Heating and Cooling Estimators

Here I've tried to use Lucy's notation for macro-atom estimators. l and u represent lower and upper levels, and κ represents the continuum level or upper ion. q is the 'absorption fraction' derived below, and q_{ul} and q_{lu} are the collisional rate coefficients.

1. Macro-atom estimators

In the macro-atom approach, we basically treat two communication pathways. bound-free transitions represent a way for radiant energy to communicate with the thermal pool and bound-bound transitions represent a way for ionization/excitation energy to communicate with the thermal pool.

The heating and cooling rates for macro-atom bound-bound transitions are the rates of collisional excitations and de-excitations - i.e. the rate at which thermal energy is converted into bound-bound excitation energy and vice versa.

$$H_{bb,matoms} = \sum_{lines} q_{ul} n_u n_e h \nu_{ul} V \tag{1}$$

$$C_{bb,matoms} = \sum_{lines} q_{lu} n_l n_e h \nu_{ul} V \tag{2}$$

For bound-free transitions, the rate at which spontaneous recombinations convert thermal and ionization energy into radiant energy is $\alpha_{sp}^E h \nu_i n_{\kappa} n_e$, where $h \nu_i$ is the potential of the b-f transition. The amount of this energy which goes into the actual thermal pool is therefore given by

$$C_{bf,matoms} = \sum_{bfjumps} (\alpha_{sp}^E - \alpha_{sp} + \alpha_{st}^E - \alpha_{st}) n_e n_{\kappa} \nu_{\kappa l}$$
(3)

where here I have also included stimulated recombination as we do in the code. For photoionizations, we write a similar expression. The rate of at which a level l absorbs energy by b-f transitions is given by $\gamma^E h \nu_{\kappa l} n_{\kappa} n_e$, but the amount $\gamma_i h \nu_{\kappa l} n_i$ goes into ionization energy, giving

$$H_{bf,matoms} = \sum_{lines} (\gamma_{sp}^E - \gamma) n_l h \nu_{\kappa l} \tag{4}$$

as the rate at which radiant energy heats the plasma via b-f transitions.

2. Simple-ion estimators

In simple-ions it is in some ways a little more complicated. First we define q which will be different for each b-b transition, following Nick's thesis, which is given by (NB: I don't actually know how to derive this)

$$q = \frac{q_{ul}n_e(1 - e^{-h\nu/kT_e})}{\beta_{ul}A_{ul} + q_{ul}n_e(1 - e^{-h\nu/kT_e})}$$
(5)

where β_{ul} is the angle-averaged escape probability. q represents the probability that an excited bound electron will collisionally de-excite. Our b-b heating rate is computed during the photon propagation and is a sum over photons which come into resonance with each line, given by

$$H_{bb,simple} = \sum_{photons} \sum_{lines} (1 - q)(1 - e^{-\tau_S}) w_{photon}$$
(6)

And our bound bound cooling rate is given by

$$C_{bb,simple} = \sum_{lines} q \left(n_l \frac{g_u}{g_l} - n_u \right) q_{ul} n_e \frac{(1 - e^{-h\nu/kT_e})}{(e^{h\nu/kT_e} - 1)} h\nu_{ul}$$
 (7)

The bound-free heating rate is given by

$$H_{bf,simple} = \sum_{photons} \sum_{bfjumps} w_{photon} e^{-\tau} \frac{\nu - \nu_{\kappa l}}{\nu}$$
 (8)

where ν here is the frequency of the photon in question. The bound-free cooling rate is then

$$C_{bf,simple} = \sum_{bfjumps} \alpha_{sp} n_e n_{\kappa} \tag{9}$$