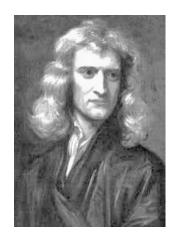
## Chapter 5 DIFFERENTIAL EQUATIONS



Isaac Newton
(1642-1726)
Math/Physics
Universal Gravity
Newtonian Mechanics
Differential Calculus

**Lecture 19** 

5.3 Series Solutions: Frobenius Method



Gottfried Wilhelm Leibniz (1646-1716) Math/Physic Integral Calculus Leibnitz Notation

## 5.2 Series Solutions: Frobenius Method

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For many problems in physics and engineering, we need to solve 10 second-order inhomogeneous ODEs:

$$\frac{d^2y}{dx^2} + P(x)\frac{dy}{dx} + Q(x) = F(x)$$
(5.12)

for which we usually use the method of separation of variables, for example, in Cartesian, cylindrical, and spherical coordinates.

$$y(x) = y_h(x) + y_p(x)$$

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$$y(x)$$

The most general solution is given by 
$$y(x) = y_h(x) + y_p(x)$$

$$y_h(x) = A_1 y_1(x) + A_2 y_2(x)$$
: Homogeneous Solutions for  $F(x) = 0$ 

$$y_p(x)$$
: Particular Solution for  $F(x) \neq 0$  (5.14)

and the homogeneous solutions can be obtained in a form of power series:

$$y_h(x) = x^k (a_0 + a_1 x + a_2 x^2 \cdots) = \sum_{n=0}^{\infty} a_n x^{k+n}, \quad a_0 \neq 0$$

[Q] What determines  $\overline{A_1}$  and  $\overline{A_2}$ 

Ex) Find the homogeneous solutions of a classical linear oscillator;

$$\frac{d^2y}{dx^2} + \omega^2 y = 0 ag{5.15}$$

Let's try a series solution:

$$y(x) = x^{k} \sum_{n=0}^{\infty} a_{n} x^{k+n} = \sum_{n=0}^{\infty} a_{n} x^{k+n} \ (a_{0} \neq 0)$$
 (5.16)

from which we have

$$y(x) = x^{k} \sum_{n=0}^{\infty} a_{n} x^{k+n} = \sum_{n=0}^{\infty} a_{n} x^{k+n} \ (a_{0} \neq 0)$$

$$\frac{d^{2}y}{dx^{2}} = \sum_{n=0}^{\infty} a_{n} (k+n)(k+n-1) x^{k+n-2}$$

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(5.17) into (5.15).

Substituting (5.16) and (5.17) into (5.15),

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$$\sum_{n=0}^{\infty} a_n (k+n)(k+n-1) x^{k+n-2} = -\omega^2 \sum_{n=0}^{\infty} a_n x^{k+n}$$

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$$= -\omega^2 \left[ a_0 x^k + a_1 x^{k+1} + a_2 x^{k+2} + \cdots \right]$$

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So we have

$$a_{n+2} = -\frac{\omega^2}{(k+n+2)(k+n+1)} a_n$$
 Recurrence Relation (5.19)

Now we note that at least the first term of LHS in (5.18) should vanish

$$\underbrace{a_0 k(k-1) = 0} \qquad \qquad 0.5.20$$

Since we initially assumed that  $a_0 \neq 0$ , we should have two cases k = 0 or k = 1

$$k = 0$$
 or  $k = 1$  (5.21)

1) 
$$k = 0$$
:  $a_{n+2} = -\frac{\omega^2}{(n+2)(n+1)} a_n \Rightarrow (-1)^n \frac{\omega^{2n}}{(2n)!} a_0$ 

$$\Rightarrow a_2 = -\frac{\omega^2}{1 \cdot 2} a_0 = -\frac{\omega^2}{2!} a_0, \quad a_4 = -\frac{\omega^2}{3 \cdot 4} a_2 = \frac{\omega^2}{4!} a_0, \quad a_6 = \frac{\omega^2}{5 \cdot 6} a_4 = -\frac{\omega^2}{4!} a_0, \quad \cdots$$

2) 
$$k = 1$$
:  $a_{n+2} = -\frac{\omega^2}{(n+3)(n+2)} a_n = (-1)^n \frac{\omega^{2n}}{(2n+1)!} a_0$ 

$$\Rightarrow a_2 = -\frac{\omega^2}{2 \cdot 3} a_0 = -\frac{\omega^2}{3!} a_0, \quad a_4 = -\frac{\omega^2}{4 \cdot 5} a_2 = \frac{\omega^2}{5!} a_0, \quad a_6 = \frac{\omega^2}{6 \cdot 7} a_4 = -\frac{\omega^2}{7!} a_0, \quad \cdots$$

## Therefore we finally have

$$y_{h}(x) = \begin{bmatrix} y_{1}(x) = y(x) \big|_{k=0} = a_{0} \left[ 1 - \frac{(\omega x)^{2}}{2!} + \frac{(\omega x)^{4}}{2!} + \cdots \right] = a_{0} \cos \omega x \\ y_{2}(x) = y(x) \big|_{k=1} = \frac{a_{0}}{\omega} \left[ \omega x - \frac{(\omega x)^{3}}{3!} + \frac{(\omega x)^{5}}{5!} + \cdots \right] = \frac{a_{0} \sin \omega x}{\omega}$$
(5.22)

## **5.3 Some Special Functions**

Screening Effect.

**Bessel** Function

$$x^{2} \frac{d^{2}y}{dx^{2}} + x \frac{dy}{dx} + (x^{2} - n^{2})y = 0$$
 (5.23)

**Legendre Function** 

$$(1-x^2)\frac{d^2y}{dx^2} + 2x\frac{dy}{dx} + n(n+1)y = 0$$
(5.24)

**Associate Legendre Function** 

$$(1-x^2)\frac{d^2y}{dx^2} + 2x\frac{dy}{dx} + \left[n(n+1) - \frac{m^2}{1-x^2}\right]y = 0$$
 (5.25)

Hermite Polynomial

$$\frac{d^2y}{dx^2} + 2x\frac{dy}{dx} + 2ny = 0$$
 (5.26)

Laguerre Polynomial

$$x\frac{d^2y}{dx^2} + (1-x)\frac{dy}{dx} + ny = 0$$
(5.27)

**Associate Laguerre Polynomial** 

$$x\frac{d^2y}{dx^2} + (m+1-x)\frac{dy}{dx} + (n-m)y = 0$$
(5.28)