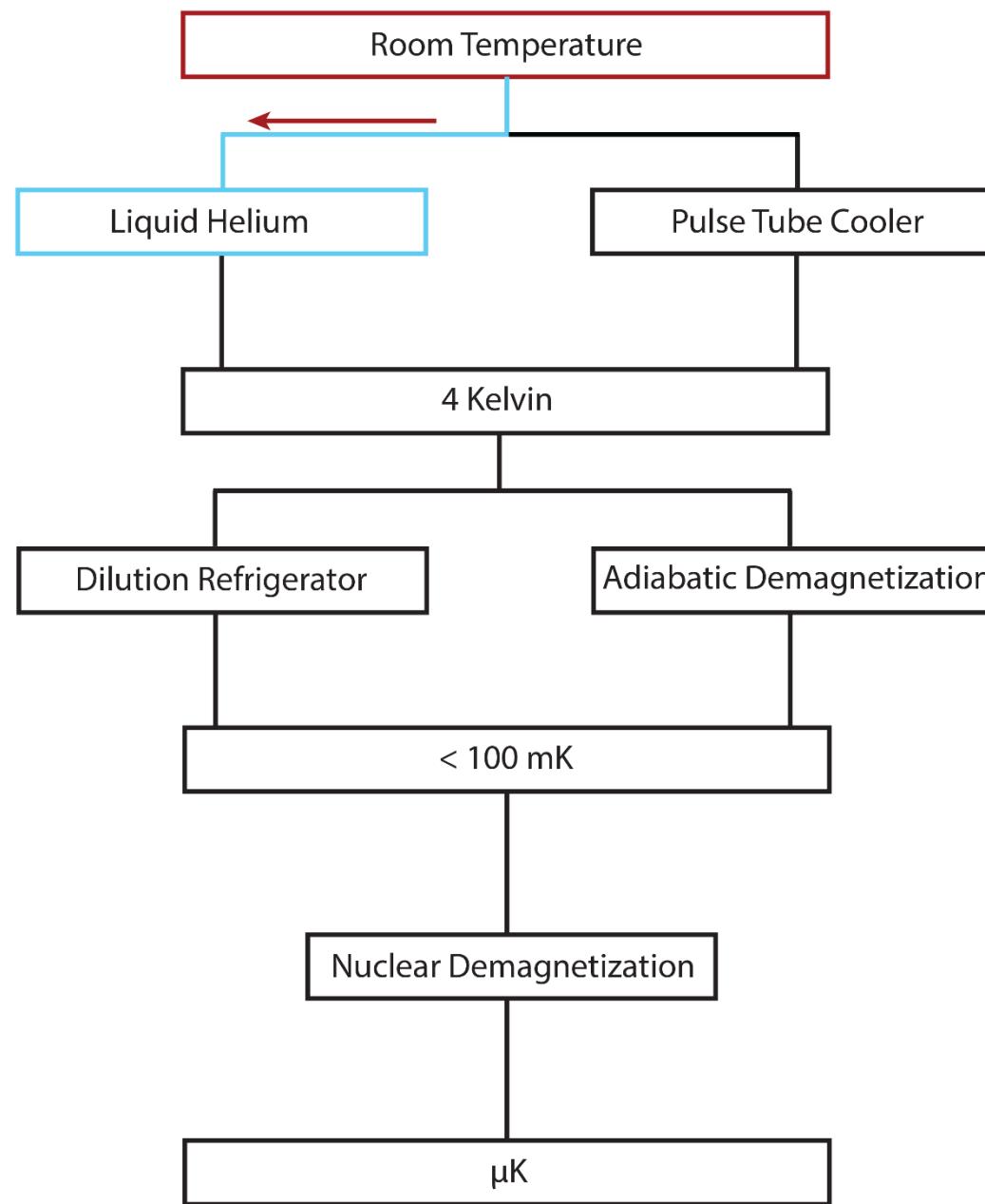
Peter Debye  
(1884 – 1966)

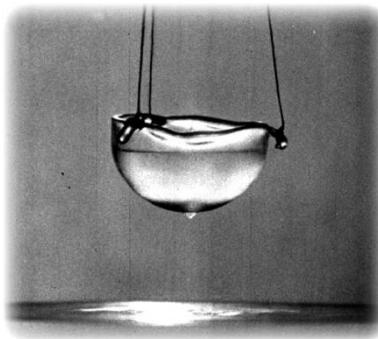


Pjotr L. Kapitza  
(1894 – 1994)



Heinz London  
(1907 – 1970)

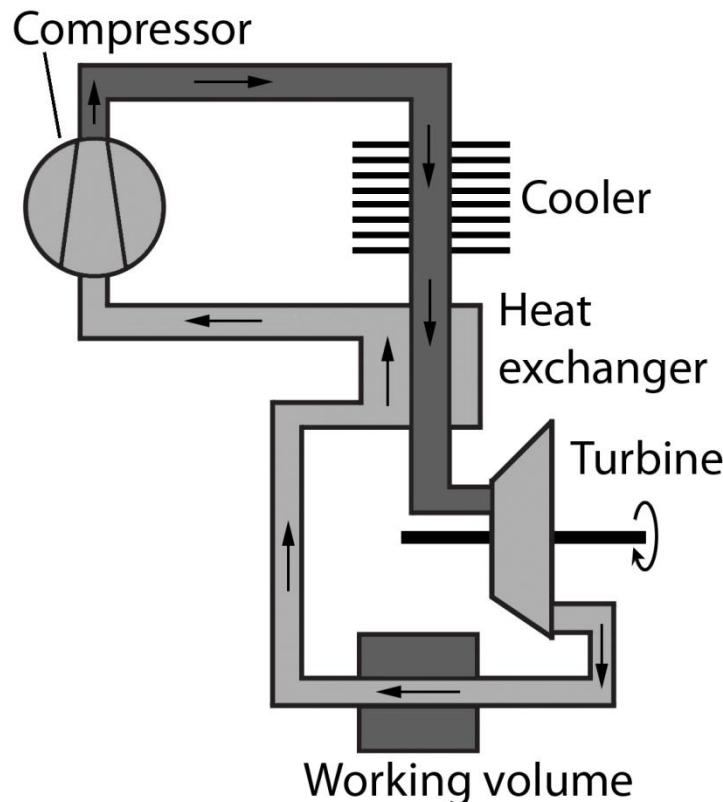




# Liquification of Cryoliquids



# Expansion Cooler



- Compression at room temperature
- Removal of excess heat at cooler
- Further cooling at heat exchanger
- Adiabatic expansion at turbine:

$$T_2 = T_1 \left( \frac{p_1}{p_2} \right)^{\frac{1-\gamma}{\gamma}} \quad \text{with} \quad \gamma = c_p/c_v$$

- Expansion: gas performs work (ideal case)

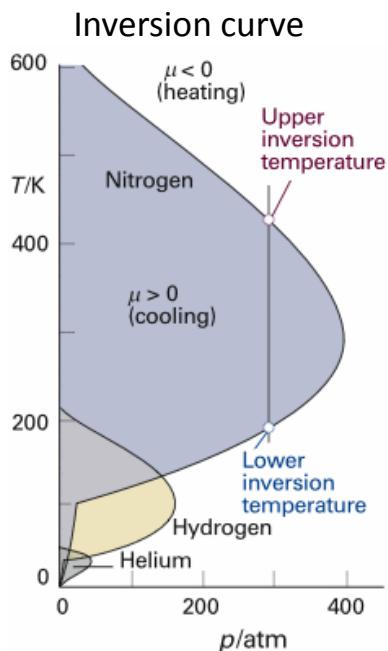
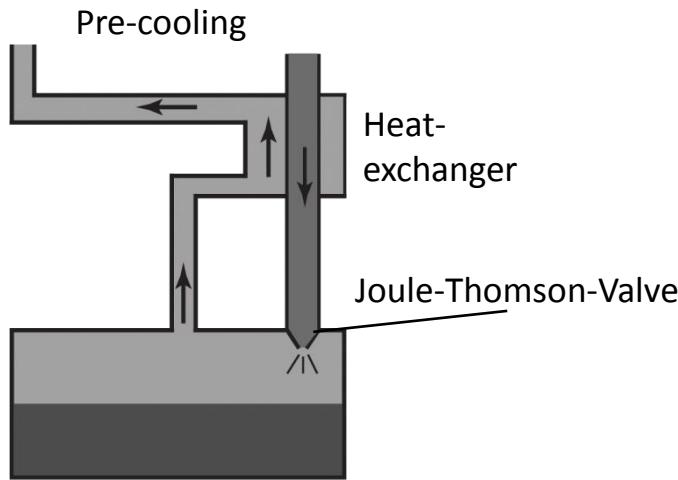
$$W = H_1 - H_2 = (U_1 + p_1 V_1) - (U_2 + p_2 V_2)$$

- Ideal gas:

$$W = \frac{5}{2} N k_B (T_1 - T_2)$$

- Gas performs work at the turbine, going up to 180.000 rpm

# Joule-Thomson-Expansion



- Expansion at nozzle, gas has to perform work against internal forces

- Change of internal energy during adiabatic expansion:

$$\Delta U = U_2 - U_1 = p_1 V_1 - p_2 V_2$$

- Joule-Thomson-Coefficient:

$$\mu_{JT} = \left( \frac{\partial T}{\partial P} \right)_H = \frac{V}{C_p} (\alpha T - 1)$$

- Cooling if  $\mu_{JT} > 0$

- Heating if  $\mu_{JT} < 0$

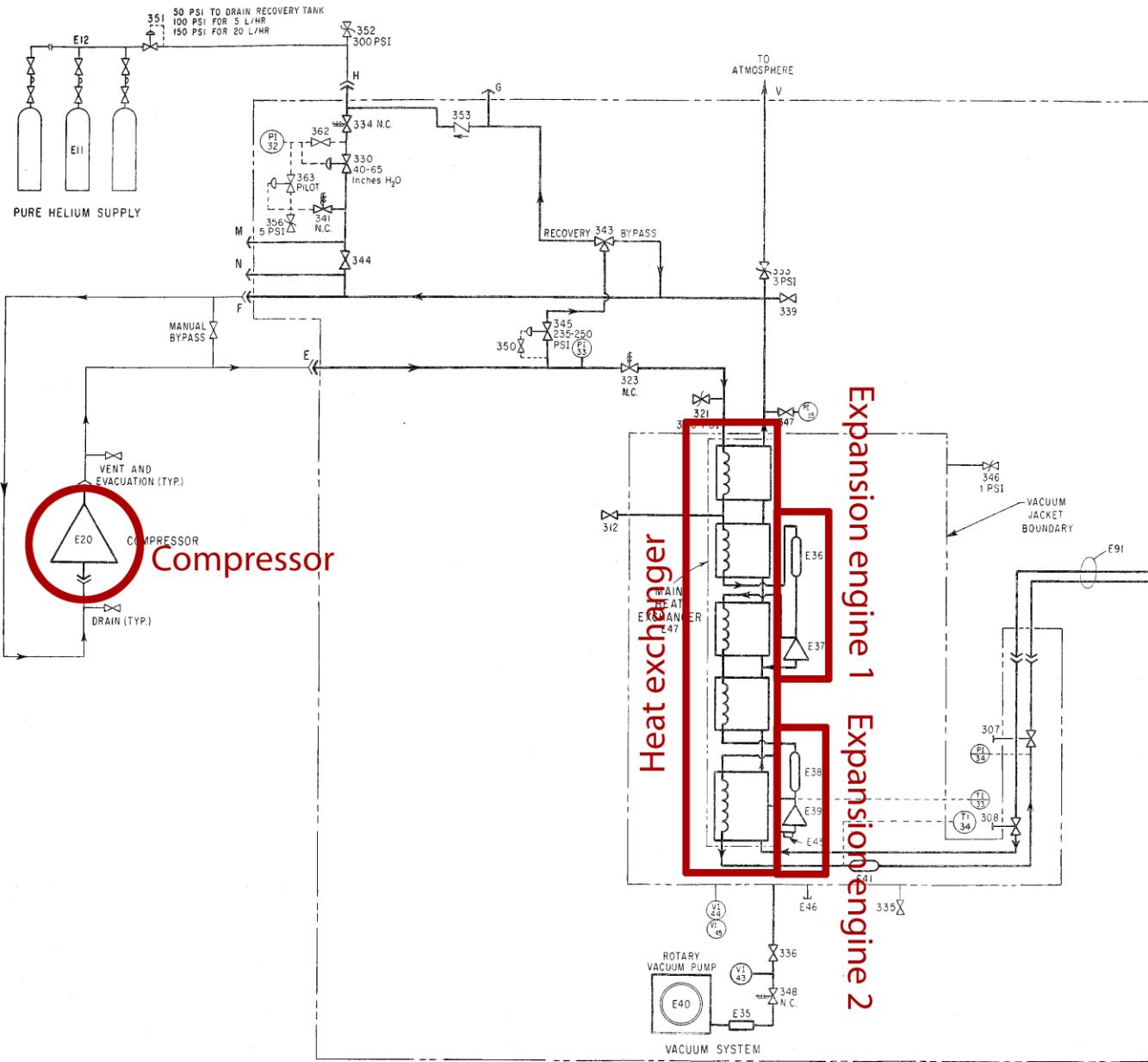
➤  $T_{inv,O_2} = 764 \text{ K}$

➤  $T_{inv,N_2} = 621 \text{ K}$

➤  $T_{inv,{}^4\text{He}} = 40 \text{ K}$

➤  $T_{inv,{}^3\text{He}} = 23 \text{ K}$

# KIP He-Liquifier



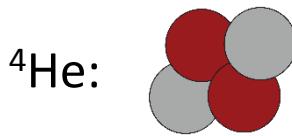
## EQUIPMENT LIST

- E11-PURE HELIUM GAS SUPPLY  
E12-PURE HELIUM GAS MANIFOLD  
E20-COMPRESSOR  
E35-VACUUM HOSE  
E36,E38-CHARCOAL FILTER  
E37,E39-EXPANSION ENGINE  
E40-ROTARY VACUUM PUMP  
E41-J-T FILTER  
E45-GETTER  
E46-VACUUM COUPLING  
E47-MAIN HEAT EXCHANGER  
E91-DELIVERY TUBE

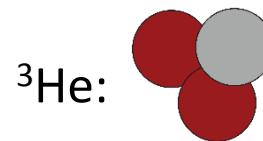
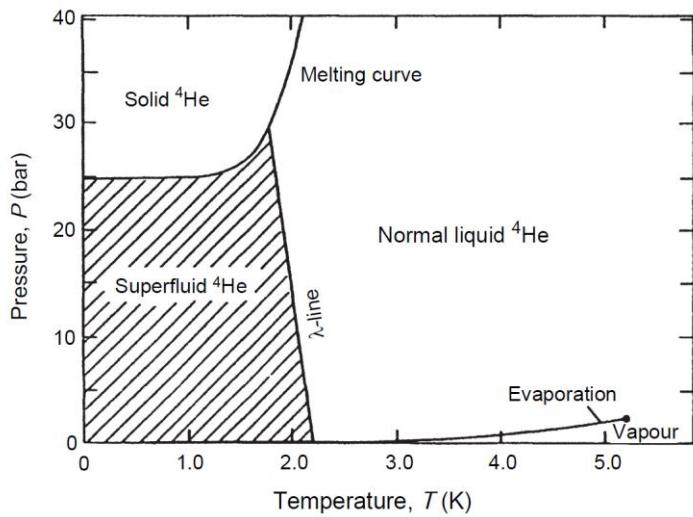
CABINET  
BOUNDARY

## JT valve

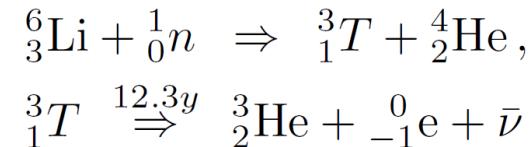
# Cryoliquid Helium



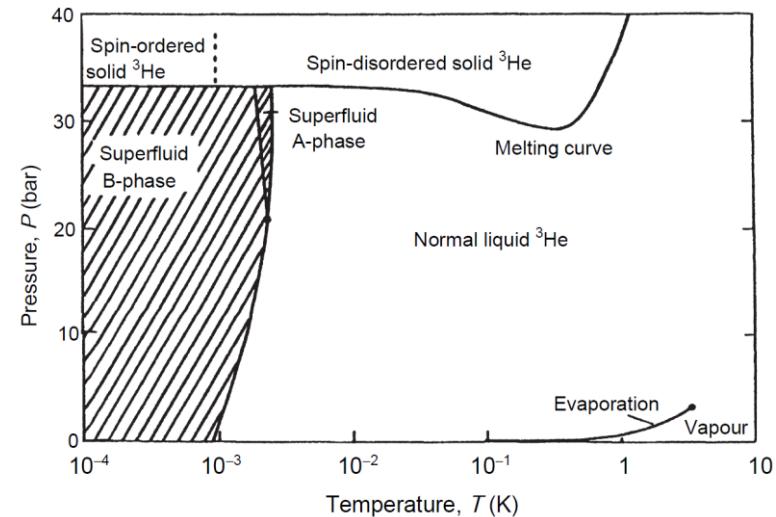
- Natural occurrence as gas: helium-rich natural gas reservoirs enriched by alpha decay
- Nuclear spin  $I = 0$  (boson)



- Byproduct of tritium fabrication in a nuclear reactor:

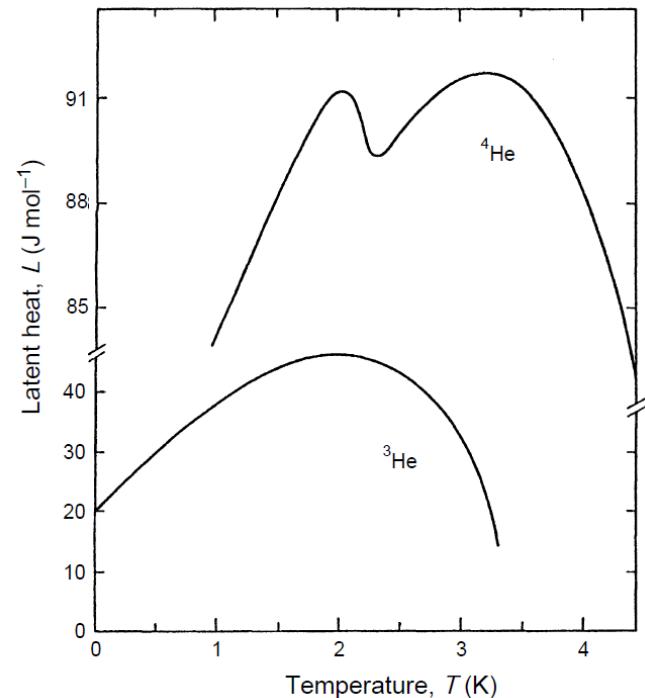
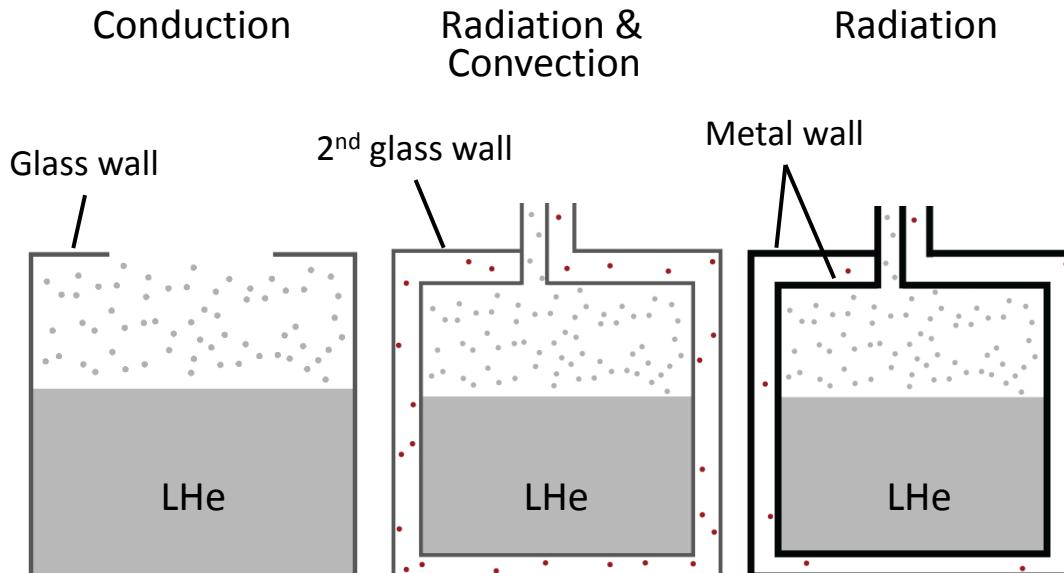


- Nuclear spin  $I = 1/2$  (fermion)



# Latent Heat of Evaporation

- Latent heat  $L$ : energy needed for the phase transition from liquid to gas
- $L$  of helium very small compared to other liquids
- Very good shielding against external parasitic heat input needed
- Energy input from:



$L$  at boiling point:

$$L_{{}^3\text{He}} = 16 \text{ J mol}^{-1}$$

$$L_{{}^4\text{He}} = 84 \text{ J mol}^{-1}$$

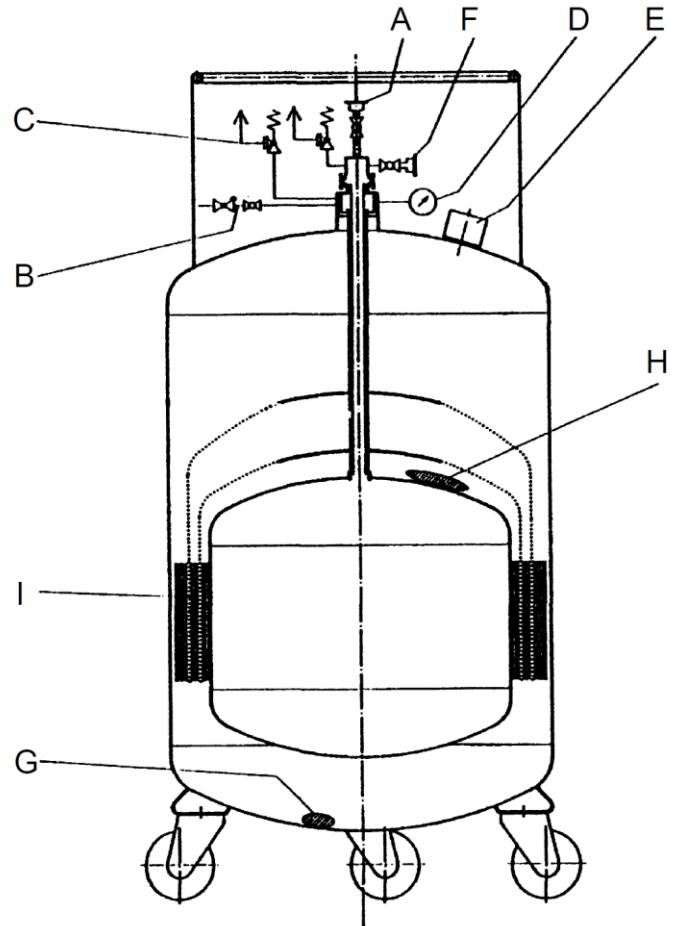
$$L_{\text{H}_2} = 449 \text{ J mol}^{-1}$$

$$L_{\text{N}_2} = 2792 \text{ J mol}^{-1}$$

$$L_{\text{Ar}} = 6447 \text{ J mol}^{-1}$$

# Transport of liquid Helium

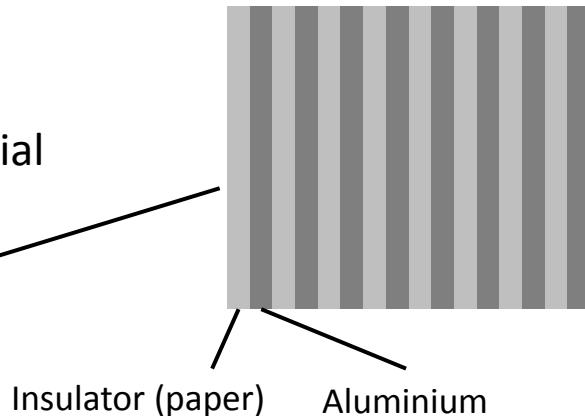
- A: Connection for transfer tube
- B: Overflow valve
- C: Safety valve
- D: Manometer
- E: Vakuum and safety valves
- F: Gas valve
- G: Getter Material
- H: Adsorbent Material
- I: Superinsulation



Rate of evaporation ~ 1 %/day

# Transport of liquid Helium

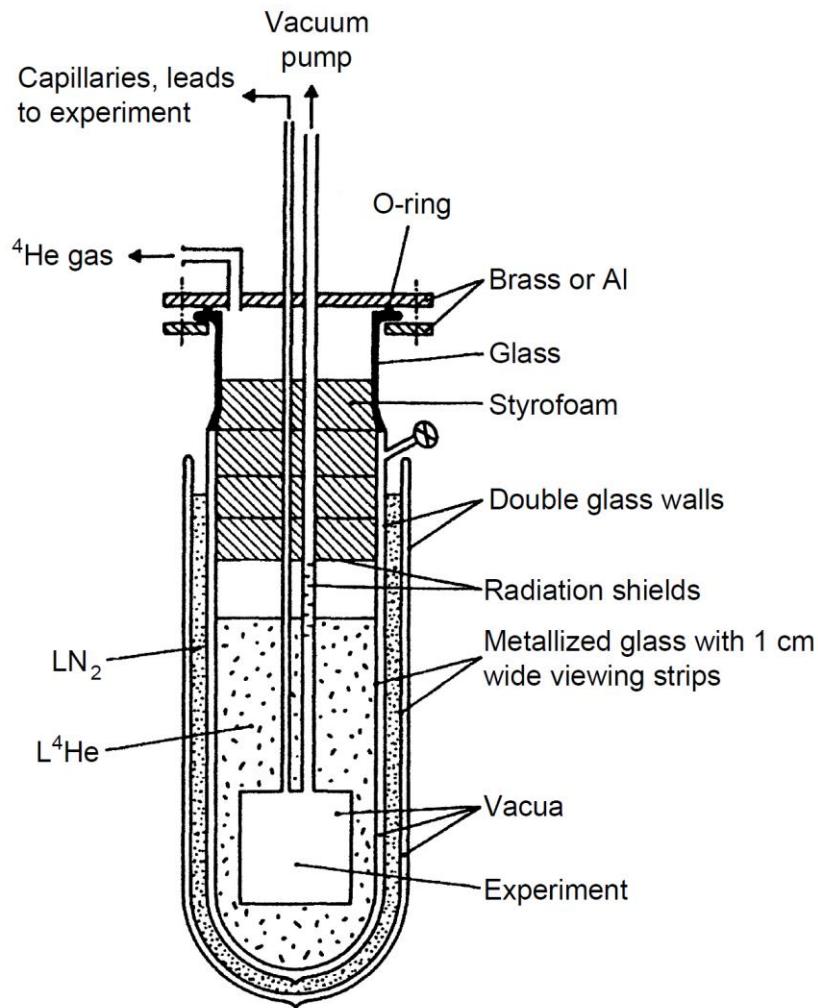
- A: Connection for transfer tube
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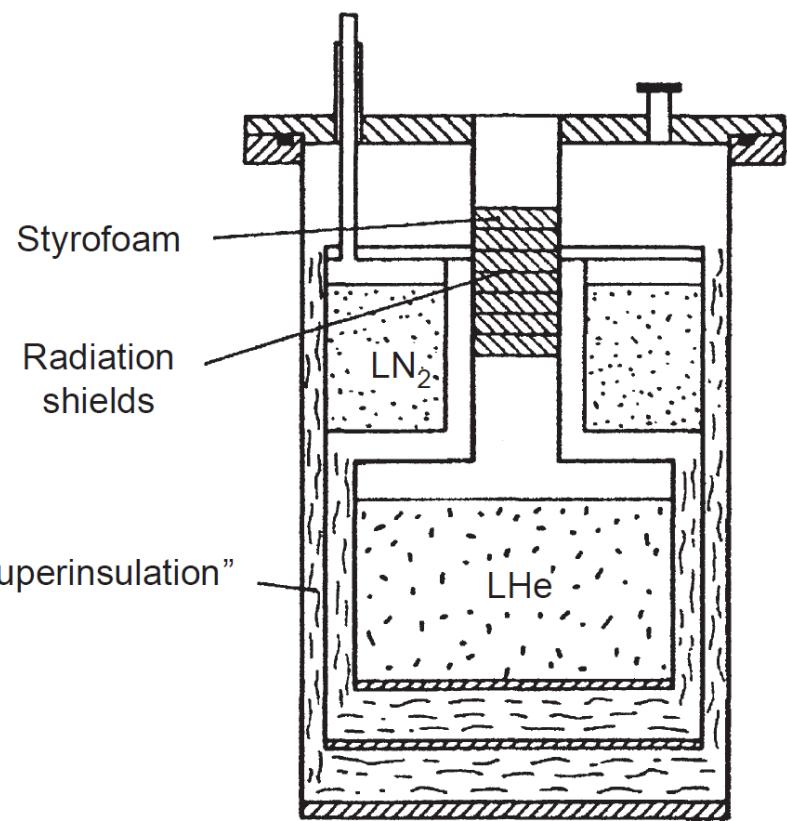
# $^4\text{He}$ Cryostat

# $^4\text{He}$ Bath Cryostat

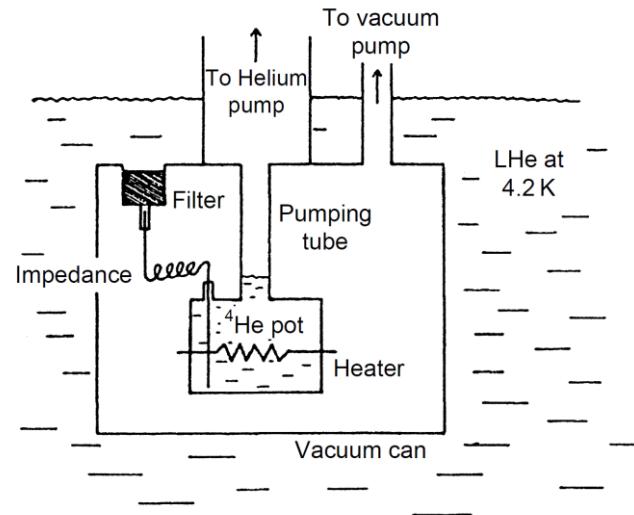
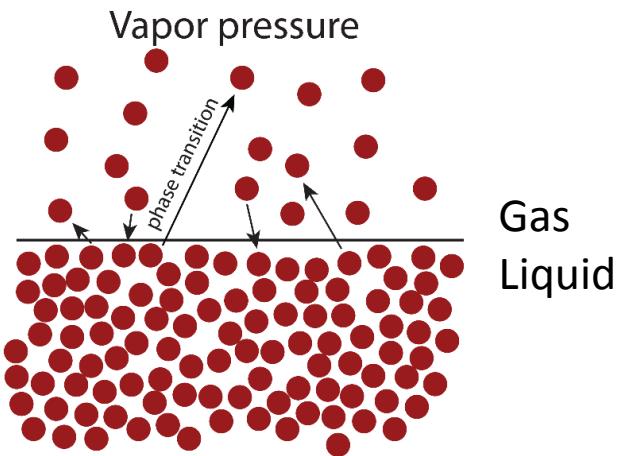
“Historic”:



Modern:



# Evaporation Cryostat



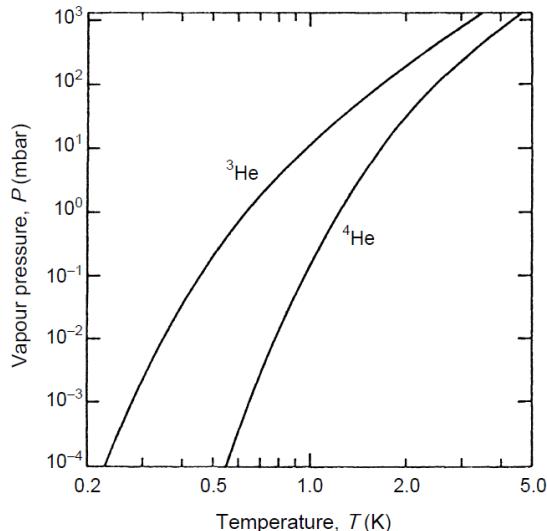
Vapor pressure:  $P_{\text{vap}} \propto e^{-L/RT}$

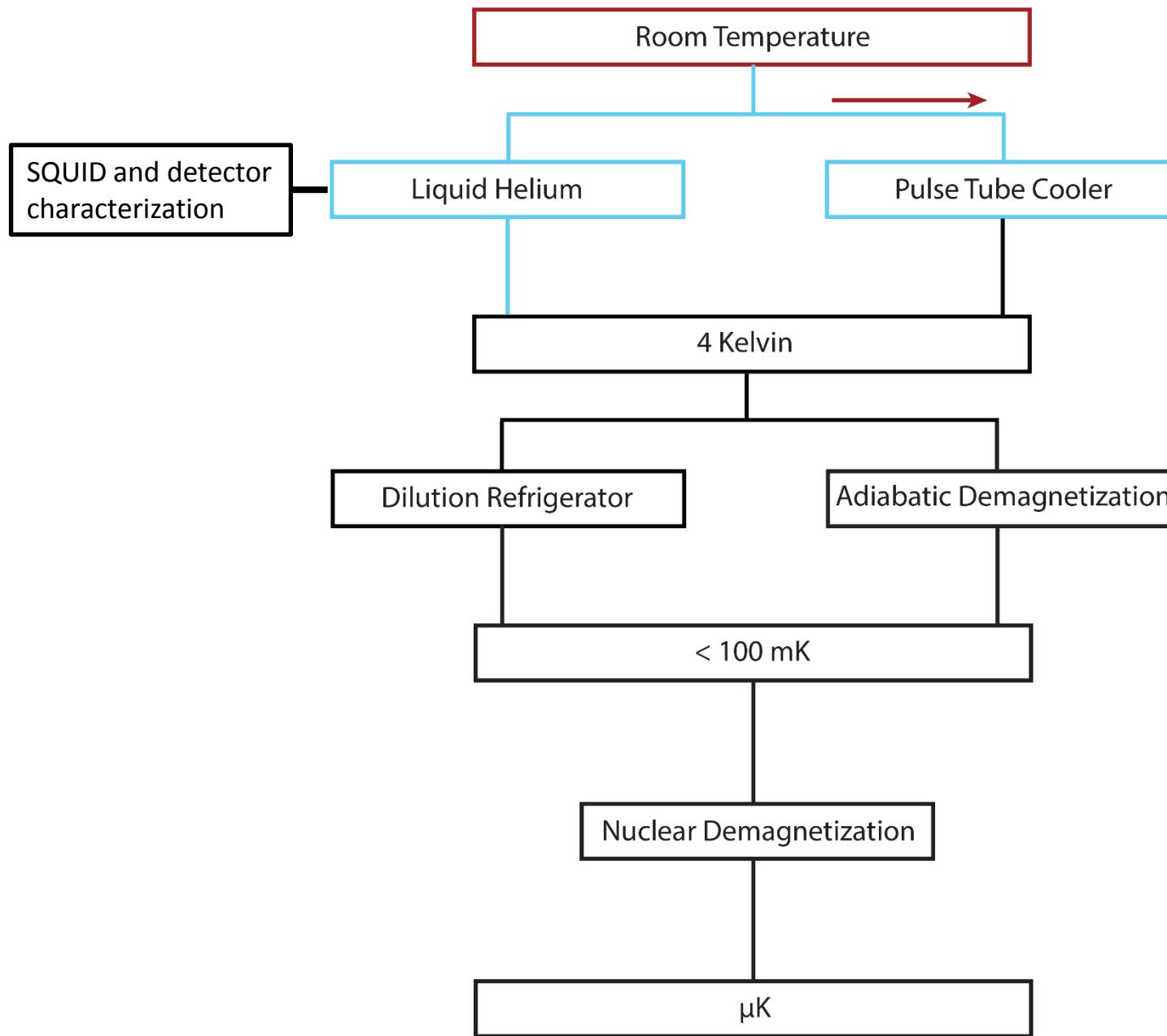
Cooling power:  $\dot{Q} = \dot{n}_g L$

Typical achieved temperatures:

${}^4\text{He}$ : 1.3 K

${}^3\text{He}$ : 0.3 K

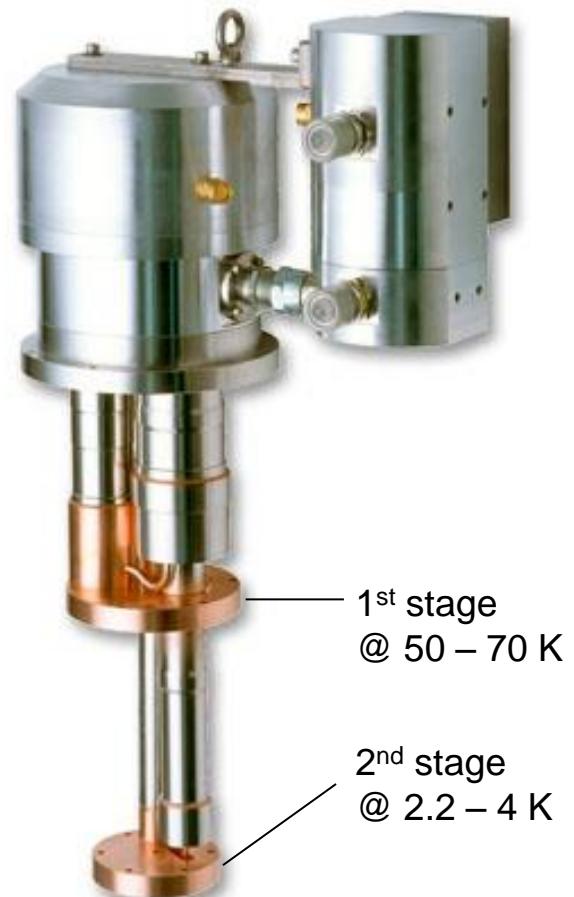




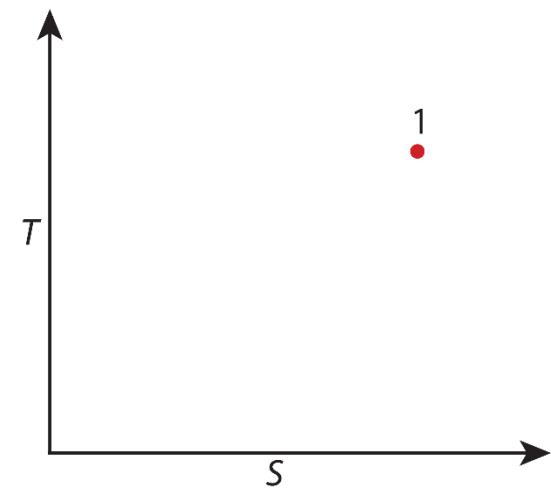
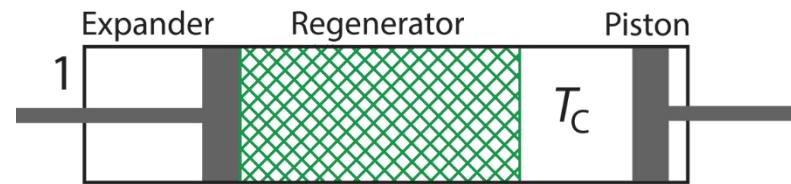
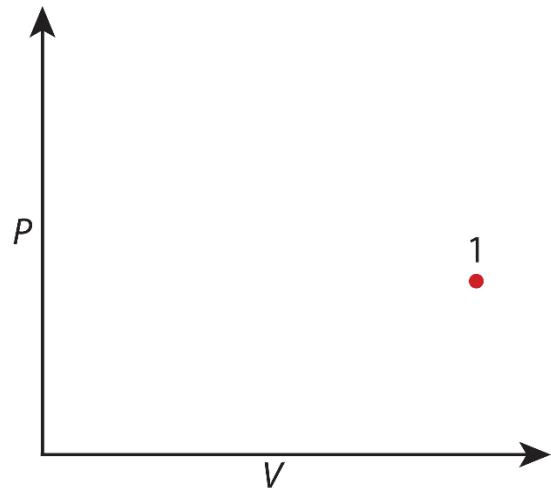
# Pulse Tube Cooler



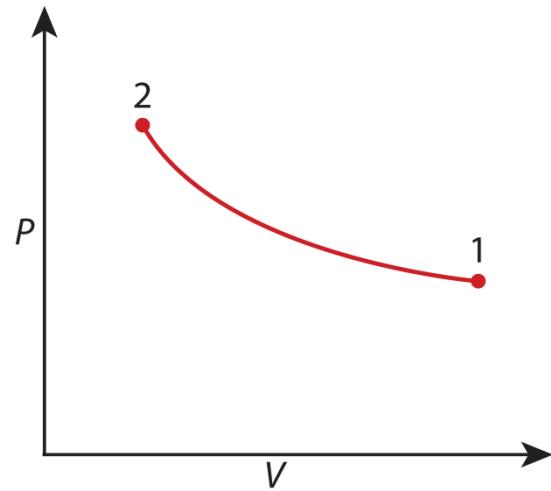
- Closed cycle refrigeration
- From 300 K to 2.2 K without cryoliquids
- Working principle: Compression, expansion and displacement of gas
- Helium as working gas under high pressure (18 -22 bar)
- Power consumption: 2 – 8 kW



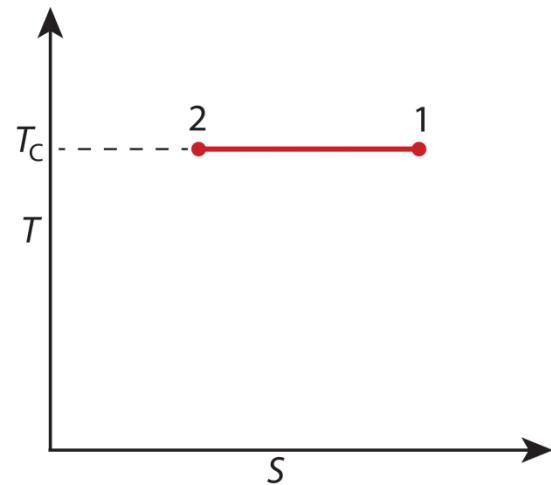
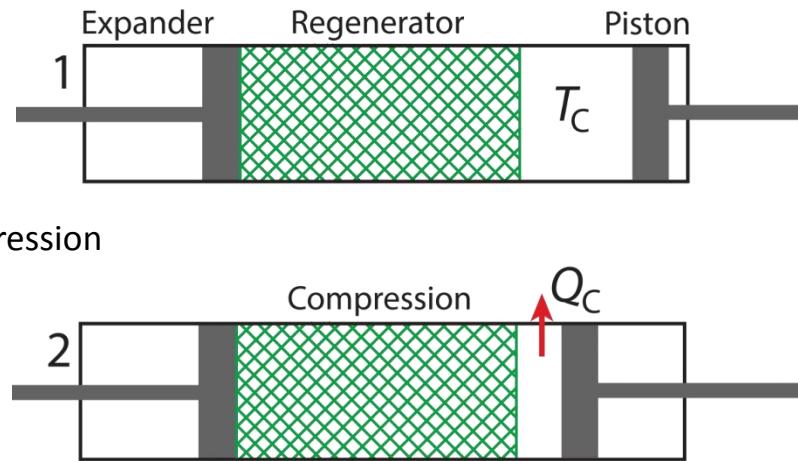
# Ideal Stirling Cycle



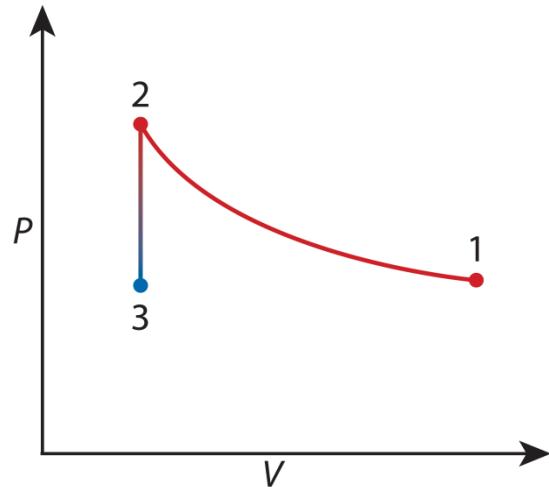
# Ideal Stirling Cycle



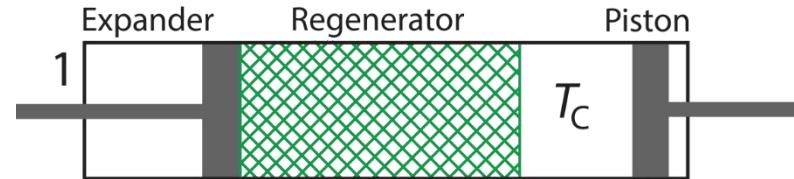
Isothermal compression



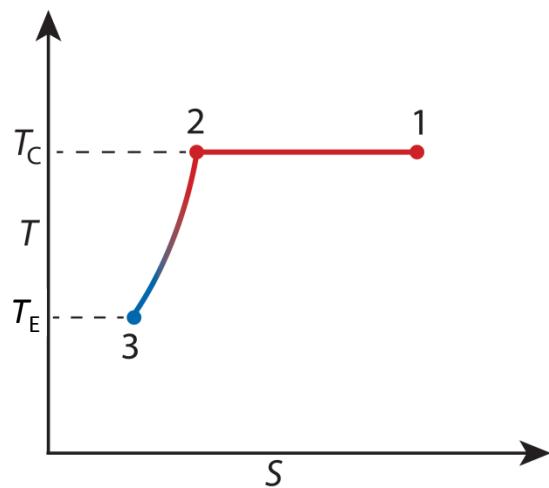
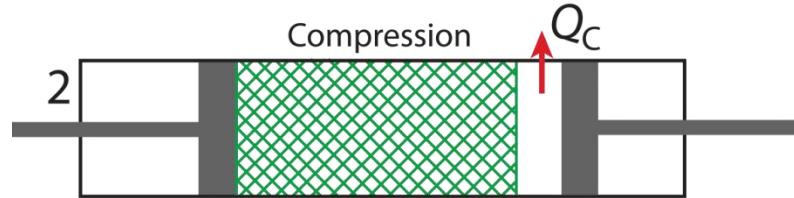
# Ideal Stirling Cycle



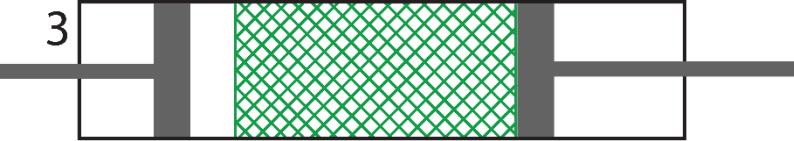
Isothermal compression



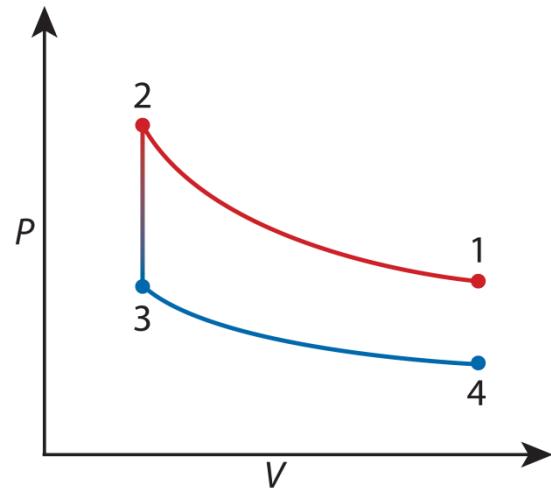
Isochoric precooling



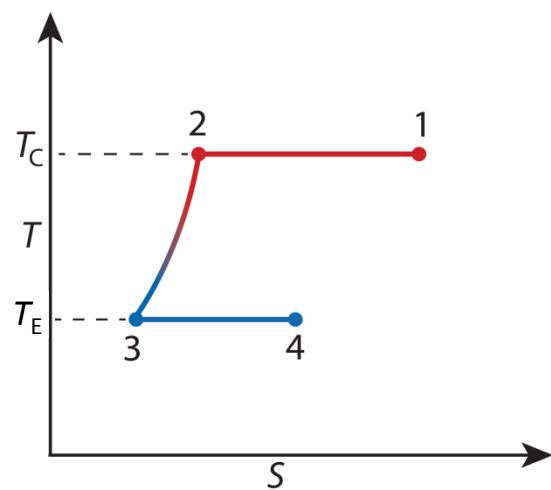
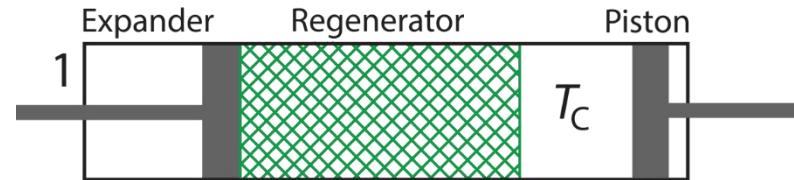
Regenerative Cooling



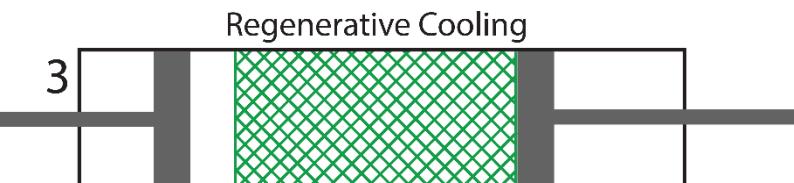
# Ideal Stirling Cycle



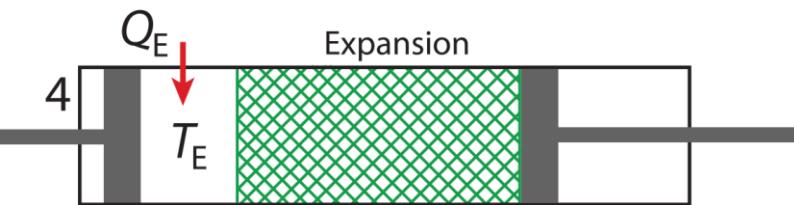
Isothermal compression



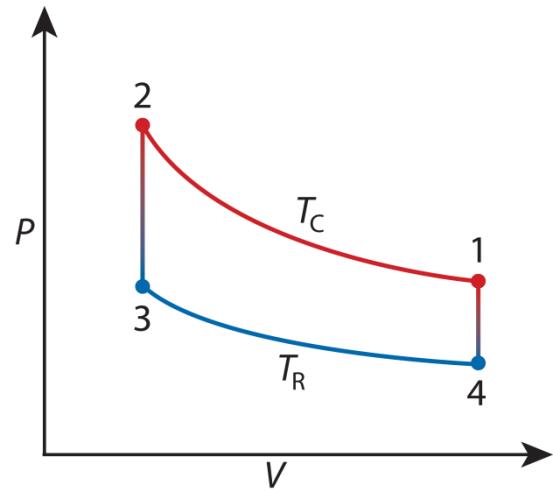
Isochoric precooling



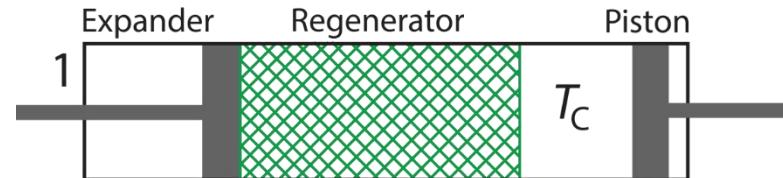
Isothermal expansion



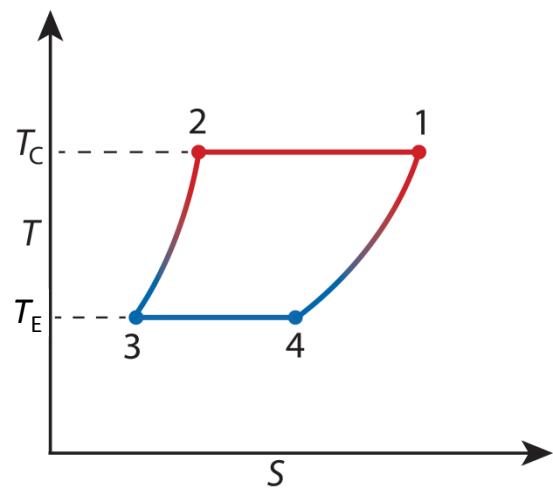
# Ideal Stirling Cycle



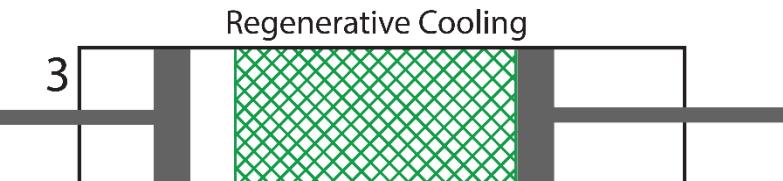
Isothermal compression



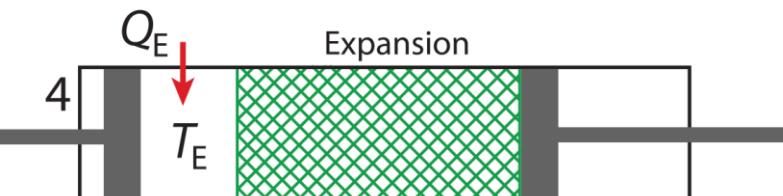
Isochoric precooling



Isothermal expansion

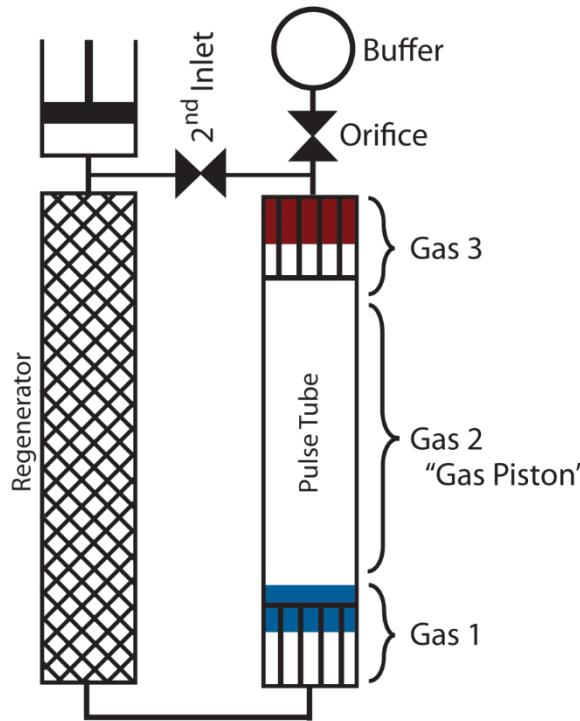


Isochoric reheating

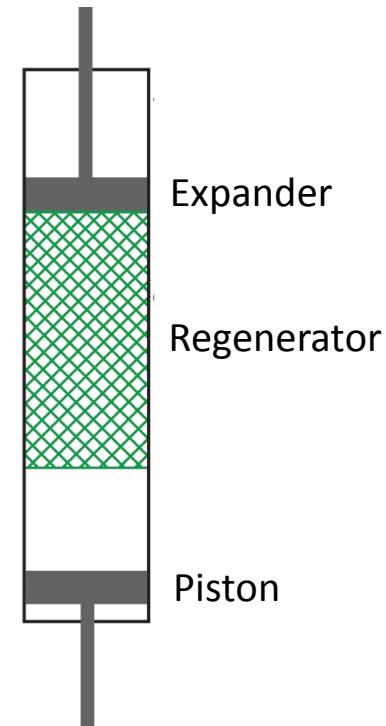


# Pulse Tube Cooler vs. Stirling

Stirling-type PTC



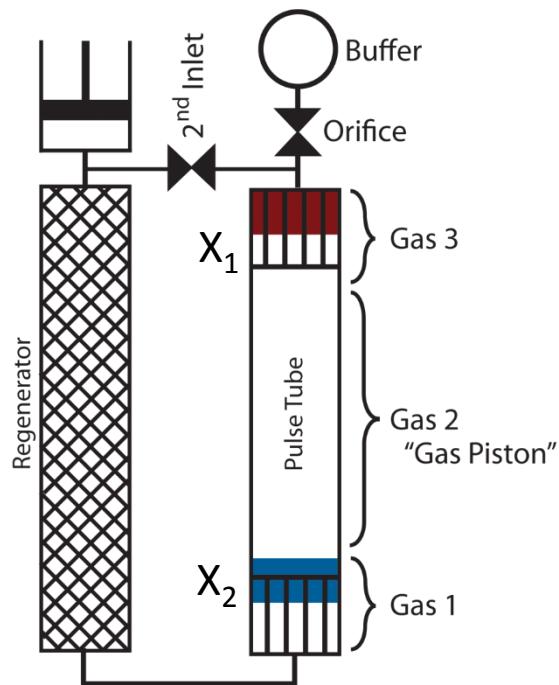
Stirling



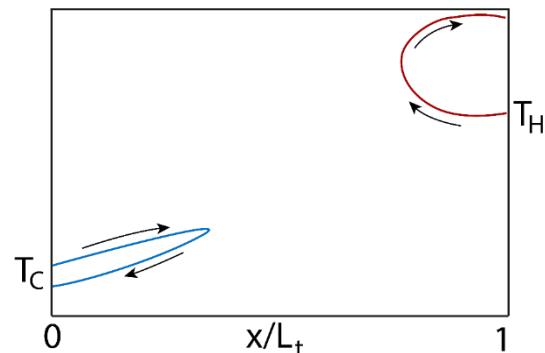
## Advantages of PTC:

- No moving parts inside the cryostat (less vibrations)
- Compressor unit and tube system can be placed far away and optimized independently

# Pulse Tube Cooler



Movement of gas inside pulse tube  
cold end    hot end



## Expansion:

- Gas 3 from buffer moves into pulse tube at  $X_1$  with temperature  $T_H$
- Gas 1 at cold end expands and moves back into the regenerator with temperature  $T < T_C$

## Compression:

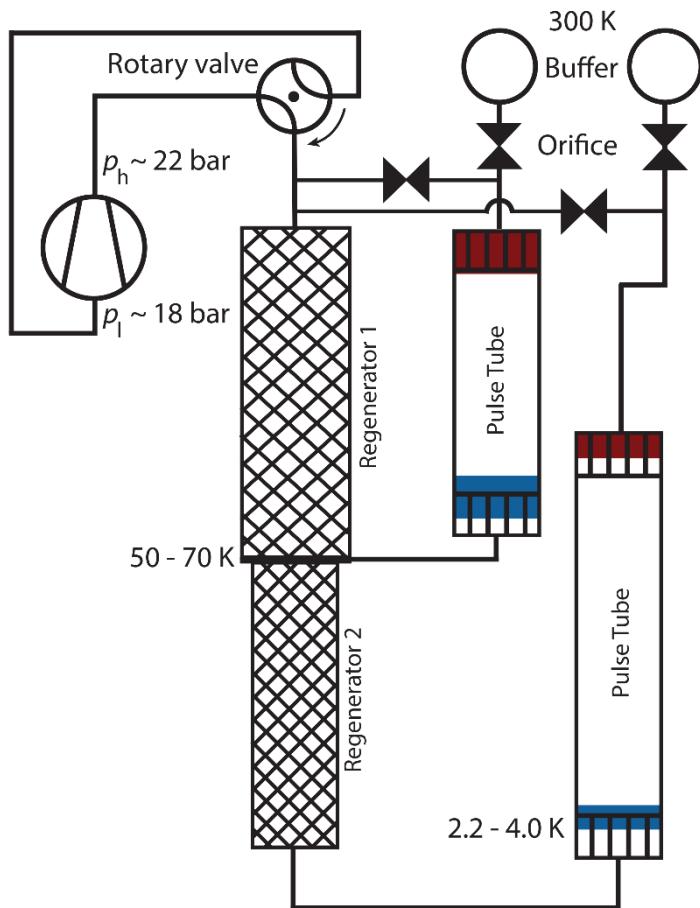
- Gas 3 at  $X_1$  is compressed and leaves the pulse tube with temperature  $T > T_H$
- Gas 1 in the regenerator is compressed and moves into  $X_2$  with  $T_C$

Gas 2 functions as a displacer and insulator for Gas 1 and Gas 3 during expansion and compression

# Two-Stage Pulse Tube Cooler

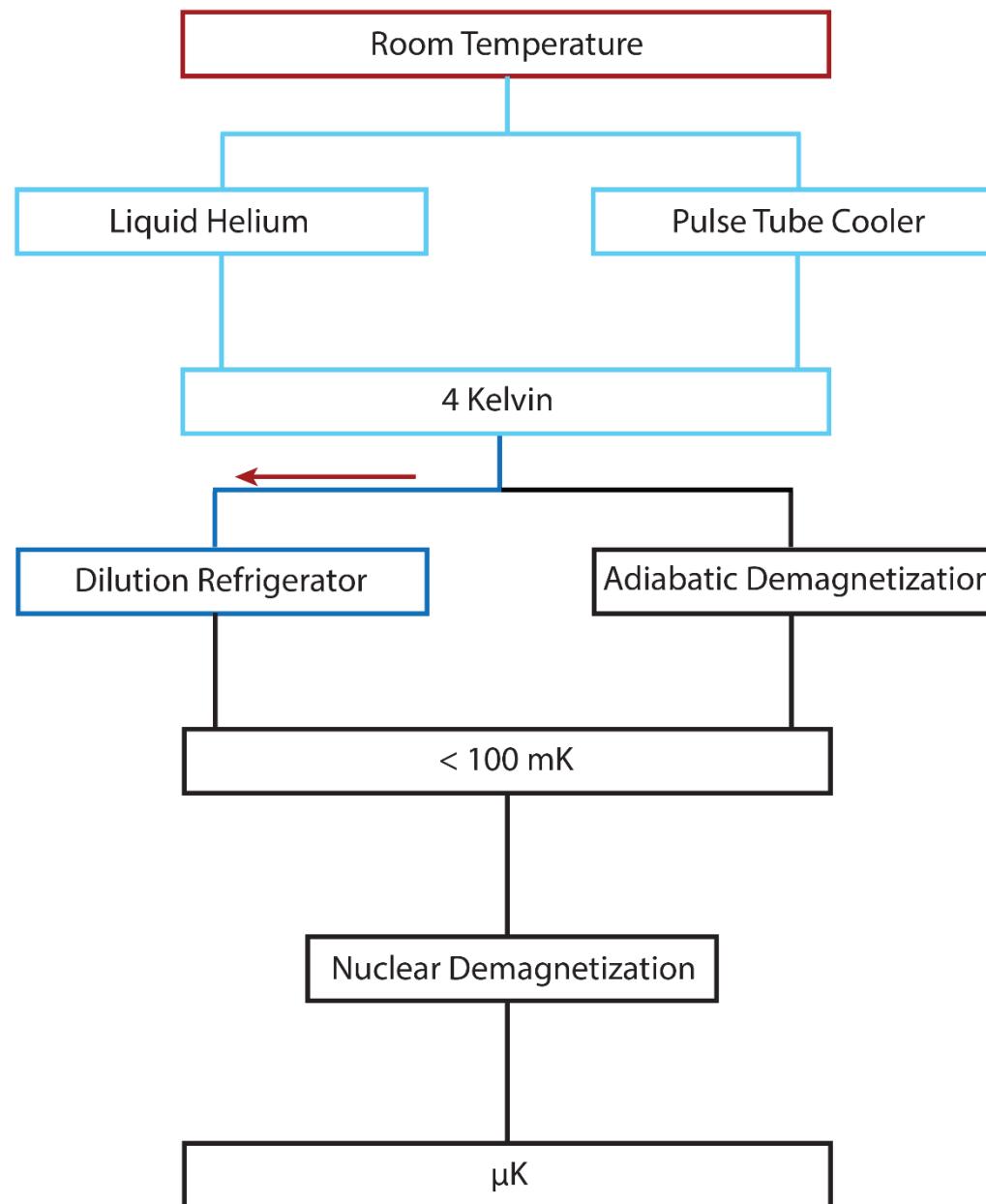


GM-type PTC

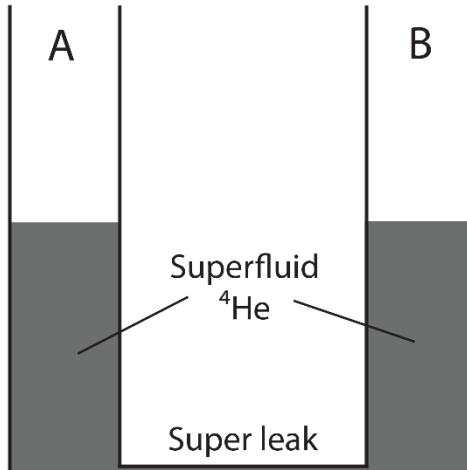


Pulse tube

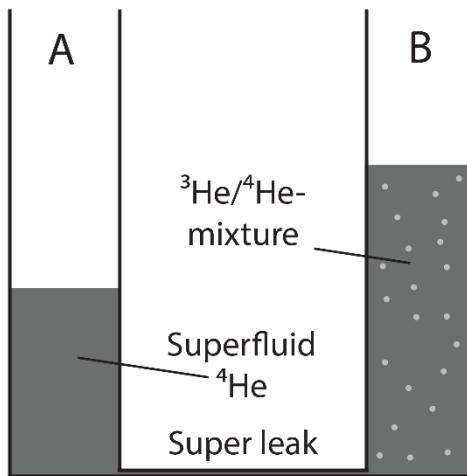




# Osmotic Pressure



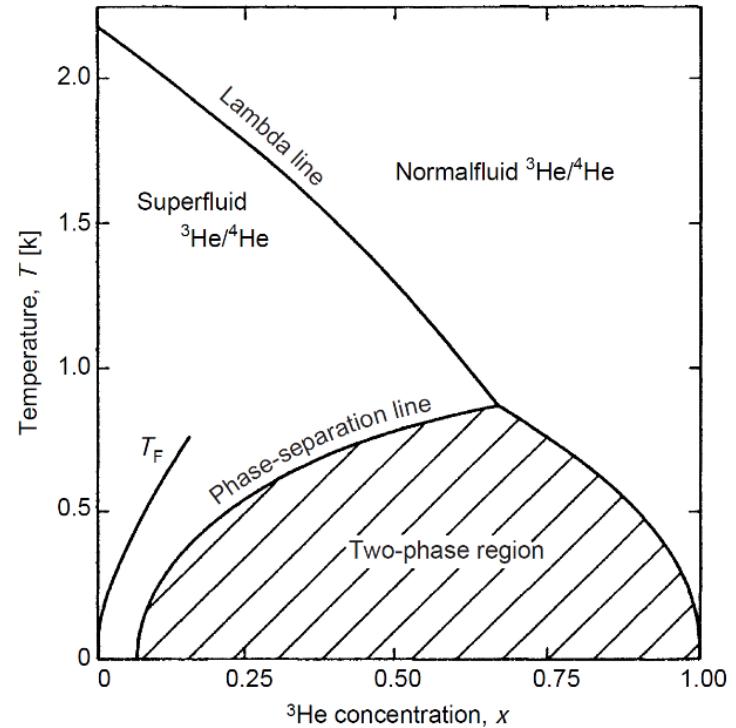
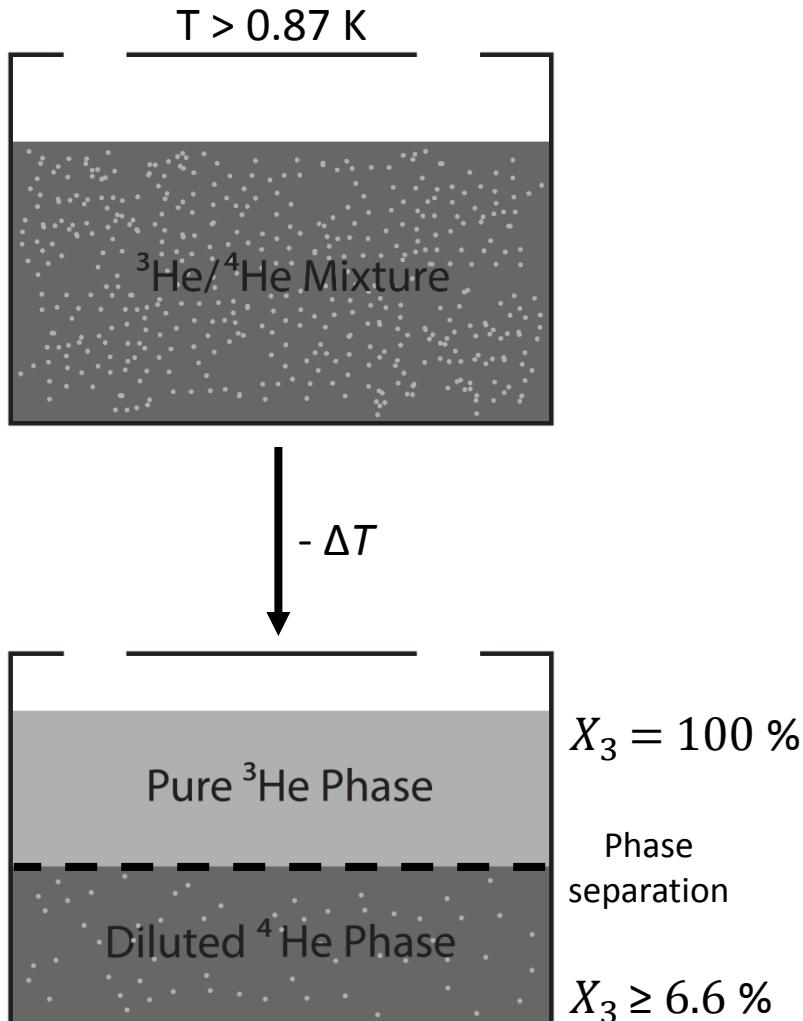
- Two columns filled with superfluid  $^4\text{He}$
- Super leak: only superfluid  $^4\text{He}$  can pass through
- $T = 1 - 2 \text{ K}$



- Adding  $^3\text{He}$  to column B creates an osmotic pressure
- $^3\text{He}$  atoms can't pass the super leak due to their non-zero viscosity
- Superfluid  $^4\text{He}$  passes from A to B trying to equalize the concentration:

$$\frac{X_{3,A}}{X_{4,A}} = \frac{X_{3,B}}{X_{4,B}}$$

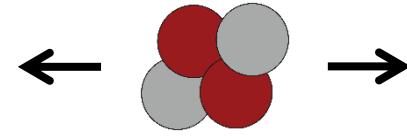
# Mixing $^3\text{He}$ and $^4\text{He}$



# Dilution Refrigerator

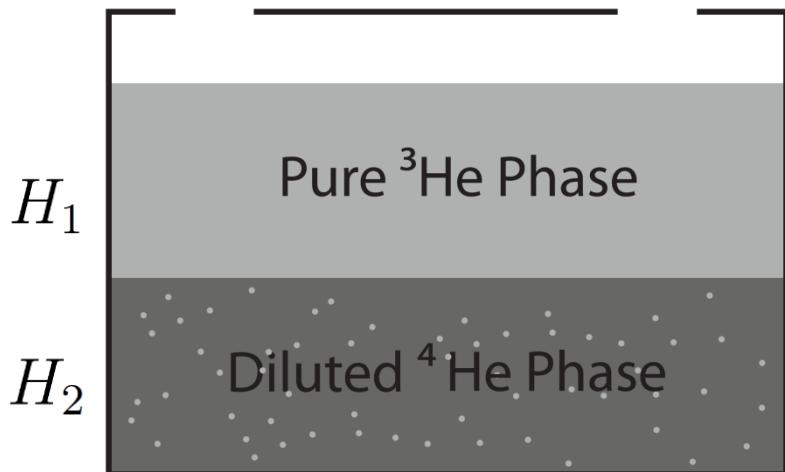
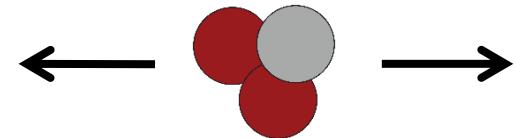
Diluted phase:

- ${}^4\text{He}$  behaves as a “superfluid background”
- ${}^3\text{He}$  can be treated as fermi gas with effective mass  $m^*$
- Zero-Point-Energy of  ${}^3\text{He}$  higher than that of  ${}^4\text{He}$
- Adding  ${}^3\text{He}$  increases kinetic energy, reducing the effective binding energy, hence limiting amount of  ${}^3\text{He}$  solved in  ${}^4\text{He}$



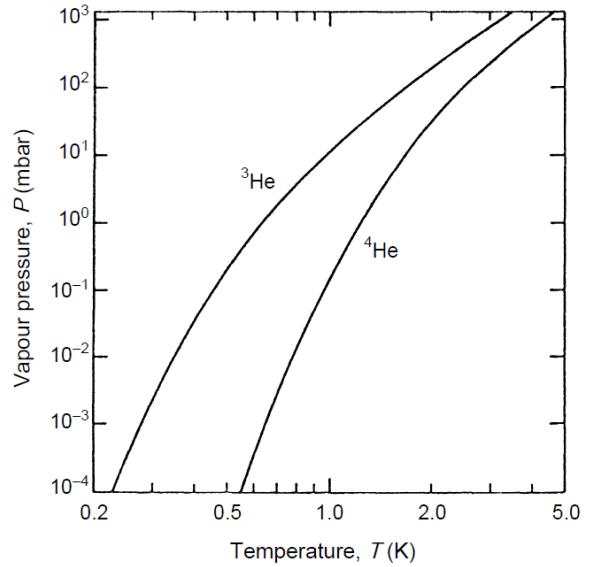
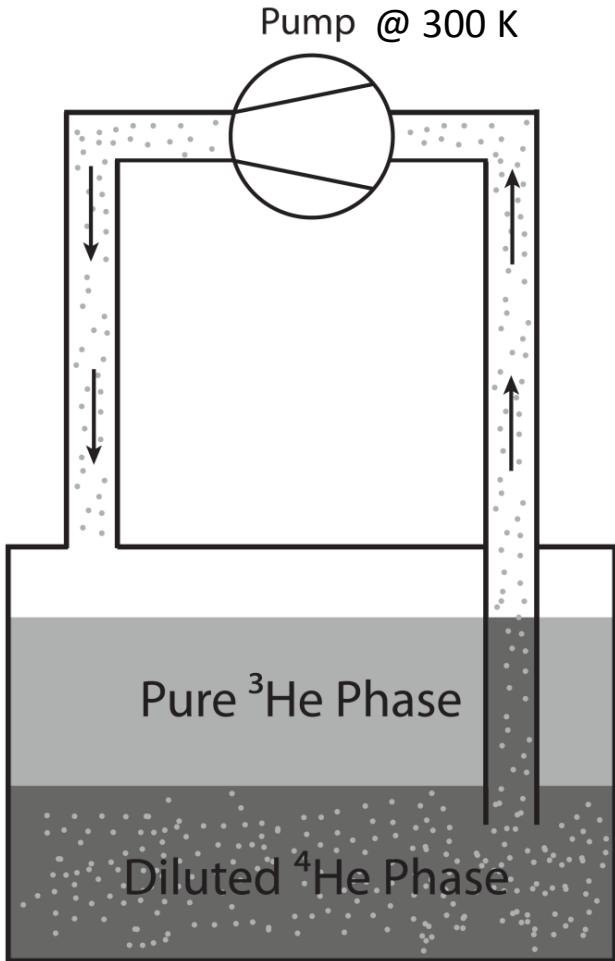
$$m^4\text{He} > m^3\text{He}$$

$$E_{Z,4\text{He}} < E_{Z,3\text{He}}$$



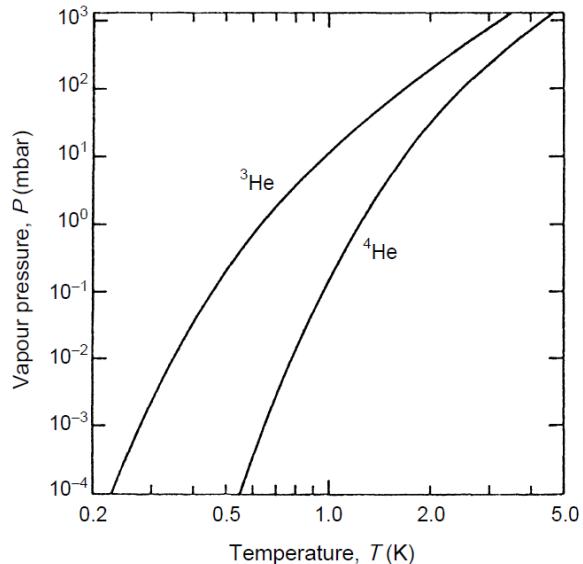
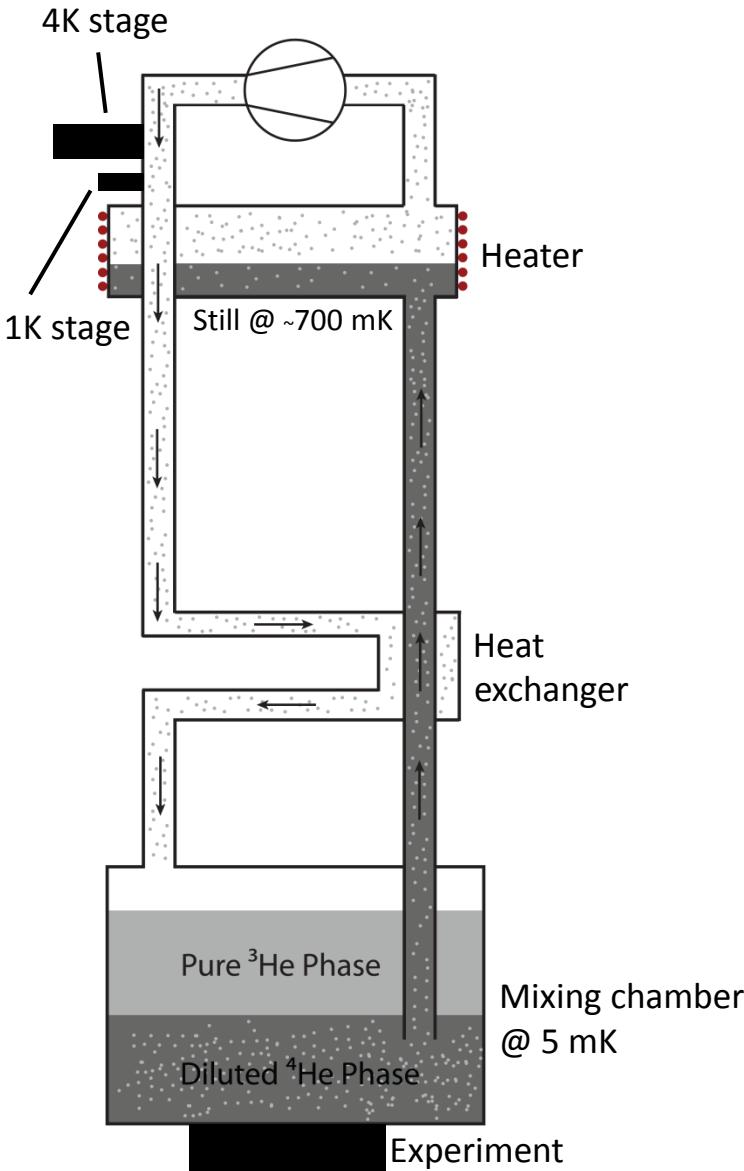
- Phases have different Enthalpy  $H_1$  and  $H_2$
- Removing  ${}^3\text{He}$  from diluted phase results in  ${}^3\text{He}$  reflow from pure phase
- ${}^3\text{He}$  transition into the diluted phase requires energy
- Cooling power:  $\dot{Q} = \dot{n}\Delta H = \dot{n}L$

# Dilution Refrigerator



- Pumping on the diluted phase removes He particles
- More  $^3\text{He}$  is pumped than  $^4\text{He}$  due to different vapour pressures
$$p^3\text{He} > p^4\text{He}$$
- Given scenario: He flowing back into the pure phase is too warm, no relevant cooling effect is achieved

# Dilution Refrigerator



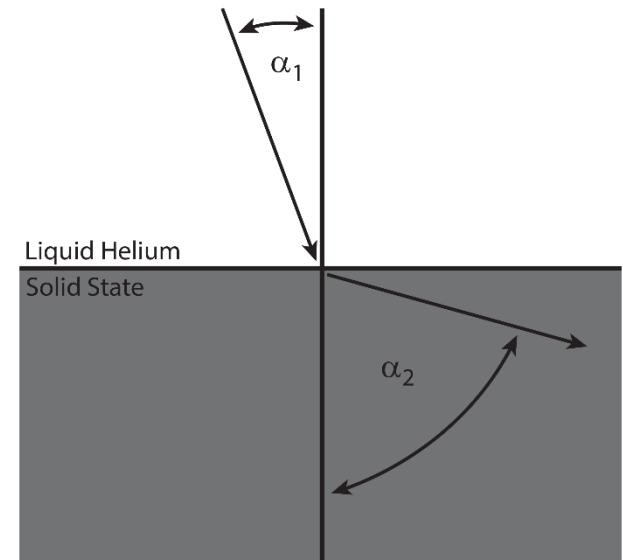
- Pre-cooling pumped helium with available cooling stages
- **Still**: kept at  $\sim 700\text{ mK}$  to keep the  $^3\text{He}$  vapor pressure high  
 $p_{^3\text{He}} \gg p_{^4\text{He}}$
- **Heat exchanger**: crucial for the final temperature of the cryostat.
- **Cooling power**:  $\dot{Q} = \dot{n}_{^3\text{He}} (95 T_M^2 - 11 T_H^2)$

$$\frac{T_H}{T_M} = 2.8 \text{ for } \dot{Q} = 0$$

# Heat Exchanger

Surface Boundary Resistance (**Kapitza-Resistance**):

- Energy carriers (electrons or phonons) are scattered at the interface of two materials
- Total reflection for  $\alpha_1 > 4^\circ$
- Phonon transmission:  $t < 10^{-3}$
- Kapitza-Resistance:  $R_K = \frac{a}{A} T^{-3}$

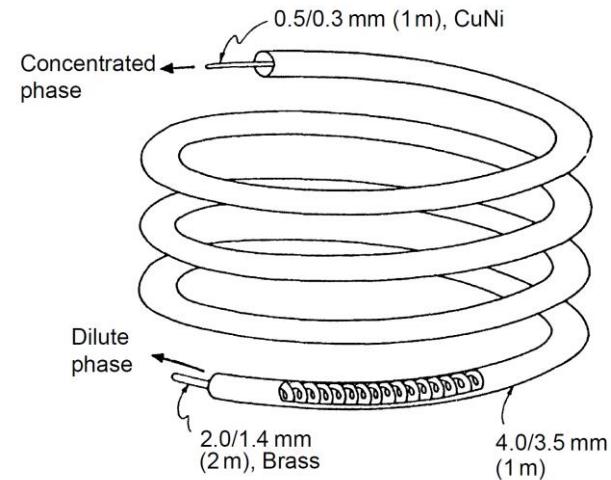


Heat exchanger and mixing chamber must have a huge surface area up to several 100 m<sup>2</sup> !

# Heat Exchanger

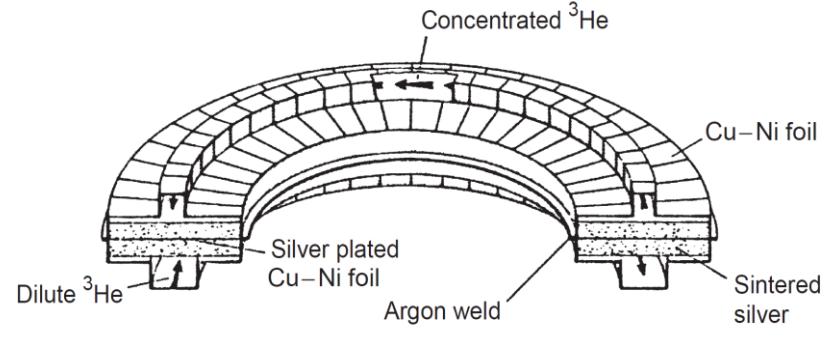
## Continuous heat exchanger:

- Two concentric capillaries
- Diluted phase moves between the tubes, concentrated phase moved in inner capillary
- Does not provide enough surface area for very low temperatures
- 30 mK can be reached with a single continuous heat exchanger

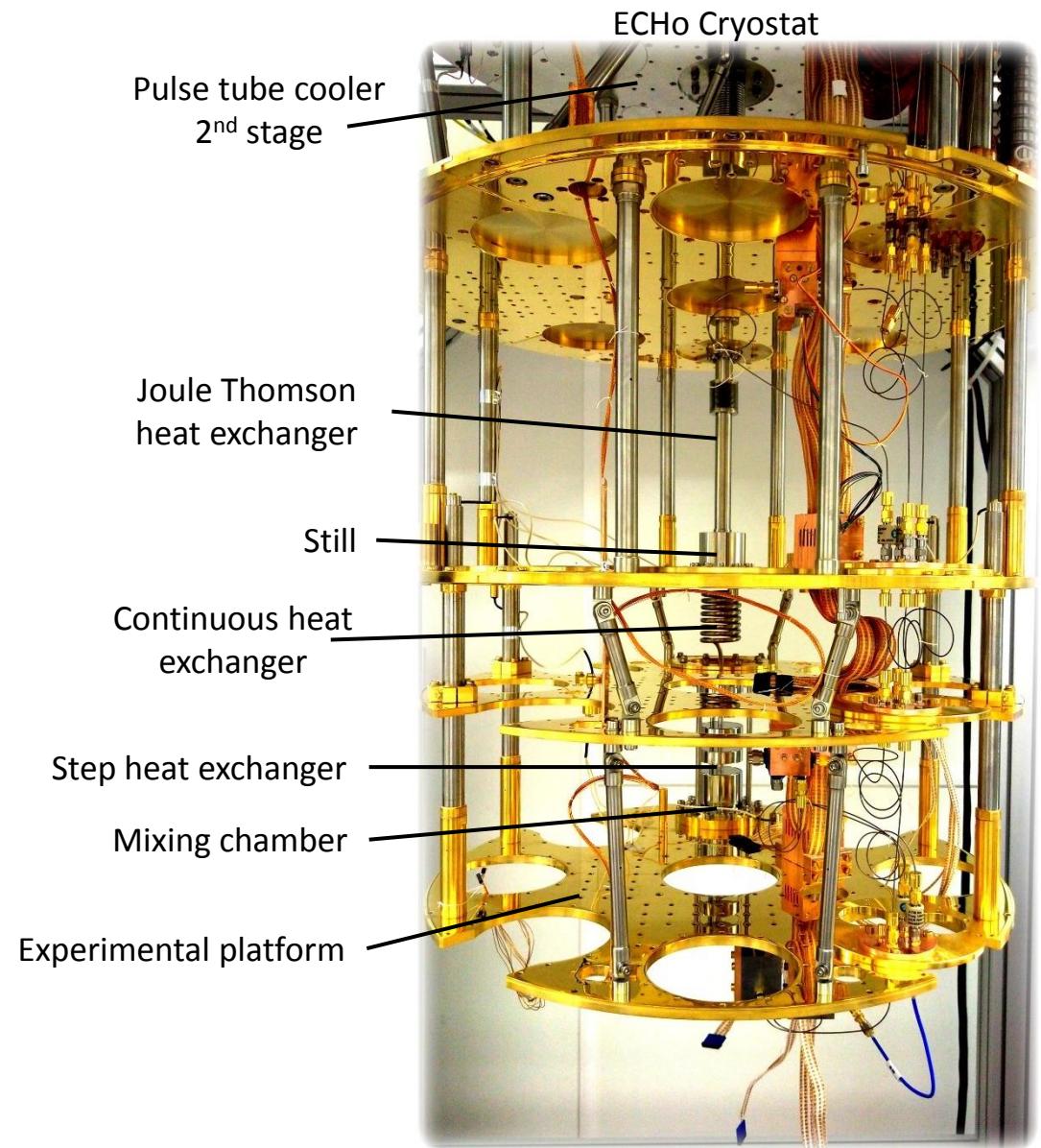
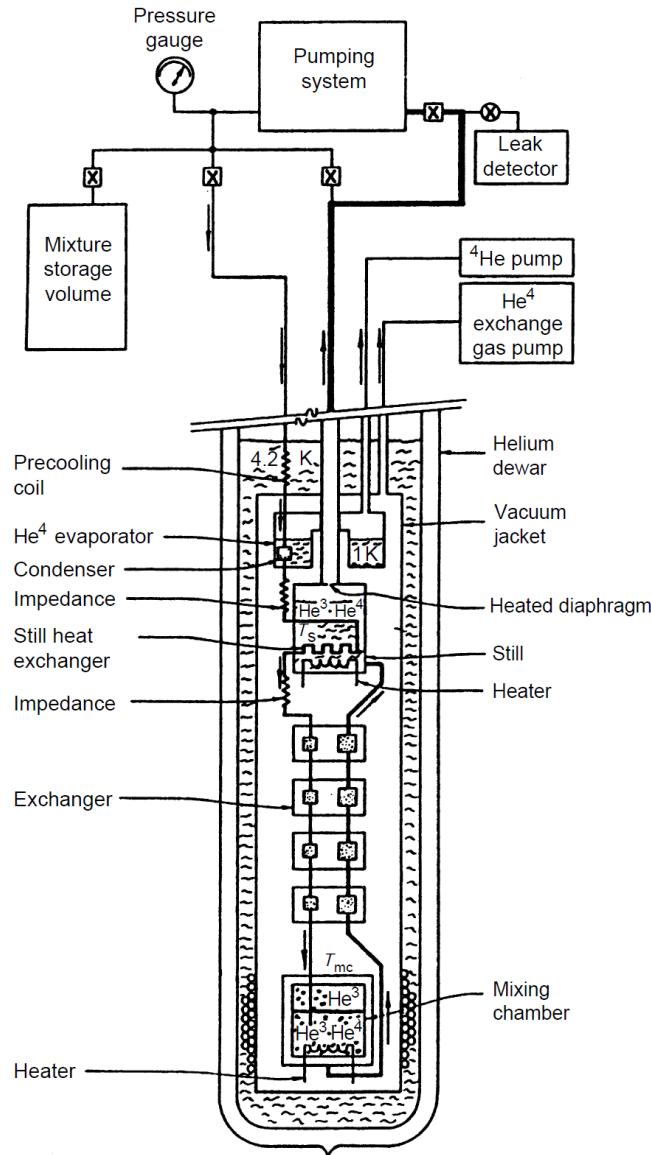


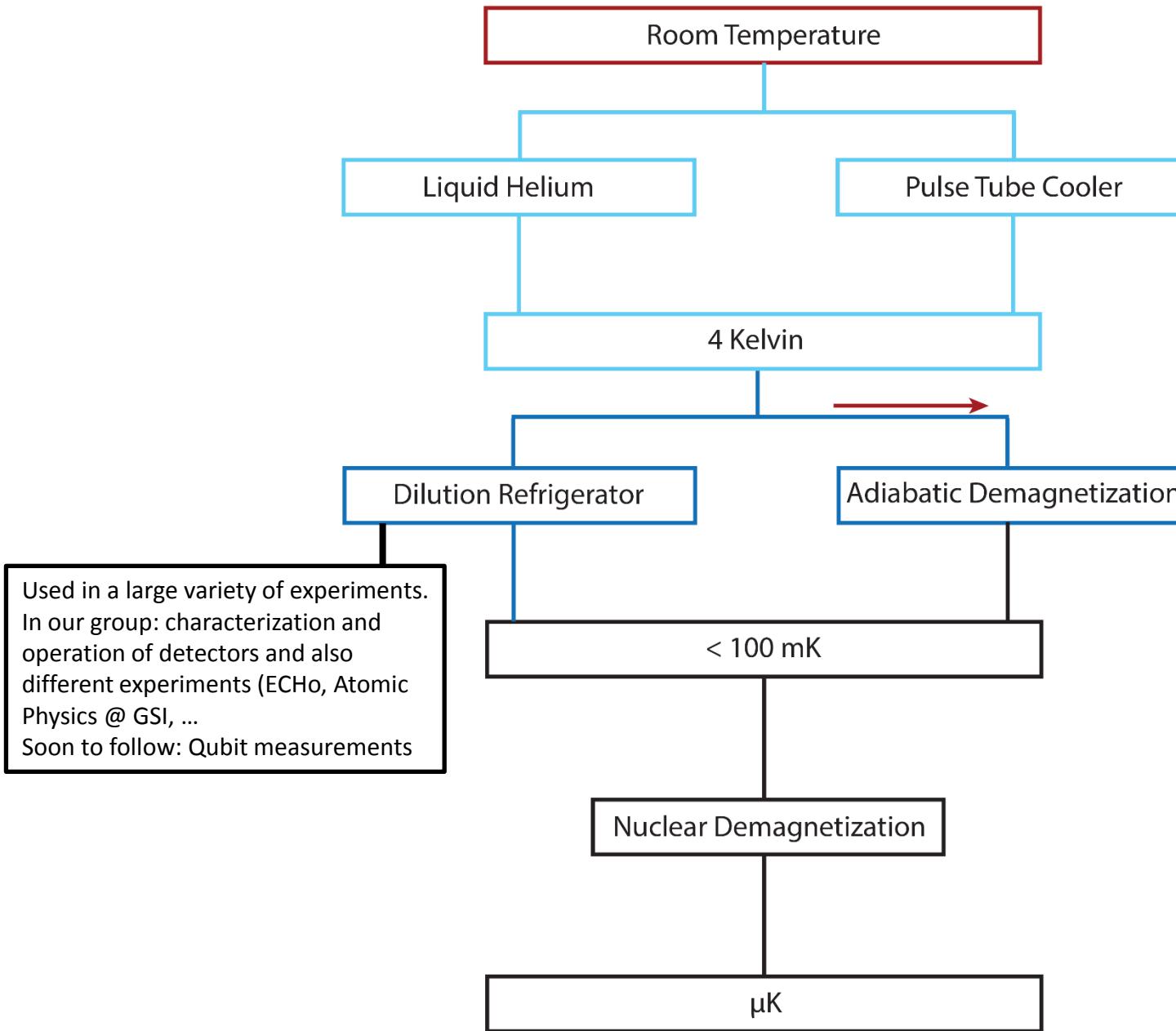
## Step heat exchanger:

- Two metal tubes welded together
- Filled with sintered metal powder to create a huge surface area
- “step” because several are connected in series, each having different geometrical and thermodynamic properties optimized for a certain temperature
- 4 mK are usually reached using several step heat exchangers in series



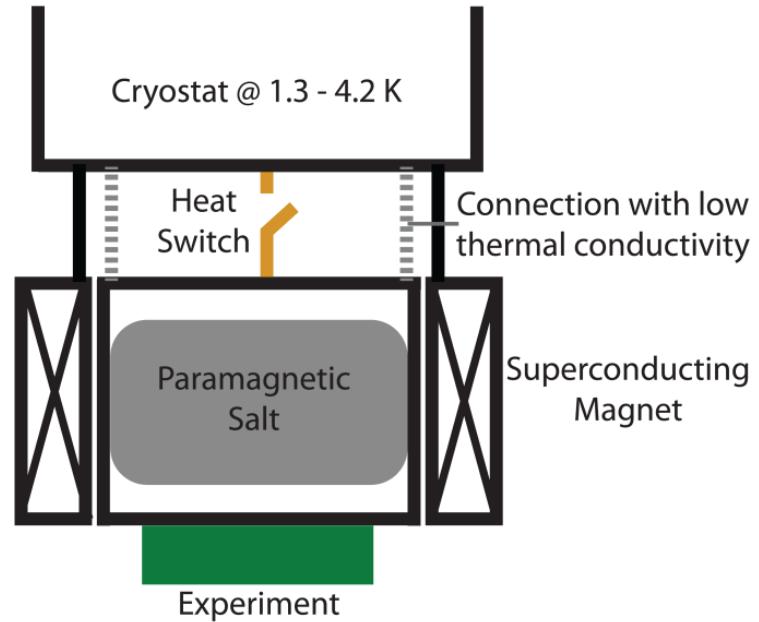
# Dilution Refrigerator



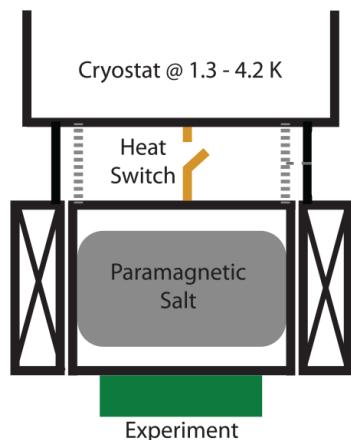
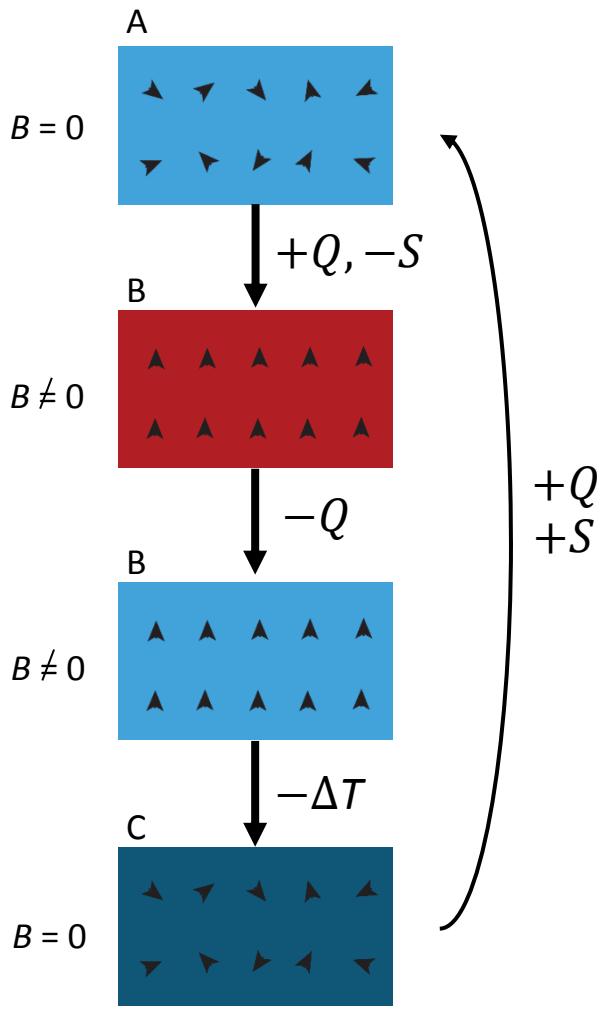


# Adiabatic Demagnetization Refrigerator

- System pre-cooled by LHe or PTC
- Paramagnetic material with unpaired electrons in shell
- Each spin carries magnetic moment:
$$\mu = -g\mu_B S$$
- Material has to be an insulator because conduction electrons influence the magnetic moments
- Paramagnetic salt in good or bad thermal contact with heat bath (depending on HS open or closed)
- B-Field created by a superconducting magnet ( $B > 5$  T)

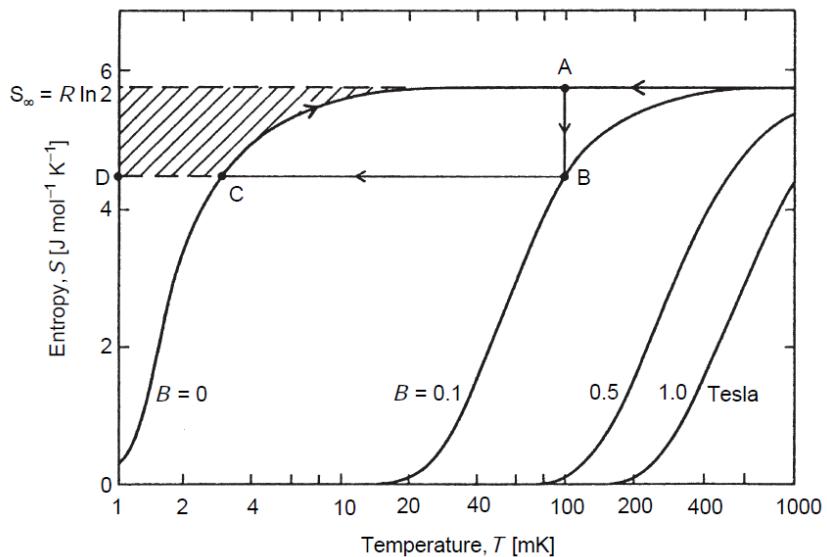


# Cooling Principle

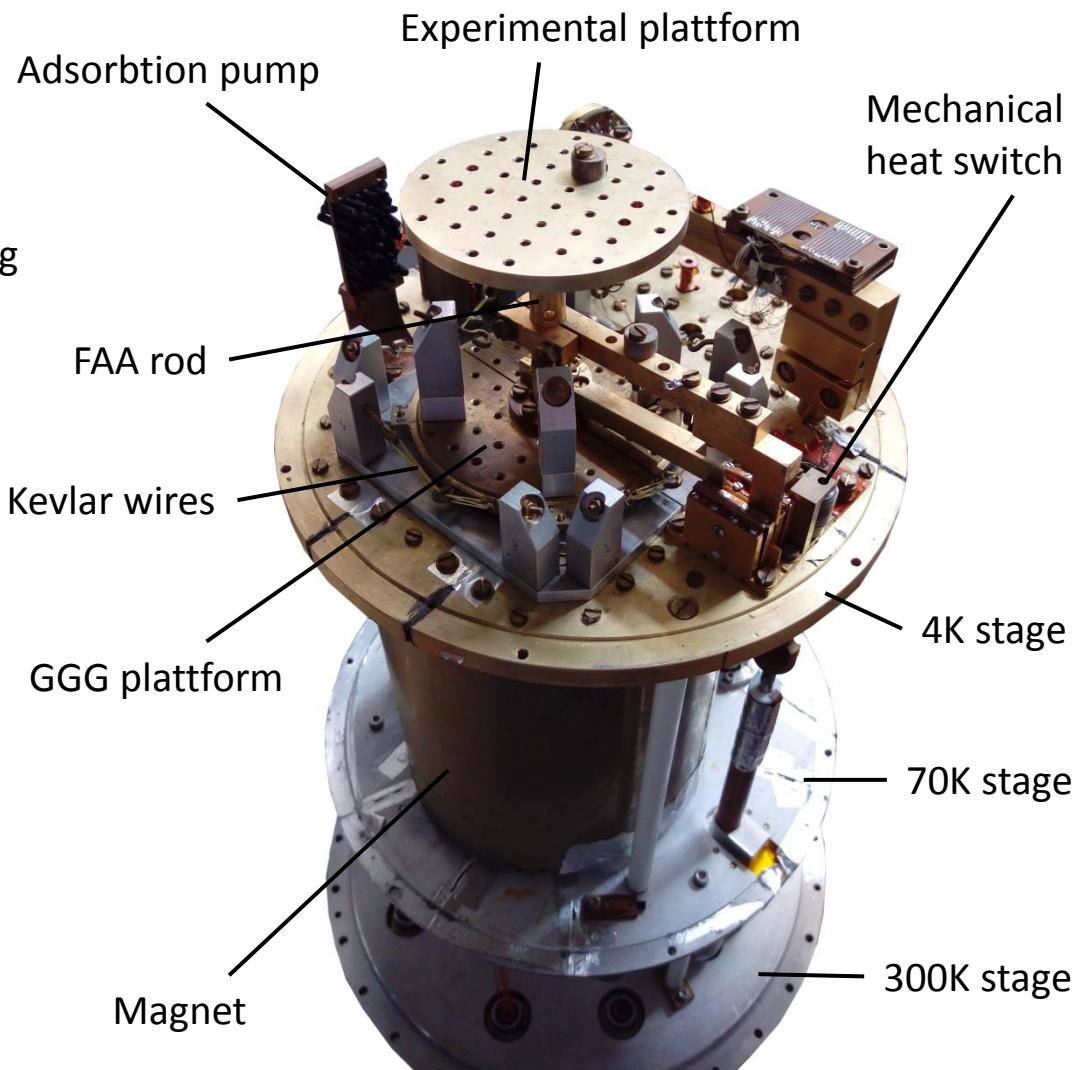
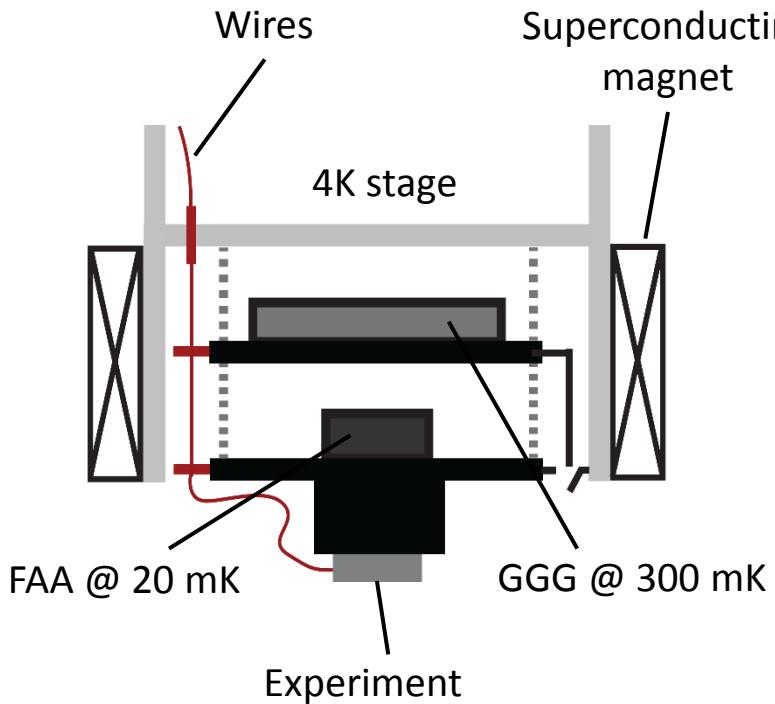


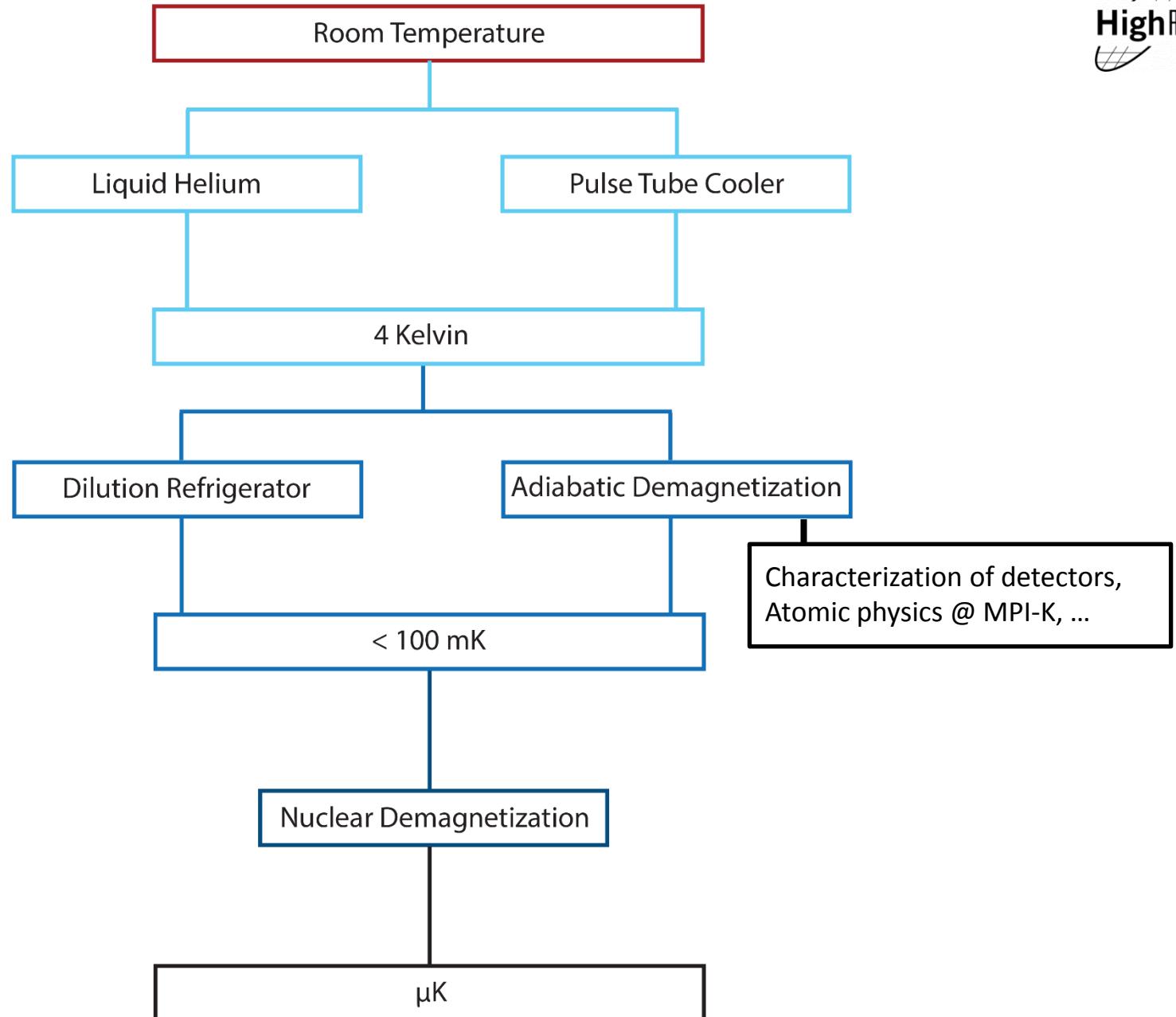
- Heat of magnetization is absorbed by  $^4\text{He}$  or PTC:  

$$Q(T_i) = nT_i[S(0, T_i) - S(B_i, T_i)]$$
- Ordering temperature limits the lowest temperature to be reached



# Two-stage ADR





# Nuclear Demagnetization

- Working principle similar to the electronic ADR

- Take material with  $\mathbf{S} = 0$  and  $\mathbf{I} \neq 0$

- Magnetic moment, electron vs. nucleus:

$$\mu_B = 9.27 \times 10^{-24} \text{ J T}^{-1}$$

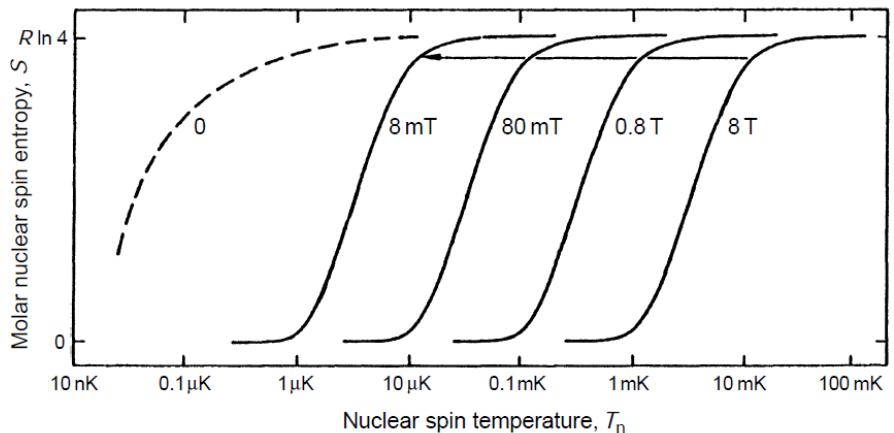
$$\mu_n = 5.05 \times 10^{-27} \text{ J T}^{-1}$$

- Small magnetic moment provides **advantages**:

- Lower ordering temperature of spins,  $T \sim 0.1 \mu\text{K}$
- Material can (must) be a metal providing high thermal conductivity
- Magnetic moment density larger in pure metals than in diluted paramagnetic salts

- Drawback:** smaller magnetic moment also means a smaller “reaction” to external magnetic fields

$$\Delta S_n \sim \frac{\Delta S_e}{1000}$$

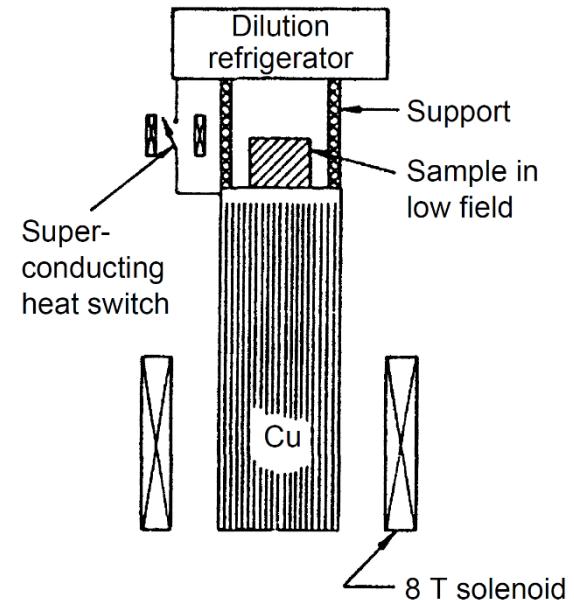


# Nuclear Demagnetization

Finding the right material with requirements:

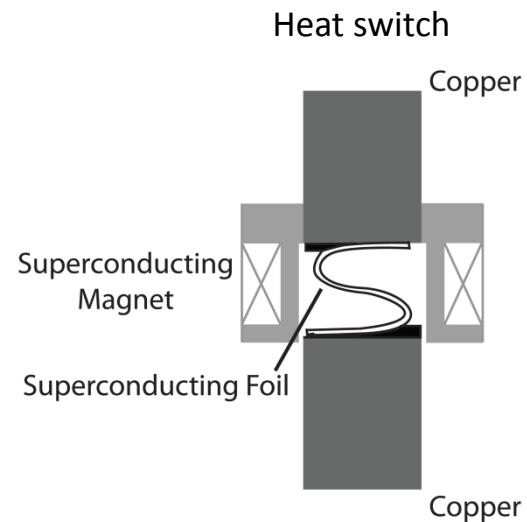
- Normal conducting metal
- High thermal conductivity
- Low spin ordering temperature
- Has no ordered electron spins
- Easy to fabricate with high purity

→ Copper or Platinum



## Superconducting heat switch

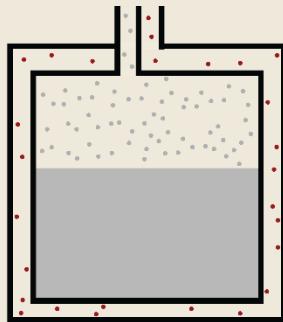
- Very low thermal conductivity in SC state  
(conduction dominated by phonons)
- Thin aluminium foil shaped like an S
  - Small mean free path of phonons
  - No normal conducting path after  $B \rightarrow 0$



# Heat Leaks in NDR

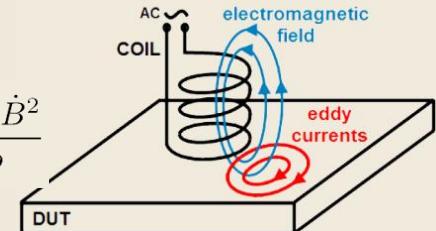
## External heat leaks:

- Conduction (mountings, wires, heat switch, ...)
- Convection (residual gas)
- Radiation (also by radioactive sources)



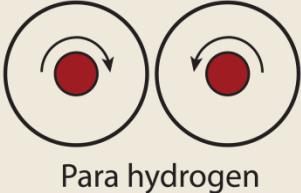
## Eddy currents:

- B-field influences electrons
- Vibrations inside B-field
- Changing B-field
- Heat input:  $\dot{Q}_e = \frac{GV\dot{B}^2}{\rho}$



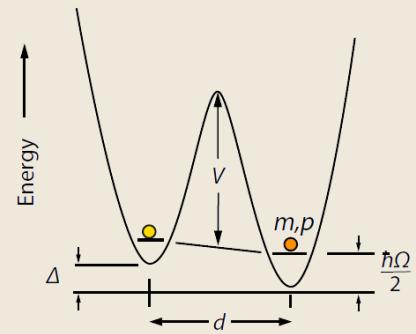
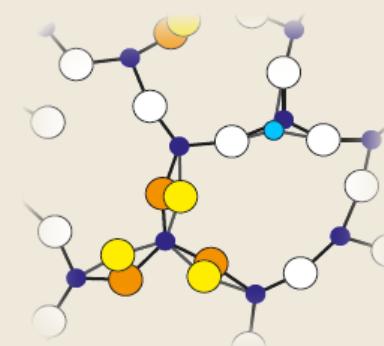
## Time dependent leaks:

- Hydrogen bubbles found in certain metals
- Ortho-Para conversion of hydrogen releases energy

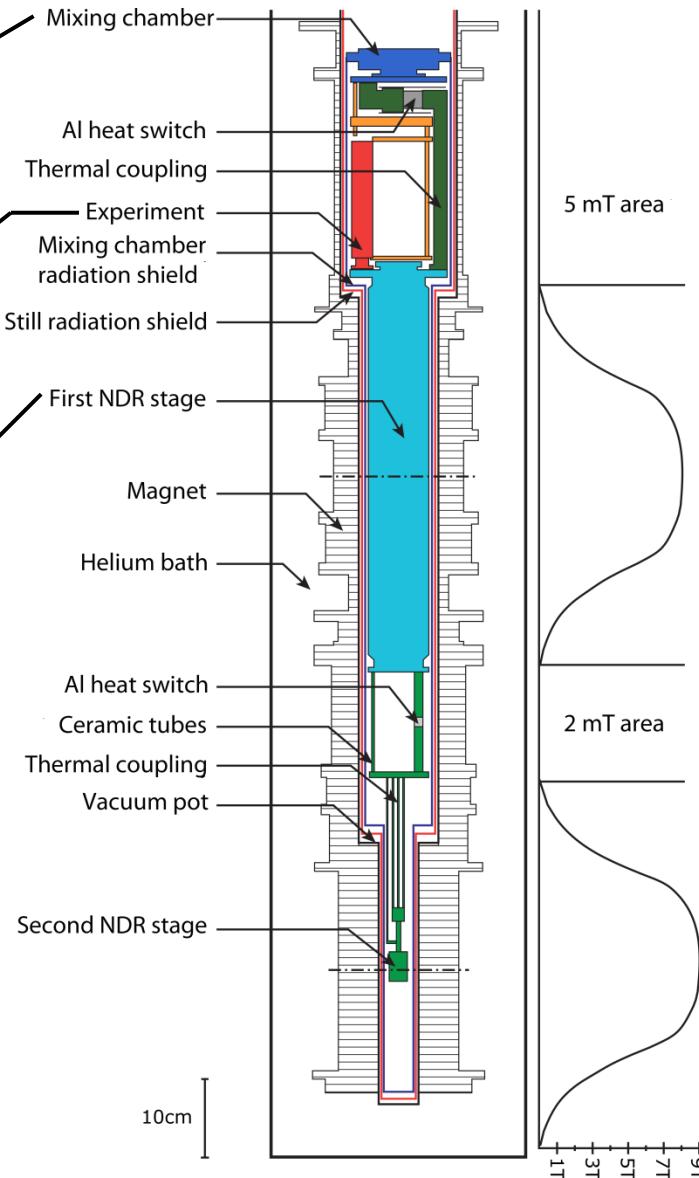


## Internal leaks:

Tunneling systems



# Nuclear Demagnetization

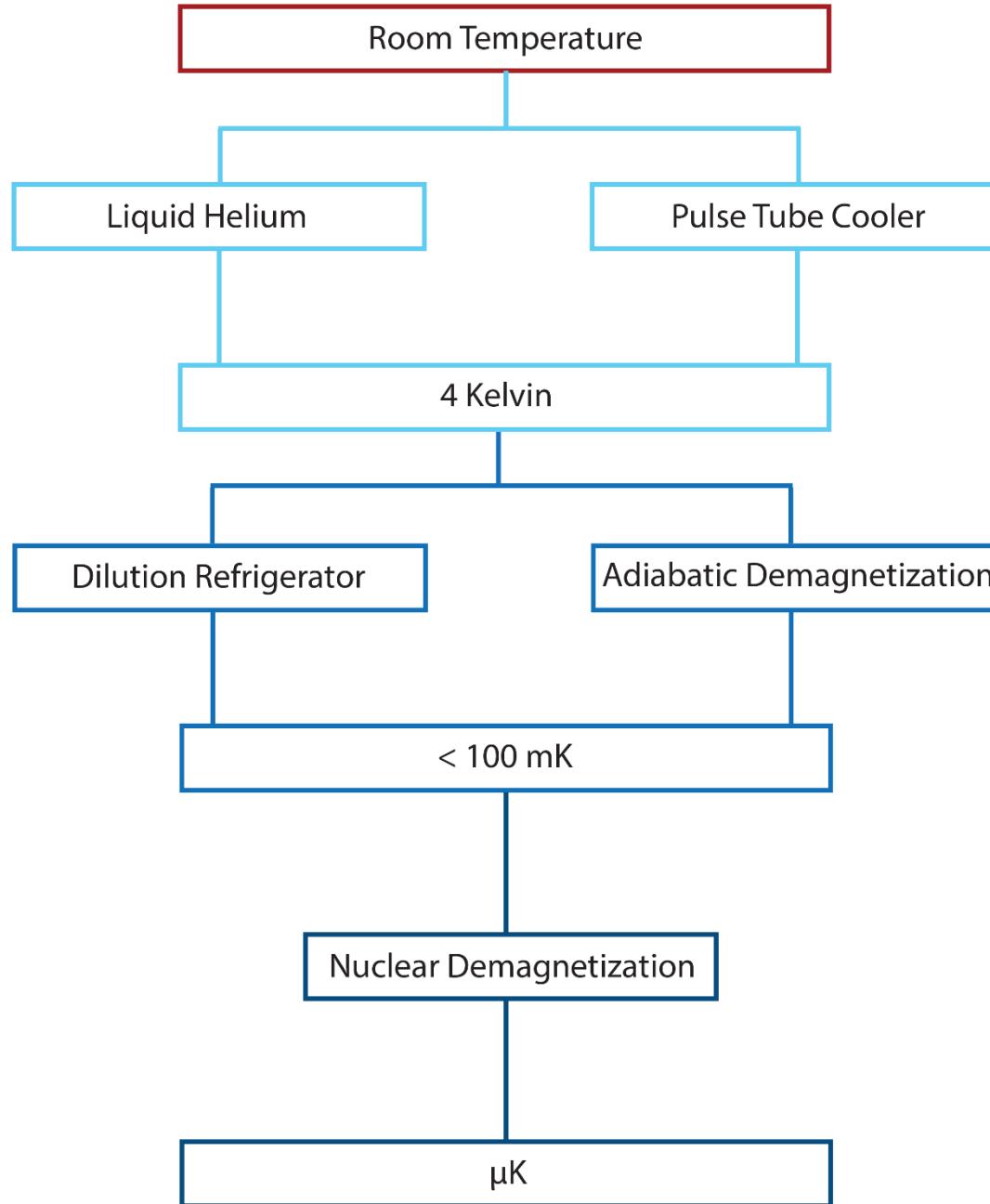


## Experiments:

- Thermometry
- Search for superconductors
- Investigation of materials
  - $^3\text{He}$
  - Glasses
  - Quantum magnets

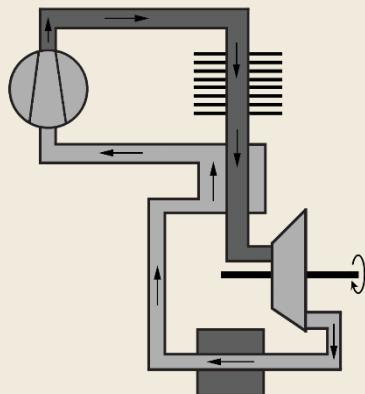
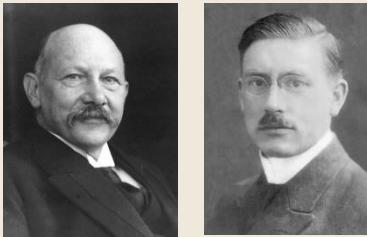
## Specific examples:

- Noise Thermometer
- Novel fractional quantum Hall States in two-dimensional electron systems
- Gravity waves on a surface of topological superfluid  $^3\text{He-B}$
- Superconductivity in Polycrystalline Boron-Doped Diamond

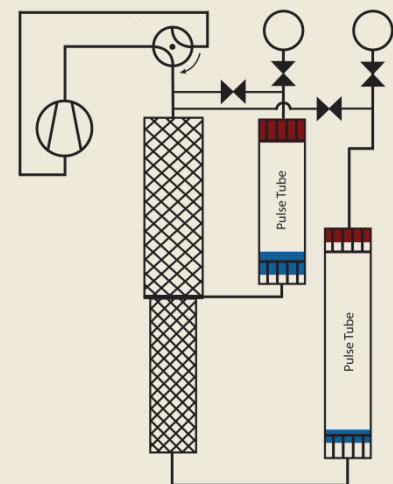
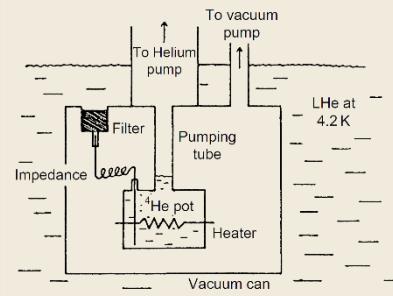


# Summary

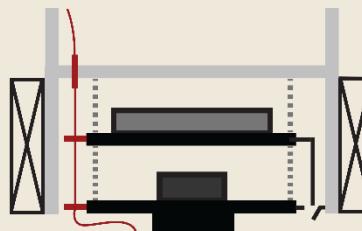
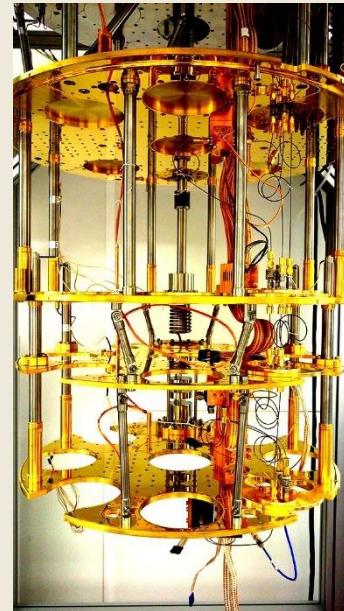
## Historic Introduction & Cryoliquids



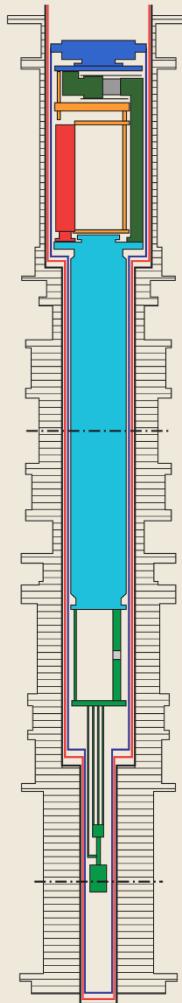
## $^4\text{He}$ Cryostat & Pulse Tube Cooler



## Dilution Refrigerator & ADR



## Nuclear Demagnetization

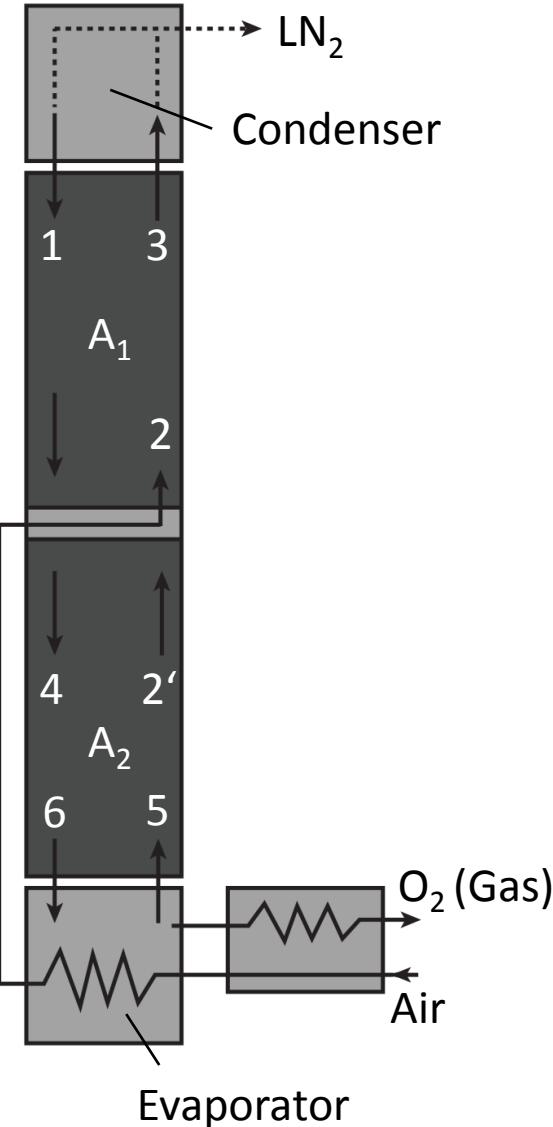


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Thank you for your attention!

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# Cryoliquid $\text{LN}_2$ @ $T = 77 \text{ K}$



## Separation of air into nitrogen and oxygen

1.  $\text{LN}_2$  (~77 K) trickles from condenser into volume  $A_1$
2. Pre-cooled air (~ 85 K) enters  $A_1$  from below and mixes with  $\text{LN}_2$
3. Nearly pure nitrogen (~ 99.5 %) reaches the condenser, part of it gets removed while the rest goes to 1
4. Droplets reaching  $A_2$  consist of 50 % nitrogen and 50 % oxygen
5. Oxygen coming from the evaporator meets droplets from 4 which in turn release their nitrogen
6. Oxygen droplets move back to the evaporator where they get evaporated again, part of the oxygen is removed

# Cooldown with Cryoliquids



## Cooldown with liquid helium:

Heat of evaporation @ 4.2 K:

$$L_{^4\text{He}} = 2.6 \text{ kJ/l}$$

Enthalpie from 4.2 K to 300 (77) K :

$$\Delta H_{^4\text{He}} = 200 \text{ (64) kJ/l}$$

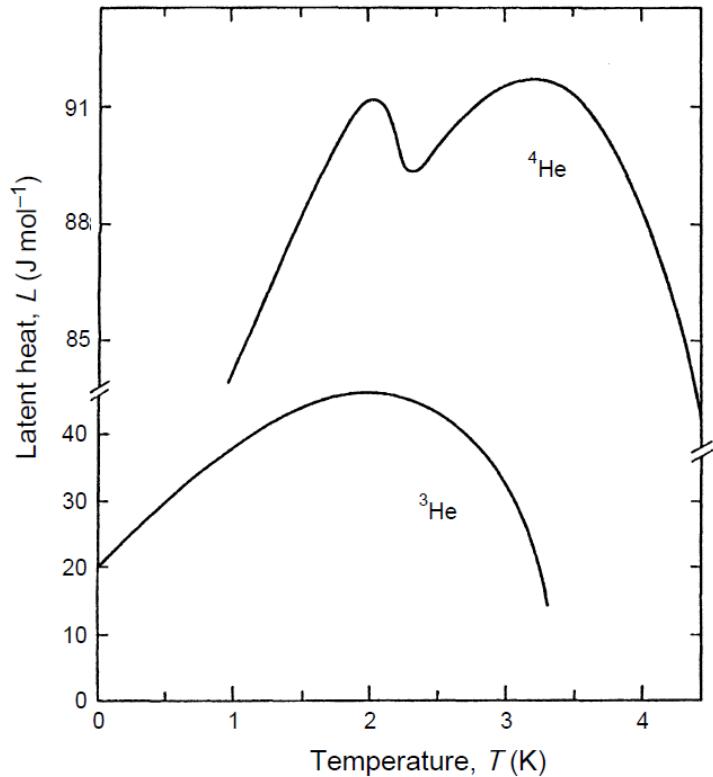
Amount of cryoliquid [l] necessary to refrigerate 1 kg of aluminum, stainless steel (SS) or copper.

cryoliquid	temperature change [K]	Al	SS	Cu
N <sub>2</sub>	300 → 77	1.0 (0.63)	0.53 (0.33)	0.46 (0.28)
<sup>4</sup> He	77 → 4.2	3.2 (0.20)	1.4 (0.10)	2.2 (0.16)
<sup>4</sup> He	300 → 4.2	66 (1.6)	34 (0.8)	32 (0.8)



- Pre-cooling of equipment with LN<sub>2</sub> to 77 K
- Use liquid <sup>4</sup>He to go from 77 K to 4.2 K
- Use the enthalpy of the helium gas, hence it should leave the cryostat as warm as possible

# Latent Heat of Evaporation



- Latent heat: energy needed for the phase transition from liquid to gas
- $L$  of helium very small compared to other liquids
- Very good shielding against external parasitic heat input needed
- Heat load from:

Conduction:

$$\text{Phonons } \dot{Q} = \frac{Ab}{4L}(T_2^4 - T_1^4)$$

$$\text{Electrons } \dot{Q} = \frac{A\kappa_0}{2L}(T_2^2 - T_1^2)$$

Radiation:

$$\dot{Q}[\text{W}] = 5.67 \times 10^{-12} A[\text{cm}^2](T_1^4 - T_2^4)$$

Gas particles:

$$\dot{Q}[\text{W}] \approx 0.02aA[\text{cm}^2]P[\text{mbar}](T_2 - T_1)[\text{K}]$$

$$L_{\text{N}_2} = 5570 \text{ J mol}^{-1}$$

# Adiabatic Demagnetization

## Paramagnetism:

- Ions with unpaired electrons in shell carry magnetic moment:

$$\mu = -g\mu_B J$$

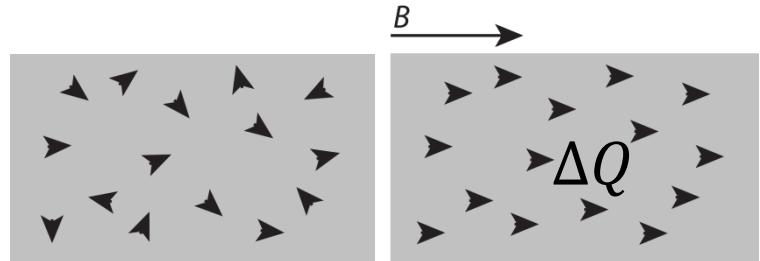
- Spins align along magnetic field lines of external origin

- Susceptibility:  $\chi = \frac{M}{H} > 1$

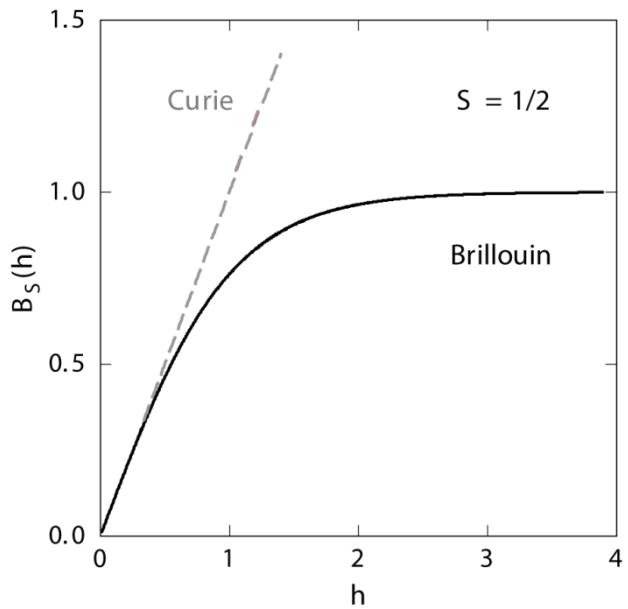
- Entropy:  $S = R \ln(2J + 1)$

- Magnetization:  $M = ng\mu_B J \mathcal{B}(h)$

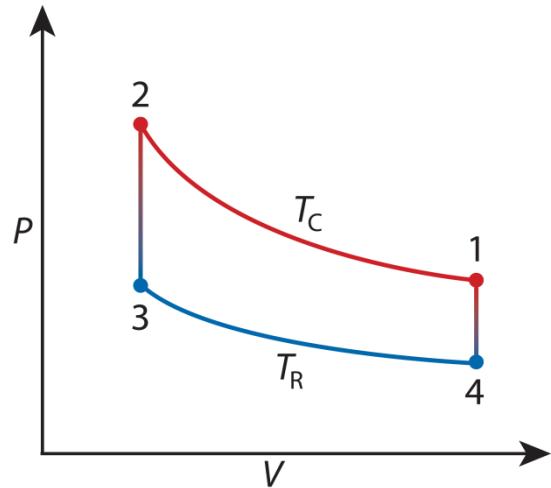
- For  $h < 1$ , Curie Law:  $\chi = \frac{M}{H} = \frac{C}{T}$



Brillouin function

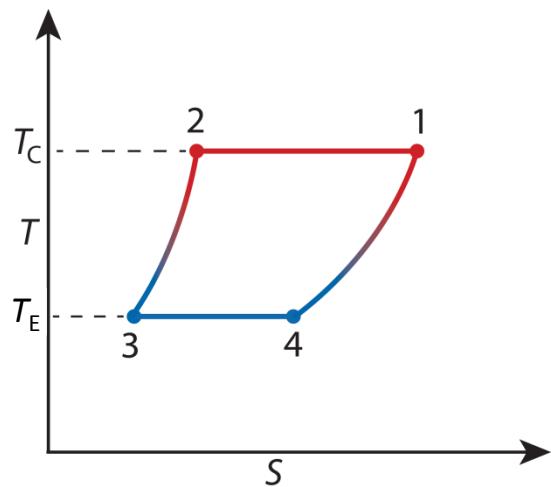


# Non-ideal Stirling Cycle



How is a real machine different from the ideal?

- Harmonic motion of pistons (instead of abrupt)
- Void volume in regenerator
- Regenerator ineffectiveness
- Pressure drop through regenerator
- Non-isothermal compression and expansion



# Nuclear Spin and Magnetic Moment

- Nuclear spin relaxation very slow (especially in insulators)

$$\frac{dM}{dt} = -\frac{(M-M_0)}{\tau_1}$$

- Nuclear spin relaxation slower than interaction between nuclear and electron spins

$$\tau_{\text{eff}} = \tau_1 \frac{C_e}{C_n + C_e}, \quad C_n \gg C_e$$

- Small magnetic moment (less interaction between nuclear moments means lower ordering temperature)

$$\mu_B = 9.27 \times 10^{-24} \text{ J T}^{-1}$$

$$\mu_n = 5.05 \times 10^{-27} \text{ J T}^{-1}$$

