

Notes on derivatives of eigenvalue and eigenvector fields

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June 2021

1 Caustic conditions

Notation: We will here study the caustic conditions corresponding to the eigen equation

$$H\mathbf{v} = \lambda\mathbf{v}, \quad (1)$$

with the deformation tensor H , the eigenvalue field λ , and the eigenvector field \mathbf{v} . Let $\mathbf{v} = (v_1, v_2)$ in 2D and $\mathbf{v} = (v_1, v_2, v_3)$ in 3D. The eigenvector field is normalized by $\mathbf{v}^T \mathbf{v} = 1$. We define the partial derivative in the q_j direction as $\partial_i = \frac{\partial}{\partial q_i}$. We use Einstein summation convention for repeated indices.

1.1 Cusp A_3

The cusps are defined by the condition

$$0 = \mathbf{v} \cdot \nabla \lambda \quad (2)$$

$$= v_i \partial_i \lambda. \quad (3)$$

Using the identity for the first order derivative of the eigenvalue field, we can write this condition as

$$0 = \mathbf{v}^T (v_i \partial_i H) \mathbf{v} \quad (4)$$

where $v_i \partial_i H$ is the partial derivative of the deformation tensor in the direction of the eigenvector field \mathbf{v} . Note that this identity enables us to replace an interpolation over the eigenvalue field by an interpolation over the deformation field.

1.2 Swallowtail A_4

The swallowtail is defined by the condition

$$0 = \mathbf{v} \cdot \nabla (\mathbf{v} \cdot \nabla \lambda) \quad (5)$$

$$= v_j \partial_j (v_i \partial_i \lambda) \quad (6)$$

$$= v_j \partial_j v_i \partial_i \lambda + v_i v_j \partial_i \partial_j \lambda. \quad (7)$$

We can interpret the first term $v_j \partial_j v_i \partial_i \lambda$ as in two steps. Firstly, we evaluate the partial derivative of \mathbf{v} in the direction \mathbf{v} . Secondly, we evaluate the directional derivative of λ in this direction. Both the first order derivative of \mathbf{v} and λ can be written in terms of first order derivatives of the deformation field H . The second term is simply the second order derivative of λ in the direction \mathbf{v} . We can express this term in terms of the second order derivative of the deformation field H .

1.3 Butterfly A_5

The butterfly caustic is defined by the condition

$$0 = \mathbf{v} \cdot \nabla(\mathbf{v} \cdot \nabla(\mathbf{v} \cdot \nabla \lambda)) \quad (8)$$

$$= v_k \partial_k (v_j \partial_j (v_i \partial_i \lambda)) \quad (9)$$

$$= v_k \partial_k (v_j \partial_j v_i \partial_i \lambda + v_j v_i \partial_i \partial_j \lambda) \quad (10)$$

$$= v_k (\partial_k v_j \partial_j v_i + v_j \partial_j \partial_k v_i) \partial_i \lambda + 3 v_i v_k \partial_k v_j \partial_i \partial_j \lambda + v_i v_j v_k \partial_i \partial_j \partial_k \lambda. \quad (11)$$

It is again possible to write the first, second, and third order derivatives of the eigenvalue and eigenvector fields in terms of first, second and third order derivatives of the deformation field.

2 Derivatives of the eigenvalue and eigenvector fields

Consider the eigenvalue λ and eigenvector fields \mathbf{v} of the deformation field H satisfying the equation

$$H\mathbf{v} = \lambda\mathbf{v}, \quad (12)$$

with the normalization $\mathbf{v}^T \mathbf{v} = 1$. When the eigenvalue fields are simple, the eigenvalue and eigenvector fields are smooth. Interestingly, we can write the derivatives of the eigenvalue and eigenvector fields in terms of derivatives of the matrix field H .

2.1 First order derivative of the eigenvalue field

Upon differentiating the definition of the eigenvalue and eigenvector fields, we obtain the identity

$$\partial_i H\mathbf{v} + H\partial_i \mathbf{v} = \partial_i \lambda \mathbf{v} + \lambda \partial_i \mathbf{v}. \quad (13)$$

If we multiply this equation with \mathbf{v}^T from the left, we obtain the identity

$$\mathbf{v}^T \partial_i H\mathbf{v} + \mathbf{v}^T H\partial_i \mathbf{v} = \partial_i \lambda \mathbf{v}^T \mathbf{v} + \lambda \mathbf{v}^T \partial_i \mathbf{v} \quad (14)$$

which simplifies to

$$\partial_i \lambda = \mathbf{v}^T \partial_i H\mathbf{v} \quad (15)$$

using the condition $\mathbf{v}^T H = \lambda \mathbf{v}^T$ and $\mathbf{v}^T \mathbf{v} = 1$.

2.2 First order derivative of the eigenvector field

Upon differentiating the definition of the eigenvalue and eigenvector fields, we obtain the identity

$$\partial_i H\mathbf{v} + H\partial_i \mathbf{v} = \partial_i \lambda \mathbf{v} + \lambda \partial_i \mathbf{v}. \quad (16)$$

We can rearrange this identity to

$$(\lambda I - H)\partial_i \mathbf{v} = (\partial_i H - \partial_i \lambda)\mathbf{v}. \quad (17)$$

In order to obtain an expression for $\partial_i \mathbf{v}$, we would like to invert $\lambda I - H$. However, this matrix is by construction singular. We for this reason multiply from the left by the Moore-Penrose inverse defined by

$$A^+ = (A^T A)^{-1} A^T \quad (18)$$

assuming $A^T A$ is invertible. This gives

$$\partial_i \mathbf{v} = (\lambda I - H)^+ \partial_i H\mathbf{v} \quad (19)$$

since $(\lambda I - H)^+ \mathbf{v} = 0$.

2.3 Second order derivative of the eigenvalue field

Upon differentiating the definition of the eigenvalue and eigenvector fields twice, we obtain the identity

$$\partial_i \partial_j H \mathbf{v} + \partial_i H \partial_j \mathbf{v} + \partial_j H \partial_i \mathbf{v} + H \partial_i \partial_j \mathbf{v} = \partial_i \partial_j \lambda \mathbf{v} + \partial_i \lambda \partial_j \mathbf{v} + \partial_j \lambda \partial_i \mathbf{v} + \lambda \partial_i \partial_j \mathbf{v}. \quad (20)$$

Again, after multiplying this equation with \mathbf{v}^T from the left, we obtain

$$\partial_i \partial_j \lambda = \mathbf{v}^T \partial_i \partial_j H \mathbf{v} + \mathbf{v}^T (\partial_i H \partial_j \mathbf{v} + \partial_j H \partial_i \mathbf{v}), \quad (21)$$

since $\mathbf{v}^T \mathbf{v} = 1$, $\mathbf{v}^T \partial_i \mathbf{v} = 0$, and $\mathbf{v}^T H = \lambda \mathbf{v}^T$. Note that the second term is a simple symmetrization over the indices i and j . We can simplify this equation using the symmetrization operator over i and j , i.e., given

$$S[A_{ij}] = \frac{1}{2}(A_{ij} + A_{ji}), \quad (22)$$

we obtain

$$\partial_i \partial_j \lambda = S[\mathbf{v}^T (\partial_i \partial_j H \mathbf{v} + 2 \partial_i H \partial_j \mathbf{v})]. \quad (23)$$

2.4 Second order derivative of the eigenvector field

Upon differentiating the definition of the eigenvalue and eigenvector fields twice, we obtain the identity

$$\partial_i \partial_j H \mathbf{v} + \partial_i H \partial_j \mathbf{v} + \partial_j H \partial_i \mathbf{v} + H \partial_i \partial_j \mathbf{v} = \partial_i \partial_j \lambda \mathbf{v} + \partial_i \lambda \partial_j \mathbf{v} + \partial_j \lambda \partial_i \mathbf{v} + \lambda \partial_i \partial_j \mathbf{v}. \quad (24)$$

After rearranging this identity and multiplying by the inverse $(H - \lambda I)^+$ from the left, we obtain

$$\partial_i \partial_j \mathbf{v} = S[(\lambda I - H)^+ (\partial_i \partial_j H \mathbf{v} + 2 \partial_i H \partial_j \mathbf{v} - 2 \partial_i \lambda \partial_j \mathbf{v})]. \quad (25)$$

2.5 Third order derivative of the eigenvalue field

Upon differentiating the definition of the eigenvalue and eigenvector fields three times, and multiplying on the left with \mathbf{v}^T we obtain the identity

$$\partial_i \partial_j \partial_k \lambda = S[\mathbf{v}^T (\partial_i \partial_j \partial_k H \mathbf{v} + 3 \partial_i \partial_j H \partial_k \mathbf{v} + 3 \partial_i H \partial_j \partial_k \mathbf{v} - 3 \partial_i \partial_j \lambda \partial_k \mathbf{v} - 3 \partial_i \lambda \partial_j \partial_k \mathbf{v})], \quad (26)$$

where the symmetrization operation over i , j , and k is defined as

$$S[A_{ijk}] = \frac{1}{6}(A_{ijk} + A_{ikj} + A_{jik} + A_{jki} + A_{kji} + A_{kij}). \quad (27)$$

Note that due to the symmetry of the argument, we can use the simpler symmetrization over cyclic permutations

$$S[A_{ijk}] = \frac{1}{3}(A_{ijk} + A_{jki} + A_{kij}). \quad (28)$$