

11-1-1 [OQU] Using the Maxwell's relation and the equation of state, determine a relation for the partial of s with respect to v at constant T for the IG model. Verify the relation using TEST at 100 kPa and 300 K.

SOLUTION:

The ideal gas equation of state can be expressed as $p = \frac{RT}{v}$. Then,

$$\left(\frac{\partial p}{\partial T} \right)_v = \frac{R}{v}$$

From the third Maxwell relation,

$$\left(\frac{\partial s}{\partial v} \right)_T = \left(\frac{\partial p}{\partial T} \right)_v \Rightarrow \left(\frac{\partial s}{\partial v} \right)_T = \frac{R}{v}$$

TEST Solution:

Launch the IG system-state TESTcalc and select Air. Evaluate State-1 from the given conditions and State-2 as a neighboring state with $p_2 = p_1 * 1.01$ (1 percent change) and $T_2 = T_1$ (holding T constant). In the I/O panel, calculate the LHS, $(s_2 - s_1)/(v_2 - v_1)$, as **0.335** unit. Using the average value of specific volume, the RHS is calculated as $2 * R_1 / (v_1 + v_2) = \mathbf{0.335}$ unit. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

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11-1-2 [OQX] Derive the relation for the slope of the $n = \text{constant}$ lines on a T - p diagram for a gas that obeys the van der Waals equation of state.

SOLUTION:

The **van der Waals equation** of state:

$\left(p + \frac{a}{v^2}\right)(v - b) = RT \left[\frac{\text{kJ}}{\text{kg}}\right]$, where, $a = \frac{27R^2T_{cr}^2}{64p_{cr}} \left[\text{m}^6\right]$; and $b = \frac{RT}{8p_{cr}} \left[\text{m}^3\right]$ are material constants.

$$T = \left(p + \frac{a}{v^2}\right) \frac{(v - b)}{R}$$

$$\Rightarrow \left(\frac{\partial T}{\partial p}\right)_v = \frac{(v - b)}{R} \frac{\partial}{\partial p} \left(p + \frac{a}{v^2}\right) = \frac{(v - b)}{R}$$

Therefore,

$$\left(\frac{\partial p}{\partial T}\right)_v = \frac{1}{\left(\frac{\partial T}{\partial p}\right)_v} = \frac{R}{(v - b)}$$

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11-1-3 [OQC] Derive the relation for the volume expansivity (β) and the isothermal compressibility (κ_T) for (a) an ideal gas and (b) a gas whose equation of state is $p(v - b) = RT$.

SOLUTION:

The volume expansivity and isothermal compressibility are expressed as

$$\beta = \frac{1}{v} \left(\frac{\partial v}{\partial T} \right)_p \quad \text{and} \quad \kappa_T = -\frac{1}{v} \left(\frac{\partial v}{\partial p} \right)_T$$

(a) For ideal gas, $v = \frac{RT}{p}$

$$\left(\frac{\partial v}{\partial T} \right)_p = \frac{R}{p}; \quad \Rightarrow \beta = \frac{1}{v} \left(\frac{R}{p} \right); \quad \Rightarrow \beta = \frac{1}{T}$$

$$\left(\frac{\partial v}{\partial p} \right)_T = -\frac{RT}{p^2}; \quad \Rightarrow \kappa_T = -\frac{1}{v} \left(-\frac{RT}{p^2} \right); \quad \Rightarrow \kappa_T = \frac{1}{p}$$

(b) For gas whose equation of state is $v = \frac{RT}{p} + b$,

$$\left(\frac{\partial v}{\partial T} \right)_p = \frac{R}{p}; \quad \Rightarrow \beta = \frac{1}{v} \left(\frac{R}{p} \right); \quad \Rightarrow \beta = \frac{R}{RT + bp}$$

$$\left(\frac{\partial v}{\partial p} \right)_T = -\frac{RT}{p^2}; \quad \Rightarrow \kappa_T = -\frac{1}{v} \left(-\frac{RT}{p^2} \right); \quad \Rightarrow \kappa_T = \frac{v - b}{pv}$$

11-1-4 [QQV] Estimate the volume expansivity (β) and the isothermal compressibility (κ_T) of refrigerant-134a at 200 kPa and 30°C.

SOLUTION:

The volume expansivity and isothermal compressibility are expressed as

$$\beta = \frac{1}{v} \left(\frac{\partial v}{\partial T} \right)_p \quad \text{and} \quad \kappa_T = -\frac{1}{v} \left(\frac{\partial v}{\partial p} \right)_T$$

Approximating differentials by differences about specific state (properties obtained using the manual approach discussed in Chapter 3 or the PC system-state TESTcalc),

$$\begin{aligned} \beta &\cong \frac{1}{v} \left(\frac{\Delta v}{\Delta T} \right)_{p=200\text{kPa}} \\ \Rightarrow \beta &= \frac{1}{v} \left(\frac{v_{40^\circ\text{C}} - v_{20^\circ\text{C}}}{(40 - 20)^\circ\text{C}} \right)_{p=200\text{kPa}} ; \quad \Rightarrow \beta = \frac{1}{0.11856} \left(\frac{0.12311 - 0.11394}{20} \right)_{p=200\text{kPa}} \\ \Rightarrow \beta &= \mathbf{0.00387 \text{ K}^{-1}} \end{aligned}$$

$$\begin{aligned} \kappa_T &\cong -\frac{1}{v} \left(\frac{\Delta v}{\Delta p} \right)_T \\ \kappa_T &= -\frac{1}{v} \left(\frac{v_{240\text{kPa}} - v_{180\text{kPa}}}{240 - 180} \right)_{T=30^\circ\text{C}} ; \quad \Rightarrow \kappa_T = -\frac{1}{0.11856} \left(\frac{0.09794 - 0.13230}{60} \right) ; \\ \Rightarrow \kappa_T &= \mathbf{0.00483 \text{ kPa}^{-1}} \end{aligned}$$

TEST Solution:

Launch the PC system-state TESTcalc and select R-134a. Evaluate State-1 at the given conditions, State-2 as a neighboring state with $p_2 = p_1$ and $T_2 = T_1 + 1$, and State-3 as another neighboring state with $T_2 = T_1$, $p_2 = 1.01 \cdot p_1$. In the I/O panel, evaluate $\beta = (1/v_1) \cdot (v_2 - v_1) / (T_2 - T_1)$ as **0.0037** and $(-1/v_1) \cdot (v_3 - v_1) / (p_3 - p_1)$ as **0.0051**. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

11-1-5 [OQQ] Estimate the Joule-Thomson coefficient (μ_J) of nitrogen at (a) 200 psia, 500°R and (b) 2000 psia, 400°R.

SOLUTION:

(a) Use the manual approach or PC system-state TESTcalc to obtain nitrogen properties.

State-1 (given $p_1 = 200 \text{ psia} = 1379 \text{ kPa}$; $T_1 = 500 \text{ R} = 277.77 \text{ K}$)

$$h_1 = -24.8 \frac{\text{kJ}}{\text{kg}}$$

State-2 (given $p_2 = (1.01)p_1$; $h_2 = h_1$)

$$T_2 = 277.81 \text{ K}$$

Using forward differencing to approximate the partial derivative,

$$\mu = \left(\frac{\partial T}{\partial p} \right)_h \cong \left(\frac{\Delta T}{\Delta p} \right)_h = \left(\frac{277.81 - 277.77}{1392.8 - 1379} \right) = 0.00258 \frac{\text{K}}{\text{kPa}} = 0.032 \frac{\text{R}}{\text{psia}}$$

(b)

State-1 (given $p_1 = 2000 \text{ psia} = 13790 \text{ kPa}$; $T_1 = 400 \text{ R} = 222.22 \text{ K}$)

$$h_1 = -129.25 \frac{\text{kJ}}{\text{kg}}$$

State-2 (given $p_2 = (1.01)p_1$; $h_2 = h_1$)

$$T_2 = 222.42 \text{ K}$$

Using forward differencing to approximate the partial derivative,

$$\mu = \left(\frac{\partial T}{\partial p} \right)_h \cong \left(\frac{\Delta T}{\Delta p} \right)_h = \left(\frac{222.42 - 222.22}{13928 - 13790} \right) = 0.00143 \frac{\text{K}}{\text{kPa}} = 0.0178 \frac{\text{R}}{\text{psia}}$$

TEST Solution:

Launch the PC flow-state TESTcalc and select N2. Evaluate State-1 at the given conditions, State-2 as a neighboring state with $p_2 = 1.01 \cdot p_1$ and $h_2 = h_1$. In the I/O panel, evaluate $\mu = (T_2 - T_1)/(p_2 - p_1)$. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

11-1-6 [OQT] For $\beta > 0$, prove that at every point of a single-phase region of an h - s diagram, the slope of a constant-pressure ($p = \text{constant}$) line is greater than the slope of a constant temperature ($T = \text{constant}$) line, but less than the slope of a constant-volume ($\beta = \text{constant}$) line.

SOLUTION:

Using Tds relation: $dh = Tds + vdp$; $\Rightarrow \frac{dh}{ds} = T + \frac{dp}{ds}$

$p = c$ line: $\left(\frac{\partial h}{\partial s}\right)_p = T$

$T = c$ line: $\left(\frac{\partial h}{\partial s}\right)_T = T + v\left(\frac{\partial p}{\partial s}\right)_T$

The 4th Maxwell relation: $\left(\frac{\partial p}{\partial s}\right)_T = -\left(\frac{\partial T}{\partial v}\right)_p$

Substituting, $\left(\frac{\partial h}{\partial s}\right)_T = T - v\left(\frac{\partial T}{\partial v}\right)_p = T - \frac{1}{\beta}$

Since $\beta > 0$, the slope of $p = c$ line is greater than the slope of $T = c$ line

$v = c$ line: $\left(\frac{\partial h}{\partial s}\right)_v = T + v\left(\frac{\partial p}{\partial s}\right)_v$ (1)

From the first Tds relation: $ds = \frac{c_v}{T}dT + \left(\frac{p}{T}\right)dv$

Divide by dp holding v constant: $\left(\frac{\partial s}{\partial p}\right)_v = \frac{c_v}{T}\left(\frac{\partial T}{\partial p}\right)_v \Rightarrow \left(\frac{\partial p}{\partial s}\right)_v = \frac{T}{c_v}\left(\frac{\partial p}{\partial T}\right)_v$ (2)

Using properties p , T , and v in the cyclic relation,

$\left(\frac{\partial p}{\partial T}\right)_v \left(\frac{\partial T}{\partial v}\right)_p \left(\frac{\partial v}{\partial p}\right)_T = -1$; $\Rightarrow \left(\frac{\partial p}{\partial T}\right)_v = -\left(\frac{\partial v}{\partial T}\right)_p \left(\frac{\partial p}{\partial v}\right)_T$;
 $\Rightarrow \left(\frac{\partial p}{\partial T}\right)_v = (-\beta v) \left(\frac{1}{-\alpha v}\right)$; $\Rightarrow \left(\frac{\partial p}{\partial T}\right)_v = \frac{\beta}{\alpha}$ (3)

Substituting (2) and (3) in (1)

$\left(\frac{\partial h}{\partial s}\right)_v = T + v\left(\frac{\partial p}{\partial s}\right)_v$; $\Rightarrow \left(\frac{\partial h}{\partial s}\right)_v = T + \frac{T\beta v}{C_v\alpha} > T$

The second term on right is always positive quantity since β is given to be positive. Therefore, the slope of $p = c$ line is less than the slope of $v = c$ line.

11-1-7 [OQY] Starting with the relation $dh = Tds + vdp$, show that the slope of a constant-pressure line a h - s diagram (a) is constant in the saturation region, and (b) increases with temperature in the superheated region.

SOLUTION:

For $p = c$ line, $dp = 0$ and the above relation reduces to $dh = Tds$, which can also be expressed as

$$\left(\frac{\partial h}{\partial s} \right)_p = T$$

Thus the slope of the $p = c$ line on an h - s diagram is equal to the temperature.

- (a) In the saturation region, $T = T_{\text{sat}@p}$ is a constant for a given pressure. Therefore, the slope is constant..
- (b) In the superheat region, the slope increases with increasing temperature since the slope is equal temperature.

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11-1-8 [OQF] Derive relation for (a) Δu , (b) Δh and (c) Δs of a gas that obeys the equation of state $(p + a/v^2)v = RT$ for an isothermal process.

SOLUTION:

(a)

$$\Delta u = u_2 - u_1 = \int_{T_1}^{T_2} c_v dT + \int_{v_1}^{v_2} \left(T \left(\frac{\partial p}{\partial T} \right)_v - p \right) dv$$

For isothermal ($T = c$) process:

$$\Delta u = \int_{v_1}^{v_2} \left(T \left(\frac{\partial p}{\partial T} \right)_v - p \right) dv$$

From the given equation of state,

$$p = \frac{RT}{v} - \frac{a}{v^2}; \Rightarrow \left(\frac{\partial p}{\partial T} \right)_v = \frac{R}{v}$$

Thus,

$$T \left(\frac{\partial p}{\partial T} \right)_v - p = \frac{RT}{v} - \frac{RT}{v} + \frac{a}{v^2}; \Rightarrow T \left(\frac{\partial p}{\partial T} \right)_v - p = \frac{a}{v^2}$$

$$\text{Substituting, } \Delta u = \int_{v_1}^{v_2} \frac{a}{v^2} dv; \Rightarrow \Delta u = a \left(\frac{1}{v_1} - \frac{1}{v_2} \right)$$

(b) The enthalpy change Δh is related to Δu through the relation

$$\Delta h = \Delta u + p_2 v_2 - p_1 v_1;$$

Thus,

$$p_2 v_2 - p_1 v_1 = \left(RT - \frac{a}{v_2} \right) - \left(RT - \frac{a}{v_1} \right); \Rightarrow p_2 v_2 - p_1 v_1 = a \left(\frac{1}{v_1} - \frac{1}{v_2} \right)$$

Substituting,

$$\Delta h = 2a \left(\frac{1}{v_1} - \frac{1}{v_2} \right)$$

(c)

$$\Delta s = s_2 - s_1;$$

$$\Delta s = \int_{T_1}^{T_2} \frac{c_v}{T} dT + \int_{v_1}^{v_2} \left(\frac{\partial p}{\partial T} \right)_v dv; \Rightarrow \Delta s = \int_{v_1}^{v_2} \left(\frac{\partial p}{\partial T} \right)_v dv$$

$$\text{Substituting } \left(\frac{\partial p}{\partial T} \right)_v = \frac{R}{v}, \text{ (derived in part (a))}$$

$$\Delta s = \int_{v_1}^{v_2} \frac{R}{v} dv; \Rightarrow \Delta s = R \ln \frac{v_2}{v_1}$$

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11-1-9 [OQD] Show that $c_v = -T(\partial v/\partial T)_s(\partial p/\partial T)_v$, and $c_p = T(\partial p/\partial T)_s(\partial v/\partial T)_p$.

SOLUTION:

From Eq. (11.38)

$$c_v = T \left(\frac{\partial s}{\partial T} \right)_v = T \left(\frac{\partial s}{\partial p} \right)_v \left(\frac{\partial p}{\partial T} \right)_v$$

Substituting the first Maxwell relation $\left(\frac{\partial s}{\partial p} \right)_v = - \left(\frac{\partial v}{\partial T} \right)_s$

$$c_v = -T \left(\frac{\partial v}{\partial T} \right)_s \left(\frac{\partial p}{\partial T} \right)_v$$

Using the definition of c_p from Eq. (11.46)

$$c_p = T \left(\frac{\partial s}{\partial T} \right)_p = T \left(\frac{\partial s}{\partial v} \right)_p \left(\frac{\partial v}{\partial T} \right)_p$$

Substituting the second Maxwell relation $\left(\frac{\partial s}{\partial v} \right)_p = \left(\frac{\partial p}{\partial T} \right)_s$

$$c_p = T \left(\frac{\partial p}{\partial T} \right)_s \left(\frac{\partial v}{\partial T} \right)_p$$

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11-1-10 [OQM] Steam is throttled from 4.5 MPa and 400°C to 3.5 MPa. Estimate the temperature change (ΔT) of the steam during this process and the average Joule-Thomson coefficient (μ_J).

SOLUTION:

Using the manual approach or the PC system-state TESTcalc, evaluate the states.

State-1 (given $p_1 = 4.5$ MPa; $T_1 = 400^\circ\text{C}$)

$$h_1 = 3204.7 \text{ kJ/kg.}$$

State-2 (given $p_2 = 3.5$ MPa; $h_2 = h_1$)

$$T_2 = 392.56^\circ\text{C}$$

Thus the temperature drop during this throttling process is

$$\Delta T = T_2 - T_1 = 392.56 - 400 = -7.44^\circ\text{C}$$

The average Joule-Thomson coefficient for this process is determined from

$$\mu = \left(\frac{\partial T}{\partial p} \right)_h; \Rightarrow \mu \approx \left(\frac{\Delta T}{\Delta p} \right)_{h=3204.7 \text{ kJ/kg}}$$
$$\Rightarrow \mu = \frac{(392.56 - 400)^\circ\text{C}}{(3.5 - 4.5) \text{ MPa}}; \Rightarrow \mu = 7.44^\circ\text{C/MPa}$$

TEST Solution:

Launch the PC flow-state TESTcalc and select H₂O. Evaluate State-1 and State-2 from the given conditions and using the fact that throttling is isenthalpic ($h_2 = h_1$). In the I/O panel, evaluate $=(T_2-T_1)/(p_2-p_1)$. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

11-1-11 [OQJ] Consider an infinitesimal reversible adiabatic compression or expansion process. By taking $s = s(P, v)$ and using the Maxwell relations, show that for this process $Pv^k = \text{constant}$, where k is isentropic expansion exponent defined as $k = (v/P)(\partial P/\partial v)_s$. Also, show that the isentropic expansion exponent k reduces to the specific heat ratio (C_p/C_v) for an ideal gas.

SOLUTION:

Using Taylor's theorem, the total differential ds of $s = s(p, v)$ can be expressed as

$$ds = \left(\frac{\partial s}{\partial p} \right)_v dp + \left(\frac{\partial s}{\partial v} \right)_p dv; \quad (1)$$

Now substituting following Maxwell relations

$$\left(\frac{\partial s}{\partial p} \right)_v = - \left(\frac{\partial v}{\partial T} \right)_s \quad \text{and} \quad \left(\frac{\partial s}{\partial v} \right)_p = \left(\frac{\partial p}{\partial T} \right)_s$$

into Eq. (1) and setting $ds = 0$ (isentropic condition), we get

$$\begin{aligned} - \left(\frac{\partial v}{\partial T} \right)_s dp + \left(\frac{\partial p}{\partial T} \right)_s dv &= 0; \\ \Rightarrow dp - \left(\frac{\partial T}{\partial v} \right)_s \left(\frac{\partial p}{\partial T} \right)_s dv &= 0; \quad \Rightarrow dp - \left(\frac{\partial p}{\partial v} \right)_s dv = 0; \\ \Rightarrow \frac{dp}{p} - \frac{1}{p} \left(\frac{\partial p}{\partial v} \right)_s dv &= 0; \end{aligned} \quad (2)$$

By definition, the isentropic expansion exponent k is:

$$k = - \frac{v}{p} \left(\frac{\partial p}{\partial v} \right)_s$$

Substituting in equation (2)

$$\frac{dp}{p} + k \frac{dv}{v} = 0$$

Taking k constant and integrating,

$$\ln p + k \ln v = c; \quad \Rightarrow \ln pv^k = c$$

Thus, $Pv^k = c$

To show that $k = \frac{C_p}{C_v}$ for an ideal gas, we write the cyclic relations for the following two groups of variables:

$$(s, T, v) \rightarrow \left(\frac{\partial s}{\partial T} \right)_v \left(\frac{\partial v}{\partial s} \right)_T \left(\frac{\partial T}{\partial v} \right)_s = -1; \Rightarrow \frac{c_v}{T} \left(\frac{\partial v}{\partial s} \right)_T \left(\frac{\partial T}{\partial v} \right)_s = -1 \quad (3)$$

$$(s, T, p) \rightarrow \left(\frac{\partial s}{\partial T} \right)_p \left(\frac{\partial p}{\partial s} \right)_T \left(\frac{\partial T}{\partial p} \right)_s = -1; \Rightarrow \frac{c_p}{T} \left(\frac{\partial p}{\partial s} \right)_T \left(\frac{\partial T}{\partial p} \right)_s = -1 \quad (4)$$

where we use Eqs. (11.38) and (11.46):

$$c_v = T \left(\frac{\partial s}{\partial T} \right)_v \quad \text{and} \quad c_p = T \left(\frac{\partial s}{\partial T} \right)_p$$

Equating the left hand sides of Eqs (3) and (4), we obtain:

$$\frac{c_p}{c_v} = \left(\frac{\partial s}{\partial p} \right)_T \left(\frac{\partial p}{\partial T} \right)_s \left(\frac{\partial v}{\partial s} \right)_T \left(\frac{\partial T}{\partial v} \right)_s = \left(\frac{\partial s}{\partial p} \frac{\partial v}{\partial s} \right)_T \left(\frac{\partial p}{\partial T} \frac{\partial T}{\partial v} \right)_s = \left(\frac{\partial v}{\partial p} \right)_T \left(\frac{\partial p}{\partial v} \right)_s;$$

For an ideal gas

$$\left(\frac{\partial v}{\partial p} \right)_T = \left(\frac{\partial \left(\frac{RT}{p} \right)}{\partial p} \right)_T; \Rightarrow \left(\frac{\partial v}{\partial p} \right)_T = -\frac{v}{p}$$

Substituting,

$$\begin{aligned} \frac{c_p}{c_v} &= -\frac{v}{p} \left(\frac{\partial p}{\partial v} \right)_s; \\ \Rightarrow k &= -\frac{v}{p} \left(\frac{\partial p}{\partial v} \right)_s = \frac{c_p}{c_v} \end{aligned}$$

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11-1-12 [QW] Consider a mixture of two gases A and B. Show that when the mass fraction x_A and x_B are known, the mole-fraction can be determined from $y_A = M_B/[M_A(1/x_A - 1) + M_B]$ and $y_B = 1 - y_A$ where M_A and M_B are the molar masses of A and B.

SOLUTION:

Starting from the definition of mole fraction:

$$\begin{aligned}
 y_A &= \frac{n_A}{n} = \frac{n_A}{n_A + n_B} = \frac{\left(\frac{m_A}{\bar{M}_A}\right)}{\left(\frac{m_A}{\bar{M}_A}\right) + \left(\frac{m_B}{\bar{M}_B}\right)} = \frac{\left(\frac{m_A}{m} \frac{1}{\bar{M}_A}\right)}{\left(\frac{m_A}{m} \frac{1}{\bar{M}_A}\right) + \left(\frac{m_B}{m} \frac{1}{\bar{M}_B}\right)} \\
 \Rightarrow y_A &= \frac{\left(x_A \frac{1}{\bar{M}_A}\right)}{\left(x_A \frac{1}{\bar{M}_A}\right) + \left(x_B \frac{1}{\bar{M}_B}\right)} = \frac{\left(\frac{x_A}{\bar{M}_A}\right)}{\left(\frac{x_A}{\bar{M}_A}\right) + \left(\frac{1-x_A}{\bar{M}_B}\right)} = \frac{(\bar{M}_A \bar{M}_B) \left(\frac{x_A}{\bar{M}_A}\right)}{(\bar{M}_A \bar{M}_B) \left[\left(\frac{x_A}{\bar{M}_A}\right) + \left(\frac{1-x_A}{\bar{M}_B}\right)\right]} \\
 \Rightarrow y_A &= \frac{\bar{M}_B x_A}{\bar{M}_B x_A + \bar{M}_A (1-x_A)} = \frac{\bar{M}_B}{\bar{M}_B + \bar{M}_A \left(\frac{1}{x_A} - 1\right)}
 \end{aligned}$$

$$n = n_A + n_B;$$

$$\Rightarrow 1 = \frac{n_A}{n} + \frac{n_B}{n};$$

$$\Rightarrow 1 = y_A + y_B;$$

$$\Rightarrow y_B = 1 - y_A$$

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11-1-13 [OTN] Nitrogen gas at 400 K and 300 kPa behaves as an ideal gas. Estimate the c_p and c_v of nitrogen at this state.

TEST Solution

Launch the IG system-state TESTcalc and select N2. Evaluate the neighboring states (1 deg-C difference) as described in the TEST-code, which can be found in the problems module of the TEST-pro site at www.thermofluids.net. Note that for State-2, the pressure is held constant ($p_2=p_1$) and for State-3 the specific volume is held constant ($v_3=v_1$). In the I/O panel, calculate c_p from $\Delta(h_2-h_1)/(T_2-T_1)$ and c_v from $\Delta(u_2-u_1)/(T_3-T_1)$ to produce

(a) 1.048 kJ/kg-K, (b) 0.7514 kJ/kg-K

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11-1-14 [OTR] A system contains oxygen (ideal gas) at 400 K and 100 kPa. As a result of some disturbance, the conditions of the gas change to 404 K and 98 kPa. (a) Estimate the change in the specific volume (Δv) of the gas using the ideal-gas relation and using Taylor's theorem. (b) Determine the exact answer using the IG equation of state.

Solution

(a) Taylor's theorem for $v = v(p, T)$ can be expressed as:

$$\Delta v \approx dv = \left(\frac{\partial v}{\partial p} \right)_T dp + \left(\frac{\partial v}{\partial T} \right)_p dT;$$

$$\text{For oxygen: } R = \frac{\bar{R}}{\bar{M}_{O_2}} = \frac{8.314}{32} = 0.260$$

The IG relation produces:

$$v = \frac{RT}{p};$$

$$\Rightarrow \left(\frac{\partial v}{\partial p} \right)_T = -\frac{RT}{p^2} = -\frac{(0.260)(400)}{99^2} = -0.01066 \text{ unit}$$

$$\text{And } \left(\frac{\partial v}{\partial T} \right)_p = \frac{R}{p} = \frac{0.260}{99} = 0.0026263 \text{ unit}$$

$$\text{Therefore, } \Delta v \approx dv = \left(\frac{\partial v}{\partial p} \right)_T dp + \left(\frac{\partial v}{\partial T} \right)_p dT = (-0.01066)(-2) + (0.0026263)(4) = 0.0318 \frac{\text{m}^3}{\text{kg}}$$

(b) The exact change in v can be calculated from the known IG relation.

$$\Delta v = v_2 - v_1 = R \left(\frac{T_2}{p_2} - \frac{T_1}{p_1} \right) = (0.260) \left(\frac{404}{100} - \frac{400}{98} \right) = 0.0318;$$

$$\Rightarrow \Delta p = \frac{1}{3.488} \left(\frac{355}{1.01} - \frac{350}{1} \right) = 0.0318 \frac{\text{m}^3}{\text{kg}}$$

TEST Solution

Launch the IG system-state TESTcalc and select O2. Evaluate the neighboring states from the given conditions as described in the TEST-code, which can be found in the problems module of the TEST-pro site at www.thermofluids.net. Evaluate the derivative in the I/O panel using central or forward differencing.

11-1-15 [OTS] Estimate the specific-heat difference ($c_p - c_v$) for liquid water at 20 MPa and 60°C.

SOLUTION:

The specific heat difference is given as

$$c_p - c_v = -T \left(\frac{\partial v}{\partial T} \right)_p^2 \left(\frac{\partial p}{\partial v} \right)_T$$

Approximating differentials by difference about the specific state and obtaining the properties using the system-state PC TESTcalc,

$$\begin{aligned} c_p - c_v &\cong -T \left(\frac{\Delta v}{\Delta T} \right)_{p=20 \text{ MPa}}^2 \left(\frac{\Delta p}{\Delta v} \right)_{T=60^\circ \text{ C}} ; \\ \Rightarrow c_p - c_v &= -(60 + 273.15) \left(\frac{v_{65^\circ \text{ C}} - v_{55^\circ \text{ C}}}{65 - 55} \right)_{p=20 \text{ MPa}}^2 \left(\frac{30000 - 10000}{v_{30 \text{ MPa}} - v_{10 \text{ MPa}}} \right)_{T=60^\circ \text{ C}} ; \\ \Rightarrow c_p - c_v &= -(333.15) (5.0 \times 10^{-7})^2 (-2.063 \times 10^9) ; \\ \Rightarrow c_p - c_v &= \mathbf{0.172 \text{ kJ/kg} \cdot \text{K}} \end{aligned}$$

TEST Solution

Launch the PC system-state TESTcalc and select H2O. Evaluate the four neighboring states as described in the TEST-code, which can be found in the problems module of the TEST-pro site at www.thermofluids.net. Calculate the partial derivatives in the I/O panel. Note that evaluating c_v directly is not possible due to high error associated with finding neighboring states at constant volume for a liquid.

11-1-16 [OTO] Estimate the specific-heat difference ($c_p - c_v$) for liquid water at 1000 psia and 150°F.

SOLUTION:

The specific heat difference is given as

$$c_p - c_v = -T \left(\frac{\partial v}{\partial T} \right)_p^2 \left(\frac{\partial p}{\partial v} \right)_T$$

Approximating differentials by difference about the specific state and obtaining the properties using the system-state PC TESTcalc (note that the pressure difference must be huge in order to create any appreciable change in the specific volume),

$$c_p - c_v \cong -T \left(\frac{\Delta v}{\Delta T} \right)_{p=1000 \text{ psia}}^2 \left(\frac{\Delta p}{\Delta v} \right)_{T=150^\circ \text{ F}} \cong -T \left(\frac{\Delta v}{\Delta T} \right)_{p=7 \text{ MPa}}^2 \left(\frac{\Delta p}{\Delta v} \right)_{T=65^\circ \text{ C}} ;$$

$$\Rightarrow c_p - c_v = -(65 + 273) \left(\frac{v_{70^\circ \text{ C}} - v_{60^\circ \text{ C}}}{70 - 60} \right)_{p=7 \text{ MPa}}^2 \left(\frac{30000 - 1000}{v_{30 \text{ MPa}} - v_{1 \text{ MPa}}} \right)_{T=65^\circ \text{ C}} ;$$

$$\Rightarrow c_p - c_v = -(338) (6.0 \times 10^{-7})^2 (-3.13 \times 10^9) ;$$

$$\Rightarrow c_p - c_v = 0.381 \text{ kJ/kg} \cdot \text{K}$$

TEST Solution

Launch the PC system-state TESTcalc and select H2O. Evaluate the four neighboring states as described in the TEST-code, which can be found in the problems module of the TEST-pro site at www.thermofluids.net. Calculate the partial derivatives in the I/O panel. Note that evaluating c_v directly is not possible due to high error associated with finding neighboring states at constant volume for a liquid.

11-1-17 [OTB] Plot the Joule-Thomson coefficient (μ_J) for nitrogen over the pressure range of 100 psia to 1500 psia at the enthalpy values of 100 Btu/lbm, 175 Btu/lbm and 225 Btu/lbm. Discuss the results.

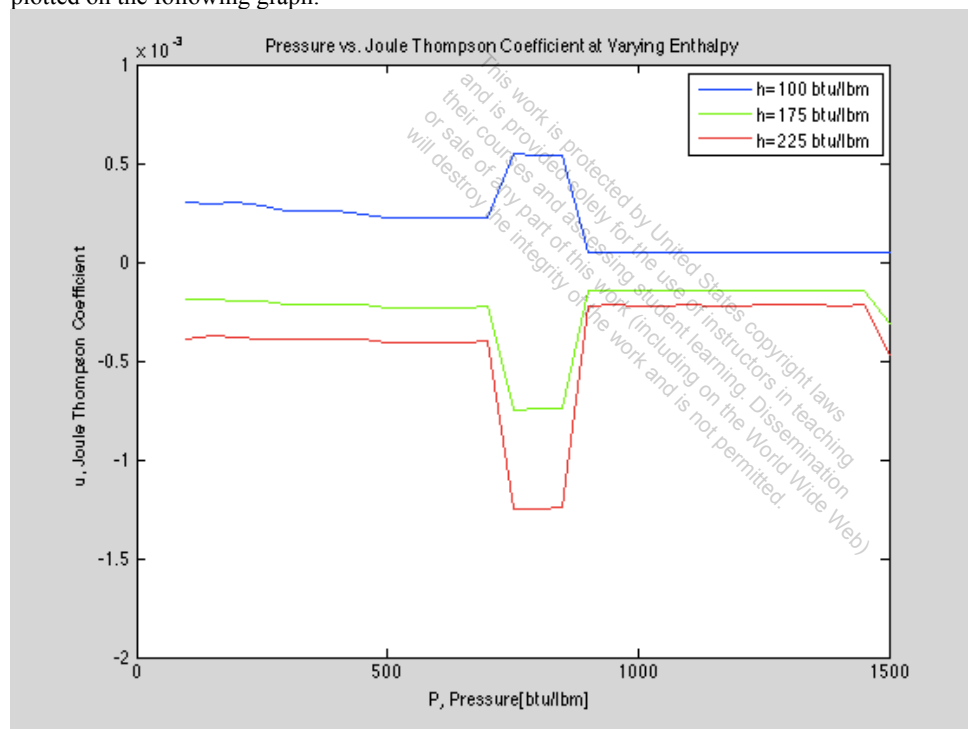
Analysis

The PC model for nitrogen is used to find the temperature under the given conditions. A one percent drop in pressure can be regarded as a differential change.

Approximating differentials by differences at specific state,

$$\mu_J \equiv \left(\frac{\partial T}{\partial p} \right)_h \approx \left[\frac{(T_2 - T_1)}{(p_2 - p_1)} \right]_h; \quad \left[\frac{^\circ\text{F}}{\text{psia}} \right]$$

Using the system-state PC TESTcalc, State-1 is found from the given pressure and enthalpy and a neighboring state at a slightly higher pressure ($p_2 = 1.01 \cdot p_1$) while holding enthalpy the same as the original state ($h_2 = h_1$). In the I/O panel, calculate the Joule-Thomson coefficient by evaluating $(T_2 - T_1)/(p_2 - p_1)$. Repeat with a different p_1 and h_1 as necessary. The results were plotted on the following graph:



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11-1-18 [OTA] Determine the enthalpy change (Δh) and the entropy change (Δs) of nitrogen per unit mole as it undergoes a change of state from 225 K and 6 MPa to 320 K and 12 MPa, (a) by assuming ideal-gas behavior, and (b) by accounting for the deviation from ideal-gas behavior through the use of generalized charts (L-K model).

SOLUTION:

From Table C-1 or the PG system-state TESTcalc, obtain the necessary material properties of

$$N_2: c_p = 1.031 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}, c_v = 0.734 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}, R = 0.2968 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

Given: $p_1 = 6 \text{ MPa}$; $T_1 = 225 \text{ K}$

Given: $p_2 = 12 \text{ MPa}$; $T_2 = 320 \text{ K}$

IG Model:

Using the manual approach described in Chapter 3 or the IG system-state TESTcalc, obtain:

$$\Delta h = h_2 - h_1 = 98.0 \frac{\text{kJ}}{\text{kg}}$$

$$\Delta s = s_2 - s_1 = 0.1575 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

RG Model (L-K):

From Table E-1, or from the RG system-state TESTcalc, obtain the necessary material properties of N_2 :

$$p_{cr} = 3.39 \text{ MPa}, T_{cr} = 126.2 \text{ K}$$

The reduced pressure and temperature for the given states are

$$p_{r1} = \frac{p_1}{p_{cr}} = 1.77, T_{r1} = \frac{T_1}{T_{cr}} = 1.78$$

$$p_{r2} = \frac{p_2}{p_{cr}} = 3.54, T_{r2} = \frac{T_2}{T_{cr}} = 2.536$$

Using the Lee-Kesler compressibility charts, Table E, we obtain the compressibility factors.

$$\Delta h = (h_2 - h_1)^{IG} - RT_{cr}(Z_{h,2} - Z_{h,1}) = 103.2 \frac{\text{kJ}}{\text{kg}}$$

$$\Delta s = (s_2 - s_1)^{IG} - R(Z_{s,2} - Z_{s,1}) = 0.180 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

TEST Solution

Launch the IG system-state TESTcalc and select N2. Evaluate the initial and final states from the known conditions and evaluate the property differences in the I/O panel. Generate the TEST-code using the Super-Calculate button. Launch the RG system-state TESTcalc in a separate tab, paste the TEST-code in the I/O panel, and click the Load button to calculate the states. Evaluate the property differences for the RG model in the I/O panel. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

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11-1-19 [OTH] Determine the enthalpy change (Δh) and the entropy change (Δs) of carbon dioxide per unit mass as it undergoes a change of state from 250 K and 7 MPa to 280 K and 12 MPa, (a) by assuming ideal-gas behavior, and (b) by accounting for the deviation from ideal-gas behavior.

SOLUTION:

From Table C-1 or the PG system-state TESTcalc, obtain the necessary material properties of

$$\text{CO}_2: R = 0.1889 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

$$\text{Given: } p_1 = 7 \text{ MPa; } T_1 = 250 \text{ K}$$

$$\text{Given: } p_2 = 12 \text{ MPa; } T_2 = 280 \text{ K}$$

IG Model:

Using the manual approach described in Chapter 3 or the IG system-state TESTcalc, obtain:

$$\Delta h = h_2 - h_1 = 25.3 \frac{\text{kJ}}{\text{kg}}$$

$$\Delta s = s_2 - s_1 = -0.0062 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

RG Model (L-K):

From Table E-1, or from the RG system-state TESTcalc, obtain the necessary material properties of CO_2 :

$$p_{cr} = 7.37 \text{ MPa, } T_{cr} = 304.1 \text{ K}$$

The reduced pressure and temperature for the given states are

$$p_{r1} = \frac{p_1}{p_{cr}} = 0.95, \quad T_{r1} = \frac{T_1}{T_{cr}} = 0.822$$

$$p_{r2} = \frac{p_2}{p_{cr}} = 1.63, \quad T_{r2} = \frac{T_2}{T_{cr}} = 0.921$$

Using the Lee-Kesler compressibility charts, Table E, we obtain the compressibility factors.

$$\Delta h = (h_2 - h_1)^{IG} - RT_{cr}(Z_{h,2} - Z_{h,1}) = 19.0 \frac{\text{kJ}}{\text{kg}}$$

$$\Delta s = (s_2 - s_1)^{IG} - R(Z_{s,2} - Z_{s,1}) = 2.28 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

TEST Solution

Launch the IG system-state TESTcalc and select CO₂. Evaluate the initial and final states from the known conditions and evaluate the property differences in the I/O panel. Generate the TEST-code using the Super-Calculate button. Launch the RG system-state TESTcalc in a separate tab, paste the TEST-code in the I/O panel, and click the Load button to calculate the states. Evaluate the property differences for the RG model in the I/O panel. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

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11-1-20 [OTE] Methane is compressed adiabatically by a steady-state flow compressor from 2 MPa and -10°C to 10 MPa and 110°C at a rate of 0.8 kg/s. Using the generalized charts, determine the required power input to the compressor.

SOLUTION:

The steady-flow energy balance equation for this compressor can be expressed as

$$\dot{W}_{\text{C,in}} = \dot{m}(h_2 - h_1)$$

From Table E-1, or from the RG system-state TESTcalc, obtain the necessary material properties of CH_4 :

$$p_{\text{cr}} = 4.6 \text{ MPa}, T_{\text{cr}} = 190.4 \text{ K}$$

The enthalpy departures of CH_4 at the specified states are determined from generalized charts to be

$$\left. \begin{aligned} T_{R1} &= \frac{T_1}{T_{\text{cr}}}; \Rightarrow T_{R1} = \frac{263}{191.1}; \Rightarrow T_{R1} = 1.376 \\ p_{R1} &= \frac{p_1}{p_{\text{cr}}}; \Rightarrow p_{R1} = \frac{2}{4.64}; \Rightarrow p_{R1} = 0.431 \end{aligned} \right\} \rightarrow Z_{h1} = 0.2$$

And

$$\left. \begin{aligned} T_{R2} &= \frac{T_2}{T_{\text{cr}}}; \Rightarrow T_{R2} = \frac{383}{191.1}; \Rightarrow T_{R2} = 2.00 \\ p_{R2} &= \frac{p_2}{p_{\text{cr}}}; \Rightarrow p_{R2} = \frac{10}{4.64}; \Rightarrow p_{R2} = 2.155 \end{aligned} \right\} \rightarrow Z_{h2} = 0.4$$

Thus,

$$\begin{aligned} h_2 - h_1 &= RT_{\text{cr}}(Z_{h1} - Z_{h2}) + (h_2 - h_1)_{\text{ideal}}; \\ \Rightarrow h_2 - h_1 &= (0.5182)(191.1)(0.2 - 0.4) + 2.2537(110 - (-10)); \\ \Rightarrow h_2 - h_1 &= 250.6 \text{ kJ/kg} \end{aligned}$$

Substituting,

$$\begin{aligned} \dot{W}_{\text{C,in}} &= (0.8 \text{ kg/s})(250.6 \text{ kJ/kg}); \\ \Rightarrow \dot{W}_{\text{C,in}} &= 201 \text{ kW} \end{aligned}$$

TEST Solution

Launch the RG single-flow, open-steady TESTcalc and select CH_4 . Evaluate the inlet and exit states from the known conditions. In the device panel, load the states, enter $\dot{Q}_{\text{dot}}=0$, and calculate $\dot{W}_{\text{dot_ext}}$. By default, the L-K model is selected, which produces a result of **198 kW**. When the N-O model is used, the answer changes to **203 kW**. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

11-1-21 [OTI] Methane gas flows through a pipeline with a mass flow rate of 110 lb/s at a pressure of 183 atm and a temperature of 56°F. Determine the volumetric flow rate (\dot{V}), in ft³/s, using (a) the ideal gas equation, (b) van der Waals equation and (c) compressibility chart.

Unknown
Deleted: \dot{V}

SOLUTION:

(a) The mass and volume flow rates are related by:

$$\dot{V} = \frac{\dot{m}}{\rho} = \dot{m}v$$

From Table C-1 or the IG flow-state TESTcalc, obtain the necessary material properties of CH₄:

$$R = 0.518 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

$$\text{Given: } p_1 = 183 \text{ atm} = 18.54 \text{ MPa}; T_1 = 56^\circ \text{F} = 286.5 \text{ K}, \dot{m} = 110 \frac{\text{lbm}}{\text{s}} = 49.9 \frac{\text{kg}}{\text{s}}$$

IG Model:

Using the manual approach described in Chapter 3 or the IG system-state TESTcalc, obtain:

$$v = \frac{RT}{p} = 0.0080 \frac{\text{m}^3}{\text{kg}}$$

$$\Rightarrow \dot{V} = \dot{m}v = 0.400 \frac{\text{m}^3}{\text{s}}$$

(b) van der Waals equation:

$$\left(p + \frac{a}{v^2}\right)(v - b) = RT$$

$$\text{where } a = \frac{27R^2T_{\text{cr}}^2}{64p_{\text{cr}}} = 0.892 \frac{\text{kPa} \cdot \text{m}^6}{\text{kg}^2}; b = \frac{RT_{\text{cr}}}{8p_{\text{cr}}} = 0.00268 \frac{\text{m}^3}{\text{kg}}$$

$$\Rightarrow \left(18540 + \frac{0.892}{v^2}\right)(v - 0.00268) = 0.518(286.5)$$

$$\Rightarrow (18540v^2 + 0.892)(v - 0.00268) = 148.4v^2$$

$$\Rightarrow v = 0.0072 \frac{\text{m}^3}{\text{kg}}$$

$$\Rightarrow \dot{V} = \dot{m}v = 0.362 \frac{\text{m}^3}{\text{s}}$$

(c) RG Model (L-K):

From Table E-1, or from the RG system-state TESTcalc, obtain the necessary material properties of CH₄:

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$$p_{cr} = 4.6 \text{ MPa}, T_{cr} = 190.4 \text{ K}$$

The reduced pressure and temperature for the given states are

$$p_{r1} = \frac{p_1}{p_{cr}} = 4.03, T_{r1} = \frac{T_1}{T_{cr}} = 1.50$$

Using the Lee-Kesler compressibility charts, Table E, we obtain the compressibility factor.

$$Z = 0.808$$

$$v = \frac{ZRT}{p} = 0.0065 \frac{\text{m}^3}{\text{kg}}$$

$$\Rightarrow \dot{V} = \dot{m}v = 0.324 \frac{\text{m}^3}{\text{s}}$$

TEST Solution

Launch the RG single-flow, open-steady TESTcalc and select CH4. Evaluate the state from the known conditions. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

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11-1-22 [OTL] Determine the specific volume (v) of water vapor at 10 MPa and 360°C, in m³/kg, using (a) the steam tables, (b) compressibility chart and (c) ideal gas equation.

SOLUTION:

From Table C-1 or the IG flow-state TESTcalc, obtain the necessary material properties of H₂O:

$$R = 0.461 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

Given: $p_1 = 10 \text{ MPa}$; $T_1 = 360^\circ\text{C} = 633 \text{ K}$

(a) PC Model:

Using the manual approach described in Chapter 3 or the PC system-state TESTcalc, obtain:

$$v = 0.0232 \frac{\text{m}^3}{\text{kg}}$$

(b) RG Model (L-K):

From Table E-1, or the RG system-state TESTcalc, obtain the necessary material properties of H₂O:

$$p_{cr} = 22.12 \text{ MPa}, T_{cr} = 647 \text{ K}$$

The reduced pressure and temperature for the given states are

$$p_{r1} = \frac{p_1}{p_{cr}} = 4.03, T_{r1} = \frac{T_1}{T_{cr}} = 1.50$$

Using the Lee-Kesler compressibility charts, Table E, we obtain the compressibility factor.

$$Z = 0.81$$

$$v = \frac{ZRT}{p} = 0.0237 \frac{\text{m}^3}{\text{kg}}$$

(c) IG Model:

Using the manual approach described in Chapter 3 or the IG system-state TESTcalc, obtain:

$$v = \frac{RT}{p} = 0.0292 \frac{\text{m}^3}{\text{kg}}$$

TEST Solution

Launch the PC system-state TESTcalc and select H₂O. Evaluate the state from the known conditions. Repeat with the IG and RG model. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

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11-1-23 [OTG] Consider refrigerant R-12 vapor at 160°F and 0.5 ft³/lb. Estimate the pressure (p) at this state, in atm, using (a) the ideal gas equation, (b) van der Waals equation and (c) compressibility chart.

From the PC system-state TESTcalc, obtain (by hovering the pointer over R-12 to display its properties on the message panel after selecting it) the necessary material properties of R-12 :

$$R = \frac{\bar{R}}{\bar{M}} = \frac{8.314}{120.93} = 0.0687 \frac{\text{kJ}}{\text{kg} \cdot \text{K}}$$

$$\text{Given: } v_1 = 0.5 \frac{\text{ft}^3}{\text{lbm}} = 0.0312 \frac{\text{ft}^3}{\text{lbm}}; T_1 = 160^\circ \text{F} = 344 \text{ K}$$

(a) IG Model:

Using the manual approach described in Chapter 3 or the IG system-state TESTcalc, obtain:

$$p = \frac{RT}{v} = \mathbf{758.2 \text{ kPa}}$$

(c) RG Model (L-K):

From Table E-1, or the RG system-state TESTcalc, obtain the necessary material properties of R-12 :

$$p_{cr} = 4.135 \text{ MPa}, T_{cr} = 385 \text{ K}$$

The reduced pressure and temperature for the given states are

$$p_{r1} = \frac{p_1}{p_{cr}} = 4.03, T_{r1} = \frac{T_1}{T_{cr}} = 1.50$$

Using the Lee-Kesler compressibility charts, Table E, we obtain the compressibility factor.

$$Z = 0.91$$

$$p = \frac{ZRT}{v} = \mathbf{692 \frac{m^3}{kg}}$$

TEST Solution

Launch the RG system-state TESTcalc and select H2O. Evaluate the state from the known conditions. Repeat with the PC model. The TEST-code for this problem can be found in the problems module of the TEST-pro site at www.thermofluids.net.

11-1-24 [OTP] For the functions $x = x(y, w)$, $y = y(z, w)$, $z = z(x, w)$, demonstrate that $(\partial x / \partial y)_w (\partial y / \partial z)_w (\partial z / \partial x)_w = 1$.

SOLUTION:

$$x = x(y, w)$$

$$dx = \left(\frac{\partial x}{\partial w} \right)_y dw + \left(\frac{\partial x}{\partial y} \right)_w dy \quad (1)$$

Similarly,

$$y = y(z, w)$$

$$dy = \left(\frac{\partial y}{\partial w} \right)_z dw + \left(\frac{\partial y}{\partial z} \right)_w dz$$

Substituting dy in equation (1)

$$dx = \left(\frac{\partial x}{\partial w} \right)_y dw + \left(\frac{\partial x}{\partial y} \right)_w \left[\left(\frac{\partial y}{\partial w} \right)_z dw + \left(\frac{\partial y}{\partial z} \right)_w dz \right]$$

$$\Rightarrow dx = \left[\left(\frac{\partial x}{\partial w} \right)_y + \left(\frac{\partial x}{\partial y} \right)_w \left(\frac{\partial y}{\partial w} \right)_z \right] dw + \left(\frac{\partial x}{\partial y} \right)_w \left(\frac{\partial y}{\partial z} \right)_w dz; \quad (2)$$

Again,

$$dx = \left(\frac{\partial x}{\partial w} \right)_z dw + \left(\frac{\partial x}{\partial z} \right)_w dz \quad (3)$$

Comparing dz component of equation (2) and (3),

$$\left(\frac{\partial x}{\partial z} \right)_w = \left(\frac{\partial x}{\partial y} \right)_w \left(\frac{\partial y}{\partial z} \right)_w$$

$$\left(\frac{\partial x}{\partial y} \right)_w \left(\frac{\partial y}{\partial z} \right)_w \left(\frac{\partial z}{\partial x} \right)_w = 1$$

11-1-25 [OTZ] The following expressions for the equation of state and the specific heat (c_p) are obeyed by a certain gas: $v = RT/p\alpha T^2$ and $c_p = A + BT + Cp$ where a, A, B, C are constants. Obtain an expression for (a) the Joule-Thomson coefficient (μ_J) and (b) the specific heat (c_v).

SOLUTION:

Given $v = \frac{RT}{p} + \alpha T^2$ and $c_p = A + BT + Cp$

From Eq. (11.67), the Joule Thomson coefficient can be written as:

$$\mu_J = -\frac{1}{c_p} \left[v - T \left(\frac{\partial v}{\partial T} \right)_p \right]; \quad \left[\frac{\text{K}}{\text{kPa}} \right],$$

From the given relation: $v = \frac{RT}{p} + \alpha T^2$

$$\left(\frac{\partial v}{\partial T} \right)_p = \frac{R}{p} + 2\alpha T$$

$$\mu_J = -\frac{1}{c_p} \left[v - T \left(\frac{\partial v}{\partial T} \right)_p \right] = -\frac{1}{c_p} \left[\frac{RT}{p} + \alpha T^2 - T \left(\frac{R}{p} + 2\alpha T \right) \right]$$

$$\mu_J = \frac{\alpha T^2}{c_p} = \frac{\alpha T^2}{A + BT + Cp}$$

Using Eq. (11.58)

$$c_v = c_p - T \left(\frac{\partial v}{\partial T} \right)_p \left(\frac{\partial p}{\partial T} \right)_v$$

$$\text{Now, } \left(\frac{\partial p}{\partial T} \right)_v = \left(\frac{\partial}{\partial T} \left(\frac{RT}{v - \alpha T^2} \right) \right)_v = \left(\frac{Rv + \alpha RT^2}{(v - \alpha T^2)^2} \right)$$

$$\text{Therefore, } c_v = c_p - T \left(\frac{R}{p} + 2\alpha T \right) \left(\frac{Rv + \alpha RT^2}{(v - \alpha T^2)^2} \right)$$

$$\Rightarrow c_v = A + BT + Cp - T \left(\frac{R}{p} + 2\alpha T \right) \left(\frac{Rv + \alpha RT^2}{(v - \alpha T^2)^2} \right)$$

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11-1-26 [OTK] The differential of pressure obtained from a certain equation of state is given by the following expression: $dp = \{2(v - b) / RT\}dv - \{(v - b)^2 / RT^2\}dT$. Determine the equation of state.

SOLUTION:

Expressing the differential in the format of Eq. (11.9),

$$M(v, T) = \frac{2(v - b)}{RT} \text{ and } N(v, T) = \frac{(v - b)^2}{RT^2}$$

For a function $p(v, T)$ to exist, the test of exactness must be satisfied. That is,

$$\left(\frac{\partial M}{\partial T}\right)_v = \left(\frac{\partial N}{\partial v}\right)_T$$

$$\text{Now, } \left(\frac{\partial M}{\partial T}\right)_v = \frac{-2(v - b)}{RT^2} \text{ and } \left(\frac{\partial N}{\partial v}\right)_T = \frac{2(v - b)}{RT^2}$$

The two expressions are not equal. Therefore, the given expression is not the differential of an equation of state.

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11-1-27 [OTC] The differential of pressure obtained from a certain equation of state is given by the following expression: $dp = \{-RT/(v-b)^2\}dv + \{R/(v-b)\}dT$. Determine the equation of state.

SOLUTION:

Expressing the differential in the format of Eq. (11.9),

$$M(v, T) = \frac{-RT}{(v-b)^2} \text{ and } N(v, T) = \frac{R}{(v-b)}$$

For a function $p(v, T)$ to exist, the test of exactness must be satisfied. That is,

$$\left(\frac{\partial M}{\partial T}\right)_v = \left(\frac{\partial N}{\partial v}\right)_T$$

$$\text{Now, } \left(\frac{\partial M}{\partial T}\right)_v = \frac{-R}{(v-b)^2} \text{ and } \left(\frac{\partial N}{\partial v}\right)_T = \frac{-R}{(v-b)^2}$$

Therefore a function $p(v, T)$ exists. Expressing its differential using Taylor's theorem:

$$dp = \left(\frac{\partial p}{\partial v}\right)_T dv + \left(\frac{\partial p}{\partial T}\right)_v dT$$

Comparing with the given differential, we obtain:

$$\left(\frac{\partial p}{\partial v}\right)_T = \frac{-RT}{(v-b)^2} \text{ and } \left(\frac{\partial p}{\partial T}\right)_v = \frac{R}{(v-b)}$$

Integrating the second equation, we obtain

$$p = \frac{RT}{(v-b)} + f(v)$$

$$\Rightarrow \left(\frac{\partial p}{\partial v}\right)_T = \frac{-RT}{(v-b)^2} + \frac{df}{dv}$$

Comparing the two expressions for this partial,

$$\frac{df}{dv} = c$$

$$\text{Therefore, } p = \frac{RT}{(v-b)} + c$$

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xx11-1-28 [OTU] Derive the relation $c_p = -T(\partial^2 g / \partial T^2)_p$

SOLUTION:

Treating T and p as independent variable

$$\begin{aligned} \text{As } c_p &= \left(\frac{\partial h}{\partial T} \right)_p, \\ h &= h(T, p) \\ \Rightarrow dh &= \left(\frac{\partial h}{\partial T} \right)_p dT + \left(\frac{\partial h}{\partial p} \right)_T dp; \\ \Rightarrow dh &= c_p dT + \left(\frac{\partial h}{\partial p} \right)_T dp; \quad (1) \end{aligned}$$

Also,

$$\begin{aligned} s &= s(T, p) \\ \Rightarrow ds &= \left(\frac{\partial s}{\partial T} \right)_p dT + \left(\frac{\partial s}{\partial p} \right)_T dp; \quad (2) \end{aligned}$$

Substituting Equation (1) and (2) in $dh = Tds + vdp$

$$c_p dT + \left(\frac{\partial h}{\partial p} \right)_T dp = T \left[\left(\frac{\partial s}{\partial T} \right)_p dT + \left(\frac{\partial s}{\partial p} \right)_T dp \right] + vdp$$

Rearranging,

$$\left[c_p - T \left(\frac{\partial s}{\partial T} \right)_p \right] dT = \left[T \left(\frac{\partial s}{\partial p} \right)_T + v - \left(\frac{\partial h}{\partial p} \right)_T \right] dp$$

As T and p are independent, fix p while varying T : $dp = 0$ and $dT \neq 0$

Therefore,

$$\begin{aligned} \left[c_p - T \left(\frac{\partial s}{\partial T} \right)_p \right] dT &= 0; \\ \Rightarrow c_p - T \left(\frac{\partial s}{\partial T} \right)_p &= 0; \\ \Rightarrow c_p &= T \left(\frac{\partial s}{\partial T} \right)_p \quad (3) \end{aligned}$$

Now,

$$g = h - Ts$$

$$\Rightarrow \left(\frac{\partial g}{\partial T} \right)_p = \left(\frac{\partial h}{\partial T} \right)_p - T \left(\frac{\partial s}{\partial T} \right)_T - s;$$

$$\Rightarrow \left(\frac{\partial g}{\partial T} \right)_p = c_p - T \left(\frac{\partial s}{\partial T} \right)_T - s; \quad (4)$$

Substituting Eq. (3) into (4)

$$\Rightarrow \left(\frac{\partial g}{\partial T} \right)_p = -s; \quad (5)$$

Differentiate Eq. (5), we obtain

$$-\left(\frac{\partial s}{\partial T} \right)_p = \left(\frac{\partial^2 g}{\partial T^2} \right)_p \quad (5)$$

Combining Equation (3) and (5)

$$c_p = -T \left(\frac{\partial^2 g}{\partial T^2} \right)_p$$

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11-1-29 [OTX] Prove that $(\partial\beta/\partial p)_T = -(\partial\kappa_T/\partial T)_p$.

SOLUTION:

Volumetric expansivity:

$$\beta = \frac{1}{v} \left(\frac{\partial v}{\partial T} \right)_p \quad (1)$$

Differentiating equation (1) w.r.t. p at constant T ,

$$\begin{aligned} \left(\frac{\partial\beta}{\partial p} \right)_T &= \frac{1}{v} \left(\frac{\partial^2 v}{\partial p \partial T} \right) + \left(\frac{\partial v}{\partial T} \right)_p \left[\frac{\partial}{\partial p} \left(\frac{1}{v} \right) \right]_T \\ \Rightarrow \left(\frac{\partial\beta}{\partial p} \right)_T &= \frac{1}{v} \left(\frac{\partial^2 v}{\partial p \partial T} \right) + \left(\frac{\partial v}{\partial T} \right)_p \left(\frac{-1}{v^2} \right) \left(\frac{\partial v}{\partial p} \right)_T \\ \Rightarrow \left(\frac{\partial\beta}{\partial p} \right)_T &= \frac{1}{v} \left(\frac{\partial^2 v}{\partial p \partial T} \right) + \beta \kappa_T \quad (2) \end{aligned}$$

Isothermal compressibility:

$$\kappa_T = -\frac{1}{v} \left(\frac{\partial v}{\partial p} \right)_T \quad (3)$$

Differentiating equation (3) w.r.t. T at constant p ,

$$\begin{aligned} \left(\frac{\partial\kappa_T}{\partial T} \right)_p &= -\frac{1}{v} \left(\frac{\partial^2 v}{\partial T \partial p} \right) - \left(\frac{\partial v}{\partial p} \right)_T \left[\frac{\partial}{\partial T} \left(\frac{1}{v} \right) \right]_p \\ \Rightarrow \left(\frac{\partial\kappa_T}{\partial T} \right)_p &= -\frac{1}{v} \left(\frac{\partial^2 v}{\partial T \partial p} \right) - \left(\frac{\partial v}{\partial p} \right)_T \left(\frac{-1}{v^2} \right) \left(\frac{\partial v}{\partial T} \right)_p \\ \Rightarrow \left(\frac{\partial\kappa_T}{\partial T} \right)_p &= -\frac{1}{v} \left(\frac{\partial^2 v}{\partial T \partial p} \right) - \beta \kappa_T \quad (4) \end{aligned}$$

As $\frac{\partial^2 v}{\partial p \partial T} = \frac{\partial^2 v}{\partial T \partial p}$, comparing Eqs. (2) and (4) we get

$$\left(\frac{\partial\beta}{\partial p} \right)_T = - \left(\frac{\partial\kappa_T}{\partial T} \right)_p$$

11-1-30 [OTV] At certain states, the p - v - T data for a particular gas can be represented as $Z = 1 - Ap/T^4$, where Z is the compressibility factor and A is a constant. Obtain an expression for the difference in specific heats ($c_p - c_v$).

SOLUTION:

From Eq. (11.59)

$$c_p - c_v = -T \left(\frac{\partial v}{\partial T} \right)_p^2 \left(\frac{\partial p}{\partial v} \right)_T;$$

$$\Rightarrow c_p - c_v = -T \frac{\left(\frac{\partial v}{\partial T} \right)_p^2}{\left(\frac{\partial v}{\partial p} \right)_T}; \quad (1)$$

the equation of state is:

$$v = Z \left(\frac{RT}{p} \right) = \frac{RT}{p} - \frac{AR}{T^3};$$

Thus

$$\left(\frac{\partial v}{\partial T} \right)_p = \frac{R}{p} + \frac{3AR}{T^4}; \quad \left(\frac{\partial v}{\partial p} \right)_T = -\frac{RT}{p^2};$$

Substituting in Eq (1),

$$c_p - c_v = -\frac{T \left(\frac{R}{p} + \frac{3AR}{T^4} \right)^2}{\frac{RT}{p^2}}; \Rightarrow c_p - c_v = R \left[1 + \frac{3Ap}{T^4} \right]^2$$

With $\frac{Ap}{T^4} = 1 - Z$

$$c_p - c_v = R[1 + 3(1 - Z)]^2;$$

$$\Rightarrow c_p - c_v = R(4 - 3Z)^2$$