

PHY1013S ELECTRICITY GAUSS'S LAW

GAUSS'S LAW

Learning outcomes:
At the end of this chapter you should be able to...

- Calculate the electric flux through a surface.
- Use Gauss's Law to calculate the electric field due to symmetric charge distributions.
- Use Gauss's Law to determine the charge distribution on hollow conductors in electric fields.

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FLUX

\hat{n} is the unit vector normal to the surface.
 $\vec{A} = A\hat{n}$ is the surface's area vector.

The amount of flow, Φ , through the surface depends on:

- the area of the surface, A ;
- the angle between \vec{v} and \vec{A} .
 (Φ is a maximum when $\theta = 0^\circ$; a minimum for $\theta = 90^\circ$.)

Hence: $\Phi = vA\cos\theta$ or just: $\Phi = \vec{v} \cdot \vec{A}$

Replacing the velocity vector with the electric field vector \vec{E} , we get: $\Phi_e = \vec{E} \cdot \vec{A}$

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FLUX

For a surface made up of small elements, each ΔA , the total flux is $\Phi = \sum \vec{E} \cdot \Delta \vec{A}$

For a closed (Gaussian) surface, in the limit $\Delta A \rightarrow 0$,

$\Phi = \oint \vec{E} \cdot d\vec{A}$ [N m²/C]

The electric flux through a Gaussian surface is proportional to the *net* number of electric field lines passing through that surface.

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$\Phi = \vec{E} \cdot \vec{A}$

$\Delta\Phi < 0$

Gaussian surface

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Gaussian surface

$\Phi = \vec{E} \cdot \vec{A}$

$\Delta\Phi = 0$

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$\Phi = \vec{E} \cdot \vec{A}$

Gaussian surface

1 2 3

$\Delta\Phi > 0$

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HRW₆ 546 SP24-1

A Gaussian surface in the form of a cylinder of radius R is immersed in a uniform electric field \vec{E} , with the cylinder axis parallel to the field.

Determine the (net) flux of the field through this surface.

$$\Phi = \oint \vec{E} \cdot d\vec{A} = \int_a \vec{E} \cdot d\vec{A} + \int_b \vec{E} \cdot d\vec{A} + \int_c \vec{E} \cdot d\vec{A}$$

$$\therefore \Phi = \int_a E(\cos 180^\circ) dA + \int_b E(\cos 90^\circ) dA + \int_c E(\cos 0^\circ) dA$$

$$\therefore \Phi = -E \int_a dA + E \int_b (0) dA + E \int_c dA$$

$$\therefore \Phi = -EA + 0 + EA \quad \therefore \Phi = 0$$

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HRW₆ 547 SP24-2

A non-uniform electric field given by $\vec{E} = 3.0x\hat{i} + 4.0\hat{j}$ pierces a Gaussian cube situated on the x -axis between $x = 1.0$ m and $x = 3.0$ m.

Determine the electric flux through the right face.

$$d\vec{A} = dA\hat{i} \quad \therefore \Phi = \int \vec{E} \cdot d\vec{A} = \int (3.0x\hat{i} + 4.0\hat{j}) \cdot (dA\hat{i})$$

$$\therefore \Phi = \int [(3.0x)(dA)\hat{i} \cdot \hat{i} + (4.0)(dA)\hat{j} \cdot \hat{i}]$$

$$\therefore \Phi = \int (3.0x dA + 0) = 3.0 \int x dA$$

$$\therefore \Phi = 3.0 \int (3.0) dA = 9.0 \int dA = (9.0 \text{ N/C})(4.0 \text{ m}^2)$$

$$\therefore \Phi = 36 \text{ Nm}^2/\text{C}$$

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GAUSS' LAW

The net electric flux through a Gaussian surface is proportional to the net charge enclosed: $\epsilon_0 \Phi = Q_{\text{in}}$

Substituting for Φ , $\epsilon_0 \oint \vec{E} \cdot d\vec{A} = Q_{\text{in}}$

Notes:

- Q_{in} is the *net, enclosed* charge.
- \vec{E} is the *total* field through the surface.
- These equations are valid only in vacuum (or air).
- Gauss's Law is both easier to use and more universal/fundamental than Coulomb's Law.

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GAUSS'S LAW: MULTIPLE CHARGES

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USING GAUSS'S LAW TO DETERMINE ELECTRIC FIELDS

- Draw the situation.
- Choose a Gaussian surface appropriate to the symmetry.
- Apply Gauss's law: $\epsilon_0 \oint \vec{E} \cdot d\vec{A} = Q_{\text{in}}$

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GAUSS \Leftrightarrow COULOMB

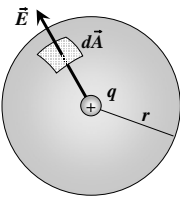
- Draw the situation...
- Choose a Gaussian surface appropriate to the symmetry...
(For a point charge, choose a concentric sphere as a Gaussian surface.)
The electric field is constant over the surface and directed radially outwards, so $\vec{E} \cdot d\vec{A} = E dA$
- Apply Gauss's law...

$$Q_{in} = \epsilon_0 \oint \vec{E} \cdot d\vec{A} = \epsilon_0 \oint E dA$$

$$\therefore Q_{in} = \epsilon_0 E \oint dA$$

$$\therefore Q_{in} = \epsilon_0 E 4\pi r^2$$

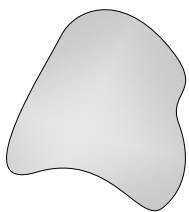
$$E = \frac{q}{4\pi\epsilon_0 r^2}$$



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ZERO FIELD



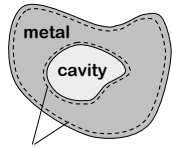
If there is no electric field *at all*, quite obviously the net flux through any Gaussian surface is also zero.

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CHARGE ON AN ISOLATED CONDUCTOR

Any excess charge added to an isolated conductor moves entirely to the external surface of that conductor.



Gaussian surfaces

- The field inside the conductor must be zero (otherwise there would be currents inside the conductor)...
- Therefore there is no flux through the Gaussian surfaces...
- \therefore there can be no charge *within* the Gaussian surfaces...
- \therefore all the charge must lie *outside* the Gaussian surfaces...
i.e. ...

All the charge lies on the external surface of the conductor.

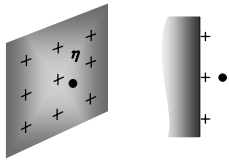
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EXTERNAL FIELD DUE TO A CHARGED CONDUCTOR

For a non-spherical conductor, the surface charge density, η , varies over the surface, and the field established around the conductor is very complex.

However, for a point just outside the surface, the adjacent section of surface is small enough to be considered as flat, and the charge density as uniform...

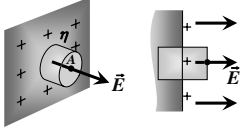


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EXTERNAL FIELD DUE TO A CHARGED CONDUCTOR

A cylindrical Gaussian surface with an end cap area of A is embedded in the surface of the conductor as shown.



The flux through the outer cap is EA .

The total charge enclosed by the cylinder is given by ηA .

Therefore, according to Gauss's law: $\epsilon_0 EA = \eta A$

and hence: $E = \frac{\eta}{\epsilon_0}$

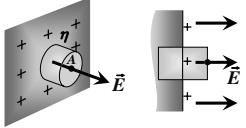
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EXTERNAL FIELD DUE TO A CHARGED CONDUCTOR

Summary:

- The electric field is zero everywhere *inside* the conductor.
- Any excess charge is all on the *surface*.
- The external field lies perpendicular to the surface and is given by $E = \frac{\eta}{\epsilon_0}$.
- On an irregularly shaped conductor the charge collects around sharp points, but $E = \frac{\eta}{\epsilon_0}$ still holds true.

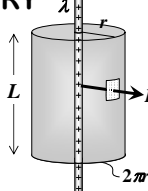


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CYLINDRICAL SYMMETRY

For a long, thin, cylindrical *insulator* with a uniform linear charge density of λ ... we choose a cylindrical Gaussian surface with radius r and height L :



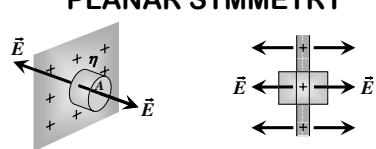
The area of the curved surface is $2\pi r L$.
By symmetry, the total flux through the surface is $E 2\pi r L$.
The total charge enclosed by the cylinder is λL .
Therefore, according to Gauss's law: $\epsilon_0 E 2\pi r L = \lambda L$

and hence: $E = \frac{\lambda}{2\pi\epsilon_0 r}$

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PLANAR SYMMETRY



For a large, flat, thin *insulating* sheet with a uniform surface charge density of η ... we choose a cylindrical Gaussian surface with end cap area A , which pierces the sheet perpendicularly.

According to Gauss's law: $\epsilon_0(E A + E A) = \eta A$

and hence: $E = \frac{\eta}{2\epsilon_0}$

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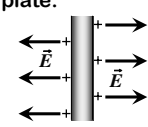
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PLANAR SYMMETRY

For a charged large, flat, thin *conducting* plate:

surface charge density = η_1

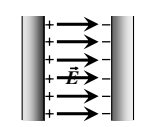
$E = \frac{\eta_1}{\epsilon_0}$



When two oppositely charged plates are brought close together, all the charge moves to the inner faces:

surface charge density $\eta = 2\eta_1$

$E = \frac{2\eta_1}{\epsilon_0} = \frac{\eta}{\epsilon_0}$



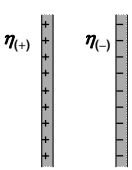
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HRW, 547 SP23-6

Two large, parallel, non-conducting sheets each carry a fixed uniform charge on one side: $\eta_{(+)} = 6.8 \mu\text{C}/\text{m}^2$ and $\eta_{(-)} = 4.3 \mu\text{C}/\text{m}^2$.

Determine the electric field (a) to the left, (b) between, and (c) to the right of the sheets.



$E_{(+)} = \frac{\eta_{(+)}}{2\epsilon_0} = 3.84 \times 10^5 \text{ N/C}$ $E_{(-)} = \frac{\eta_{(-)}}{2\epsilon_0} = 2.43 \times 10^5 \text{ N/C}$

Left: $\vec{E}_{(+)} \leftarrow$, $\vec{E}_{(-)} \leftarrow$ Between: $\vec{E}_{(+)} \rightarrow$, $\vec{E}_{(-)} \rightarrow$ Right: $\vec{E}_{(+)} \rightarrow$, $\vec{E}_{(-)} \leftarrow$

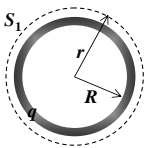
$\vec{E}_L = 1.4 \times 10^5 \text{ N/C, left}$ $\vec{E}_B = 6.3 \times 10^5 \text{ N/C, right}$ $\vec{E}_R = 1.4 \times 10^5 \text{ N/C, right}$

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SPHERICAL SYMMETRY

A spherical shell of radius R and charge q is surrounded by a concentric spherical Gaussian surface (S_1) with radius r (where $r \geq R$).



According to Gauss's law: $\epsilon_0 E 4\pi r^2 = q$

and hence, for S_1 : $E = \frac{q}{4\pi\epsilon_0 r^2}$

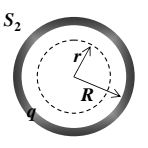
i.e. A uniform spherical shell of charge acts, on all charges outside it, as if all its charge were concentrated at its centre. [Shell theorem 1]

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SPHERICAL SYMMETRY

S_2 is a concentric spherical Gaussian surface with radius r lying *within* the spherical shell of charge q (i.e. $r \leq R$).



According to Gauss, for S_2 : $E = 0$

i.e. A uniform spherical shell of charge exerts no electrostatic force on a charged particle located inside it. [Shell theorem 2]

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