

Spring 2019



EECE 588
Lecture 9

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Power Density, Radiation Intensity, and Radiation Resistance

- The average radiation intensity can be calculated using:

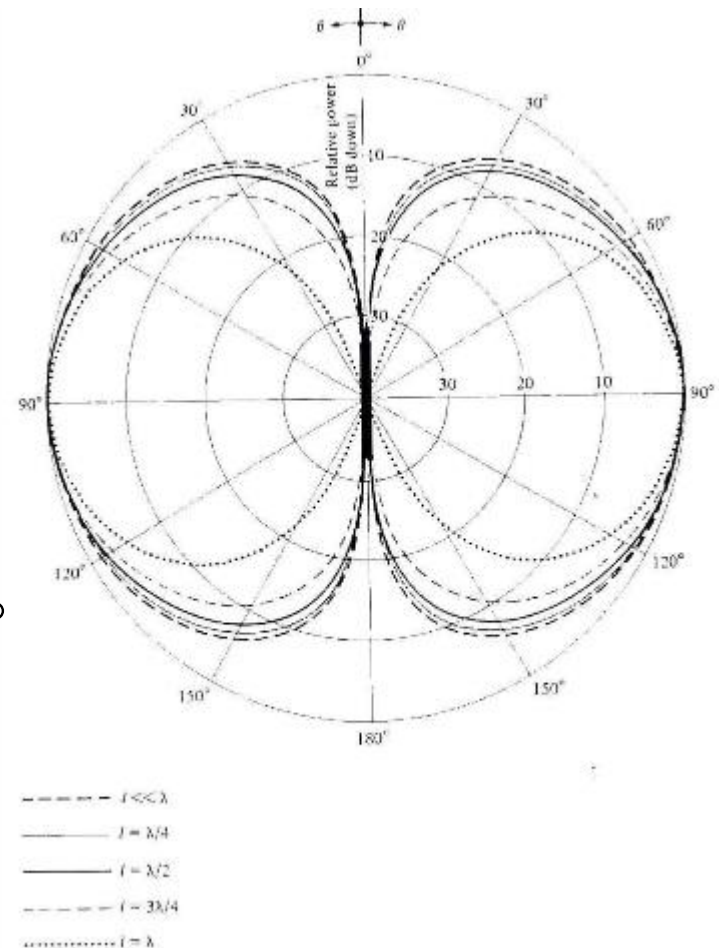
$$\vec{W}_{av} = \frac{1}{2} \text{Re}[\vec{E} \times \vec{H}^*] = \frac{1}{2} \text{Re}[\hat{\theta} E_{\theta} \times \hat{\phi} H_{\phi}^*] = \frac{1}{2} \text{Re}\left[\hat{\theta} E_{\theta} \times \hat{\phi} \frac{E_{\theta}^*}{\eta}\right]$$

$$\vec{W}_{av} = \hat{r} W_{av} = \hat{r} \frac{1}{2\eta} |E_{\theta}|^2 = \hat{r} \eta \frac{|I_0|^2}{8\pi^2 r^2} \left\{ \cos\left(\frac{kl}{2} \cos\theta\right) - \cos\left(\frac{kl}{2}\right) \right\}^2 / \sin^2 \theta$$

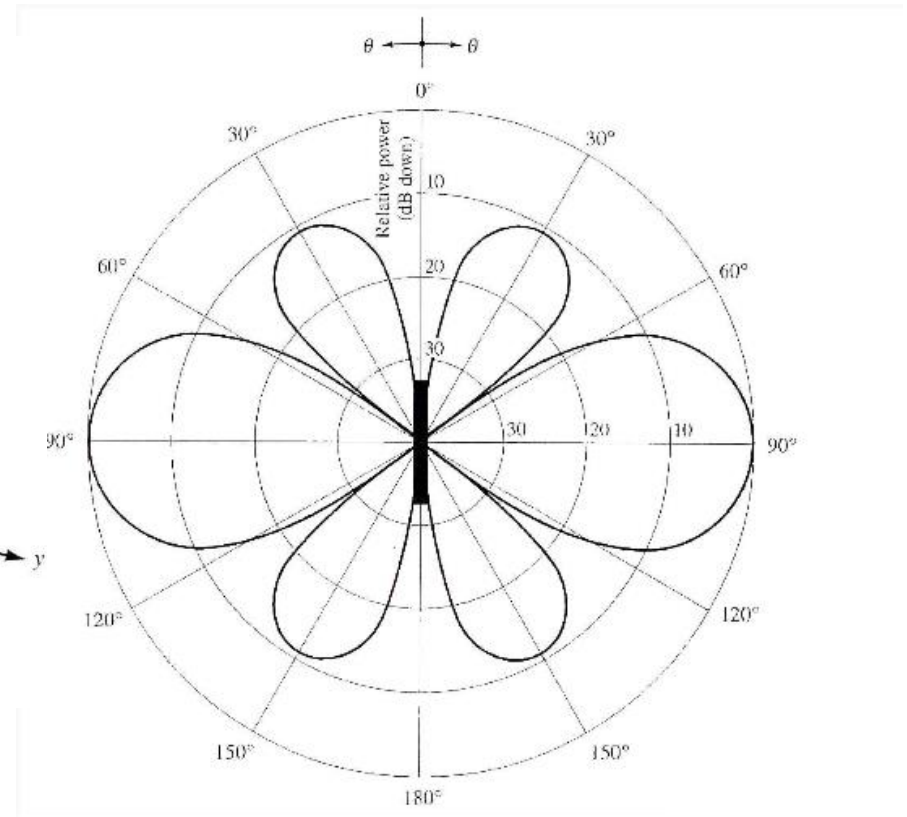
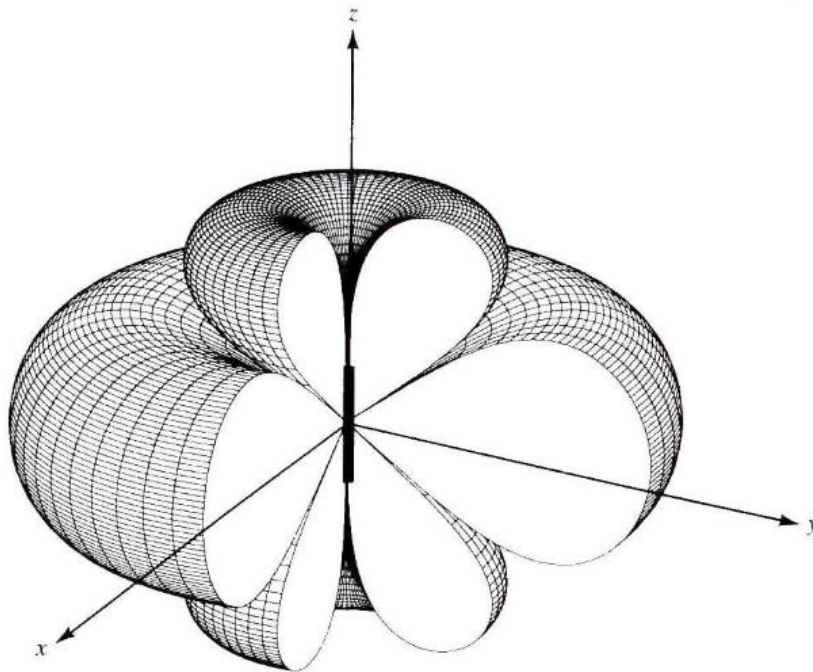
$$U = r^2 W_{av} = \eta \frac{|I_0|^2}{8\pi^2} \left[\frac{\cos\left(\frac{kl}{2} \cos\theta\right) - \cos\left(\frac{kl}{2}\right)}{\sin \theta} \right]^2$$

Radiation Patterns

- $l \ll \lambda$ → 3dB Beam width = 90°
- $l = \lambda/4$ → 3dB Beam width = 87°
- $l = \lambda/2$ → 3dB Beam width = 78°
- $l = 3\lambda/4$ → 3dB Beam width = 64°
- $l = \lambda$ → 3dB Beam width = 47.8°



Radiation Patterns for $l=1.25\lambda$



Radiated Power

$$P_{rad} = \oint_S \vec{W}_{av} \cdot d\vec{S} = \int_0^{2\pi} \int_0^\pi \hat{r} W_{av} \cdot \hat{r} r^2 \sin \theta \, d\theta \, d\varphi = \int_0^{2\pi} \int_0^\pi W_{av} r^2 \sin \theta \, d\theta \, d\varphi$$

$$P_{rad} = \int_0^{2\pi} \int_0^\pi W_{av} r^2 \sin \theta \, d\theta \, d\varphi = \eta \frac{|I_0|^2}{4\pi} \int_0^\pi \frac{\left[\cos\left(\frac{kl}{2} \cos \theta\right) - \cos\left(\frac{kl}{2}\right) \right]^2}{\sin \theta} d\theta$$

$$P_{rad} = \eta \frac{|I_0|^2}{4\pi} \left[C + \ln(kl) - C_i(kl) + \frac{1}{2} \sin(kl) \{ S_i(2kl) - 2S_i(kl) \} + \right. \\ \left. \frac{1}{2} \cos(kl) \{ C + \ln(kl/2) + C_i(2kl) - 2C_i(kl) \} \right]$$

Radiated Power

- $C=0.5772$

$$C_i(x) = -\int_x^{\infty} \frac{\cos y}{y} dy = \int_{\infty}^x \frac{\cos y}{y} dy \qquad S_i(x) = -\int_0^x \frac{\sin y}{y} dy$$

$$R_{rad} = \frac{2P_{rad}}{|I_0|^2} = \frac{\eta}{2\pi} \left[C + \ln(kl) - C_i(kl) + \frac{1}{2} \sin(kl) \{ S_i(2kl) - 2S_i(kl) \} + \right. \\ \left. \frac{1}{2} \cos(kl) \{ C + \ln(kl/2) + C_i(2kl) - 2C_i(kl) \} \right]$$

- Derivations for X_m are from Chapter 8 of text:

$$X_m = \frac{\eta}{4\pi} \left[2S_i(kl) + \cos(kl) [2S_i(kl) - S_i(2kl)] \right. \\ \left. - \sin(kl) [2C_i(kl) - C_i(2kl) - C_i\left(\frac{2ka^2}{l}\right)] \right]$$

Directivity

$$D_0 = 4\pi \frac{F(\theta, \varphi)_{\max}}{\int_0^{2\pi} \int_0^\pi F(\theta, \varphi) \sin \theta \, d\theta \, d\varphi}$$

■ $U = B_0 F(\theta, \varphi)$

$$F(\theta, \varphi) = F(\theta) = \left[\frac{\cos\left(\frac{kl}{2} \cos \theta\right) - \cos\left(\frac{kl}{2}\right)}{\sin \theta} \right]^2 \quad B_0 = \eta \frac{|I_0|^2}{8\pi^2}$$

$$D_0 = \frac{2F(\theta)|_{\max}}{\int_0^\pi F(\theta) \sin \theta \, d\theta} \quad \longrightarrow \quad D_0 = \frac{2F(\theta)|_{\max}}{Q}$$

$$Q = \left\{ C + \ln(kl) - C_i(kl) + \frac{1}{2} \sin(kl) [S_i(2kl) - 2S_i(kl)] \right. \\ \left. + \frac{1}{2} \cos(kl) [C + \ln(kl/2) + C_i(2kl) - 2C_i(kl)] \right\}$$

Input Resistance

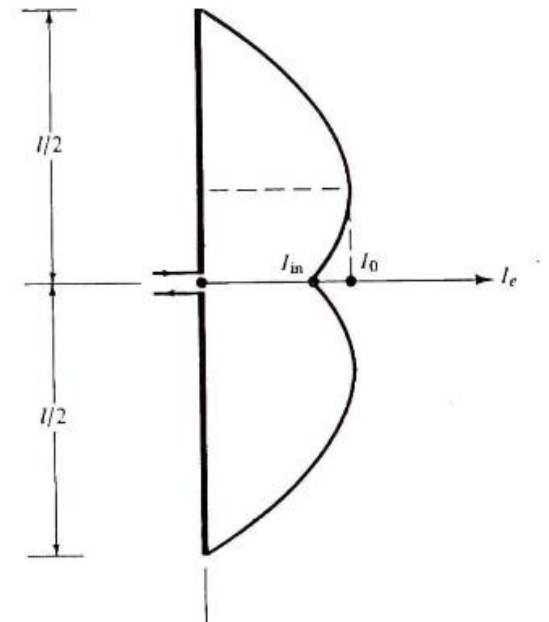
- We defined the input impedance as: the ratio of the voltage to the current and the input terminals of antennas.
- For a lossless antenna, the input resistance, which is the real part of the input impedance is related to the radiation resistance of the antenna.
- So far, we have calculated the radiation resistance of electrically small antennas as well as dipoles with sinusoidal current distribution. So far, we have used a definition, in which the radiation resistance is referred to the maximum current.
- For certain dipole antennas, this maximum does not occur at the input terminals (e.g., $l = \lambda$).

Input Resistance

- We can refer the radiation resistance to the input terminals of the antenna.
- Let's assume the antenna is lossless (i.e. $R_L = 0$).
- Then, we can equate the power at the input terminals of the antenna to the power at the current maximum.

$$\frac{|I_{in}|^2}{2} R_{in} = \frac{|I_0|^2}{2} R_r$$

$$R_{in} = \left[\frac{I_0}{I_{in}} \right]^2 R_r$$

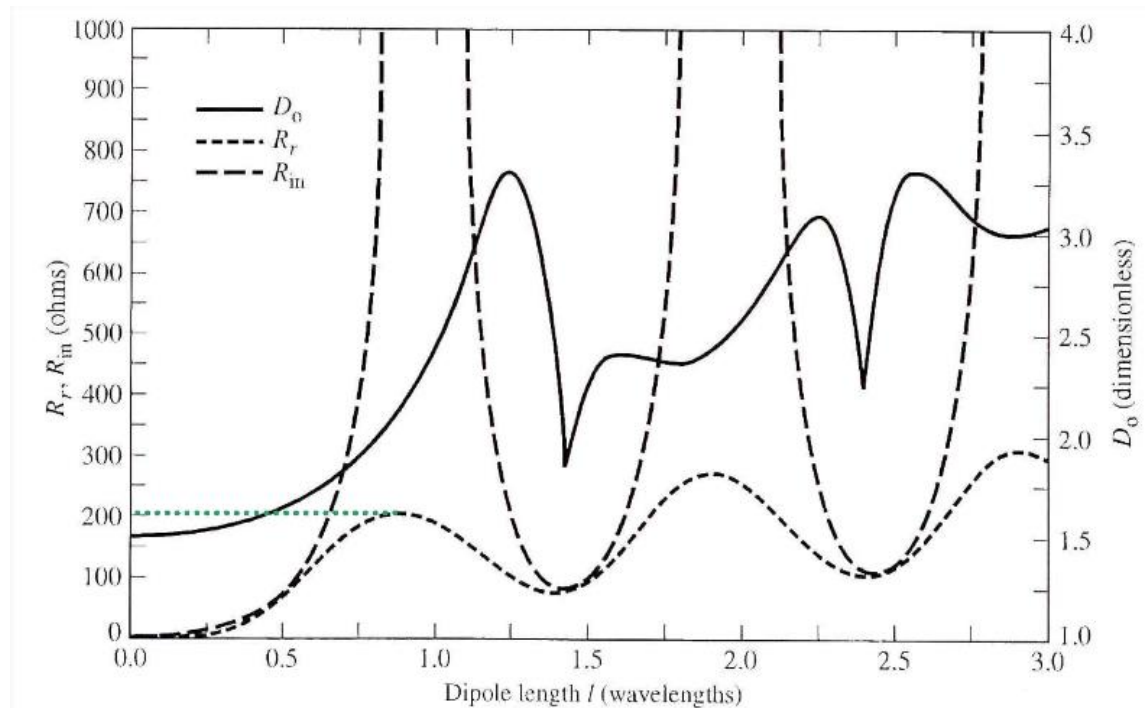


Input Resistance

- For a dipole of length l , the current at the input terminals (I_{in}) is related to the current maximum (I_0):

$$I_{in} = I_0 \sin(kl/2)$$

$$R_{in} = \frac{R_r}{\sin^2(kl/2)}$$



Half-Wavelength Dipole

$$E_{\theta} \approx j\eta \frac{I_0 e^{-jkr}}{2\pi r} \left[\frac{\cos\left(\frac{\pi}{2} \cos\theta\right)}{\sin\theta} \right]$$

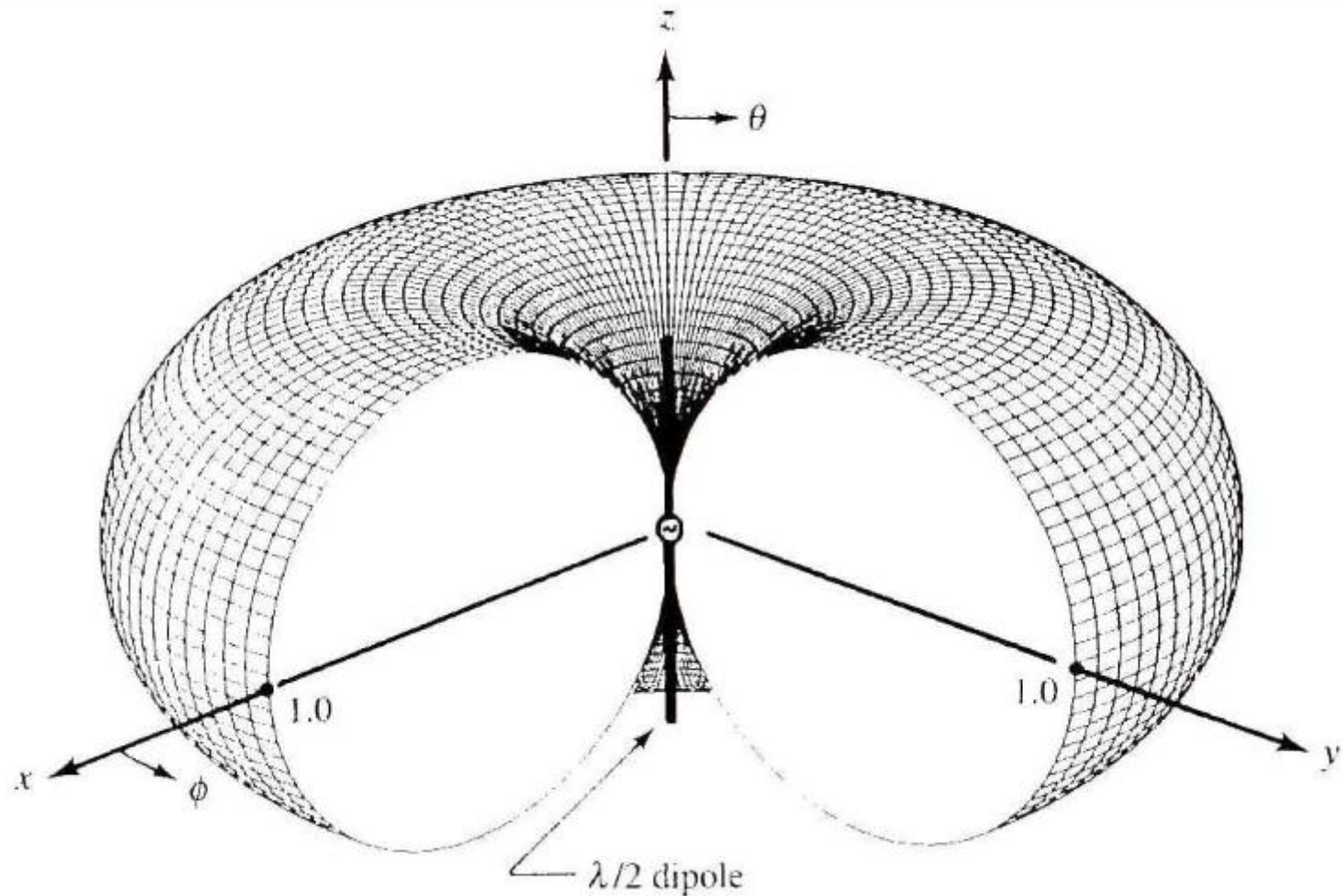
$$H_{\phi} \approx j \frac{I_0 e^{-jkr}}{2\pi r} \left[\frac{\cos\left(\frac{\pi}{2} \cos\theta\right)}{\sin\theta} \right]$$

$$W_{av} = \eta \frac{|I_0|^2}{8\pi^2 r^2} \left[\frac{\cos\left(\frac{\pi}{2} \cos\theta\right)}{\sin\theta} \right]^2 \approx \eta \frac{|I_0|^2}{8\pi^2 r^2} \sin^3 \theta \quad \longrightarrow \quad U = \eta \frac{|I_0|^2}{8\pi^2} \sin^3 \theta$$

$$P_{rad} = \eta \frac{|I_0|^2}{4\pi} \int_0^{\pi} \frac{\cos^2\left(\frac{\pi}{2} \cos\theta\right)}{\sin\theta} d\theta = \eta \frac{|I_0|^2}{8\pi} C_{in}(2\pi)$$

$$C_{in}(2\pi) = \int_0^{2\pi} \left(\frac{1 - \cos y}{y} \right) dy$$

Radiation Patterns of a Half-Wavelength Dipole



Half-Wavelength Dipole

- $C_{in}(2\pi)=2.435$.

$$D_0 = 4\pi U_{\max} / P_{rad} = 4\pi \frac{U|_{\theta=\pi/2}}{P_{rad}} = \frac{4}{C_{in}(2\pi)} \approx 1.643 = 2.15 \text{ dBi}$$

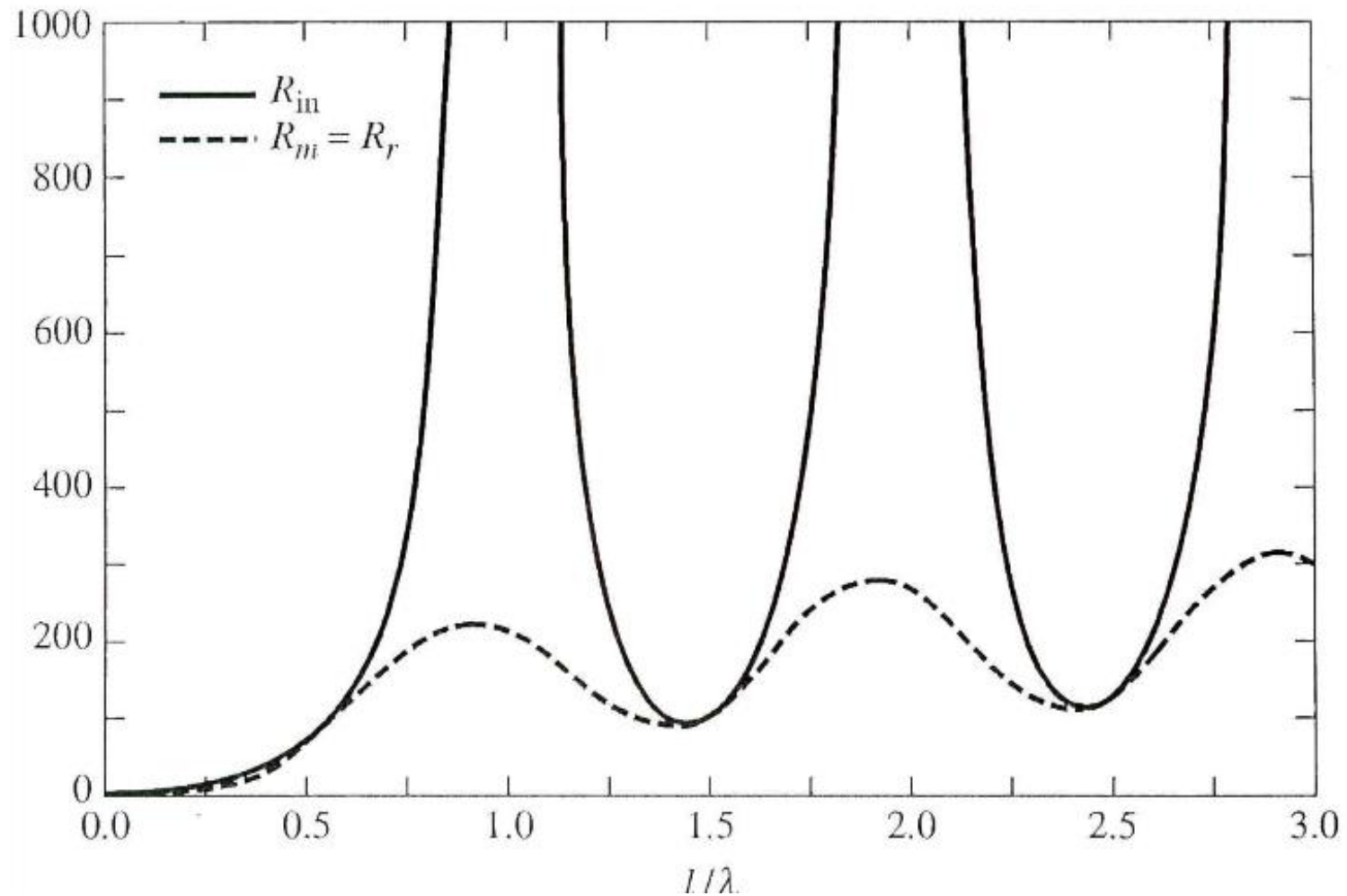
$$A_{em} = \frac{\lambda^2}{4\pi} D_0 \rightarrow A_{em} = \frac{\lambda^2}{4\pi} (1.64) \approx 0.13 \lambda^2$$

$$R_r \approx \frac{2P_{rad}}{|I_0|^2} = \frac{\eta}{4\pi} C_{in}(2\pi) = 30(2.43) \approx 73 \Omega$$

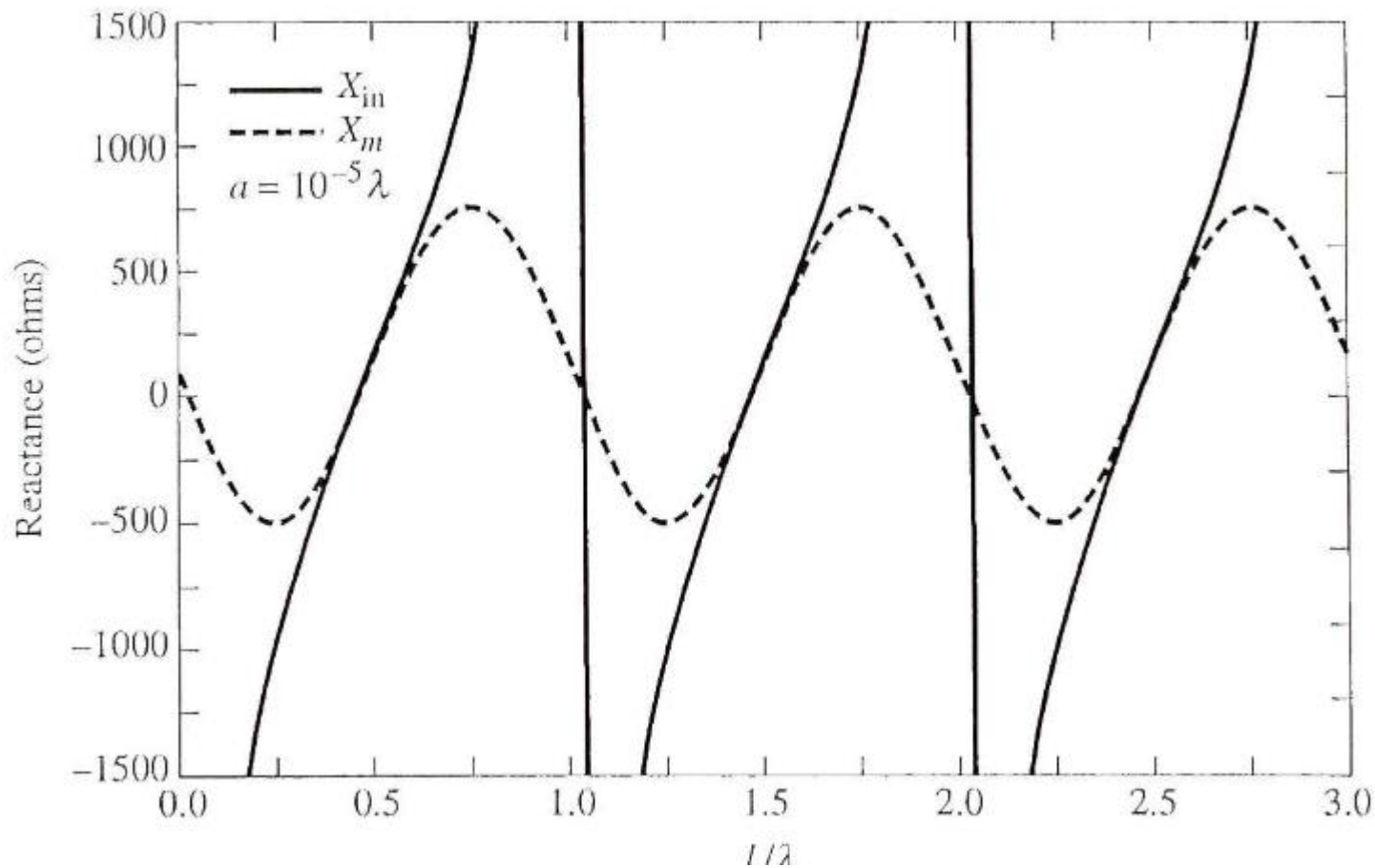
- The input impedance of a half-wavelength dipole antenna is:

$$Z_{in} = 73 + j43 \Omega$$

Dipole Antennas' Input Resistance



Dipole Antennas' Input Reactance



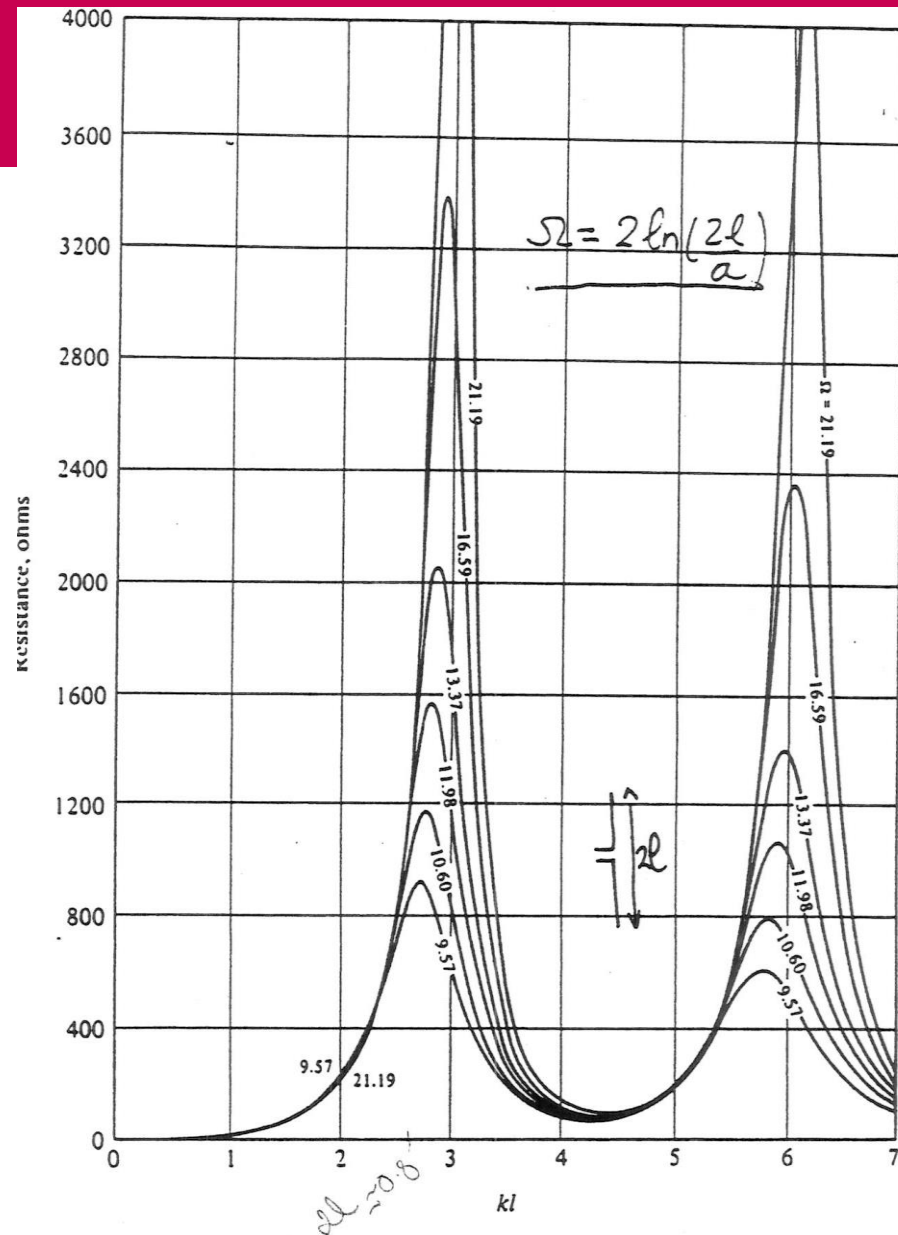


Fig. 7.15a Hallén's Curves of Resistance of a Center-Fed Cylindrical Dipole versus kl and Ω (Reprinted from E. Hallén, Cruft Laboratory Report No. 46, 1946, Courtesy of Harvard University.)

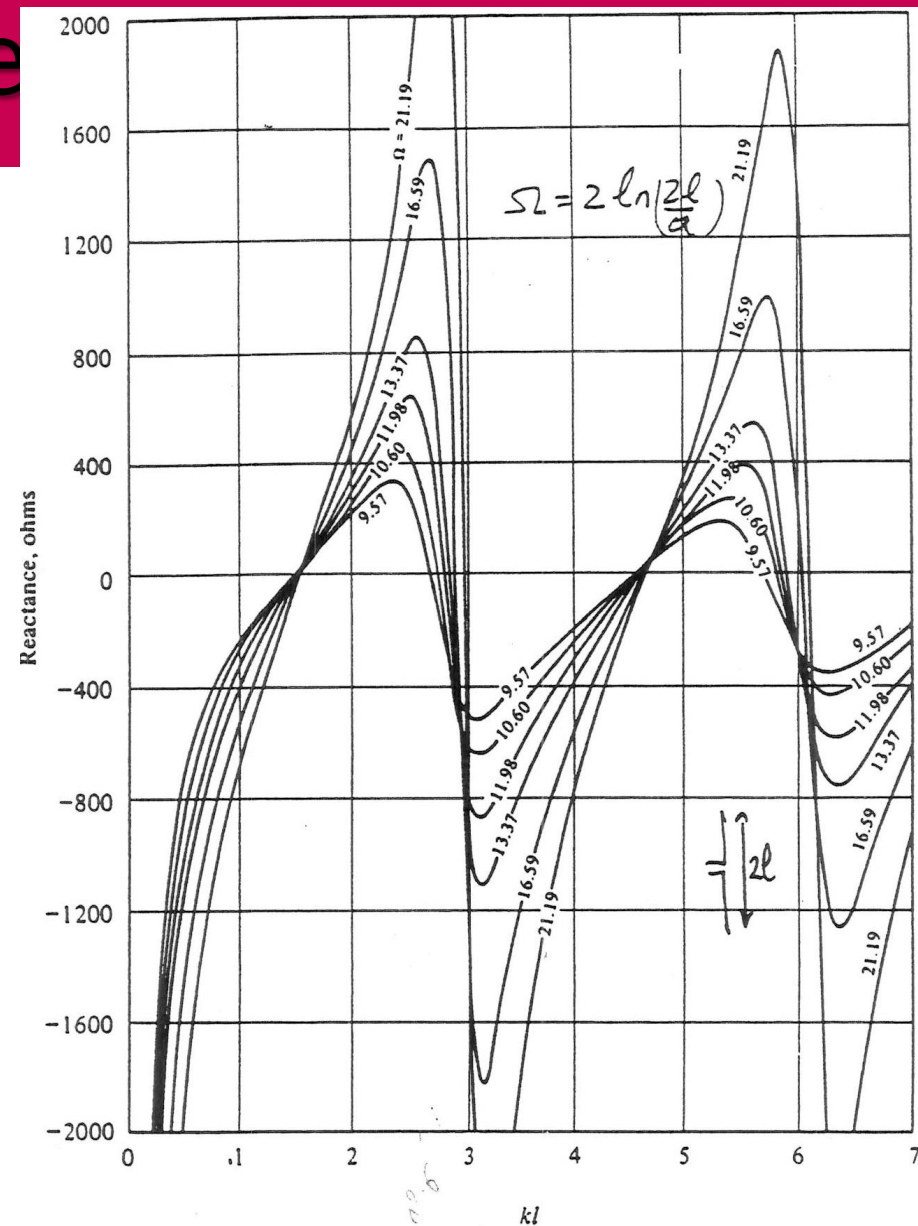


Fig. 7.15b Hallén's Curves of Reactance of a Center-Fed Cylindrical Dipole versus kl and Ω (Reprinted from E. Hallén, Cruft Laboratory Report No. 46, 1946, Courtesy of Harvard University.)

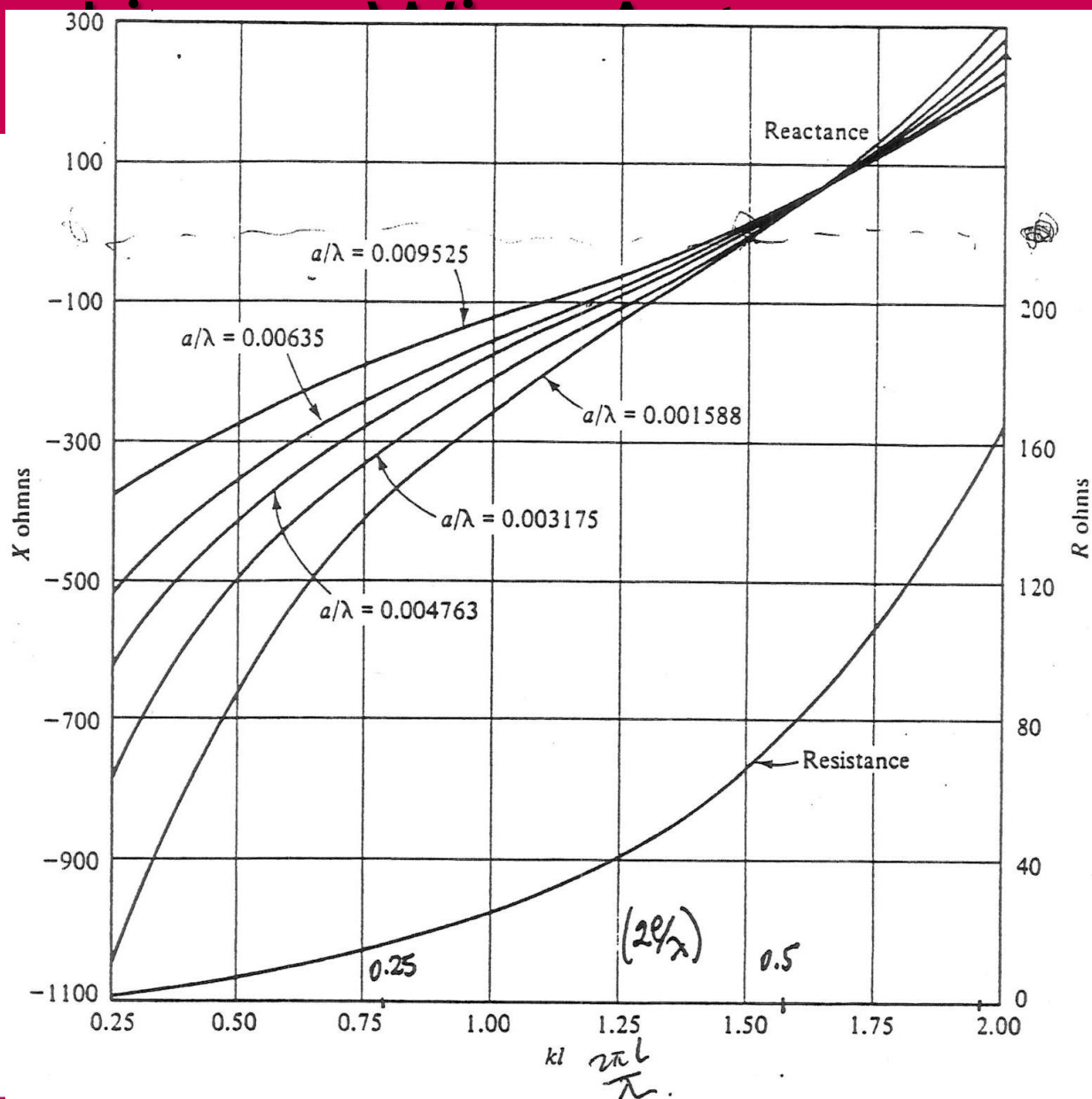


Fig. 7.11 The Resistance and Reactance of a Center-Fed Dipole versus kl and a/λ ; Values Computed by the Induced EMF Method Using Equation 7.65

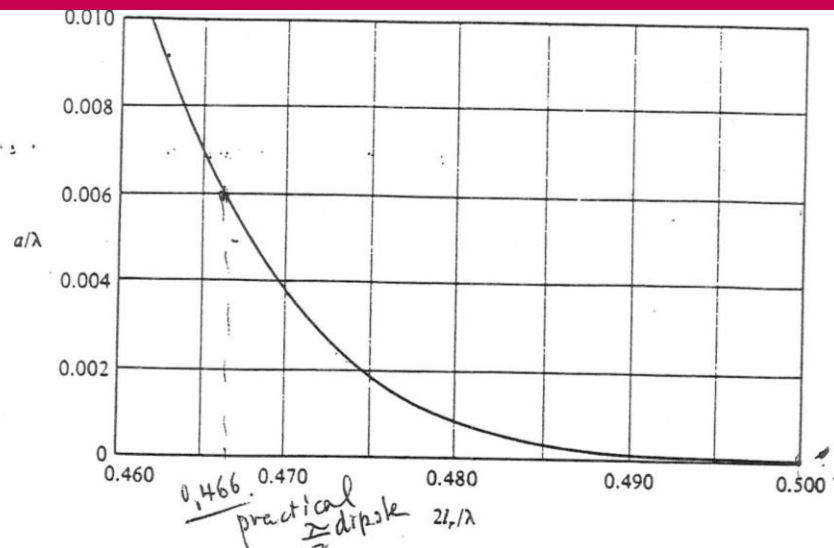


Fig. 7.13 Resonant Length versus Radius for Center-Fed Cylindrical Dipoles

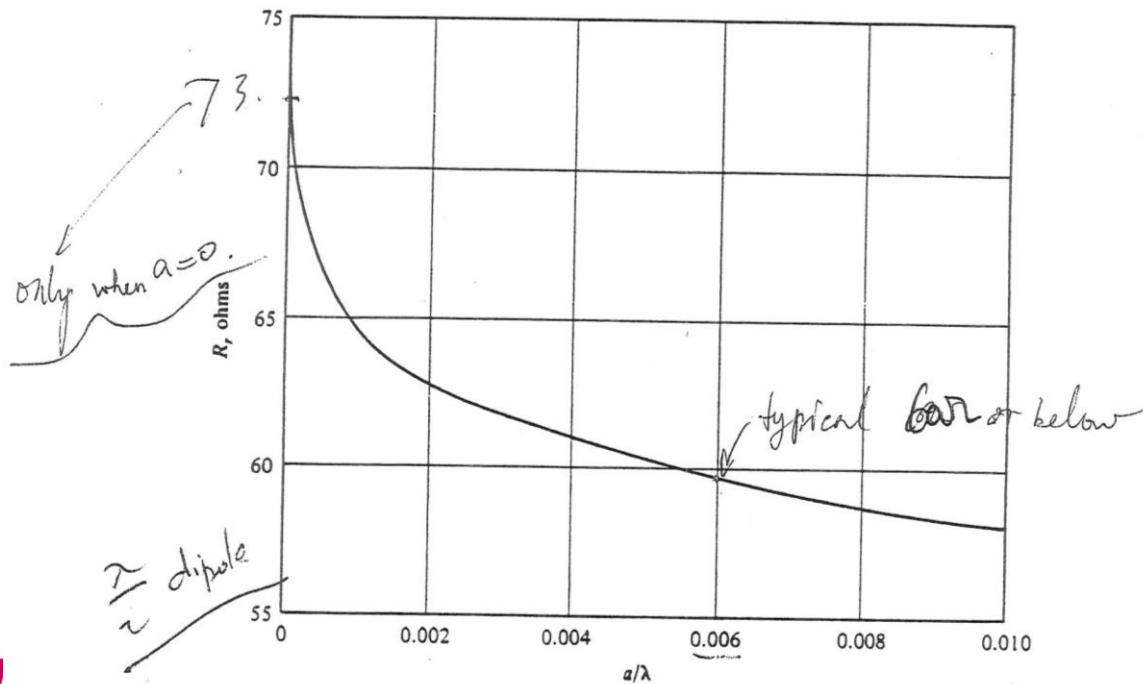


Fig. 7.14 Resonant Resistance versus Radius for Center-Fed Cylindrical Dipoles

Image Theory

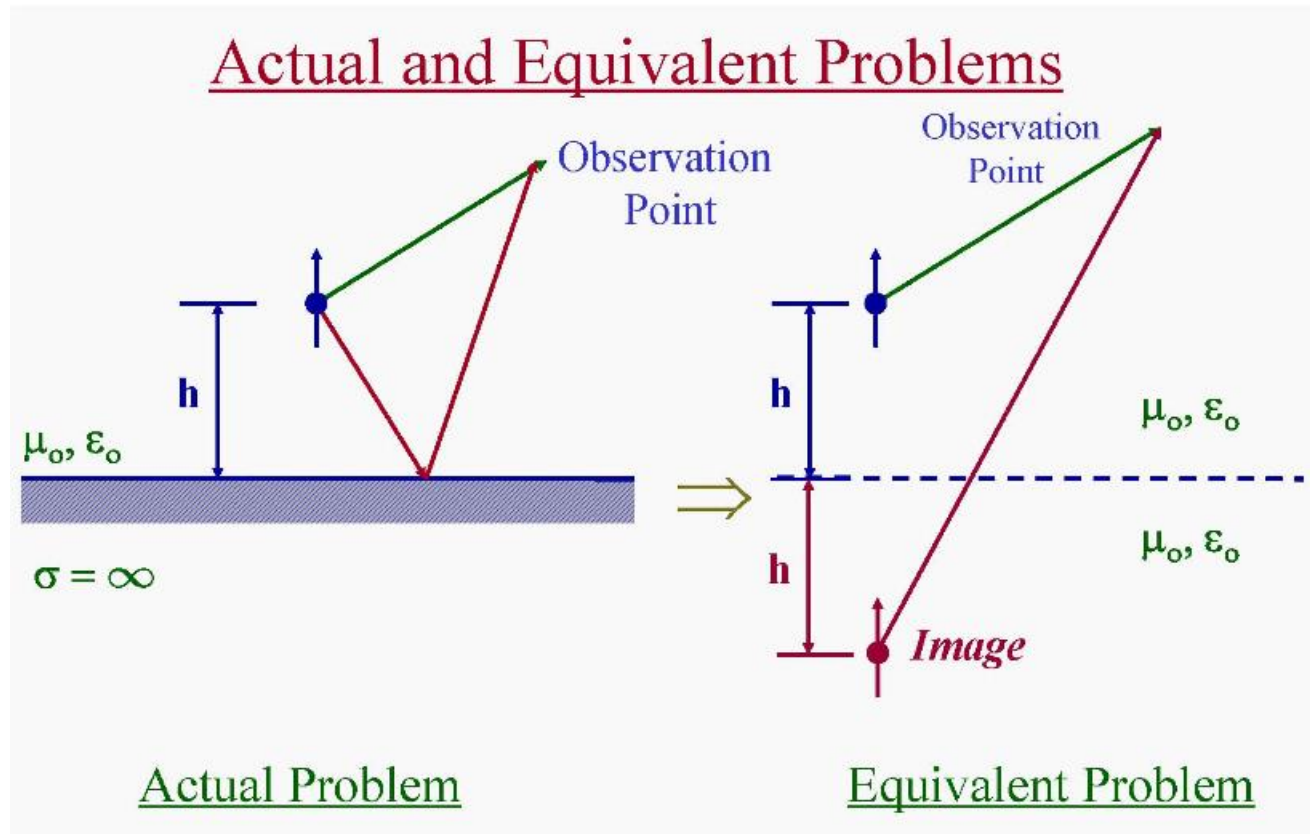


Image Theory: Perfect Electric Conductor (PEC)

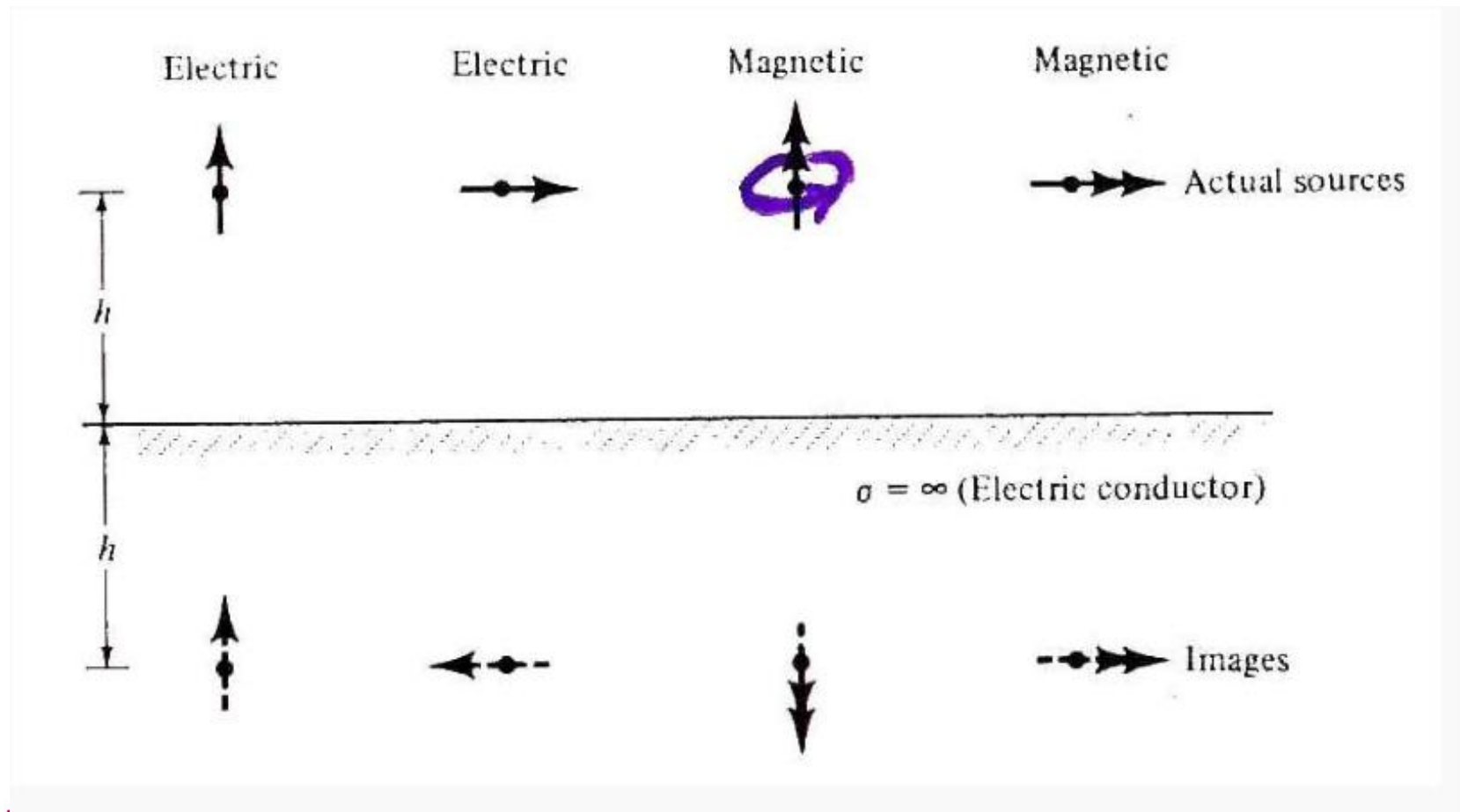
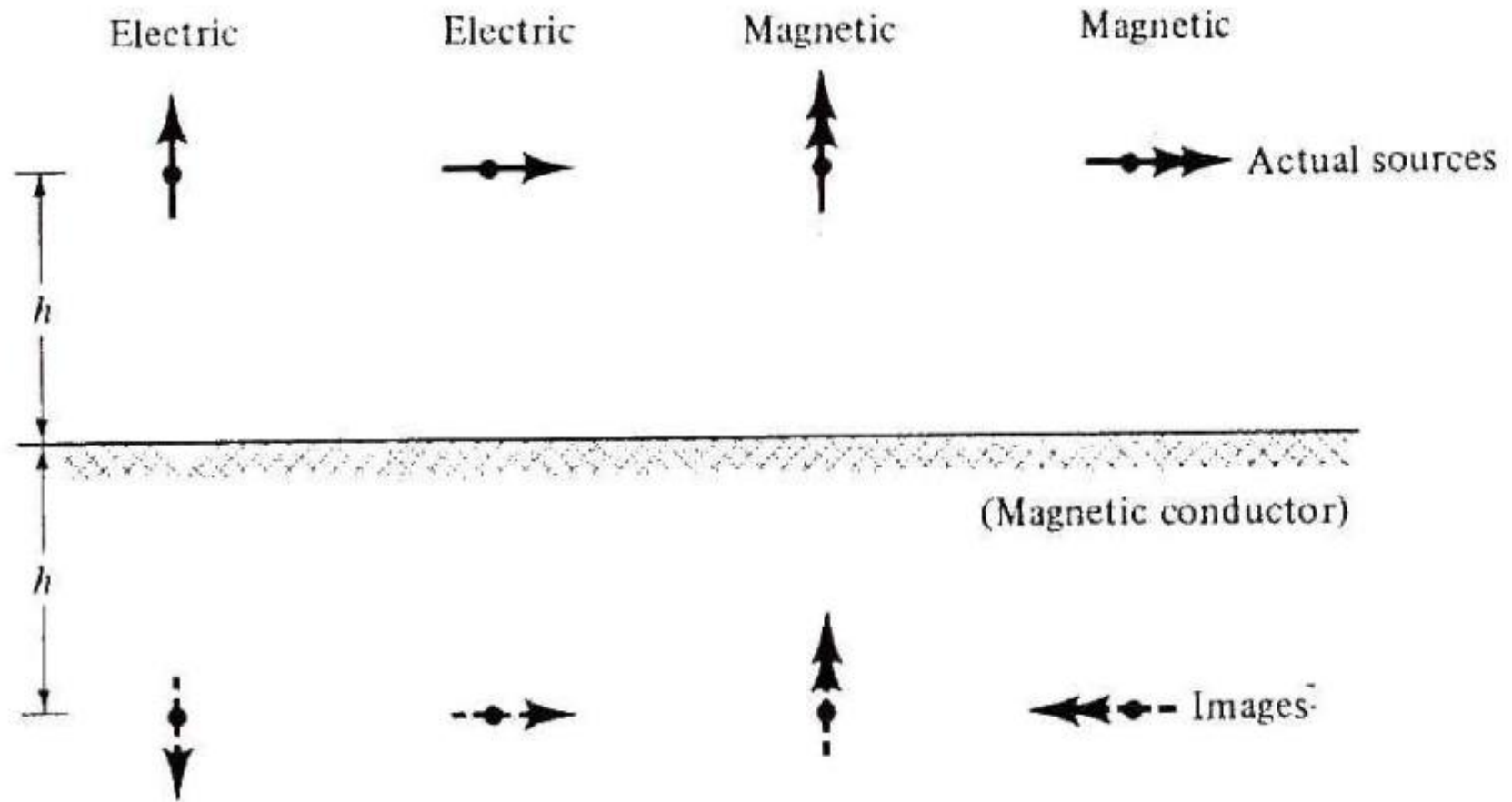
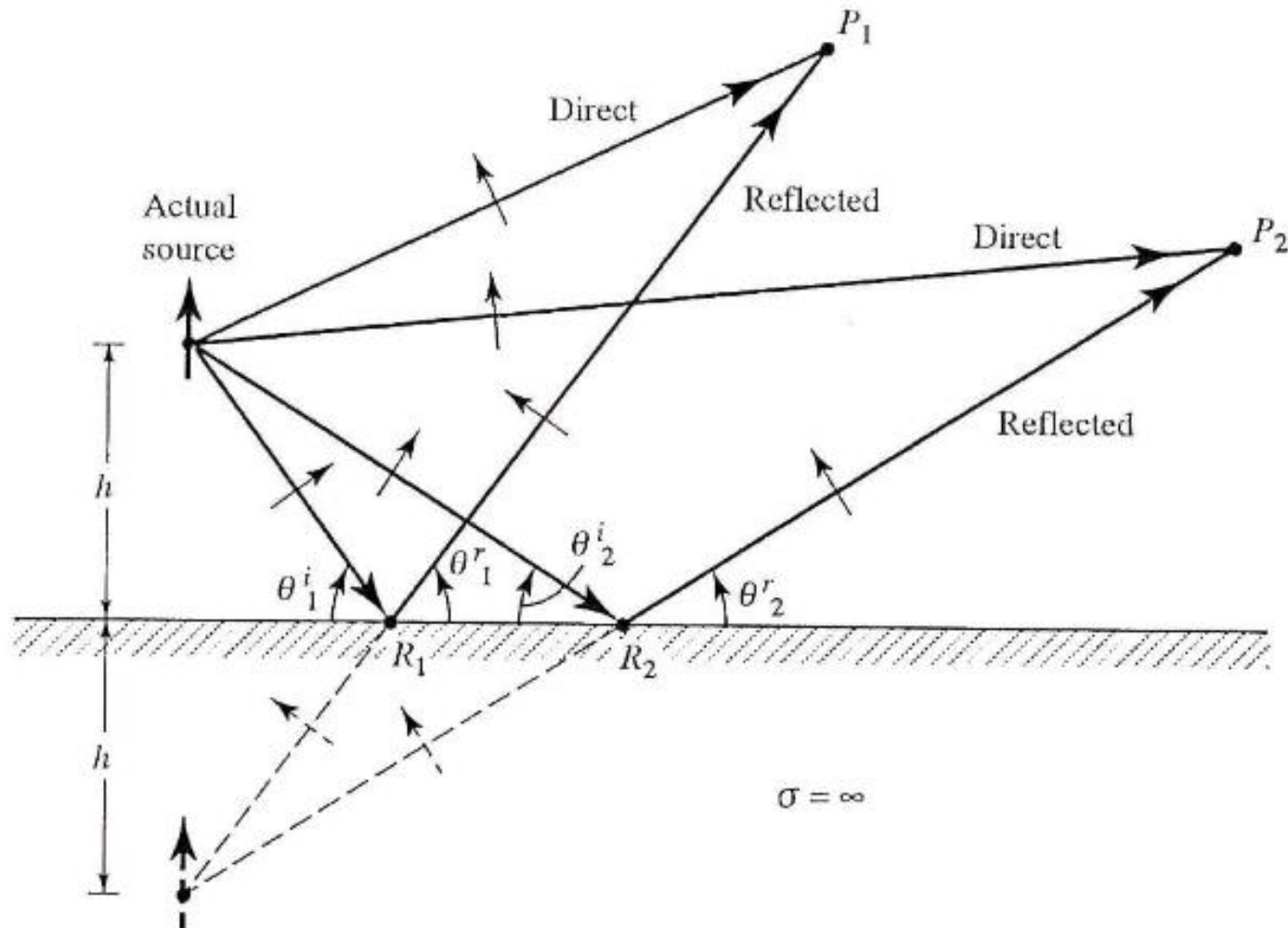


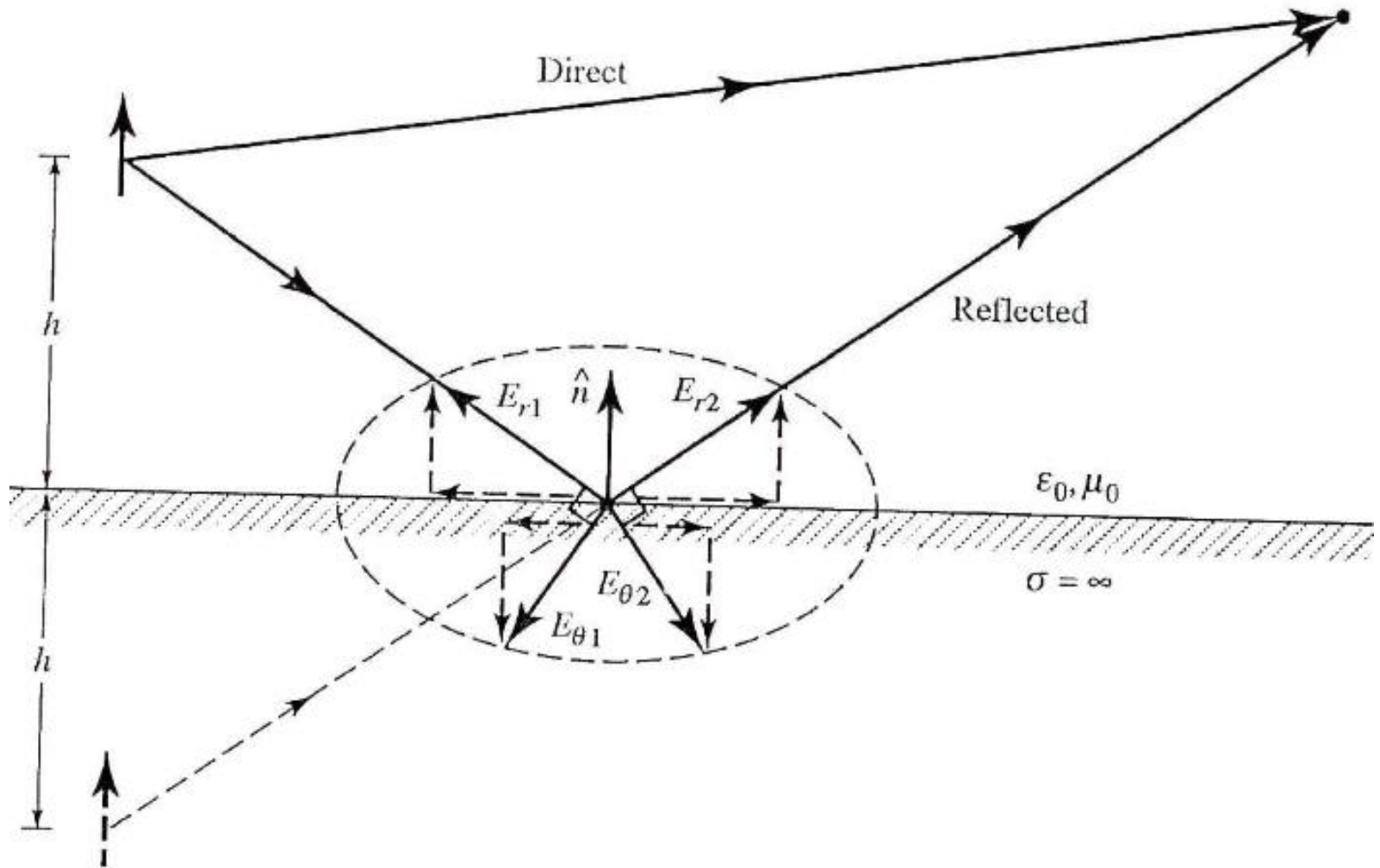
Image Theory: Perfect Magnetic Conductor



Vertical Electric Dipole and Image Theory (PEC)



Field Components at Point of Reflection

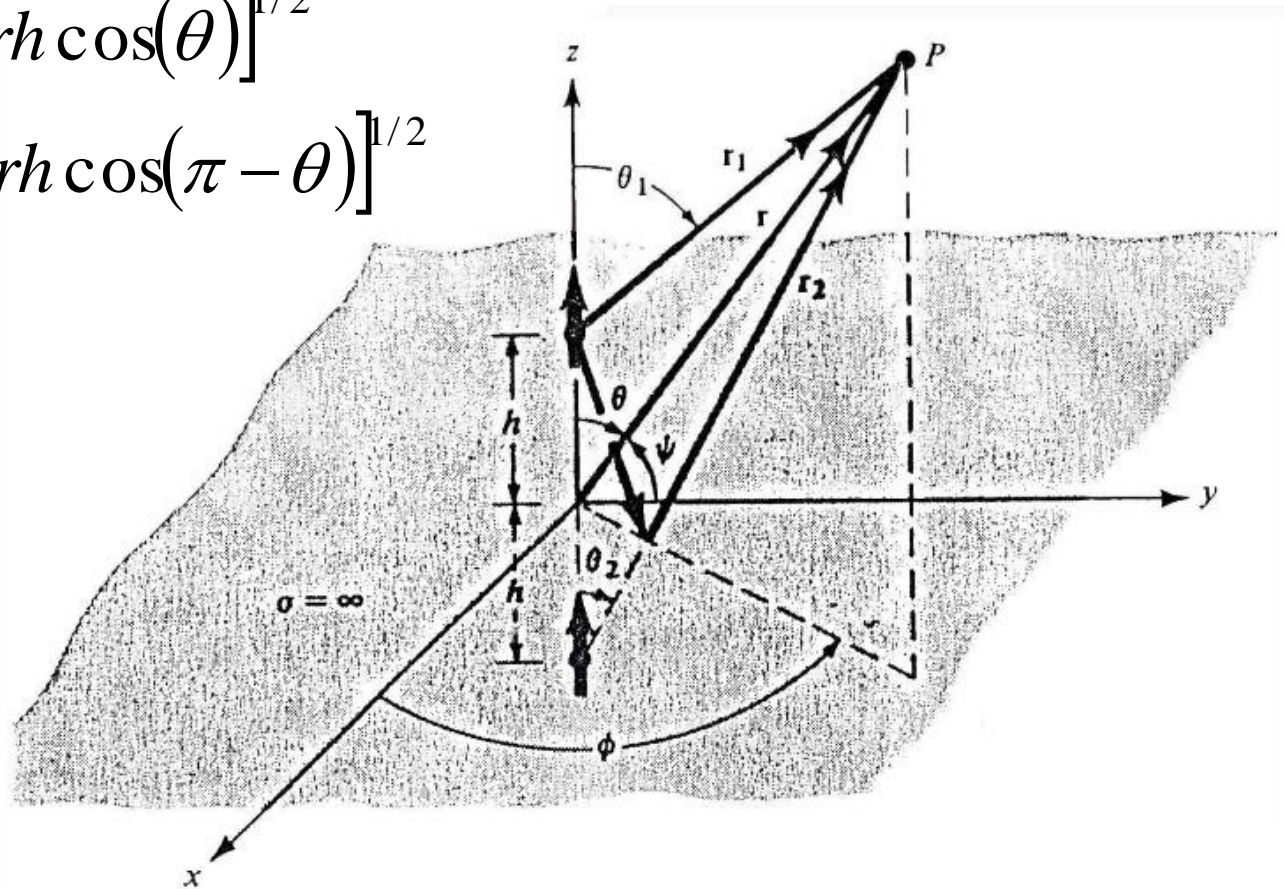


Analyzing the Vertical Electric Dipole Above PEC

- We use the image theory in conjunction with superposition:

$$r_1 = \left[r^2 + h^2 - 2rh \cos(\theta) \right]^{1/2}$$

$$r_2 = \left[r^2 + h^2 - 2rh \cos(\pi - \theta) \right]^{1/2}$$

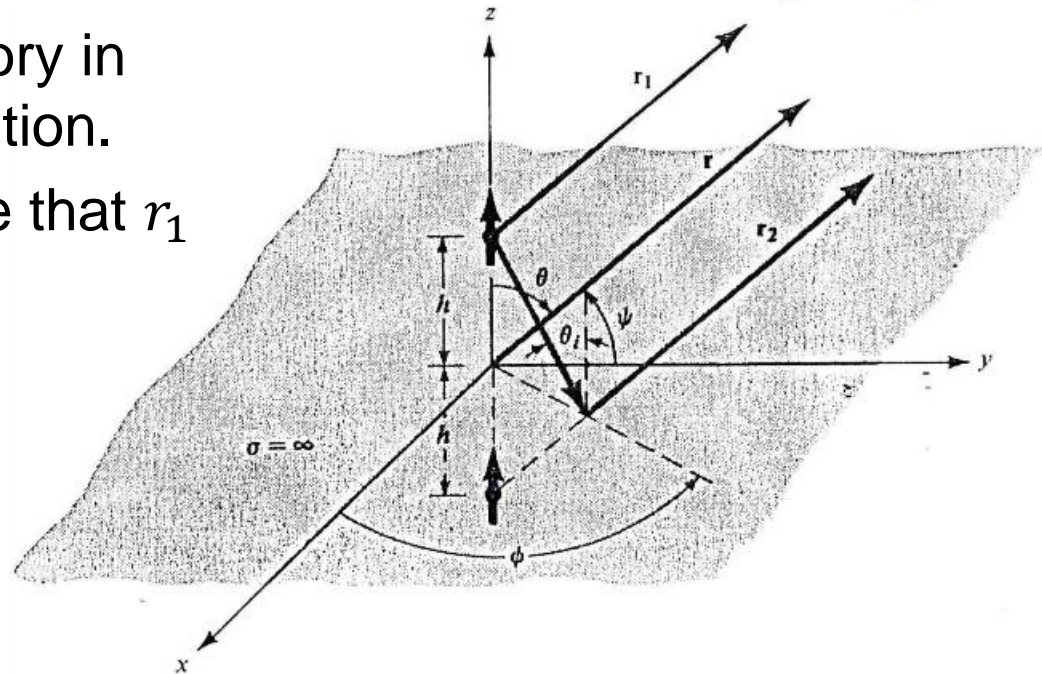


Analyzing the Vertical Electric Dipole Above PEC

- We will use the image theory in conjunction with superposition.
- In far field, we can assume that r_1 and r_2 are in parallel:

$$r_1 \approx r - h \cos(\theta)$$

$$r_2 \approx r + h \cos(\theta)$$



$$E_{\theta}^d = j\eta \frac{k l I_0 e^{-jk r_1}}{4\pi r_1} \sin \theta_1$$

$$E_{\theta}^r = j\eta \frac{k l I_0 e^{-jk r_2}}{4\pi r_2} \sin \theta_2$$