

CHAPTER 26

TENSORS

Tensors provide a mathematical framework for describing physical quantities that are independent of the choice of coordinate system. They generalize scalars and vectors and play a central role in continuum mechanics, electromagnetism, relativity, and modern physics.

26.1 SCALARS AND VECTORS

26.1.1 SCALARS

A scalar is a quantity that is fully described by a single number and remains unchanged under coordinate transformations.

Examples include mass, temperature, and energy.

26.1.2 VECTORS

A vector is a quantity with magnitude and direction. In a Cartesian coordinate system, a vector \mathbf{A} is represented by components

$$\mathbf{A} = (A^1, A^2, A^3)$$

Under a change of coordinates, the components transform linearly.

26.2 COORDINATE TRANSFORMATIONS

Consider two coordinate systems related by a linear transformation

$$x^{i'} = \frac{\partial x^{i'}}{\partial x^j} x^j$$

The transformation matrix is defined as

$$\Lambda^{i'}_j = \frac{\partial x^{i'}}{\partial x^j}$$

26.3 DEFINITION OF A TENSOR

A tensor is defined by how its components transform under a change of coordinates.

26.3.1 CONTRAVARIANT VECTORS

A contravariant vector V^i transforms as

$$V^{i'} = \frac{\partial x^{i'}}{\partial x^j} V^j$$

26.3.2 COVARIANT VECTORS

A covariant vector W_i transforms as

$$W_{i'} = \frac{\partial x^j}{\partial x^{i'}} W_j$$

26.4 METRIC TENSOR

The metric tensor defines distances and angles in a space.

In a general coordinate system, the squared line element is

$$ds^2 = g_{ij} dx^i dx^j$$

The metric tensor g_{ij} is symmetric

$$g_{ij} = g_{ji}$$

26.5 RAISING AND LOWERING INDICES

The metric tensor allows conversion between covariant and contravariant components.

Lowering an index is performed by

$$V_i = g_{ij} V^j$$

Raising an index is performed using the inverse metric g^{ij}

$$V^i = g^{ij} V_j$$

26.6 GENERAL TENSORS

A tensor of type (p, q) has p contravariant indices and q covariant indices.

Its components transform as

$$T^{i'_1 \dots i'_p}_{j'_1 \dots j'_q} = \frac{\partial x^{i'_1}}{\partial x^{i_1}} \dots \frac{\partial x^{i'_p}}{\partial x^{i_p}} \frac{\partial x^{j_1}}{\partial x^{j'_1}} \dots \frac{\partial x^{j_q}}{\partial x^{j'_q}} T^{i_1 \dots i_p}_{j_1 \dots j_q}$$

26.7 TENSOR OPERATIONS

26.7.1 ADDITION AND SCALAR MULTIPLICATION

Tensors of the same type may be added componentwise.

26.7.2 TENSOR PRODUCT

The tensor product of two tensors A and B produces a tensor of higher rank

$$(A \otimes B)^{ij} = A^i B^j$$

26.7.3 CONTRACTION

Contraction reduces the rank of a tensor by summing over one upper and one lower index

$$T^i_i$$

26.8 IMPORTANT PHYSICAL TENSORS

26.8.1 STRESS TENSOR

In continuum mechanics, the stress tensor σ_{ij} relates force to area

$$F_i = \sigma_{ij} n^j$$

where n^j is the normal vector.

26.8.2 MOMENT OF INERTIA TENSOR

The moment of inertia tensor is defined as

$$I_{ij} = \sum_k m_k (\delta_{ij} r_k^2 - x_{k,i} x_{k,j})$$

26.8.3 ELECTROMAGNETIC FIELD TENSOR

In relativistic electrodynamics, the electromagnetic field is represented by a rank-2 tensor.

26.9 TENSOR CALCULUS

26.9.1 PARTIAL DERIVATIVES

The partial derivative of a tensor is generally not a tensor.

26.9.2 COVARIANT DERIVATIVE

To preserve tensorial character, the covariant derivative is introduced

$$\nabla_k V^i = \partial_k V^i + \Gamma_{kj}^i V^j$$

where Γ_{kj}^i are the Christoffel symbols.

26.10 CLOSING REMARKS

Tensors provide a coordinate-independent language for physical laws. Their systematic use unifies diverse areas of physics and enables the formulation of laws valid in arbitrary coordinate systems.