## Heating and Cooling Estimators in Python: Macro-atoms and Simple-ions

Here I've tried to use Lucy's notation for macro-atom estimators. Take a three level system, in which l and u represent lower and upper levels, and  $\kappa$  represents the continuum level or upper ion. q is the 'absorption fraction' derived below, and  $q_{ul}$  and  $q_{lu}$  are the collisional rate coefficients.

## 1. Macro-atom estimators

In the macro-atom approach, we basically treat two communication pathways. bound-free transitions represent a way for radiant energy to communicate with the thermal pool and bound-bound transitions represent a way for excitation energy to communicate with the thermal pool.

The heating and cooling rates for macro-atom bound-bound transitions are the rates of collisional excitations and de-excitations - i.e. the rate at which thermal energy is converted into bound-bound excitation energy and vice versa.

$$C_{bb,matoms} = \sum_{lines} q_{lu} n_l n_e h \nu_{ul} V \tag{1}$$

$$H_{bb,matoms} = \sum_{lines} q_{ul} n_u n_e h \nu_{ul} V \tag{2}$$

For bound-free transitions, we define the normal photoionization and recombination rate coefficients  $\gamma$  and  $\alpha$ , where  $\alpha$  includes stimulated recombination as we do in the code. Note this differs to the approach in Lucy (2003), where it is instead included as a negative photoionization term, hence the notation  $\tilde{\gamma}$ . We also need to define two 'modified rate coefficients' which are the rates at which b-f transitions add and remove energy to the radiation field. These are denoted  $\gamma^E$  and  $\alpha^E$ .

The rate at which recombinations convert thermal and ionization energy into radiant energy is then  $\alpha^E h \nu_{\kappa l} n_{\kappa} n_e$ , where  $h \nu_{\kappa l}$  is the potential of the b-f transition, or the energy difference between continuum  $\kappa$  and the level l we are recombining too. The amount of this energy which is removed from the actual thermal pool therefore needs a quantity  $\alpha h \nu_{\kappa l} n_{\kappa} n_e$  subtracted from it, giving

$$C_{bf,matoms} = \sum_{bfjumps} (\alpha^E - \alpha) n_e n_\kappa \nu_{\kappa l} V$$
(3)

where here I have also included stimulated recombination as we do in the code. Note this differs to the approach in Lucy (2003), where it is instead included as a negative photoionization term, hence the notation  $\tilde{\gamma}$ . For photoionizations, we write a similar expression. The rate of at which a level l absorbs energy by b-f transitions is given by  $\gamma^E h \nu_{\kappa l} n_{\kappa} n_e$ , but the amount  $\gamma h \nu_{\kappa l} n_l$  goes into ionization energy, giving

$$H_{bf,matoms} = \sum_{bfjumps} (\gamma^E - \gamma) n_l h \nu_{\kappa l} V \tag{4}$$

as the rate at which radiant energy heats the plasma via b-f transitions.

## 2. Simple-ion estimators

In simple-ions it is in some ways a little more complicated. First we define q which will be different for each b-b transition, following Nick's thesis, which is given by (NB: I don't actually know how to derive this)

$$q = \frac{q_{ul}n_e(1 - e^{-h\nu/kT_e})}{\beta_{ul}A_{ul} + q_{ul}n_e(1 - e^{-h\nu/kT_e})}$$
(5)

where  $\beta_{ul}$  is the angle-averaged escape probability. q represents the probability that an excited bound electron will collisionally de-excite. Our b-b heating rate is computed during the photon propagation and is a sum over photons which come into resonance with each line, given by

$$H_{bb,simple} = \sum_{photons} \sum_{lines} (1 - q)(1 - e^{-\tau_S}) w_{photon}$$
 (6)

And our bound bound cooling rate is given by

$$C_{bb,simple} = \sum_{lines} q \left( n_l \frac{g_u}{g_l} - n_u \right) q_{ul} n_e \frac{(1 - e^{-h\nu/kT_e})}{(e^{h\nu/kT_e} - 1)} h\nu_{ul}$$
 (7)

The bound-free heating rate is given by

$$H_{bf,simple} = \sum_{photons} \sum_{bfjumps} w_{photon} e^{-\tau} \frac{\nu - \nu_0}{\nu}$$
(8)

where  $\nu$  here is the frequency of the photon in question, and  $\nu_0$ . The bound-free cooling rate is then

$$C_{bf,simple} = ?? (9)$$