Class notes: Advanced Topics in Macroeconomics

Topic: Three Near-Linear Methods (cont)

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In last class, we introduced the third method for solving near-linear methods, but did not complete it. I repeat the notes here and then describe how to implement all for the most simple problem.

Method 3 (repeated)

Next, we use the insights of Vaughan (1970) to To avoid slow iteration on the Riccati equation, Vaughan (1970) exploits certain properties of the first-order conditions of the LQ problem defined above. (See his paper which is posted on the website.) Vaughan assumes no discounting or cross-product terms, so we'll map the variables and coefficients to  $\tilde{X}$ ,  $\tilde{u}$ ,  $\tilde{A}$ ,  $\tilde{B}$ , and  $\tilde{Q}$  as shown earlier. Also note that because the solution does not depend on the variances and covariances of  $\epsilon$ , we can abstract from the uncertainty for now. Writing out the Lagrangian, we have

$$\mathcal{L} = \sum_{t=0}^{\infty} \{ \tilde{X}_t' \tilde{Q} \tilde{X}_t + \tilde{u}_t' R \tilde{u}_t - \lambda_{t+1}' \left( X_{t+1} - \tilde{A} \tilde{X}_t - \tilde{B} \tilde{u}_t \right) \}$$
 (1)

Taking derivatives with respect to  $\tilde{u}_t$ ,  $\tilde{X}_{t+1}$ , and  $\lambda_{t+1}$ , we obtain the following first-order conditions

$$2R\tilde{u}_t + B'\lambda_{t+1} = 0$$

$$\tilde{Q}\tilde{X}_{t+1} - \lambda_{t+1} + \tilde{A}'\lambda_{t+2} = 0$$

$$\tilde{X}_{t+1} - \tilde{A}\tilde{X}_t - \tilde{B}\tilde{u}_t = 0$$
(2)

for  $t \geq 0$ , where  $\{\lambda_t\}$  is a sequence of Lagrange multipliers. Eliminating  $\tilde{u}_t$  and letting  $\tilde{\lambda}_t = 1/2\lambda_t$ , we have:

$$\begin{bmatrix} \tilde{X}_t \\ \tilde{\lambda}_t \end{bmatrix} = \begin{bmatrix} \tilde{A}^{-1} & \tilde{A}^{-1} \tilde{B} R^{-1} \tilde{B}' \\ \tilde{Q} \tilde{A}^{-1} & \tilde{Q} \tilde{A}^{-1} \tilde{B} R^{-1} \tilde{B}' + \tilde{A}' \end{bmatrix} \begin{bmatrix} \tilde{X}_{t+1} \\ \tilde{\lambda}_{t+1} \end{bmatrix}.$$

Let  $\mathcal{H}$  be the coefficient matrix on the right hand side. Vaughan showed that this matrix can be decomposed and used directly to obtain the Riccati matrix P (and hence the solution to the LQ problem); that is, he showed that

$$\mathcal{H} = \begin{bmatrix} V_{11} & V_{12} \\ V_{21} & V_{22} \end{bmatrix} \begin{bmatrix} \Lambda & 0 \\ 0 & \Lambda^{-1} \end{bmatrix} \begin{bmatrix} V_{11} & V_{12} \\ V_{21} & V_{22} \end{bmatrix}^{-1},$$

where the eigenvalues of  $\Lambda$  are outside of the unit circle. Notice that the eigenvalues come in reciprocal pairs. This is an important property that implies a unique stable solution, one that satisfies the transversality condition and ensures a bounded return.

Using the fact that the Lagrange multiplier is the derivative of the value function  $(\tilde{\lambda}_t = P\tilde{X}_t)$ , it is easy to figure out how to set P so as to get a stationary dynamical system for X. Let  $W = V^{-1}$ . In this case, it is easy to show that:

$$\tilde{X}_{t+1} = \{V_{11}\Lambda^{-1}(W_{11} + W_{12}P) + V_{12}\Lambda(W_{21} + W_{22}P)\}\tilde{X}_{t}.$$

Since  $\Lambda$  has roots outside the unit circle, it must be the case that  $P = -W_{22}^{-1}W_{21}$ . Note that since  $W = V^{-1}$ , this is equivalent to setting  $P = V_{21}V_{11}^{-1}$ .

In the case that  $\tilde{A}$  is not invertible, we can modify the method slightly and use generalized eigenvalues with the following alternative system:

$$\begin{bmatrix} \tilde{A} & 0 \\ -\tilde{Q} & I \end{bmatrix} \begin{bmatrix} \tilde{X}_t \\ \tilde{\lambda}_t \end{bmatrix} = \begin{bmatrix} I & \tilde{B}R^{-1}\tilde{B}' \\ 0 & \tilde{A}' \end{bmatrix} \begin{bmatrix} \tilde{X}_{t+1} \\ \tilde{\lambda}_{t+1} \end{bmatrix}.$$

Let  $\mathcal{H}_1$  be the coefficient matrix for the state and costate in t+1, and let  $\mathcal{H}_2$  be the coefficient matrix for the state and costate in t. Then, instead of taking eigenvalues of  $\mathcal{H}$  as above, we take generalized eigenvalues with the pair  $(\mathcal{H}_1, \mathcal{H}_2)$ .

Once we have a steady-state solution to the Riccati matrix, we can use the earlier formula to compute F and the law of motion for the state variables:

$$X_{t+1} = (A - BF)X_t + C\epsilon_{t+1}$$
(3)

Furthermore, given an initial condition for the states,  $X_0$ , and a realization of the shocks,  $\epsilon_t$ ,  $t \geq 0$ , we can generate time-series for  $X_t$  and  $u_t$ .

## A Simple Example

In class, three students are asked to go to the board and solve the following problem:

$$\max_{c_t} E_0 \ln (c_t)$$
 subject to  $c_t + k_{t+1} = Ak_t^{\theta}$ 

by applying the three different methods without any computer assistance. The student applying dynamic programming makes a guess of  $V_0(k)$  and iterates backwards. Things go well if they use  $V_0(k) = 0$  as an initial guess. The student applying an LQ method constructs a second-order Taylor expansion of the return function and a first-order Taylor expansion of the constraints. Things go well if they substitute out for  $c_t$  and use the fact that  $k_{t+1}$  is the control and equal to the state  $k_t$  in one period when writing out the linear constraints. The student applying Vaughan's method can see all of the insight of Vaughan by working with the second order difference equation in  $k_{t+2}$ ,  $k_{t+1}$ , and  $k_t$ . In particular, it is obvious that roots come in reciprocal (or  $\beta$ -reciprocal) pairs.