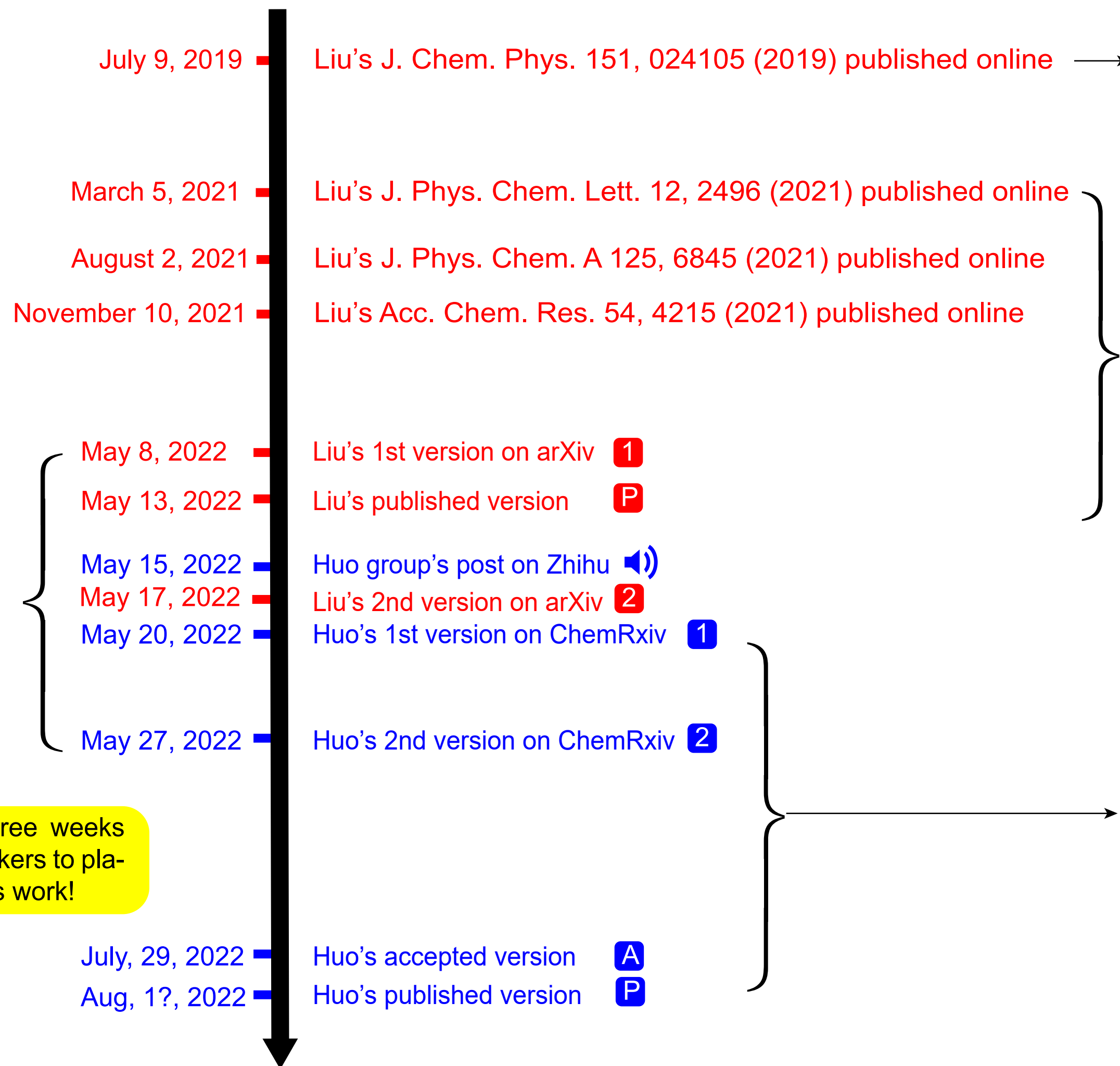


Evidence 2: PLAGIARISM

Timeline of Events

Statements



In this paper, we first pointed out that the **sum of the electron populations is equal to 1** is a natural constraint for the Meyer-Miller mapping variables for **general F -state systems**, when parameter $\gamma=0$:

$$\sum_{n=1}^F \frac{1}{2} \left[\left(x^{(n)} \right)^2 + \left(p^{(n)} \right)^2 \right] = 1$$

In these papers, we further extended the relations to any $\gamma > -1/F$ for **general F -state systems**:

$$\sum_{n=1}^F \frac{1}{2} \left[\left(x^{(n)} \right)^2 + \left(p^{(n)} \right)^2 \right] = 1 + F\gamma$$

Plagiarized from

The constraint of the Meyer-Miller mapping variables:

$$\sum_{n=1}^N \frac{1}{2} \left(p_n^2 + q_n^2 - \gamma \right) = 1$$

Evidence #2:

In the last paragraph of page 7 of Version 1 of Huo and his coworkers (first released online to ChemRxiv on May 20, 2022):

“

In the $\mathfrak{su}(N)$ mapping formalism, the total population constraint on the $2N$ -dimensional phase space comes naturally from the normalization of the generalized spin coherent state^{53,57} as follows

$$\langle \Omega | \Omega \rangle = \sum_n c_n^* c_n = \frac{1}{2r_s} \sum_{n=1}^N (q_n^2 + p_n^2) = 1, \quad (53)$$

which properly enforces the total electronic diabatic population to be one for these MMST mapping variables

$$\sum_{n=1}^N \frac{1}{2} (q_n^2 + p_n^2 - \gamma) = 1. \quad (54)$$

Alternatively, one can obtain this condition from the basic property of the SW transform that preserves the trace of the electronic identity operator (Eq. 28) as follows

$$[\hat{\mathcal{T}}]_s = \sum_{n=1}^N [|n\rangle\langle n|]_s = 1 - r_s + \sum_n \frac{1}{2} (q_n^2 + p_n^2) = 1.$$

”

as well as the lower left part of page 9 of Version 2 (first released online to ChemRxiv on May 27, 2022):

“

In the $SU(N)$ mapping formalism, the total population constraint on the $2N$ -dimensional phase space comes *naturally* from the normalization of the generalized spin coherent states^{69,75} as follows

$$\langle \Omega | \Omega \rangle = \sum_n c_n^* c_n = \frac{1}{2r_s} \sum_{n=1}^N (q_n^2 + p_n^2) = 1, \quad (56)$$

which properly enforces the total electronic diabatic population to be one (see Eq. 51) for these MMST mapping variables

$$\sum_{n=1}^N \frac{1}{2} (q_n^2 + p_n^2 - \gamma) = 1. \quad (57)$$

Alternatively, one can obtain this condition from the basic property of the S-W transform that preserves the trace of the electronic identity operator (Eq. 30) as follows

$$[\hat{\mathcal{I}}]_s = \sum_{n=1}^N [|n\rangle\langle n|]_s = 1 - r_s + \sum_n \frac{1}{2}(q_n^2 + p_n^2) = 1.$$

Note that the recent work of the eCMM is developed based on *manually* adding an extra total population constraint (as described in Eq. 57) on the MMST mapping oscillator phase space. Historically, it was realized⁷⁵ that a mapping from the quantum Schrödinger's equation to $2N$ classical phase space Hamilton's EOMs is incorrect,

”

and the left part of page 24 of Version 3 (accepted on **July 29, 2022**):

“

In the $SU(N)$ mapping formalism, the total population constraint on the $2N$ -dimensional phase space comes *naturally* from the normalization of the generalized spin coherent states^{69,108} as follows

$$\langle \Omega | \Omega \rangle = \sum_n c_n^* c_n = \frac{1}{2r_s} \sum_{n=1}^N (q_n^2 + p_n^2) = 1, \quad (\text{D11})$$

which properly enforces the total electronic diabatic population to be one (see Eq. D6) for these MMST mapping variables

$$\sum_{n=1}^N \frac{1}{2}(q_n^2 + p_n^2 - \gamma) = 1. \quad (\text{D12})$$

Alternatively, one can obtain this condition from the basic property of the S-W transform that preserves the trace of the electronic identity operator (Eq. 30) as follows

$$[\hat{\mathcal{I}}]_s = \sum_{n=1}^N [|n\rangle\langle n|]_s = 1 - r_s + \sum_n \frac{1}{2}(q_n^2 + p_n^2) = 1.$$

Note that the recent work of the eCMM is developed based on *manually* adding an extra total population constraint (as described in Eq. D12) on the MMST mapping oscillator phase space. Historically, it was realized¹⁰⁸ that a mapping from the quantum Schrödinger's equation

”

Huo and his coworkers mentioned the constraint that the sum over electronic population is equal to 1, but neglected the fact that this normalization/constraint had been first proposed in eq (28) of *J. Chem. Phys.* 151, 024105 (2019) [submitted on **May 1, 2019**, accepted on **June 11, 2019** and published on **July 9, 2019**]:

“

The conservation of the total population Eq. (16) is implicitly used for the kinetic energy term in Eq. (23). (This is a new observation that is essentially the key point of this paper.) The classical Hamiltonian for both nuclear and electronic DOFs is then

”
“

$$\sum_{n=1}^F P_n(0) = 1 \text{ and } 0 \leq P_n(0) \leq 1 \quad (\forall n). \quad (28)$$

”

In the paper *J. Chem. Phys.* 151, 024105 (2019), the constraint phase space manifold had been firstly utilized by us to implement phase space mapping for nonadiabatic dynamics, which in principle could be parameterized by either Meyer-Miller variables or (Stratonovich) angle variables. Using this constraint, the Q, W, or P versions of Stratonovich phase space used in spin mapping methods proposed in *J. Chem. Phys.* 152, 084110 (2020) were only three special cases of the constraint coordinate-momentum phase space in classical mapping model(CMM) methods. Later, the relationship was further clarified again in our paper *Wiley Interdiscip. Rev. Comput. Mol. Sci.* e1619 (2022) [submitted on **February 5, 2022**, released on **arXiv on May 8, 2022** and officially published on **May 13, 2022**]. Neither Version 1 nor Version 2 of Huo and his coworkers cited this article.

In our *J. Phys. Chem. Lett.* 12, 2496–2501 (2021) [submitted on **January 22, 2021**, accepted on **February 19, 2021** and published online on **March 5, 2021**], we had clearly clarified the one-to-one mapping (i.e., mapping and inversed mapping) between identity operator and 1:

“

Supporting Information.) As expected, the kernel and its inverse are properly normalized

$$\begin{aligned} \text{Tr}_e[\hat{K}(\mathbf{x}, \mathbf{p})] &= \text{Tr}_e[\hat{K}^{-1}(\mathbf{x}, \mathbf{p})] = 1 \\ \int_{S(\mathbf{x}, \mathbf{p})} d\mu(\mathbf{x}, \mathbf{p}) \hat{K}(\mathbf{x}, \mathbf{p}) &= \int_{S(\mathbf{x}, \mathbf{p})} d\mu(\mathbf{x}, \mathbf{p}) \hat{K}^{-1}(\mathbf{x}, \mathbf{p}) = \hat{I}_e \end{aligned} \quad (9)$$

where \hat{I}_e is the identity operator in F -dimensional Hilbert space

”

It is natural that the constraint of the manifold corresponds to the one-to-one mapping conditions, thus it is never a “*manually*” added constraint. Huo and his coworkers interpreted *our* ideas in **their** way and claimed **they** firstly got these ideas.