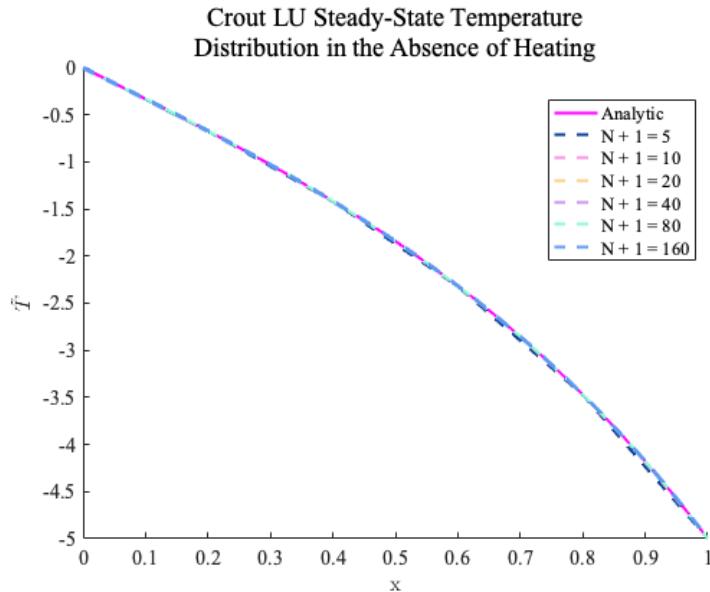


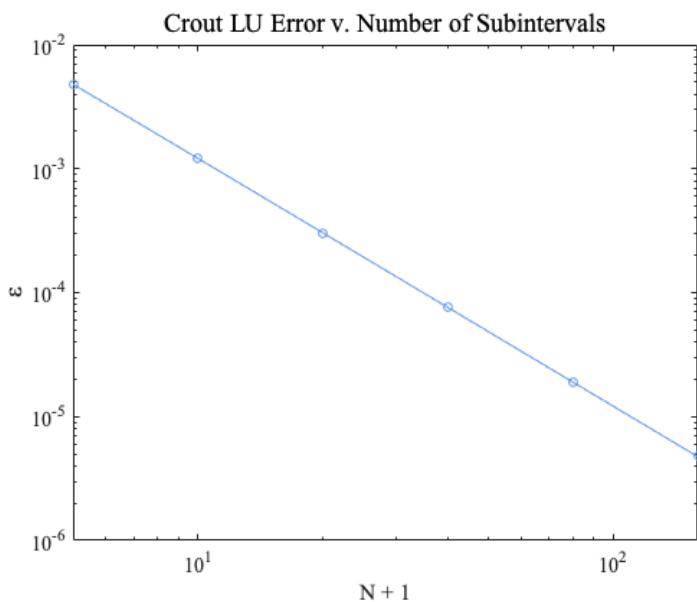
# Lab 8

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A.



**Figure 1:** Steady-state temperature distribution in the absence of heating, using Crout LU decomposition to solve a system of finite difference equations for  $N$  interior nodes.

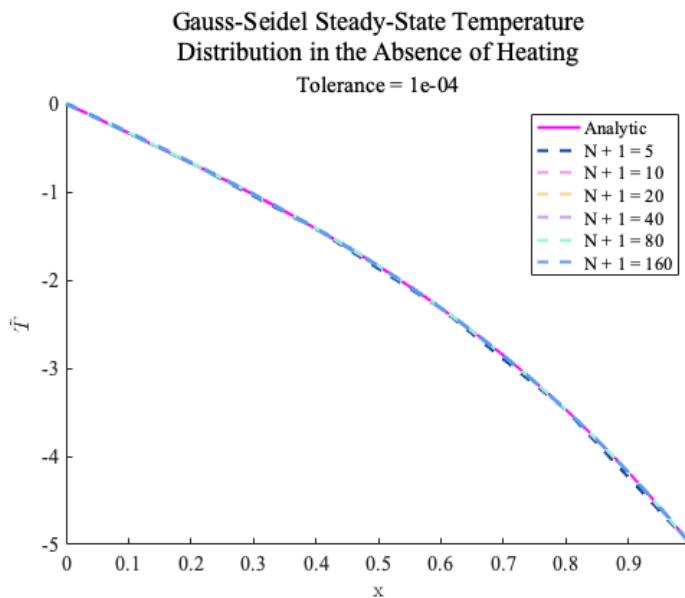


<b>N+1</b>	<b>Max. Error</b>
5	0.048
10	0.0012
20	3.0349e-04
40	7.5931e-05
80	1.8998e-05
160	4.7496e-06

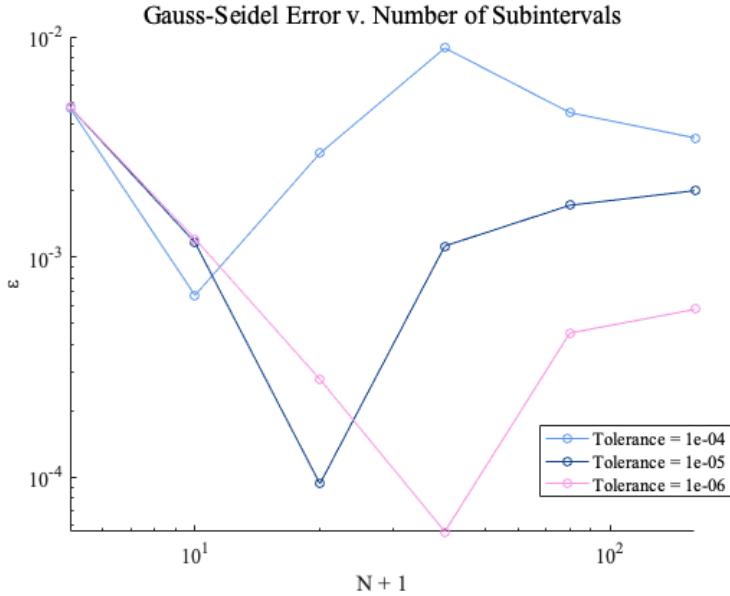
**Figure 2:** Maximum absolute error  $\varepsilon = \max | \bar{T}_i - \bar{T}_i^a |$  as a function of  $1/\Delta x$  on a log-log scale.

The Crout LU decomposition results agree with the left division results from Lab 7. Error versus the inverse of step size still confirms that this method is  $O(h^2)$ , since halving step size reduces error by a factor of 4.

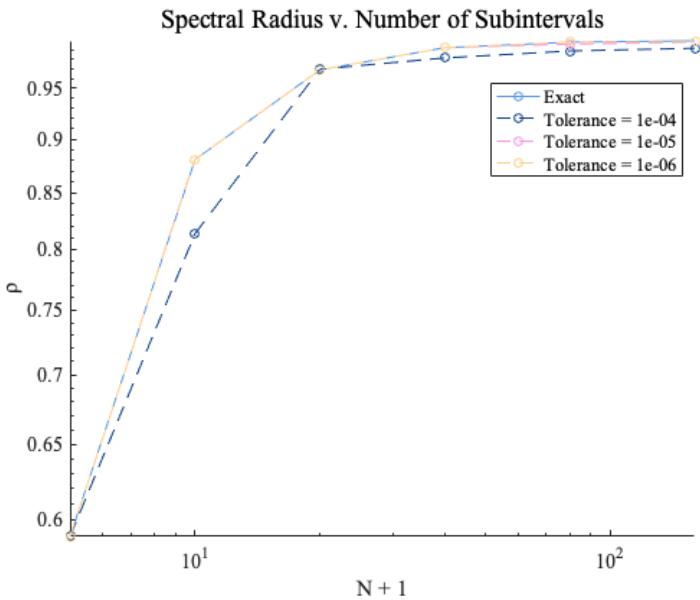
B.



**Figure 1:** Steady-state temperature distribution using a Gauss-Seidel iteration to solve the same system of finite difference equations. For the initial guess, I added a small random offset to the analytic solution for each  $N$  (otherwise, most runs would converge in 1 iteration of Gauss-Seidel). I used the  $L_\infty$  relative norm of at least  $10^{-4}$  as the stopping criterion. Although the graph above shows the solution specifically for a tolerance of  $10^{-4}$ , tolerances of  $10^{-5}$  and  $10^{-6}$  gave similar results, but required more iterations. The number of iterations also increased as  $N$  increased. For a tolerance of  $10^{-4}$ , the solution for  $N + 1 = 5$  subintervals converged after 11 iterations, while  $N + 1 = 160$  converged after 194. With the same initial guesses, the solutions for  $N + 1 = 5$  and  $N + 1 = 160$  with a tolerance of  $10^{-6}$  required 19 and 1486 iterations respectively to converge.



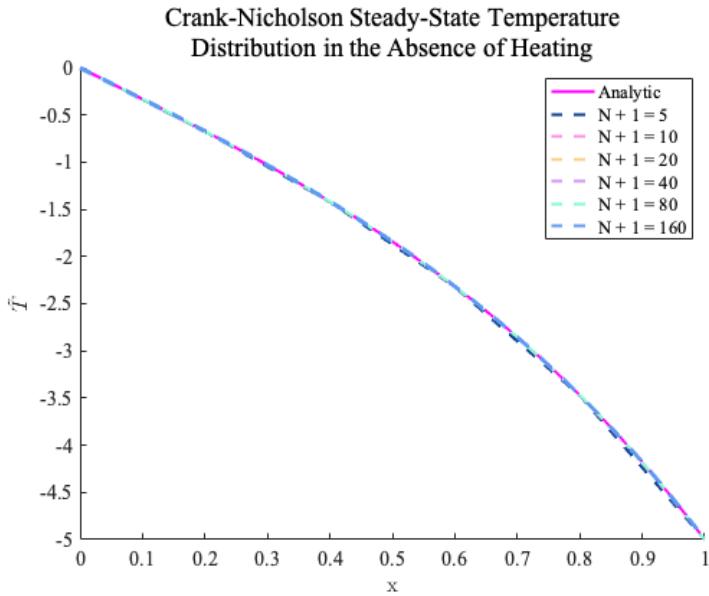
**Figure 2:** The Gauss-Seidel error behavior differs from the behavior of the error associated with the solution from the Crout LU decomposition. Plotting the error versus  $N + 1$  for various tolerances suggests that the maximum absolute error between the numerical and analytic solutions is generally smaller for lower tolerances. Whereas the Crout LU error consistently decreases as  $N$  increases, the Gauss-Seidel error first decreases to some minimum between the smallest and largest values of  $N$ , then increases. This means that unlike Crout LU, where the largest  $N$  (smallest step size) clearly led to the most accurate solution, the most accurate Gauss-Seidel solution for all tolerances is for an intermediate value of  $N$ , which seems to increase as tolerance decreases.



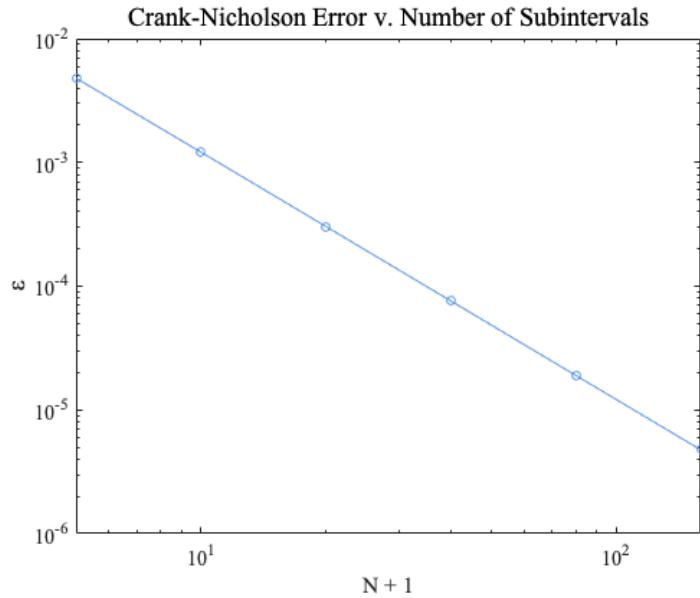
**Figure 3:** Estimates of the spectral radius of A, the matrix describing the system of finite difference equations, for various N and tolerances. These estimates were computed as  $\|\delta^{k+1}\|_2 / \|\delta^k\|_2$ , where  $\delta^k = x^k - x^{k-1}$ . As N increases, the spectral radius approaches but never reaches 1 (since the spectral radius must be less than 1 for the iteration to converge). This is consistent with the Gauss-Seidel iteration converging slower for larger values of N. Based on the graph, lowering the tolerance seems to result in more accurate estimates of the spectral radius. Tolerances of  $10^{-5}$  and  $10^{-6}$  produce estimates that visually overlap with the exact spectral radii, whereas the largest tolerance,  $10^{-4}$ , deviates but still follows the same trend.

### C.

$c_b p_t \frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} - m(T - T_a)$  bioheat eqn, no heating (time-dependent).  
 $\frac{c_b p_t}{k} \frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} - \frac{m}{k}(T - T_a)$   
 $\frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} - \frac{m}{k}(T - T_a)$   
 $\frac{\partial T}{\partial x}|_{x=0} = -5$   
 $\tilde{T}(x=0, t) = T_c - T_a$   
 $\tilde{T}(x=L, t) = T_b - T_a$   
 $\tilde{T}(x=L, t=0) = -5$   
 $\sim \frac{\partial^2 T}{\partial t} = \frac{\partial^2 T}{\partial x^2} - \lambda^2 T, \quad \tilde{T} = T - T_a, \quad \lambda^2 = \frac{m}{k}$   
 $\frac{\frac{\partial^2 T_i^{k+1} - T_i^k}{\Delta t}}{\Delta t} = \frac{1}{\Delta t^2} \left( [T_{i+1}^{k+1} - 2T_i^{k+1} + T_{i-1}^{k+1}] - [T_{i+1}^k - 2T_i^k + T_{i-1}^k] \right) - \frac{\lambda^2}{\Delta t} (T_i^{k+1} - T_i^k)$   
 $2\Delta t T_i^{k+1} - \Delta t (T_{i+1}^{k+1} - 2T_i^{k+1} + T_{i-1}^{k+1}) + \frac{1}{\Delta t} \Delta t \lambda^2 T_i^{k+1} = 2\Delta t T_i^k + \Delta t (T_{i+1}^k - 2T_i^k + T_{i-1}^k) - \frac{1}{\Delta t} \Delta t \lambda^2 T_i^k$   
 $-\Delta t T_{i+1}^{k+1} + (2\Delta t^2 + 2\Delta t + \frac{1}{\Delta t} \Delta t \lambda^2) T_i^{k+1} - \Delta t T_{i-1}^{k+1} = \Delta t T_{i+1}^k + (2\Delta t^2 - 2\Delta t - \frac{1}{\Delta t} \Delta t \lambda^2) T_i^k + \Delta t T_{i-1}^k$   
 Want  $A \vec{T}^{k+1} = B \vec{T}^k + c$ .  
 $r = 2\Delta t^2 + 2\Delta t + \frac{1}{\Delta t} \Delta t \lambda^2$        $a_0 = 2\Delta t^2 - 2\Delta t - \frac{1}{\Delta t} \Delta t \lambda^2$   
 $A = \begin{bmatrix} r & -\Delta t & 0 \\ -\Delta t & r & -\Delta t \\ 0 & -\Delta t & r \end{bmatrix}, \quad \vec{T}^k = \begin{bmatrix} -\Delta t T(0, t) \\ 0 \\ -\Delta t T(L, t) \end{bmatrix}, \quad \vec{T}^{k+1} = \begin{bmatrix} q & \Delta t & 0 \\ \Delta t & q & \Delta t \\ 0 & \Delta t & q \end{bmatrix} \vec{T}^k + \begin{bmatrix} \Delta t T(0, t) \\ 0 \\ \Delta t T(L, t) \end{bmatrix}$   
 $A \vec{T}^{k+1} = B \vec{T}^k + \begin{bmatrix} 2\Delta t T(0, t) \\ 0 \\ 2\Delta t T(L, t) \end{bmatrix}, \quad \vec{T}^0 = \begin{bmatrix} T(x_0) \\ \vdots \\ T(x_N) \end{bmatrix} = \begin{bmatrix} -5 \\ \vdots \\ -5 \end{bmatrix}$

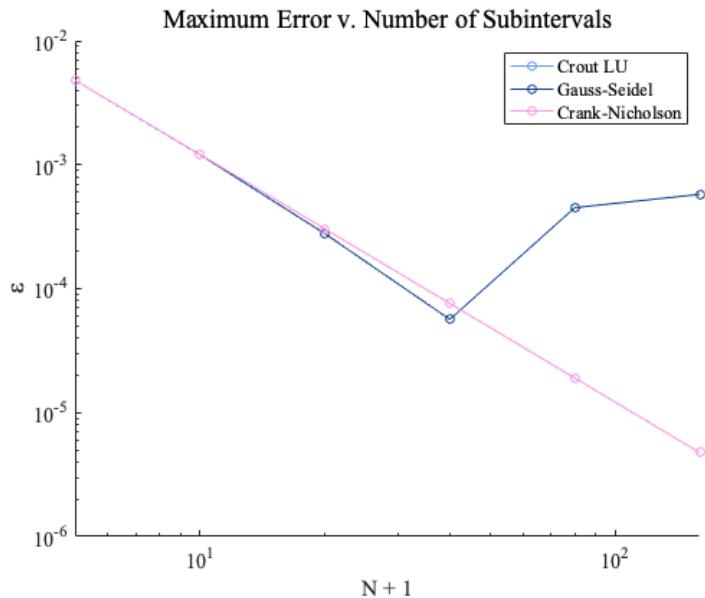


**Figure 1:** The transient or steady-state solution to the bioheat equation with the time derivative term, using the Crank-Nicholson time-stepping algorithm. To ensure stability, I used  $\Delta t = 1/(4n) \ll h/2$ . I used an  $L_\infty$  absolute norm between the current previous solutions with a tolerance of  $10^{-8}$  as my stopping criterion. Larger values of  $N$  not only required a smaller  $\Delta t$ , but also more time-steps to reach steady-state. The solution for  $N + 1 = 5$  required 27 iterations with  $\Delta t = 1/16$ , fewer iterations but a larger step size than the solution for  $N + 1 = 160$ , which used 803 iterations and  $\Delta t = 1/640$ .



**Figure 2:** The error behavior of the steady-state Crank-Nicholson solution agrees with the error behavior of the finite difference solution using Crout LU decomposition.

## Additional Figures



**Figure 1:** The Crout LU and Crank-Nicholson errors overlap. The dark blue curve is the error associated with the Gauss-Seidel iteration with a tolerance of  $10^{-6}$ , which clearly differs from Crout LU and Crank-Nicholson. For larger  $N$ , Crout LU and Crank-Nicholson lead to more accurate solutions than Gauss-Seidel, but this is somewhat expected. Crout LU and Crank-Nicholson both solve an approximating system exactly, whereas the Gauss-Seidel solution is an approximation of an approximating system.