#### Spring 2019



EECE 588 Lecture 12

**Prof. Wonbin Hong** 

- Now, we want to consider the radiation from a loop antenna with  $C \cong \lambda$  but we still assume that the current is a uniform one.
- Note that this is inherently a false assumption.
- Even though this is a false assumption, the results will be helpful in our real calculations.

$$R = \sqrt{r^2 + a^2 - 2ar\sin\theta\cos\phi'} \approx \sqrt{r^2 - 2ar\sin\theta\cos\phi'} \quad \text{for } r >> a$$

Using binomial expansion

Phase 
$$R \approx r \sqrt{1 - 2\frac{a}{r}\sin\theta\cos\varphi'} \approx r - a\sin\theta\cos\varphi' = r - a\cos\psi_o$$

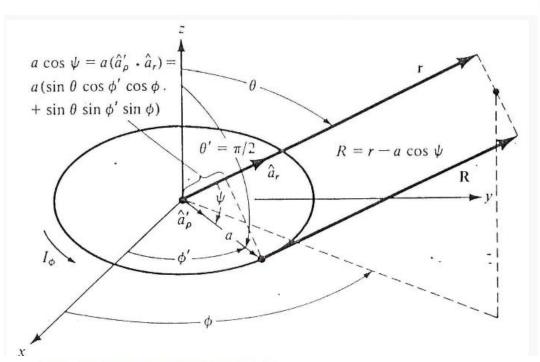
**Amplitude** 

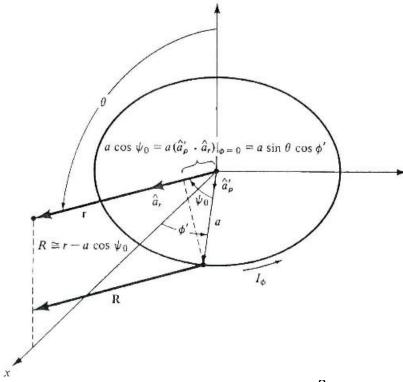




Note that angle  $\psi$  is the angle between r and r'.

•  $\psi_0$  is the angle between r and r' when  $\varphi = 0$ .







$$A_{\varphi} = \frac{a\mu I_0}{4\pi} \int_{0}^{2\pi} \cos\varphi' \frac{e^{-jk\sqrt{r^2 + a^2 - 2ar\sin\theta\cos\varphi'}}}{\sqrt{r^2 + a^2 - 2ar\sin\theta\cos\varphi'}} d\varphi'$$

$$A_{\varphi} \approx \frac{a\mu I_0 e^{-jkr}}{4\pi r} \int_0^{2\pi} \cos\varphi' e^{jka\sin\theta\cos\varphi'} d\varphi'$$

$$A_{\varphi} \approx \frac{a\mu I_{0}e^{-jkr}}{4\pi r} \int_{0}^{2\pi} \cos\varphi' e^{jka\sin\theta\cos\varphi'} d\varphi'$$

$$A_{\varphi} \approx \frac{a\mu I_{0}e^{-jkr}}{4\pi r} \left[ \int_{0}^{\pi} \cos\varphi' e^{jka\sin\theta\cos\varphi'} d\varphi' + \int_{\pi}^{2\pi} \cos\varphi' e^{+jka\sin\theta\cos\varphi'} d\varphi' \right]$$

$$\varphi' = \varphi'' + \pi$$

$$A_{\varphi} \approx \frac{a\mu I_{0}e^{-jkr}}{4\pi r} \left[ \int_{0}^{\pi} \cos\varphi' e^{jka\sin\theta\cos\varphi'} d\varphi' - \int_{0}^{\pi} \cos\varphi'' e^{-jka\sin\theta\cos\varphi''} d\varphi'' \right]$$



We know 
$$\pi j^n J_n(z) = \int_0^\pi \cos(n\varphi) e^{jz\cos\varphi} d\varphi$$

 $J_n(z)$  is the Bessel function of the first kind of order n

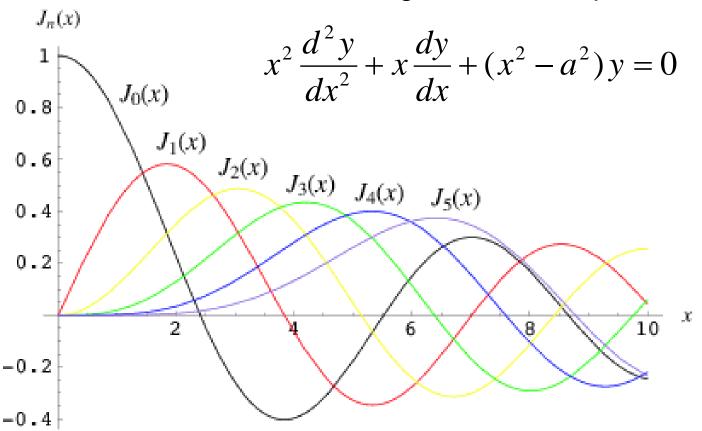
$$A_{\varphi} \approx \frac{a\mu I_0 e^{-jkr}}{4\pi r} \left\{ \pi j J_1(ka\sin\theta) - \pi j J_1(-ka\sin\theta) \right\}$$

$$J_n(z) = \sum_{m=0}^{\infty} \frac{(-1)^m (z/2)^{n+2m}}{m!(n+m)!}$$

$$J_n(-z) = (-1)^n J_n(z) \qquad A_{\varphi} \approx j \frac{a\mu I_0 e^{-jkr}}{2r} J_1(ka\sin\theta)$$

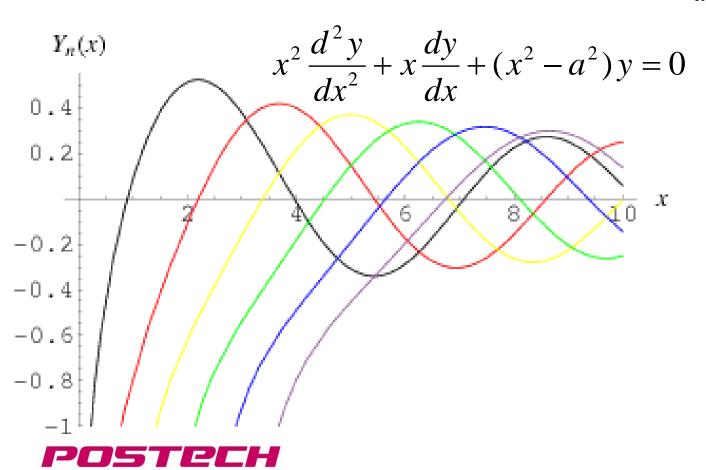


Note that the Bessel functions of the first kind  $J_a(x)$  are solutions to the following differential equation:





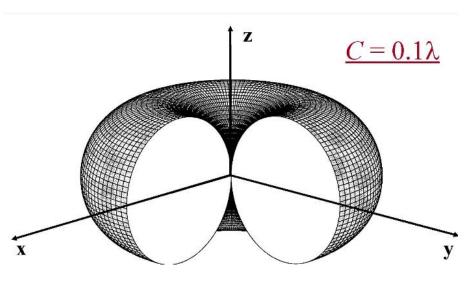
■ This equation has another solution called the Bessel function of the second kind or the Neumann function  $Y_a(x)$ :

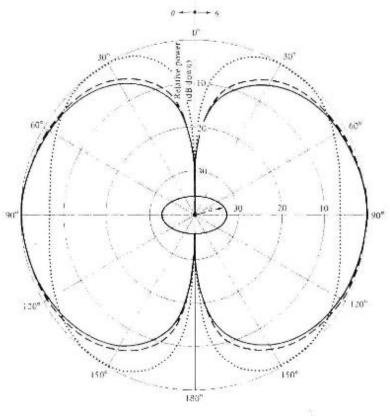


$$\begin{split} E_r &\approx E_\theta = 0 \\ E_\varphi &\approx \frac{ak\eta I_0 e^{-jkr}}{2r} J_1(ka\sin\theta) \\ H_r &\approx H_\varphi = 0 \end{split}$$

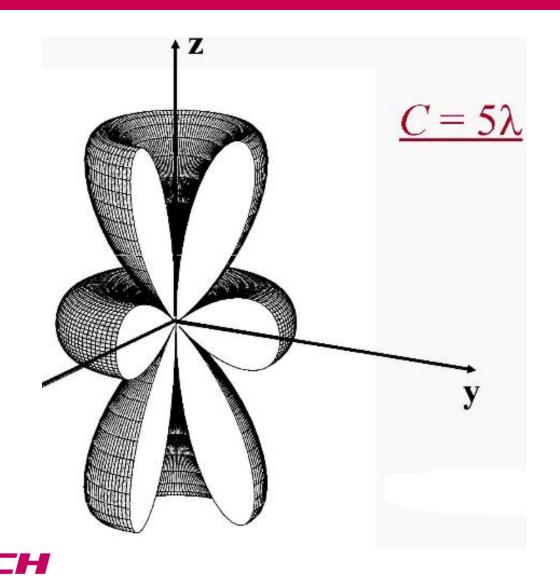
$$H_{\theta} = -E_{\varphi}/\eta = -\frac{akI_0e^{-jkr}}{2r}J_1(ka\sin\theta)$$













$$P_{rad} = \iint_{S} \vec{W}_{av} \cdot d\vec{s} = \pi (a\omega\mu)^{2} |I_{0}|^{2} / 4\eta \int_{0}^{\pi} J_{1}^{2} (ka\sin\theta) \sin\theta d\theta$$

$$1/2\int_{0}^{\pi} J_{1}^{2}(ka\sin\theta)\sin\theta d\theta = Q_{11}^{(1)}(ka) = \frac{1}{ka}\sum_{m=0}^{\infty} J_{2m+3}(2ka)$$



# Large Loop Approximation (a≥λ/2)

$$\int_{0}^{\pi} J_{1}^{2}(ka\sin\theta) \sin\theta d\theta = \frac{1}{ka} \int_{0}^{2ka} J_{2}(x) dx \approx \frac{1}{ka}$$
$$P_{rad} \approx \pi (a\omega\mu)^{2} |I_{0}|^{2} / (4\eta ka)$$

$$U|_{\max} = \frac{(a\omega\mu)^2 |I_0|^2}{8\eta} J_1(ka\sin\theta) \Big|_{ka\sin\theta=1.84} = \frac{(a\omega\mu)^2 |I_0|^2}{8\eta} (0.582)^2$$

$$R_{rad} = \frac{2P_{rad}}{\left|I_0\right|^2} = 60\pi^2 \left(\frac{C}{\lambda}\right)^2$$

$$D_0 = 0.677 \left(\frac{C}{\lambda}\right)$$



# Intermediate Loop Approximation $(\lambda/6\pi \le a < \lambda/2)$

$$R_{rad} = \frac{2P_{rad}}{|I_0|^2} = \eta \pi (ka)^2 Q_{11}^{(1)}(ka)$$

$$D_0 = \frac{4\pi U_{\text{max}}}{P_{\text{rad}}} = \frac{F_m(ka)}{Q_{11}^{(1)}(ka)}$$

$$|F_m(ka) = J_1^2(ka\sin\theta)|_{\max} = \begin{bmatrix} J_1^2(1.84) = 0.339 & ka > 1.84 \\ J_1^2(ka) & ka < 1.84 \end{bmatrix}$$



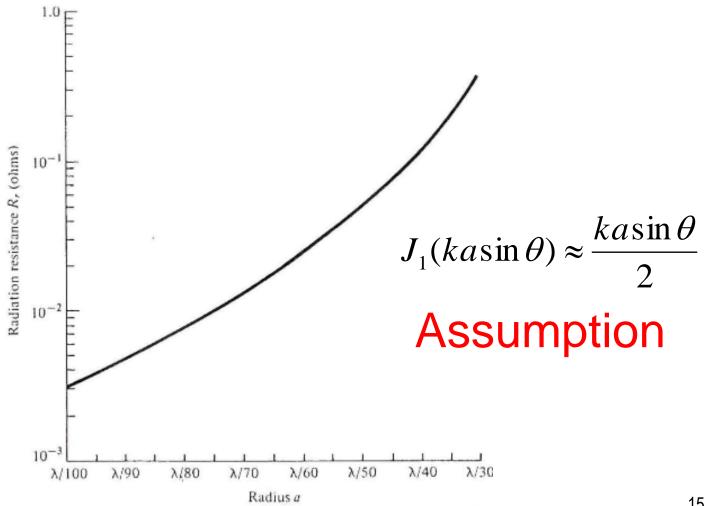
#### Small Loop Approximation (a < $\lambda$ /6 $\pi$ )

$$J_1(ka\sin\theta) \approx \frac{ka\sin\theta}{2}$$

With this approximation, the experssions will be those we obtained for the electrically small loop.

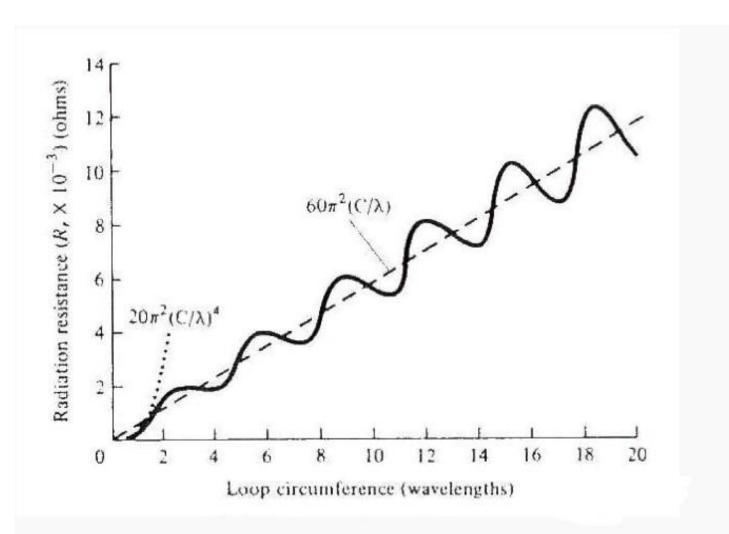


### Radiation Resistance for a Constant Current Loop Antenna



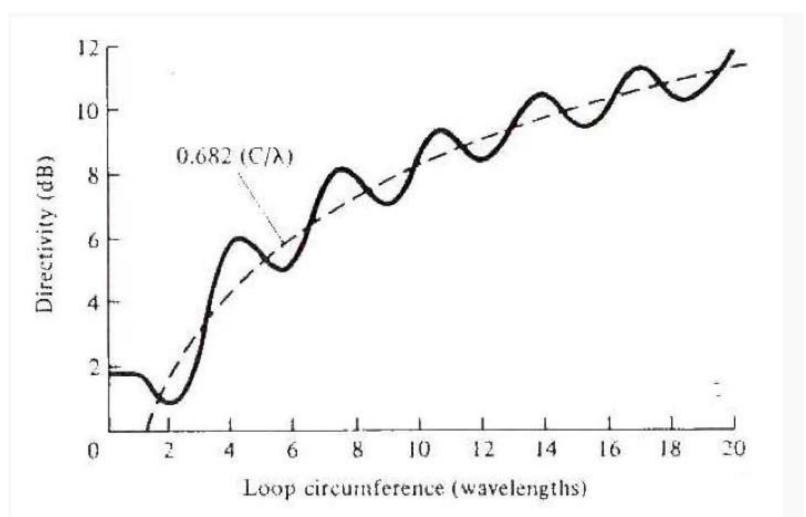


# Radiation Resistance of a Circular Loop with Constant Current





# Directivity of a Circular Loop with Constant Current





#### Circular Loop with Non Uniform Current Distribution

- As the dimensions of the loop increase the current variations along the circumference of the loop must be taken into account.
- Generally, a common assumption for the current distribution is a sinusoidal variation.
- This assumption is good. However, it is not valid close to the feed point.
- A better approximation is representing the current distribution by a Fourier Series.



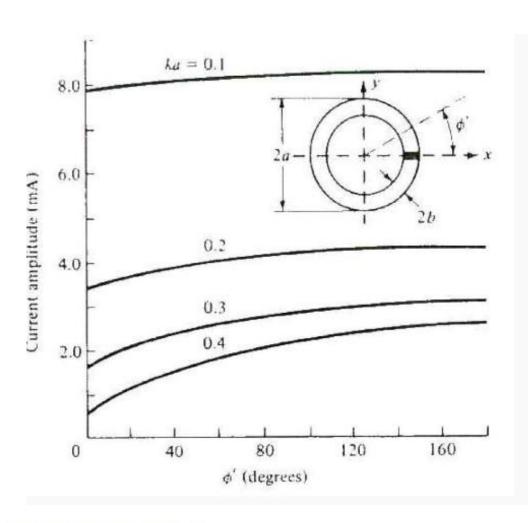
# Circular Loop with Non Uniform Current Distribution

$$I(\varphi') = I_0 + 2\sum_{n=1}^{M} I_n \cos(n\varphi')$$

- $\phi'$  is measured from the feed point of the loop along the circumference.
- Using this approximation, the problem can be solved analytically.
- This, however, is cumbersome and we do not attempt it here.



# Magnitude of the Current Distribution as a Function of Location along the Loop

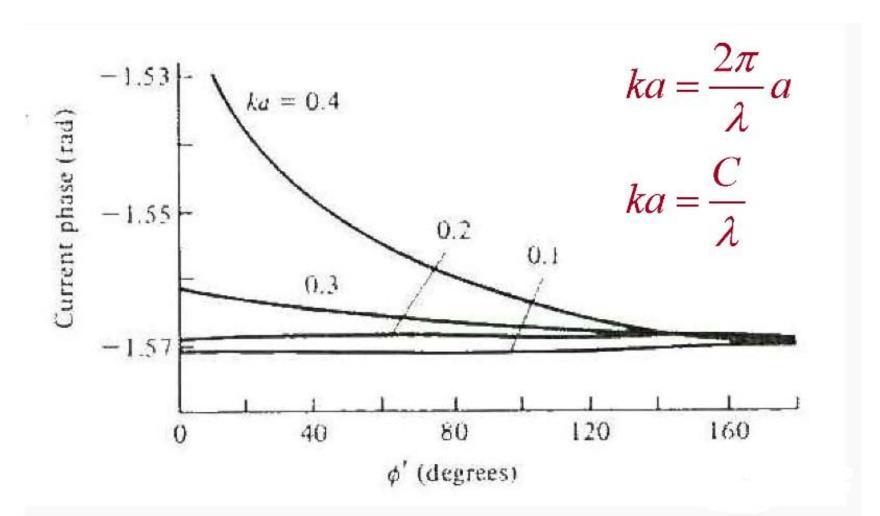


$$ka = \frac{2\pi}{\lambda}a$$

$$ka = \frac{C}{\lambda}$$



# Phase of the Current Distribution as a Function of Location along the Loop

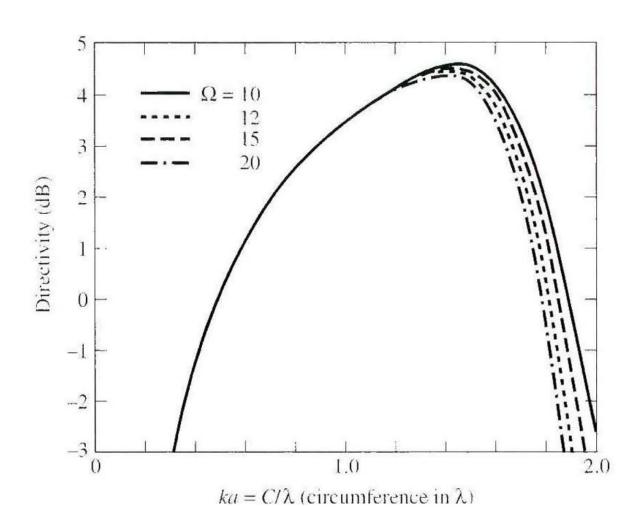




# Directivity of the Loop for $\theta=0^{\circ}$

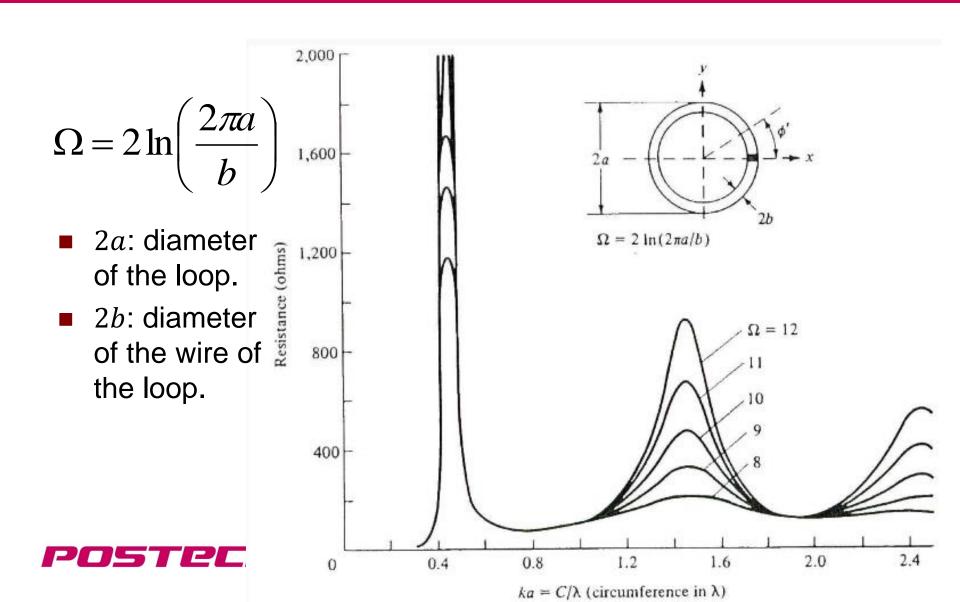
$$\Omega = 2 \ln \left( \frac{2\pi a}{b} \right)$$

- 2a: diameter of the loop.
- 2b: diameter of the wire of the loop.

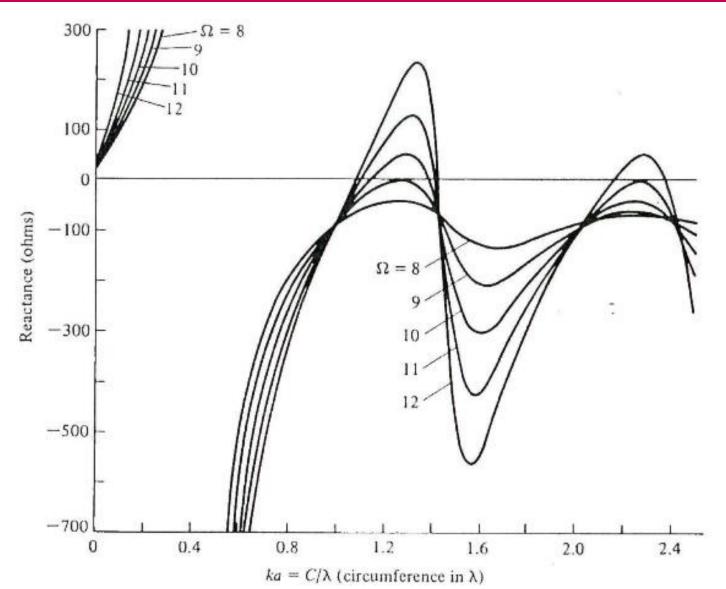




# Input Impedance of a Circular Loop



# Input Impedance of a Circular Loop





### Ferrite Loops

- Ferrite materials can be placed in the loop to increase the magnetic flux density of the loop.
- This will increase radiation resistance and hence the efficiency of the antenna.

$$\frac{R_f}{R_r} = \left(\frac{\mu_{ce}}{\mu_0}\right)^2$$

- $\blacksquare$   $R_f$ =radiation resistance of ferrite loop.
- $\blacksquare$   $R_r$ =radiation resistance of air core loop.
- $\mu_{ce}$ =effective permeability of the ferrite core.
- $\mu_0$ = free space permeability.



# Ferrite Loops

For a single turn small ferrite loop:

$$R_f = 20\pi^2 \left(\frac{C}{\lambda}\right)^4 \left(\frac{\mu_{ce}}{\mu_0}\right)^2 = 20\pi^2 \left(\frac{C}{\lambda}\right)^4 \mu_{cer}^2$$

$$R_f = 20\pi^2 \left(\frac{C}{\lambda}\right)^4 \left(\frac{\mu_{ce}}{\mu_0}\right)^2 N^2 = 20\pi^2 \left(\frac{C}{\lambda}\right)^4 \mu_{cer}^2 N^2$$

Relative effective permeability of the ferrite core is related to the relative intrinsic permeability of the unbounded ferrite material:



# Ferrite Loops

D is the demagnetization factor which is a function of the

shape of the core

