

# Enhanced positive water vapor feedback associated with tropical deep convection: New evidence from Aura MLS

Hui Su,<sup>1</sup> William G. Read,<sup>2</sup> Jonathan H. Jiang,<sup>2</sup> Joe W. Waters,<sup>2</sup> Dong L. Wu,<sup>2</sup> and Eric J. Fetzer<sup>2</sup>

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[1] Recent simultaneous observations of upper tropospheric (UT) water vapor and cloud ice from the Microwave Limb Sounder (MLS) on the Aura satellite provide new evidence for tropical convective influence on UT water vapor and its associated greenhouse effect. The observations show that UT water vapor increases as cloud ice water content increases. They also show that, when sea surface temperature (SST) exceeds  $\sim 300$  K, UT cloud ice associated with tropical deep convection increases sharply with increasing SST. The moistening of the upper troposphere by deep convection leads to an enhanced positive water vapor feedback, about 3 times that implied solely by thermodynamics. Over tropical oceans when SST greater than  $\sim 300$  K, the ‘convective UT water vapor feedback’ inferred from the MLS observations contributes approximately 65% of the sensitivity of the clear-sky greenhouse parameter to SST. **Citation:** Su, H., W. G. Read, J. H. Jiang, J. W. Waters, D. L. Wu, and E. J. Fetzer (2006), Enhanced positive water vapor feedback associated with tropical deep convection: New evidence from Aura MLS, *Geophys. Res. Lett.*, 33, L05709, doi:10.1029/2005GL025505.

## 1. Introduction

[2] Water vapor is one of the dominant greenhouse gases in Earth’s atmosphere and can have strong feedbacks on climate change. A prevailing hypothesis is that the atmosphere tends to maintain a fixed relative humidity profile and a constant lapse rate as the surface temperature increases [Manabe and Wetherald, 1967]. The resulting exponential increase in water vapor with temperature, based on the Clausius-Clapeyron relation, would trap more longwave radiation and thus cause further increase in surface temperature [Manabe and Wetherald, 1967; Stephens, 1990]. A clear-sky greenhouse parameter  $g$  is defined as  $g = (\sigma T_s^4 - OLR_{clr})/\sigma T_s^4$  [Raval and Ramanathan, 1989, hereinafter referred to as RR], where  $T_s$  is the sea surface temperature (SST),  $\sigma = 5.67 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$  is the Stefan-Boltzmann constant, and  $OLR_{clr}$  is the clear-sky outgoing longwave radiation (OLR). It has been found that  $g$  increases approximately linearly with SST for  $T_s < 298$  K, and at a much faster rate for  $T_s > 298$  K [RR]. This phenomenon is called the tropical ‘super greenhouse effect’, as the rate at which longwave trapping increases

with SST is larger than that for surface emission [Ramanathan and Collins, 1991, hereinafter referred to as RC]. Various mechanisms can affect the super greenhouse effect over tropical warm oceans, including atmospheric lapse rate change and tropospheric water vapor variation [Minschwaner and McElroy, 1992; Inamdar and Ramanathan, 1994]. In these studies, the contribution of upper tropospheric humidity (UTH) to the ‘super greenhouse effect’ is debatable, partly due to the lack of accurate measurements of UTH at that time.

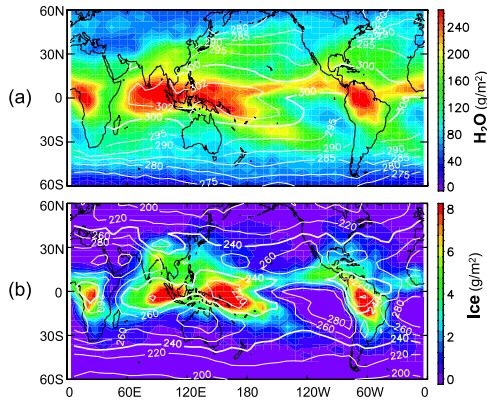
[3] It has been recognized that tropical deep convection and associated UT clouds are closely related to the variability of UTH. Evaporation of detrained upper-level condensates can be an important source for UTH [Betts, 1990; Sun and Lindzen, 1993]. Previous studies have demonstrated that average relative humidity in the 550–200 hPa layer tends to increase with the amount of deep convective clouds [Udelhofen and Hartmann, 1995, hereinafter referred to as UH; Soden and Fu, 1995, hereinafter referred to as SF]. The relationship between high-altitude clouds and SST is also of great interest. In particular, whether cirrus clouds could act as a ‘thermostat’ to regulate tropical SST has been intensely debated [RC; Wallace, 1992; Fu et al., 1992; Hartmann and Michelsen, 1993; Pierrehumbert, 1995; Sun and Liu, 1996]. Due to the lack of direct measurements of UT cloud profile, OLR, infrared brightness temperature or ‘highly reflective cloud’ (HRC) fraction based on visible satellite mosaics have been used as surrogates for cloud information [e.g., UH; SF; Collins et al., 1996; Waliser et al., 1993]. These indices give the area fraction and/or occurrence frequency of convective clouds, but not the mass of condensates (liquid water and ice). Our study utilizes direct measurements of UTH and cloud ice water content (IWC) from Aura MLS and examines the relationships among UTH, clouds, SST and the greenhouse effect of UT water vapor.

## 2. Data

[4] The Earth Observing System (EOS) MLS on Aura, launched on 15 July 2004, provides unprecedented simultaneous measurements of UT water vapor, IWC, and temperature [Schoeberl et al., 2006; Waters et al., 2006]. The MLS data (Version 1.5) include useful  $\text{H}_2\text{O}$  at altitudes 316 hPa and above, and IWC at 215 hPa and above with a vertical resolution about  $\sim 3$  km [Livesey et al., 2005; Froidevaux et al., 2006]. The MLS  $\text{H}_2\text{O}$  single measurement precision varies from  $\sim 25\%$  at 316 hPa to  $\sim 10\%$  at 100 hPa, with an expected accuracy of  $\sim 10\%$  [Livesey et al., 2005]. For IWC, the sensitivity of single measurement ranges from  $\sim 4 \text{ mg m}^{-3}$  at 215 hPa,  $\sim 1 \text{ mg m}^{-3}$  at 147 and  $\sim 0.4 \text{ mg m}^{-3}$  at 100 hPa [Livesey et al., 2005]. The

<sup>1</sup>Skillstorm/Jet Propulsion Laboratory, California Institute of Technology, Pasadena, California, USA.

<sup>2</sup>Jet Propulsion Laboratory, California Institute of Technology, Pasadena, California, USA.



**Figure 1.** Maps of the annual mean MLS-observed (a)  $\text{H}_2\text{O}$  and (b) Cloud ice. The observed SST (in K) and OLR (in  $\text{W m}^{-2}$ ) averaged over the same period are shown as white contours in Figures 1a and 1b, respectively.

uncertainty of the MLS IWC is dominated by a possible scaling error [Wu *et al.*, 2006], which is within a factor of 2. MLS cannot detect thin cirrus below the instrument sensitivity. It may also underestimate cloud ice with IWC greater than  $\sim 50 \text{ mg m}^{-3}$  due to saturation in the ice signal [Wu *et al.*, 2006]. The underestimated cloud ice is estimated to be less than 15% of total ice mass at 215 hPa and less than 5% at 147 hPa and higher levels.

[5] The MLS data have a horizontal resolution of  $\sim 200 \text{ km}$ , and measurements are made at  $\sim 1:30 \text{ am}$  and  $\sim 1:30 \text{ pm}$  local time at the low latitudes studied here. Effects of small scale variability are reduced by using monthly means in  $8^\circ$  longitude  $\times$   $4^\circ$  latitude grid boxes, from which annual averages are constructed. Each monthly average in a grid box contains 30–60 individual measurements. The estimated precisions of monthly-means in such grid boxes are within 5% for  $\text{H}_2\text{O}$ , and about 0.3 and  $2 \text{ mg m}^{-3}$  for IWC at 147 hPa and 215 hPa, respectively.

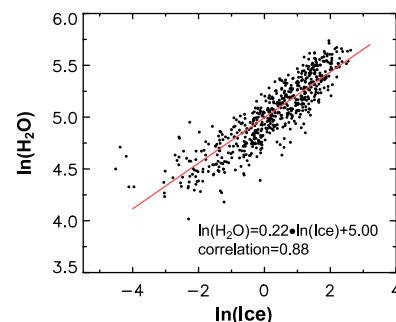
[6] We use the monthly optimum interpolated Reynolds SST (Version 2) [Reynolds and Smith, 1994] for the same period as for MLS data, and the SST data are re-gridded from  $1^\circ \times 1^\circ$  grids to the  $8^\circ \times 4^\circ$  MLS grids. Similar regridding is performed for the land-ocean mask associated with the Reynolds SST. We use total-sky OLR from the National Centers for Environmental Prediction (NCEP)/National Center for Atmospheric Research (NCAR) reanalysis provided by the Climate Diagnostics Center. For clear-sky OLR, we use data from the Atmospheric Infrared Sounder (AIRS) on board Aqua [Olsen, 2005]. They have been found to agree well with Clouds and the Earth's Radiant Energy System (CERES) clear-sky OLR [Suskind *et al.*, 2003]. The CERES OLR data concurrent with MLS observations were not available at the time of analysis. Our analysis is over oceanic regions between  $30^\circ\text{S}$  and  $30^\circ\text{N}$ . Here, we present results using annual means, based on monthly averages from August 2004 to July 2005. Similar results for the relationships among UTH, IWC and SST are found for the separate seasons throughout the year.

### 3. UT $\text{H}_2\text{O}$ and Cloud Ice

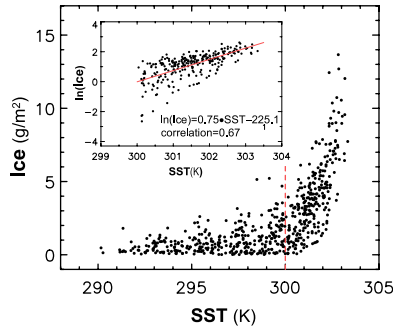
[7] We analyze the MLS  $\text{H}_2\text{O}$  data vertically integrated from 316 to 147 hPa, a layer that we will refer to as UT. For

cloud ice, we use the vertically integrated IWC from 215 to 147 hPa because the MLS IWC at 316 hPa is not a robust product in V1.5. The horizontal distributions of annual mean UT  $\text{H}_2\text{O}$  and cloud ice are shown in Figures 1a and 1b, respectively, with the annual mean SST and total-sky OLR for the same period plotted in white contours. The horizontal variation of cloud ice measured by MLS is generally consistent with results from general circulation models (GCMs) and from ECMWF analyses [Li *et al.*, 2005]. As shown in Figure 1b, the spatial distribution of MLS cloud ice over the tropics resembles that of OLR. Ice values greater than  $\sim 4 \text{ g m}^{-2}$  are coincident with OLR less than  $\sim 240 \text{ W m}^{-2}$ , indicative of deep convection [Graham and Barnett, 1987]. The pattern of UT  $\text{H}_2\text{O}$  shows remarkable similarity to that of cloud ice, with maxima over the inter-tropical convergence zone (ITCZ), including the western Pacific warm pool, and monsoonal regions of central Africa and South America. The high ice and low OLR over ocean are within the SST contour of  $\sim 300 \text{ K}$ . The eastern Pacific cold SST region is observed to have less cloud ice, and is drier than surrounding regions. Over the subtropics, both  $\text{H}_2\text{O}$  and cloud ice exhibit low values, consistent with relatively infrequent convection there.

[8] Figure 2 shows the scatter plot of annual mean UT  $\text{H}_2\text{O}$  versus cloud ice. Each point represents a grid-box ( $8^\circ \times 4^\circ$ ) over tropical oceans between  $30^\circ\text{S}$  and  $30^\circ\text{N}$ . The total number of points is 518. It is evident that UT  $\text{H}_2\text{O}$  is positively correlated with UT cloud ice. The correlation seen in Figure 2 is not a measurement artifact because the MLS  $\text{H}_2\text{O}$  and IWC measurements do not interact spectrally. The spatial correlation between  $\ln(\text{H}_2\text{O})$  and  $\ln(\text{Ice})$  is 0.88, above the 99% significance level with a minimum degree of freedom of 5, based on two-sided student's *t*-test. The variability of deep convective clouds explains approximately 77% of the variations of UT  $\text{H}_2\text{O}$ . The slope of the linear fit to the data is about  $0.22 \pm 0.005(2\%)$  and is not affected by any scaling error in the MLS data. Additional error of the slope associated with IWC measurement sensitivity is estimated to be less than 2%. Similar positive correlation between  $\text{H}_2\text{O}$  and IWC is found for the individual layers at 215 and 147 hPa, with correlation coefficient 0.85 and 0.76, respectively. Relative humidity with respect to ice (RH<sub>i</sub>) at these levels also increases with IWC (not shown), with correlation coefficients around 0.8 for both levels. These results are consis-



**Figure 2.** Scatter plot of MLS-observed UT  $\text{H}_2\text{O}$  and cloud ice (both in  $\text{g m}^{-2}$ , plotted in natural logarithmic scales), with the least squares linear fit in red.



**Figure 3.** Scatter plot of cloud ice versus SST. The inset is the scatter plot of the natural logarithm of cloud ice versus SST for  $T_s > 300$  K, with the least squares linear fit in the red solid line. The red dashed line marks  $T_s = 300$  K.

tent with previous studies and support the view that the upper-troposphere is moistened by vertical transport of moisture within convective towers and by evaporation of detrained ice [UH; SF]. The positive correlation between UT  $H_2O$  and cloud ice is also found for tropical land regions, with similar values of correlation.

#### 4. Deep Convection and SST

[9] It is known that tropical deep convection and associated clouds are closely related to SST, SST gradient and low-level moisture convergence [Lin *et al.*, 2006, and references therein]. Here, we examine the connection between annual mean UT cloud ice and SST. Figure 3 shows the annual mean tropical cloud ice within the layer between 215 hPa and 147 hPa, scattered against SST. A notable feature is the sharp increase in cloud ice at SST greater than  $\sim 300$  K. This is consistent with previous studies that find a convective threshold at SST values of 300–301 K [e.g., Graham and Barnett, 1987; Waliser *et al.*, 1993].

[10] Most importantly, Figure 3 shows that UT cloud ice associated with deep convection (Ice greater than  $\sim 4$  g  $m^{-2}$ ) increase sharply as SST increases for  $T_s > \sim 300$  K. The logarithm of cloud ice varies approximately linearly with SST above 300 K (see the Figure 3 inset) with a correlation coefficient of 0.67, which is statistically significant at the 95% level with a minimum degree of freedom of 7. The slope of  $\ln(\text{Ice})$  versus SST is about  $0.75 \pm 0.04$  (5%)  $K^{-1}$ .

[11] Figure 3 gives the first direct observational evidence that UT IWC associated with deep convection increases with SST for  $T_s > \sim 300$  K. This increase in IWC within a grid box of  $8^\circ \times 4^\circ$  for a given period could be due to the increase of occurrence frequency of clouds and/or the increase in IWC for individual convective events. This exponential increase of cloud ice with SST is consistent with conventional wisdom that tropical deep convection is an increasing function of SST (for  $T_s > \sim 300$  K) [RC; Waliser *et al.*, 1993; Collins *et al.*, 1996]. Moreover, Figure 3 suggests that UT cloud formation efficiency also increases with SST [Lin *et al.*, 2006]. In the study of Waliser *et al.* [1993], the HRC index decreases with extremely high SST for  $T_s > \sim 303$  K. This is not shown in our analysis, probably

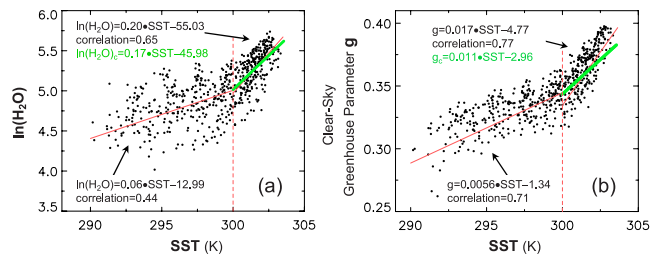
due to the relatively coarse spatial and temporal scales examined here.

#### 5. Enhanced UT $H_2O$ Feedback

[12] Figure 4a shows the scatter plot of the natural logarithm of UT  $H_2O$  versus SST. For  $T_s < 300$  K, the  $\ln(H_2O)$  versus SST slope is  $0.060 K^{-1}$ , close to the theoretical slope based on the Clausius-Clapeyron equation,  $0.064 K^{-1}$  [Stephens, 1990]. For  $T_s > 300$  K, the  $\ln(H_2O)$  versus SST slope is 3.3 times larger,  $0.20 K^{-1}$ . This implies that for  $T_s > 300$  K dynamic control of UT  $H_2O$  dominates thermodynamic control. Combining the observed correlations between UT  $H_2O$  and cloud ice (Figure 2), and between cloud ice and SST for  $T_s > 300$  K (Figure 3), we can infer the convection-induced UT  $H_2O$  sensitivity to SST, denoted as  $(d \ln H_2O / dT_s)_c$ . We define

$$\left( \frac{d \ln H_2O}{dT_s} \right)_c = \frac{\partial \ln H_2O}{\partial \ln IWC} \cdot \frac{\partial \ln IWC}{\partial T_s}, \quad (1)$$

where  $\partial \ln H_2O / \partial \ln IWC = 0.22$ , and  $\partial \ln IWC / \partial T_s = 0.75 K^{-1}$  based on the least squares linear fit slopes shown in Figures 2 and 3 inset. This yields a *convective* UT  $H_2O$  sensitivity to SST  $(d \ln H_2O / dT_s)_c \approx 0.17 K^{-1}$  (the green line in Figure 4a), about 85% of the observed UT  $H_2O$  sensitivity to SST for  $T_s > 300$  K. It is 2.8 times as large as the slope for  $T_s < 300$  K. Considering the propagation of estimate errors for the two terms on the r.h.s. of equation (1), the error of the inferred  $(d \ln H_2O / dT_s)_c$  is about 10%. The likely MLS underestimate of cloud ice for IWC greater than  $50 \text{ mg m}^{-3}$  may lead to a potential underestimate of the slope  $(\partial \ln IWC / \partial T_s)$ , implying an even stronger convection-induced UT water vapor sensitivity to SST than the rate presented here (up to 15% stronger). The closeness of the observed  $d \ln H_2O / dT_s$  and the inferred *convective*  $(d \ln H_2O / dT_s)_c$  ( $0.20 K^{-1}$  versus  $0.17 K^{-1}$ ) is strong evidence that the sharp increase in UT  $H_2O$  for  $T_s > 300$  K is primarily caused by deep convection. The correlation coefficients for observed  $\ln(H_2O)$  and SST are 0.44 for  $T_s < 300$  K and 0.65 for  $T_s > 300$  K, above the 95% significance



**Figure 4.** Scatter plots (a) of the natural logarithm of UT  $H_2O$  (in  $g m^{-2}$ ) versus SST; and (b) of greenhouse parameter  $g$  versus SST. The red solid lines are the separate least squares fits for  $T_s < 300$  K and  $T_s > 300$  K. The green lines are the *convective* UT  $H_2O$  sensitivity to SST (Figure 4a) and sensitivity of  $g$  to SST (Figure 4b), as inferred from MLS data (see text for details). The red dashed lines mark  $T_s = 300$  K.



level with a minimum degree of freedom of 19 and 8, respectively.

[13] The UT mean RH<sub>i</sub> also bears an interesting relation with SST (figure not shown). Over regions of SST lower than 300 K, UT RH<sub>i</sub> is nearly constant, in line with the fixed relative humidity hypothesis [Manabe and Wetherald, 1967]. However, UT RH<sub>i</sub> increases rapidly with SST for  $T_s > \sim 300$  K.

[14] Using the clear-sky OLR data from AIRS, we compute the greenhouse parameter  $g$  as defined in RR. It is found  $g$  varies approximately linearly with the natural logarithm of UT H<sub>2</sub>O from MLS, with a slope  $0.063 \pm 0.001$  (2%) (figure not shown). The correlation coefficient between the two is 0.90, above the 99% significance level with a minimum degree of freedom of 5.

[15] Figure 4b shows the scatter plot of  $g$  versus SST. It is clear that there are two different slopes for two regimes of SST. For  $T_s < 300$  K, the best fit slope of  $g$  is  $5.6 \times 10^{-3} \text{ K}^{-1}$ . For  $T_s > 300$  K,  $g$  varies with SST at a rate about  $17 \times 10^{-3} \text{ K}^{-1}$ , nearly 3 times of that for  $T_s < 300$  K. The difference in the sensitivity of  $g$  to SST in the two regimes is similar to that of the UT H<sub>2</sub>O sensitivity to SST shown in Figure 4a, suggesting the ‘super greenhouse effect’ in the high SST range could be largely due to increasing UT water vapor associated with increasing UT clouds. As done for the convection-induced UT H<sub>2</sub>O sensitivity to SST, we can infer the convection-induced sensitivity of  $g$  to SST,  $(dg/dT_s)_c$ . We write

$$\left(\frac{dg}{dT_s}\right)_c = \frac{\partial g}{\partial \ln H_2O} \left(\frac{d \ln H_2O}{dT_s}\right)_c. \quad (2)$$

With the value of  $(d \ln H_2O/dT_s)_c$  derived before and the slope of  $g$  versus  $\ln(H_2O)$ , we obtain  $11.0 \times 10^{-3} \text{ K}^{-1}$  as the convective sensitivity of  $g$  to SST (shown in green in Figure 4b), with an error about 15% (10% error associated with the least squares linear fit slope estimates and 5% error from IWC retrieval precision). This convective sensitivity of  $g$  versus SST is approximately 2 times the sensitivity of  $g$  to SST for  $T_s < 300$  K and accounts for about 65% of the observed clear-sky greenhouse parameter sensitivity to SST. Other factors, such as lower tropospheric water vapor variations and atmospheric lapse rate change, can also impact the sensitivity of  $g$  to SST, but their contributions are smaller than that from the UT H<sub>2</sub>O.

## 6. Implications for Climate Change

[16] The dominance of the convection-induced  $(dg/dT_s)_c$  for the observed  $dg/dT_s$  is a strong evidence of ‘convective UT water vapor feedback’ as an important mechanism for the ‘super greenhouse effect’ over tropical oceans for  $T_s > \sim 300$  K. The key elements in this feedback are (1) the rapid increase of cloud ice with SST for  $T_s > \sim 300$  K (Figure 3) and (2) the increase of UTH with cloud ice (Figure 2). This ‘convective UT water vapor feedback’ is about 3 times larger than the feedback expected solely from thermodynamics. Thus it warrants great attention for accurate simulations of climate change.

[17] On the other hand, tropical SST has been relatively stable over the past 100 million years [Crowley and North, 1991]. There must be a strong regulating mechanism to

keep the tropical climate stable. It is still an open question whether cloud feedbacks or large-scale dynamics are predominant in limiting tropical SST. Nevertheless, global observations of cloud profiles and related parameters, such as those presented here from MLS on Aura, will hopefully contribute to identifying the mechanisms and helping improve climate modeling.

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- E. J. Fetzer, J. H. Jiang, W. G. Read, H. Su, J. W. Waters, and D. L. Wu, M/S 183-701, Jet Propulsion Laboratory, California Institute of Technology, 4800 Oak Grove Drive, Pasadena, CA 91109-8099, USA. (Hui.Su@jpl.nasa.gov)