

Beyond Newtonian gravity: tests of GR

Let us use the Newtonian metric we derived last time,

$$ds^2 = -(1+2\phi)dt^2 + (1-2\phi)(dx^2 + dy^2 + dz^2)$$

where ϕ is the Newtonian gravitational potential.

Now let us consider the path of a photon through this geometry. We will choose the parameter λ for its trajectory

$$\text{so } \frac{dx^\mu}{d\lambda} = p^\mu.$$

The path should be a Minkowski-space geodesic (i.e. a straight line) + a perturbation depending on ϕ ,

$$x^\mu(\lambda) = \underbrace{x^{(0)\mu}(\lambda)}_{\substack{\downarrow \\ \text{background} \\ \text{path}}} + \underbrace{x^{(1)\mu}(\lambda)}_{\substack{\downarrow \\ \text{perturbation}}}$$

Subtlety: we want to evaluate quantities along the background path, & solve for $x^{(1)\mu}(\lambda)$

But \Rightarrow need to keep true path close to background path, so ϕ is similar on both

Approach: split path into short segments, approximate each one by a straight line (background path) when integrating along true path

$$\text{Let } k^\mu \equiv \frac{dx^{(0)\mu}}{d\lambda}, \quad L^\mu \equiv \frac{dx^{(1)\mu}}{d\lambda}$$

Photons follow null trajectories, $g_{\mu\nu}(k^\mu + L^\mu)(k^\nu + L^\nu) = 0$

Expand in h, L .

$$\text{0th order: } \eta_{\mu\nu} k^\mu k^\nu = 0. \quad \text{Let } k \equiv k^0 = \sqrt{|\vec{k}|^2}$$

$$\text{1st order: } 2\eta_{\mu\nu} k^\mu L^\nu + h_{\mu\nu} k^\mu k^\nu = 0$$

$$\begin{aligned} \Rightarrow -kL^0 + \vec{L} \cdot \vec{k} &= -\frac{1}{2} h_{\mu\nu} k^\mu k^\nu \\ &= -\frac{1}{2} [(-2\phi)k^2 + (-2\phi)|\vec{k}|^2] \\ &= 2k^2\phi \end{aligned}$$

Arrows now indicate
3-vectors, not general
vectors

The Christoffel symbols for our metric are

$$\Gamma^0_{oi} = \Gamma^i_{oo} = \partial_i \phi$$

$$\Gamma^i_{jke} = \delta_{jk} \partial_i \phi - \delta_{ik} \partial_j \phi - \delta_{ij} \partial_k \phi$$

Expanding the geodesic equation:

$$\frac{d}{d\lambda} (k^\mu + L^\mu) + \Gamma^\mu_{\rho\sigma} (k^\rho + L^\rho) (k^\sigma + L^\sigma) = 0$$

0th order: $\frac{dk^\mu}{d\lambda} = 0$, i.e. $k^\mu = \text{constant}$

1st order: $\frac{dL^\mu}{d\lambda} = -\Gamma^\mu_{\rho\sigma} k^\rho k^\sigma$

$$\Rightarrow \frac{dL^0}{d\lambda} = -\Gamma^0_{oi} k^o k^i - \Gamma^0_{io} k^o k^i$$

$$= -2k (\partial_i \phi) k^i$$

$$\frac{dL^i}{d\lambda} = -\Gamma^i_{oo} k^2 - \Gamma^i_{jk} k^j k^k$$

$$= -\partial_i \phi k^2 - (k^2 \partial_i \phi - \vec{k} \cdot \nabla \phi k^i - \vec{k} \cdot \nabla \phi k^i)$$

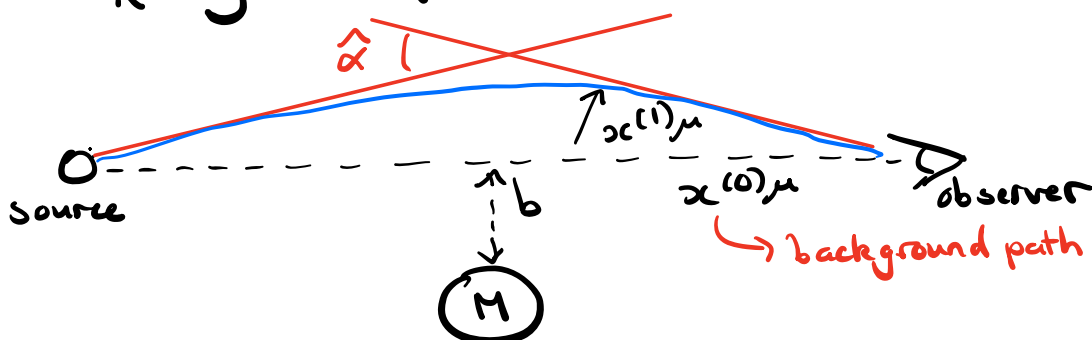
$$\frac{d\vec{L}}{d\lambda} = -2k^2 \left[\nabla \phi - \vec{k} \cdot \nabla \phi \frac{\vec{k}}{k^2} \right]$$

∇ here means the 3D grad operator, not covariant derivative

This is the gradient transverse to the path, $\nabla_{\perp} \phi$

Note $(\nabla_{\perp} \phi) \cdot \vec{k} = \vec{k} \cdot \nabla \phi - \vec{k} \cdot \nabla \phi = 0$.

Thus \vec{L} is perturbed in a direction perpendicular to \vec{k} by the potential ϕ .



We define the deflection angle as

$$\hat{\alpha} \equiv -\frac{\Delta \vec{L}}{k}, \text{ with } \Delta \vec{L} = \int \frac{d\vec{L}}{d\lambda} d\lambda$$

$$\Rightarrow \Delta \vec{L} = -2k^2 \int \nabla_{\perp} \phi d\lambda$$

If the source of ϕ is a point mass,

$$\phi = -\frac{GM}{r} = \frac{-GM}{\sqrt{x^2 + b^2}} \quad \text{for } \vec{k} \text{ in the } x\text{-direction, \& placing the mass at } x=0$$

$$\begin{aligned} \Rightarrow \nabla_{\perp} \phi &= \nabla \phi - (\nabla \phi)_x \hat{x} \\ &= \frac{-GM \cdot b}{(b^2 + x^2)^{3/2}} \hat{b} \\ &= \frac{+GM \vec{b}}{(b^2 + x^2)^{3/2}} \end{aligned}$$

$$\Rightarrow \hat{\alpha} = 2 \int \left(\frac{+GM \vec{b}}{(b^2 + x^2)^{3/2}} \right) \underbrace{k d\lambda}_{= dx, \text{ so } \frac{dx}{d\lambda} = k}$$

$$\begin{aligned} &= 2GM \vec{b} \int_{-\infty}^{\infty} \frac{dx}{(b^2 + x^2)^{3/2}} \quad \text{for path length } \gg b \\ &= \frac{4GM}{b^2} \vec{b} \end{aligned}$$

$$|\hat{\alpha}| = 4GM/b$$

Measured in 1919 using the Sun as the mass, & observing the positions of background stars during a solar eclipse.

For the Sun, $GM/c^2 = 1.48 \times 10^5 \text{ cm}$, &

b> Solar radius $R \approx 6.96 \times 10^{10} \text{ cm}$,
so $|\hat{\alpha}| \leq \frac{4 \times 1.48 \times 10^5}{6.96 \times 10^{10}} \approx 1.75 \text{ arcsec}$

Other classic tests of GR are gravitational redshift and the perihelion precession of Mercury.
the gravitational redshift follows directly from the equivalence principle & occurs even in uniform gravity, but the perihelion precession probes spacetime curvature and can be derived from the Schwarzschild (black hole) metric.